

Core–Collapse Supernova Cooling in the Dark Z Model

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Based on the ongoing work
with Eung Jin Chun, Kwang Sik Jeong, Dong–Won Jung, and Seokhoon Yun

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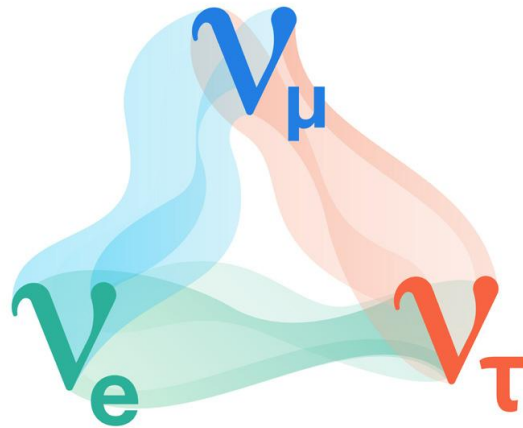
1. Dark Z Model

Physics Beyond the Standard Model

- Standard Model: Excellent explanation of nature up to the TeV ($\sim 10^{-18}$ m) scale; unresolved problems require BSM physics

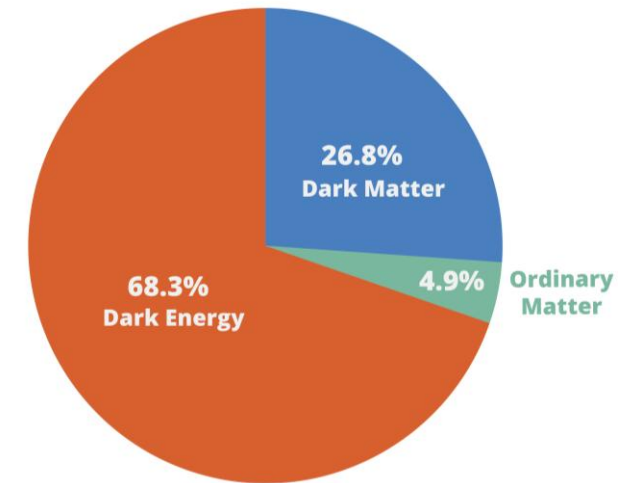


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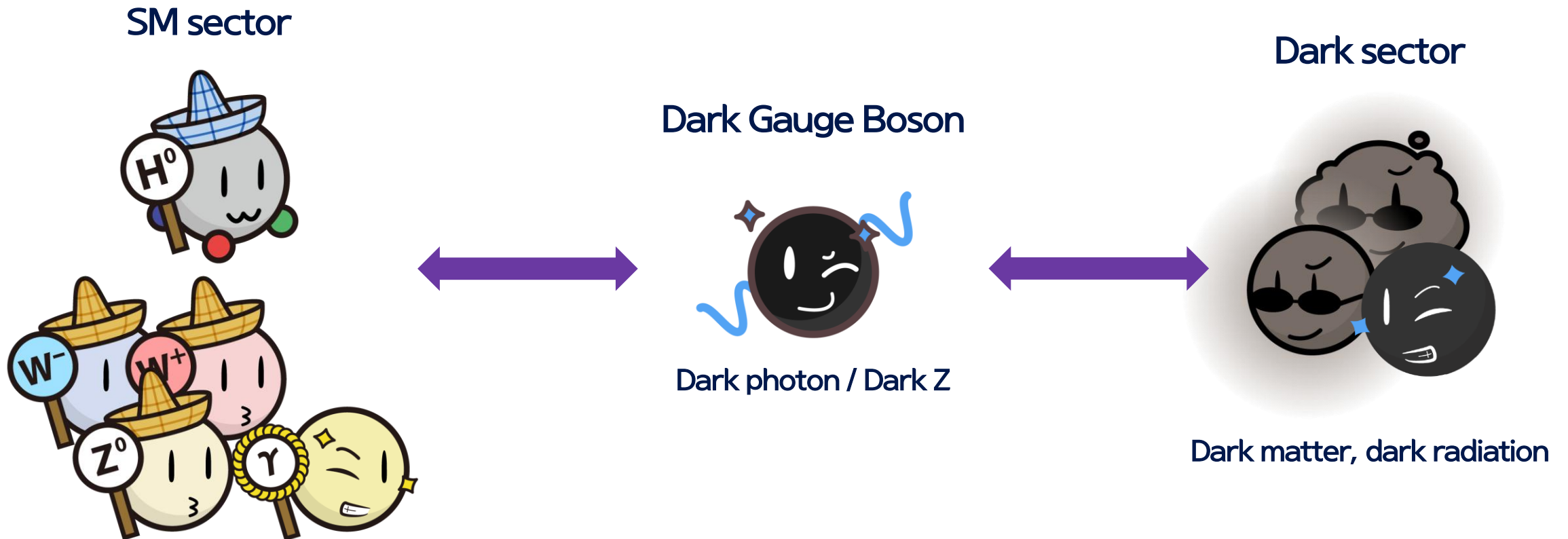
Estimated matter-energy content of the Universe



Planck 2018

Dark Sector

- Dark sector: hidden sector that are neutral under the SM gauge symmetries
- Dark gauge boson: may act as a mediator between the SM sector and dark sector



Dark Z Model

- Dark Z (Z'): Gauge boson of a hidden $U(1)_D$ gauge symmetry
- Couples to the SM currents via kinetic and mass mixing

Kinetic mixing

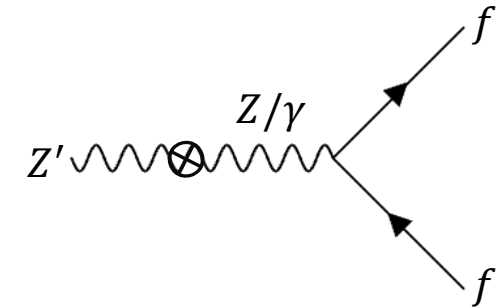
$$\frac{\varepsilon_0}{2} B^{\mu\nu} X_{\mu\nu} - \frac{\varepsilon_3}{2} W_3^{\mu\nu} X_{\mu\nu}$$

($\varepsilon_0, \varepsilon_3$: Kinetic mixing parameters)

Mass mixing

$$m_{ZX}^2 Z^\mu X_\mu$$

(m_{ZX}^2 : Mass mixing parameter)



$$\mathcal{L}_{\text{int}} = Z'_\mu (e\theta_A J_{\text{EM}}^\mu + g_Z\theta_Z J_{\text{NC}}^\mu)$$

- Low mass limit: $m_{Z'} \ll m_Z$

$$J_{\text{NC}}^\mu \text{ coupling } \theta_Z \sim (\varepsilon_0 - \varepsilon_3) \frac{m_{Z'}^2}{m_Z^2} - \frac{m_{ZX}^2}{m_Z^2}$$

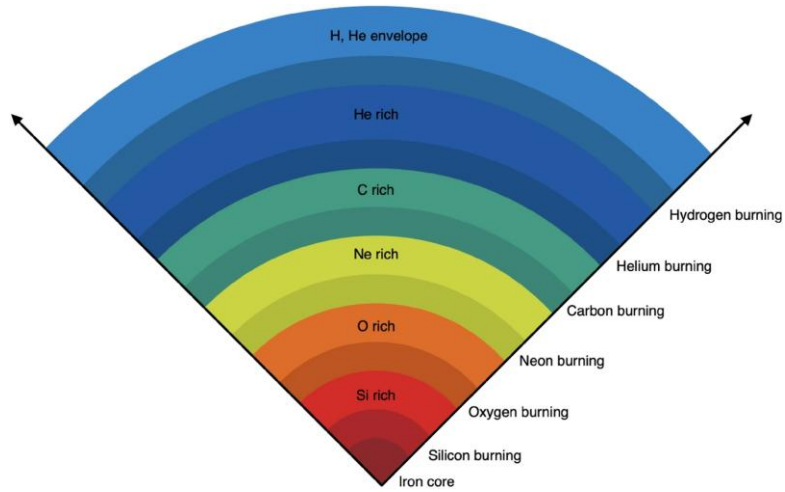
Suppressed

No $m_{Z'}/m_Z$ suppression

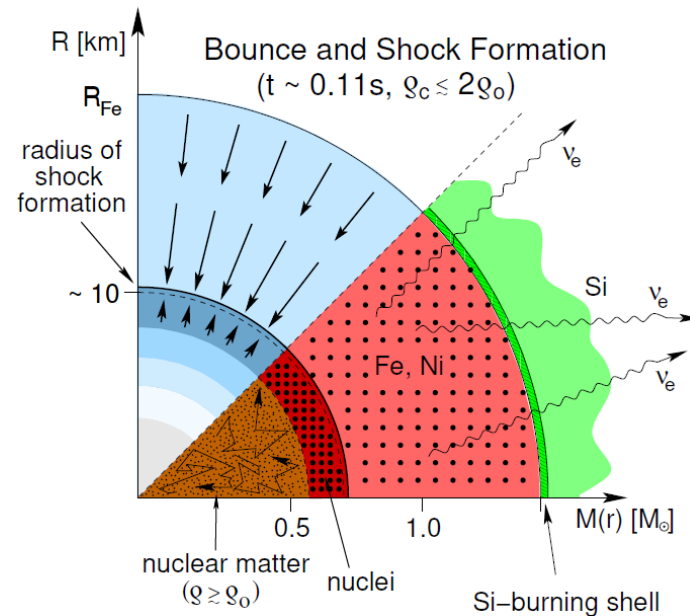
2. Core–Collapse Supernova

Initial Phase of Core-Collapse Supernova

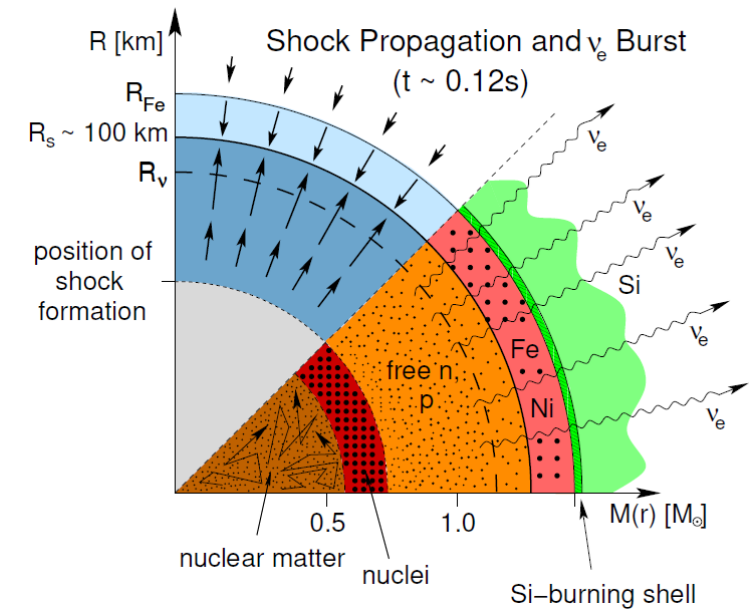
- At the final stage of stellar evolution, an Fe core forms and nuclear fusion ceases
- Upon reaching nuclear saturation density, core bounce occurs and a shock forms
- Shock dissociates heavy nuclei and generates a ν_e burst



A. Arcones and F. Thielemann, 2022



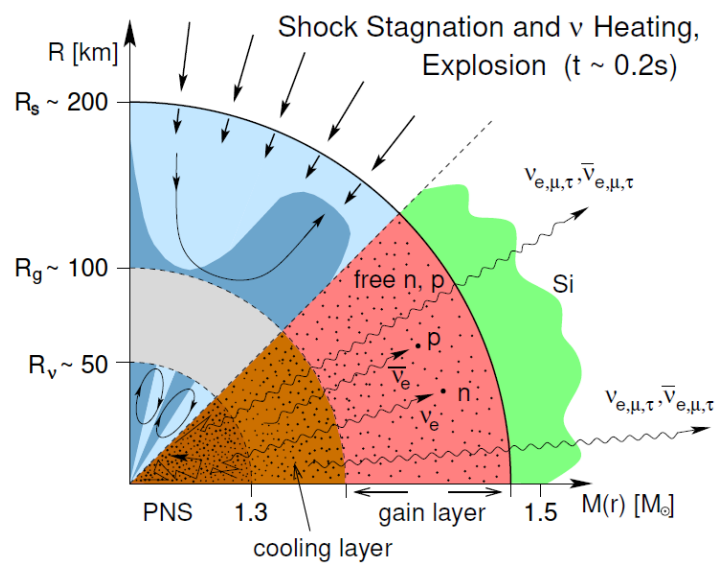
H.T. Janka et al., 2007



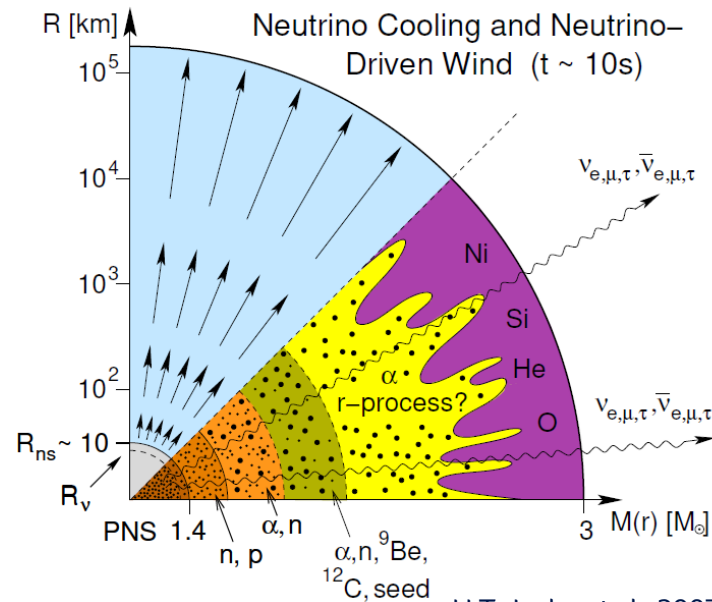
H.T. Janka et al., 2007

PNS Formation and Shock Revival

- Hot and dense proto-neutron star (PNS) formed; $T \sim \mathcal{O}(10) \text{ MeV}$, $\rho \sim \mathcal{O}(10^{14}) \text{ g/cm}^3$
- Neutrinos from the PNS deposit energy into the shock
 - Revival success $\rightarrow$ Neutron star
 - failure $\rightarrow$ Black hole

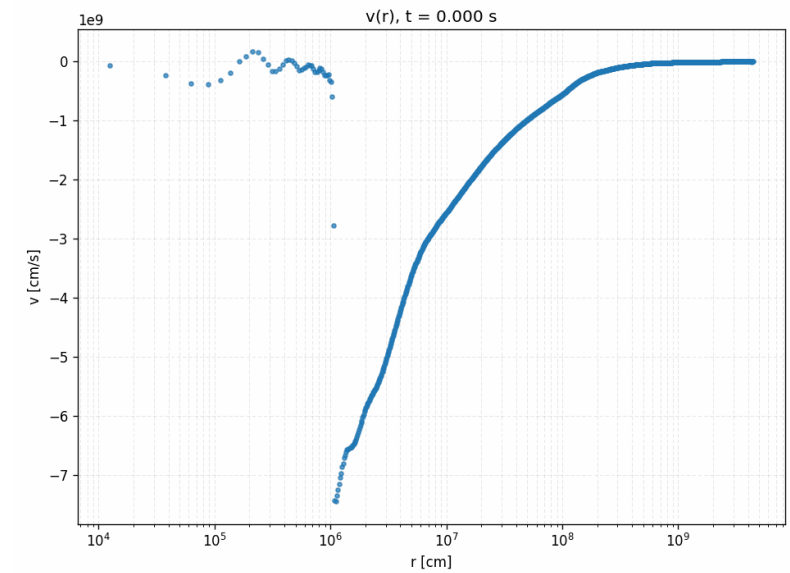


H.T. Janka et al., 2007



H.T. Janka et al., 2007

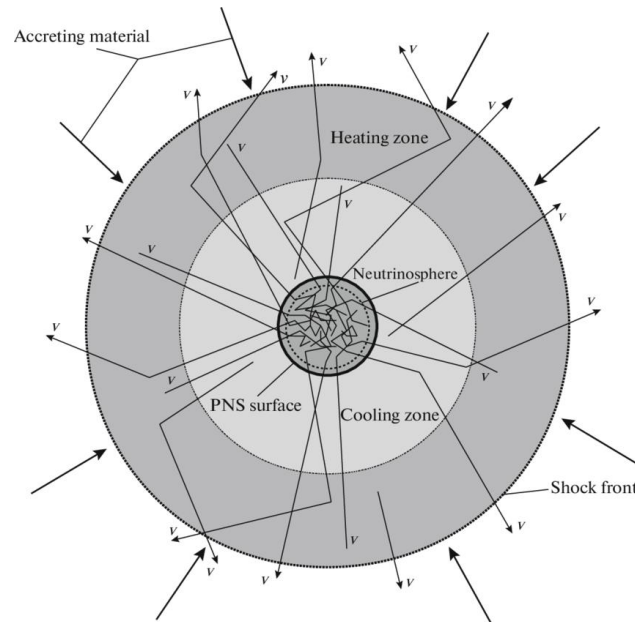
Simulation result: velocity profile



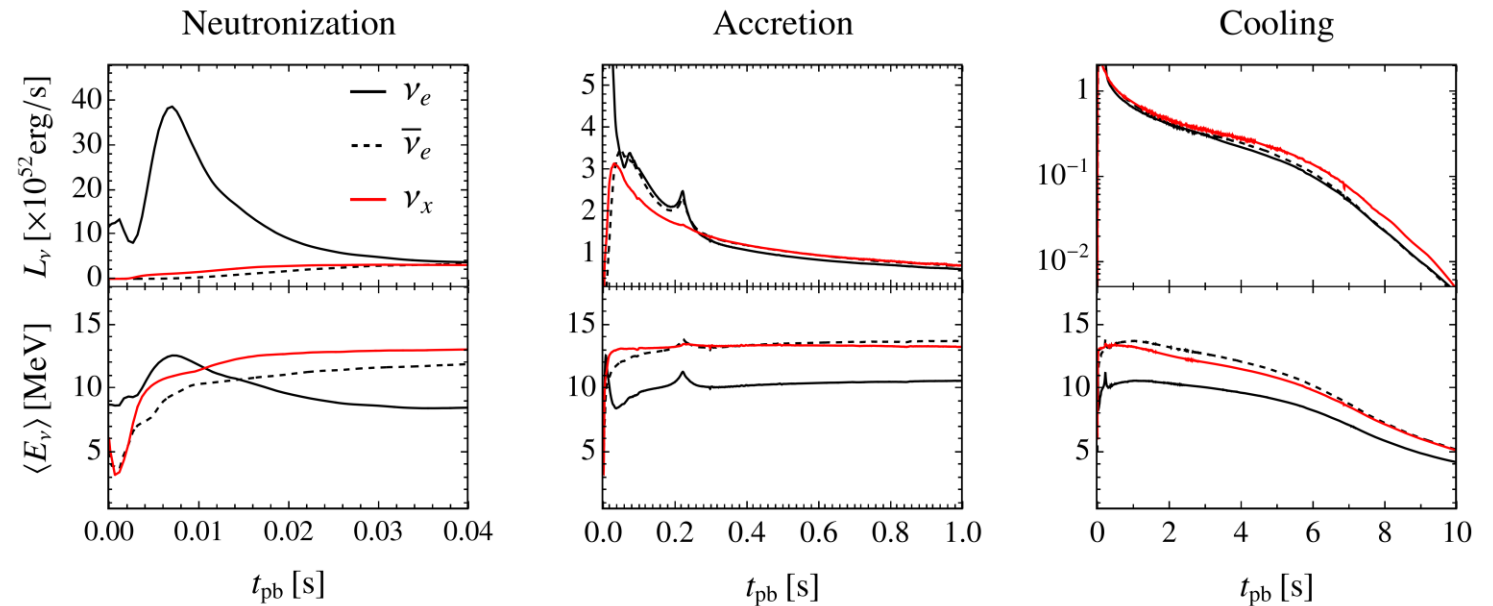
Neutrino Cooling

- About 99% of the gravitational binding energy is released through neutrinos of all flavors

$$L_\nu \sim 10^{52} \text{ erg/s}, \quad \langle E_\nu \rangle \sim \mathcal{O}(10) \text{ MeV}, \quad t_{\text{diff}} \sim 10 \text{ s}$$



Y. Krivosheyev and G. Kogan, 2018

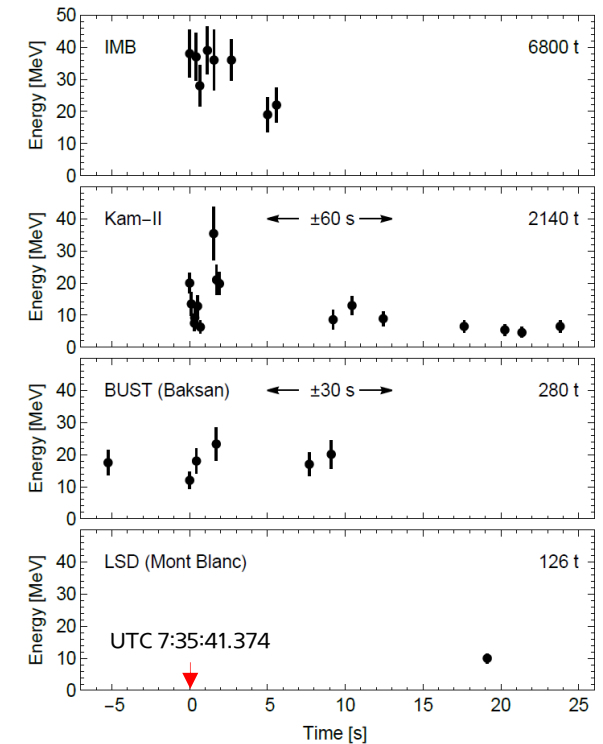
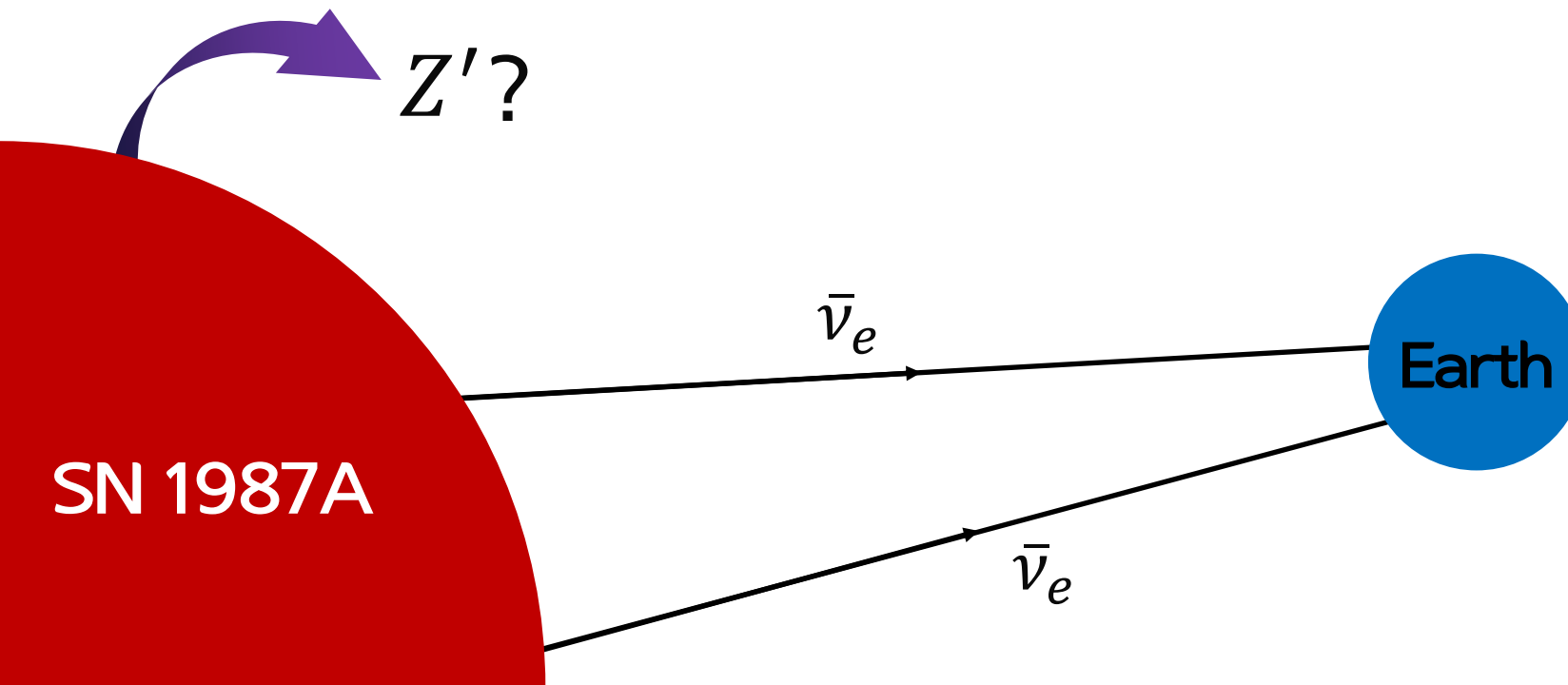


P. Carenza et al., 2024

3. Supernova Cooling

SN 1987A Cooling Bounds

- Electron antineutrinos were observed from SN 1987A
- Excessive energy loss through Z' emission would shorten the observed neutrino signal
- This requires $L_{Z'} \lesssim \mathcal{O}(10^{52})$ erg/s at $t_{\text{pb}} \sim 1$ s



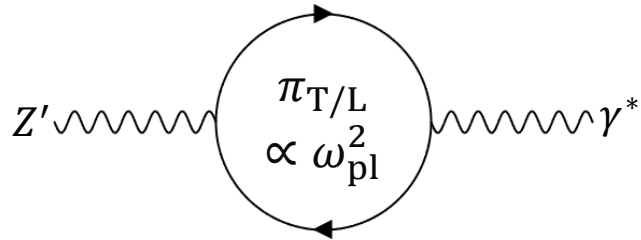
D. F. G. Fiorillo et al, 2023

Dark Z in the PNS Medium

- PNS environment: Strongly degenerate electrons

$$\omega_{\text{pl}} \sim \mathcal{O}(10) \text{ MeV}$$

- In-medium plasma effect induced by $Z' e^-$ coupling

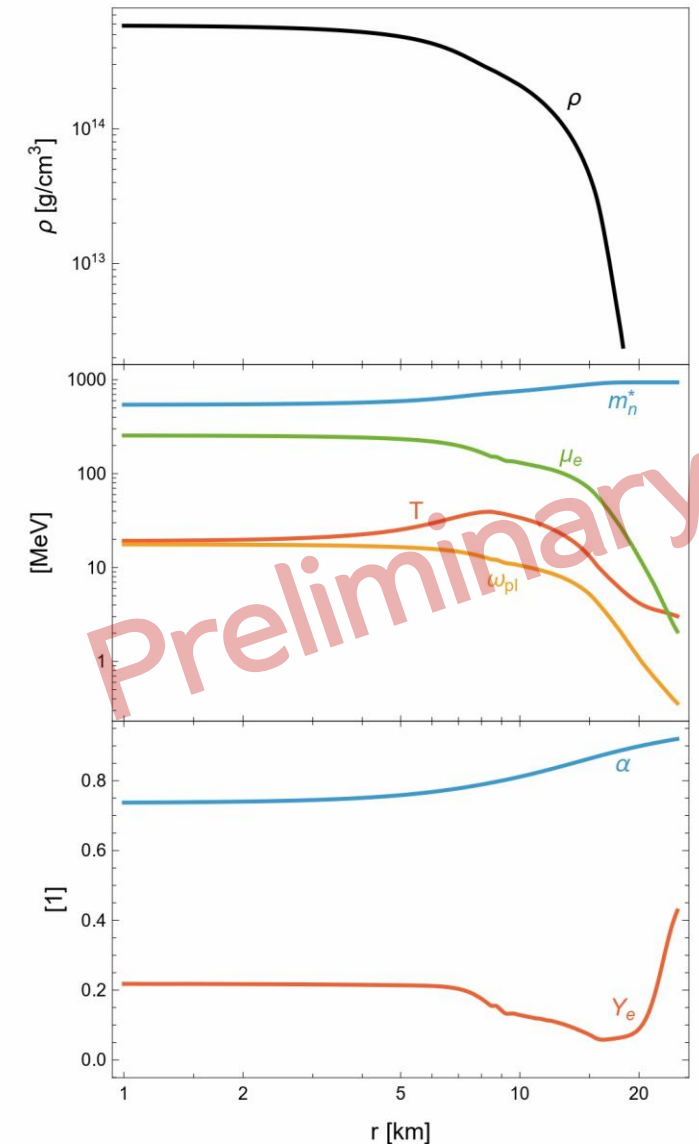


In-medium mixing with the photon

$$q_{\text{eff, T/L}}^V = q^V \frac{m_{Z'}^2}{m_{Z'}^2 - \pi_{\text{T/L}}}$$

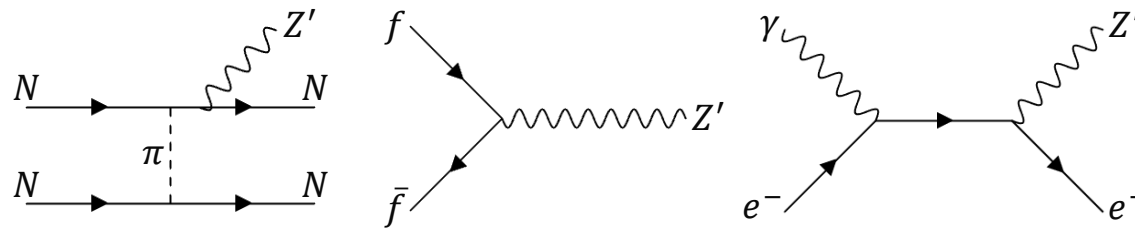
Vector coupling is modified for charged particles

- T and L modes are modified differently
- Resonance for $m_{Z'}^2 \simeq \text{Re } \pi_{\text{T/L}}$
- **Axial coupling is unaffected**



Dark Z Production inside the PNS

■ Production channels



Axial

nn brem.
dominant

ν inverse decay
subdominant

Vector

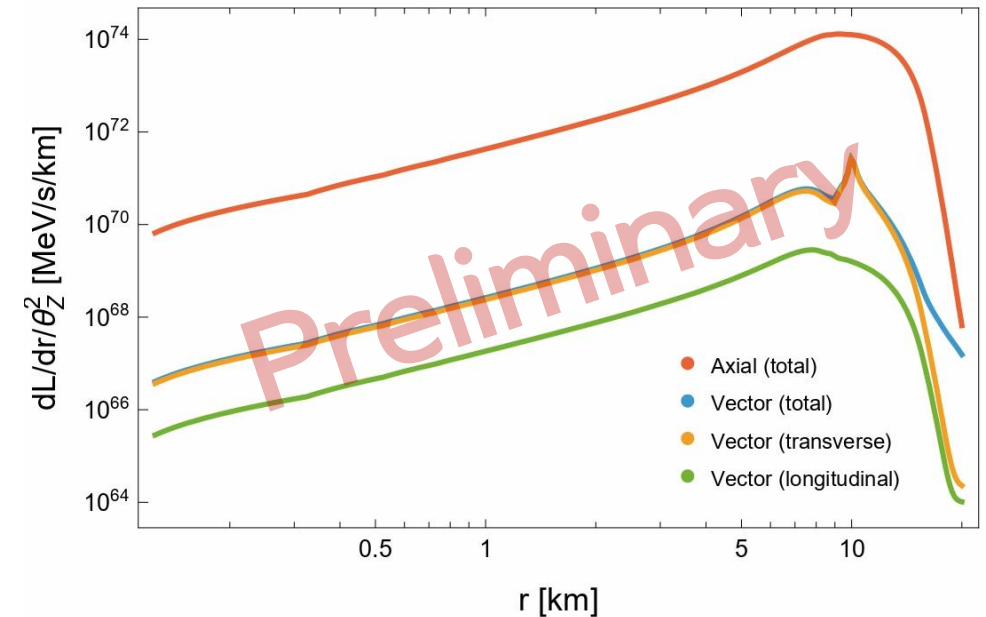
pn brem.
dominant

ν/e inverse decay
subdominant

Sizable near
resonance

- NN bremsstrahlung channels are modified by multiple coherent scatterings in dense medium

Differential luminosity for $m_{Z'} = 9.8$ MeV



(ν inverse decays are omitted in vector T/L)

Dark Z Absorption and Transport

- Absorption occurs via the inverse processes of production

$$\tau = \int_{\text{path}} dl \frac{\Gamma_{\text{abs}}}{v_{Z'}} \quad \begin{array}{l} \tau \gg 1: \text{Trapped inside the PNS} \\ \tau \ll 1: \text{Freely-streams} \end{array}$$

$$\langle P_{\text{esc}} \rangle_{\mu} = \langle e^{-\tau} \rangle_{\mu}$$

- NN bremsstrahlung gives a dominant contribution, while decays into neutrinos are sizable for massive Z'

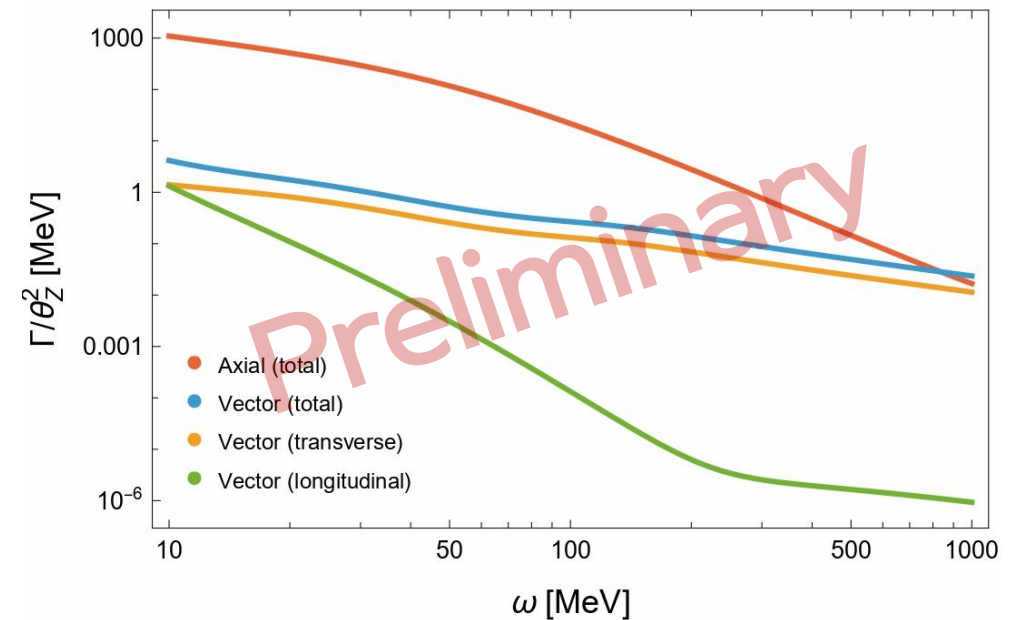
- Gravitational effects

Gravitational trapping: low-energy Z' cannot escape

Gravitational redshift: energy decreases during propagation

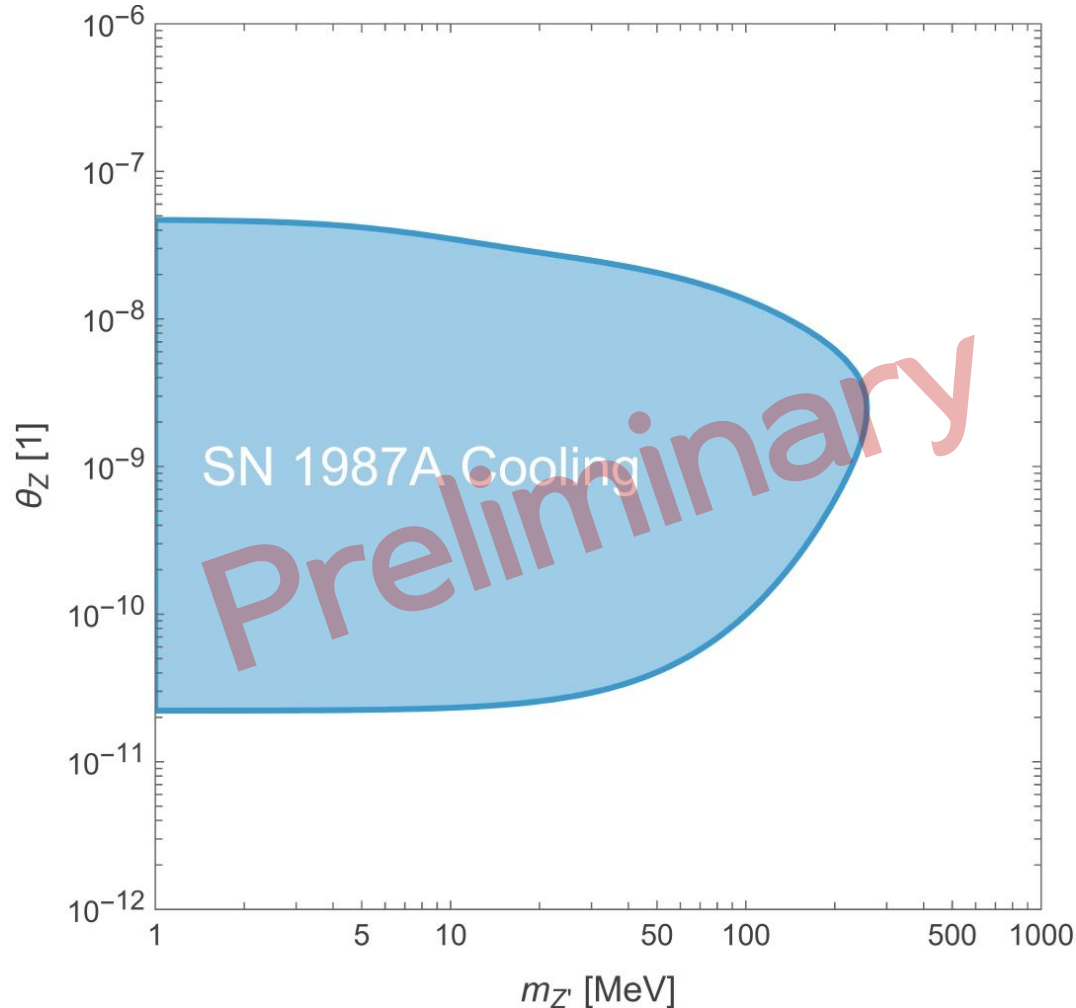
Time dilation: luminosity at infinity is reduced

Absorption rate at 10 km for $m_{Z'} = 9.8$ MeV



(ν inverse decays are omitted in vector T/L)

SN 1987A Cooling Bounds on the Z'



- Cooling bound

$$L_{Z'}^\infty = 4\pi \int_0^{R_\nu} dr r^2 \alpha^2(r) \int_{m_{Z'}/\alpha(r)}^\infty d\omega \left(\frac{dQ_{Z'}}{d\omega} \right)_{T/L} \langle e^{-\tau_{T/L}} \rangle_\mu$$
$$\leq L_\nu \simeq 4.4 \times 10^{52} \text{ erg/s}$$

- Cooling bound is mainly controlled by the axial coupling
- Vector contribution is subdominant despite resonant plasma enhancement

4. Conclusions

Conclusions

- Dark Z boson can couple to the SM neutral current through mass mixing
- PNS plasma modifies the vector interaction differently for T/L modes, while the axial interaction does not receive this in-medium mixing correction
- Dark Z production is dominated by axial nn bremsstrahlung although vector channels receive enhanced contributions near resonance
- SN 1987A cooling constrains the dark Z parameter space, the resulting bound is strongly controlled by the axial coupling