

## Axion dark matter search with the storage ring EDM method

➤ This work can be found in the arXiv: <http://arxiv.org/abs/1710.05271>

**Seongtae Park**

*Center for Axion and Precision Physics(CAPP)  
of Institute for Basic Science (IBS), South Korea*

On behalf of  
Seung Pyo Chang, Selcuk Haciomeroglu, On Kim,  
Soohyung Lee, Seongtae Park, Yannis K. Semertzidis

**IBS conference on Dark World  
Daejeon, Korea  
Oct. 30-Nov. 3, 2017**

# Outline

- 1. About storage ring EDM experiment**
- 2. About axion-coupled EDM**
- 3. srEDM method for Axion search**
- 4. About sensitivity of the experiment**
- 5. Summary and conclusion**

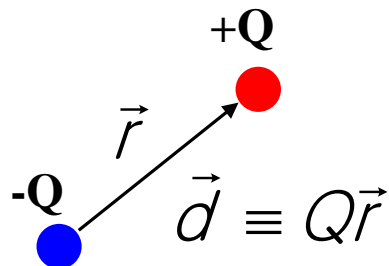
# Introduction to EDM

## ❖ Motivation

- ✓ Strong CP problem,  $\Theta_{QCD}$
- ✓ Matter-antimatter asymmetry (Baryogenesis)

## ➤ P, T violation due to an EDM

### Electric Dipole Moment



$$H = -\mu \vec{\sigma} \cdot \vec{B} - d \vec{\sigma} \cdot \vec{E}$$

$$T(\vec{B}, \vec{\sigma} \text{ sign change}) : H = -\mu \vec{\sigma} \cdot \vec{B} + d \vec{\sigma} \cdot \vec{E}$$

$$P(\vec{E} \text{ sign change}) : H = -\mu \vec{\sigma} \cdot \vec{B} + d \vec{\sigma} \cdot \vec{E}$$

$$\mu = g \frac{q}{2m}, d = \eta \frac{q}{2mc}$$

A nonzero particle EDM violates P, T, and assuming CPT conservation, also CP violation.

# The first direct measurement of EDM in neutron and it's sensitivity improvement

PHYSICAL REVIEW

VOLUME 108, NUMBER 1

OCTOBER 1, 1957

## Experimental Limit to the Electric Dipole Moment of the Neutron

J. H. SMITH,\* E. M. PURCELL, AND N. F. RAMSEY

*Oak Ridge National Laboratory, Oak Ridge, Tennessee, and Harvard University, Cambridge, Massachusetts*

(Received May 17, 1957)

1957

An experimental measurement of the electric dipole moment of the neutron by a neutron-beam magnetic resonance method is described. The result of the experiment is that the electric dipole moment of the neutron equals the charge of the electron multiplied by a distance  $D = (-0.1 \pm 2.4) \times 10^{-20}$  cm. Consequently, if an electric dipole moment of the neutron exists and is associated with the spin angular momentum, its magnitude almost certainly corresponds to a value of  $D$  less than  $5 \times 10^{-20}$  cm.

PRL 97, 131801 (2006)

PHYSICAL REVIEW LETTERS

week ending  
29 SEPTEMBER 2006

## Improved Experimental Limit on the Electric Dipole Moment of the Neutron

C. A. Baker,<sup>1</sup> D. D. Doyle,<sup>2</sup> P. Geltenbort,<sup>3</sup> K. Green,<sup>1,2</sup> M. G. D. van der Grinten,<sup>1,2</sup> P. G. Harris,<sup>2</sup> P. Iaydjiev,<sup>1,\*</sup> S. N. Ivanov,<sup>1,†</sup> D. J. R. May,<sup>2</sup> J. M. Pendlebury,<sup>2</sup> J. D. Richardson,<sup>2</sup> D. Shiers,<sup>2</sup> and K. F. Smith<sup>2</sup><sup>1</sup>*Rutherford Appleton Laboratory, Chilton, Didcot, Oxon OX11 0QX, United Kingdom*<sup>2</sup>*Department of Physics and Astronomy, University of Sussex, Falmer, Brighton BN1 9QH, United Kingdom*<sup>3</sup>*Institut Laue-Langevin, BP 156, F-38042 Grenoble Cedex 9, France*

(Received 9 February 2006; revised manuscript received 29 March 2006; published 27 September 2006)

2006

An experimental search for an electric dipole moment (EDM) of the neutron has been carried out at the Institut Laue-Langevin, Grenoble. Spurious signals from magnetic-field fluctuations were reduced to insignificance by the use of a cohabiting atomic-mercury magnetometer. Systematic uncertainties, including geometric-phase-induced false EDMs, have been carefully studied. The results may be interpreted as an upper limit on the neutron EDM of  $|d_n| < 2.9 \times 10^{-26} e \text{ cm}$  (90% C.L.).

DOI: 10.1103/PhysRevLett.97.131801

PACS numbers: 13.40.Em, 07.55.Ge, 11.30.Er, 14.20.Dh

# Current EDM bounds and plan

## ❖ Current EDM bounds

- ✓ SM predicts non-vanishing EDM
  - $|d_e| < 10^{-38} \text{ e.cm}$
  - $|d_{n,p}| < 10^{-31} \text{ e.cm}$
  - Beyond current experiment limit
- ✓ SUSY prediction:  $10^{-25} \sim 10^{-28} \text{ e} \cdot \text{cm}$  (nEDM limit)
- ✓ Neutron EDM bound:  $|d_n| < 2.9 \times 10^{-26} \text{ e.cm}$  ('06, ultracold neutrons)
- ✓ Proton EDM bound:  $|d_p| < 7.9 \times 10^{-25} \text{ e.cm}$  ('09,  $^{179}\text{Hg}$ )
- ✓ Electron EDM bound:  $|d_e| < 8.7 \times 10^{-29} \text{ e. cm}$  ('14,  $^{205}\text{Tl}$ )

## ❖ Target sensitivity level in the storage ring pEDM experiment

- ✓ High statistics ( $10^{11}$  protons/store) is achievable using storage ring
- ✓ Goal  **$10^{-29} \text{ e} \cdot \text{cm}$**  (statistical limit in about one year)
- ✓  $\rightarrow 10^{-30} \text{ e} \cdot \text{cm}$  (with an upgrade)

## ❖ Physics reach $> 10^3 \text{ TeV}$

# EDM and spin precession

## ❖ Spin dynamics (with EDM and MDM) in magnetic+electric field (T-BMT equation)

$$\frac{d\vec{s}}{dt} = \vec{\mu} \times \vec{B} + \vec{d} \times \vec{E} = \vec{s} \times (\vec{\omega}_s + \vec{\omega}_{edm}) \quad (\text{for particle at rest})$$

$$\mu = (ge/2m)s = ge\hbar/4m, \quad d = (\eta e/2mc)s = \eta e\hbar/4mc$$

$$\text{with } \vec{\beta} \cdot \vec{E} = \vec{\beta} \cdot \vec{B} = 0$$

(T-BMT equation: for moving particle)

$$\vec{\omega} = -\frac{e}{m} \left[ a\vec{B} - \left( a - \frac{m}{p} \right) \frac{\vec{\beta} \times \vec{E}}{c} + \frac{\eta}{2} \left( \frac{\vec{E}}{c} + \vec{\beta} \times \vec{B} \right) \right], \quad a = \frac{g-2}{2}$$

MDM in B-field

MDM in induced B-field

EDM term

$\vec{\omega}_{edm}$

$\vec{\omega}_a$

# Storage ring technique for EDM search

g-2 precession in **pure electric ring**

$$\Rightarrow \vec{\omega}_a = \frac{e}{m} \left[ a - \left( \frac{m}{p} \right)^2 \right] \frac{\vec{\beta} \times \vec{E}}{c}$$

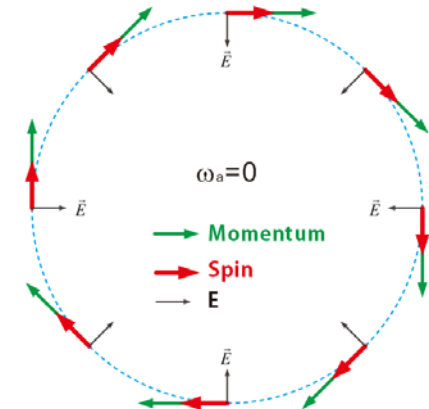
for the proton ( $a=1.792847357(23)$ )  $>0$

$$\vec{\omega}_a = 0 \quad \text{at} \quad p = \frac{m}{\sqrt{a}} = 0.700740 \text{ GeV}/c : \text{magic momentum}$$

## ❖ Use frozen spin method (static EDM measurement)

- ✓ Spin freezes to the momentum direction
- ✓  $a>0$  particles (p,e, etc.): use magic momentum
- ✓  $a<0$  (deuteron) : use E+B field
- ✓ spin precesses only on the vertical plane!
- ✓ No precession on the ring plane

$$\vec{\omega}_{EDM} = -\frac{e\eta}{2m} \frac{\vec{E}}{c}$$



## ❖ Storage ring EDM collaboration is trying to establish the experiment at CERN

✓ **EDM Kick-off meeting at CERN in Mar. 2017**

✓ About 50 participants



# Polarimeter and asymmetry

## ❖ Use asymmetrical proton scattering on Carbon target

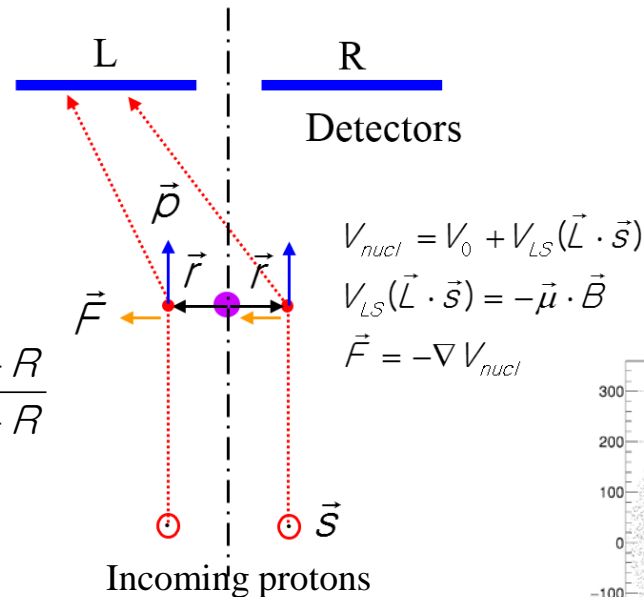
- Hadronic elastic scattering (spin-orbit interaction)
- Asymmetrical proton hit distribution on the detector plane
- L/R (U/D) asymmetry for vertical (horizontal) component of proton polarization

➤ For spin 1/2 particle

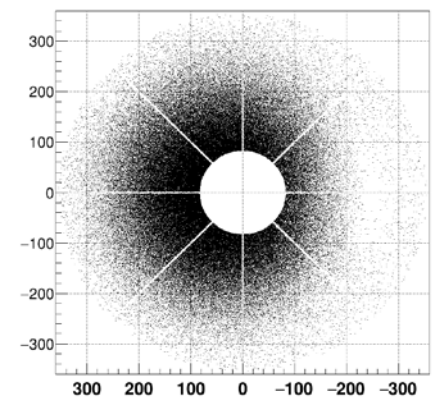
$$\sigma(\theta) = \sigma_{unpol}(\theta) [1 + P_y A_y]$$

$$L/R \text{ asymmetry } \varepsilon_{LR} = P_y A_y = \frac{L - R}{L + R}$$

$$P_y = \frac{1}{A_y} \frac{L - R}{L + R}$$



Hit Positions on GEM (Primary Protons)

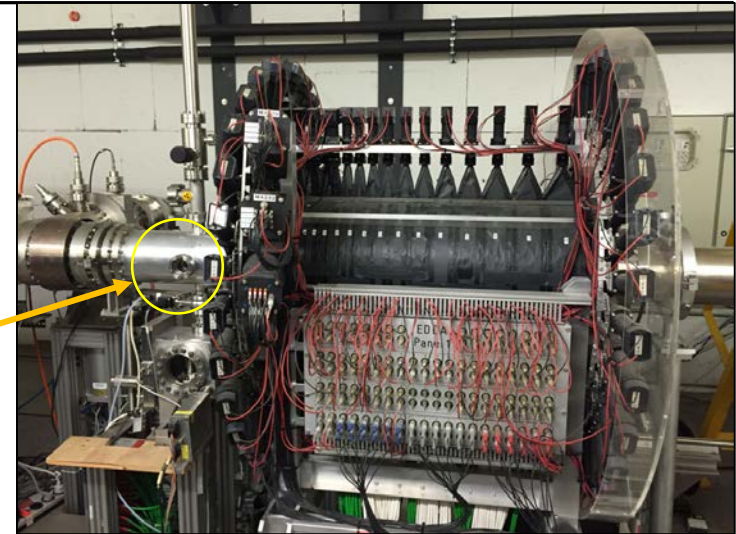
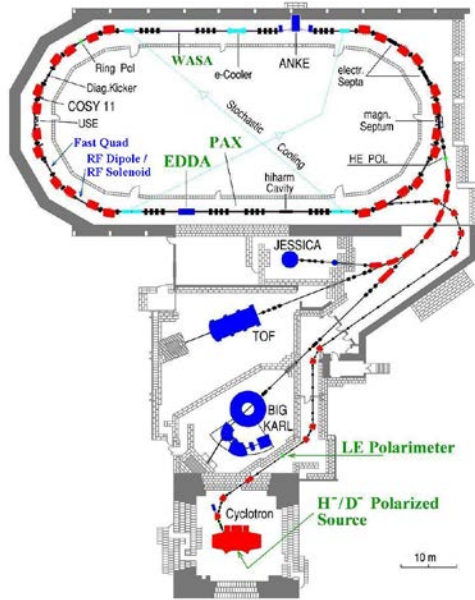


Simulation by Hoyong Jeong

- $P_y$  is calculated from the asymmetry with known  $A_y$
- $P_y$  changes in time due to the precession in E field.



# COSY ring and EDDA polarimeter

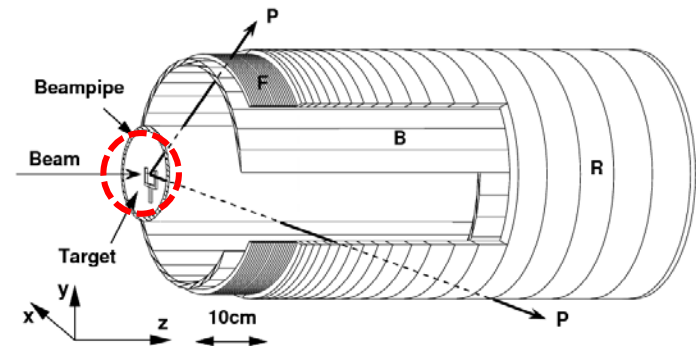


COSY carbon tube target

## COSY ring (Juelich, Germany)



Carbon block target (17 mm thick)



COSY EDDA detector  
FZJ, Juelich, Germany

# About Axion-coupled oscillating EDM

❖ The oscillating axion field is coupled with

✓ Photon

$$P_{sig} = \eta_{g_{a\gamma\gamma}}^2 \left( \frac{\rho_a}{m_a} \right) B_0^2 V C Q_L \quad a \rightarrow \gamma\gamma \quad \text{Searched by resonance microwave cavity experiment (IBS/CAPP, ADMX)}$$

✓ **gluons**, fermions, nucleon, etc. → Oscillating EDM

$$d_n = 2.4 \times 10^{-16} \frac{a}{f_a} e \cdot cm \approx 9 \times 10^{-35} \cos(m_a t) [e \cdot cm]$$

$$a(t) = a_0 \cos(m_a t) \Rightarrow d(t) = d_{DC} + d_{AC} \cos(m_a t + \varphi_x)$$

# Axion coupled EDM measurement using storage ring method

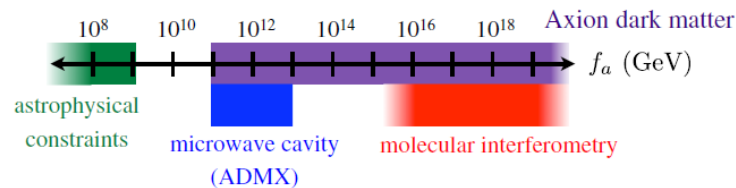
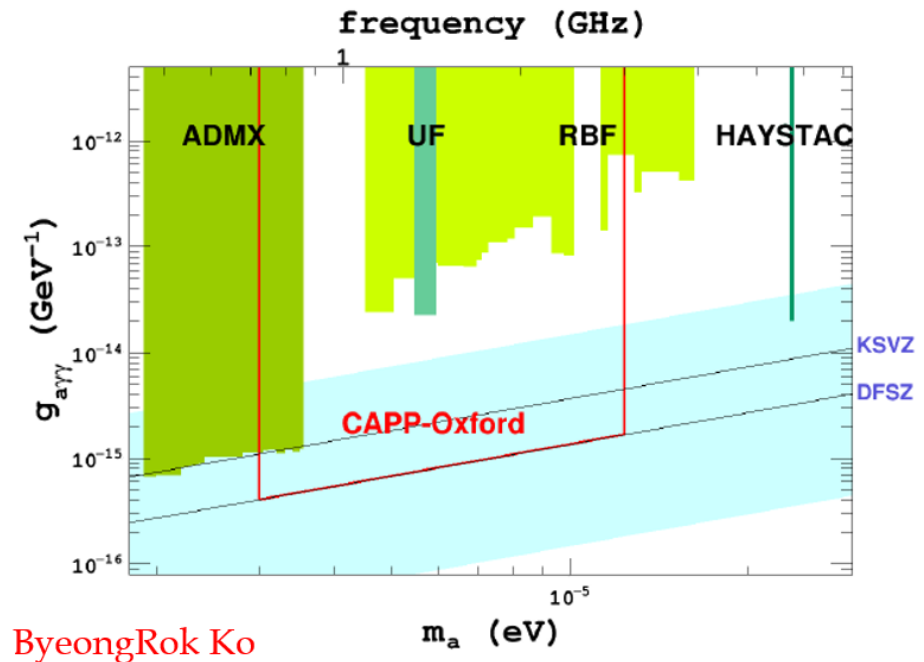


FIG. 1 (color online). (Color Online) The parameter space of the axion in  $f_a$  (GeV). Values of  $f_a < 10^9$  GeV are ruled out by astrophysical constraints (green). Values of  $f_a \geq 10^{11}$  GeV allow the axion to be the dark matter (purple). The (blue) region labeled “microwave cavity” shows the region of parameter space that is potentially observable with microwave cavity experiments, e.g. ADMX. The (red) region labeled “molecular interferometry” shows the range of  $f_a$  which is potentially observable with the experiments proposed here. The lower limit on this region may in fact be lower than shown, depending on technological advances.

Peter W. Graham, Surjeet Rajendran,  
Phys. Rev. D 84, 055013 (2011).

“Use molecular interferometry”



Experiments on axion search

## ❖ Our proposal: “Use storage ring”

- Scan frequency range up to 0.1 kHz~100 MHz using storage ring method.
- axion parameter space:  $10^{13} \text{ GeV} \leq f_a \leq 10^{19} \text{ GeV}$

# Axion + Storage ring EDM method

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Using storage ring EDM method for Axion search

# Spin tracking for AxionEDM search

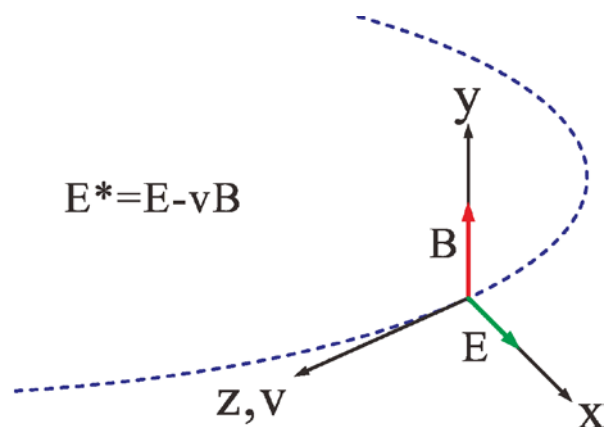
In particle rest frame

$$\frac{d\vec{s}}{dt} = \vec{\mu} \times \vec{B} + \vec{d} \times \vec{E} = \vec{s} \times (\vec{\omega}_a + \vec{\omega}_{edm})$$

$$\mu = (ge/2m)s = ge\hbar/4m, \quad d = (\eta e/2mc)s = \eta e\hbar/4mc$$

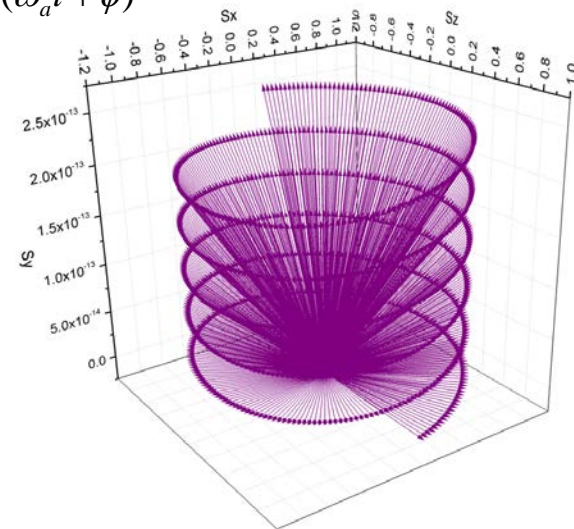
$$\frac{d\vec{\beta}}{dt} = \frac{e}{\gamma m} \left[ \vec{\beta} \times \vec{B} + \frac{\vec{E}}{c} - \vec{\beta} \frac{\vec{\beta} \cdot \vec{E}}{c} \right]$$

$$\frac{d\vec{s}}{dt} = \frac{e}{m} \vec{s} \times \left[ \left( \frac{g}{2} - \frac{\gamma-1}{\gamma} \right) \vec{B} - \left( \frac{g-2}{2} \frac{\gamma}{\gamma+1} \right) (\vec{\beta} \cdot \vec{B}) \vec{\beta} - \left( \frac{g}{2} - \frac{\gamma}{\gamma+1} \right) \frac{\vec{\beta} \times \vec{E}}{c} + \frac{\eta}{2} \left( \frac{\vec{E}}{c} - \frac{\gamma}{\gamma+1} \left( \frac{\vec{\beta} \cdot \vec{E}}{c} \right) \vec{\beta} + \vec{\beta} \times \vec{B} \right) \right]$$

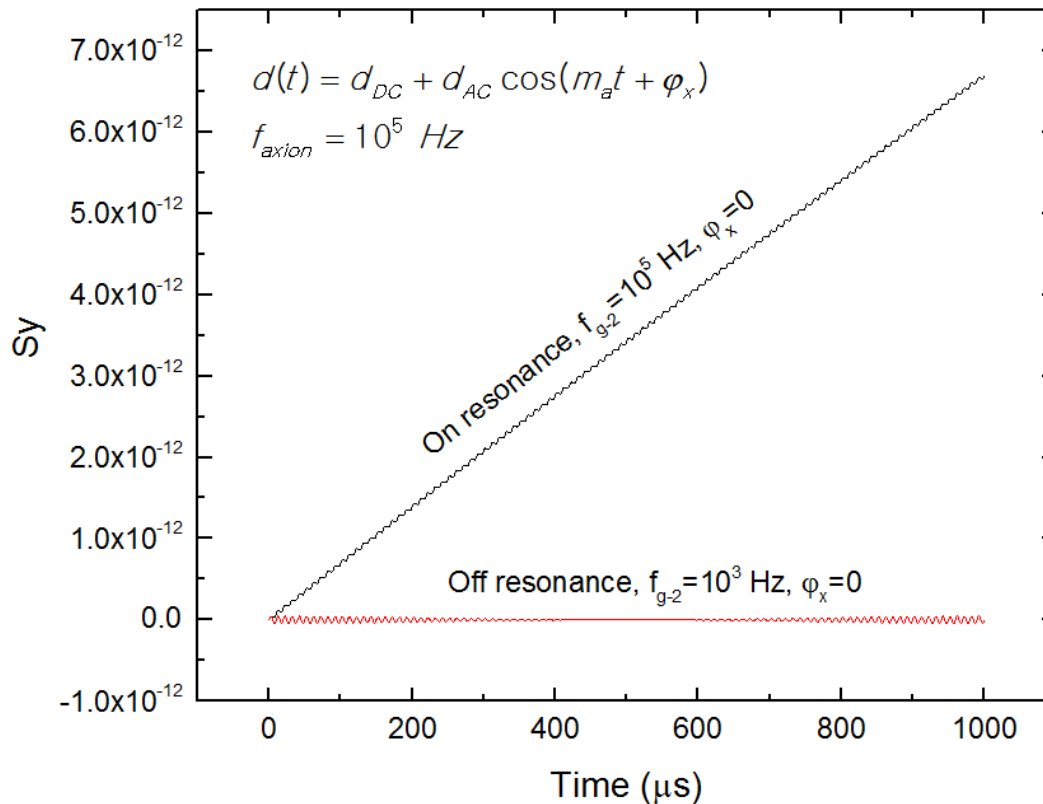


$$d(t) = d_{DC} + d_{AC} \cos(\omega_a t + \phi)$$

$$\eta(t) = \frac{d(t)4mc}{e\hbar}$$



# Resonance of Axion induced oscillating EDM with g-2 precession in storage ring



$$\frac{d\vec{s}}{dt} = \vec{\mu} \times \vec{B} + \vec{d} \times \vec{E} = \vec{s} \times (\vec{\omega}_a + \vec{\omega}_d)$$

$$\vec{\omega}_a = -\frac{e}{m} \left[ a\vec{B} - \left( a - \frac{1}{\gamma^2 - 1} \right) \frac{\vec{\beta} \times \vec{E}}{c} \right]$$

$$\vec{\omega}_d = -\frac{e}{m} \left[ \frac{\eta}{2} \left( \frac{\vec{E}}{c} + \vec{\beta} \times \vec{B} \right) \right]$$

$$\omega_d = \frac{d\theta}{dt} = -\frac{2d}{\hbar} E^*, \quad E^* = E + c\beta B$$

Benefit from large  
effective E field!

# Static EDM and AxionEDM experiments

## ❖ Static EDM measurement

- ✓ Use frozen spin method
- ✓ Spin precesses only on the vertical plane
- ✓ No precession on the ring plane

$$\vec{\omega}_{EDM} = -\frac{e\eta}{2m} \frac{\vec{E}}{c}$$

## ❖ AxionEDM measurement uses non-zero g-2 precession

- ✓ Benefit from larger effective E-field ( $\vec{v} \times \vec{B}$ )
- ✓ More sensitive

$$\omega_d = \frac{d\theta}{dt} = -\frac{2d}{\hbar} E^*, \quad E^* = E + c\beta B \quad \leftarrow \text{Also used for muon g-2/EDM experiment}$$

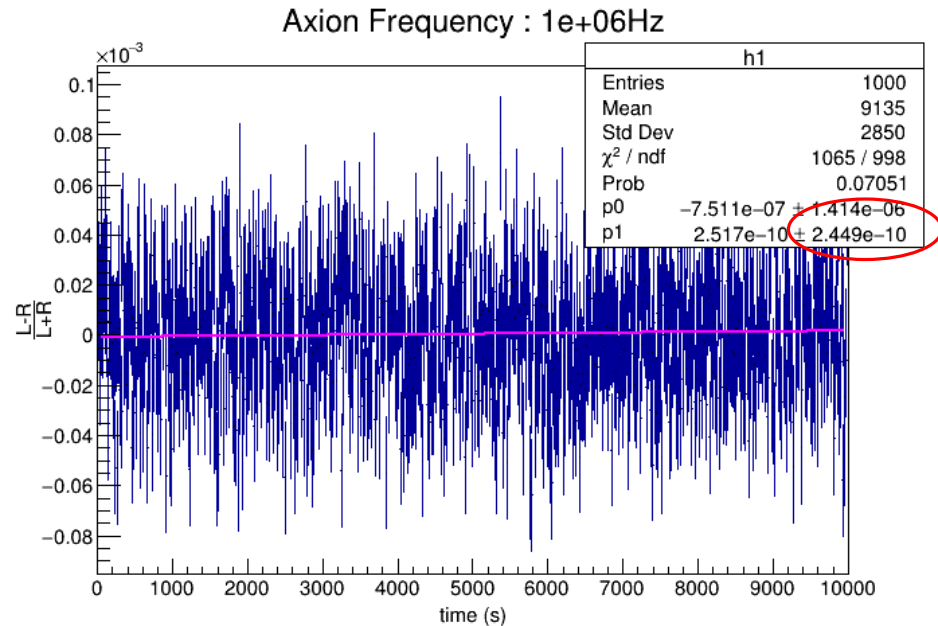
# Asymmetry simulation and sensitivity calculation

Error for the EDM d

$$\omega_d = \frac{d\theta}{dt} = -\frac{2d}{\hbar} E^*$$

$$E^* = E + c\beta B \quad \leftarrow E, B=?$$

$$\sigma_d = \frac{\hbar}{2E^*} \sigma_{\omega_d}$$



Simulated L/R asymmetry with linear fit

$$L/R \text{ asymmetry } \varepsilon_{LR} = P_y A_y = \frac{L - R}{L + R}$$

$$P_y = \frac{1}{A_y} \frac{L - R}{L + R}$$



# EDM precession and statistical error

## Statistical error for EDM $d$ (sensitivity) (static (DC) EDM measurement)

$$\sigma_d = \frac{2\hbar}{PAE\sqrt{N_{\text{tot}}fT_{\text{tot}}\tau_p}}$$

P: Degree of polarization (0.8)

A: Analyzing power (0.6)

E\*: Effective electric field (8MV/m)

$N_{\text{tot}}$  = Number of particles/storage ( $5 \times 10^{10}$ )

$f$ : Effective detection efficiency of polarimeter (0.011)

$T_{\text{tot}}$ : Total running time ( $10^7$  s/year)

$\tau_p$ : Spin coherence time ( $10^3$  s)

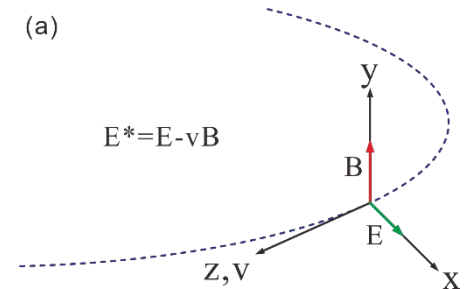
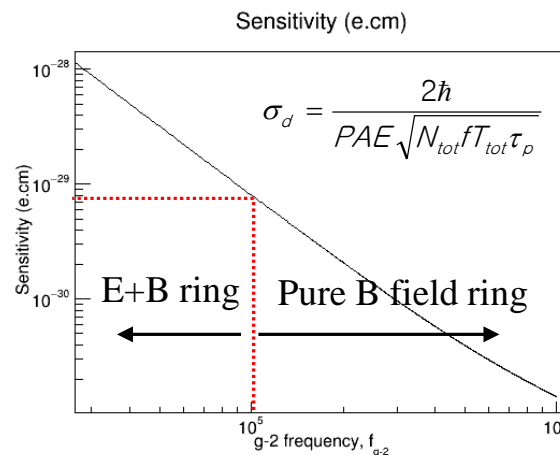
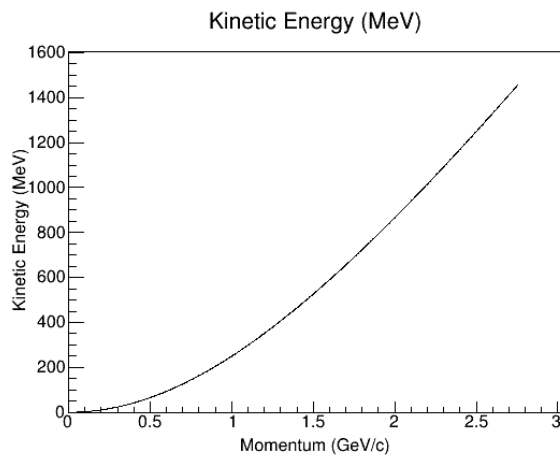
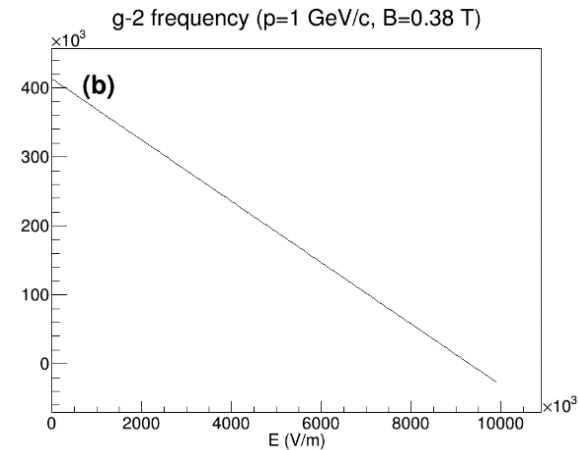
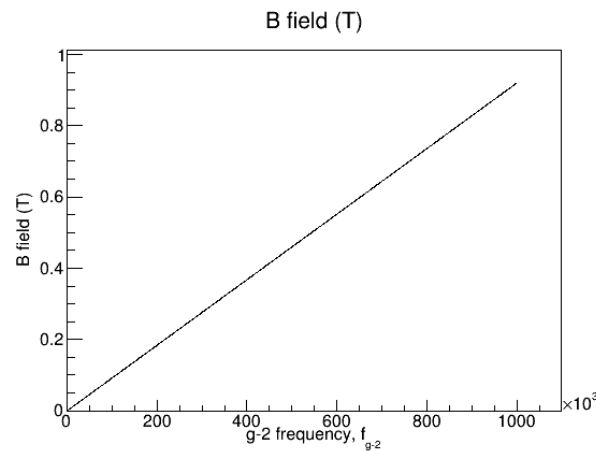
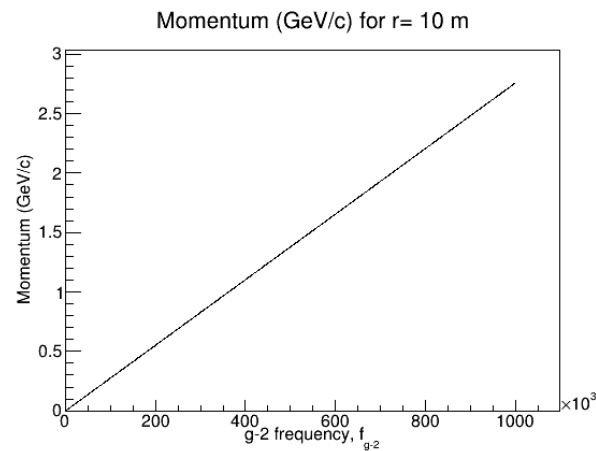
for storage ring proton EDM experiment

$\rightarrow \sim 10^{-29} \text{ e}\cdot\text{cm}$

# Determination of storage ring parameters

Pure magnetic ring for deuteron,  $a=-0.14$

Ring bending radius  $r=10\text{m}$



Use E+B combined ring  
to improve the sensitivity

$$E^* = E + c\beta B$$

# Sensitivity calculation results for **deuteron**

B (T)	P (GeV/c)	$f_G$ (Hz)	$E_r$ (V/m)	$E^*$ (V/m)	Sensitivity (e·cm)		Ring
					a	b	
0.38	0.9429	$10^2$	$8.82 \times 10^6$	$4.23 \times 10^7$	$1.9 \times 10^{-31}$	$1.9 \times 10^{-31}$	E/B ring ( $r = 10$ m)
0.38	0.9433	$10^3$	$8.80 \times 10^6$	$4.24 \times 10^7$	$6.0 \times 10^{-31}$	$1.9 \times 10^{-31}$	
0.38	0.9473	$10^4$	$8.65 \times 10^6$	$4.27 \times 10^7$	$1.9 \times 10^{-30}$	$1.9 \times 10^{-31}$	
<b>0.38</b>	<b>0.988</b>	<b><math>10^5</math></b>	<b><math>7.05 \times 10^6</math></b>	$4.60 \times 10^7$	$5.5 \times 10^{-30}$	<b><math>1.8 \times 10^{-31}</math></b>	
0.38	1.035	$2 \times 10^5$	$5.06 \times 10^6$	$5.00 \times 10^7$	$7.2 \times 10^{-30}$	$1.6 \times 10^{-31}$	
0.38	1.133	$4 \times 10^5$	$3.47 \times 10^5$	$5.86 \times 10^7$	$8.7 \times 10^{-30}$	$1.4 \times 10^{-31}$	
0.38	1.239	$6 \times 10^5$	$-5.47 \times 10^6$	$6.83 \times 10^7$	$9.1 \times 10^{-30}$	$1.2 \times 10^{-31}$	
0.38	1.355	$8 \times 10^5$	$-1.26 \times 10^7$	$7.93 \times 10^7$	$9.1 \times 10^{-30}$	$1.0 \times 10^{-31}$	
0.38	1.484	$10^6$	$-2.14 \times 10^7$	$9.21 \times 10^7$	$8.8 \times 10^{-30}$	$8.8 \times 10^{-31}$	
0.80	2.513	$10^6$	$-9.13 \times 10^6$	$2.01 \times 10^8$	$4.0 \times 10^{-30}$	$4.0 \times 10^{-31}$	B ring ( $r = 10$ m)
0.9198	2.7574	$10^6$	0	$2.28 \times 10^8$	$3.5 \times 10^{-30}$	$3.5 \times 10^{-31}$	
9.1977	27.574	$10^7$	0	$2.75 \times 10^9$	$9.3 \times 10^{-31}$	$9.3 \times 10^{-31}$	
a : Axion $Q = 10^6$ , Polarimeter Efficiency = 0.02, Initial polarization = 0.8, Analyzing power A=0.36, SCT = $10^4$ s. b : Axion $Q = 10^{10}$ , Polarimeter Efficiency = 0.02, Initial polarization = 0.8, Analyzing power A=0.36, SCT = $10^4$ s.							

➤ g-2 frequency tune by E-field

# Sensitivity calculation results for **proton**

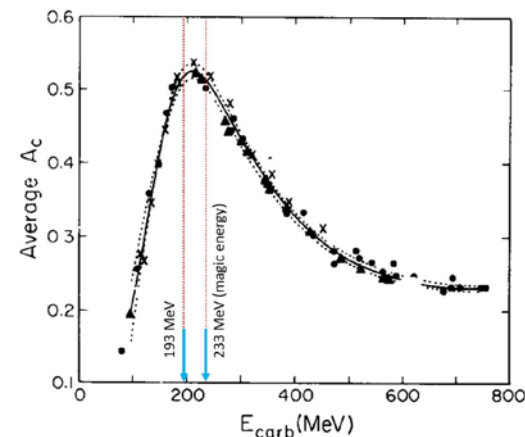
B (T)	P (GeV/c)	$f_G$ (Hz)	$E_r$ (V/m)	$E^*$ (V/m)	Sensitivity (e-cm)		Ring
					a	b	
0.00010	0.6984	$10^2$	$-8.00 \times 10^6$	$8.02 \times 10^6$	$1.0 \times 10^{-30}$	$1.0 \times 10^{-30}$	E/B ring ( $r = 52$ m)
0.00008	0.6982	$10^3$	$-8.00 \times 10^6$	$8.01 \times 10^6$	$3.2 \times 10^{-30}$	$1.0 \times 10^{-30}$	
-0.00017	0.6964	$10^4$	$-8.00 \times 10^6$	$7.97 \times 10^6$	$1.0 \times 10^{-29}$	$1.0 \times 10^{-30}$	
-0.00243	0.6747	$10^5$	$-8.00 \times 10^6$	$7.57 \times 10^6$	$3.4 \times 10^{-29}$	$1.1 \times 10^{-30}$	
-0.00495	0.6519	$2 \times 10^5$	$-8.00 \times 10^6$	$7.15 \times 10^6$	$5.0 \times 10^{-29}$	$1.1 \times 10^{-30}$	
-0.01523	0.7103	$4 \times 10^5$	$-1.10 \times 10^7$	$8.24 \times 10^6$	$6.2 \times 10^{-29}$	$9.8 \times 10^{-31}$	
-0.02002	0.6711	$6 \times 10^5$	$-1.10 \times 10^7$	$7.51 \times 10^6$	$8.3 \times 10^{-29}$	$1.1 \times 10^{-30}$	
-0.02666	0.6643	$8 \times 10^5$	$-1.20 \times 10^7$	$7.38 \times 10^6$	$9.8 \times 10^{-29}$	$1.1 \times 10^{-30}$	
-0.03327	0.6583	$10^6$	$-1.30 \times 10^7$	$7.27 \times 10^6$	$1.1 \times 10^{-28}$	$1.1 \times 10^{-30}$	
0.36587	1.0968	$10^7$	0	$8.33 \times 10^7$	$3.1 \times 10^{-29}$	$3.1 \times 10^{-31}$	B ring ( $r = 10$ m)
3.65868	10.9684	$10^8$	0	$1.09 \times 10^9$	$7.4 \times 10^{-30}$	$7.4 \times 10^{-32}$	

a : Axion  $Q = 10^6$ , Polarimeter Efficiency = 0.02, Initial polarization = 0.8, SCT =  $10^4$  s.  
b : Axion  $Q = 10^{10}$ , Polarimeter Efficiency = 0.02, Initial polarization = 0.8, SCT =  $10^4$  s.  
Analyzing power  $A$  :  $A = 0.6$  for E/B ring,  $A = 0.25$  for B ring

## ➤ g-2 frequency tune by B-field

$$\vec{\omega}_a = -\frac{e}{m} \left[ a\vec{B} - \left( a - \frac{1}{\gamma^2 - 1} \right) \frac{\vec{\beta} \times \vec{E}}{c} \right]$$

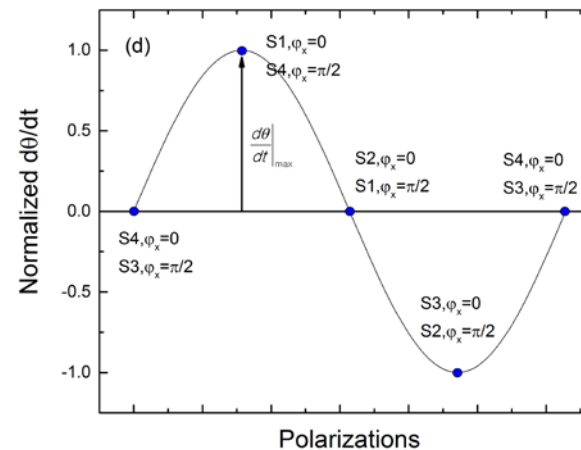
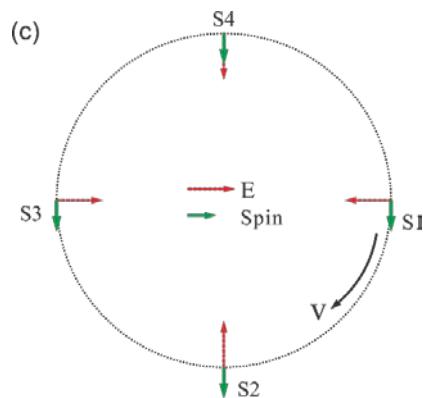
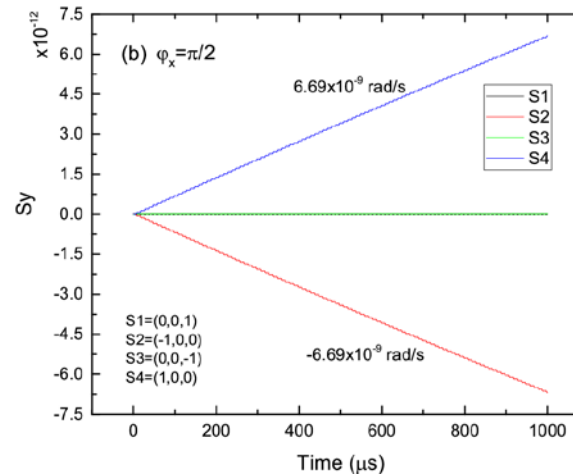
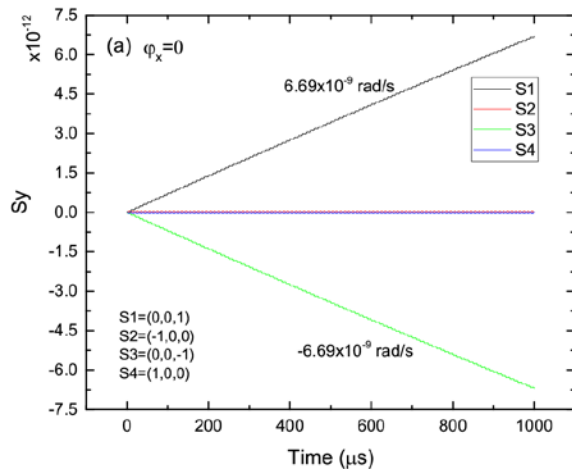
- $a_d = -0.14$ ,  $a_p = 1.79$
- $m_d \sim 2m_p$



p-C analyzing power

# Axion phase effect

$d(t) = d_{DC} + d_{AC} \cos(\omega_a t + \varphi_x) \rightarrow \varphi_x=0$  means  $d$  is parallel to  $s$ , positive and max.



- Axion phase is not an experimentally controllable parameter

# Axion Q and sensitivity

## ❖ Resonance cavity experiment

- ✓ No benefit from the high axion  $Q$  value if the cavity  $Q_L$  is smaller than the axion  $Q$  value.

### ➤ P. Sikivie's Haloscope

Axion conversion power ( $\sim 10^{-24} \text{W}$ )

$$P_{a \rightarrow \gamma\gamma} = g_{a\gamma\gamma}^2 \frac{\rho_a}{m_a} B^2 V C_{mnp} \min(Q_L, Q_a)$$

## ❖ Storage ring method

- ✓ No limit with axion  $Q_a$  value
- ✓ Longer measurement time with larger axion  $Q_a \rightarrow$  Higher sensitivity!

$E^*$ (V/m)	Sensitivity (e·cm)		Ring
	a	b	
$4.23 \times 10^7$	$1.9 \times 10^{-31}$	$1.9 \times 10^{-31}$	E/B rin ( $r = 10$ )
$4.24 \times 10^7$	$6.0 \times 10^{-31}$	$1.9 \times 10^{-31}$	
$4.27 \times 10^7$	$1.9 \times 10^{-30}$	$1.9 \times 10^{-31}$	
$4.60 \times 10^7$	$5.5 \times 10^{-30}$	$1.8 \times 10^{-31}$	
$5.00 \times 10^7$	$7.2 \times 10^{-30}$	$1.6 \times 10^{-31}$	
$5.86 \times 10^7$	$8.7 \times 10^{-30}$	$1.4 \times 10^{-31}$	
$6.83 \times 10^7$	$9.1 \times 10^{-30}$	$1.2 \times 10^{-31}$	
$7.93 \times 10^7$	$9.1 \times 10^{-30}$	$1.0 \times 10^{-31}$	
$9.21 \times 10^7$	$8.8 \times 10^{-30}$	$8.8 \times 10^{-31}$	
$2.01 \times 10^8$	$4.0 \times 10^{-30}$	$4.0 \times 10^{-31}$	
$2.28 \times 10^8$	$3.5 \times 10^{-30}$	$3.5 \times 10^{-31}$	B ring ( $r = 10$ )
$2.75 \times 10^9$	$9.3 \times 10^{-31}$	$9.3 \times 10^{-31}$	

$\epsilon = 0.02$ ,

with  $A=0.36$ . SCT =  $10^4$  s.

# Easy systematic error control

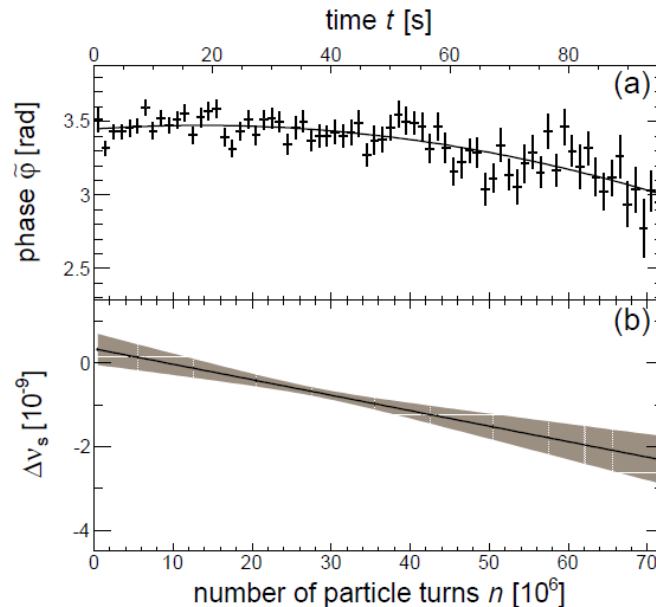


FIG. 3. (a): Phase  $\bar{\varphi}$  as a function of turn number  $n$  for all 72 turn intervals of a single measurement cycle for  $\nu_s^{\text{fix}} = -0.160975407$ , together with a parabolic fit. (b): Deviation  $\Delta\nu_s$  of the spin tune from  $\nu_s^{\text{fix}}$  as a function of turn number in the cycle. At  $t \approx 38$  s, the interpolated spin tune amounts to  $\nu_s = (-16097540771.7 \pm 9.7) \times 10^{-11}$ . The error band shows the statistical error obtained from the parabolic fit, shown in panel (a).

D. Eversmann et al. *New method for a continuous determination of the spin tune in storage rings and implications for precision experiments*. Phys. Rev. Lett., 115, 094801, Aug 2015.

➤ Experiment at COSY (Juelich, Germany)

Spin tune  $\nu_s$

$$\nu_s = \gamma G = \gamma(g - 2)/2$$

➤ Precise g-2 frequency tune is possible!

# Summary and conclusion

- ❖ We propose using the storage ring method to measure the axion induced oscillating EDM at the resonance conditions between the axion frequency and g-2 precession frequency.
- ❖ Calculated experimental conditions (E, B-field etc.) for deuteron and proton.
- ❖ The resulting sensitivity reaches up to  $\sim 10^{-32}$  e.cm for an oscillating EDM which have never been implemented so far. This assumes  $0.3 \text{ GeV/cm}^3$  local axion dark matter density. The EDM amplitude depends on the square root of the local axion density. The local dark matter density may be enhanced by several orders of magnitude due to focusing effects due to planet motion.
- ❖ A wide range of frequency (0.1 kHz~100 MHz) of axion dark matter can be tested using both deuterons and protons in the same storage ring.



# IBS conference on Dark World

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Thank you!