

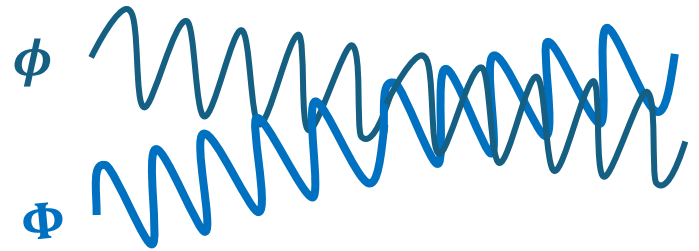
Mixed Wave Dark Matter and Implications in Neutrino Physics

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Based on arXiv: 2604.14480 (Accepted in PRD)

Collaborated with Hye-Sung Lee (KAIST)



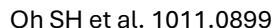
with Ultra-light Mass Scale



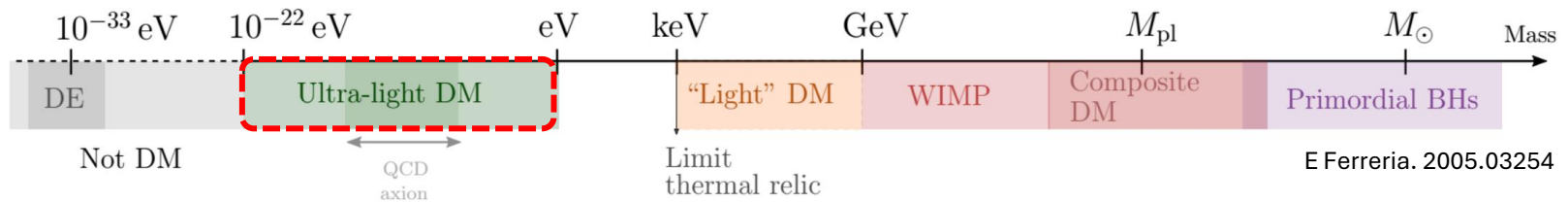
$$m_a \sim \frac{\Lambda^2}{f_a} \sim \frac{m_\pi f_\pi}{f_a}$$

QCD axion or ALP

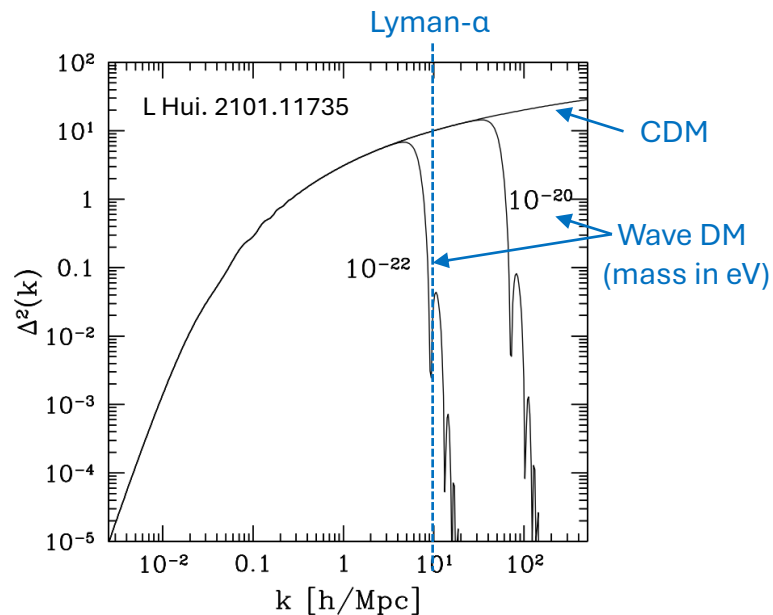
: Motivations of ultra-light mass of DM



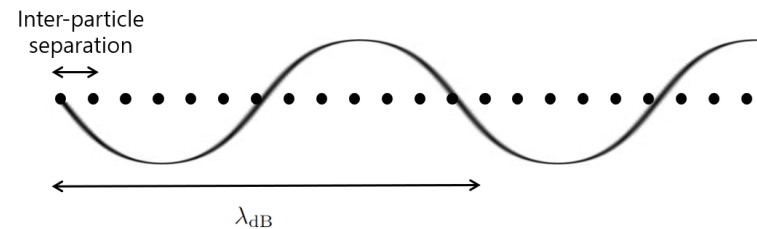
Wave Dark Matter with Ultra-light Mass Scale



Lower bound from structure formation



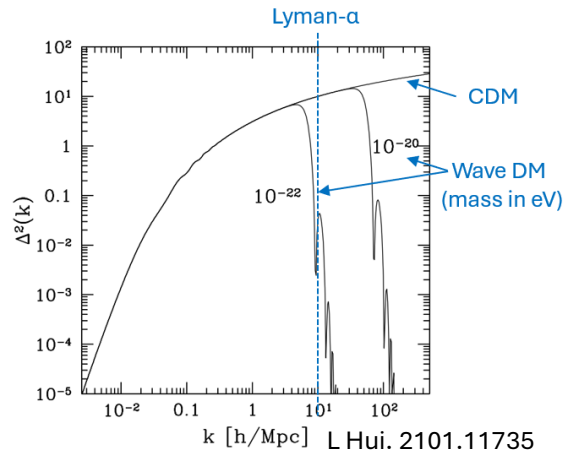
Upper bound from wave-like behavior



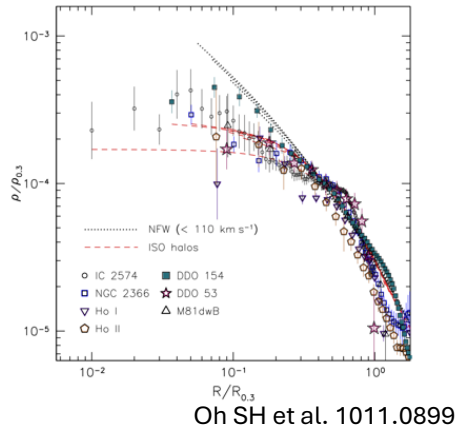
$$\mathcal{N} = n \lambda_{\text{dB}}^3 = \left(\frac{30 \text{ eV}}{m} \right)^4 \left(\frac{10^{-4}}{v} \right)^3 \gg 1$$

Observational / Experimental Probes of Wave Dark Matter

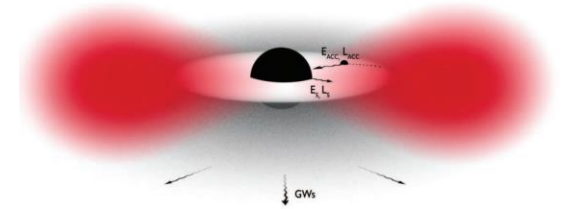
Matter power spectrum



Galactic structure

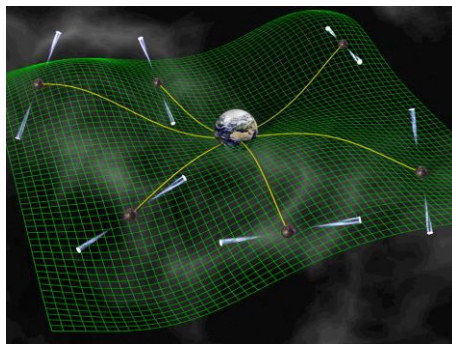


BH superradiance



R. Brito et al. 1501.06570

Pulsar time array



D. Champion/Max Planck Institute

Modulation signal

$$\alpha(t) = \alpha \left[1 + d_e \kappa \phi_0 \cos(m_\phi t) \right]$$

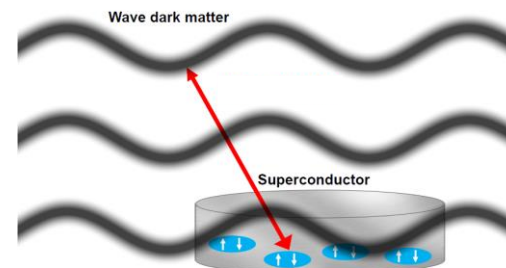
Sr Clock

$$f_{Sr} \propto \alpha^{2+\epsilon} m_e c^2$$

$$\epsilon = 0.06$$

$$\frac{\delta(f_{Sr}/f_c)}{f_{Sr}/f_c} = 1.06 \frac{\delta\alpha}{\alpha}$$

CJ. Kennedy et al. 2008.08773



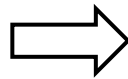
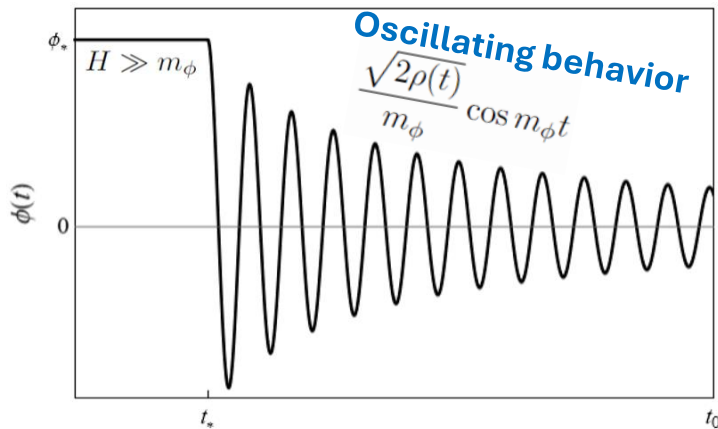
YL ChoeJo, YK, HS Lee, J Lee, J Yi.
2509.22892

Oscillating Wave Dark Matter to other sectors

$$S = \int d^4x \sqrt{-g} \left[-\frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - \frac{1}{2} m_\phi^2 \phi^2 \right]$$

$$\ddot{\phi} + 3H\dot{\phi} + m_\phi^2 \phi = 0$$

$$\phi(t) = \begin{cases} \phi_* & H > m_\phi \\ \frac{\sqrt{2\rho(t)}}{m_\phi} \cos m_\phi t & H < m_\phi \end{cases}$$



Coupled to gauge boson

A Arvanitaki et al. 1405.2925
K Van Tilburg et al. 1503.06886
J Guo et al. 2206.14221

$$\mathcal{L} \supset -\frac{1}{4} d_e \kappa \phi F_{\mu\nu} F^{\mu\nu}$$

$$\alpha(t) = \alpha \left[1 + d_e \kappa \phi_0 \cos(m_\phi t) \right]$$

$$\mathcal{L} \supset -\frac{1}{4} d_g \kappa \phi G_{\mu\nu}^a G^{a\mu\nu}$$

$$\Lambda_{\text{QCD}}(t) = \Lambda_{\text{QCD}} \left[1 + d_g \kappa \phi_0 \cos(m_\phi t) \right]$$

Coupled to SM fermion

A Arvanitaki et al. 1405.2925
M Filzinger et al. 2312.13723
YL ChoeJo, YK, HS Lee, J Lee, J Yi.
2509.22892

$$\mathcal{L} \supset -d_{m_e} \kappa \phi m_e \bar{e} e$$

$$m_e(t) = m_e \left[1 + d_{m_e} \kappa \phi_0 \cos(m_\phi t) \right]$$

$$\mathcal{L} \supset -d_{m_q} \kappa \phi m_q \bar{q} q$$

$$m_q(t) = m_q \left[1 + d_{m_q} \kappa \phi_0 \cos(m_\phi t) \right]$$

Coupled to neutrino

A Berlin. 1608.01307
G Krnjaic et al. 1705.06740
YL ChoeJo, YK, HS Lee. 2305.16900

$$\mathcal{L} \supset -g \phi \bar{N} N$$

$$M_N(t) = M_N + g \phi_0 \cos(m_\phi t)$$

Single vs. Mixed Wave Dark Matter

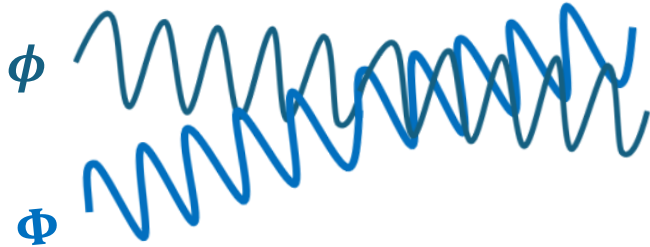
Single wave dark matter



Harmonic oscillation

$$\phi(t) = \phi_a \cos m_\phi t$$

Mixed wave dark matter



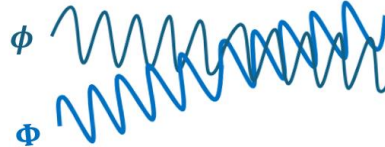
Non-trivial time-dependent structure
(non-monochromatic, symmetry-broken...)

$$\phi(t) = \phi_0(t) + \phi_a(t)f(m_\phi t)$$

Mixed Wave Dark Matter and parametric structure

$$V(\phi, \Phi) = \frac{1}{2}m^2\phi^2 + \frac{1}{2}M^2\Phi^2 + \boxed{\frac{1}{2}\kappa\phi^2\Phi^2}$$

Wave DM mixing



For $M^2 \gg \kappa\phi^2$

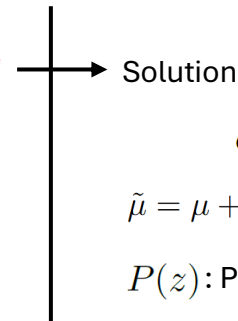
$$\Phi(t) \simeq \Phi_0 \sin(Mt)$$

EoM of ϕ

$$\ddot{\phi} + 3H\dot{\phi} + (m^2 + \kappa\Phi^2)\phi = 0$$

$$\boxed{\phi''(z) + [A - 2q \cos(2z)]\phi(z) = 0}$$

Mathieu-type equation



Solution

$$\phi(z) = e^{\tilde{\mu}(A,q)z} P(z)$$

$\tilde{\mu} = \mu + i\nu$: Mathieu characteristic

$P(z)$: Periodic function

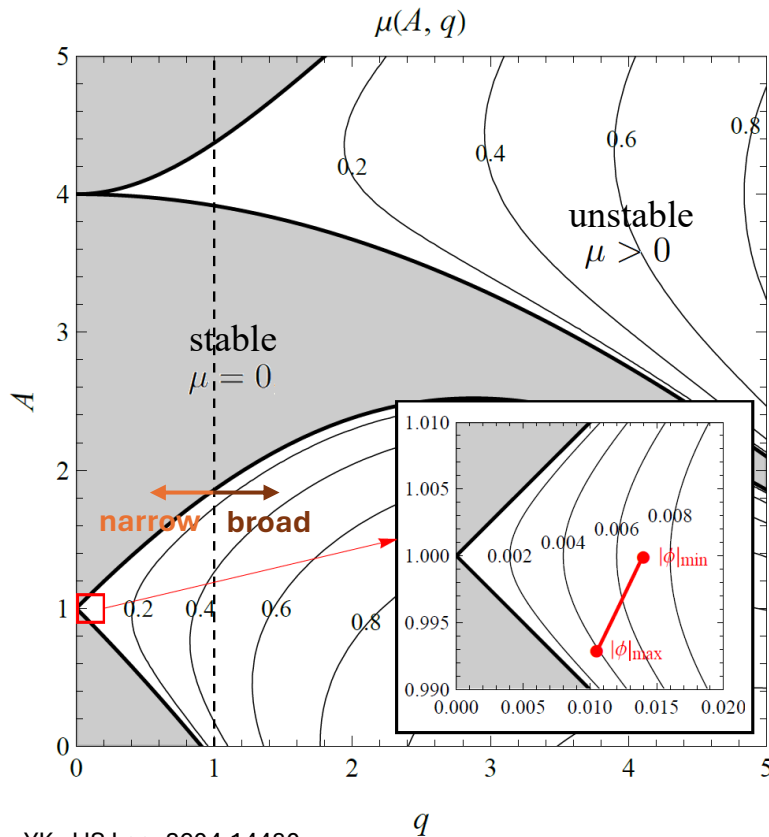
where $z \equiv Mt$ $' \equiv d/dz$

$$q \equiv \frac{\kappa\Phi_0^2}{4M^2}, \quad A \equiv \frac{m^2}{M^2} + 2q$$

Mixed Wave Dark Matter and parametric structure

$$\text{EoM} \quad \phi''(z) + [A - 2q \cos(2z)]\phi(z) = 0$$

$$\text{Sol} \quad \phi(z) = e^{\tilde{\mu}(A,q)z} P(z) \quad \tilde{\mu} = \mu + i\nu$$



YK, HS Lee. 2604.14480

Broad resonance $q \gg 1$

- Instability regions are wide
- μ is large
- **Exponential growth** $|\phi(z)| \propto \exp(\mu z)$

Narrow resonance $q \ll 1$

- Instability regions are narrow

$$A \simeq n^2 \quad (n = 1, 2, \dots)$$

$$|A - 1| \lesssim q \quad \text{for } n = 1$$

- μ is small

$$\mu \simeq \frac{1}{2} \sqrt{q^2 - (A - 1)^2} \quad \text{for } n = 1$$

- **Growth is suppressed**

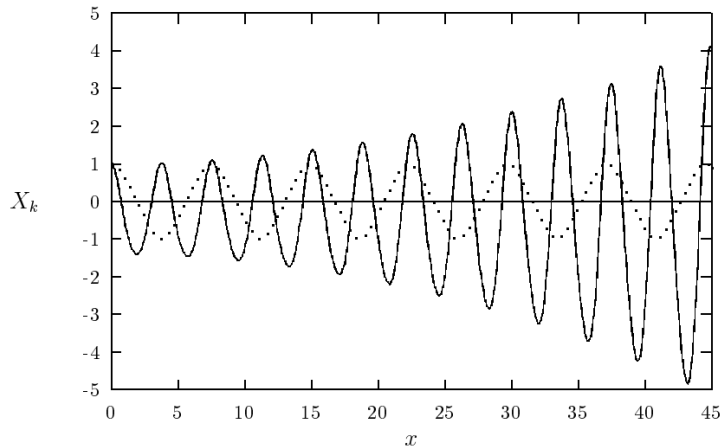
Mixed Wave Dark Matter and parametric structure

Broad resonance $q \gg 1$

- Parametric resonance
- i.e., Preheating after inflation
- **Amplified production** of particle

$$X_k(t) = a(t)\chi_k(t)$$

x : dimensionless conformal time



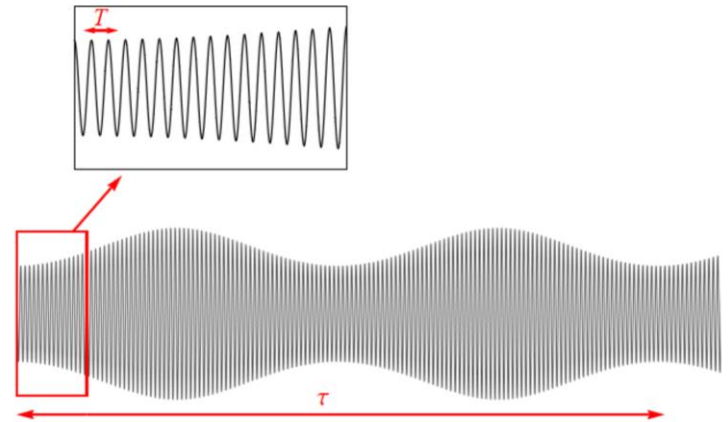
P Greene et al. (1997)

Narrow resonance $q \ll 1$

$$V(\phi, \Phi) = \frac{1}{2}m^2\phi^2 + \frac{1}{2}M^2\Phi^2 + \frac{1}{2}\kappa\phi^2\Phi^2 \quad q \equiv \frac{\kappa\Phi_0^2}{4M^2}$$

Subdominant
mixing term

Amplification is insufficient → Growth stops



YK, HS Lee. 2604.14480

Mixed Wave Dark Matter to Wave-envelope DM

$e^{\mu z}$ with $z = Mt$

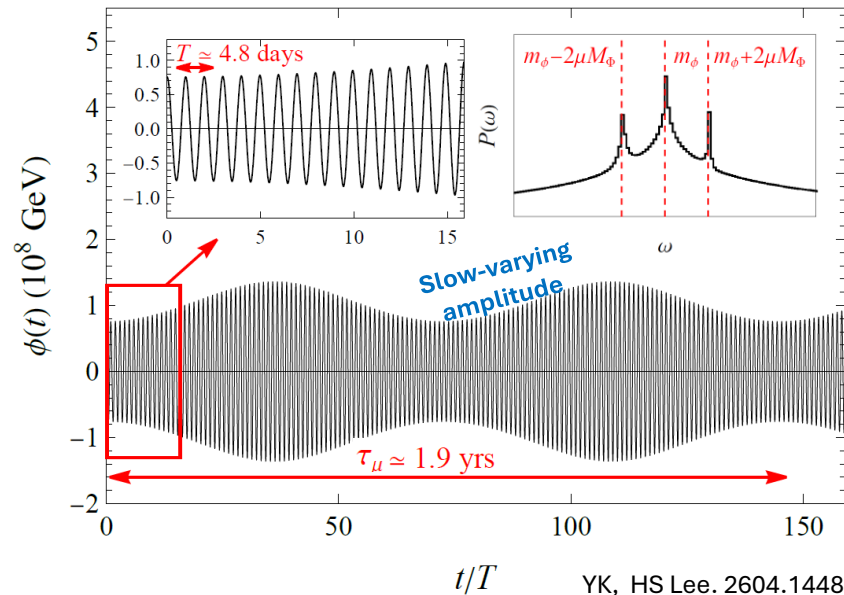
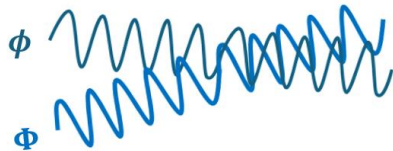
New time scale $\tau_\mu \simeq \frac{2\pi}{\mu M} \gg T \simeq \frac{2\pi}{M}$ for $\mu \ll 1$

 Slow oscillation from mixing Wave DM oscillation

Narrow resonance

"Wave-envelope DM"

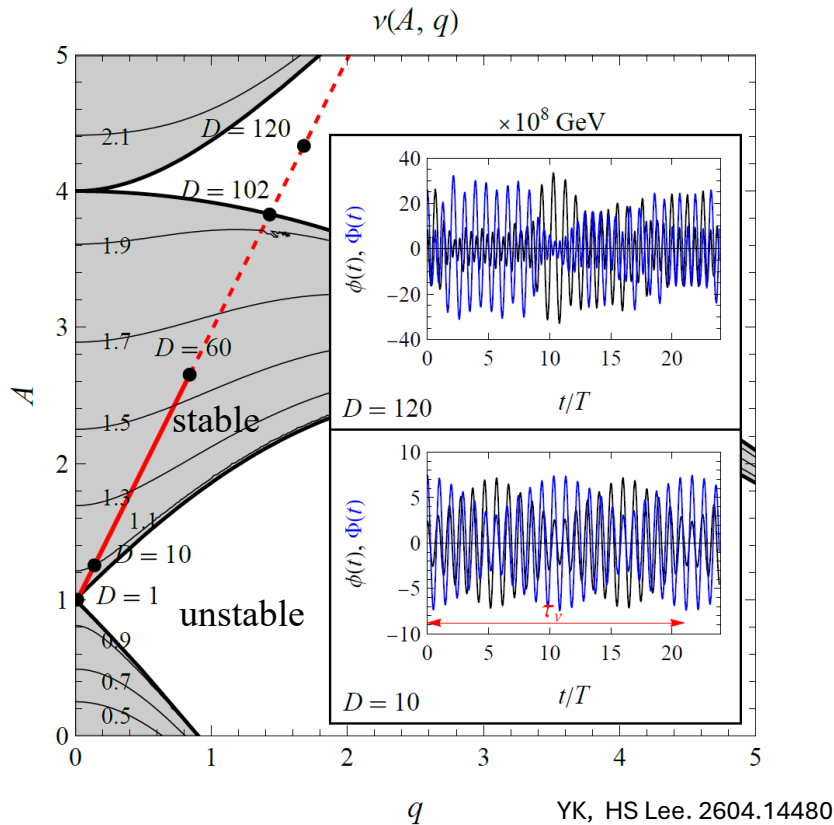
Mixed wave DM



Mixed Wave Dark Matter to Wave-envelope DM

EoM $\phi''(z) + [A - 2q \cos(2z)]\phi(z) = 0$

Sol $\phi(z) = e^{\tilde{\mu}(A,q)z} P(z)$ $\tilde{\mu} = \mu + i\nu$



Stable region

$$\mu = 0, \nu = n + \delta$$

$$\tau_\nu \simeq \frac{4\pi}{|\nu - n|M}$$

Unstable region

$$\mu > 0, \nu = 0$$

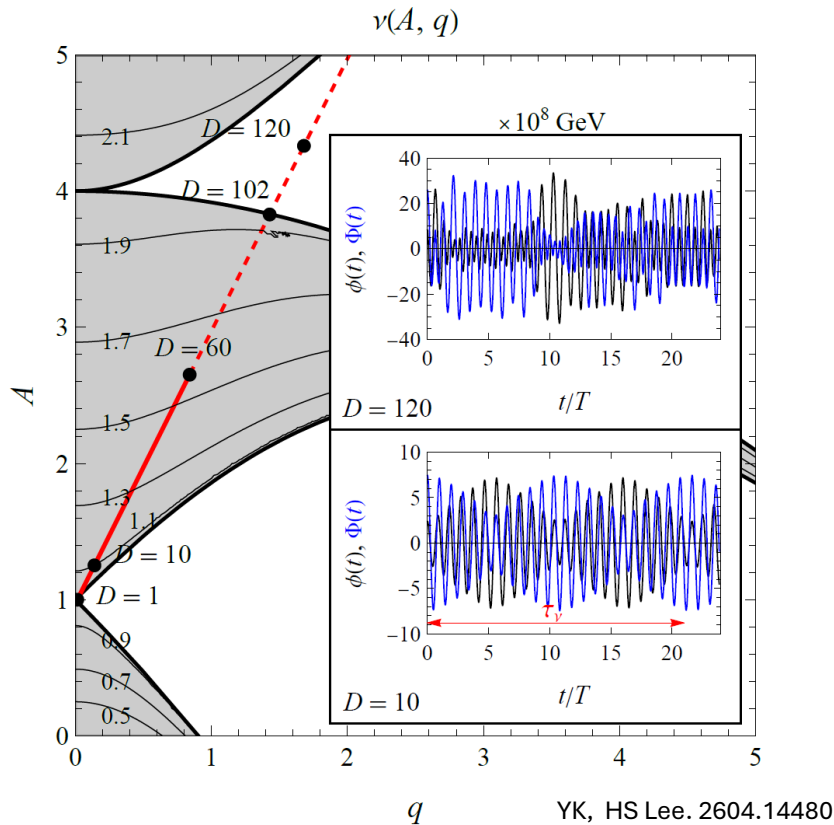
$$\tau_\mu \simeq \frac{2\pi}{\mu M}$$

Mixed Wave Dark Matter to Wave-envelope DM

Dense environment with enhancement factor D

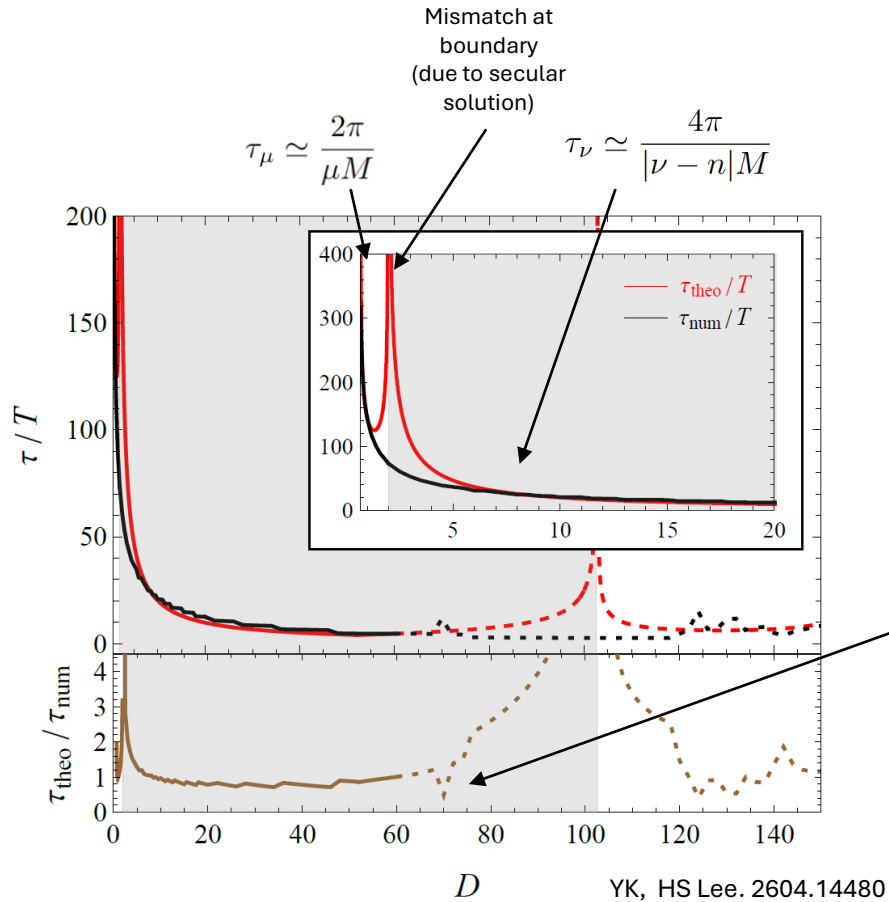
$$q \equiv \frac{\kappa \Phi_0^2}{4M^2}, \quad A \equiv \frac{m^2}{M^2} + 2q$$

$$q(D) = Dq_0, \quad A(D) = A_0 + 2q_0(D - 1)$$



- $D = 1$ 1st instability band, narrow resonance
- $D = 2$ Enter stable region
- $D = 10$ Example plot 1 (envelope in stable region)
- $D = 60$ Break down of Mathieu approximation
Due to backreaction
- $D = 102$ Get out stable region
- $D = 120$ Example plot 2 (irregular pattern)

Mixed Wave Dark Matter to Wave-envelope DM



Theoretical vs Numerical envelope timescales

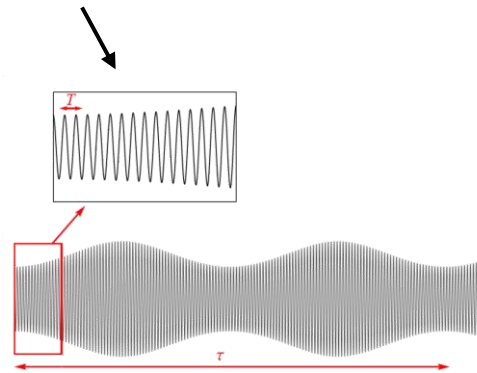
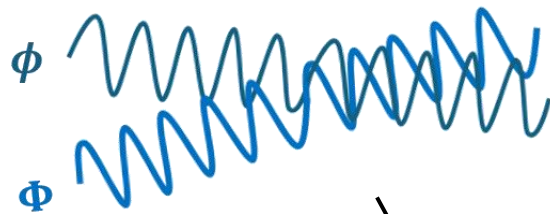
$$\tau_{\text{theo}} \begin{cases} \tau_{\mu} \simeq \frac{2\pi}{\mu M} \\ \tau_{\nu} \simeq \frac{4\pi}{|\nu - n|M} \end{cases}$$

$$\begin{aligned} \ddot{\phi} + 3H\dot{\phi} + (m^2 + \kappa\Phi^2)\phi &= 0 \\ \ddot{\Phi} + 3H\dot{\Phi} + (M^2 + \kappa\phi^2)\Phi &= 0 \end{aligned}$$

$\downarrow$
 τ_{num}

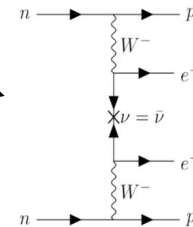
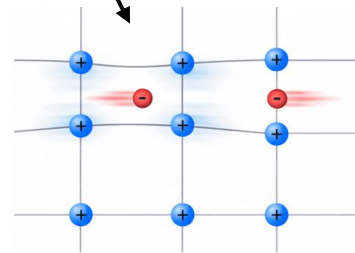
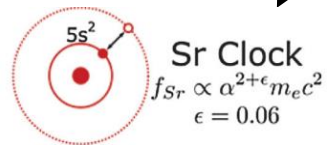
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- Due to backreaction

Mixed Wave Dark Matter



“Wave-envelope DM”

Beyond harmonic modulation

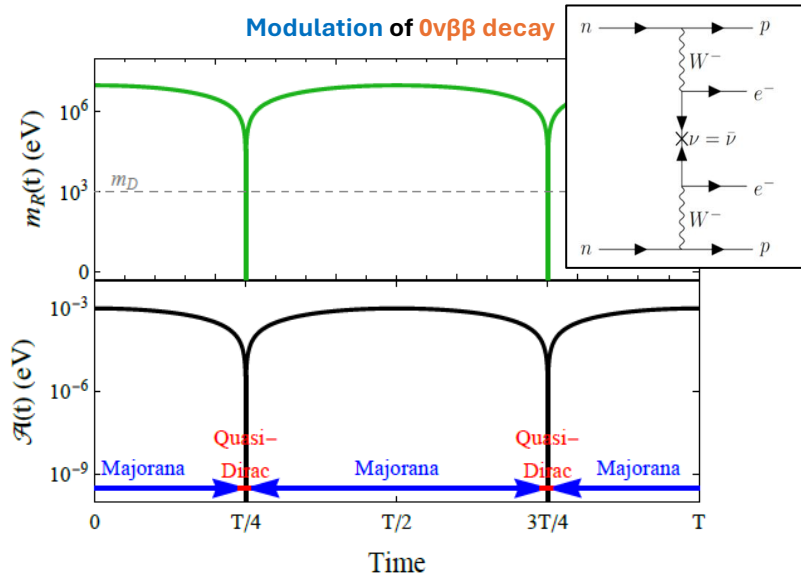


Implications in Neutrino Physics $\mathcal{L} \supset -\frac{1}{2}g\phi\bar{\nu}_M\nu_M$

Single wave dark matter

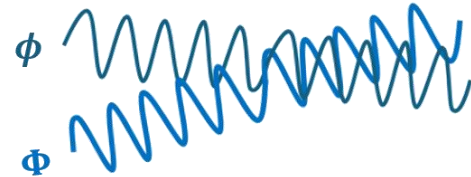


$$m_R(t) = m_{R,0} \cos m_\phi t$$

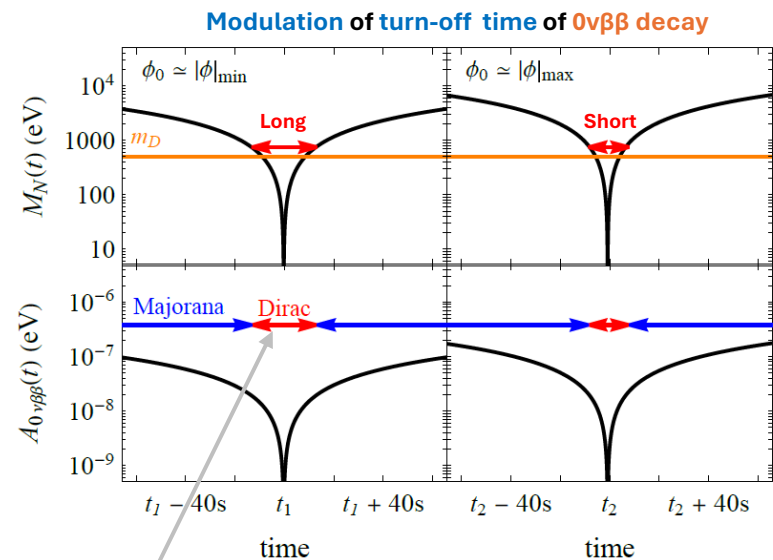


YL ChoeJo, YK, HS Lee. 2305.16900

Mixed wave dark matter



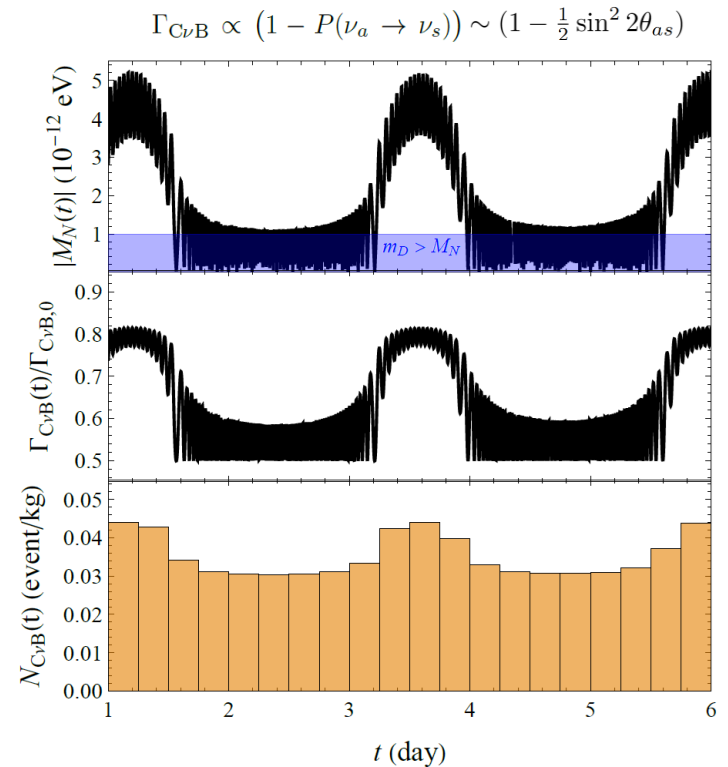
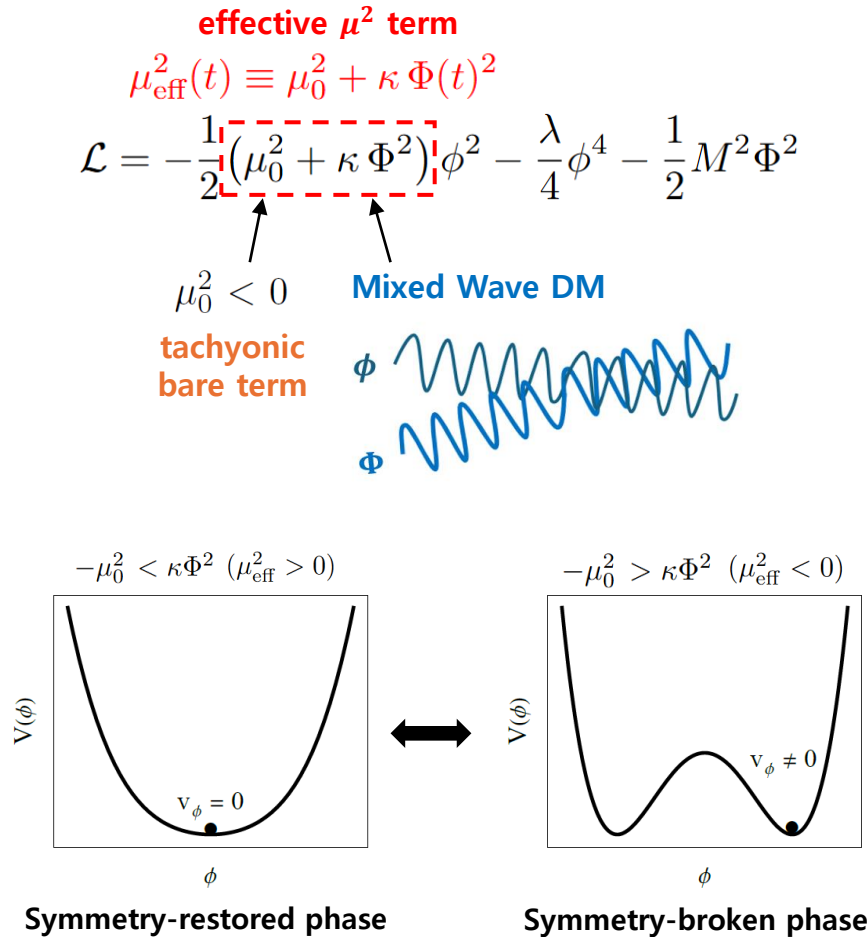
$$m_R(t) = m_{R,0}(t) \cos m_\phi(t)t$$



$$\Delta T_{\text{Dirac}} \simeq \frac{T}{\pi} \sin^{-1} \left(\frac{m_D}{g\phi_0} \right)$$

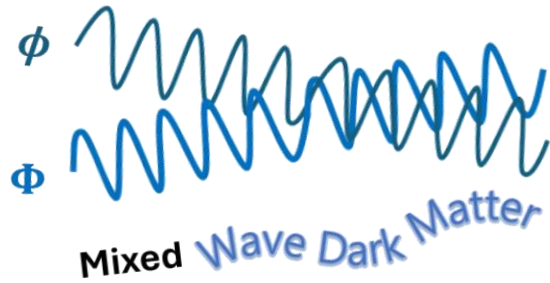
YK, HS Lee. 2604.14480

Implications in Neutrino Physics $\mathcal{L} \supset -\frac{1}{2}g\phi\bar{\nu}_M\nu_M$

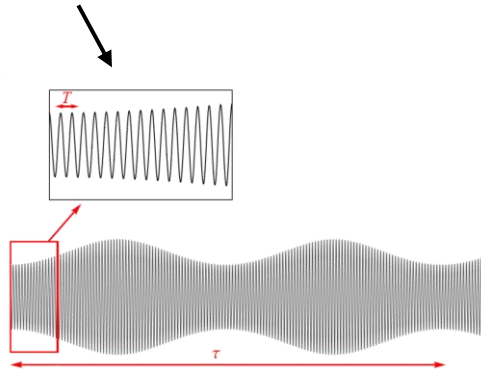


YK, HS Lee. 2503.04949

Summary



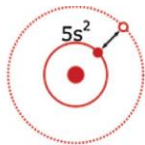
- **Mixed wave DM** can lead to the non-trivial time-dependent structure.



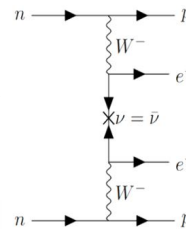
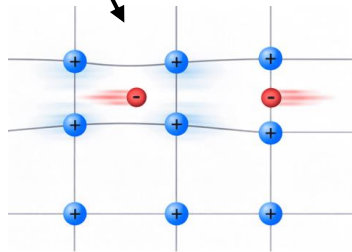
"Wave-envelope DM"

- **Wave-envelope DM** can arise with slow-varying amplitude from parametric dynamics.

Beyond harmonic modulation



Sr Clock
 $f_{Sr} \propto \alpha^{2+\epsilon} m_e c^2$
 $\epsilon = 0.06$



- This can induce **modulation signals** in various observables.

Thank you!

Backup slides

Mixed Wave Dark Matter to Wave-envelope DM

$$V(\phi, \Phi) = \frac{1}{2}m^2\phi^2 + \frac{1}{2}M^2\Phi^2 + \frac{1}{2}\kappa\phi^2\Phi^2$$

$$M^2 \gg \kappa\phi^2$$

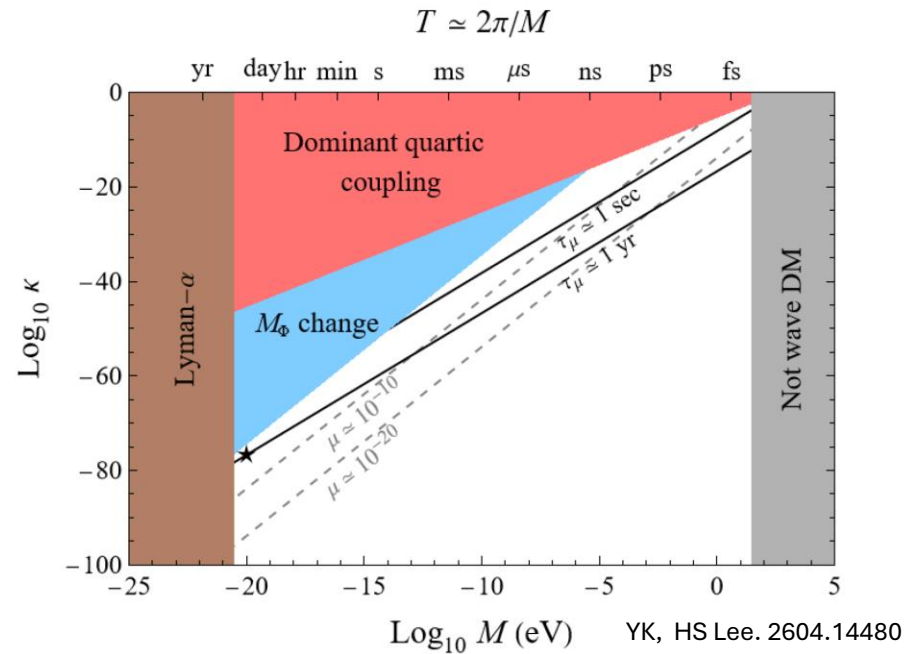
$$q \equiv \frac{\kappa\Phi_0^2}{4M^2}, \quad A \equiv \frac{m^2}{M^2} + 2q$$

$$\mu \simeq \frac{1}{2}\sqrt{q^2 - (A-1)^2} \quad \text{for } A \simeq 1$$

$$\tau_\mu \simeq \frac{2\pi}{\mu M} \gg T \simeq \frac{2\pi}{M}$$

Slow oscillation
from mixing

Wave DM
oscillation



Mixed Wave Dark Matter to Wave-envelope DM

EoM $\phi''(z) + [A - 2q \cos(2z)]\phi(z) = 0$

Sol $\phi(z) = e^{\tilde{\mu}(A,q)z} P(z)$
period π

Stable region

$\mu = 0, \nu = n + \delta \quad |\delta| \ll 1$

$$\phi = e^{i\nu z} \sum_{l=-\infty}^{\infty} c_l e^{2ilz}$$

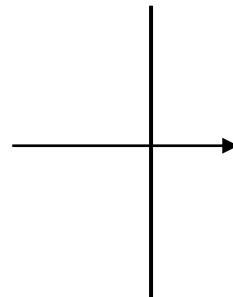
$$-(\nu + 2l)^2 c_l + A c_l - q(c_{l-1} + c_{l+1}) = 0$$

Near the n -th band with $(A, q) \simeq (n^2, 0)$

Dominant modes $l = 0$ and $l = -n$

Frequencies ν and $\nu - 2n$

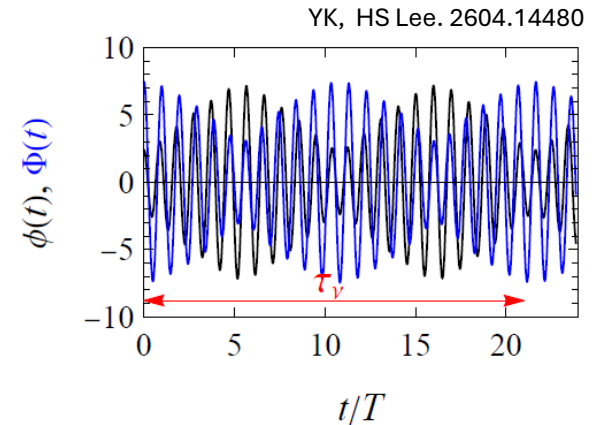
$\nu = n + \delta \quad n + \delta$ and $n - \delta$



adjacent to the n -th band

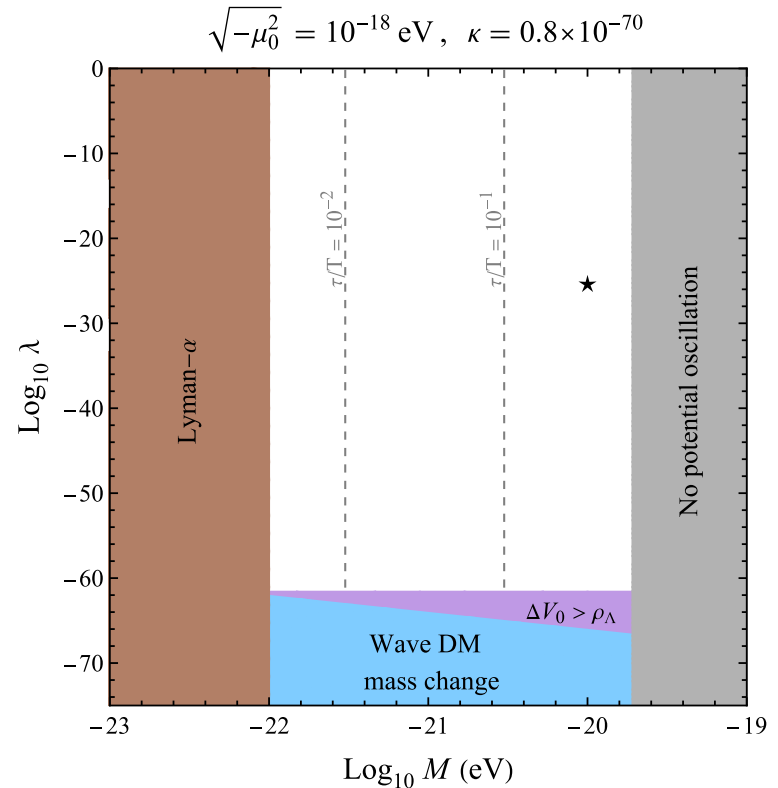
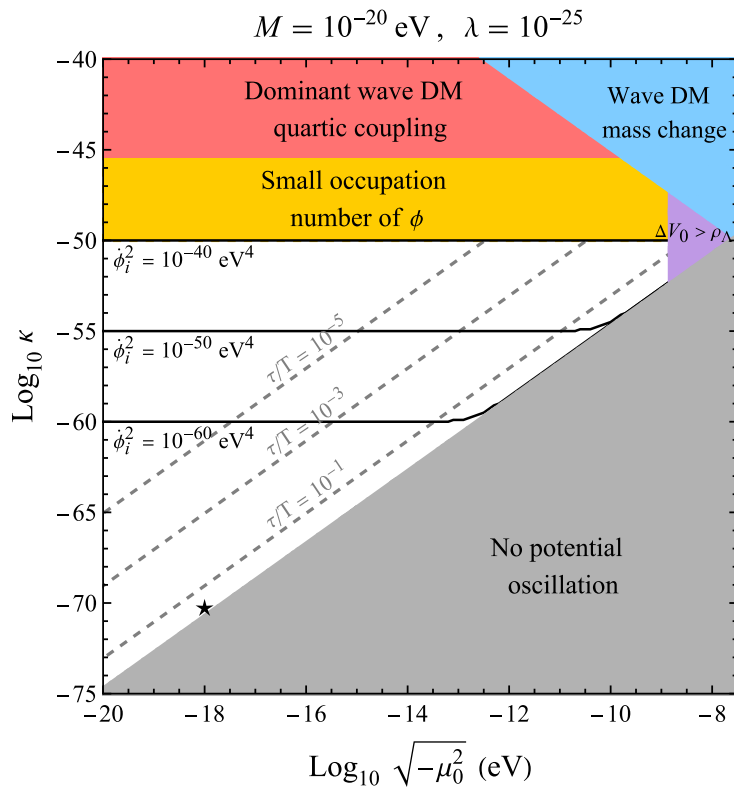
$$\tau_\nu \simeq \frac{4\pi}{|\nu - n|M} \gg T \simeq \frac{2\pi}{M}$$

$$|\delta| = |\nu - n| \ll 1$$



Potential oscillation constraint

$$\mathcal{L} = -\frac{1}{2}(\mu_0^2 + \kappa \Phi^2) \phi^2 - \frac{\lambda}{4} \phi^4 - \frac{1}{2} M^2 \Phi^2$$



τ : time interval of symmetry-broken phase in one period

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Implications in Neutrino Physics

Dirac



vs

Majorana



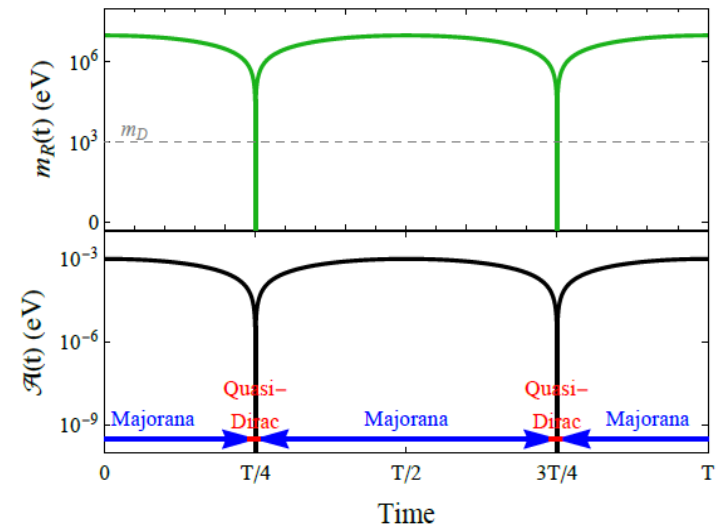
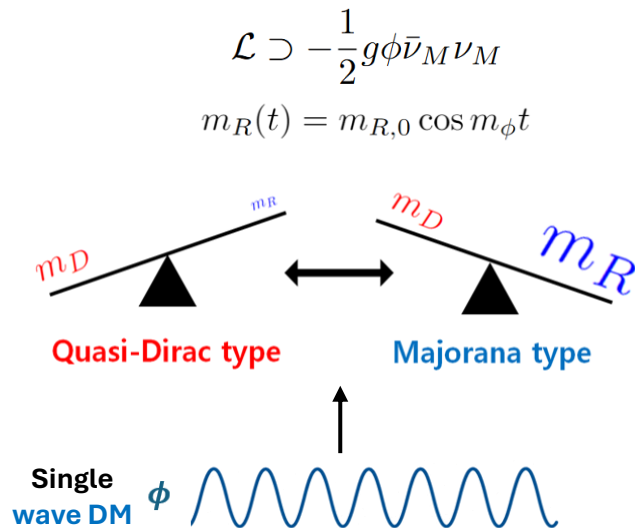
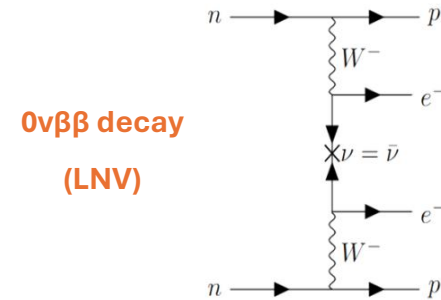
Dirac mass

$\mathcal{L} = -m_D \bar{\nu}_D \nu_D$

Majorana mass

$\mathcal{L} = -\frac{1}{2} m_R \bar{\nu}_M \nu_M$

$\nu_D = \nu_L + \nu_R$
 $\nu_M = \nu_R + \nu_R^c$



YL ChoeJo, YK, HS Lee. 2305.16900

Mixed Wave Dark Matter Scenario II: Potential Oscillation

YK, HS Lee. 2503.04949

effective μ^2 term

$$\mu_{\text{eff}}^2(t) \equiv \mu_0^2 + \kappa \Phi(t)^2$$

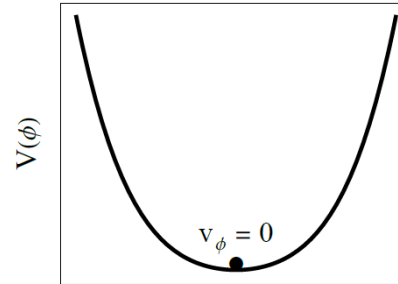
$$\mathcal{L} = -\frac{1}{2}(\mu_0^2 + \kappa \Phi^2)\phi^2 - \frac{\lambda}{4}\phi^4 - \frac{1}{2}M^2\Phi^2$$

$\mu_0^2 < 0$ **Mixed Wave DM**

tachyonic bare term

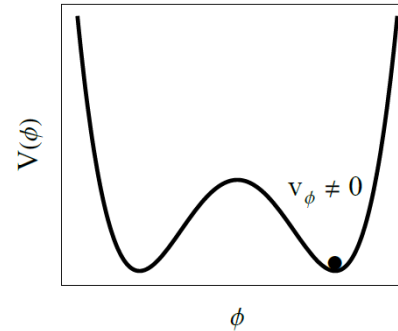
ϕ 

Φ 



$-\mu_0^2 < \kappa \Phi^2$ ($\mu_{\text{eff}}^2 > 0$)

Symmetry-restored phase
w/ zero VEV



$-\mu_0^2 > \kappa \Phi^2$ ($\mu_{\text{eff}}^2 < 0$)

Symmetry-broken phase
w/ nonzero VEV