

Unbound states in ^{17}C probed via single-neutron removal from ^{18}C at 245 MeV/u

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Rare Isotope Science Project

S.Kim, Thesis:

“Spectroscopy of ^{17}C via one-neutron knockout reaction”

c.f.) J.W.Hwang, S.Kim et al., PLB769(2017)503.

“Single-neutron knockout from ^{20}C and the structure of ^{19}C ”

Issues:

- A) The cross-shell states in light neutron-rich nuclei are known to be notoriously difficult to describe theoretically, because they involve transitions encompassing two major shells on a proton-neutron asymmetric system.
- B) Obtaining a reliable relevant effective interaction in such cases remains of a special difficulty.
- C) The cross-shell states might have implications to astrophysical radiative neutron captures processes, since they are strongly influenced by the presence of dipole resonances.

This presentation:

- An experimental attempt to enlarge the known realm of such dipole states in N-rich C isotopes.
- Capability of recently available shell-model calculations on dipole and other states examined.

RIKEN SAMURAI Day-1 collaborators

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Nuclear structure of very-neutron-rich nuclei

- Exotic properties are anticipated due to the large excess of neutrons and weak binding of valence neutrons

1. Nucleon density distribution may change drastically.

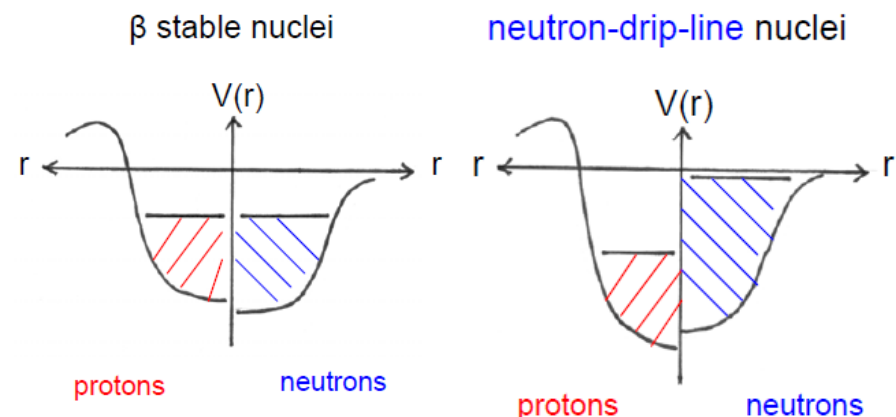
neutron halo, cluster

2. Mean field can change accordingly.

modification of shell structure, melting of magicity

3. Deformability of nuclei is modified accordingly.

4. Large Fermi energy difference between neutrons and protons will modify the effects of Pauli blocking and enhance the coupling with continuum states.



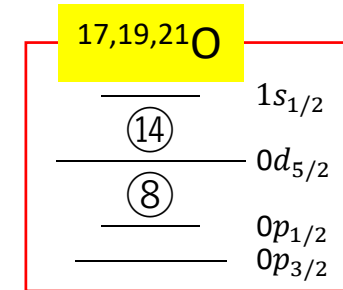
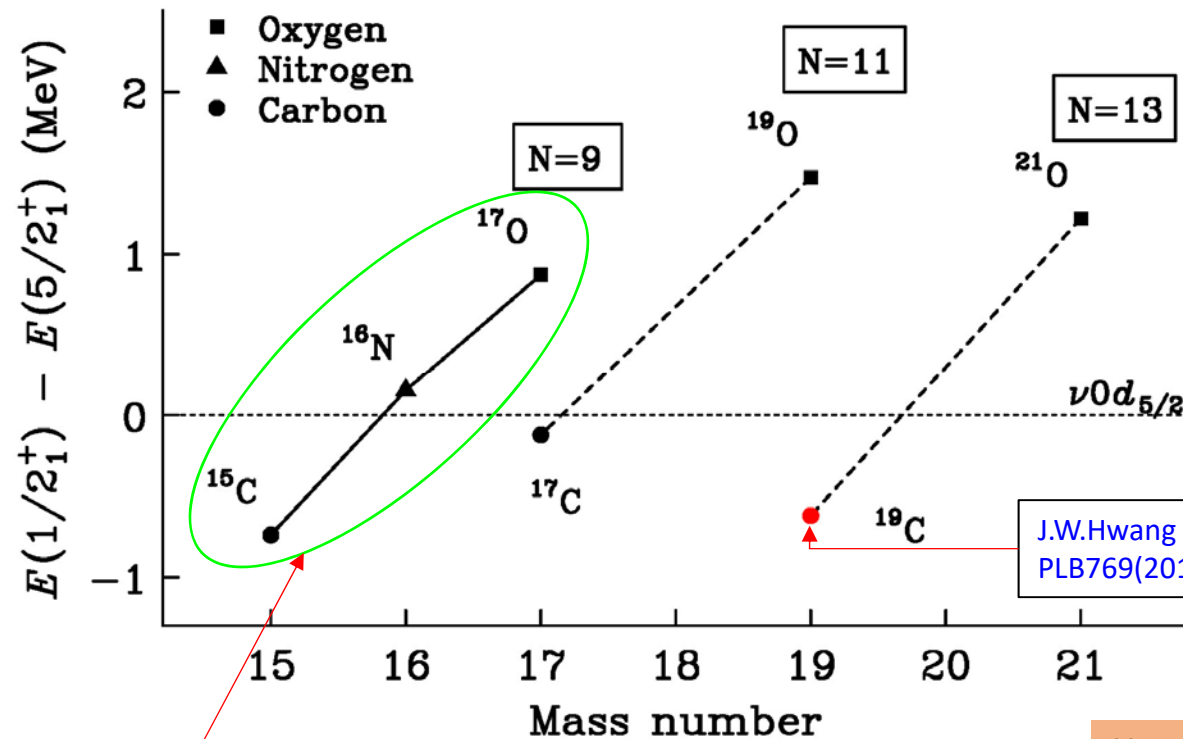
Why study Carbon isotopes

1. In light elements ($Z < 10$) the number of isotopes is the largest (14); suited for systematic study.
2. They offer a basis to study intricate effects of the nuclear forces.
 - Quenched $N=14$ shell gap
 - Proton core polarization effects
3. The region ($A \lesssim 20$) represents a frontier of current *ab initio* calculations including continuum and three-body forces. Advanced shell-model interactions also exist.
4. Possible implications of properties of light neutron-rich nuclei (including C) to r-process nucleosynthesis in neutrino-driven winds have been suggested. Dipole resonances ($L=1$) have impact on the neutron capture rates.

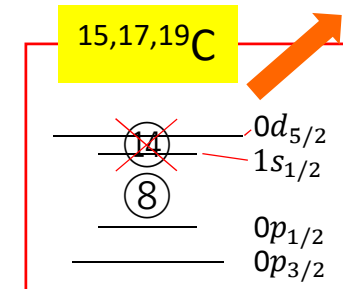
M.Terasawa et al., ApJ562(2001)470. T.Sasaqui et al., ApJ634(2005)1173.

S.Goriely, PLB436(1998)10.

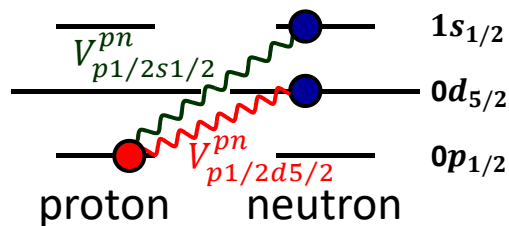
M.Stanoiu et al., PRC78(2008)034315.



Deformation



◆ I.Talmi, I.Unna, PR4(1960)469.

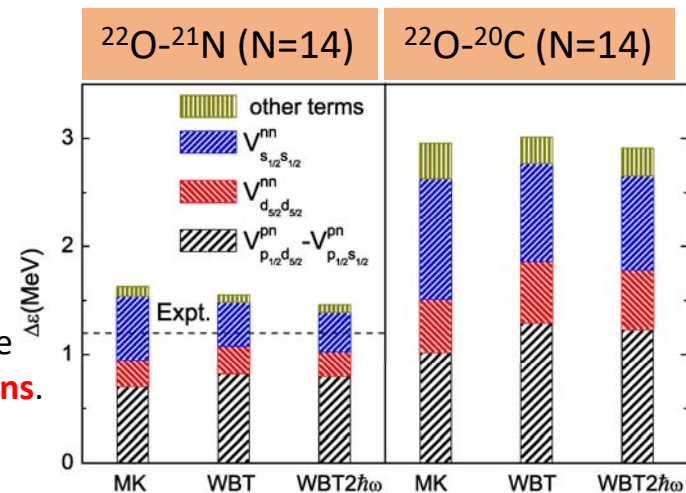


$$2 \left(V_{p1/2d5/2}^{pn} - V_{p1/2s1/2}^{pn} \right)$$

◆ C.X.Yuan et al., Nucl. Phys. A 883 (2012) 25.

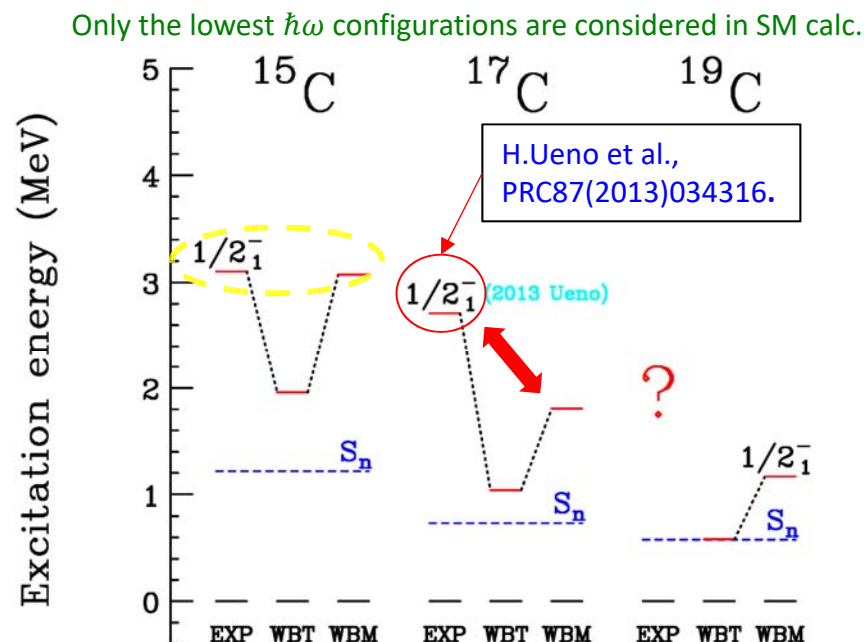
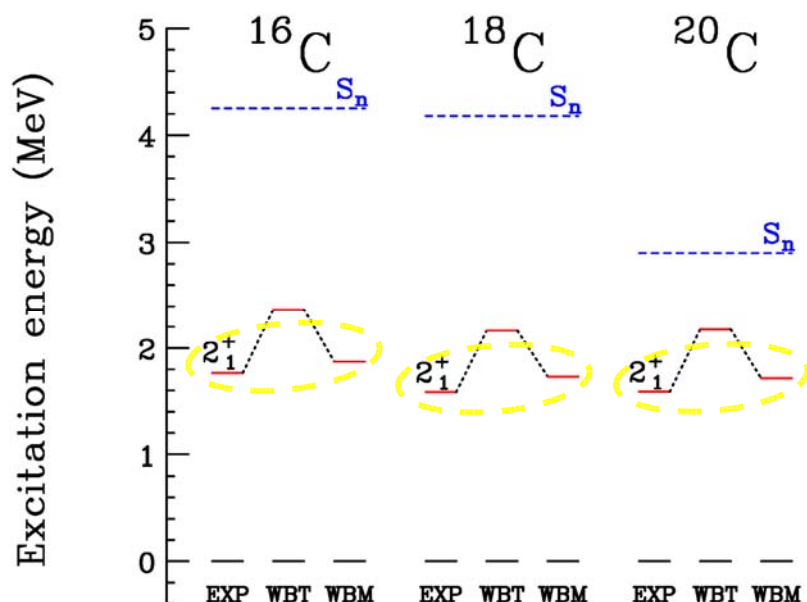
N=14 shell gap evolution

- Proton-neutron int.
 - Attractive neutron-neutron int. due to the $\Delta\epsilon$ (MeV)
- many-body correlations.**



Proton core polarization effects

- Responsible for reduced neutron-neutron matrix elements in the sd shell in Carbon
- WBM provides a remedy for 2_1^+ in $^{16,18,20}\text{C}$ and $1/2_1^-$ in ^{15}C .
- To what extent WBM is applicable?



Overprediction of $E(2_1^+)$ with WBT.

- Core polarization effects** (should be mitigated in C).

K.Sieja, F.Nowacki, NPA857(2011)9.

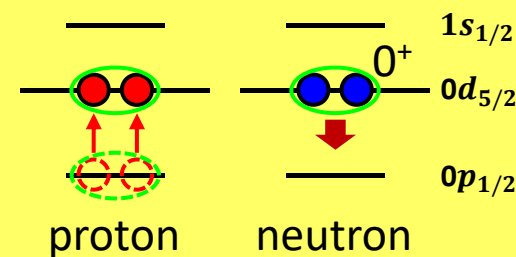
- Extended radial wave function.

A.Signoracci et al., PRC83(2011)024315.

WBT: E.K Warburton, B.A.Brown, PRC46(1992)923.

WBM: WBP with USD part reduced by 25%.

Proton core polarization effects



T.T.S.Kuo, G.E.Brown, NP85(1966)40.

Shell-model interactions examined

Code (Model space)	Interactions
NuShell@MSU (OXBASH) (<i>spsdpf</i>)	WBT : E.K Warburton, B.A.Brown, PRC46(1992)923. • $(0 - 1)\hbar\omega$ WBM : WBP with USD part reduced by 25%. M.Stanoiu et al., PRC78(2008)034315.
NuShellX@MSU (<i>psd</i> , $\hbar\omega$ unrestricted) B.A.Brown, W.D.M.Rae, Nucl. Data. Sheets 120 (2014) 115.	WBT_X : WBT as calculated with NuShellX. WBM_X : WBM as calculated with NuShellX. YSOX : C.Yuan et al., PRC85(2012)064324. • $(0 - 3)\hbar\omega$ • $\langle pp V pp\rangle$ SFO • $\langle sd V sd\rangle$ SDPF-M • $\langle psd V psd\rangle$ and $\langle pp V sdsd\rangle$ V_{MU} T.Otsuka et al., PRL104(2010)012501. • V_{MU} strengths adjusted to reproduce N-rich B,C,O masses. <div data-bbox="1653 678 1944 869"> <p>2p2h excitations</p> </div>
Ab Initio Coupled-Cluster Effective Interaction CCEI : G.R.Jansen et al., PRL113(2014)142502. <ul style="list-style-type: none"> Nonperturbative calculation. Based on NN and NNN forces from chiral effective field theory. No account for the particle continuum. 	

$$\hat{H} = \sum_{i < j} \left(\frac{(p_i - p_j)^2}{2mA} + \hat{V}_{NN}^{(i,j)} \right) + \sum_{i < j < k} \hat{V}_{3NF}^{(i,j,k)}$$

Objectives

- To furnish spectroscopic information on ^{17}C , including cross-shell states.

- **Confirm $1/2_1^-$ (and $3/2_1^-$) in ^{17}C reported in a βn study.**

- **New information on other states in ^{17}C .**

H.Ueno et al.,
PRC87(2013)034316.

- **Test theories against new spectroscopy.**

(1) enlarged $\hbar\omega$ model space, (2) reduced sd neutron-neutron matrix elements, (3) new SM interactions.

- To develop a new experimental technique involving neutron knockout which complements the βn spectroscopy.

- **Newly commissioned SAMURAI spectrometer.**

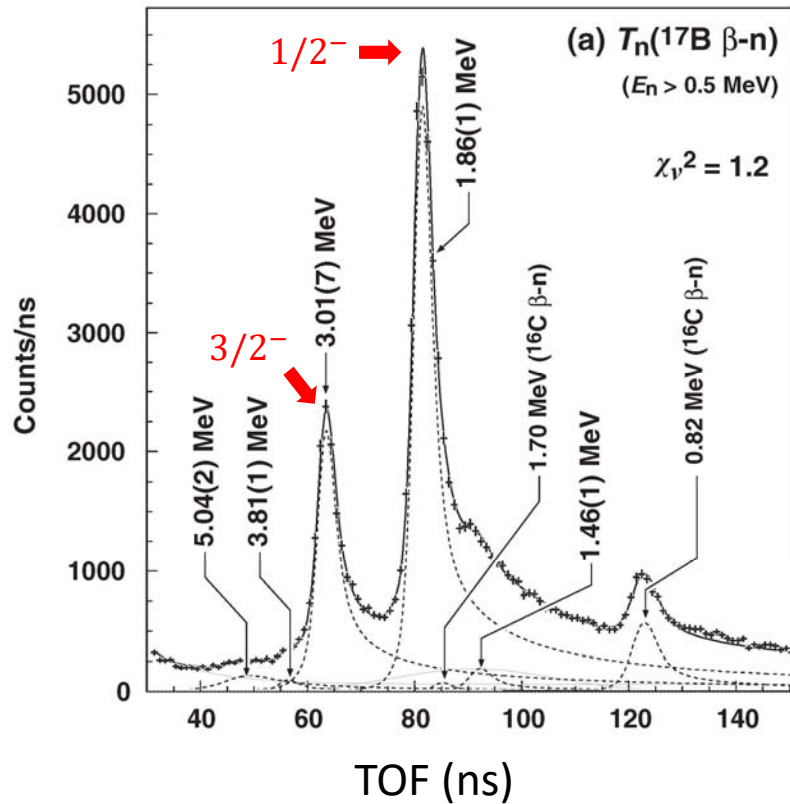
T.Kobayashi et al.,
NIMB317(2013)294.

- **Reaction:** $^{12}\text{C}(^{18}\text{C}, ^{17}\text{C}^* \rightarrow ^{16}\text{C} + n)$ @ 245 MeV/nucleon

- **Observables:** σ_{-1n} , $d\sigma/dp_{\parallel}$, $d\sigma/dp_T$.

β -delayed neutron and γ -ray spectroscopy of ^{17}C utilizing spin-polarized ^{17}B

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 N. Takahashi,³ and K. Yoneda¹

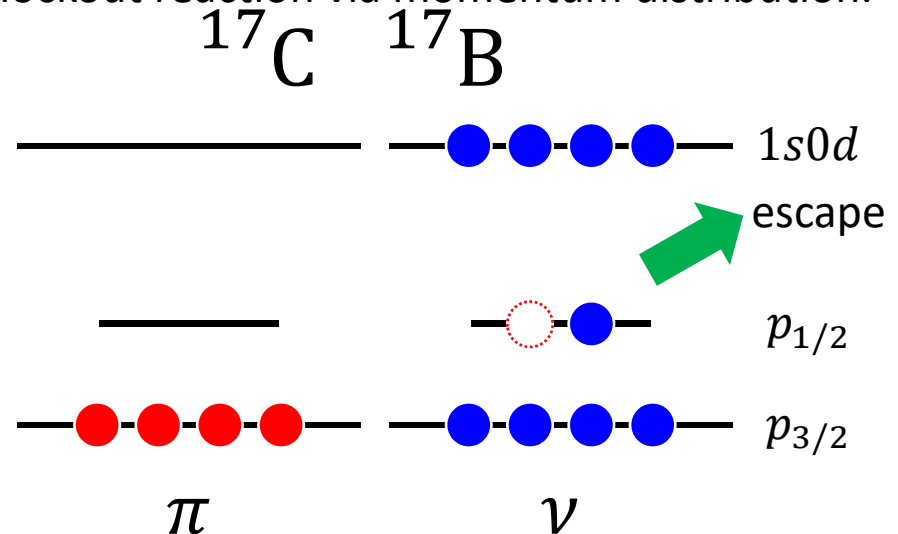


E_n (MeV)	E_x (MeV)	Width (MeV)	$\text{Log } ft$	J^π
1.86(1)	2.71(2)	0.04(1)	4.8(1)	$1/2^-$
3.01(1)	3.93(2)	0.16(4)	4.9(1)	$3/2^-$
1.86(1)	4.05(2)	0.06(6)	6.0(1)	$(5/2^-)$

- Allowed transitions do not involve parity change.

$$B(GT) = \frac{1}{2I_i + 1} \left| \left\langle f \left\| \sum_k \sigma_k \tau_{\pm}^k \right\| i \right\rangle \right|^2$$

- Resulting ν p-shell hole states can be populated by knockout reactions.
- The angular momentum, and thus parity, of the state can be determined experimentally in the knockout reaction via momentum distribution.

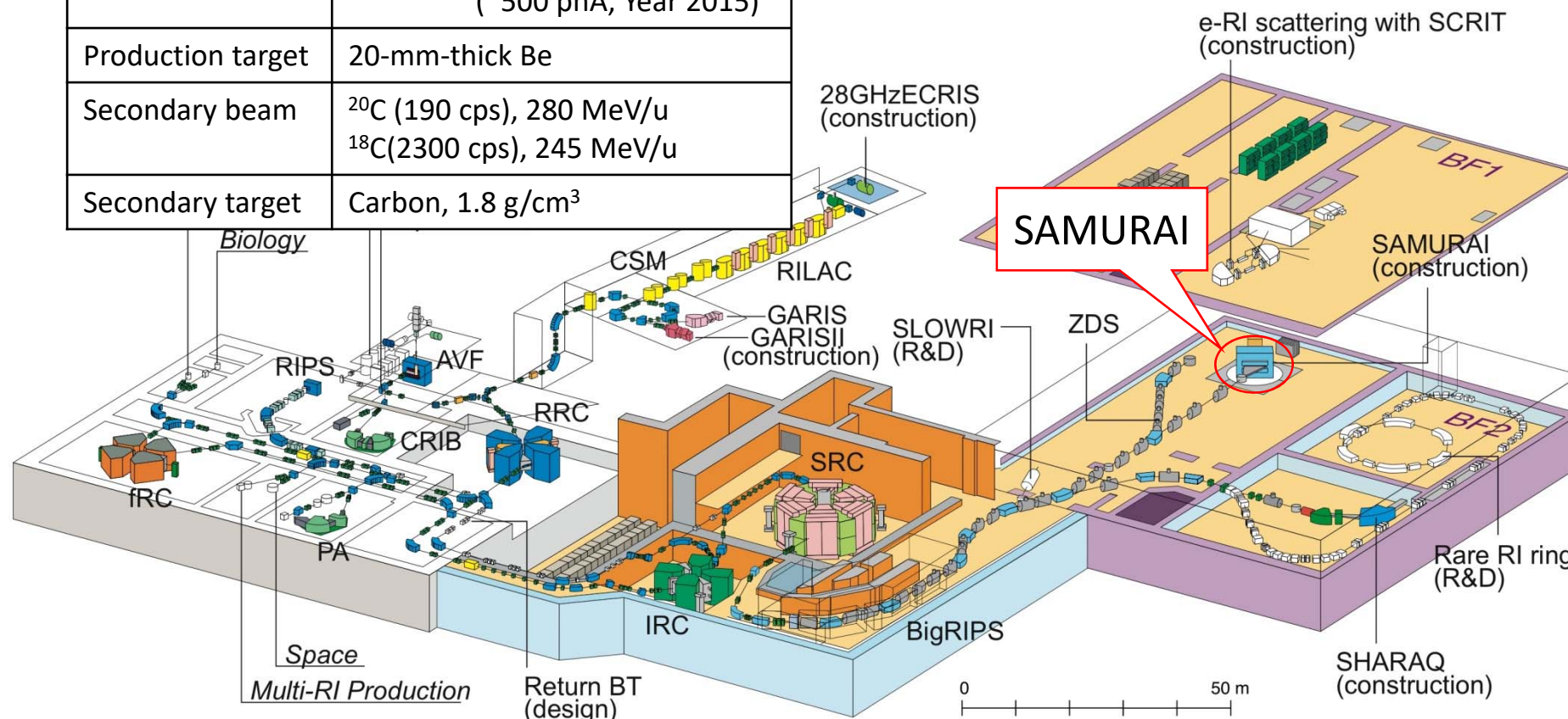


RIKEN RIBF : Radioactive Isotope Beam Factory

RIBF: Completed in 2007

New-Generation RI-Beam Facility

Primary beam	^{48}Ca , 345 MeV/u, 250 pnA MAX (~500 pnA, Year 2015)
Production target	20-mm-thick Be
Secondary beam	^{20}C (190 cps), 280 MeV/u ^{18}C (2300 cps), 245 MeV/u
Secondary target	Carbon, 1.8 g/cm ³



SAMURAI Day-1 campaign 2012

- I. Coulomb Breakup of ^{19}B and ^{22}C , Nakamura et al.
- II. Study of ^{18}B , ^{21}C , and excited states of ^{19}B , ^{22}C , Orr et al.
- III. Structure of Unbound Oxygen Isotopes ^{25}O , ^{26}O , Kondo et al.

Y.Kondo et al., PRL116(2016)102503.

Y.Togano et al., PLB761(2016)412.

J.W.Hwang et al., PLB769(2017)503.

SAMURAI spectrometer

Superconducting Analyzer for Multi-particles from RAdioIsotope beams

T.Kobayashi et al, NIMB317,294(2013).

➤ Invariant mass method



$$M_{\text{inv}} = \sqrt{(E_C + E_n)^2 - |\mathbf{P}_C + \mathbf{P}_n|^2}$$

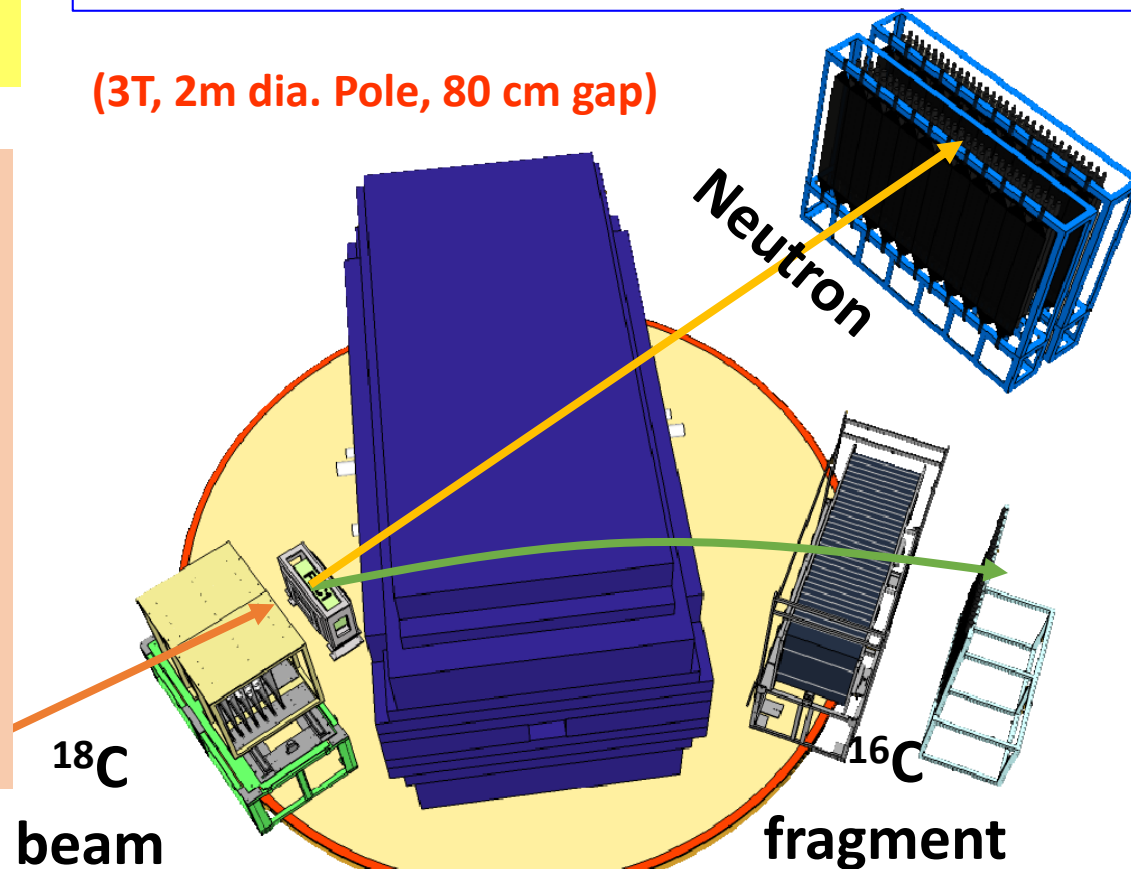
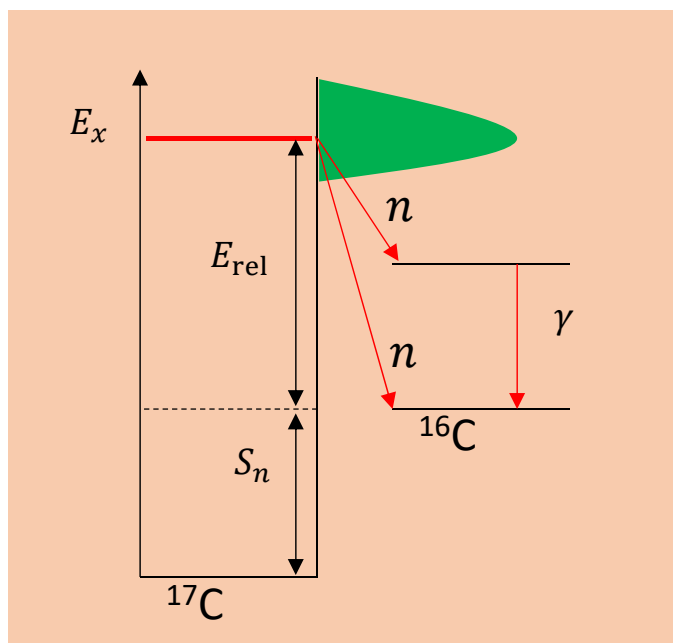
$$E_{\text{rel}} = M_{\text{inv}} - M_C - M_n$$

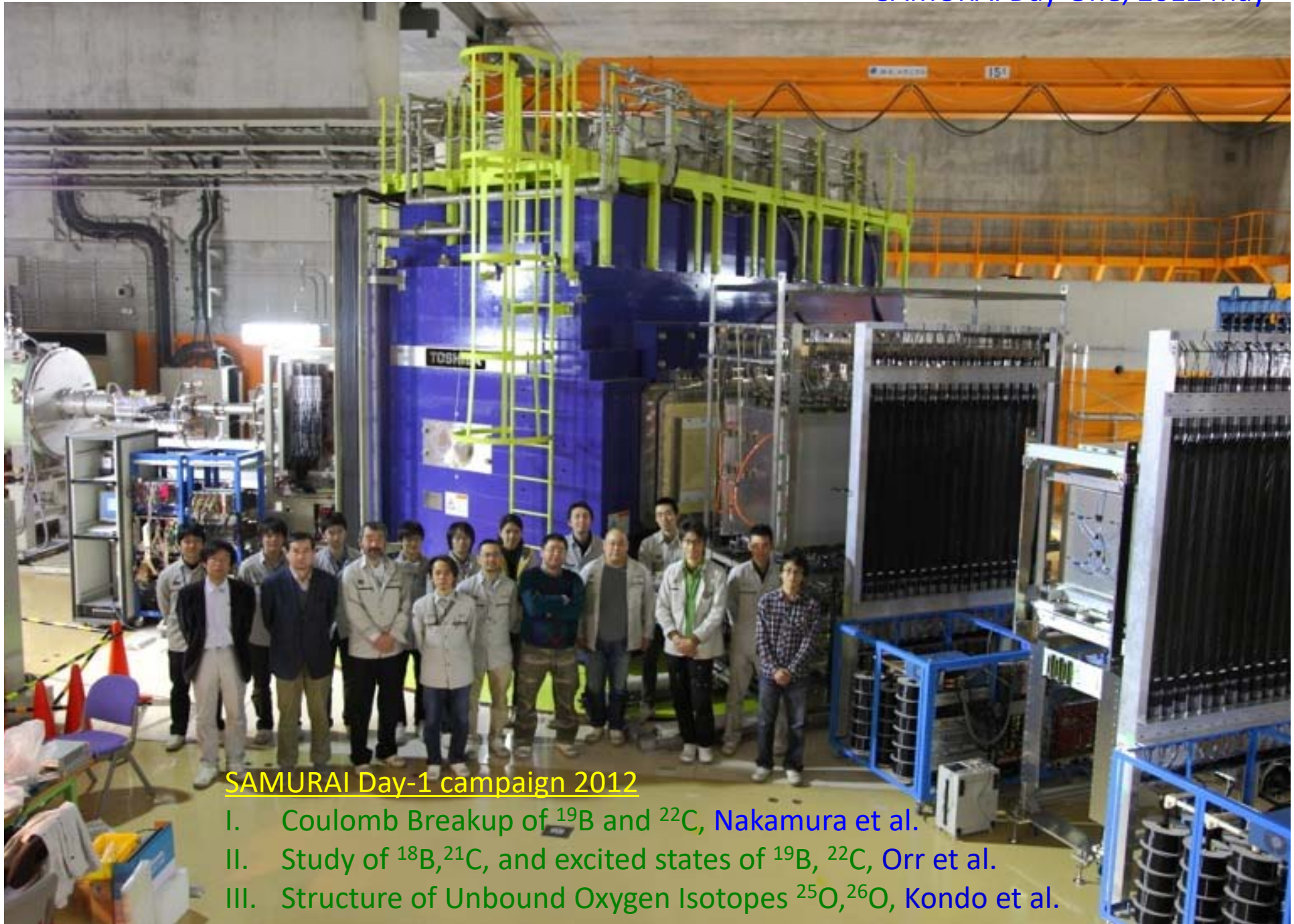
$$E_x = E_{\text{rel}} + S_n$$

- (1) High detection efficiency for decay neutrons. **Large ε_n**
- (2) High coverage of decay phase space due to kinematical focusing.
- (3) Large acceptance near $E_{\text{rel}} \sim 0$ MeV due to magnetic separation of neutrons and charged particles.
- (4) Good E_{rel} resolution suited for spectroscopy.

$$\Delta E_{\text{rel}} \propto \sqrt{E_{\text{rel}}(E_{\text{in}}/A)}(\Delta P/P)$$

(3T, 2m dia. Pole, 80 cm gap)





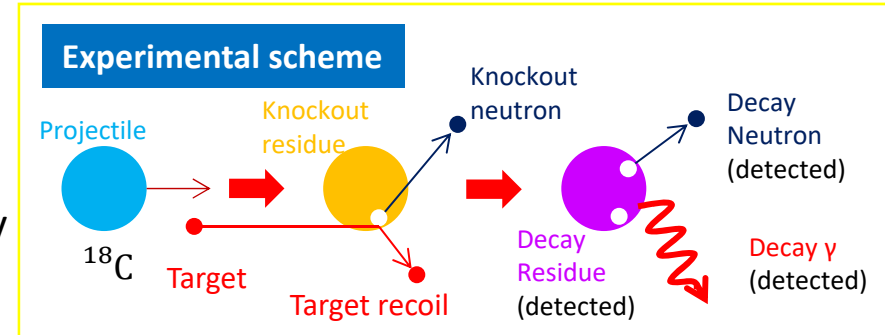
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- II. Study of ^{18}B , ^{21}C , and excited states of ^{19}B , ^{22}C , Orr et al.
- III. Structure of Unbound Oxygen Isotopes ^{25}O , ^{26}O , Kondo et al.

One-neutron knockout reaction

P.G.Hansen and J.A.Tostevin, Annu. Rev. Nucl. Part. Sci. 53, 219 (2003).

- Exhibit large cross sections ~ 100 mb.
- Knocked-out nucleon is democratically selected from valence orbits.
- Momentum distributions of the residue carry information on L, within the approximation the nucleon is suddenly removed.



Stripping cross section :

$$\frac{d\sigma_{\text{str}}}{d^3k_c} = \frac{1}{(2\pi)^3} \frac{1}{2l+1} \sum_m \int d^2b_v [1 - |S_v(b_v)|^2] \times \left| \int d^3r e^{-ik_c \cdot r} S_c(b_c) \psi_{lm}(\mathbf{r}) \right|^2$$

Inclusive -1n cross section:

$$\sigma_{-1n} = \sum_{nlj} \left(\frac{A}{A-1} \right)^N \cdot C^2 S(J^\pi, nlj) \cdot \sigma_{sp}(nlj, S_n^{\text{eff}})$$

Longitudinal momentum p_{\parallel} distribution:

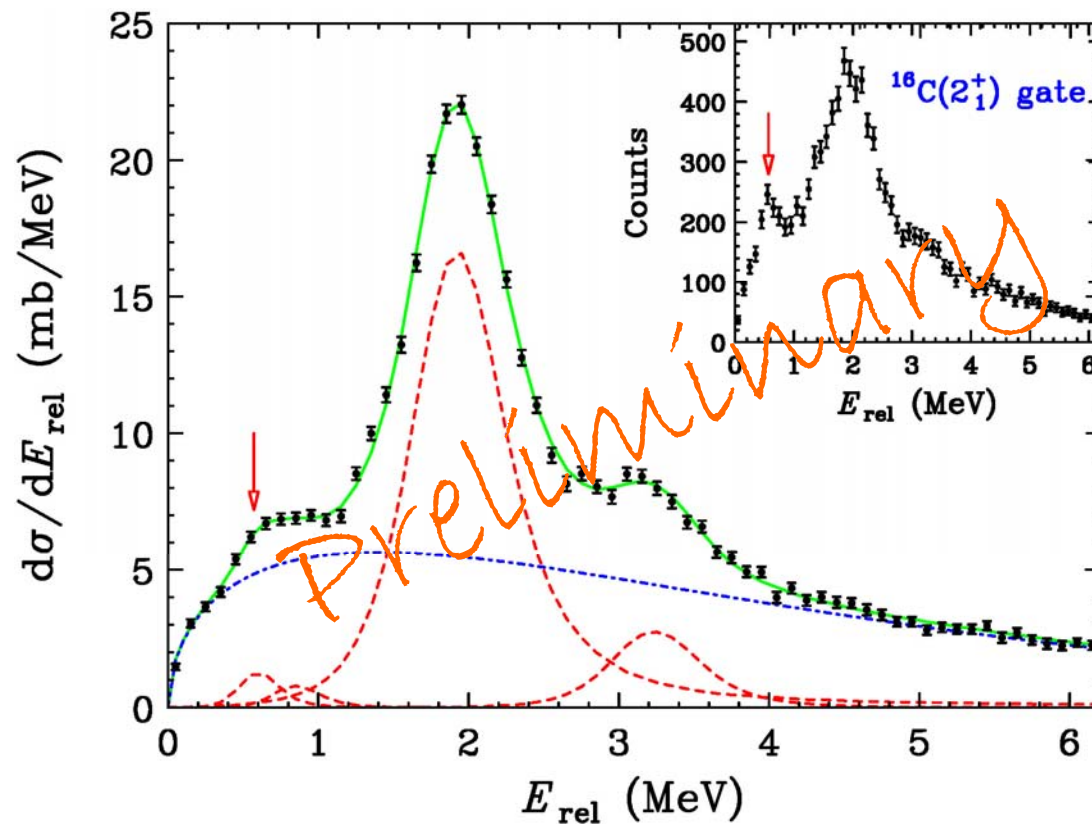
$$\frac{d\sigma_{\text{str}}}{dk_z} = \frac{1}{(2\pi)^3} \frac{1}{2l+1} \sum_m \int d^2b_v [1 - |S_v(b_v)|^2] \int_0^\infty d^2\rho |S_c(b_c)|^2 \left| \int_{-\infty}^\infty dz e^{-ik_z z} \boxed{\psi_{lm}(\mathbf{r})} \right|^2$$

Transverse momentum p_{\perp} distribution:

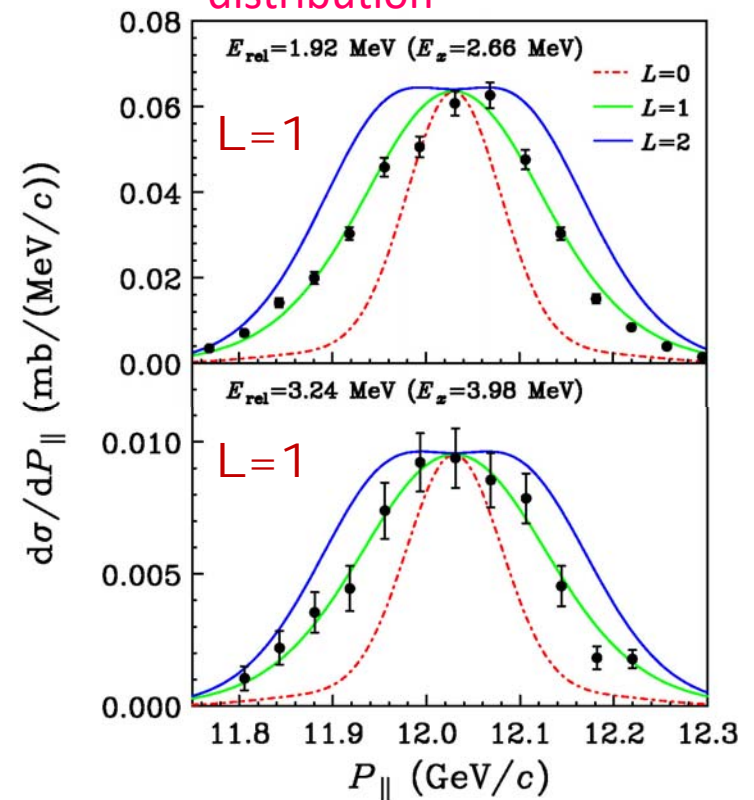
$$\frac{d\sigma_{\text{str}}}{d^2k_{\perp}} = \frac{1}{2\pi} \frac{1}{2l+1} \int_0^\infty d^2b_v [1 - |S_v(b_v)|^2] \sum_{m,p} \int_{-\infty}^\infty dz \left| \int d^2\rho \exp(-i\mathbf{k}_{\perp}^{\perp} \cdot \boldsymbol{\rho}) S_c(b_c) \boxed{\psi_{lm}(\mathbf{r})} \right|^2$$

Single-particle wave function

Results: $^{12}\text{C}(^{18}\text{C}, ^{17}\text{C}^* \rightarrow ^{16}\text{C} + n) @ 245 \text{ MeV/u}$

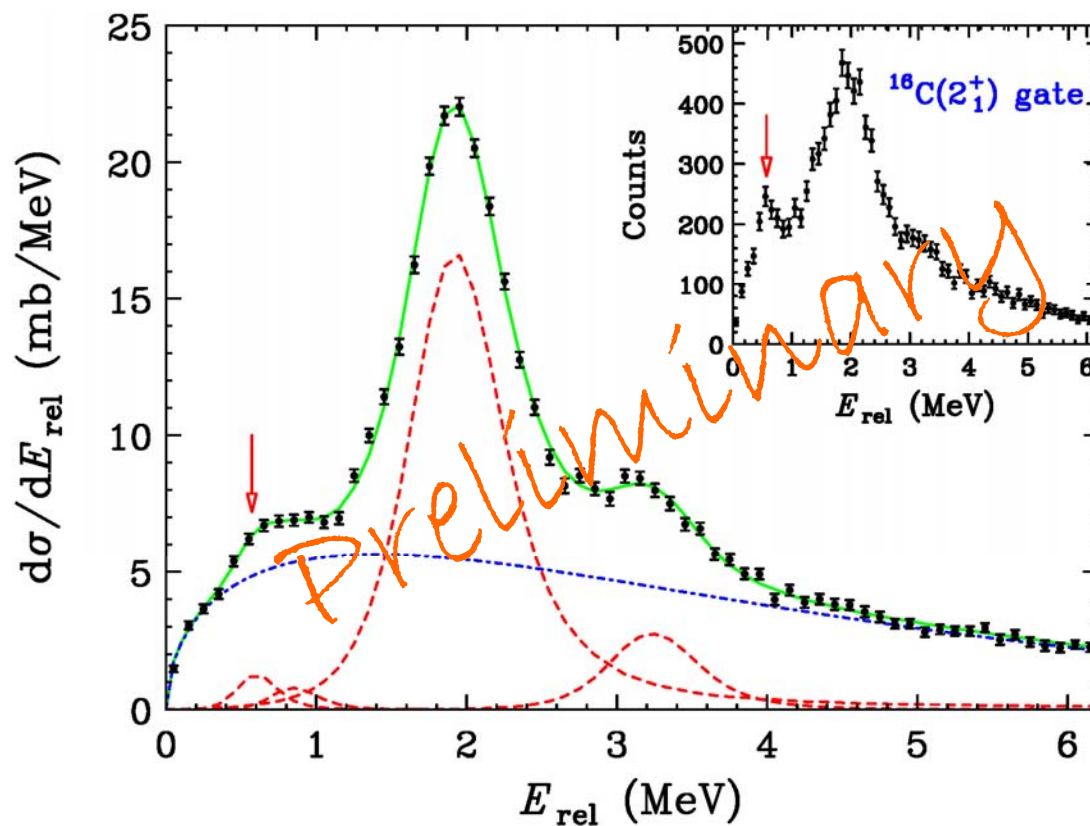


Longitudinal momentum distribution

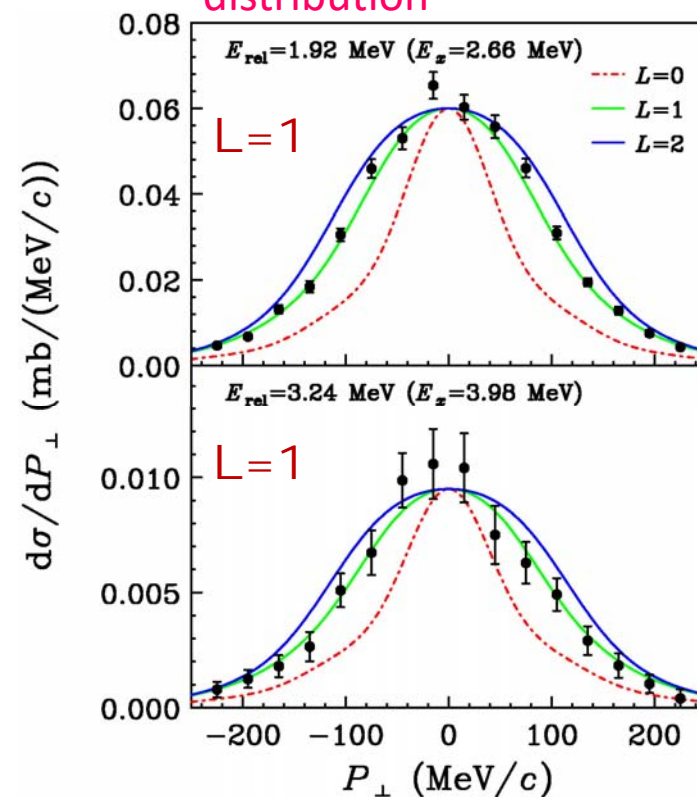


Momentum distributions were calculated by MOMDIS: C.A.Bertulani, A.Gade, Comput. Phys. Commun. 175(2006)372.

Results: $^{12}\text{C}(^{18}\text{C}, ^{17}\text{C}^* \rightarrow ^{16}\text{C} + n) @ 245 \text{ MeV/u}$



Transverse momentum distribution



Momentum distributions were calculated by **MOMDIS**: [C.A.Bertulani, A.Gade, Comput. Phys. Commun. 175\(2006\)372.](#)

Summary of parameters of observed resonances in $^{17}\text{C}^*$

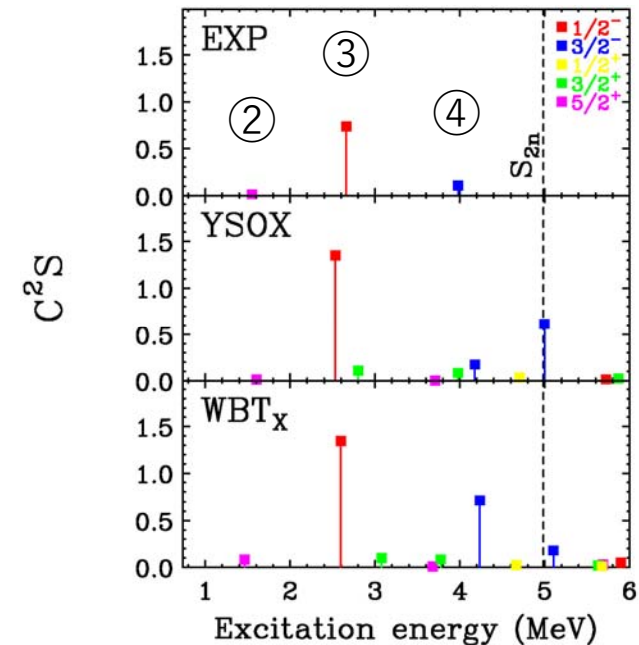
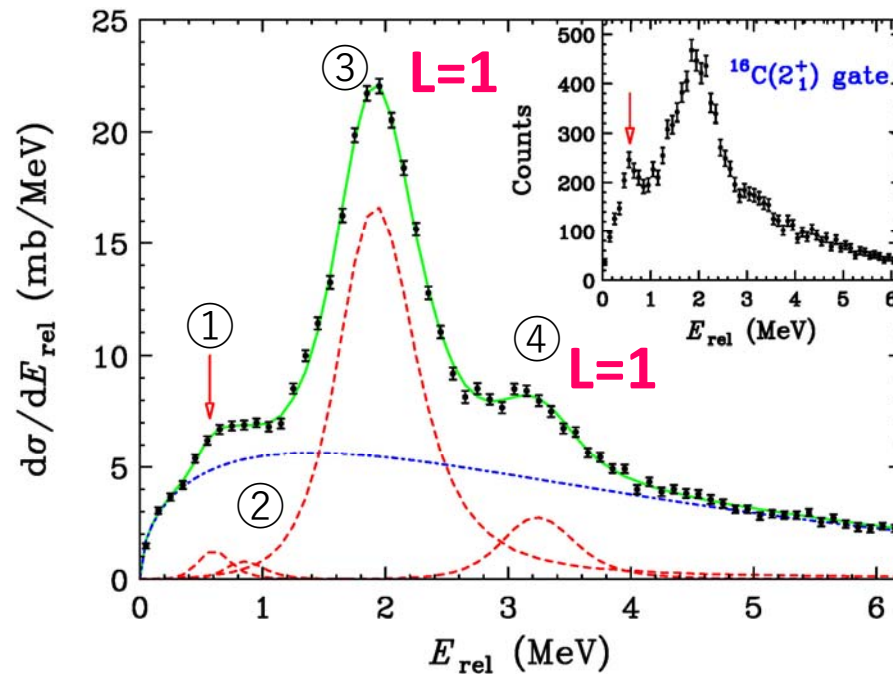
No.	E _{rel} (MeV)	Ex (MeV)	Γ (MeV)	L	$\sigma_{-1n}^{\text{exp}}$ (mb)	σ_{-1n}^{th} (mb)	C^2S^{exp}	$C^2S^{\text{th (b)}}$	$E_x^{\text{th (b)}}$ (MeV)	J^π
①	0.57(4) ^(a)	3.07(5)	0 (fixed)	---	0.38(7)	---	---	---	3.07	$9/2_1^+$ ^(c)
②	0.81(5)	1.55(6)	0 (fixed)	---	0.37(8)	0.36	0.015	0.015	1.60	$(5/2_2^+)$
③	1.92(1)	2.66(2)	0.32(1)	1	15.8(10)	29.0	0.736	1.350	2.53	$1/2_1^-$ ^(d)
④	3.24(2)	3.98(2)	0.03(6)	1	2.20(16)	3.58	0.107	0.174	4.18	$3/2_1^-$ ^(d)

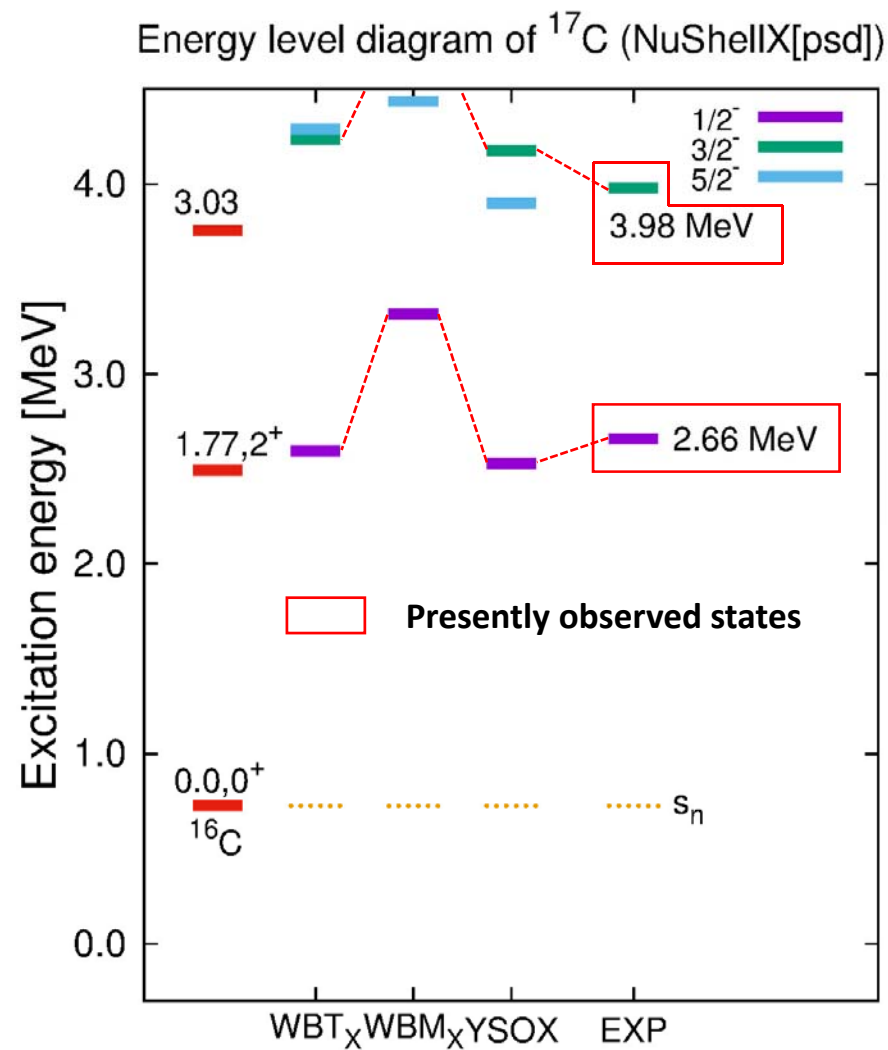
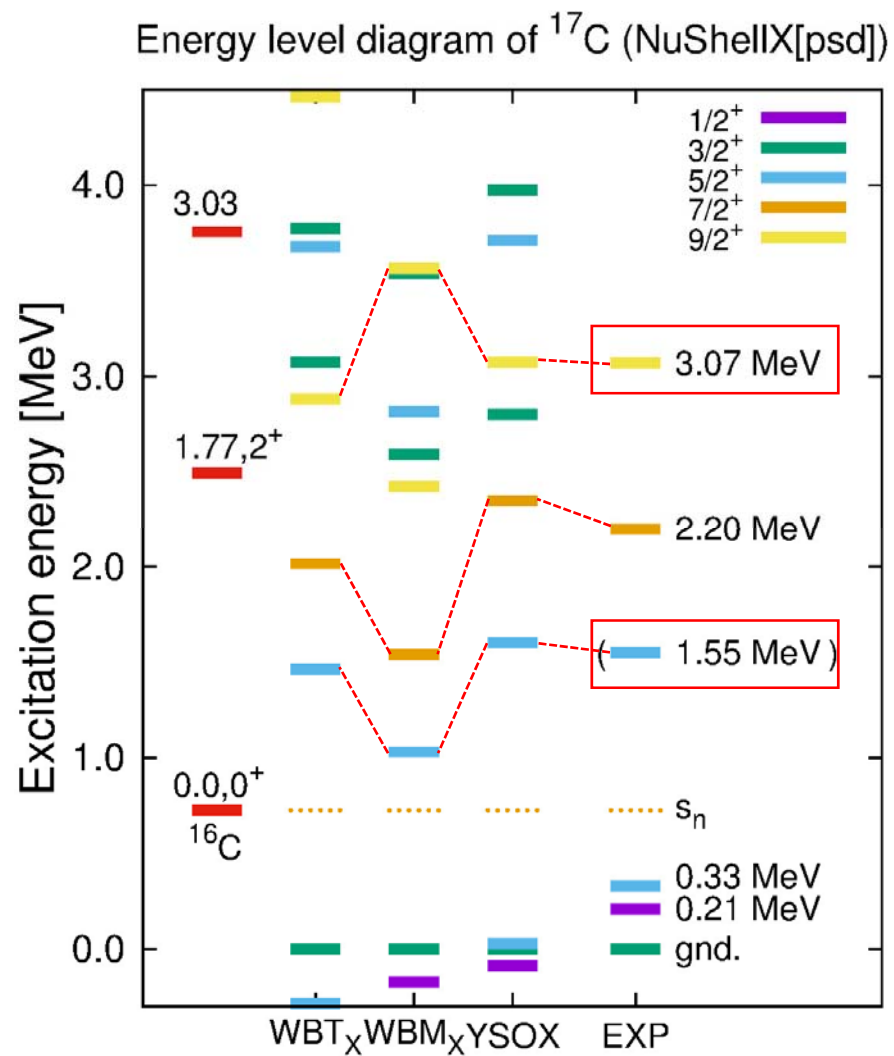
(a) γ -ray coincidence is observed.

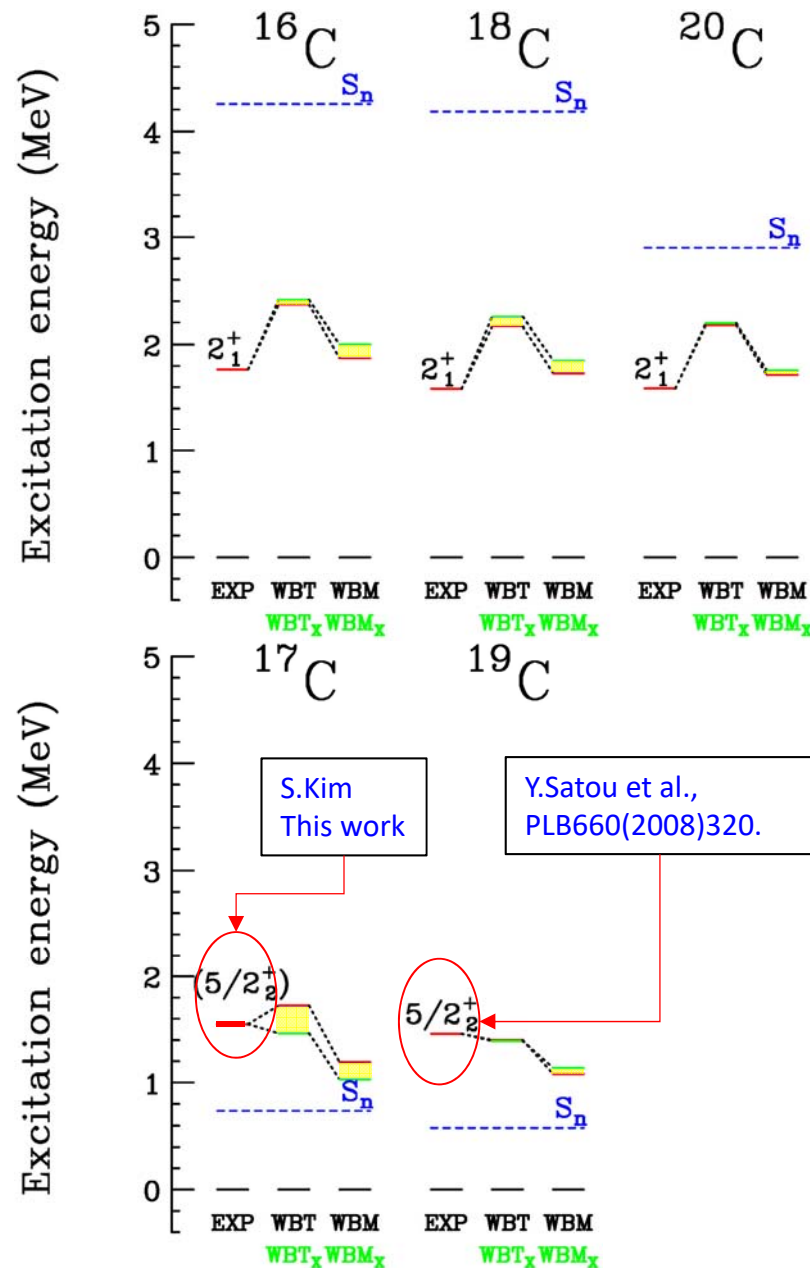
(b) From YSOX.

(c) H.G.Bohlen et al., EPJA31(2007)279. Y.Satou et al., PLB660(2008)320.

(d) H.Ueno et al., PRC87(2013)034316. Ex=2.71(2) MeV for $1/2_1^-$ and 3.93(2) MeV for $3/2_1^-$.





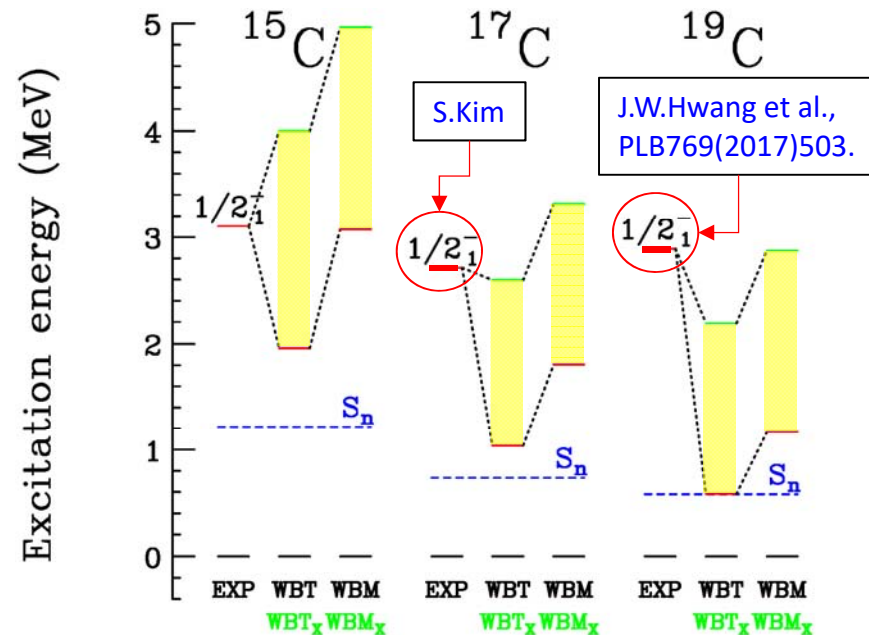


Discussion-1:

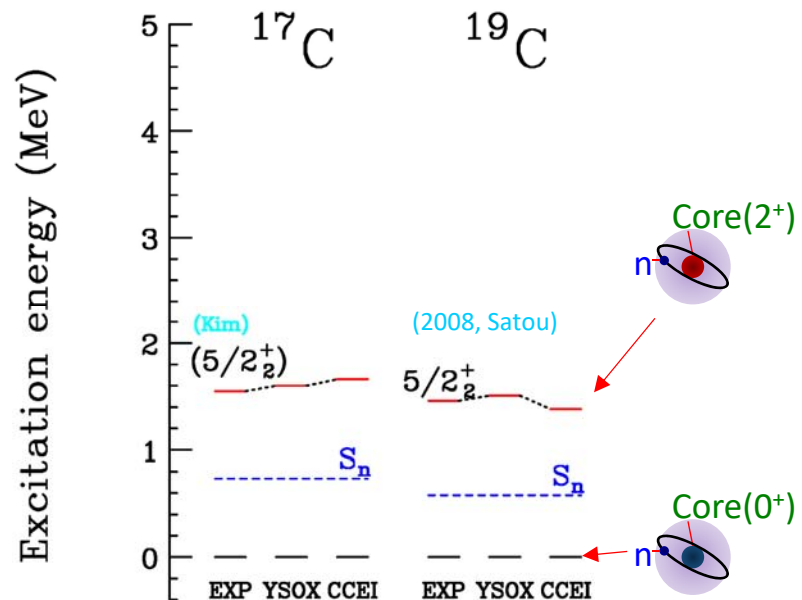
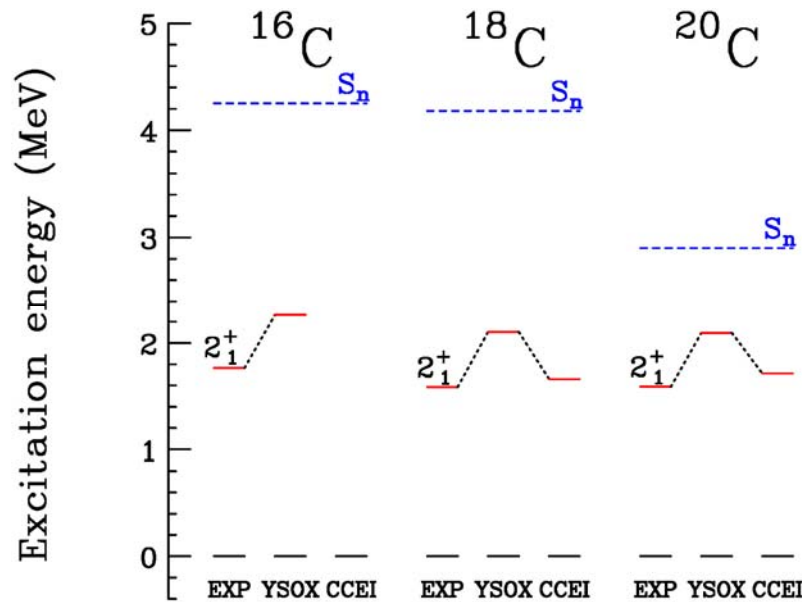
Spectroscopy of lowest-lying states,

- 2_1^+ in even **C** and
 - $5/2_2^+$ and $1/2_1^-$ in odd **C**,
- in comparison to **WBT_(x)**, **WBM_(x)** SM.

- The effects of enlarged $\hbar\omega$ model space by using **NuShellX** are small (large) for positive (negative) parity states.
- WBM** gives a deteriorated description of $5/2_2^+$.



Excitation energy is from $E(3/2_1^+)$ and $E(1/2_1^+)$ for ^{17}C and ^{19}C , respectively.

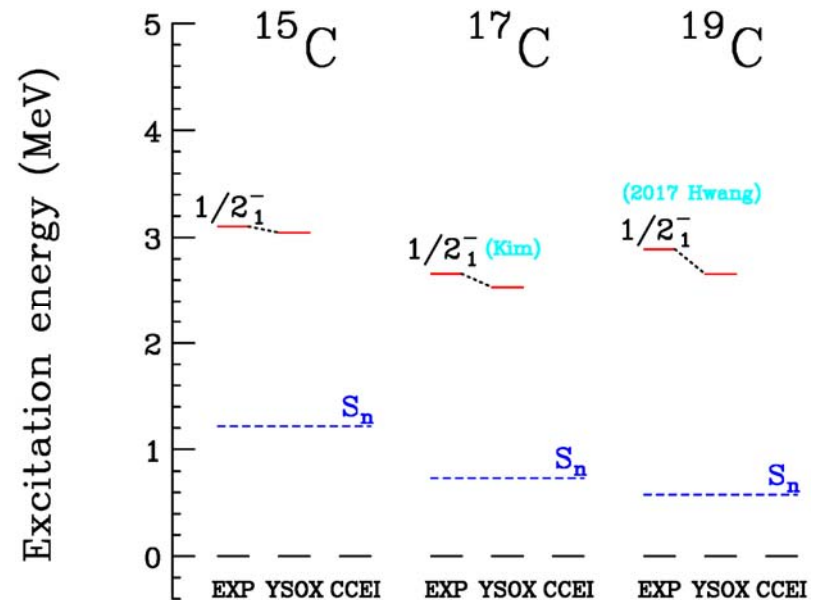


Discussion-2:

Spectroscopy of lowest-lying states,

- 2_1^+ in even **C** and
 - $5/2_2^+$ and $1/2_1^-$ in odd **C**,
- in comparison to **YSOX**, **CCEI** SM.

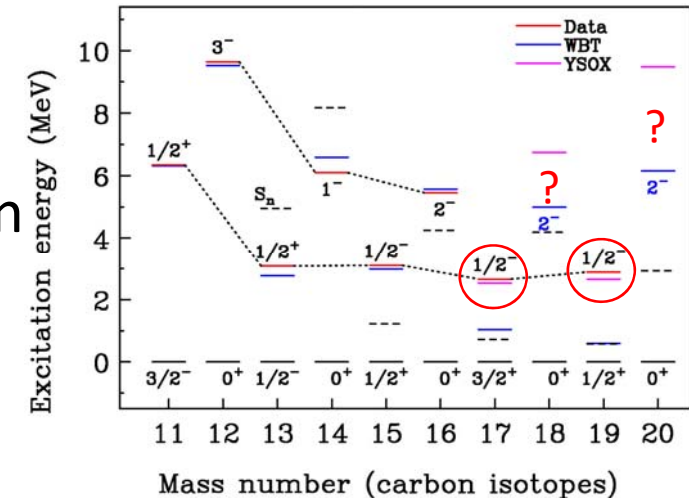
- **YSOX** does not solve the overprediction problem of 2_1^+ , while it provides a good description of $5/2_2^+$ and $1/2_1^-$.
- **CCEI** works well for both positive parity states, 2_1^+ and $5/2_2^+$.



Excitation energy is from $E(3/2_1^+)$ and $E(1/2_1^-)$ for ^{17}C and ^{19}C , respectively.

Summary

- Knockout reaction can give information which complements the βn study.
- First $1/2_1^-$ cross-shell states in C now extend to ^{19}C .
- New spectroscopy is used to test SM interactions.
 - **YSOX** describes cross shell states well.
 - **YSOX** overpredicts 2_1^+ in $^{16,18,20}\text{C}$, but **CCEI** not.
 - **YSOX** and **CCEI** describe the $5/2_2^+$ location in $^{17,19}\text{C}$ well.
- Light nuclear systems, as C isotopes, continues to pose challenges to our attempt to understand nuclear structures.



Deformations and electromagnetic moments in carbon and neon isotopesH. Sagawa,^{*} X. R. Zhou,[†] and X. Z. Zhang[‡]*Center for Mathematical Sciences, University of Aizu, Aizu-Wakamatsu, Fukushima 965-8560, Japan*

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(Received 2 July 2004; published 18 November 2004)

Electromagnetic observables will provide useful information to study the structure of nuclei, not only ground states but also excited states. Namely, these observables are expected to pin down precise information of deformations and unknown spin parities of both stable and unstable nuclei since the deformation is intimately related to observables such as Q moments and $E2$ transitions. The magnetic moments will also give empirical information of deformation in comparison with the single-particle value (Schmidt value) in the spherical limit. The isotope dependence of the deformation is an interesting subject to study in relation to a manifestation of the spontaneous symmetry breaking effect from the beginning of one closed shell to the next closed shell. To this end, carbon and neon might be promising in future experiments within the next few years. The effect of spontaneous symmetry breaking effect is a general phenomenon known in many fields of physics. In molecular physics, spontaneous symmetry breaking was discovered by Jahn and Teller in 1937 [6]. There is a very similar fundamental mechanism in nuclear physics and molecular physics, which is responsible for causing deformations in the ground state. The coupling to the quadrupole vibration is the main origin of the static deformation in atomic nuclei [7]. On the other

N.Imai et al., PRL92(2004)062501. "Anomalously hindered E2 strength $B(E2;2^{+} \rightarrow 0^{+})$ in ^{16}C "

Z.Elekes et al., PLB586(2004)34. "Decoupling of valence neutrons from the core in ^{16}C "

Z.Elekes et al., PLB614(2005)174. "Low-lying excited states in $^{17,19}\text{C}$ "

H.J.Ong et al., PRC73(2006)024610. "Neutron-dominant quadrupole collective motion in ^{16}C "

J.G.Bohlen et al., EPJA31(2007)279. "Spectroscopy of ^{17}C and (sd) 3 structures in heavy carbon isotopes"

H.J.Ong et al., PRC78(2008)014308. "Lifetime measurements of first excited states in $^{16,18}\text{C}$ "

M.Stanoiu et al., PRC78(2008)034315. "Disappearance of the N=14 shell gap in the carbon isotopic chain"

Y.Satou et al., PLB660(2008)320. "Unbound excited states in $^{19,17}\text{C}$ "

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