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Sampling Non-Abelian Lattice Gauge Theories with Diffusion Models

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Thomas Ranner

Institute for Theoretical Physics, TU Wien
thomas.ranner@tuwien.ac.at

Based on

arXiv > hep-lat > arXiv:2601.19552

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High Energy Physics - Lattice

[Submitted on 27 Jan 2026]

Generalizable Equivariant Diffusion Models for Non-Abelian Lattice Gauge Theory

Gert Aarts, Daa E. Habibi, Andreas Ipp, David I. Müller, Thomas R. Ranner, Lingxiao Wang, Wei Wang, Qianteng Zhu

We demonstrate that gauge equivariant diffusion models can accurately model the physics of non-Abelian lattice gauge theory using the Metropolis-adjusted annealed Langevin algorithm (MAALA), as exemplified by computations in two-dimensional $U(2)$ and $SU(2)$ gauge theories. Our network architecture is based on lattice gauge equivariant convolutional neural networks (L-CNNs), which respect local and global symmetries on the lattice. Models are trained on a single ensemble generated using a traditional Monte Carlo method. By studying Wilson loops of various size as well as the topological susceptibility, we find that the diffusion approach generalizes remarkably well to larger inverse couplings and lattice sizes with negligible loss of accuracy while retaining moderately high acceptance rates.

Motivation

- Lattice quantum field theory as alternative to perturbative methods in QFT
- Very slow (calculations may take weeks on HPCs)
- Solving the path integral to calculate observables

$$\langle O \rangle = \frac{1}{Z} \int DU O(U) e^{-S_E(U)}, \quad Z = \int DU e^{-S_E(U)}$$

- By using Monte Carlo integration:

$$\langle O \rangle \approx \frac{1}{M} \sum_{i=1}^M O(U_i), \quad U_i \sim p_0(U) = \frac{e^{-S_E(U)}}{Z}$$

- Develop diffusion models for these problems to possibly speed up calculations in the future.
 - For now, demonstrate that they can be applied

Introduction to diffusion models

- Generative machine learning method
- Most famous for image generation
- Based on stochastic processes
- Draw random samples from complicated probability distribution

This summer, he got invited to a Discord chat server where people were testing Midjourney, **which uses a complex process known as “diffusion”** to turn text into custom images. Users type a series of words in a message to Midjourney; the bot spits back an image seconds later.

An A.I.-Generated Picture Won an Art Prize. Artists Aren't Happy.

“I won, and I didn’t break any rules,” the artwork’s creator says.

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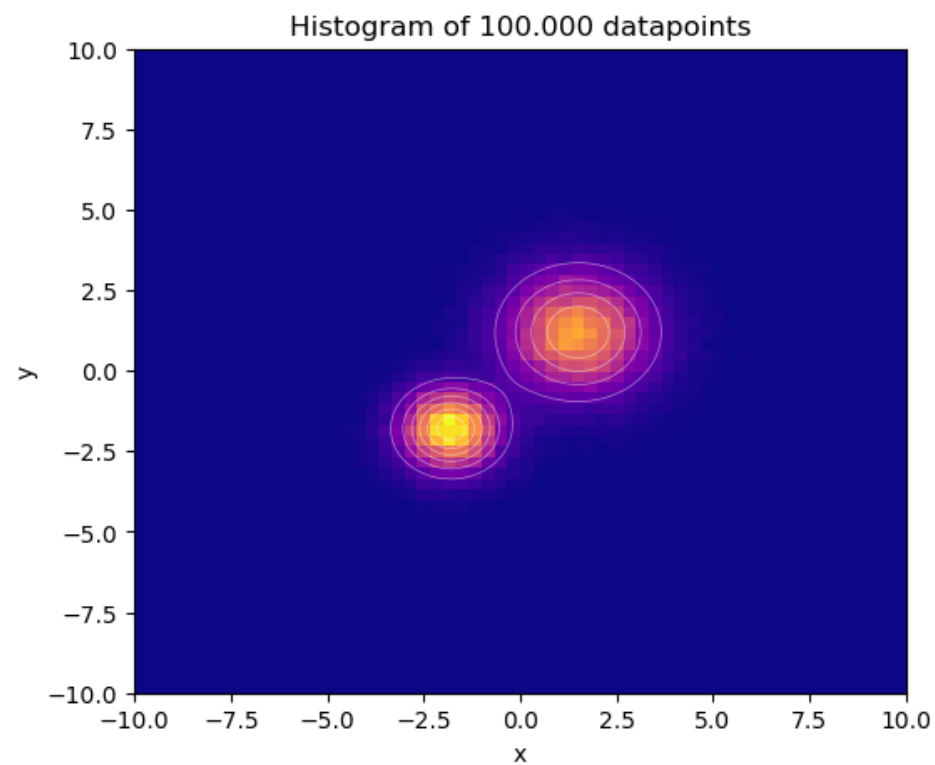
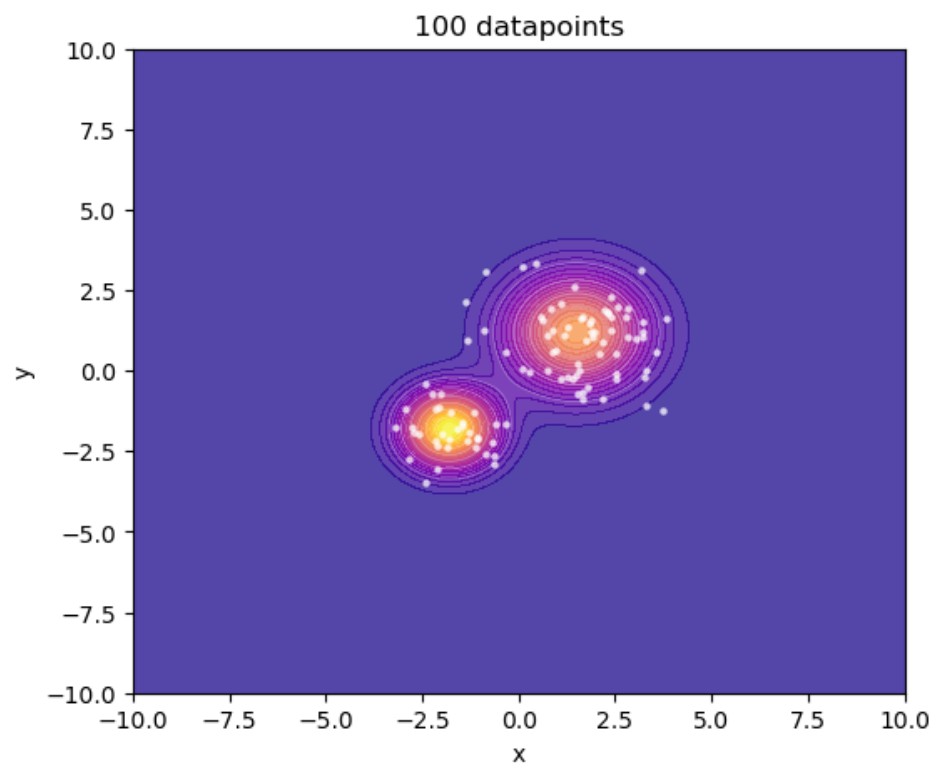
Jason Allen's A.I.-generated work, “Théâtre D’opéra Spatial,” took first place in the digital category at the Colorado State Fair. via Jason Allen



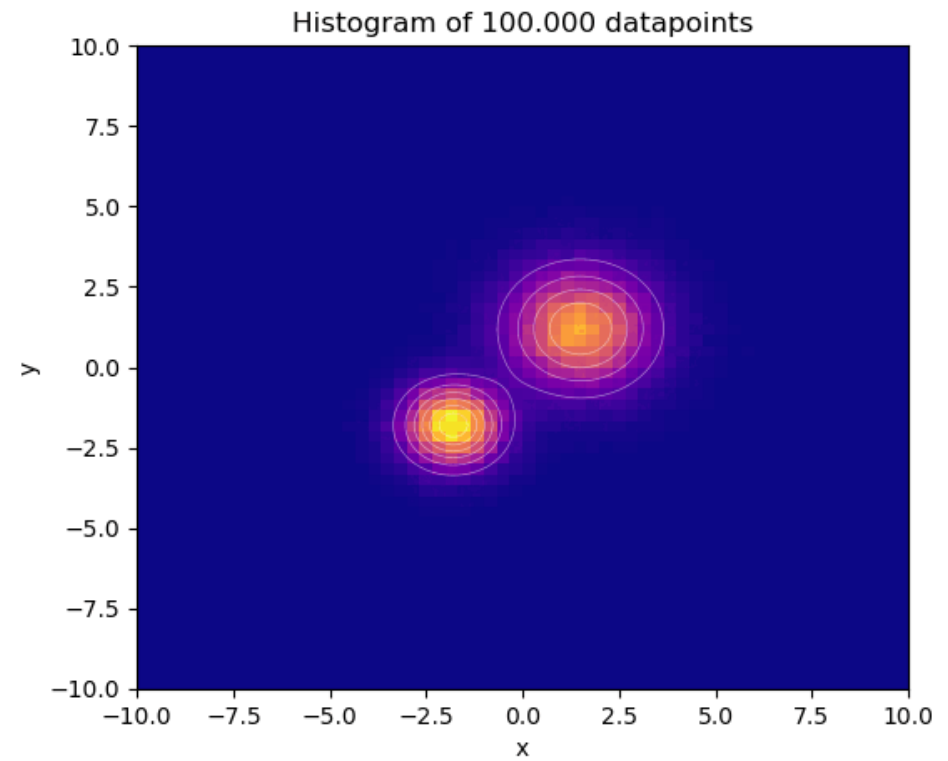
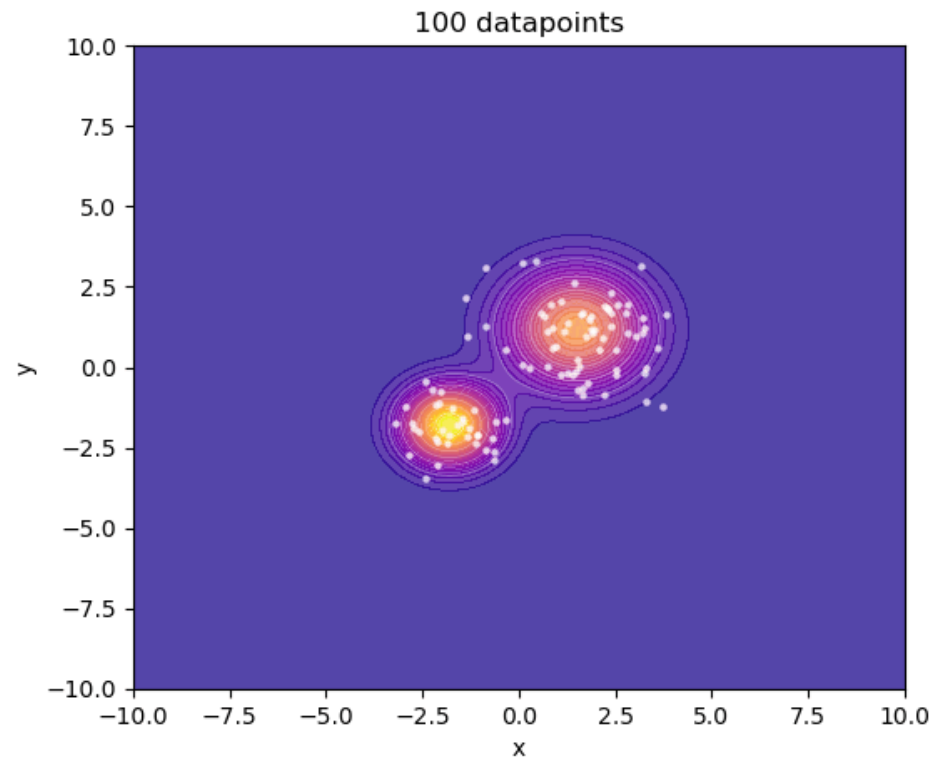
By Kevin Roose

Sept. 2, 2022

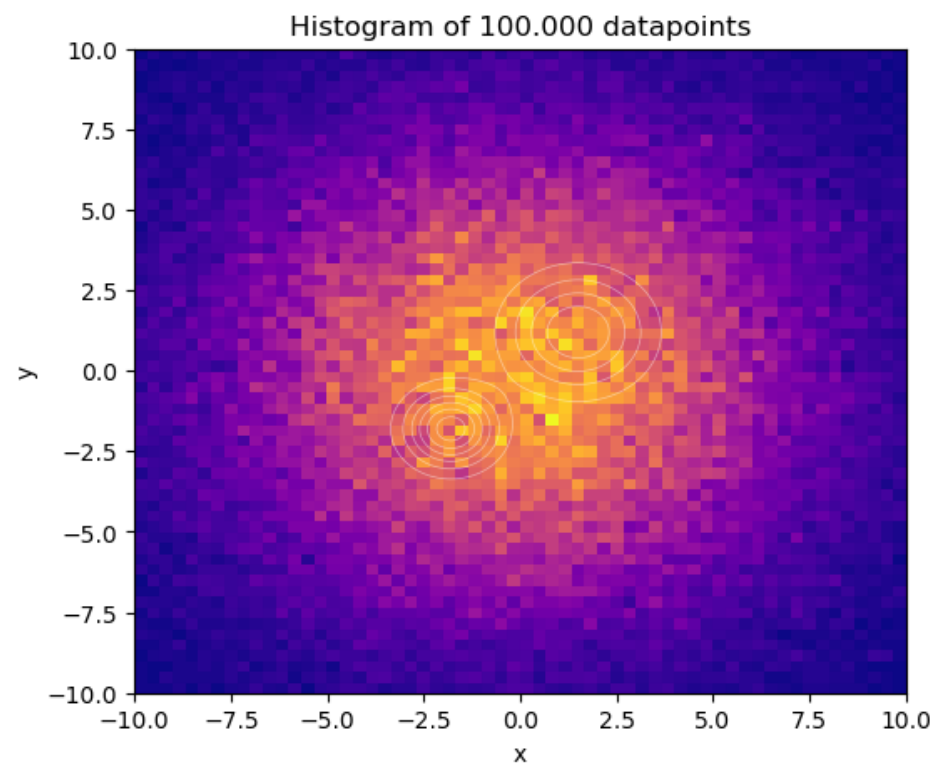
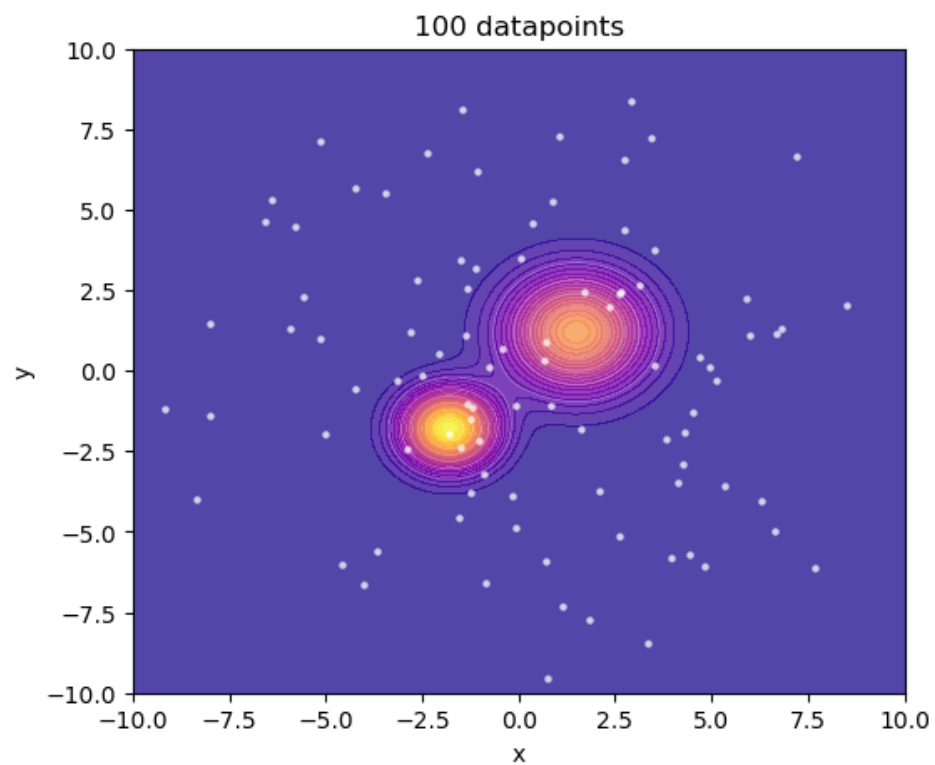
Introduction to diffusion models



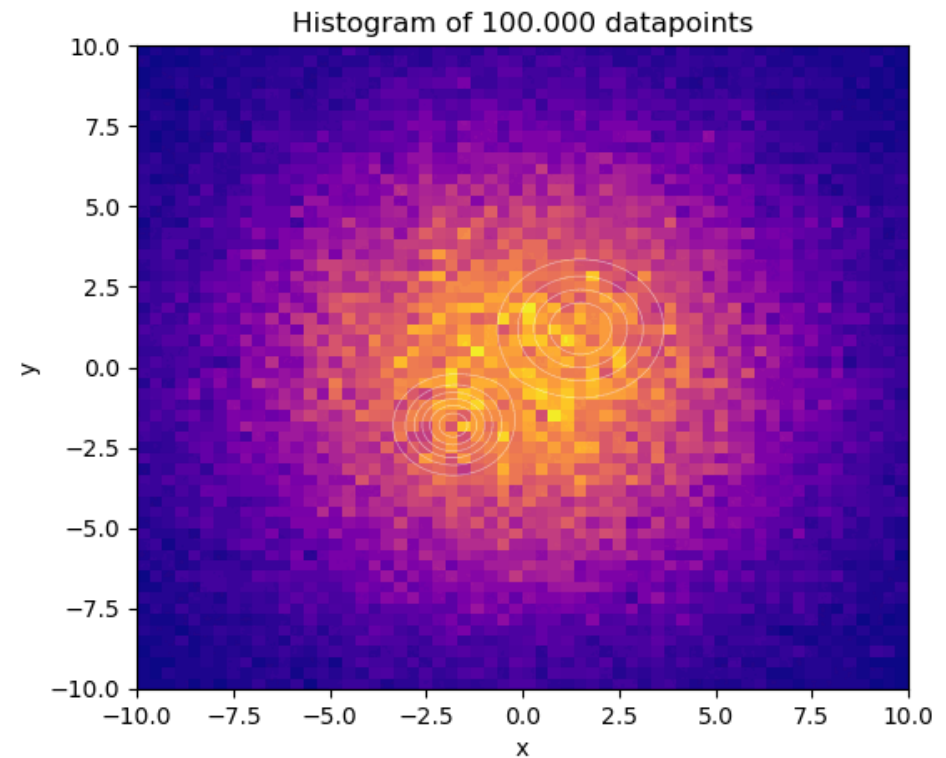
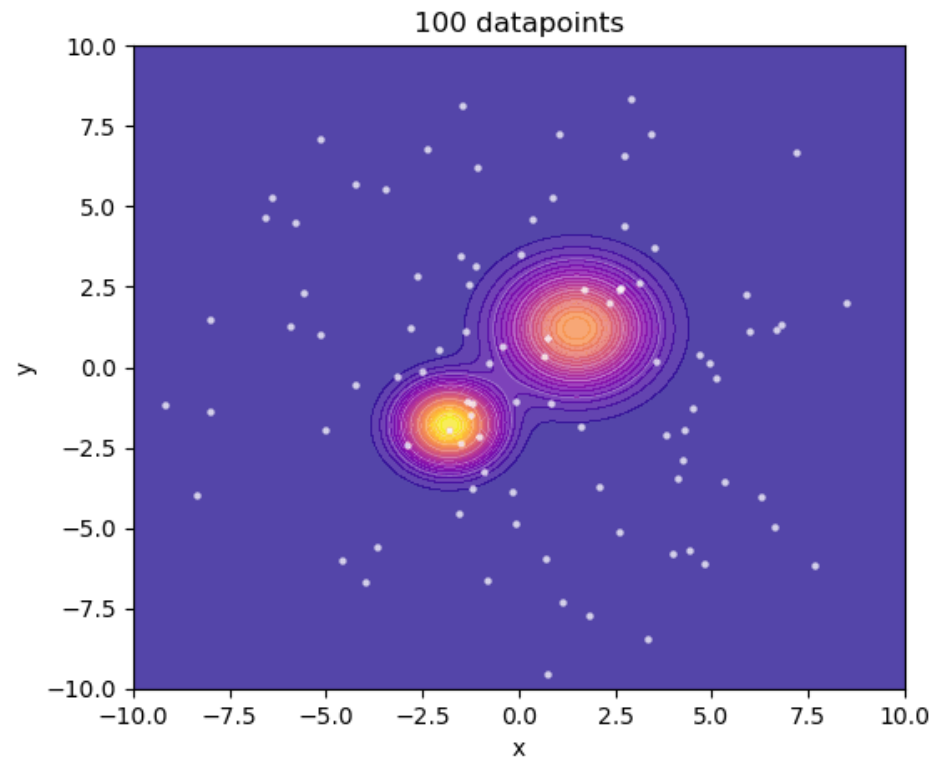
Introduction to diffusion models



Introduction to diffusion models



Introduction to diffusion models



Diffusion models for non-Abelian gauge theories

- Diffusion on group elements of $U(N)$ or $SU(N)$
- Complex unitary $N \times N$ matrices

$$U = \exp \left(i \sum_j a_j T_j \right)$$

Diagram illustrating the components of the equation:

- U is labeled as the **Group element**.
- a_j is labeled as **Scalar coefficients**.
- T_j is labeled as **Generators of the group**.

e.g. for $U(2)$:

$$\begin{aligned} & \frac{1}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \\ & \frac{1}{2} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \\ & \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \\ & \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \end{aligned}$$

- Construct diffusion process that stays in the group

Diffusion models for non-Abelian gauge theories

- Stratonovich stochastic differential equation (SDE):

$$dU_t = i \sum_a T^a (K^a(U_t)dt + g(t)dW_t^a) \circ U_t$$

Drift term

Diffusion term

Derivation via
Fokker-Planck
equation

- Reverse process:

$$dU_t = i \sum_a T^a [(K^a(U^{(t)}) - g(t)^2 D^a \ln p_t(U^{(t)}))dt + g(t)d\tilde{W}_t^a] \circ U_t$$

Score function
In general unknown, will be
learned by a neural network.

Trained using samples from
the target distribution

Non-Abelian lattice gauge theory

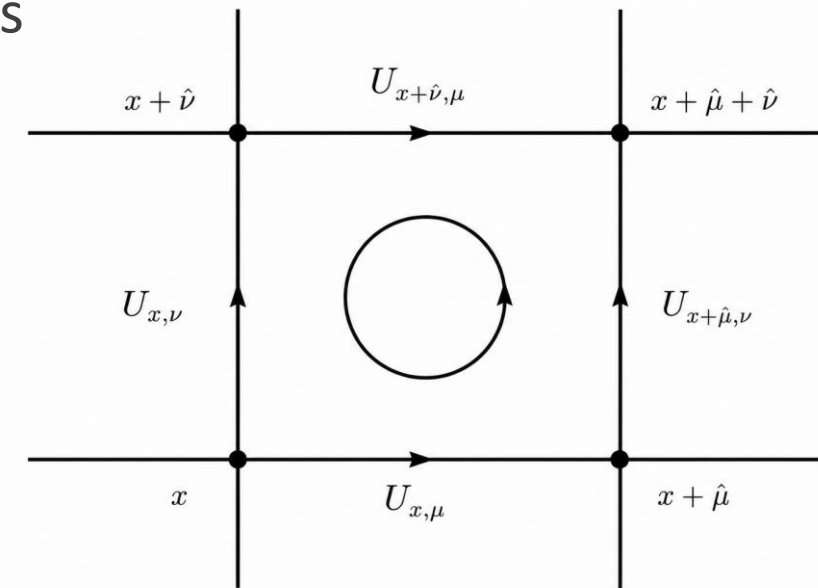
- Discretization of the gauge field on a Hypercubic lattice
- Link variables $U_{x,\mu}$ as connections between lattice points
- Links are elements of the gauge group
- Plaquette, elemental closed loop:

$$P_{x,\mu\nu} = U_{x,\mu} U_{x+\hat{\mu},\nu} U_{x+\hat{\nu},\mu}^\dagger U_{x,\nu}^\dagger$$

- Wilson action:

$$S_E(U) = -\beta \sum_{x,\mu < \nu} \text{Re Tr}[P_{x,\mu\nu}]$$

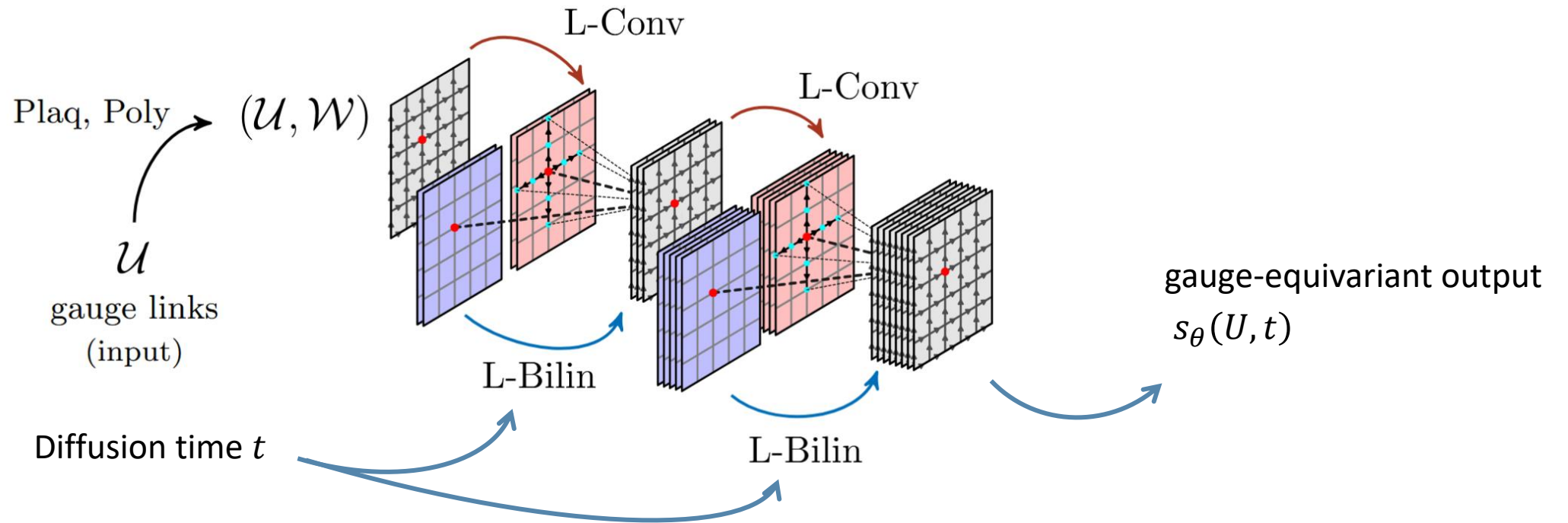
- β ... inverse coupling constant



Based on: Gattringer and Lang, Quantum chromodynamics on the lattice

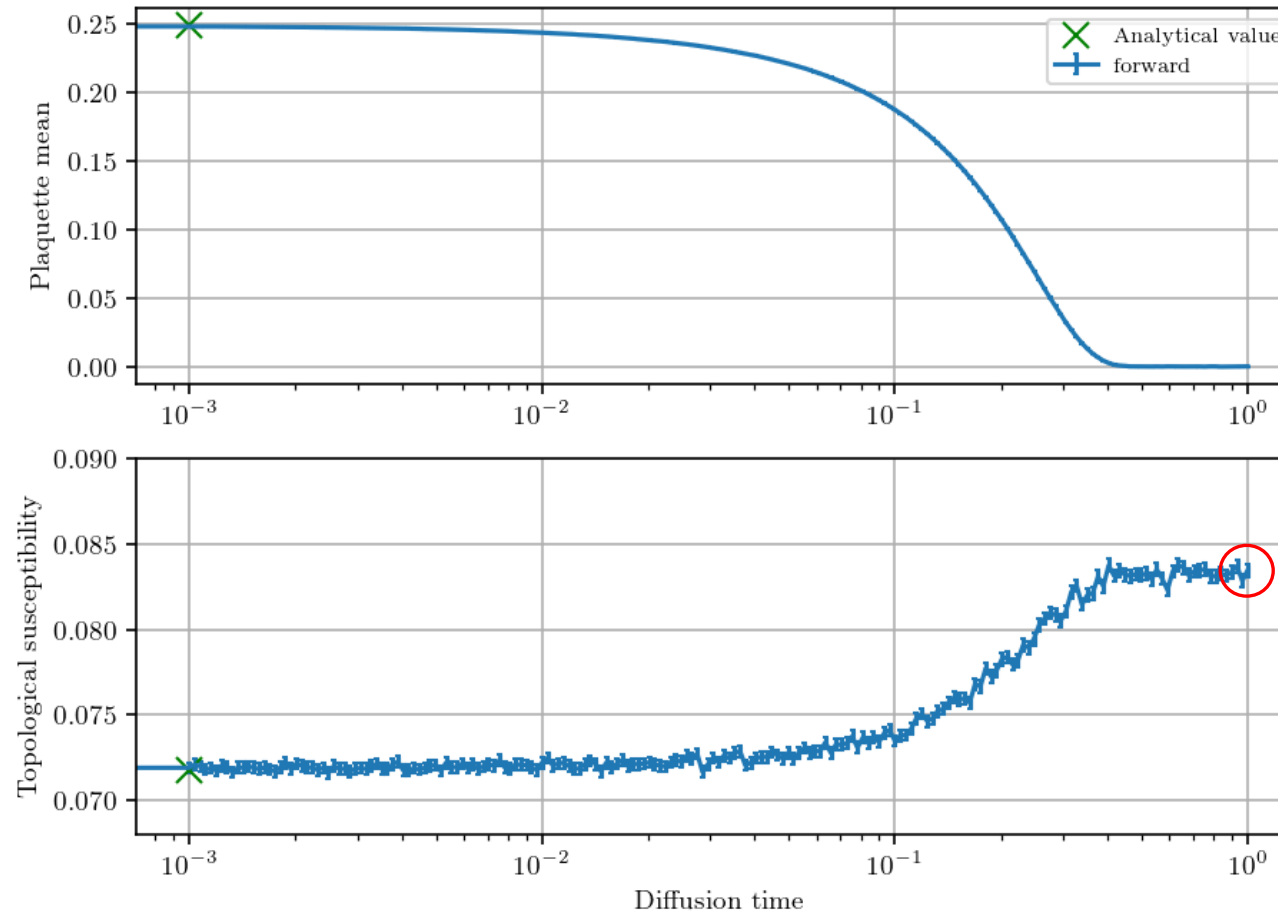
Lattice gauge-equivariant convolutional neural networks (L-CNNs)

- Adaption of convolutional NNs to lattice gauge theories



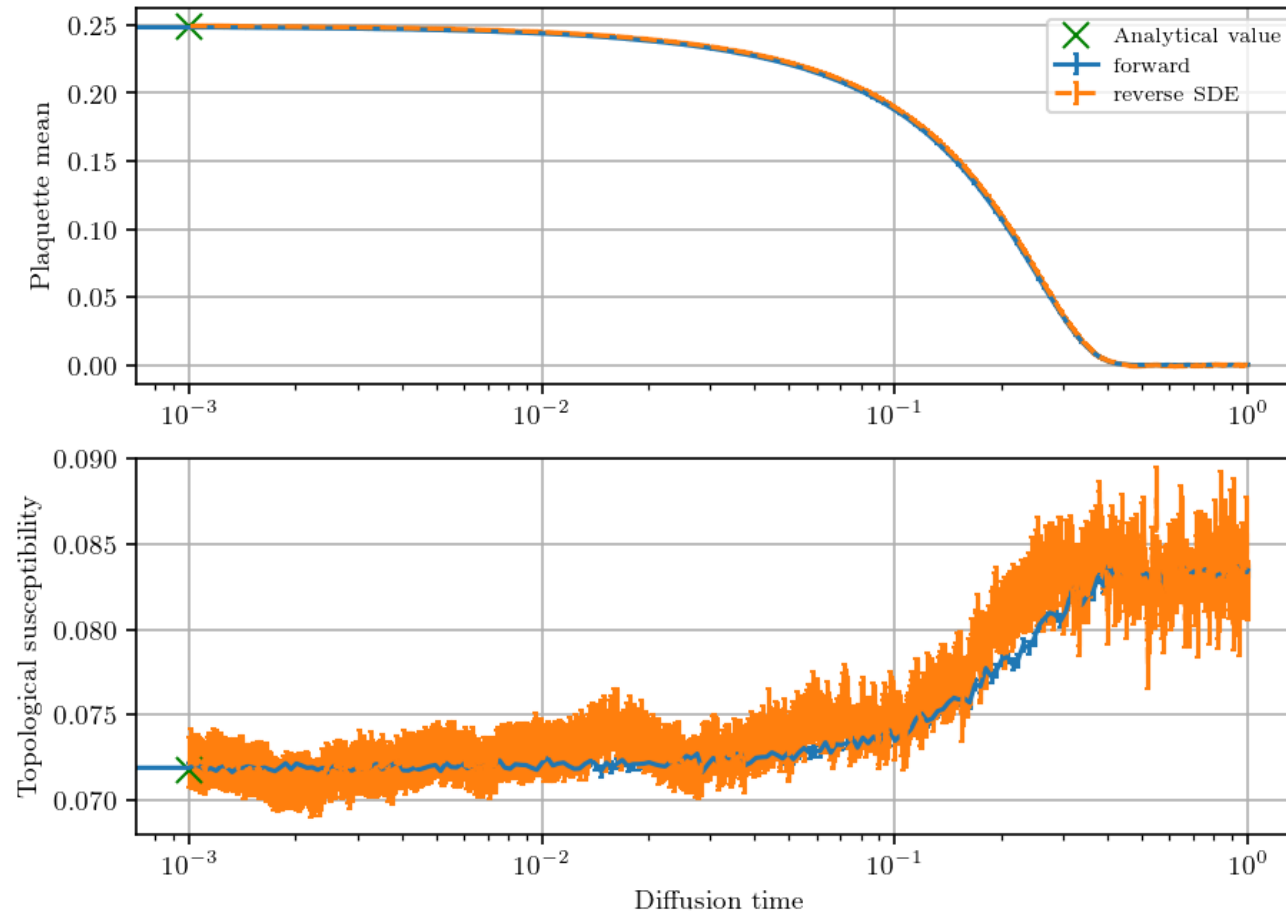
Source: Favoni et al., "Lattice Gauge Equivariant Convolutional Neural Networks"

Reverse SDE results for 2D U(2)



$\frac{1}{12}$...strong coupling limit

Reverse SDE results for 2D U(2)



But why?

- Goal is to draw configurations from a probability distribution
- But we need configurations to train the diffusion model
- Generate more configurations faster
 - For images it has been shown that the DM learns the underlying probability distribution
- Extrapolate to different lattice sizes
- Interpolate/extrapolate to different couplings
 - Two options: Learn the coupling dependence (not done yet)
 - Alternative sampling method MAALA – advantage: Metropolis adjustment for exact results

For 5.000 samples:
Reverse SDE: 6min
HMC: 55min

Stochastic Quantization (SQ)

- Langevin equation:

$$dU_\tau = i \sum_a T^a \left(-D^a S_E(U_\tau) d\tau + \sqrt{2} dW_\tau^a \right) \circ U_\tau$$

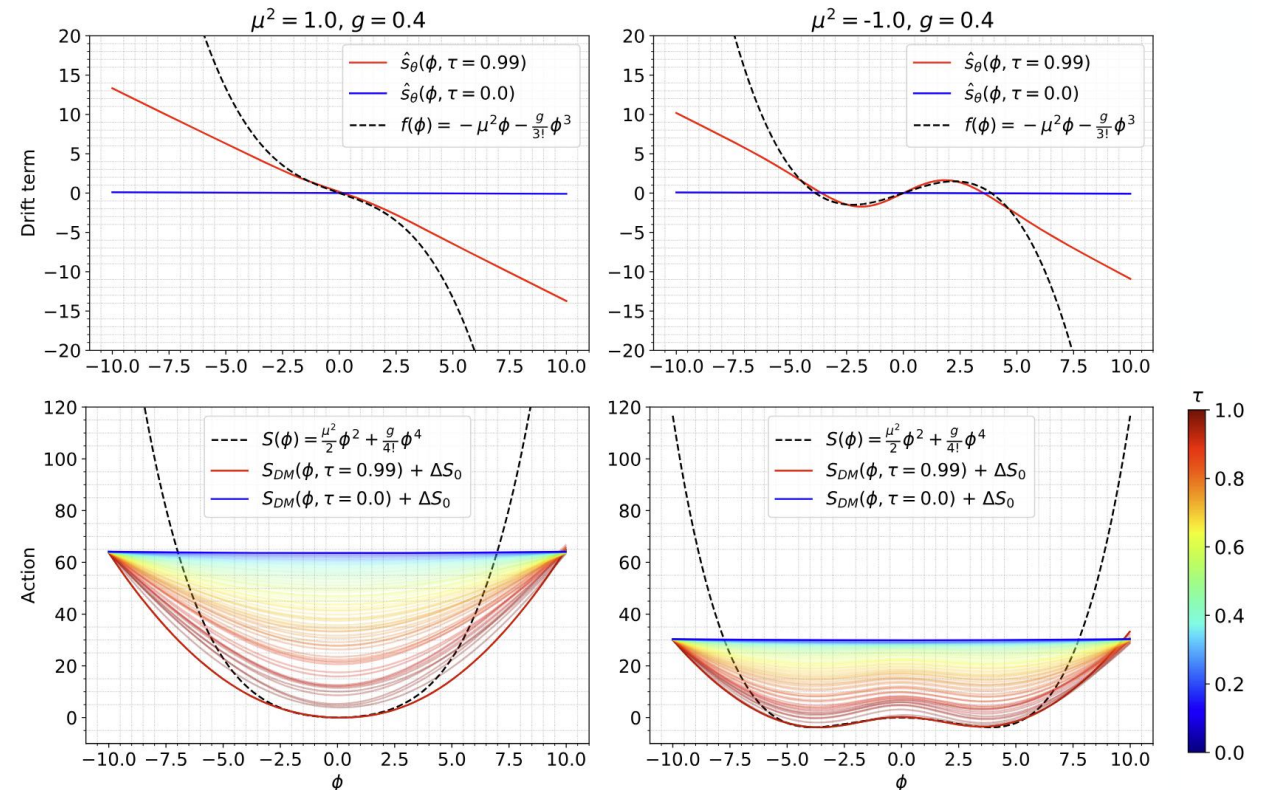
- For $\tau \rightarrow \infty$ configurations U_τ will follow

$$p(U) = \frac{1}{Z} e^{-S_E(U)} \xrightarrow{D^a \ln} -D^a S_E(U)$$

- Score function $s^a(U) = D^a \ln(p(U))$ is the drift of SQ

Score function

- Time dependent score-function
- Learns drift of some evolving action along the diffusion process
- Run SQ for a number of intermediate actions
- Annealed sampling
 - Gradually changing the coupling constant during sampling

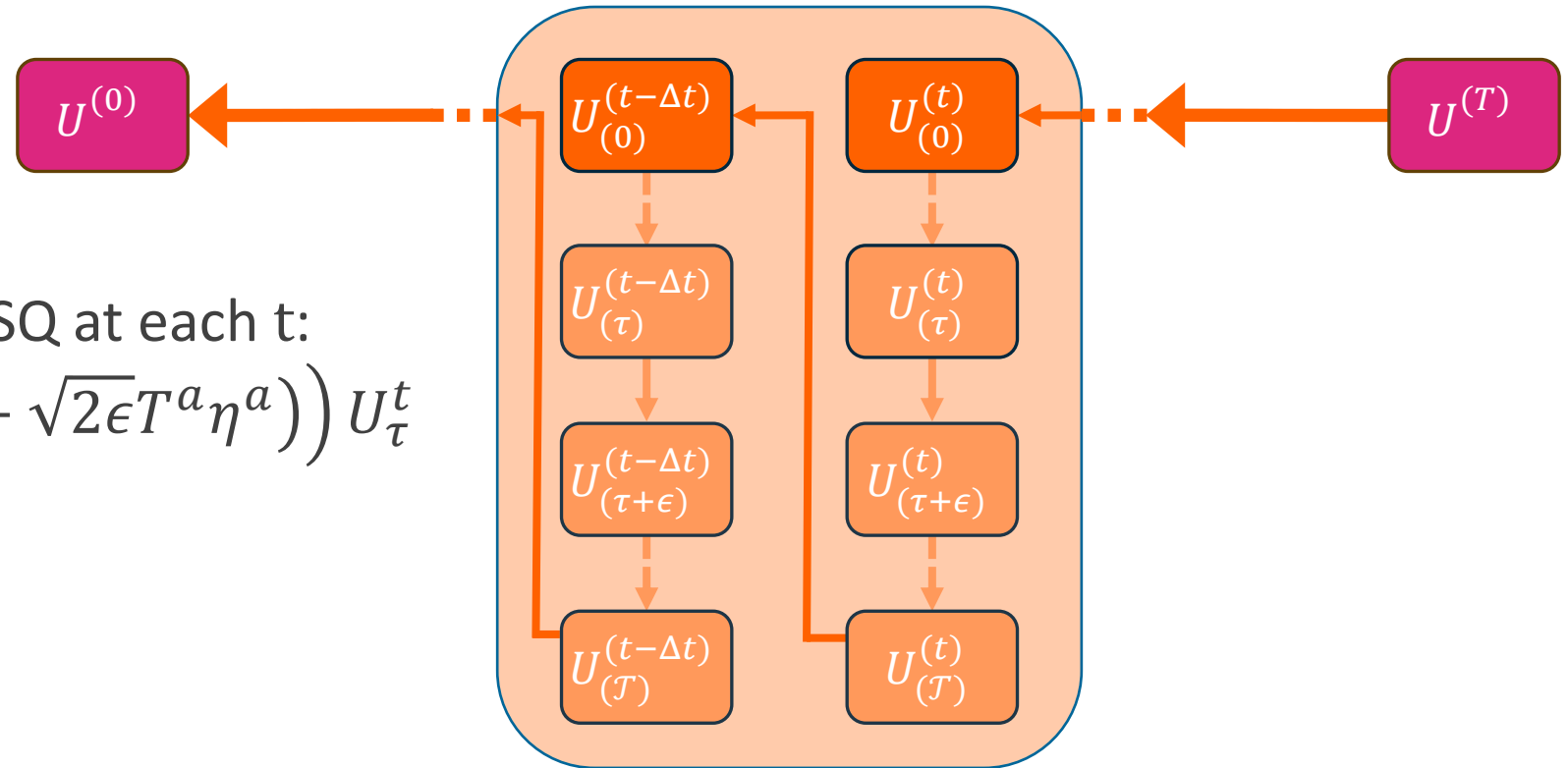


Source: Wang, Aarts, and Zhou, "Diffusion models as stochastic quantization in lattice field theory"

Metropolis adjusted annealed Langevin (MAALA)

- On a diffusion time grid, run SQ at each t :

$$U_{\tau+\epsilon}^t = \exp \left(i \left(-\epsilon s_{\theta}(U_{\tau}^t, t) + \sqrt{2\epsilon} T^a \eta^a \right) \right) U_{\tau}^t$$



Metropolis adjusted annealed Langevin (MAALA)

- This diffusion process traces the same path as the reverse SDE
- What is the advantage?
- Scaling with beta

$$U_{\tau+\epsilon}^t = \exp \left(i \left(-\epsilon \frac{\beta}{\beta_0} s_{\theta}(U_{\tau}^t, t) + \sqrt{2\epsilon} T^a \eta^a \right) \right) U_{\tau}^t$$

$$\frac{\beta}{\beta_0} s_{\theta}^a(U, t=0) = -\frac{\beta}{\beta_0} D^a S_E(\beta_0; U) = -D^a S_E(\beta; U)$$

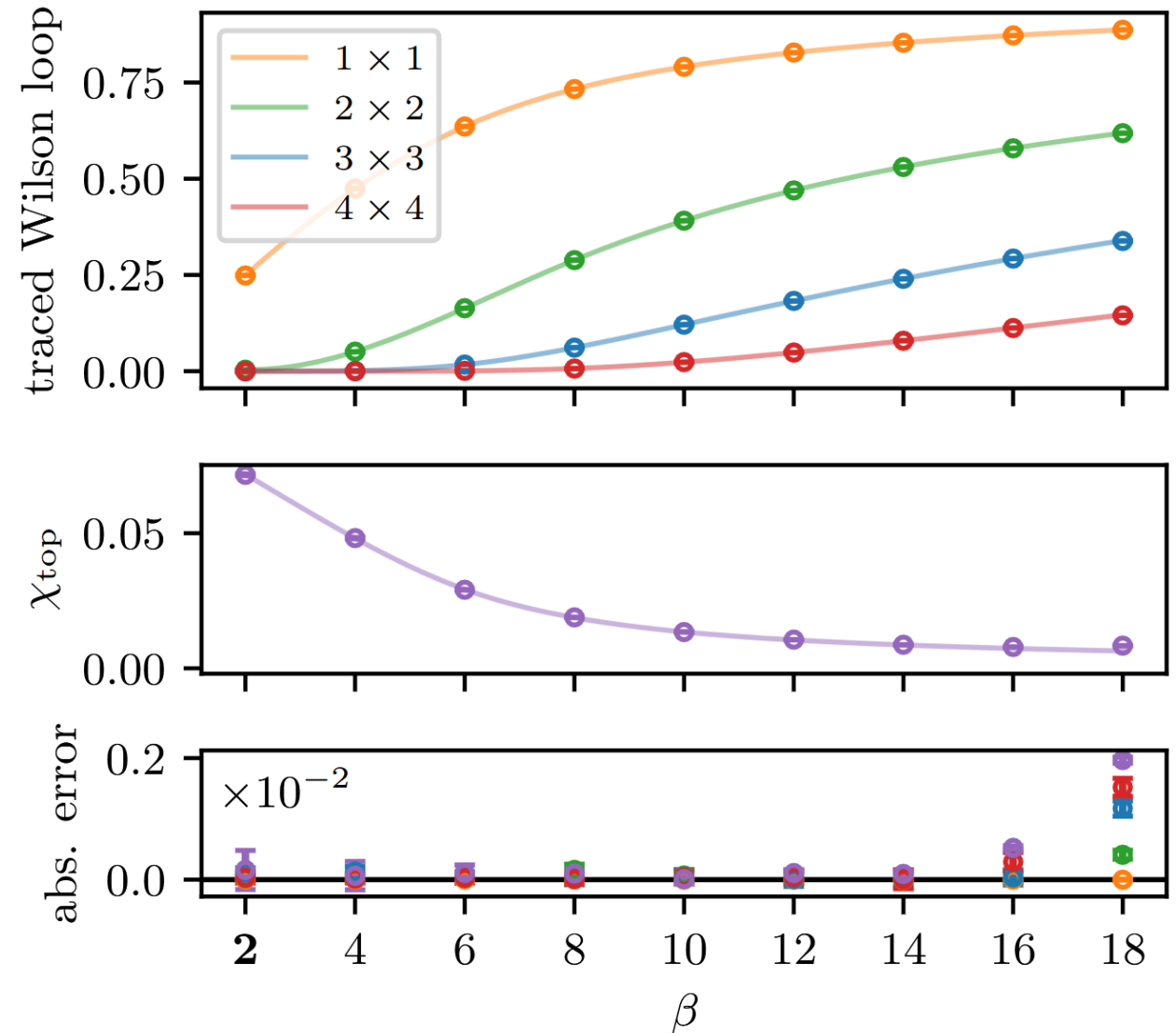
- Metropolis adjustment

Training setup for 2D U(2)

- 16×16 lattice
- Using HMC, generate 100.000 training configurations at $\beta = 2$
- L-CNN with two bilinear-convolution layers (15.266 trainable parameters)
- Generate 10.000 configurations at several β 's

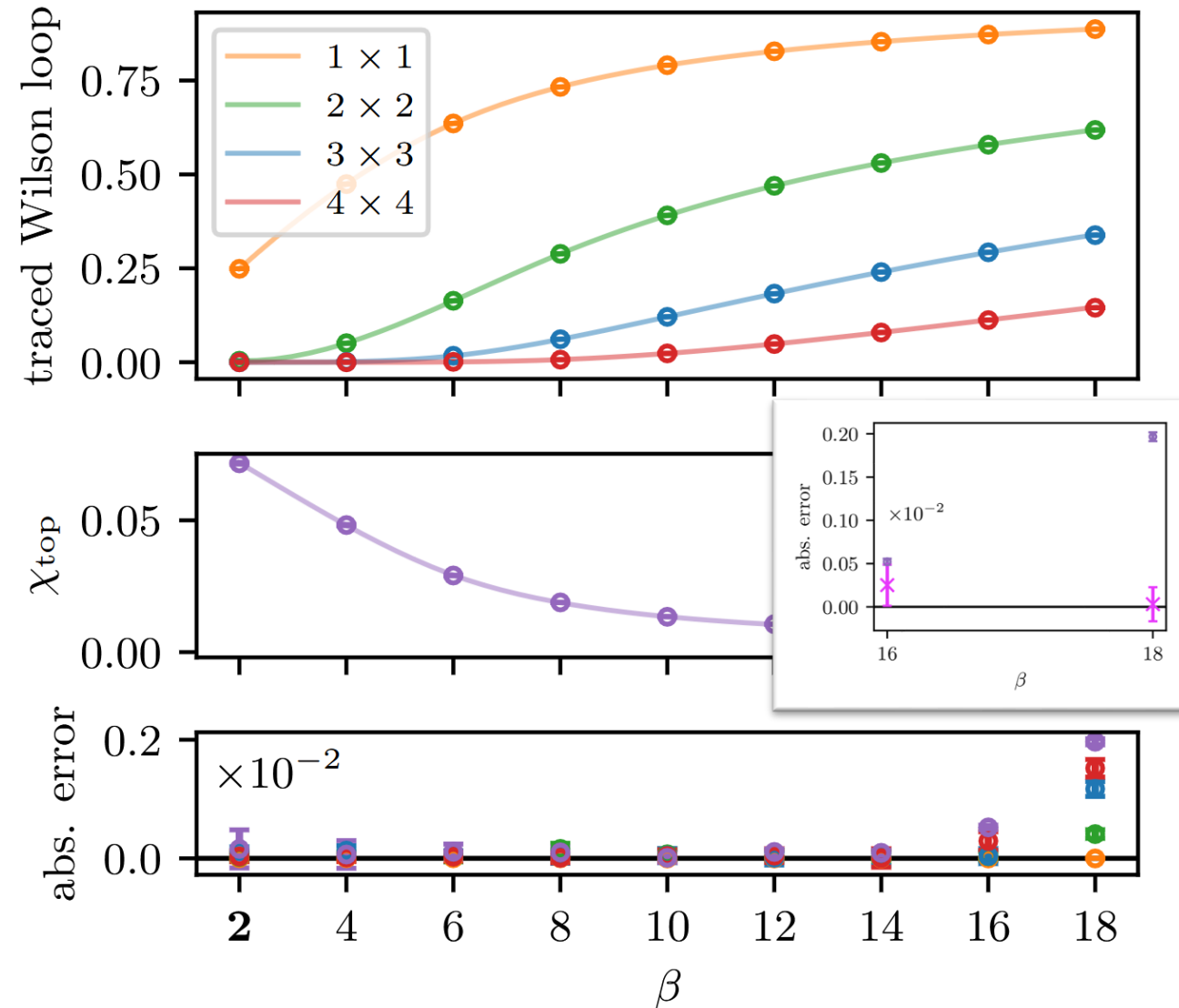
Results for 2D U(2)

- Extrapolation to higher β
- Fixed number of diffusion steps
- Excellent agreement up to $\beta = 14$
- Extrapolation to bigger lattices ($L = 64$) shows similar performance



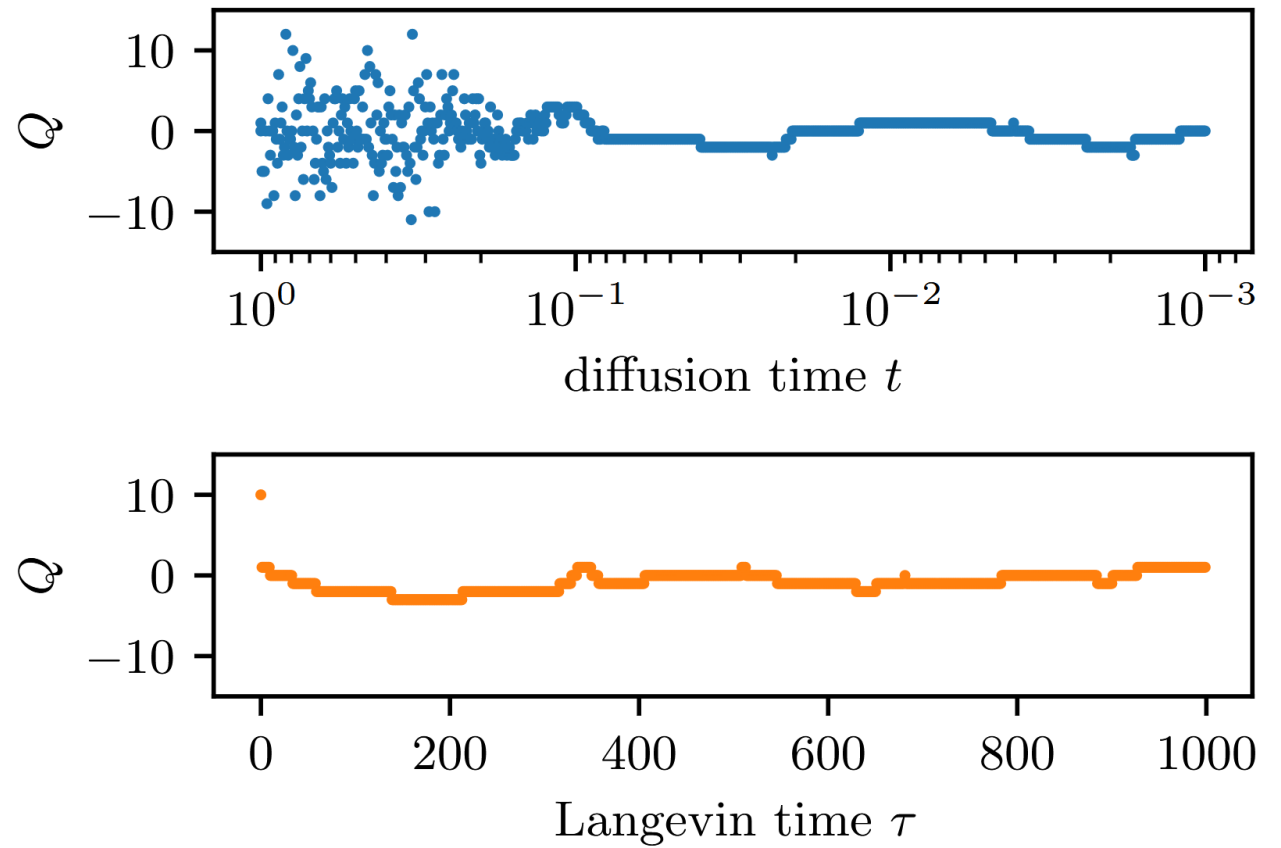
Results for 2D U(2)

- Extrapolation to higher β
- Fixed number of diffusion steps
- Excellent agreement up to $\beta = 14$
- Extrapolation to bigger lattices ($L = 64$) shows similar performance
- Perfect results at $\beta = 18$ require 10 times longer diffusion time



Results for 2D U(2)

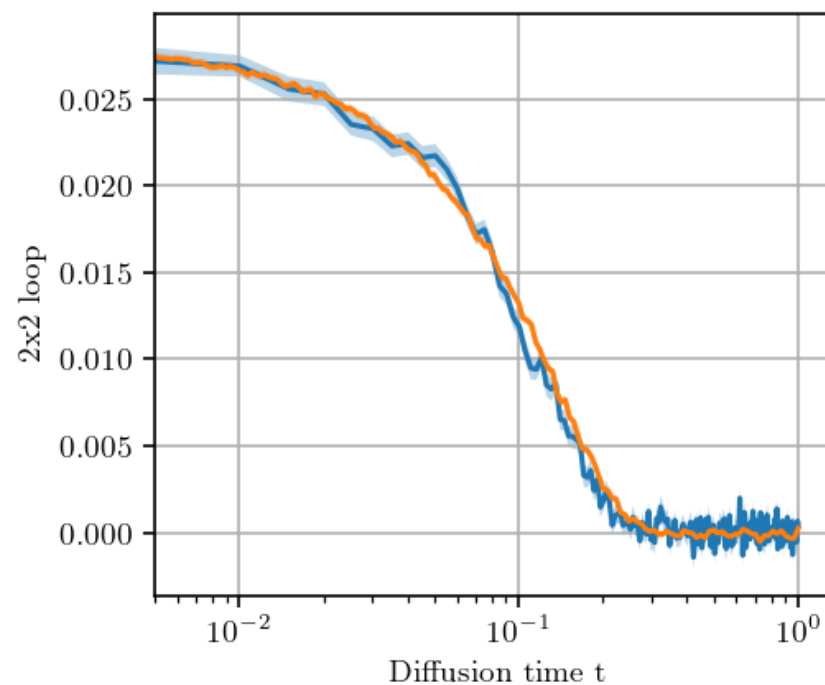
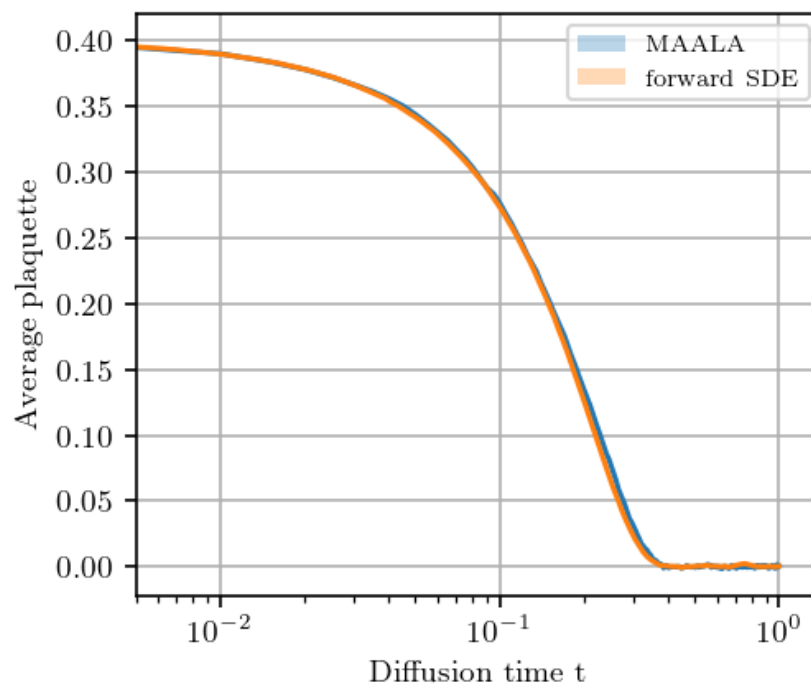
- Topological freezing
 - Markov Chain getting stuck in one topological sector
 - Appears towards high β
 - Hinders correct sampling over the whole phase space
- Annealing allows for more exploration in the initial sampling phase



First results for 4D SU(3)

- On a 4^4 lattice
- $\beta = 5$

PRELIMINARY



Summary

- Adaptation of diffusion models to non-Abelian lattice gauge theories
- Alternative sampling method MAALA
 - Extrapolation to different couplings
 - Metropolis adjustment
- Excellent extrapolation performance for 2D $U(2)$
- First promising results for 4D $SU(3)$

Thank you for your attention!

Backup slides

Introduction to diffusion models

- Stochastic differential equation describing the diffusion process:

$$dx = f(x, t)dt + g(t)dw$$

Drift term

Diffusion term

- Corresponding reverse process:

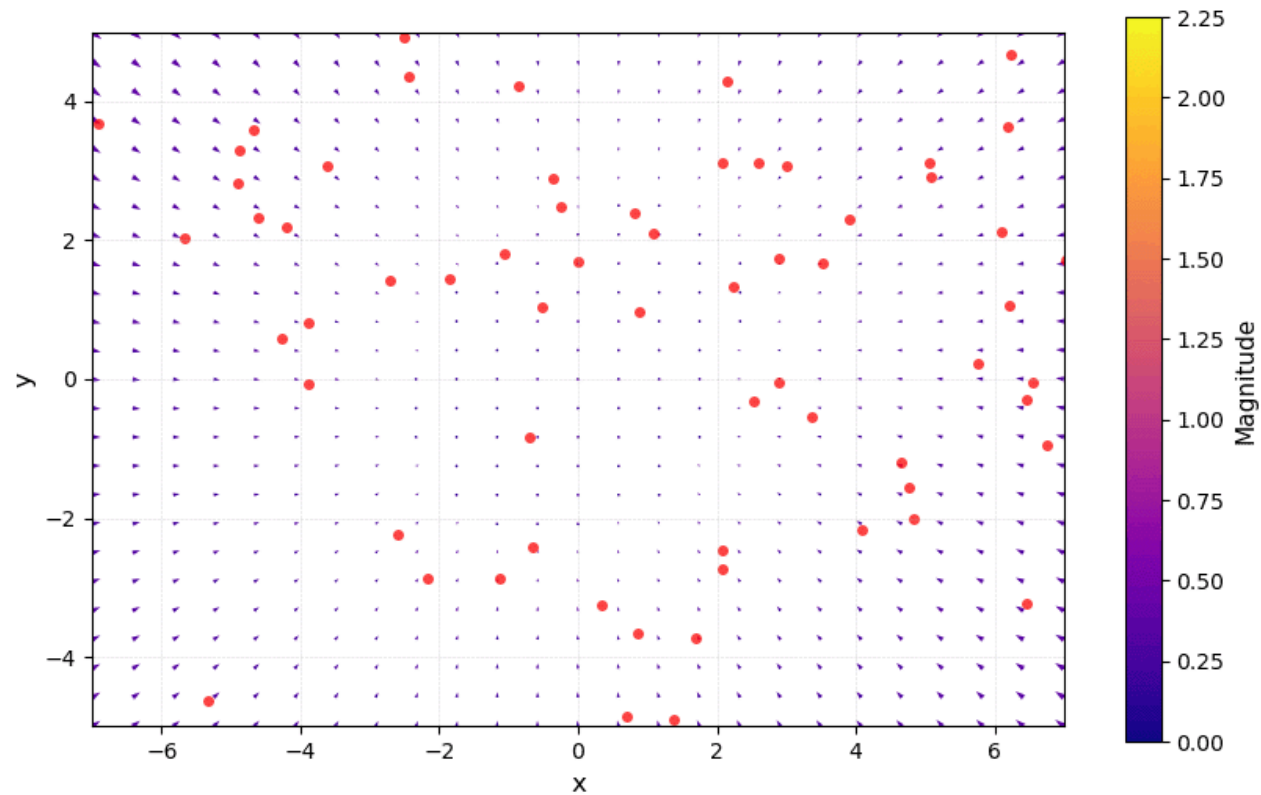
$$dx = [f(x, t) - g(t)^2 \nabla_x \log p_t(x)]dt + g(t)d\bar{w}$$

Score function
In general unknown,
will be learned by a
neural network.

Derivation via
Fokker-Planck
equation

Introduction to diffusion models

- Score function:



Non-Abelian lattice gauge theory

- Observables we look at:
- Average traced square loops:

$$W_{1 \times 1} = \left\langle \frac{1}{N} \text{Re Tr}(P_{x,\mu\nu}) \right\rangle$$

$n \times n$ loops
analogously

- Topological susceptibility (for 2D N=2):

$$\chi_{\text{top}} = \frac{\langle Q^2 \rangle}{V}, \quad Q = \frac{1}{2\pi} \sum_x \arg(\det P_{x,01})$$

Diffusion models for non-Abelian gauge theories

- Solving the SDE numerically:

$$U_{t+\epsilon} = \exp \left(i \sum_a T^a [\epsilon K^a(U_t) + \sqrt{\epsilon} g(t) \eta^a] \right) U_t$$

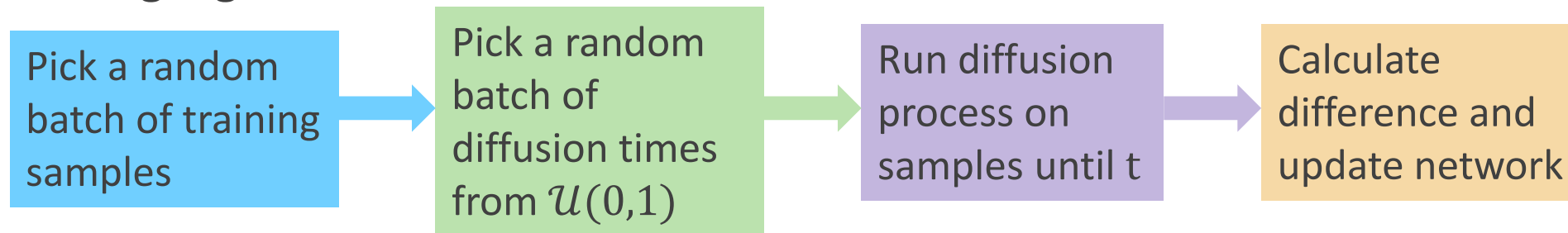
$$\eta^a \sim \mathcal{N}(0,1)$$

Diffusion models for non-Abelian gauge theories

- To actually solve the reverse SDE:
 - approximate the score function with a neural network $s_{\theta}^a(U, t)$
- Train with samples from the target distribution, minimizing the loss function

$$\int_0^T dt E_{p_0(U_0)} E_{p_{0t}(U_t|U_0)} \sum_a [s_{\theta}^a(U_t, t) - D^a \ln p_{0t}(U_t|U_0)]^2$$

- Training algorithm:



Diffusion models for non-Abelian gauge theories

- The actual SDE we use:

$$dU_t = i \sum_a T^a s^t dW_t^a \circ U_t, \quad s = 25$$

- Time interval $t \in [0,1]$
- Trick to solve it in one step:

$$U_t = \exp\left(i\sigma(t) \sum_a T^a \eta^a\right) U_0, \quad \sigma(t) = \sqrt{(s^{2t} - 1)/(2 \ln(s))}$$

- Analytical expression for the grad log of the transition probability:

$$D^a \ln p_{0t}(U_t|U_0) = -\frac{1}{\sigma(t)} \eta^a$$

Metropolis adjustment

- At t close to zero, apply Metropolis adjustment

- Langevin step proposes $\widetilde{U}_\tau^t$ given U_τ^t , accept with

$$p_{\text{accept}} = \min \left\{ 1, \frac{p(\widetilde{U}_\tau^t)}{p(U_\tau^t)} \frac{q(U_\tau^t | \widetilde{U}_\tau^t)}{q(\widetilde{U}_\tau^t | U_\tau^t)} \right\}$$

p ... target distribution

q ... transition probability of the diffusion process

$$\exp \left(S_E(U_\tau^t) - S_E(\widetilde{U}_\tau^t) \right)$$

$$\exp \left(\sum_{x,\mu} \left[\text{Tr}(\widehat{\eta_{x,\mu}^2}) - \text{Tr}(\widehat{\eta_{x,\mu}^{\prime 2}}) \right] \right)$$

$$\widehat{\eta_{x,\mu}} = \frac{1}{i\sqrt{2\epsilon}} \ln \left(\widetilde{U_{\tau;x,\mu}^t} U_{\tau;x,\mu}^{t\dagger} \right) - \sqrt{\frac{\epsilon}{2}} \sum_a T^a s_{\theta;x,\mu}(U_{\tau;x,\mu})$$

Lattice gauge-equivariant convolutional neural networks (LCNNs)

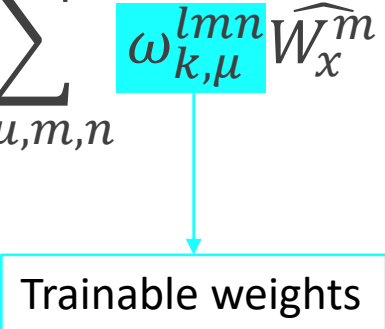
- Input: Links U and Plaquettes $P \rightarrow W$
 - Plus their complex conjugates and unit matrices
- Bilinear convolution layer:

$$\widehat{W}_x'^l = \sum_{k,\mu,m,n} \omega_{k,\mu}^{lmn} \widehat{W}_x^m U_{x,k\hat{\mu}} \widehat{W}_{x+k\hat{\mu}}^n U_{x,k\hat{\mu}}^\dagger$$

Ordered product of links starting at x and ending at $x + k\hat{\mu}$

Lattice gauge-equivariant convolutional neural networks (LCNNs)

- Input: Links U and Plaquettes $P \rightarrow W$
 - Plus their complex conjugates and unit matrices
- Bilinear convolution layer:

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A blue box highlights the weight term $\omega_{k,\mu}^{lmn}$ in the equation. A blue arrow points from this box down to a separate box labeled "Trainable weights".

Trainable weights

Lattice gauge-equivariant convolutional neural networks (LCNNs)

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 - Plus their complex conjugates and unit matrices
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$$\widehat{W}_x^{\prime l} = \sum_{k, \mu, m, n} \omega_{k, \mu}^{lmn} \widehat{W}_x^m U_{x, k\hat{\mu}} \widehat{W}_{x+k\hat{\mu}}^n U_{x, k\hat{\mu}}^\dagger$$

$l \dots$ output channels

Lattice gauge-equivariant convolutional neural networks (LCNNs)

- Input: Links U and Plaquettes $P \rightarrow W$
 - Plus their complex conjugates and unit matrices
- Bilinear convolution layer:

$$\widehat{W}_x'^l = \sum_{k,\mu,\mathbf{m},\mathbf{n}} \omega_{k,\mu}^{l\mathbf{m}\mathbf{n}} \widehat{W}_x^{\mathbf{m}} U_{x,k\hat{\mu}} \widehat{W}_{x+k\hat{\mu}}^{\mathbf{n}} U_{x,k\hat{\mu}}^\dagger$$

$\mathbf{m}, \mathbf{n} \dots$ input channels

Lattice gauge-equivariant convolutional neural networks (LCNNs)

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 - Plus their complex conjugates and unit matrices
- Bilinear convolution layer:

$$\widehat{W}_x'^l = \sum_{\mathbf{k}, \mu, m, n} \omega_{\mathbf{k}, \mu}^{lmn} \widehat{W}_x^m U_{x, \mathbf{k} \hat{\mu}} \widehat{W}_{x + \mathbf{k} \hat{\mu}}^n U_{x, \mathbf{k} \hat{\mu}}^\dagger$$

$$-(K - 1) \leq k \leq K - 1$$

$K \dots$ Kernel size

Lattice gauge-equivariant convolutional neural networks (LCNNs)

- Input: Links U and Plaquettes $P \rightarrow W$
 - Plus their complex conjugates and unit matrices
- Bilinear convolution layer:

$$\widehat{W}_x'^l = \sum_{k, \mu, m, n} \omega_{k, \mu}^{lmn} \widehat{W}_x^m U_{x, k\hat{\mu}} \widehat{W}_{x+k\hat{\mu}}^n U_{x, k\hat{\mu}}^\dagger$$

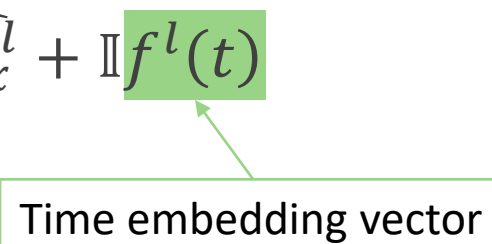
μ ... Lattice direction

Lattice gauge-equivariant convolutional neural networks (LCNNs)

- Time embedding by random Fourier features

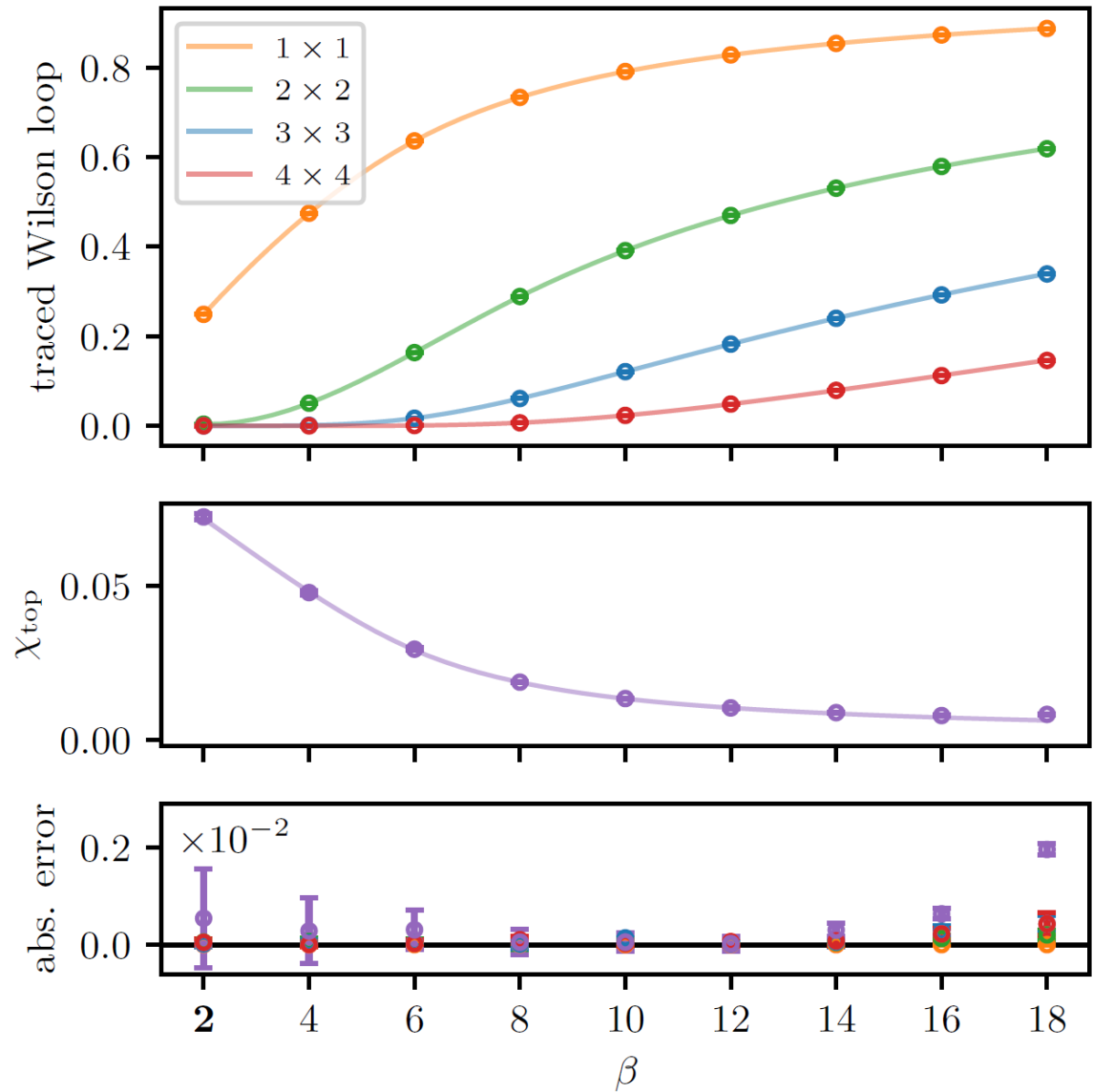
$$\widehat{W}_x'^l = \widehat{W}_x^l + \mathbb{I} f^l(t)$$

Time embedding vector



Results for 2D U(2)

- Extrapolation to a larger lattice
- Model trained at $L = 16$
- Generating configurations at $L = 64$



Results for 2D U(2)

- Metropolis acceptance rates

