

Fermilab



Istituto Nazionale di Fisica Nucleare
Sezione di Padova



Istituto Nazionale di Fisica Nucleare
Laboratori Nazionali di Legnaro

Status and Prospects of QUaere AXion

at Laboratori Nazionali di Legnaro

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Status and Prospects of QUAX-PD/LNL

- Introduction: Axion Search at QUAX-PD/LNL
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- Summary

Introduction

Strong CP Problem and Composition of Dark Matter

1. The strong CP problem

$$\mathcal{L}_\theta = \bar{\theta} \frac{g_s^2}{32\pi^2} G_{\mu\nu}^a \tilde{G}^{a\mu\nu}$$

Weinberg (1975); 't Hooft (1976)

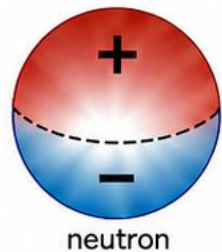
- This term violates CP in QCD.
- It induces a neutron electric dipole moment.
- Experiments imply $|\bar{\theta}| \lesssim 10^{-10}$.

$$d_n \approx (1 - 3) \times 10^{-16} \bar{\theta} e \cdot \text{cm}$$

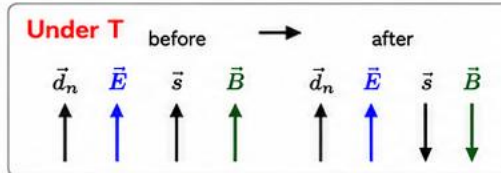
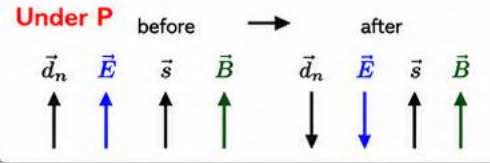
representative hadronic estimate

Crewther et al. (1979); Pospelov & Ritz (2005)

Why does a permanent neutron EDM violate P and T?



neutron



$$\Delta E = -\vec{d}_n \cdot \vec{E} - \vec{\mu}_n \cdot \vec{B}$$

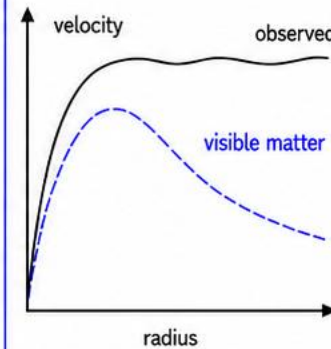
Baker et al. (2006); Abel et al. (2020)

Why is $\bar{\theta}$ so small? This is the strong CP problem.

2. What is dark matter made of?

Dark matter certainly exists observationally, but its microscopic particle nature is still unknown.

Galaxy rotation curves



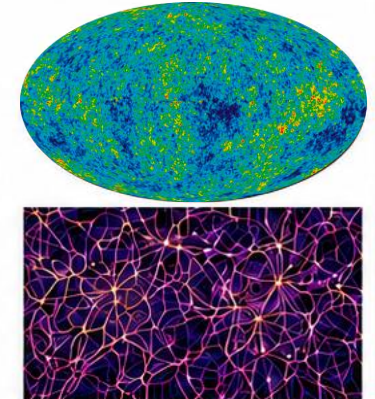
Rubin, Thonnard & Ford (1980)

Gravitational lensing



Clowe et al. (2006)

CMB / large-scale structure

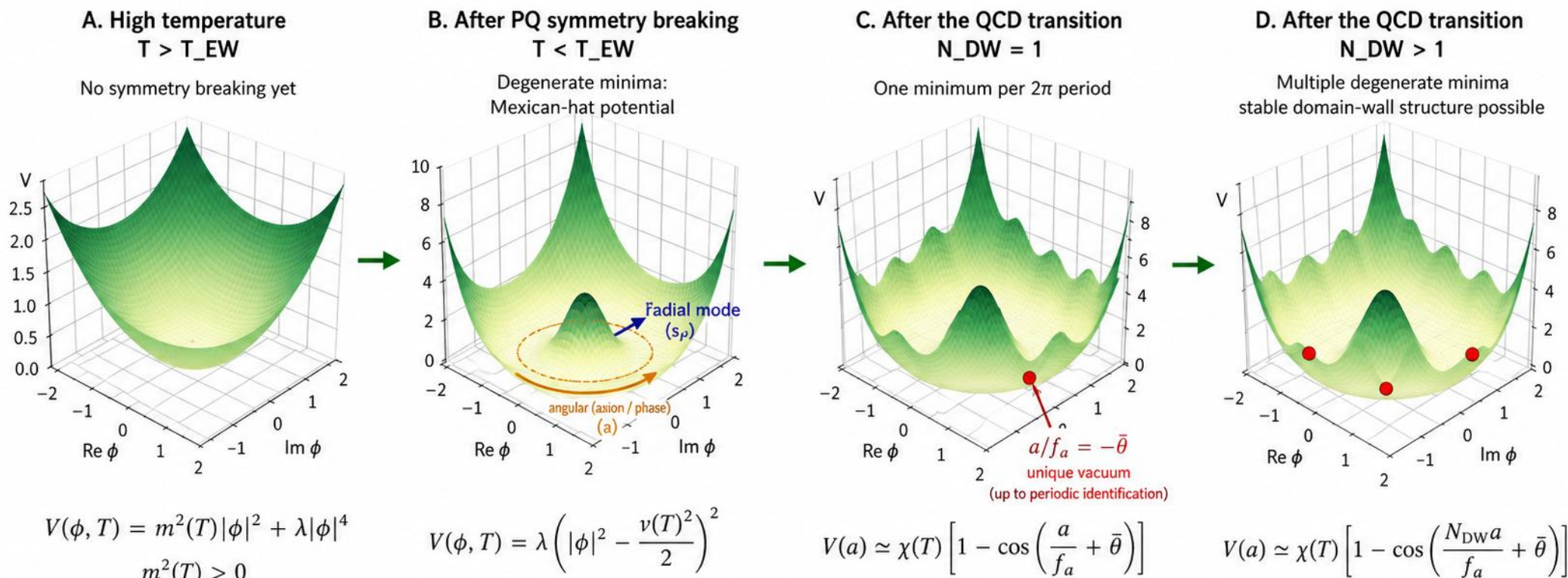


Planck Collaboration (2020)

We know dark matter is there, but we still do not know which particle it is made of.

Introduction

Axion can solve strong CP problem



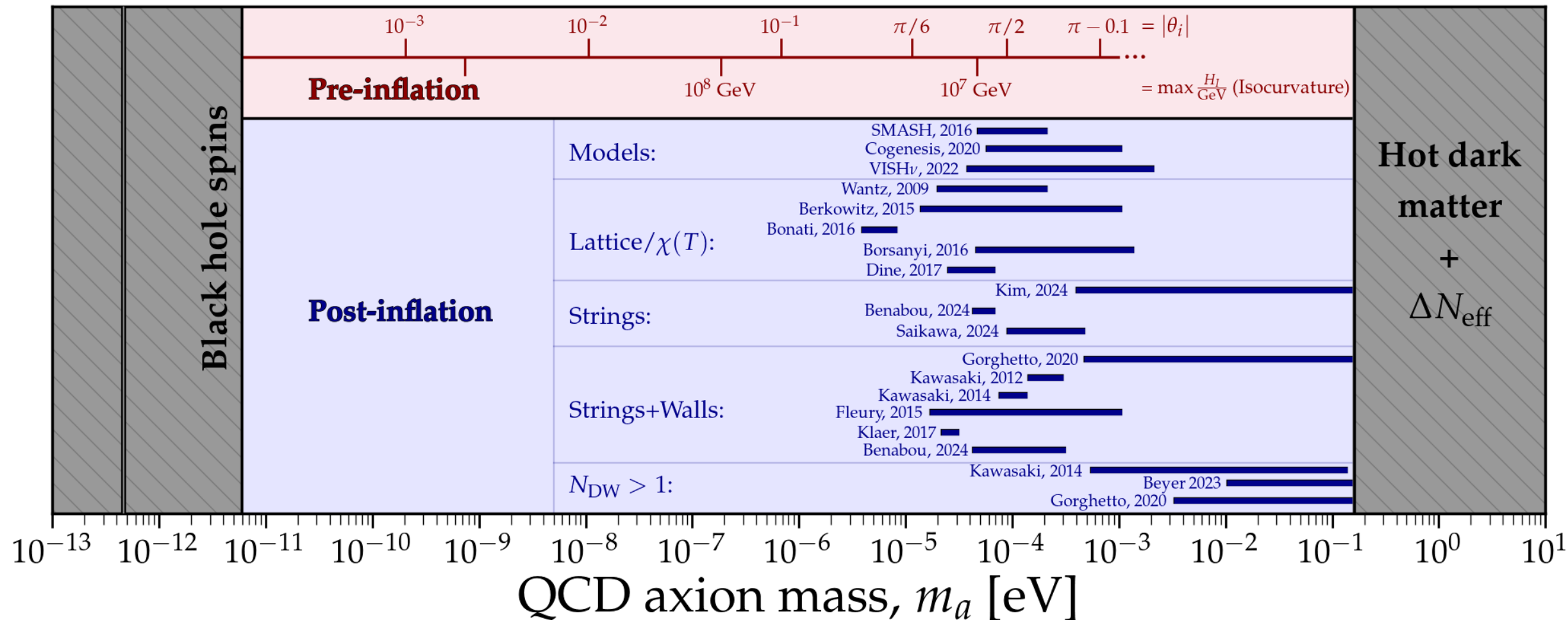
QCD generates the axion potential. For $N_{DW} = 1$ the vacuum is unique and the axion relaxes the effective $\bar{\theta}$ angle, solving the strong CP problem. For $N_{DW} > 1$ several degenerate vacua appear.

Refs. R.D. Peccei & H.R. Quinn, Phys. Rev. Lett. 38, 1440 (1977); S. Weinberg, Phys. Rev. Lett. 40, 223 (1978); F. Wilczek, Phys. Rev. Lett. 40, 279 (1978); L. Di Luzio et al., Phys. Rept. 870, 1 (2020).

Introduction

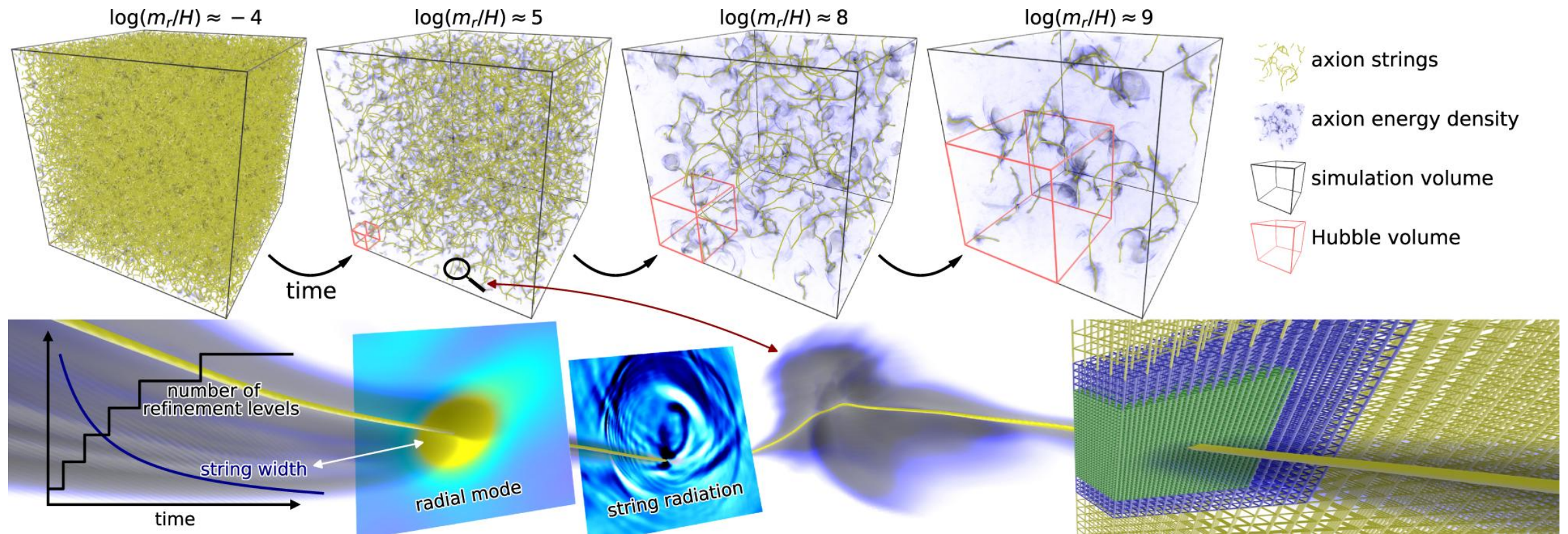
Axion is a dark matter candidate

<https://cajohare.github.io/AxionLimits/>



Introduction

Axion is a dark matter candidate

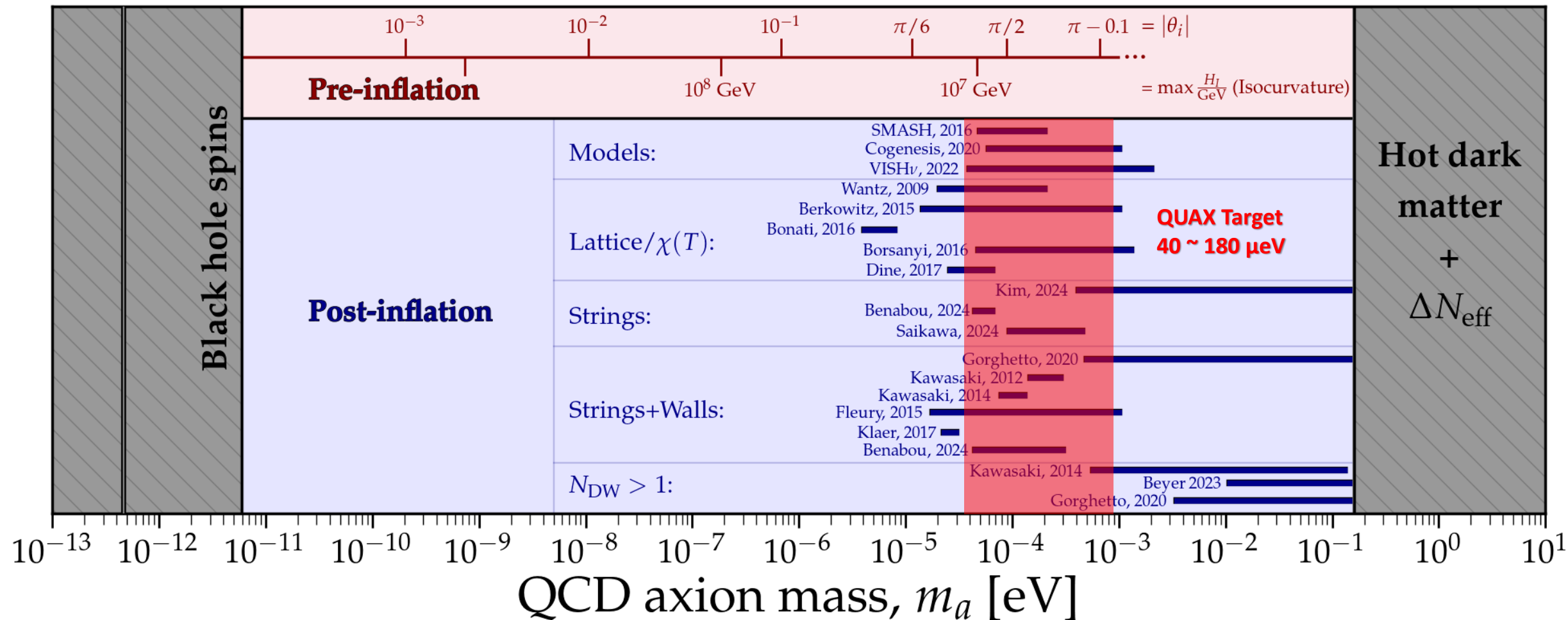


Buschmann *et al.* (2022)

Introduction

Axion is a dark matter candidate

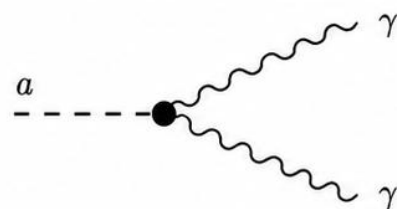
<https://cajohare.github.io/AxionLimits/>



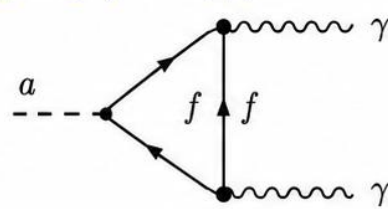
Introduction

Axion in the low energy regime

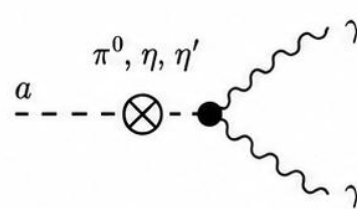
A. Axion – Photon Coupling ($a \rightarrow \gamma\gamma$)



effective vertex



charged-fermion triangle



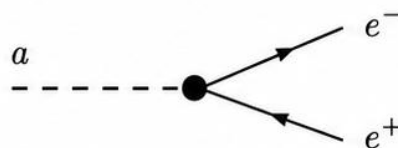
meson mixing (QCD axion)

$$g_{a\gamma\gamma} = \frac{\alpha}{2\pi f_a} \left(\frac{E}{N} - C_{\text{QCD}} \right) = \frac{\alpha}{\pi f_a} g_\gamma$$

Model	E/N	g_γ ($C_{\text{QCD}} = 1.92(4)$)	g_γ ($C_{\text{QCD}} \approx 1.77(8)$)
KSVZ (minimal)	0	−0.96	−0.885
DFSZ (canonical)	8/3	+0.373	+0.448

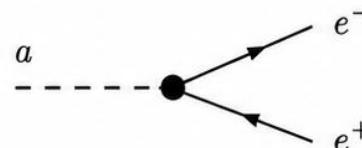
lattice QCD update: $C_{\text{QCD}} \approx 1.77(8)$ (vs standard 1.92(4))

B. Axion – Electron Coupling ($a \rightarrow e^+e^-$)



derivative coupling

$$\frac{C_e}{2f_a} (\partial_\mu a) \bar{e} \gamma^\mu \gamma^5 e$$



pseudoscalar form

$$-i g_{ae} a \bar{e} \gamma^5 e$$

$$g_{ae} = C_e \frac{m_e}{f_a}$$

KSVZ

$C_e = 0$ (tree level)

RG-induced: $C_e \sim 0.8 \times 10^{-3}$

$g_{ae} \sim 4.1 \times 10^{-19}$ ($f_a = 10^{12}$ GeV)

DFSZ

DFSZ-I: $C_e = +\frac{1}{3} \sin^2 \beta$

DFSZ-II: $C_e = -\frac{1}{3} \cos^2 \beta$

β

$\tan \beta = v_u/v_d$

$\tan \beta = 1 \Rightarrow \sin^2 \beta = \cos^2 \beta = \frac{1}{2}$

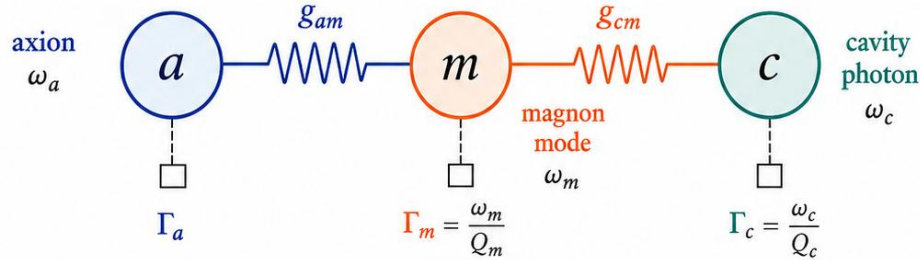
$C_e = \pm \frac{1}{6}$

$g_{ae} = \pm 8.52 \times 10^{-17}$ ($f_a = 10^{12}$ GeV)

Refs: di Cortona et al. (2016); Bauer et al. (2021); Choi et al. (2021); Jaeckel & Thormaehlen (2019).

Introduction

Ferromagnetic Haloscope - QUAX-g_{ae}



Coupled oscillator / multimode Hamiltonian

$$\hat{H} = \sum_i \hbar \left(\omega_{c,i} - i \frac{\gamma_{c,i}}{2} \right) a_i^\dagger a_i + \sum_j \hbar \left(\omega_{m,j} - i \frac{\gamma_{m,j}}{2} \right) m_j^\dagger m_j + \sum_{i,j} \hbar g_{ij} (a_i^\dagger m_j + a_i m_j^\dagger)$$

Example values used in the B_a and P_a estimates

m_a	=	42 μeV
g_{aee}	=	1.0×10^{-15} (benchmark, model dependent)
v_a	=	220 km/s $\approx 7.3 \times 10^{-4} c$
ρ_a	=	0.45 GeV/cm ³
n_a	=	$\rho_a / (m_a c^2) \approx 1.1 \times 10^{13} \text{ cm}^{-3}$
n_s	=	$1.0 \times 10^{22} \text{ cm}^{-3}$ (YIG spin density)
V_m	=	$4.8 \times 10^{-2} \text{ cm}^3$ (EG spheres, diameter 2.1 mm)
N_s	=	$n_s V_m \approx 4.8 \times 10^{20}$
ω_a	=	$6.4 \times 10^{10} \text{ s}^{-1}$ ($\approx 2\pi \times 10.2 \text{ GHz}$)
τ_s	=	$1.0 \times 10^{-6} \text{ s}$
B_a	$\approx$	$4 \times 10^{-23} \text{ T}$

Values are benchmark inputs for order-of-magnitude estimates.

Axion signal power

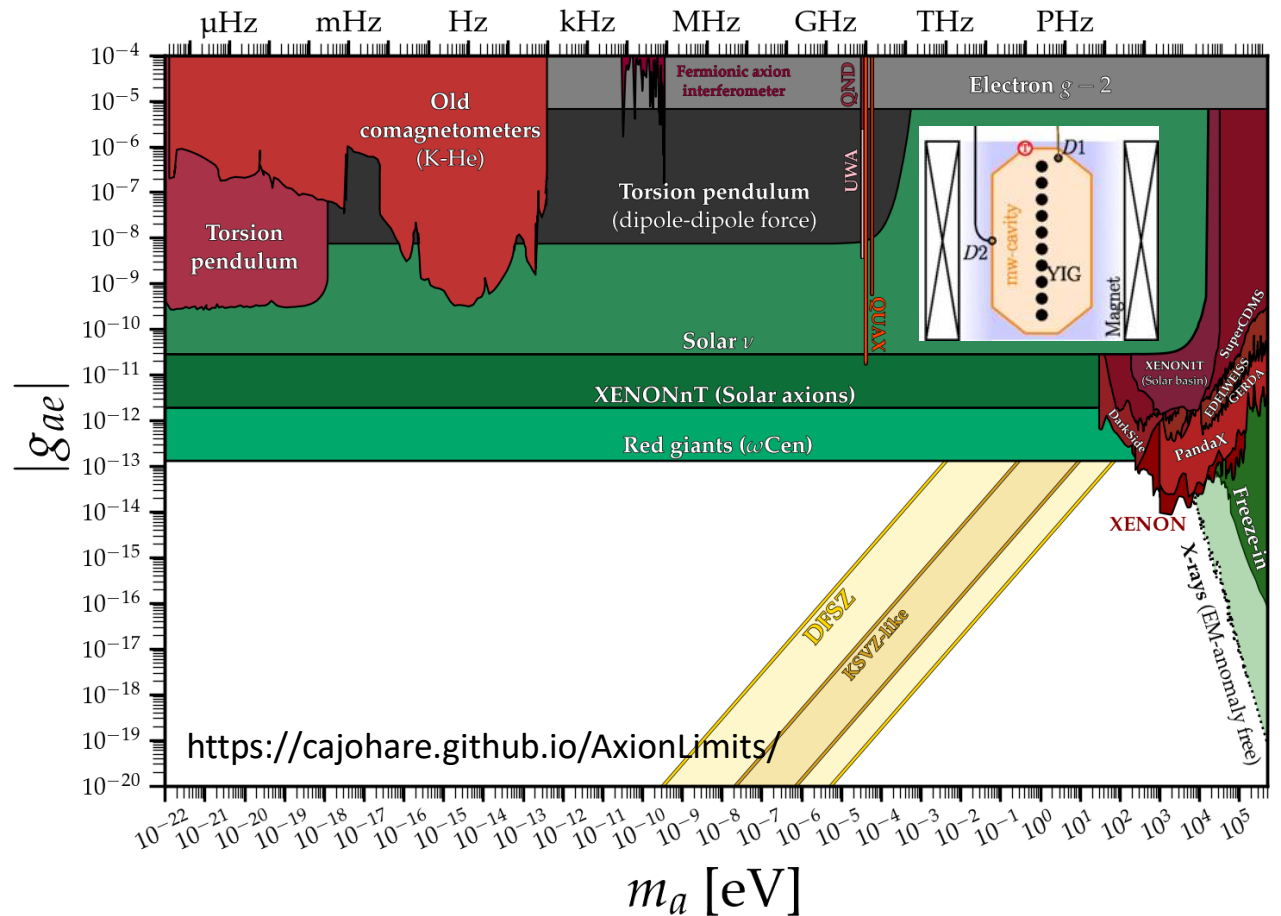
$$\mathcal{L}_{ae} \simeq -2\mu_B \boldsymbol{\sigma} \cdot \left(\frac{g_{aee}}{2e} \nabla a \right) \equiv -2\mu_B \boldsymbol{\sigma} \cdot \mathbf{B}_a$$

$$B_a = \frac{g_{aee}}{2e} \sqrt{\frac{n_a \hbar}{m_a c}} m_a v_a \simeq 4 \times 10^{-23} \left(\frac{m_a}{42 \mu\text{eV}} \right) \text{ T}$$

$$P_a = \gamma_e \mu_B N_s \omega_a B_a^2 \tau_s$$

Example estimate: $P_a \approx 1 \times 10^{-31} \text{ W}$

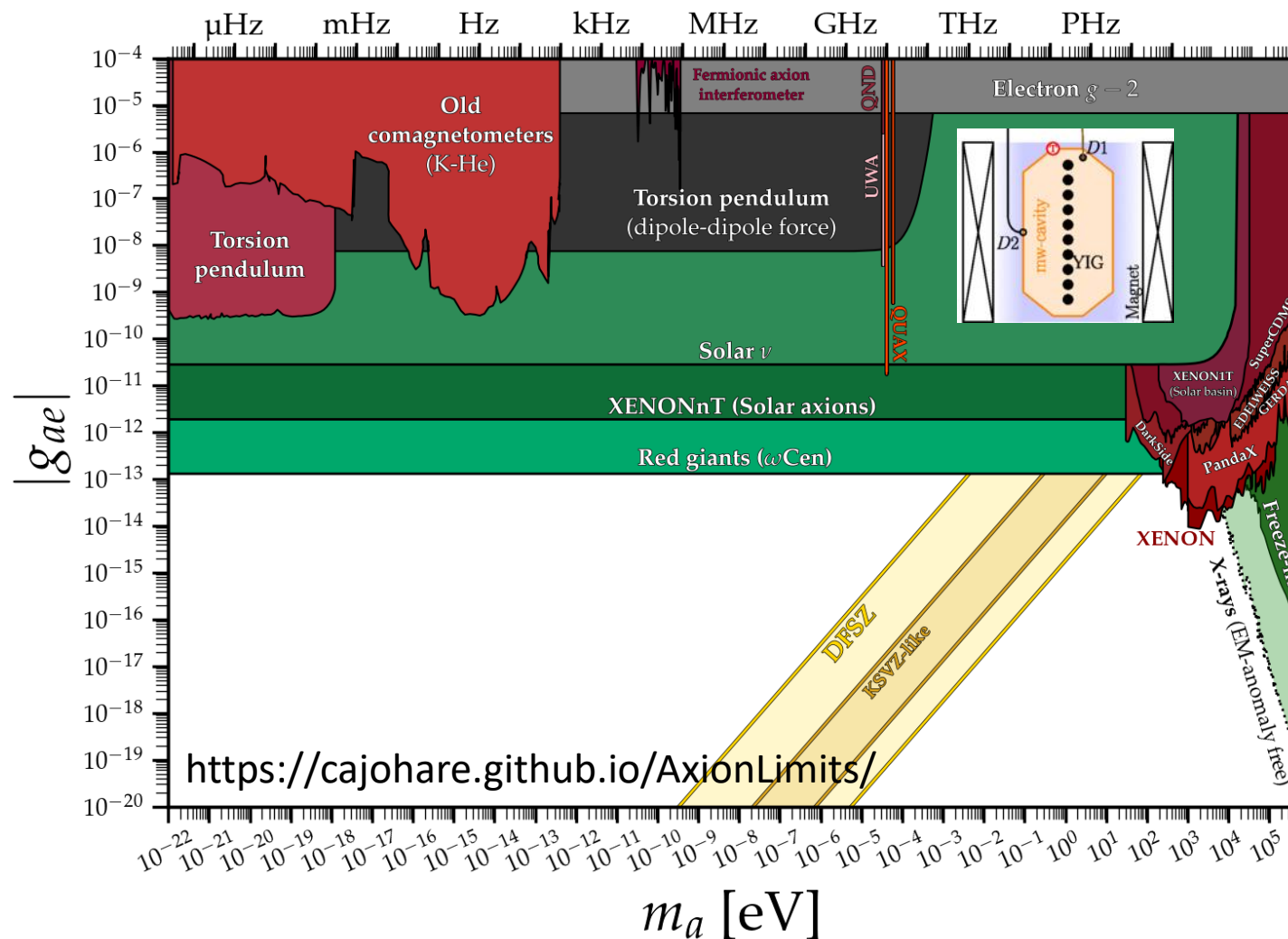
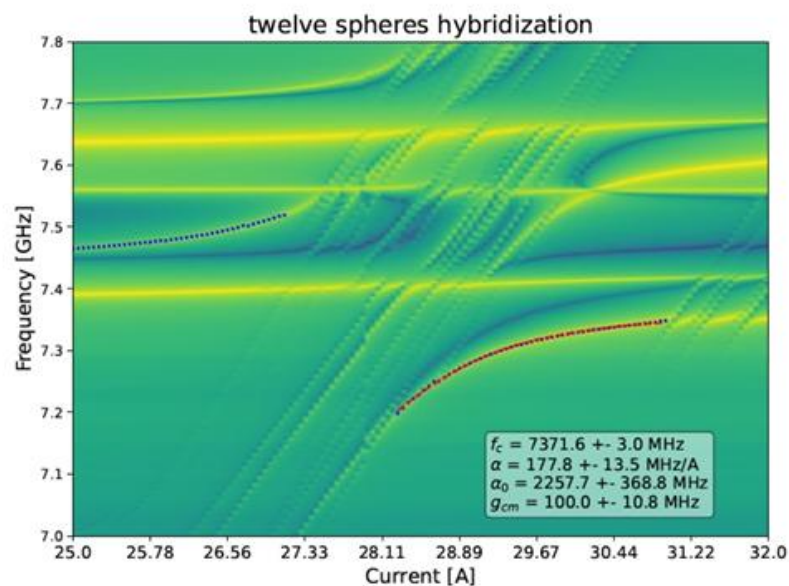
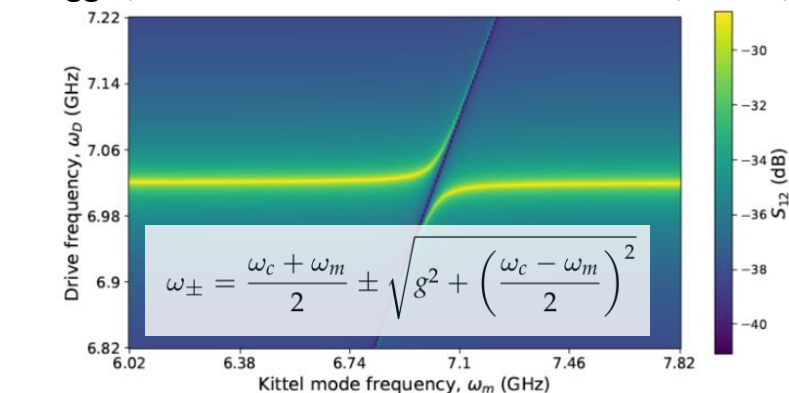
References: Crescini et al., *Phys. Dark Universe* 15, 135–141 (2017); *Phys. Rev. Lett.* 124, 171801 (2020); *Commun. Phys.* 3, 164 (2020); *Phys. Rev. B* 104, 064426 (2021).



Introduction

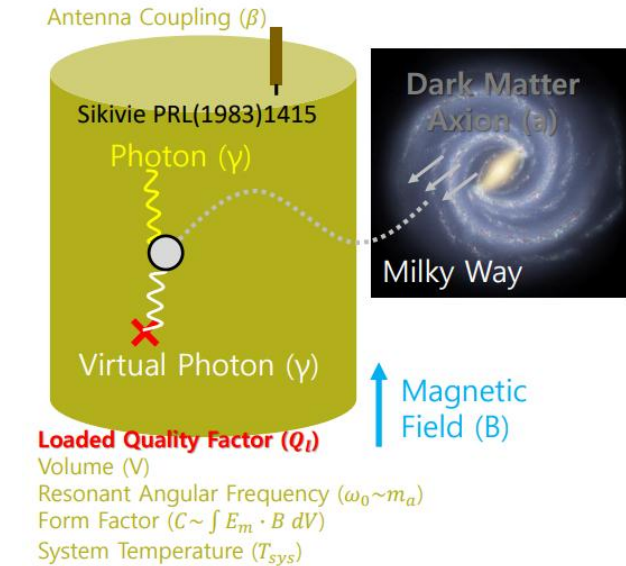
Ferromagnetic Haloscope - QUAX- g_{ae}

C. Braggio, Presented at the STDM Conference, Rome, Italy



Introduction

Axion Haloscope - QUAX- $g_{a\gamma}$



Signal Power

$$P_{sig} = \frac{\beta}{1+\beta} g_{a\gamma\gamma}^2 \frac{\rho_a}{m_a^2} B^2 V \omega_0 C \frac{Q_a Q_l}{Q_a + Q_l}$$

Coupling Constant

Dark Matter Axion Density

Axion Mass

Axion Quality Factor

Kim et al. JCAP03(2020)066

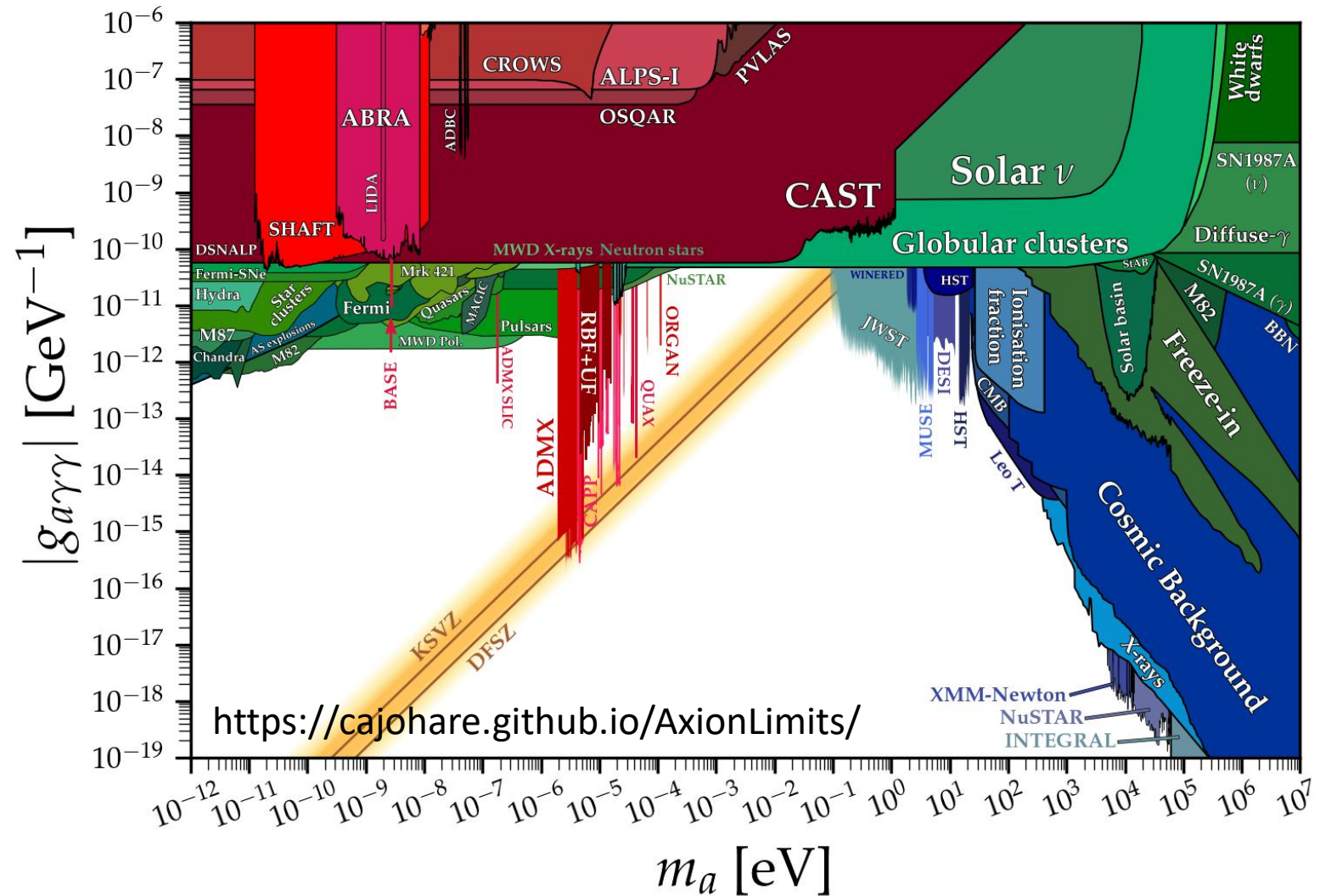
Scanning speed

$$\frac{df}{dt} \propto \frac{B^4 V^2 C^2}{k_B^2 T_{sys}^2} Q_l Q_a$$

$Q_l \gg Q_a \sim 10^6$

System Noise Temperature

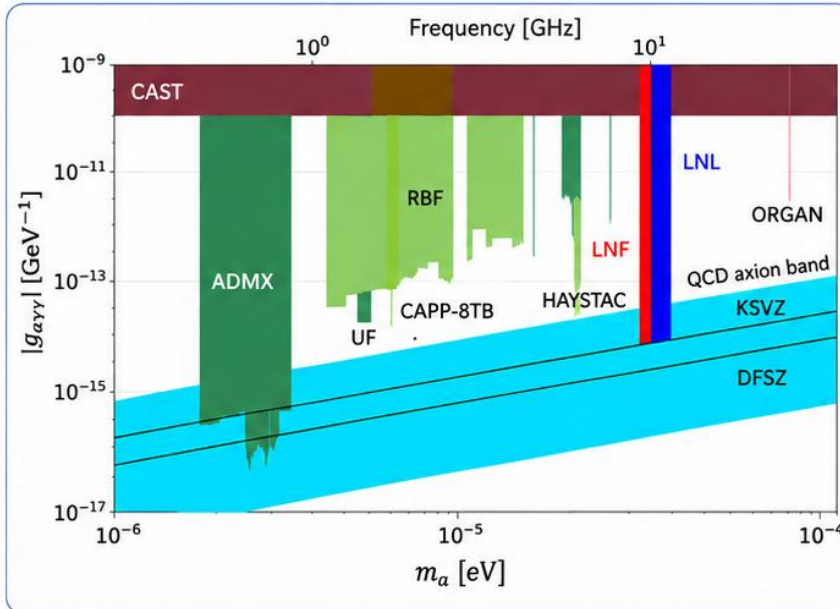
Kim, Patras Workshop, 2023



Introduction

QUaerere AXion (QUAX)

- In 2020 **INFN** has financed the **QUAX** experiment to run an observatory for searching axion via the **axion-photon coupling in the unexplored region around 10 GHz**
- QUAX** is based on **two haloscopes**: one in Legnaro/Padova (LNL) and the other in Frascati (LNF)



G. Ruoso, Presented at the STDM Conference, Rome, Italy

	LNF (Frascati)	LNL (Legnaro/Padova)
Magnetic field	9 T	14 T
Magnet length	40 cm	50 cm
Magnet inner diameter	9 cm	12 cm
Frequency range	8.5 – 10 GHz	9.5 – 11 GHz
Cavity type	Hybrid SC	Dielectric
Scanning type	Inserted rod	Mobile cylinder
Number of cavities	7	1
Cavity length	0.3 m	0.4 m
Cavity diameter	25.5 mm	58 mm
Cavity mode	TM010	pseudo-TM030
Single volume	$1.5 \cdot 10^{-4} \text{ m}^3$	$1.5 \cdot 10^{-4} \text{ m}^3$
Total volume	$7 \times 0.15 \text{ liters}$	0.15 liters
Q_0	300 000	1 000 000
Single scan bandwidth	630 kHz	30 kHz
Axion power	$7 \times 1.2 \cdot 10^{-23} \text{ W}$	$0.99 \cdot 10^{-22} \text{ W}$
Preamplifier	TWJPA/INRIM	DJJAA/Grenoble
Operating temperature	30 mK	30 mK

- The **LNL/PD haloscope** will be based on dielectric loaded cavities, traveling wave parametric amplifiers and **14 T magnet**

Introduction

QUAX-PD/LNL Infrastructure



QUAX-BLU

10 mK

9 T Magnet

Single Microwave Photon Counter

Axion/Ferromagnetic Haloscope

QUAX-LNL

Sub-100 mK

9 T Magnet

Quantum-Limited Noise Amplifier

Axion Haloscope

QUAX-RED

Sub-100 mK

14 T Magnet

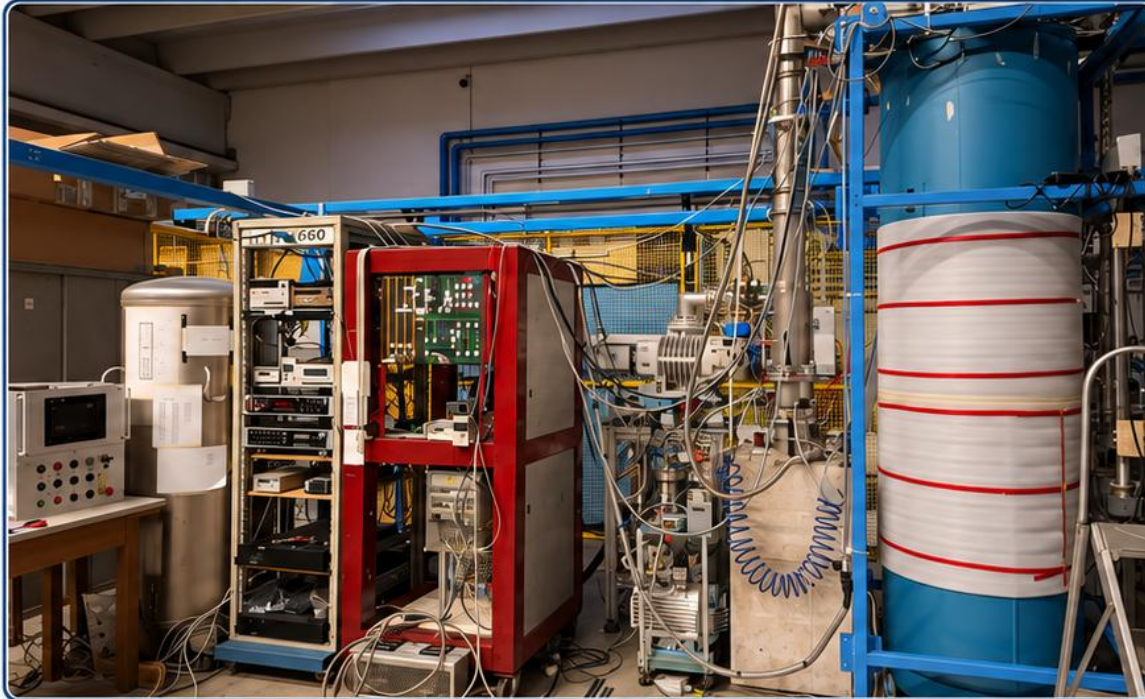
Quantum-Limited Noise Amplifier

Axion Haloscope, Dark Photon

QUAX-LNL Haloscope

Dilution Refrigerator

QUAX main criostat with gas handling system and safety control



- Main dewar: 2800 mm height, 500 mm inner diameter
- Two weeks operation: ~3000 liters of liquid helium
- Refurbished **wet dilution unit** from AURIGA experiment
- Base temperature @ MC 60 mK
- Cavity and electronics kept @ 120 mK



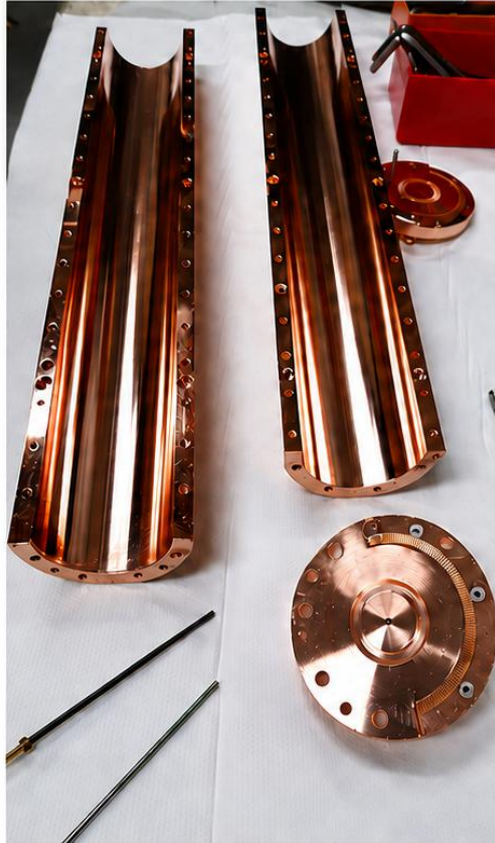
Gold-plated
heat shield

G. Ruoso, Presented at the STDM Conference, Rome, Italy

QUAX-LNL Haloscope

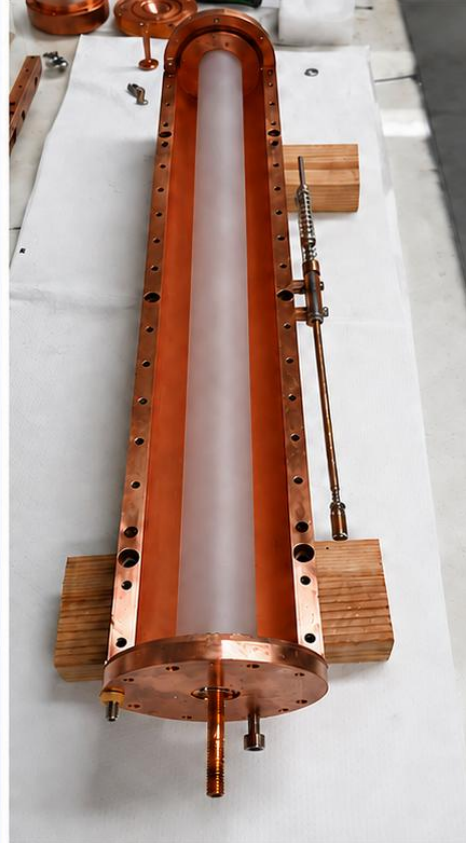
Dielectrically Loaded Tunable Cavity

Cavity parts



Endcap with RF sliding contacts and sapphire holding grooves

Clamshell mechanism for tuning



Open cavity with sapphire cylinder

Concept and specifications

- Right cylindrical hybrid cavity
- TM₀₃₀ mode for large volume
- Copper shell
- Hollow sapphire cylinder inside

- **LAMED TC:** tunable cavity
“LAMED” from the Phoenician letter meaning “goad”

- **Base frequency:** 10.2 GHz
- **Inner length:** 413.4 mm
- **Inner diameter:** 60.5 mm
- **Volume:** 1.17 liters

RF antenna



QUAX-LNL Haloscope

Dielectrically Loaded Tunable Cavity

LAMED-TC

name derived from the
Phoenician letter meaning 'goad'

L



End cap details

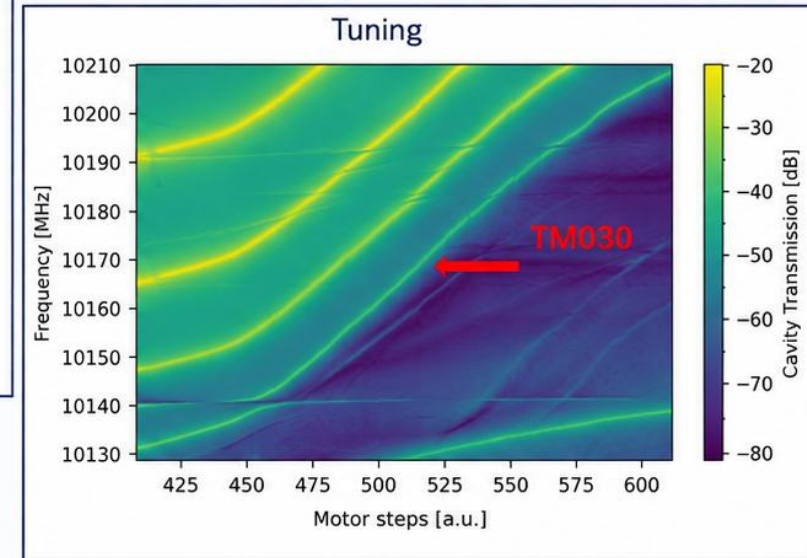


Pantograph



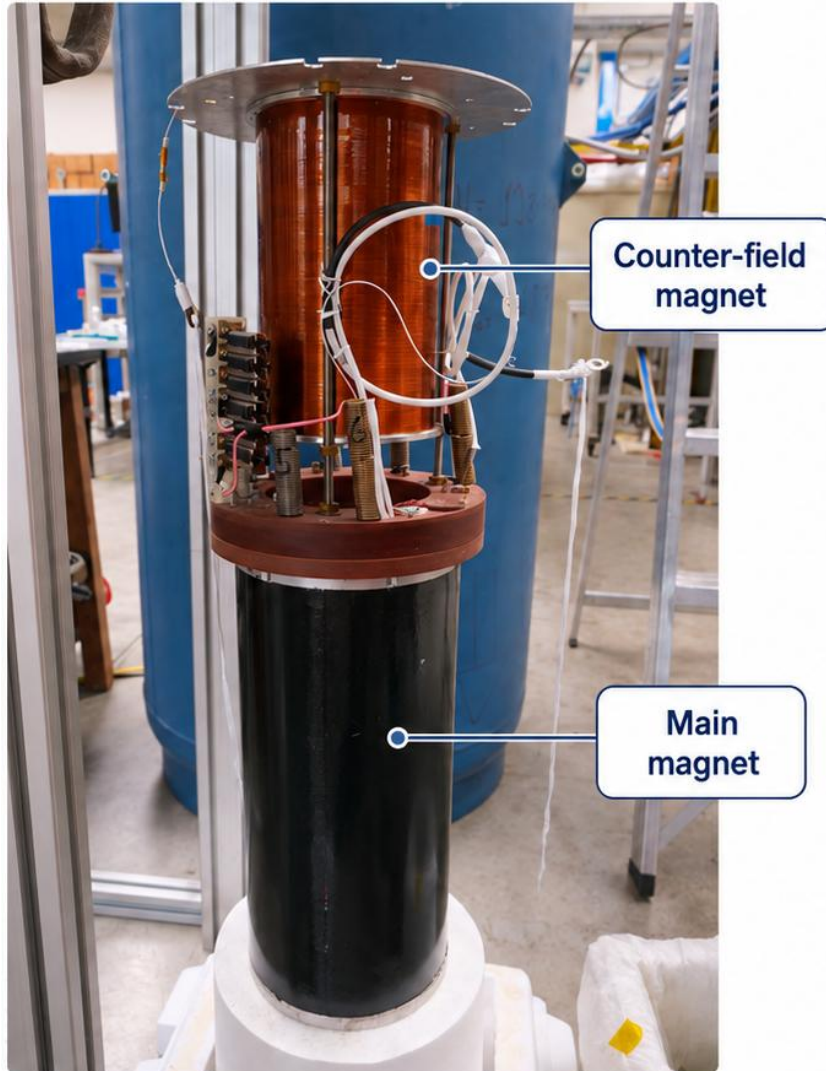
- Inner length = **413.4 mm**
- Inner diameter = **60.5 mm**
- Volume = **1.17 liters**
- $C_{030} = 0.43$

- Frame OFHC copper
- L = 420 mm sapphire cylinder (from Armenia)
- Movable side wall – clamshell mechanism
- Single-body endcaps
- Endcaps with 5.3 mm deep grooves for sapphire holding
- Copper RF sliding contacts
- Motor-actuated pantograph-like aperture mechanism
- Teflon 1.1 mm diameter to keep sapphire centered
- **Science mode TM030**
- Tuning 85 MHz
- $Q_0 \sim 80,000$



QUAX-LNL Haloscope

Magnet



Counter-field magnet

- NbTi superconducting coil
- Inner diameter **150 mm**, height **250 mm**
- Same driving current as the main magnet
- Reduces stray field below **40 mT** on sensitive electronics (also passively shielded)



Main magnet

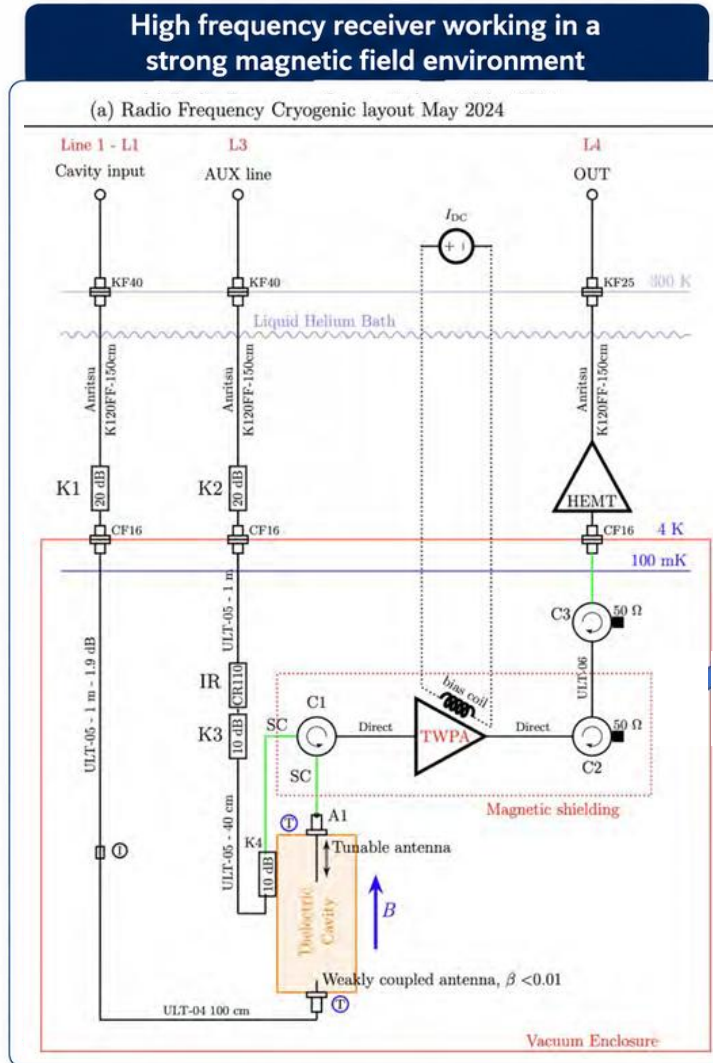
- NbTi superconducting coil
- Length **450 mm**, inner diameter **150 mm**
- Peak field **8.0 T @ 92 A** current
- Square field integral **50.2 T² m**

G. Ruoso, Presented at the STDM Conference, Rome, Italy

IBS Dark Matter Axion Group Seminar

QUAX-LNL Haloscope

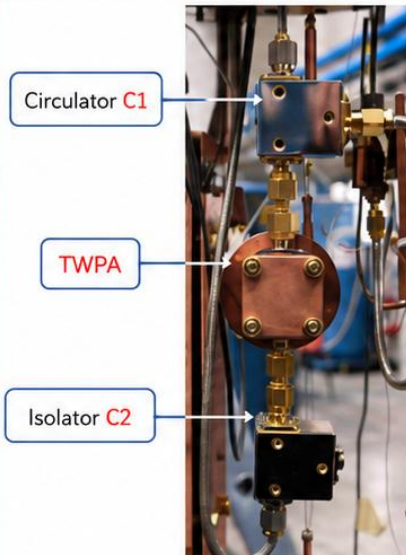
RF Receiver



The 10 GHz low noise receiver is based on a Travelling Wave Parametric Amplifier (TWPA - from N. Roch, Grenoble)

A. Ranadive, et al., *Nat. Commun.* 13, 1737 (2022)

RF chain in magnetic shielding box



Magnetic shielding (grey box)

Detail of DU insert



TWPA is made by superconducting Aluminum Josephson Junction arrays, very sensitive to ambient magnetic field

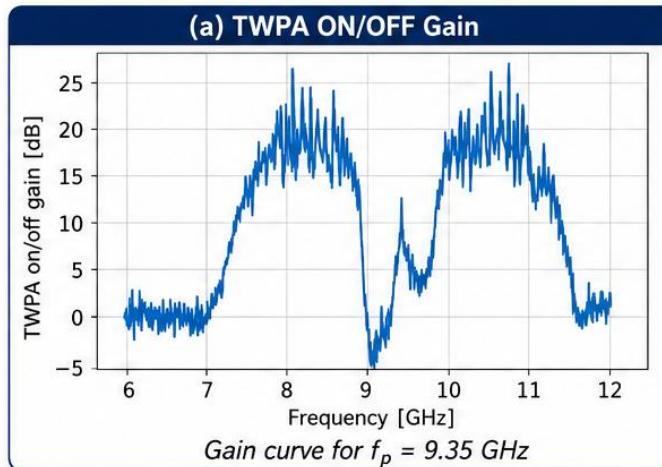
Two-layers passive shielding: **CRYOPERM (inside) + lead**

G. Ruoso, Presented at the STDM Conference, Rome, Italy

QUAX-LNL Haloscope

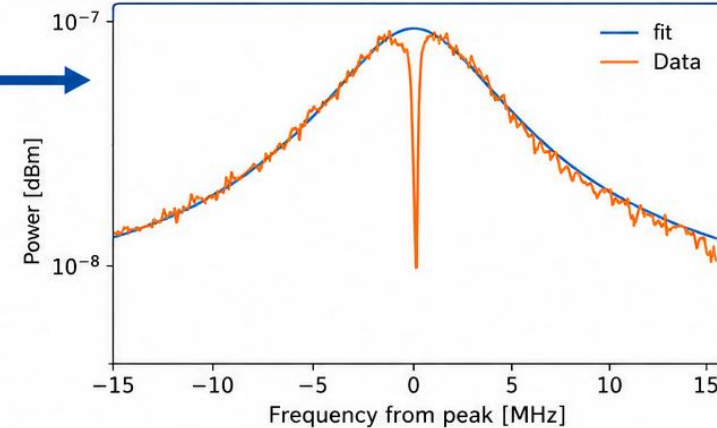
Travelling Wave Parametric Amplifier (TWPA) Operation

- **four wave mixing TWPA** with a phase matching process based on the sign reversal of the third order (Kerr) nonlinearity of a series of Josephson Junctions arrays (A. Ranadive, et al., *Nat. Commun.* **13**, 1737 (2022))
- TWPA amplification controlled by **intensity I_p** and **frequency f_p** of the pump tone and **current I_b** of a biasing magnetic field

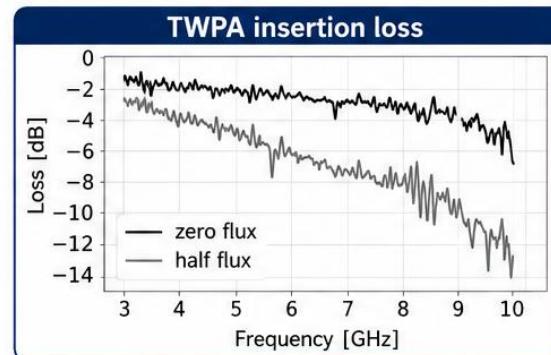


Single ripple

In order to maximize SNR one has to tune the gain profile of the TWPA amplifier



High frequency operation has to cope with large IL that reduces effective gain



Measured Noise Temperature

$$T_{\text{sys}} = 1.1 \text{ K @ } 10.2 \text{ GHz}$$

$$T_{\text{sys}} = 0.8 \text{ K @ } 7.9 \text{ GHz}$$

-> ~ 2.3 Photons

QUAX-LNL Haloscope

Run Control

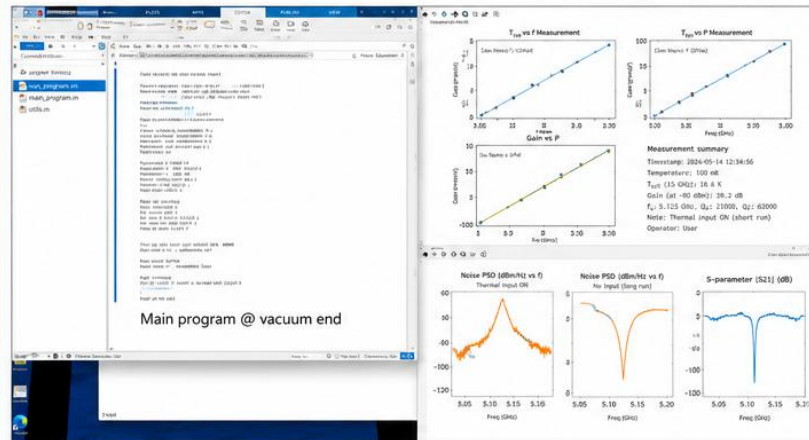
Run steps:

1. Cavity tuning (5 m)
2. TWPA amplifier set ()
3. Measurement of T_{sys} and Gain (2 m)
4. Measurement of f_c , Q_L , Q_0 , gain profile (2 m)
5. Short data acquisition with thermal input (2 m)
6. Repeat steps 3-4 (4 m)
7. Long data acquisition with no input (65 m)
8. Repeat steps 3-4 (4 m)
9. Start over

Time for a single point:

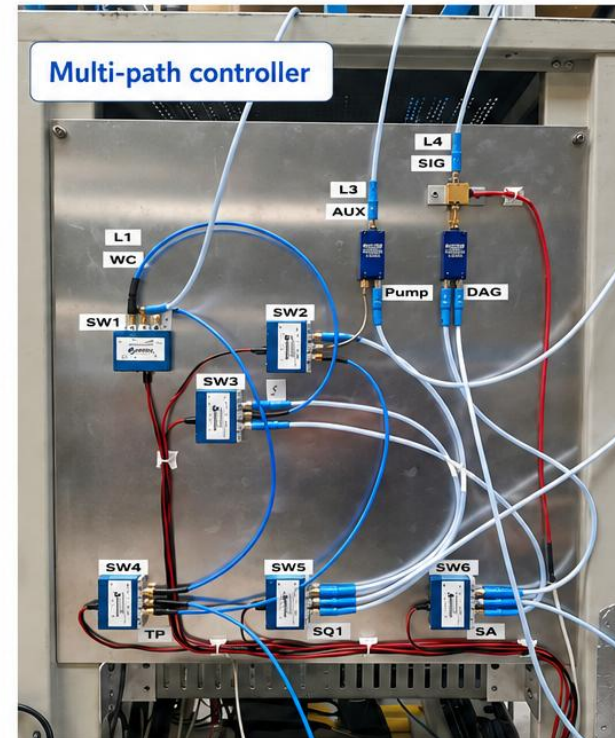
80 – 100 minutes

Slow control



All measured parameters, data and spectra saved to logfiles

Noise Temperature Measurement @ every step



Original idea published in:

A haloscope amplification chain based on a travelling wave parametric amplifier –
Rev. Sci. Instrum. 2022

QUAX-LNL Haloscope

QUAX Run 2024

Run summary

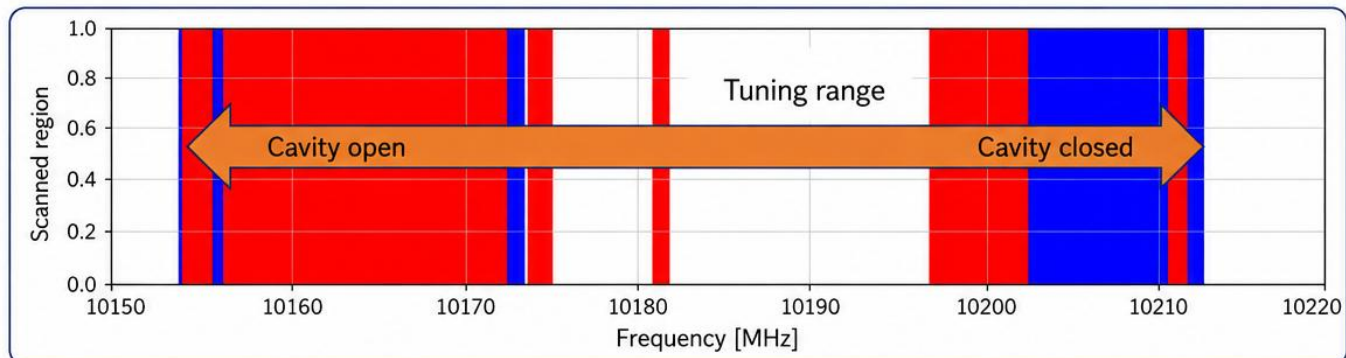
- Magnetic field ON
- Cavity frequency $f_c = 10.15\text{--}10.21\text{ GHz}$
- Noise temperature $T_{\text{sys}} = 1.1\text{--}1.5\text{ K}$
- Quality factor $Q_0 = 60000\text{--}80000$
- Antenna coupling $\beta = 1.4\text{--}1.8$
- Cavity Volume $V = 1.16\text{ liters}$
- Estimated efficiency $C_{030} = 0.4$
- Effective field $B^2 = 50.89\text{ T}^2$
- Axion mass $m_a = 44.9\text{ }\mu\text{eV}$
- Typical Integration time $t_m = 3800\text{ s}$



Expected axion power in a 10 kHz window $P_a = 6.0\text{e-}24\text{ W}$



Expected sensitivity (Dicke) in a 10 kHz window $\sigma_p = 2.5\text{e-}23\text{ W}$



Spring session

Autumn session

QUAX-LNL Haloscope

QUAX Run 2024

Spring session

Total run time 3 weeks

2 separate weeks for data taking

May 28th to May 30th - 48 h of field ON

June 11th to June 14th - 90 h of field ON

Autumn session

Total run time 2 weeks

Nov 7th to Nov 13th - 138 h of field ON

Break due to power failure

Nov 19th to Nov 21th - 42 h of field ON

Covered span:

~40 MHz with 225 h vacuum data taking

Maximum tuning: 58.45 MHz

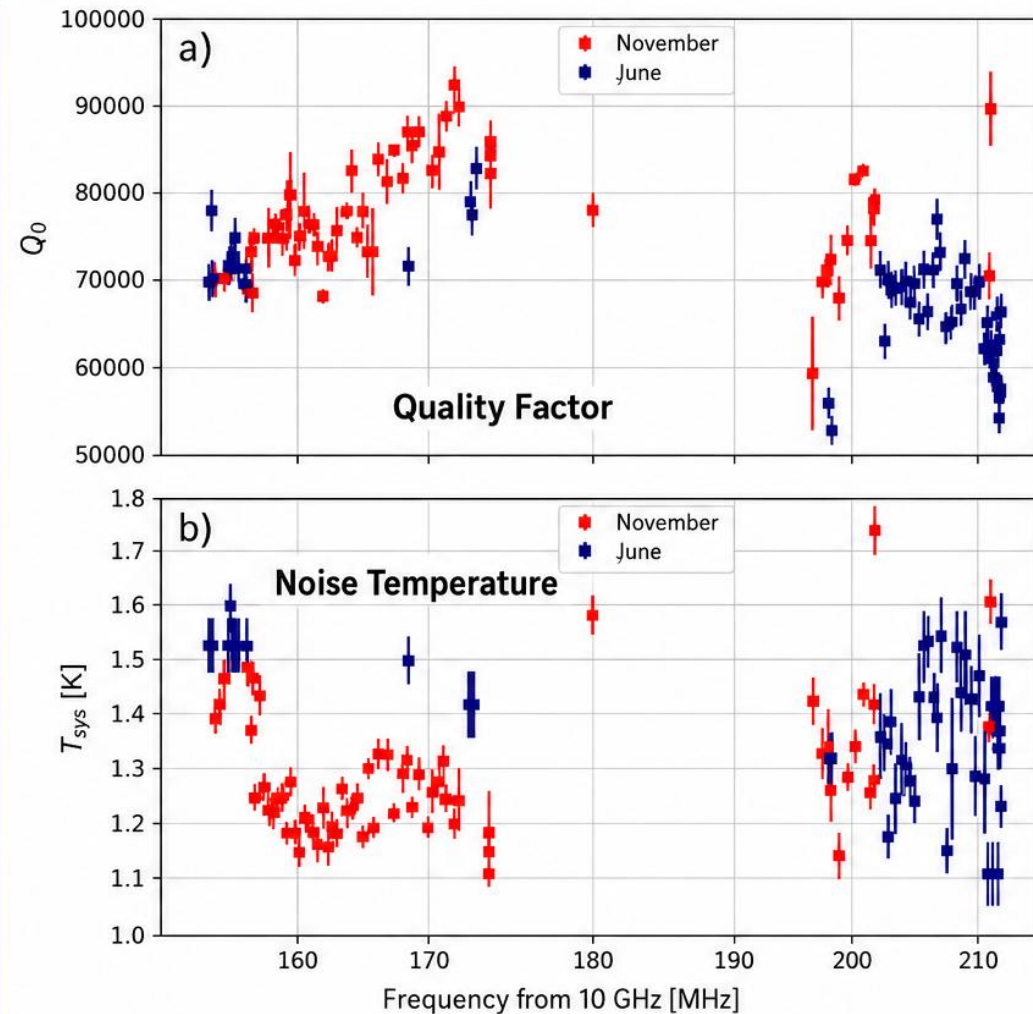
Ratio: 65% of available scan

Duty cycle 50%

Effective scan rate about

100 kHz/hour

2.5 MHz/day



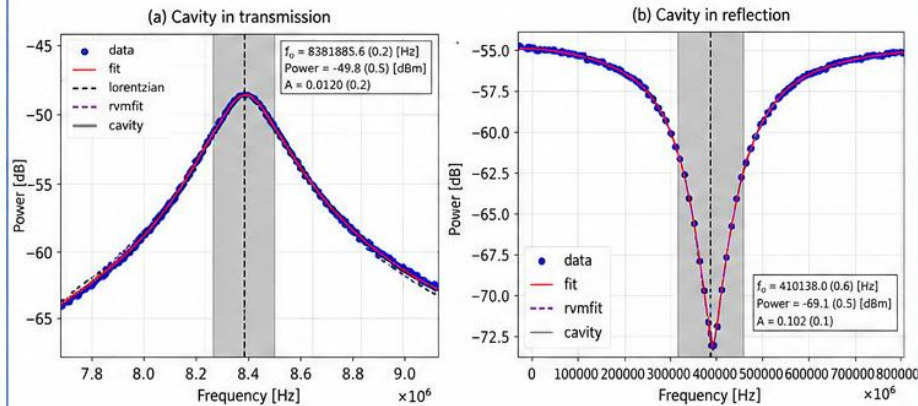
G. Ruoso, Presented at the STDM Conference, Rome, Italy

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QUAX-LNL Haloscope

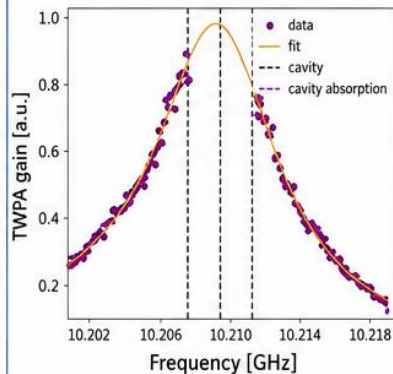
Data Analysis

Parameter evaluation (β , f_c , Q_0 , Q_L)

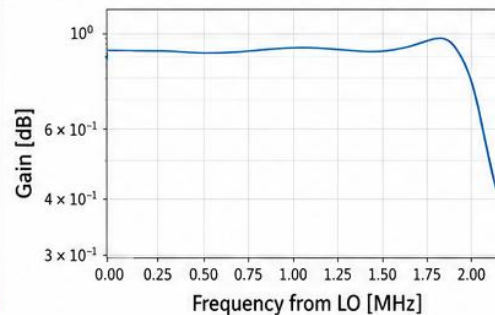


Remove freq-dependent

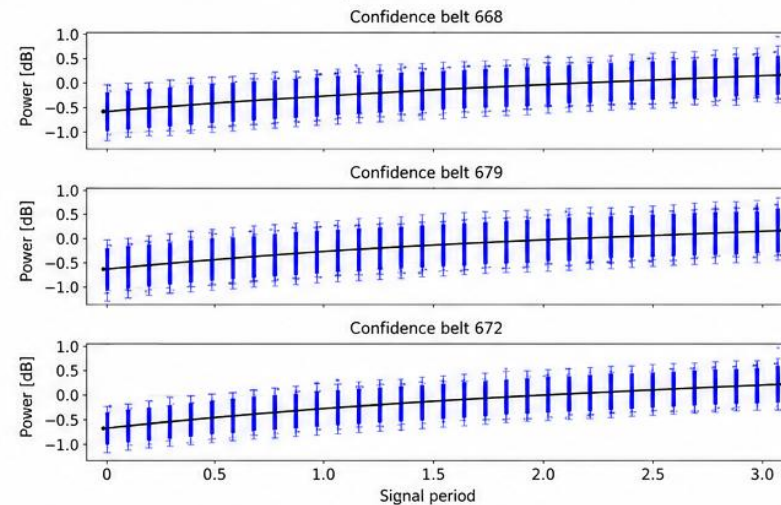
TWPA



Filters



- Average FFT of single frequency steps (bin = 268 Hz)
- Baseline removal with Savitzky Golay filtering
- Monte Carlo simulations to build confidence belts and look for excess power



Axion power

$$P_{\text{in}}^a(\nu, \nu_a) = g_{a\gamma\gamma}^2 \frac{\hbar^3 c^3 \rho_a}{m_a^2} \frac{2\pi\nu_a}{\mu_0} B^2 V C_{030} f_a(\nu, \nu_a)$$

Maxwell-Boltzmann distribution

$$f_a(\nu, \nu_a) = \frac{2}{\sqrt{\pi}} \nu \sqrt{\nu - \nu_a} \left(\frac{3}{1.7\nu_a \langle \beta_a^2 \rangle} \right)^{3/2} e^{-\frac{3(\nu - \nu_a)}{1.7\nu_a \langle \beta_a^2 \rangle}}$$

G. Ruoso, Presented at the STDM Conference, Rome, Italy

QUAX-LNL Haloscope

QUAX Run 2024 Result



No candidate signal have been detected with **threshold 4.5σ**

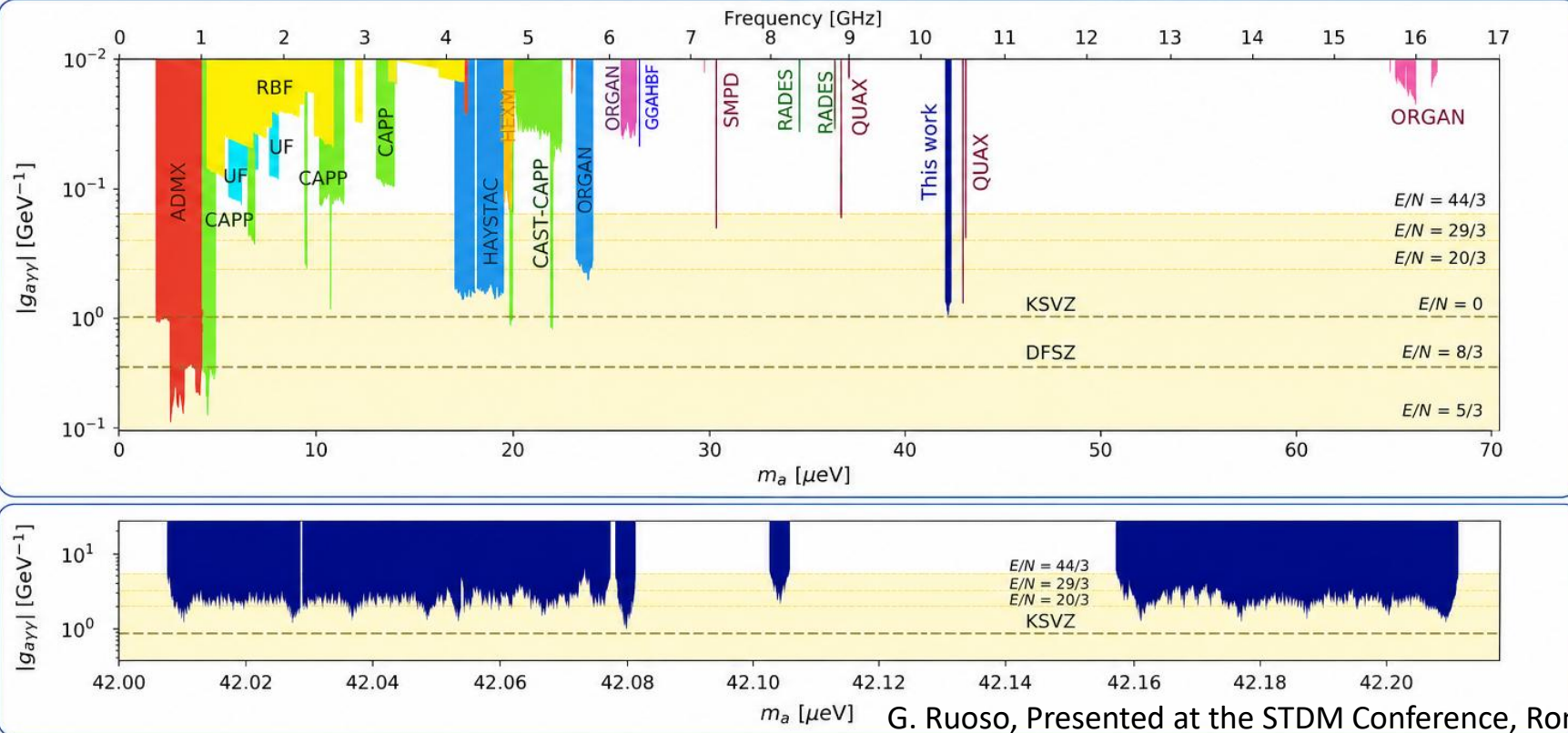
- **Axion-photon limits at 90% C.L.**

PHYSICAL REVIEW LETTERS 135, 211002 (2025)

Search for Postinflationary QCD Axions with a Quantum-Limited Tunable Microwave Receiver

G. Sardo Infirri^{1,2}, D. Alesini³, C. Braggio^{1,2}, G. Cappelli⁴, G. Carugno^{1,2}, D. D'Agostino³, A. D'Elia³, D. Di Gioacchino³, R. Di Vora⁶, M. Esposito^{1,2}, P. Falferi^{5,6}, U. Gambardella⁵, A. Garlondato², C. Gatti³, C. Ligi³, G. Lilli⁵, A. Lombardi⁶, G. Maccarrone⁵, D. Maiello^{1,2}, A. Ortolan⁵, A. Ranadive^{1,5}, A. Rettaroli³, N. Roch⁴, S. Tocci³ and G. Russo^{6,*}

(QUAX Collaboration)

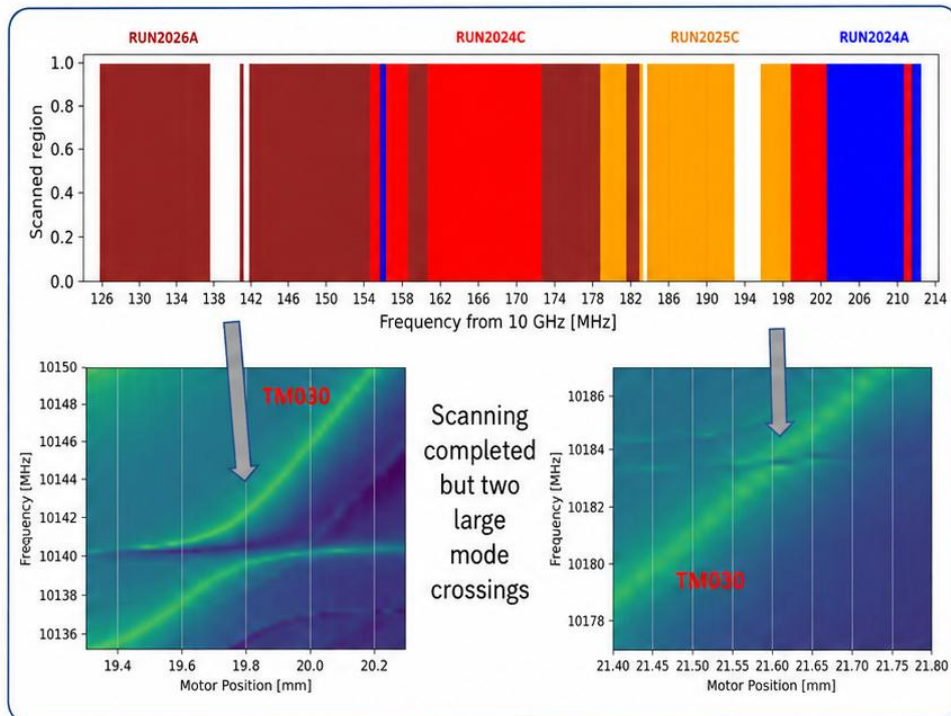


G. Russo, Presented at the STDM Conference, Rome, Italy

QUAX-LNL Haloscope

Measurement Sessions 2025-2026

- Major problems with DU unit, reduced machine availability
→ **Problems now solved**
- No changes on the RF hardware
- Improvement on tuning mechanism to exploit full cavity range
- Improvements on the RF and acquisition system to employ almost full **automation**



2025C session

Total run time 8 days

Nov 13th to Nov 17th - 92 h of field ON
Scanned about 15 MHz



2026A session

Total run time 11 days

April 28th to May 6th - 188 h of field ON
Scanned about 34 MHz



Covered span 2025 - 2026:

~49 MHz with 280 h magnetic field ON



Effective scan rate about

180 kHz/hour

4.2 MHz/day



Data analysis on the way !

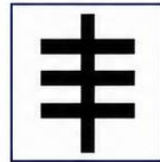
QUAX-LNL Haloscope

Future Plan

- Data analysis and check for re-scanning with current cavity
- **Install new cavity with larger tuning and increased Q factor**

SAMEK 3G

(Phoenician: Pillar)



Preliminary results:

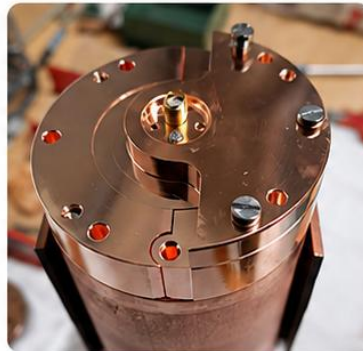
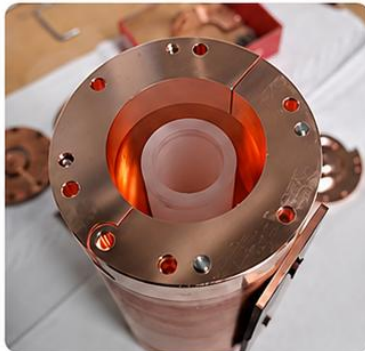
tuning range

170 MHz (10.04 – 10.21 GHz)

with

Q0 ~ 100 000 – 130 000

- New endcaps – two sliding parts
- No rf sliding contacts
- New aperture mechanism



G. Ruoso, Presented at the STDM Conference, Rome, Italy

IBS Dark Matter Axion Group Seminar

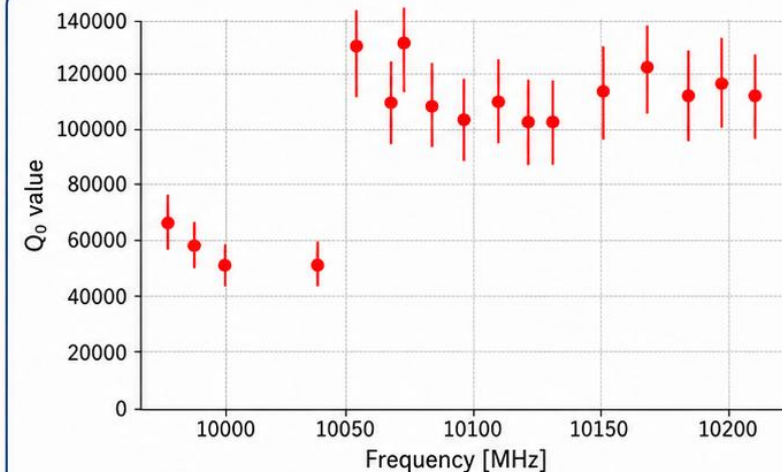
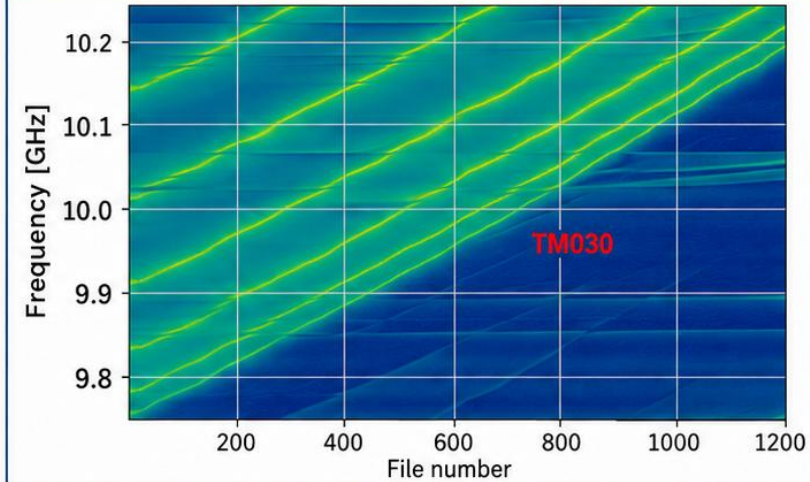
QUAX-LNL Haloscope

Future Plan

- Short prototype for testing
- Length = 243.4 mm
- Inner diameter 60.5 mm
- Sapphire from Lithuania
- Science mode TM030



Tested @ 4 K



G. Russo, Presented at the STDM Conference, Rome, Italy

QUAX-LNL Haloscope

Future Plan

1 Installation of a 14 T magnet

Expected summer 2026



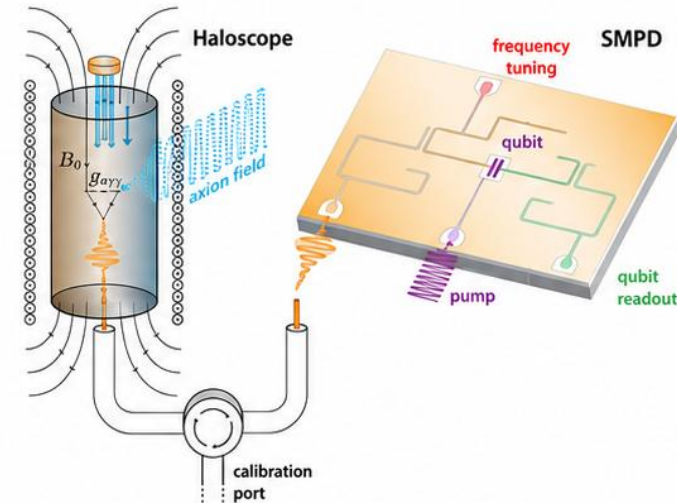
Model	JMTA-14T103
Coil configuration	Solenoid coil
Maximum central field	14 Tesla
Free bore diameter	103 mm
Outer diameter of the solenoid	400 mm excluding protrusion (tentative)
Total height of the solenoid	500 mm excluding protrusion (tentative)
Total weight (approximately)	250 kg (tentative)
Field direction	Vertical two-way
Operating current (nominal)	185 A (tentative)
Inductance	64 Henries (tentative)
Maximum sweep rate	14 Tesla / 60 minutes (tentative)
Field homogeneity (designed value)	More than 9 Tesla in a cylindrical area of 60mm diameter and 420 mm height



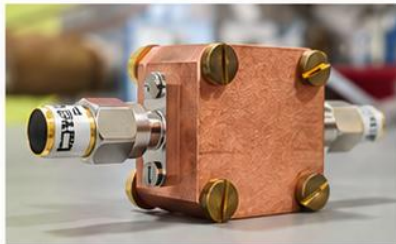
Photo 1 Appearance of JMTA-14T103.

2 Single Photon Microwave Detector (SMPD)

Help of Quantronics Group Paris



3



New TWPA

amplifier from
Grenoble optimized
for 10 – 11 GHz
band

4 Installation of new cryogenic system

- Wet assembly for magnet
- Dry unit for cavity and rf detection

This will allow for longer and safer operation.

QUAX-LNL Haloscope

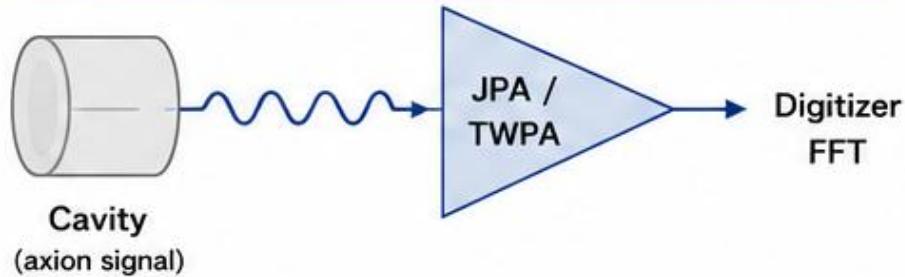
Future Plan



Single Microwave Photon Detector Readout

Beyond Standard Quantum Limit

1. SQL Linear Amplifier (measures amplitude)



What is measured?

Field quadratures X and P (amplitude & phase)

Heisenberg uncertainty: $\Delta X \Delta P \geq 1/2$

⇒ Vacuum fluctuations (1/2 photon) are always added.

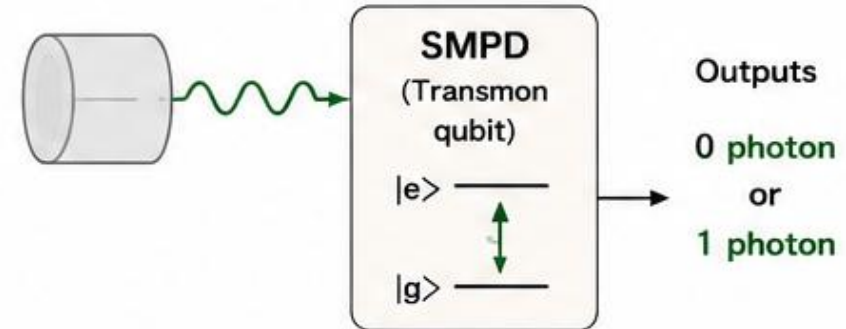
Noise (SQL)

Noise-equivalent power:

$$P_{\text{SQL}} = h\nu \quad \text{SNR}_{\text{SQL}} = \frac{P_a}{h\nu} \sqrt{\frac{t}{\Delta\nu_a}}$$

P_a : axion signal power, $\Delta\nu_a$: axion linewidth, t : integration time

2. SMPD (measures photon number)



What is measured?

Photon number $n = a^\dagger a$ (YES / NO)

Noise (SMPD)

Counts obey Poisson statistics.

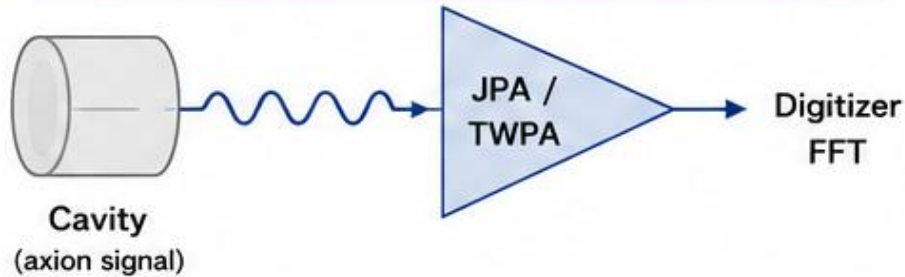
Total count rate: $\Gamma_{\text{tot}} = \Gamma_s + \Gamma_{\text{dc}}$
 $\downarrow$ signal $\downarrow$ dark counts (thermal + intrinsic)

$$\text{SNR}_{\text{SMPD}} \approx \frac{\Gamma_s}{\sqrt{\Gamma_{\text{dc}}}} \sqrt{t} \quad (\text{weak-signal limit})$$

Single Microwave Photon Detector Readout

Beyond Standard Quantum Limit

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What is measured?

Field quadratures X and P (amplitude & phase)

Heisenberg uncertainty: $\Delta X \Delta P \geq 1/2$

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Noise (SQL)

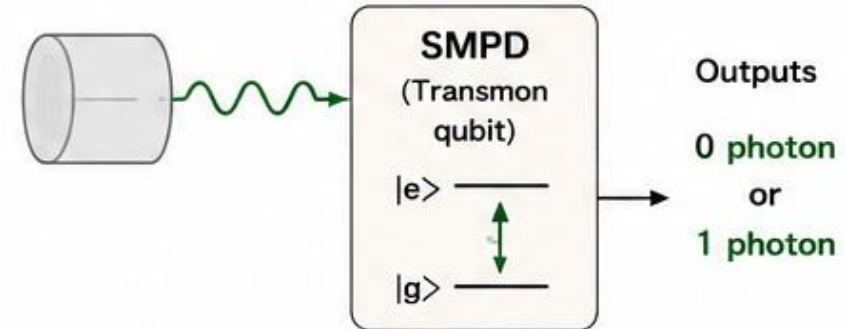
Noise-equivalent power:

$$P_{\text{SQL}} = h\nu$$

$$\text{SNR}_{\text{SQL}} = \frac{P_a}{h\nu} \sqrt{\frac{t}{\Delta\nu_a}}$$

P_a : axion signal power, $\Delta\nu_a$: axion linewidth, t : integration time

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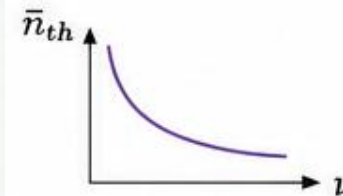


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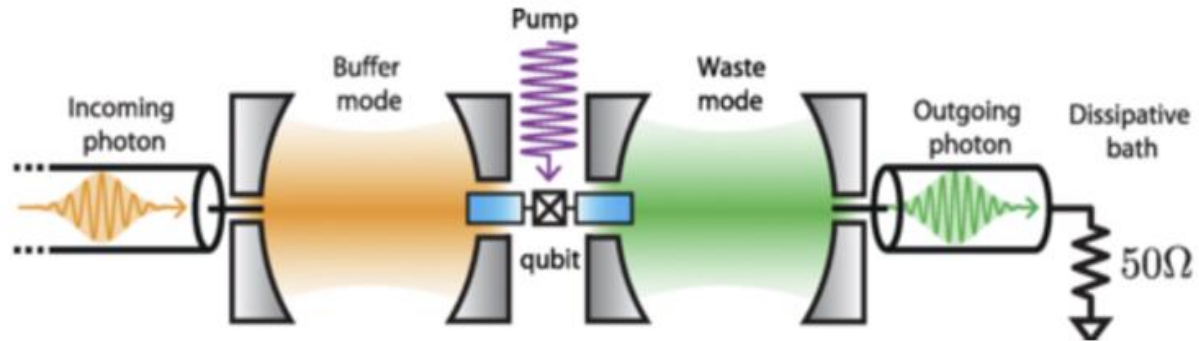
$$\bar{n}_{th} = \frac{1}{e^{h\nu/k_B T} - 1}$$

Thermal photons decrease exponentially with frequency.

Single Microwave Photon Detector Readout

Working Principle

Detector principle



A pump-activated four-wave mixing interaction **coherently converts** a buffer photon into a qubit excitation while emitting a photon into the waste mode.

4WM resonance condition

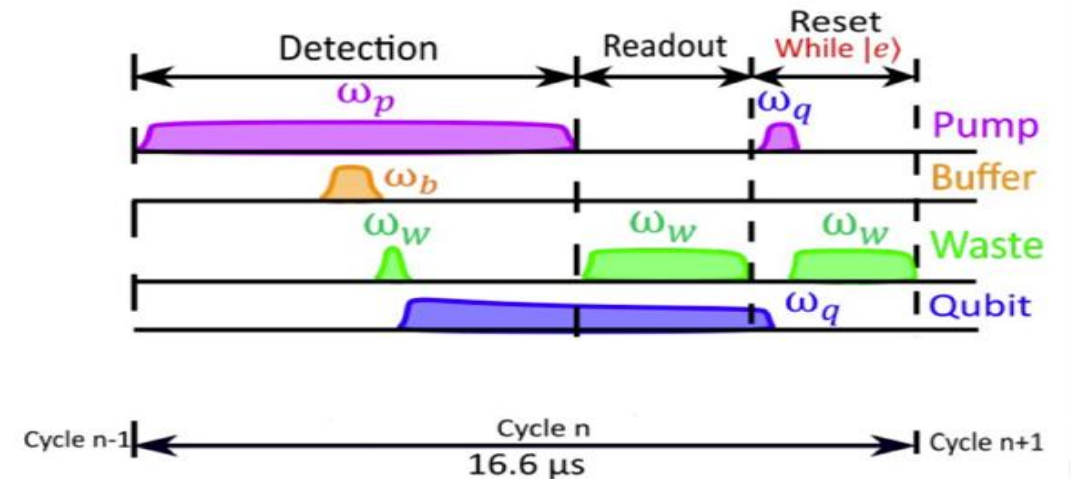
$$\omega_b + \omega_p = \omega_q + \omega_w - \chi_w$$

(χ_w : anharmonicity of the waste mode)

- ω_b buffer mode frequency
- ω_p pump frequency
- ω_q qubit frequency
- ω_w waste mode frequency

Cyclic operation of one detection cycle

The detector operates in repeated **detection-readout-reset** cycles.



Outcome of each cycle:

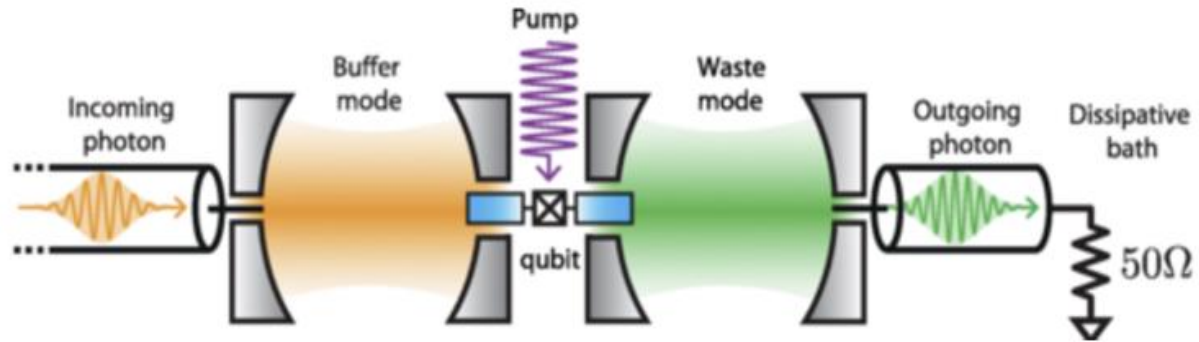


No click ($|g\rangle$)
Click ($|e\rangle$)

Single Microwave Photon Detector Readout

Working Principle

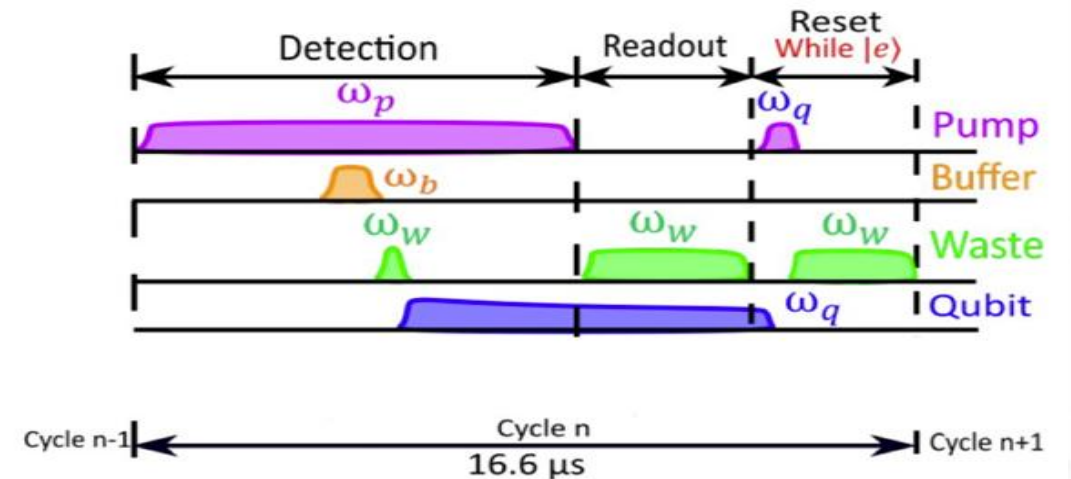
Detector principle



Phys. Rev. X 10, 021038 (2020) $\leftarrow$ 1.3 counts/ms
Nature 600, 434–438 (2021) $\leftarrow$ spin fluorescence detection
Nature 619, 276–281 (2023) $\leftarrow$ single spin flip
Phys. Rev. Appl. 21, 014043 (2024) $\leftarrow$ 85 counts/s
more at <https://iramis.cea.fr/en/spec/gq/>

Cyclic operation of one detection cycle

The detector operates in repeated **detection**–**readout**–**reset** cycles.



Outcome of each cycle:



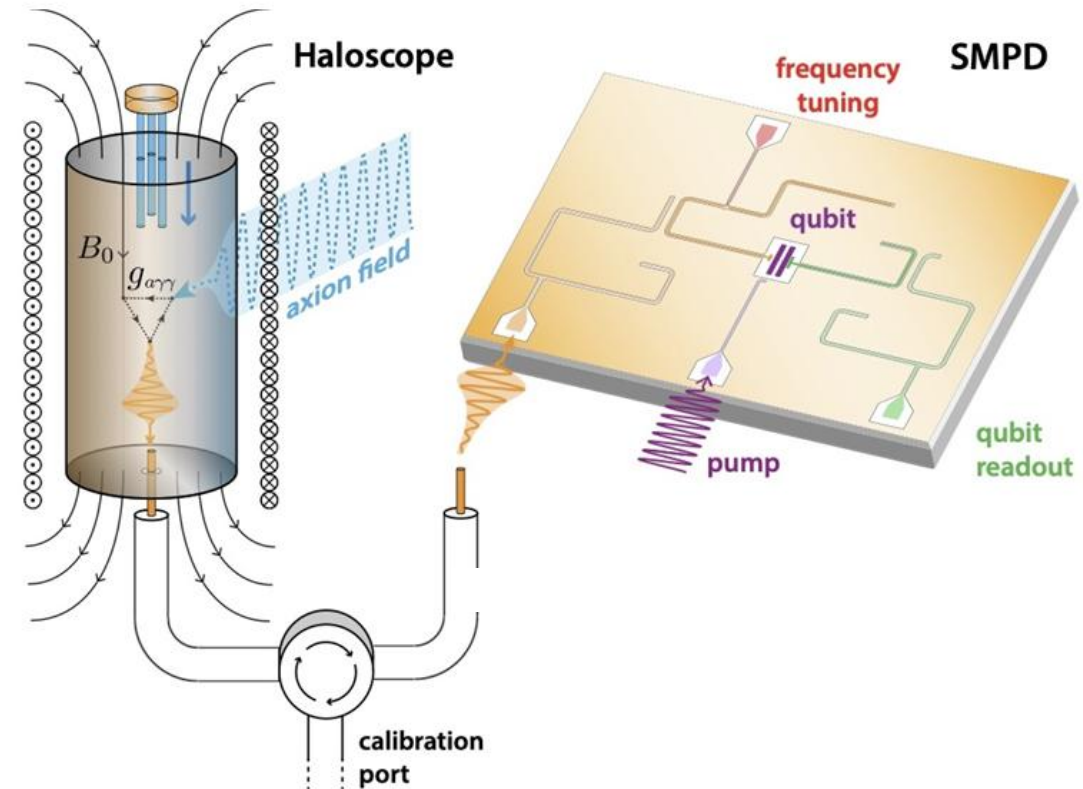
No click ($|g\rangle$)
Click ($|e\rangle$)

Single Microwave Photon Detector Readout

Pathfinder

- a **transmon qubit-based single microwave photon detector (SMPD)** is used to readout the cavity mode
 - **TWPA** for dispersive readout of the qubit state
 - hybrid (normal-superconducting) cavity TM_{010} at 7.37 GHz
tunable by a triplet of rods
 $Q_0 = 9 \times 10^5$ at **2 T-field**
 - **T=14 mK**
@ fridge Quantronics lab (CEA, Saclay)
- investigated the background,
and set a limit to $g_{a\gamma\gamma}$ [0.5 MHz band]
X20 compared to SQL

see also: Kuo *et al* Phys. Rev. D 111, 123018 (2025)



Phys. Rev. X 15, 021031 (2025)

C. Braggio, Presented at the STDM Conference, Rome, Italy

Single Microwave Photon Detector Readout Pathfinder

C. Braggio, Presented at the STDM Conference, Rome, Italy

- ◉ a **transmon-based single microwave photon detector (SMPD)** is used to readout the cavity mode
 - ◉ **TWPA** for dispersive readout of the qubit state
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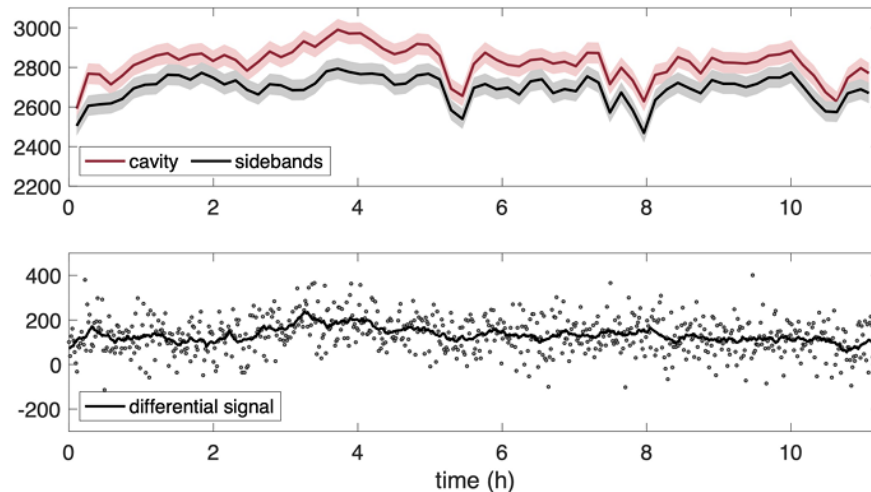
SMPD (top) and cavity

SC magnet

Single Microwave Photon Detector Readout

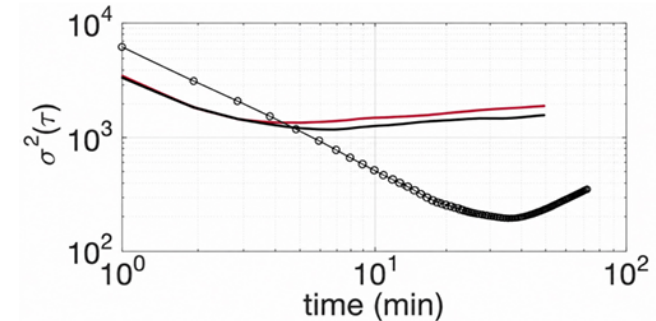
Background Limited Search: Dark Counts

C. Braggio, Presented at the STDM Conference, Rome, Italy



Phys. Rev. X 15, 021031 (2025)

- counts at $\omega_b = \omega_c$ registered in 28.6 s
 $\iff$ **average dark count rate** $\alpha \sim 90$ counts/s
- excess** at $\omega_c \leftrightarrow$ the cavity sits at a slightly higher temperature (few mK)
- both the counts at resonance and on sidebands $\omega_b = \omega_c \pm 1, 2$ MHz vary **beyond statistical uncertainty** expected for poissonian counts



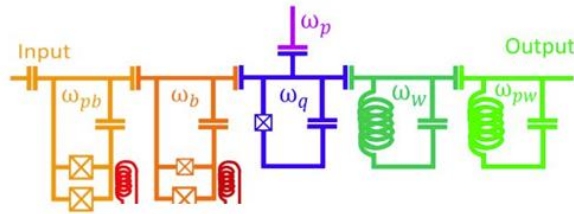
Allan variance $\rightarrow$ max integ. time ~ 10 min.

$\rightarrow \alpha = \alpha_q + \alpha_{4WM} + \alpha_{th}$, where $\alpha_{th} \propto \bar{n}_{th} = \frac{1}{e^{h\nu/kT} - 1}$
 finite flux of *thermal* photons at the SMPD input

Single Microwave Photon Detector Readout

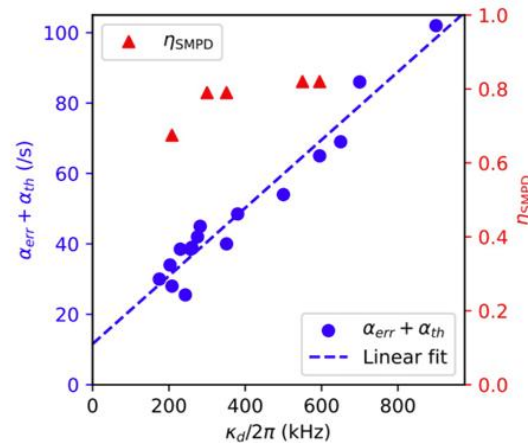
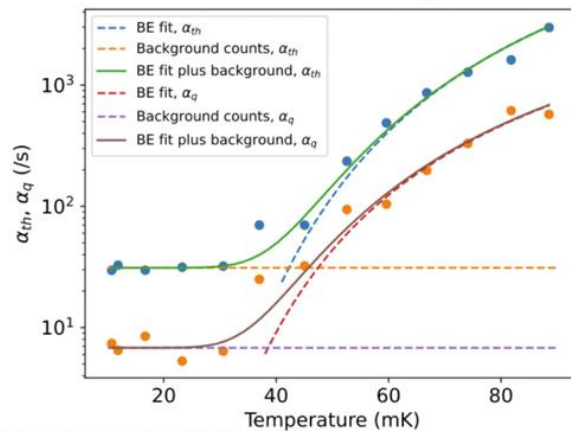
New Generation SMPD in Our Lab

C. Braggio, Presented at the STDM Conference, Rome, Italy



- Purcell resonator for **tuning buffer bandwidth**
- DC flux lines for tuning ω_{pb} and ω_b

$$\alpha_{th} = \frac{\bar{n}_{th,b} \eta_{SMPD}(\omega_b) \kappa_d}{4}$$

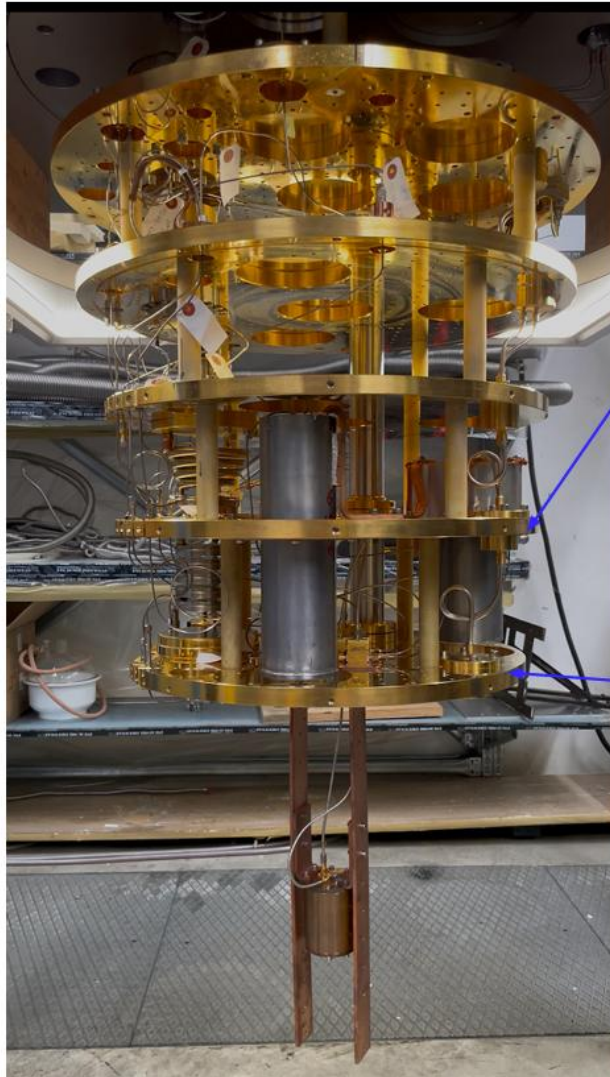


<https://arxiv.org/pdf/2501.07354>

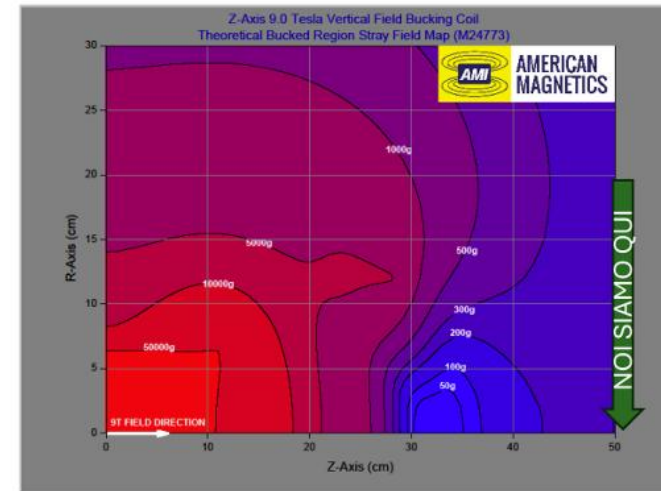
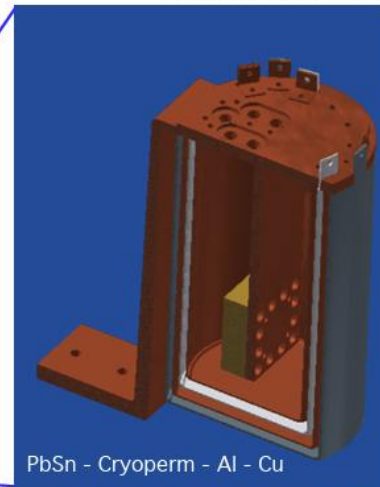


Single Microwave Photon Detector Readout

Experimental Apparatus



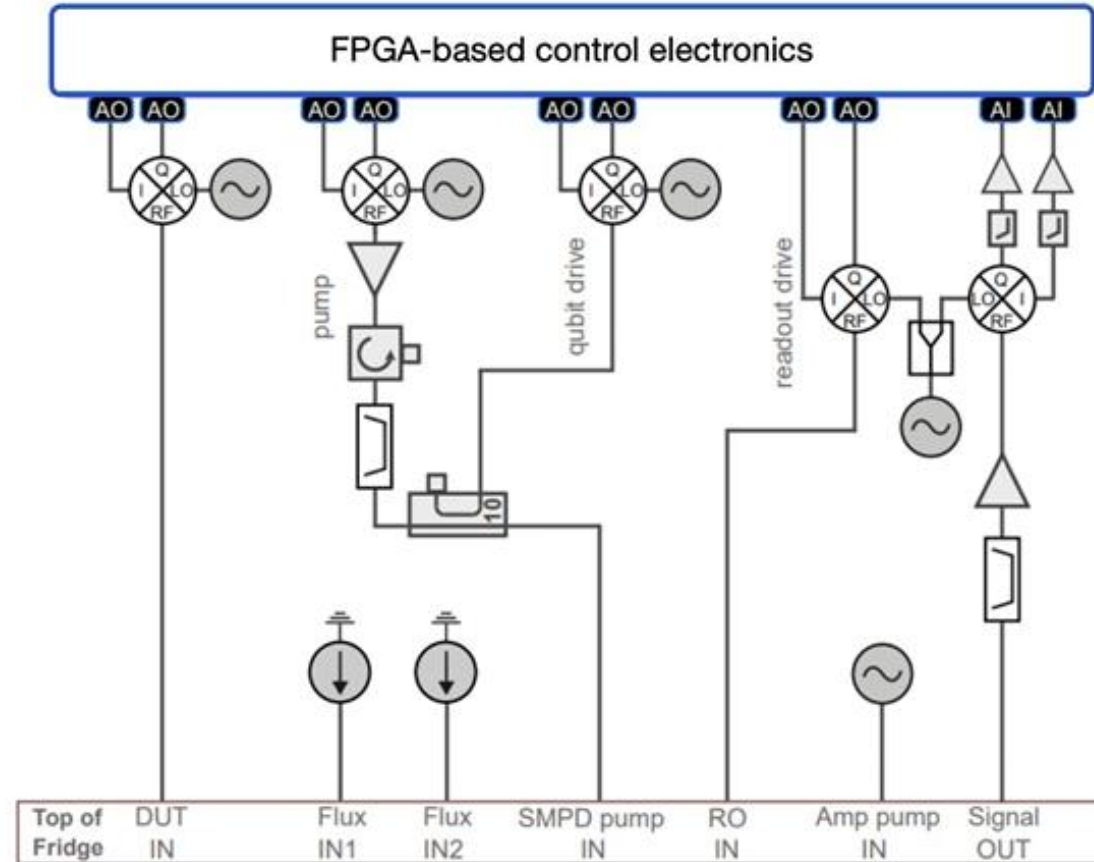
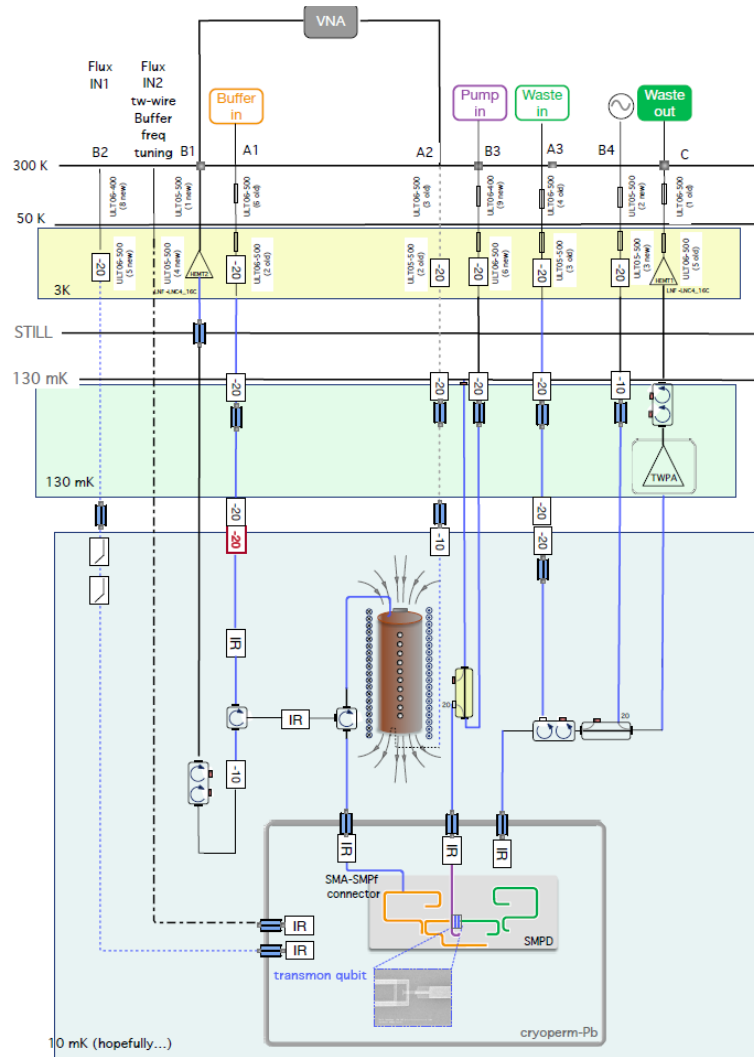
inside of SMPD box



C. Braggio, Presented at the STDM Conference, Rome, Italy

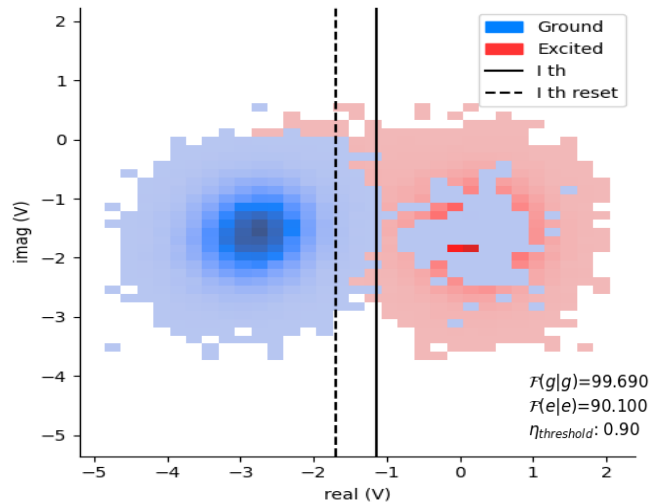
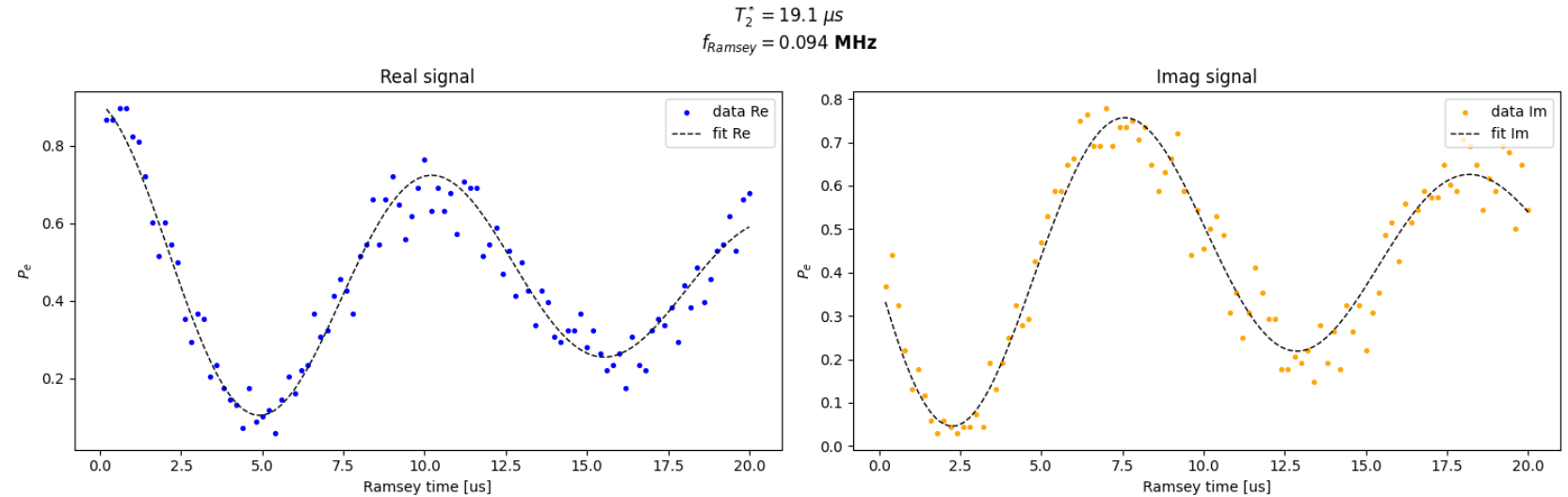
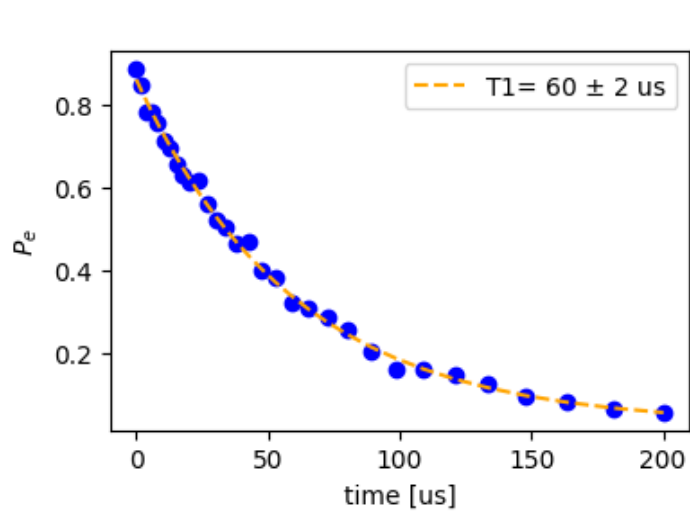
Single Microwave Photon Detector Readout

Experimental Apparatus



Single Microwave Photon Detector Readout

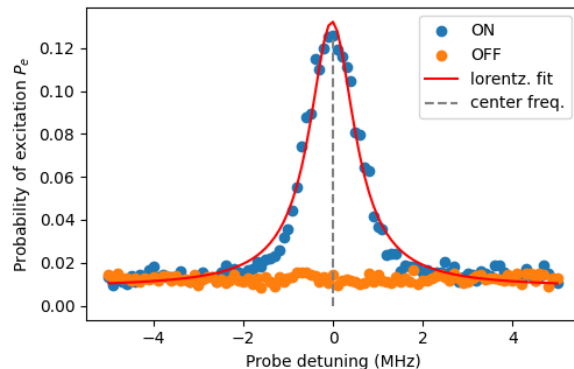
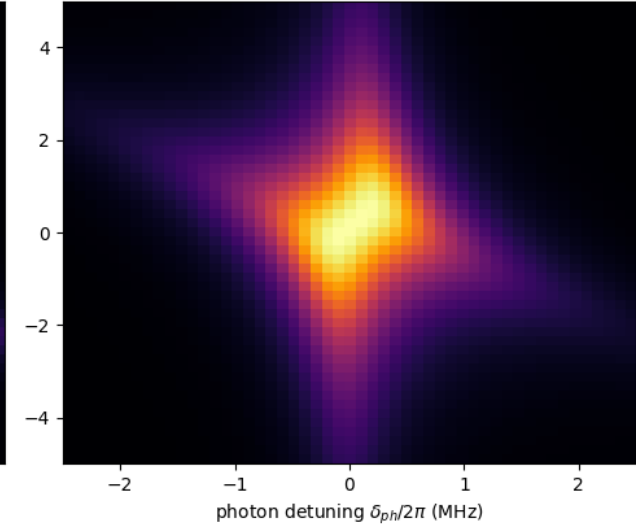
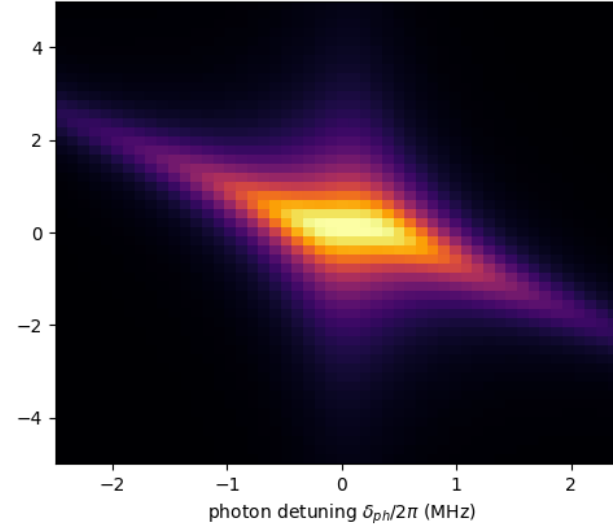
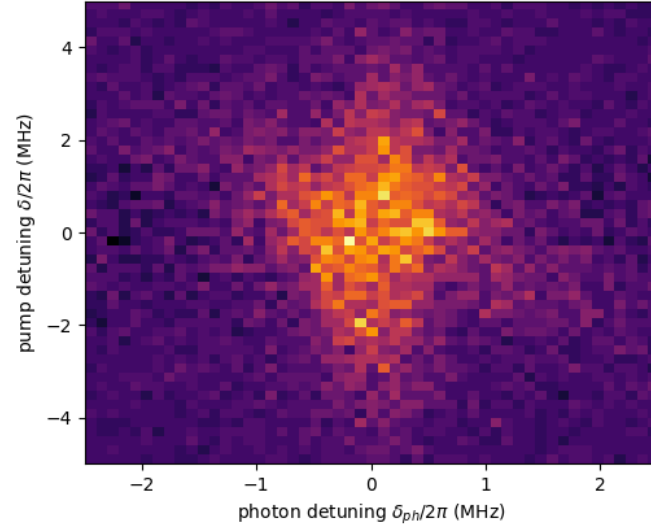
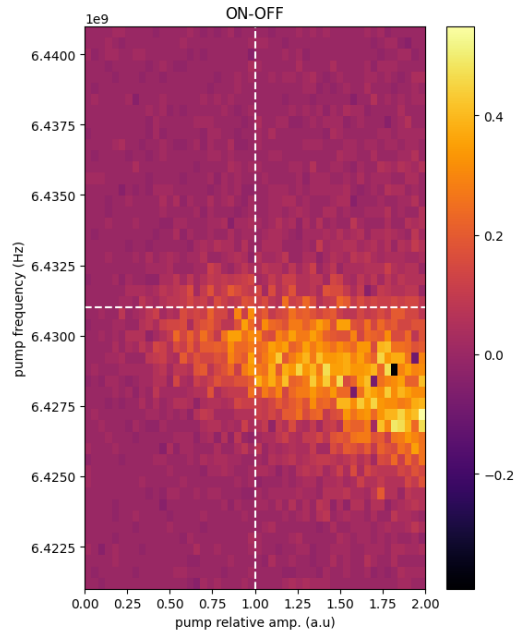
Qubit Characterization & Photon Counting with 4 Wave Mixing



- VNA: TWPA optimization, buffer/waste resonator response
 FPGA-based Qubit Measurement : T_1 , T_2^* , $T_2\text{Echo}$, Single-Shot, Qubit Reset
 Photon Counting
- (1) Calibrate Photon Flux (Stark-Shift)
 - (2) Four-Wave Mixing Calibration (pump amplitude, frequency & detuning)
 - (3) Buffer Detection Bandwidth
 - (4) Photon Counting with Cyclic Sequence (efficiency & dark count)

Single Microwave Photon Detector Readout

Qubit Characterization & Photon Counting with 4 Wave Mixing



VNA: TWPA optimization, buffer/waste resonator response

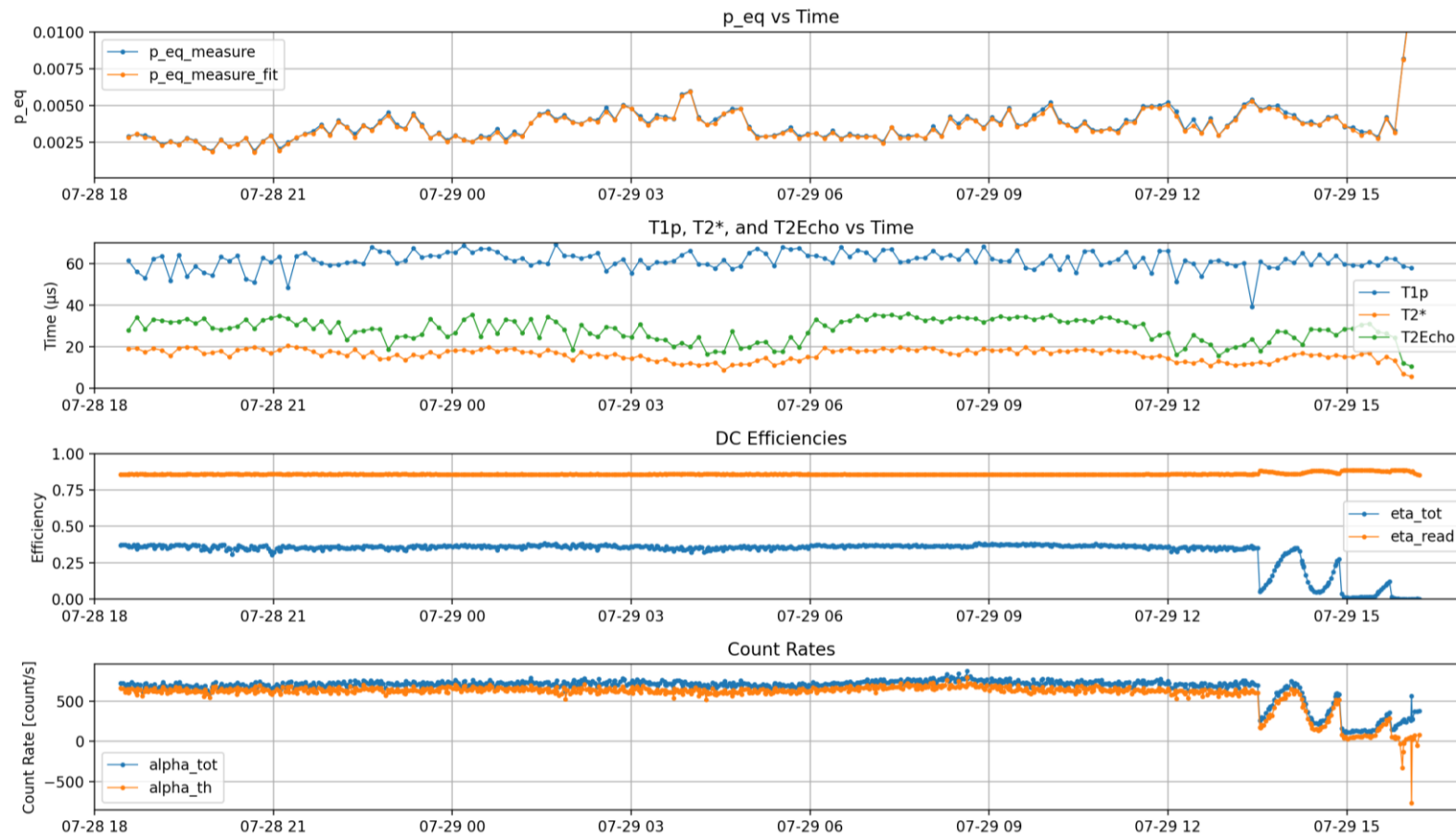
FPGA-based Qubit Measurement : T1, T2*, T2Echo, Single-Shot, Qubit Reset

Photon Counting

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Single Microwave Photon Detector Readout

New Generation SMPD in Our Lab

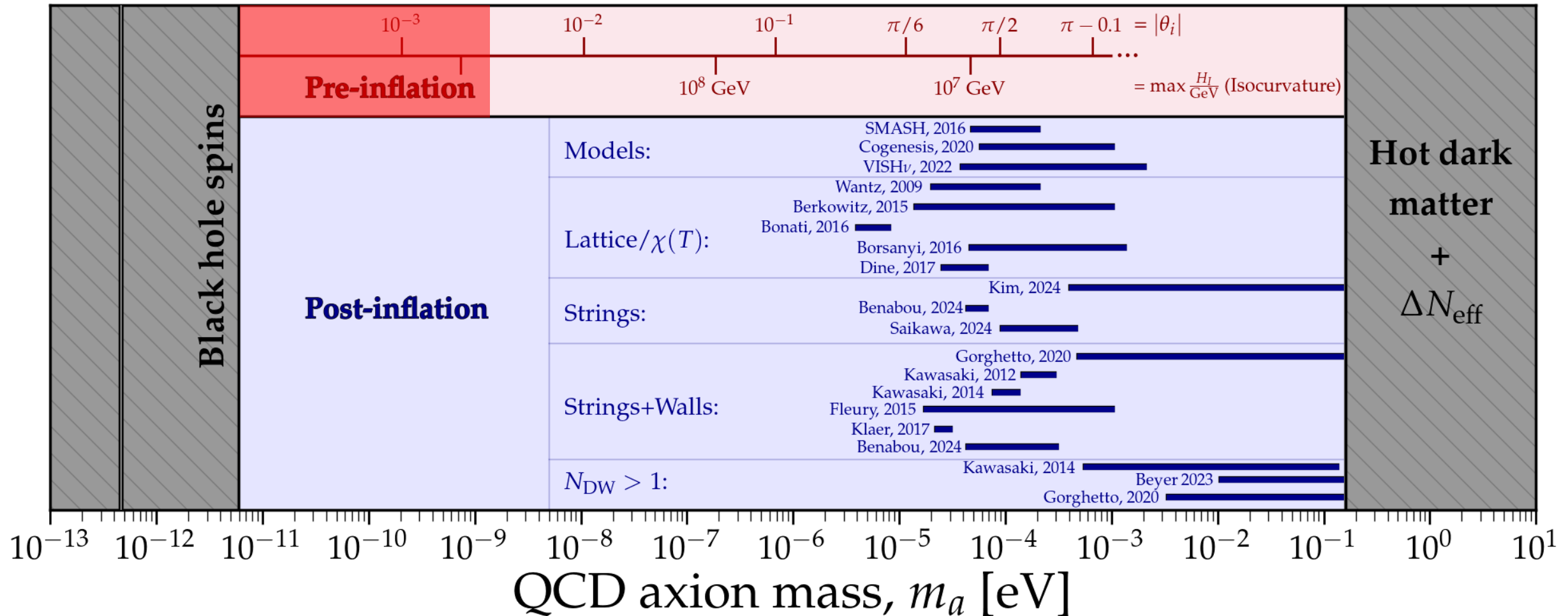


D. Ahn, INFN-SQMS PostDoc (formerly at CAPP)
G. Sardo Infrirri, UniPD PhD student

SQUID Readout

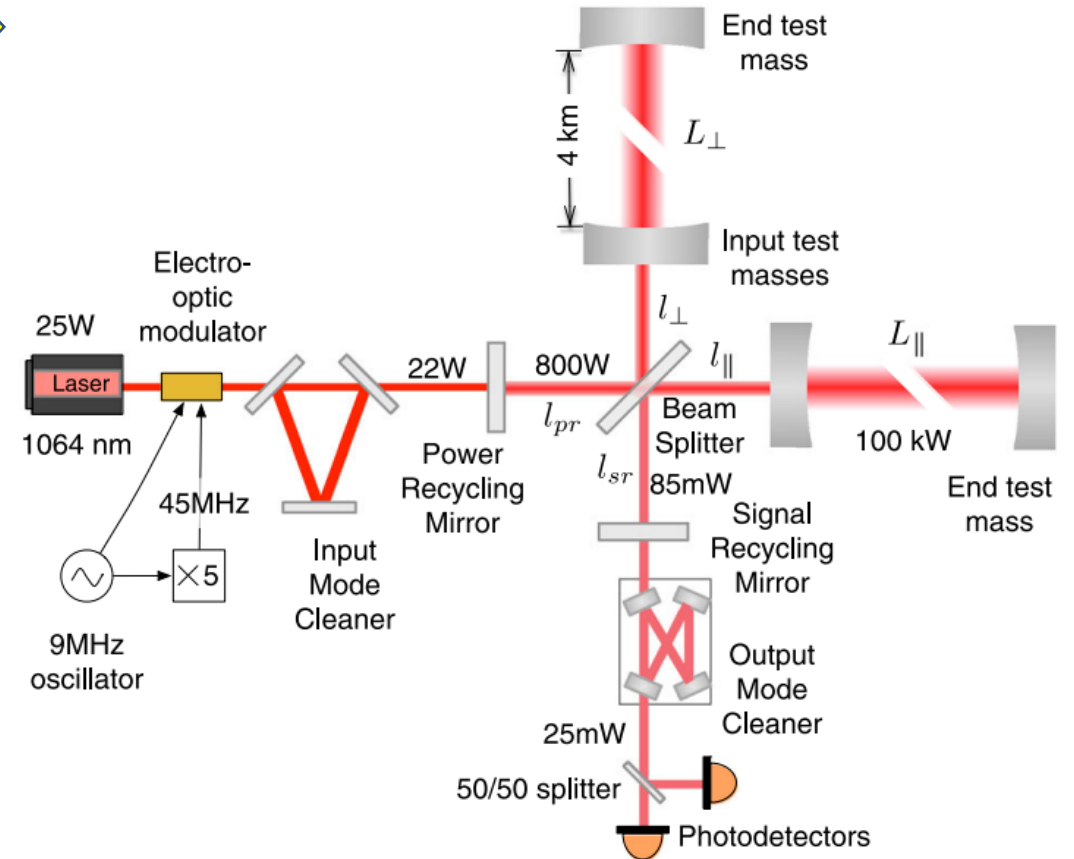
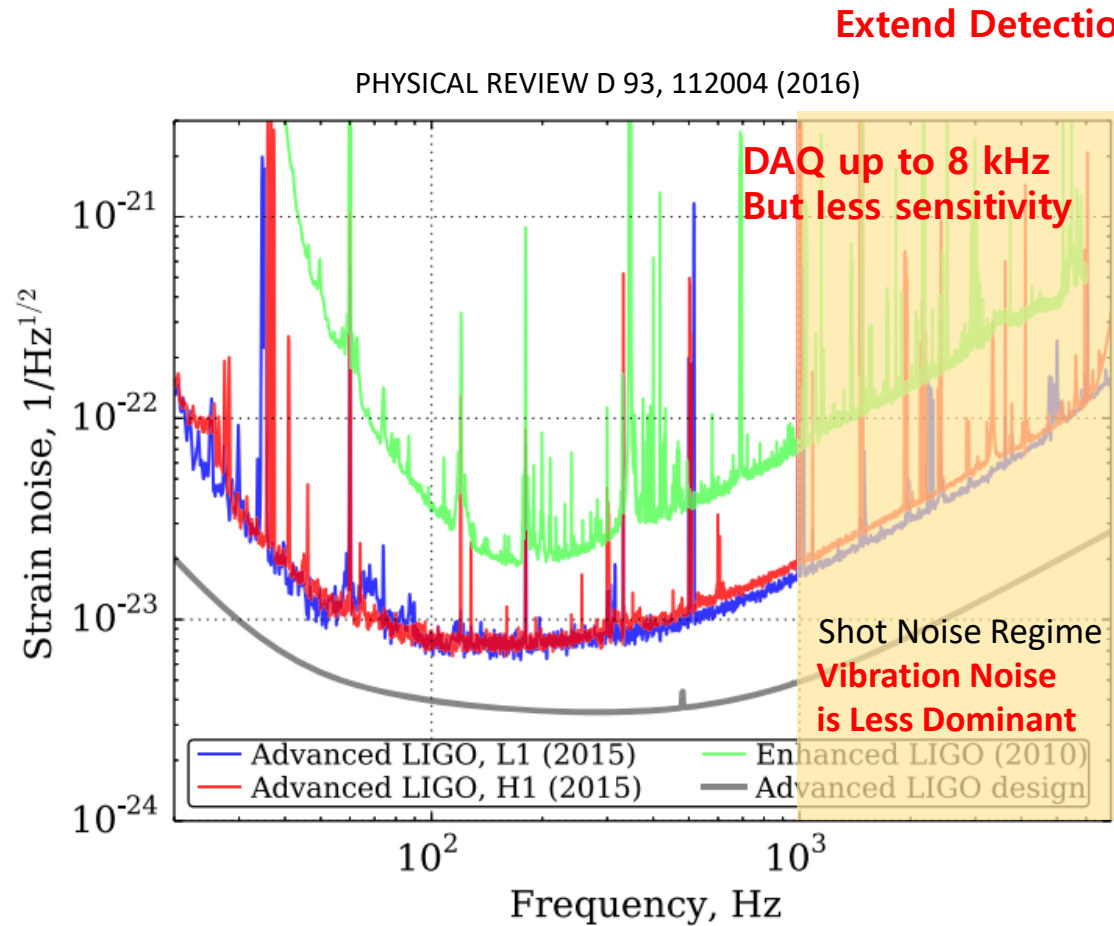
Motivation: kHz Axion Dark Matter Detection

<https://cajohare.github.io/AxionLimits/>



SQUID Readout

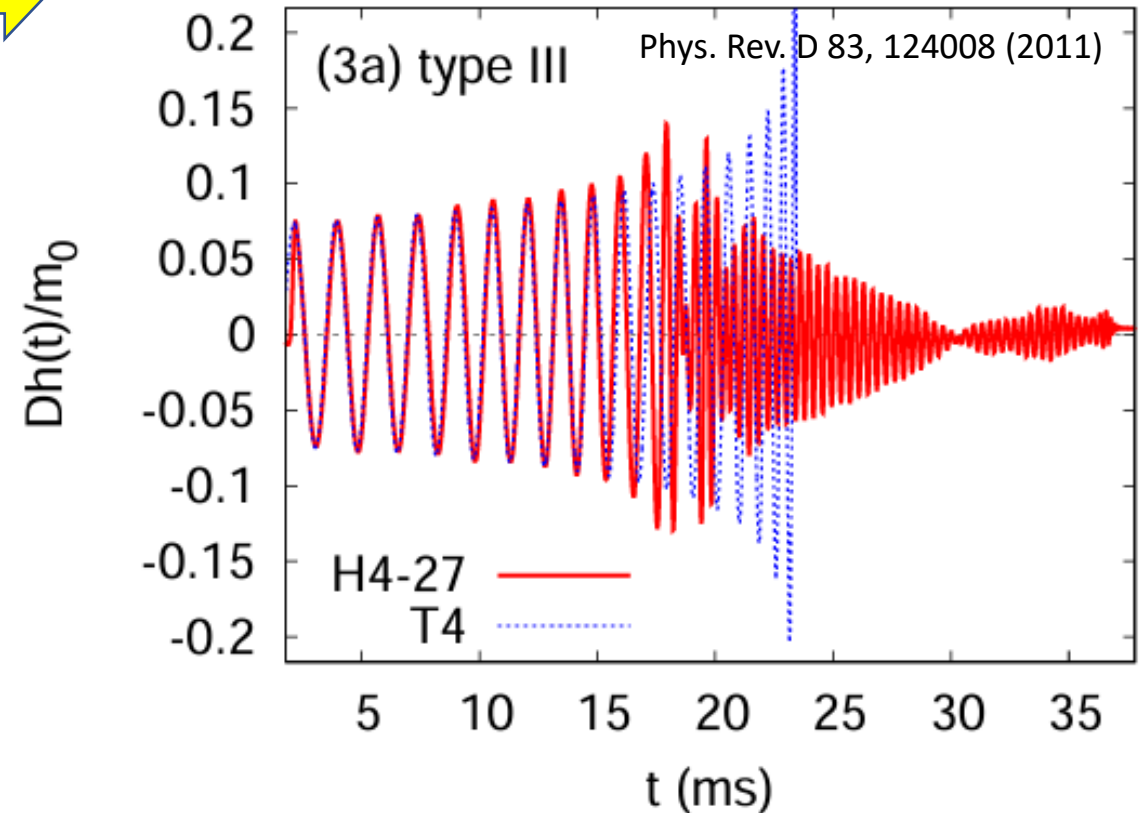
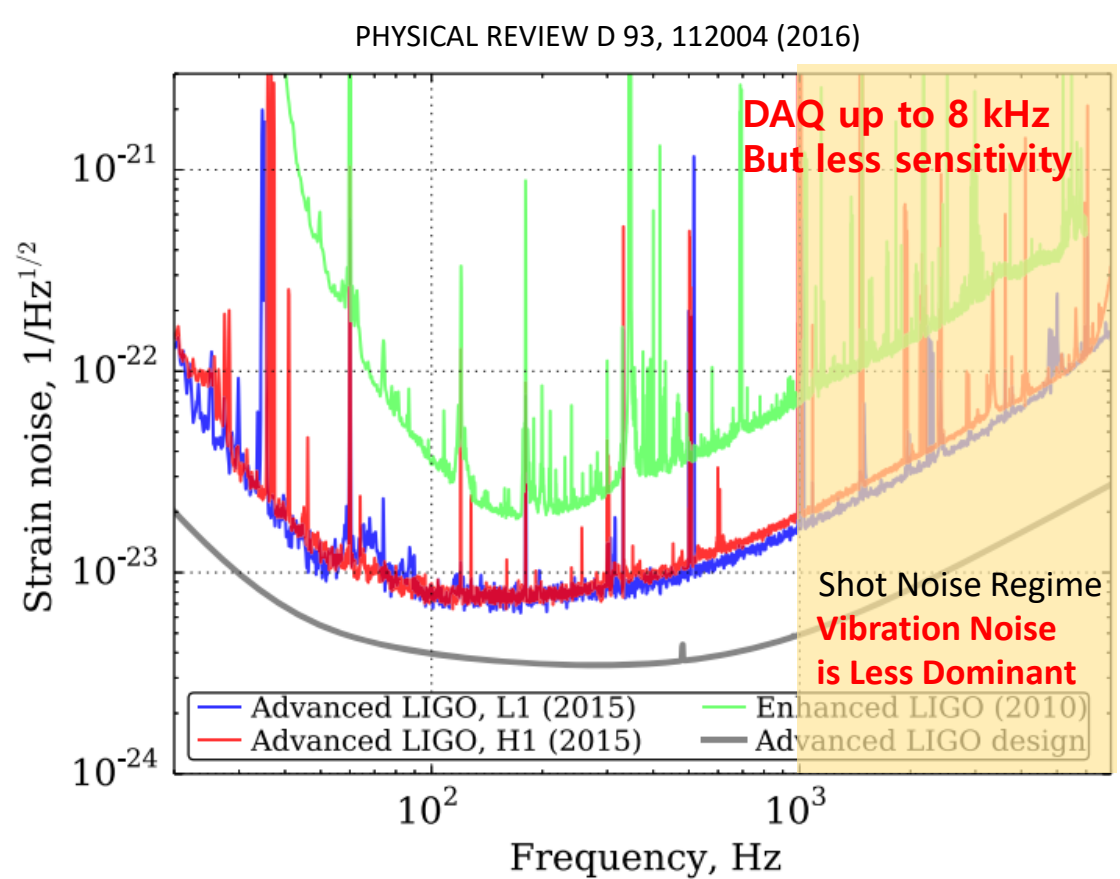
Motivation: Gravitational Wave Detection



SQUID Readout

Motivation: Gravitational Wave Detection

Extend Detection Frequency $O(10 \text{ kHz})$

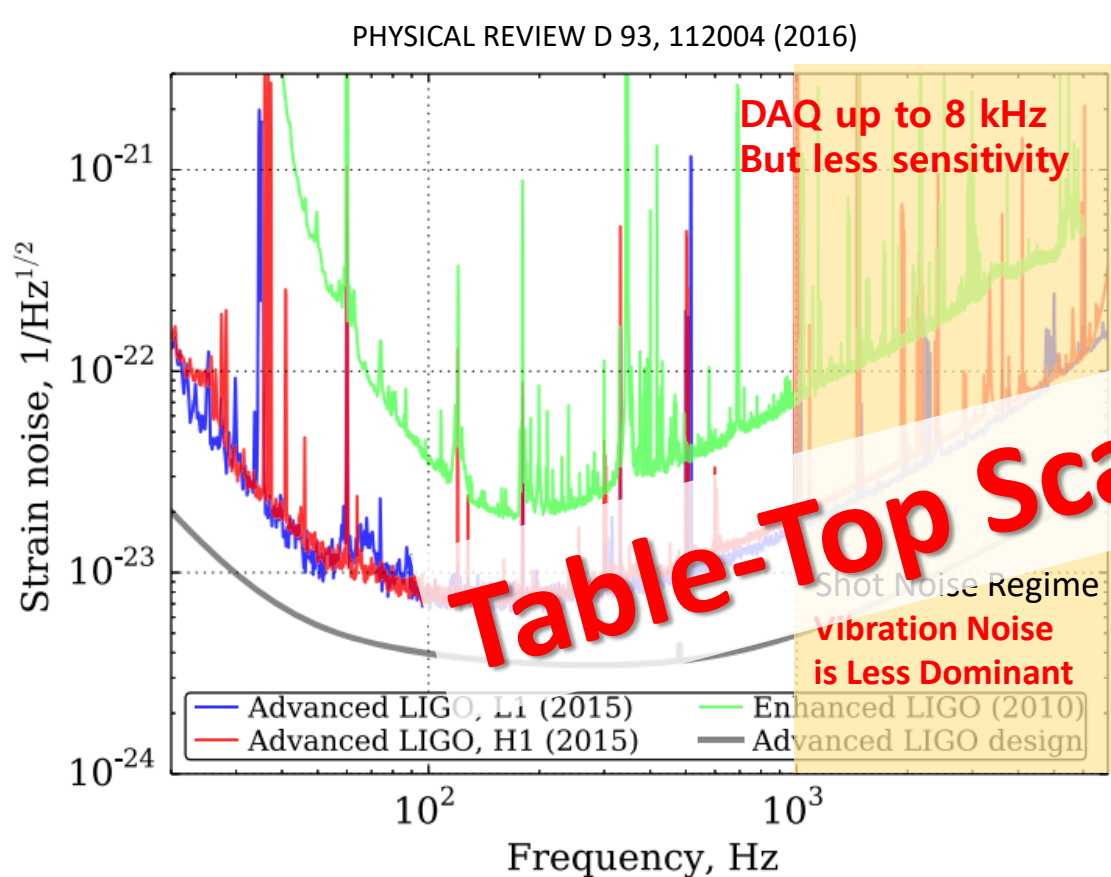


Neutron Star Binary Signal Can Give kHz Gravitational Wave Signal

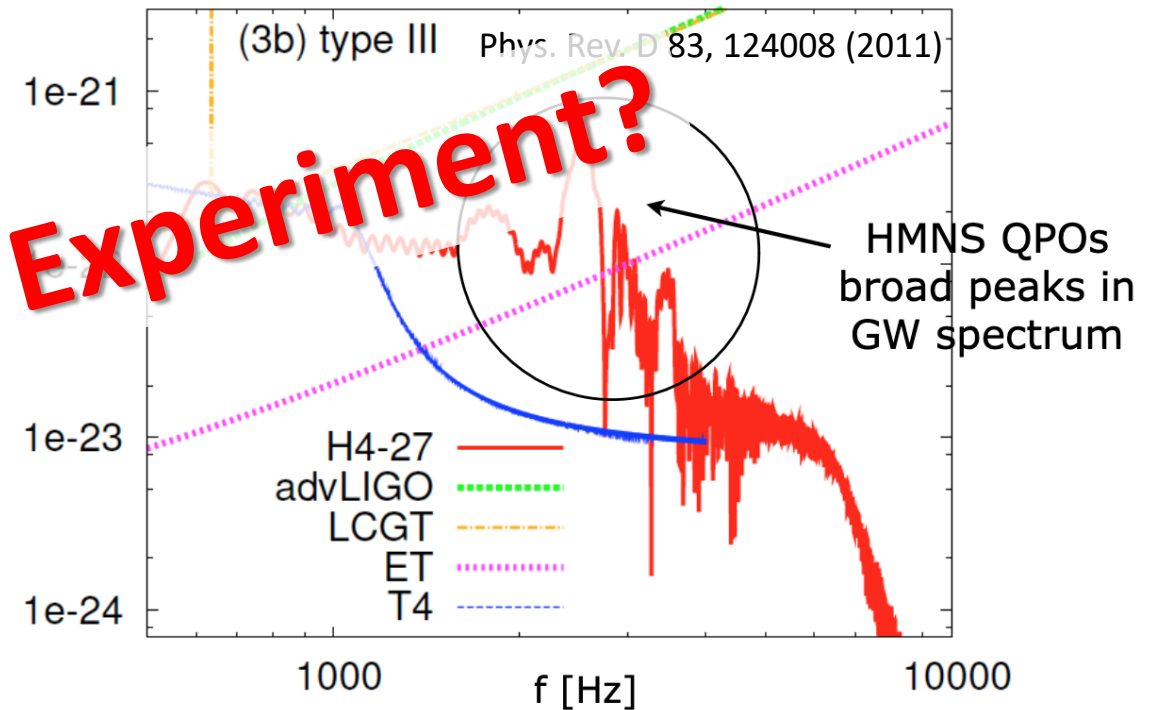
SQUID Readout

Motivation: Gravitational Wave Detection

Extend Detection Frequency $O(10 \text{ kHz})$



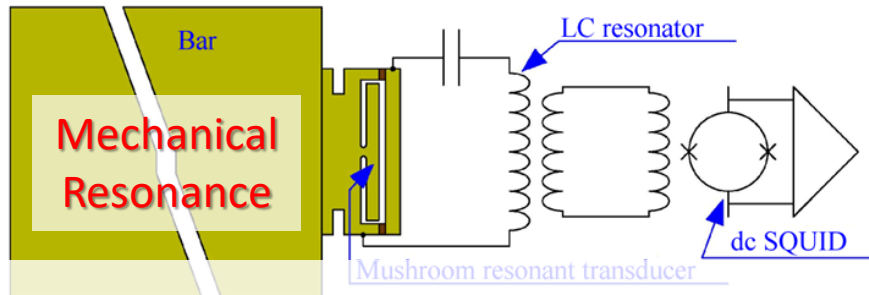
Fourier spectrum (h_{eff} @ 100 Mpc)



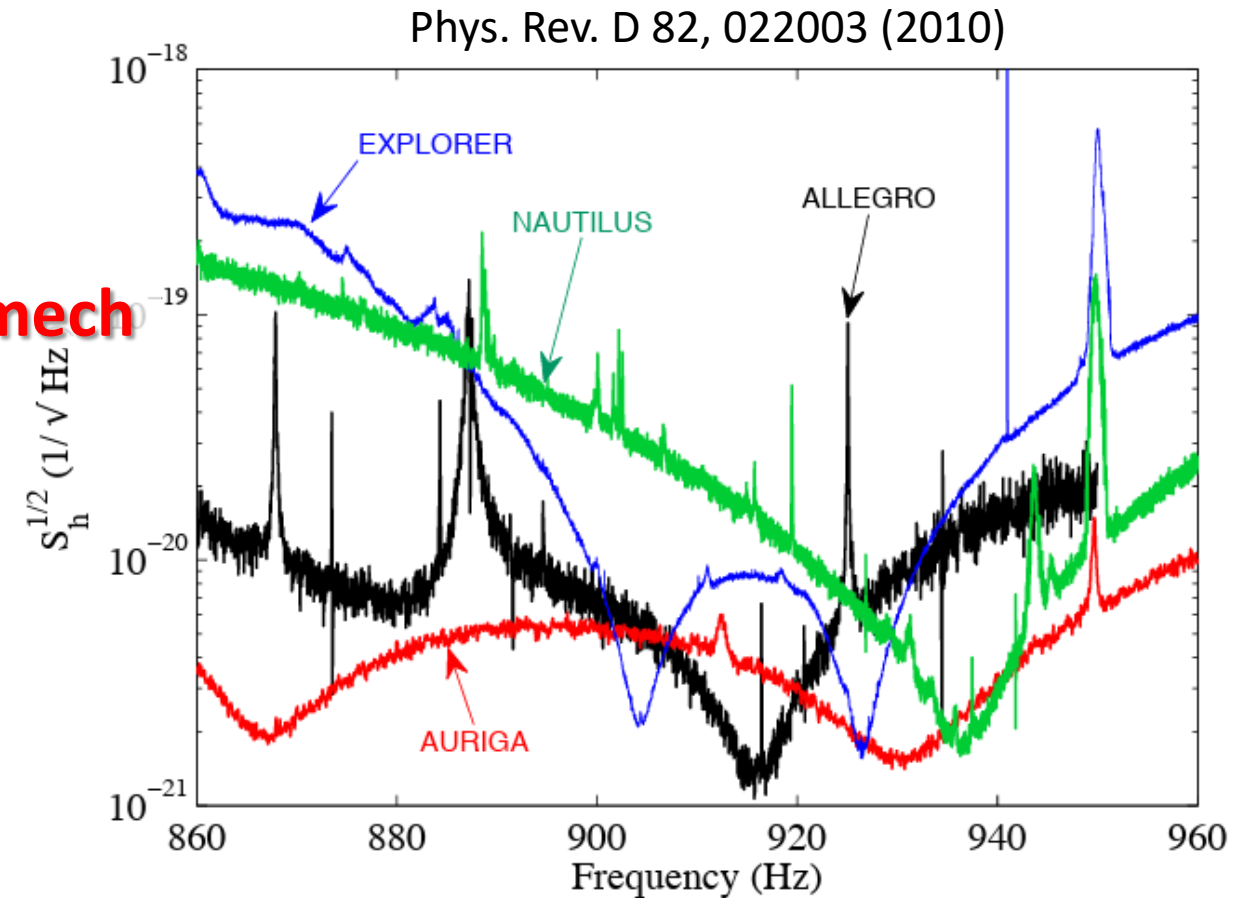
Higher Frequency Signals up to 10 kHz

SQUID Readout

Motivation: Bar Detector with SQUID Readout

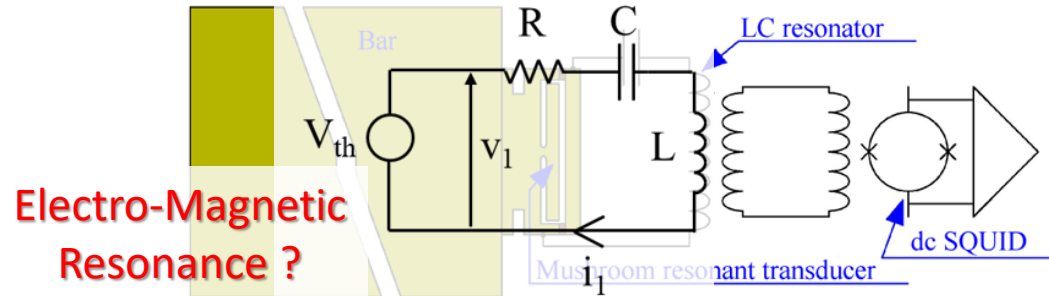


Signal $\propto$ Mass x Length x Q_{mech}

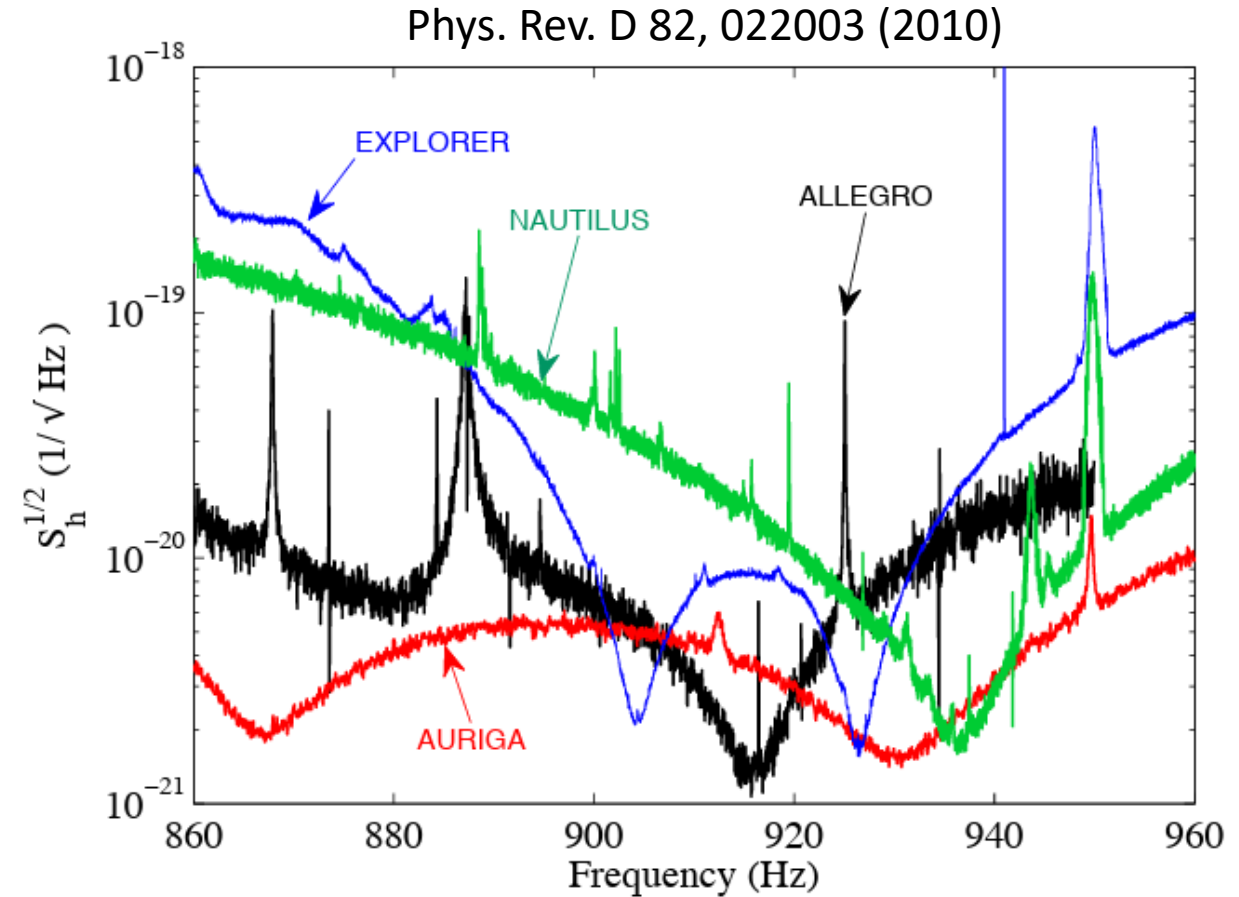


SQUID Readout

Magnetic Field Detector with SQUID Readout

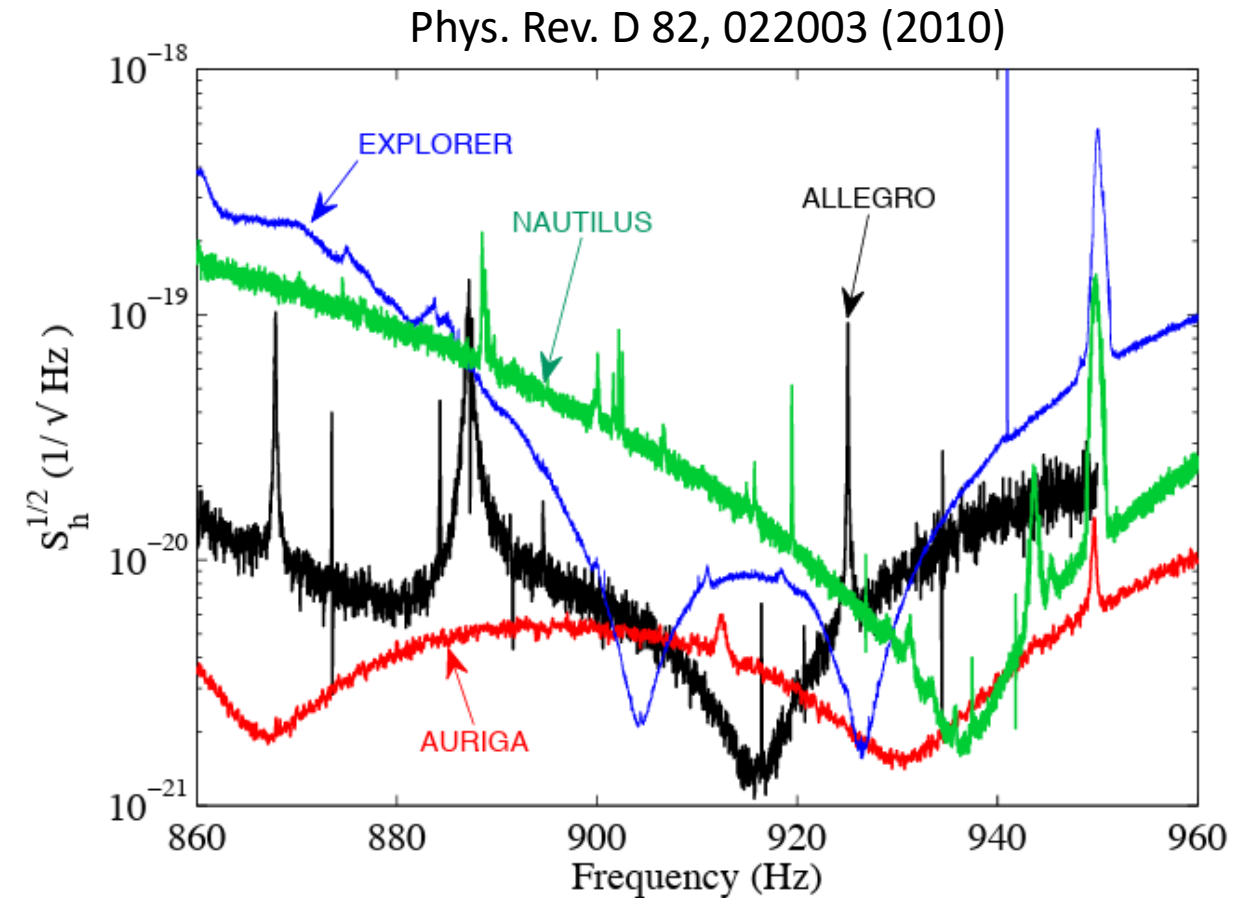
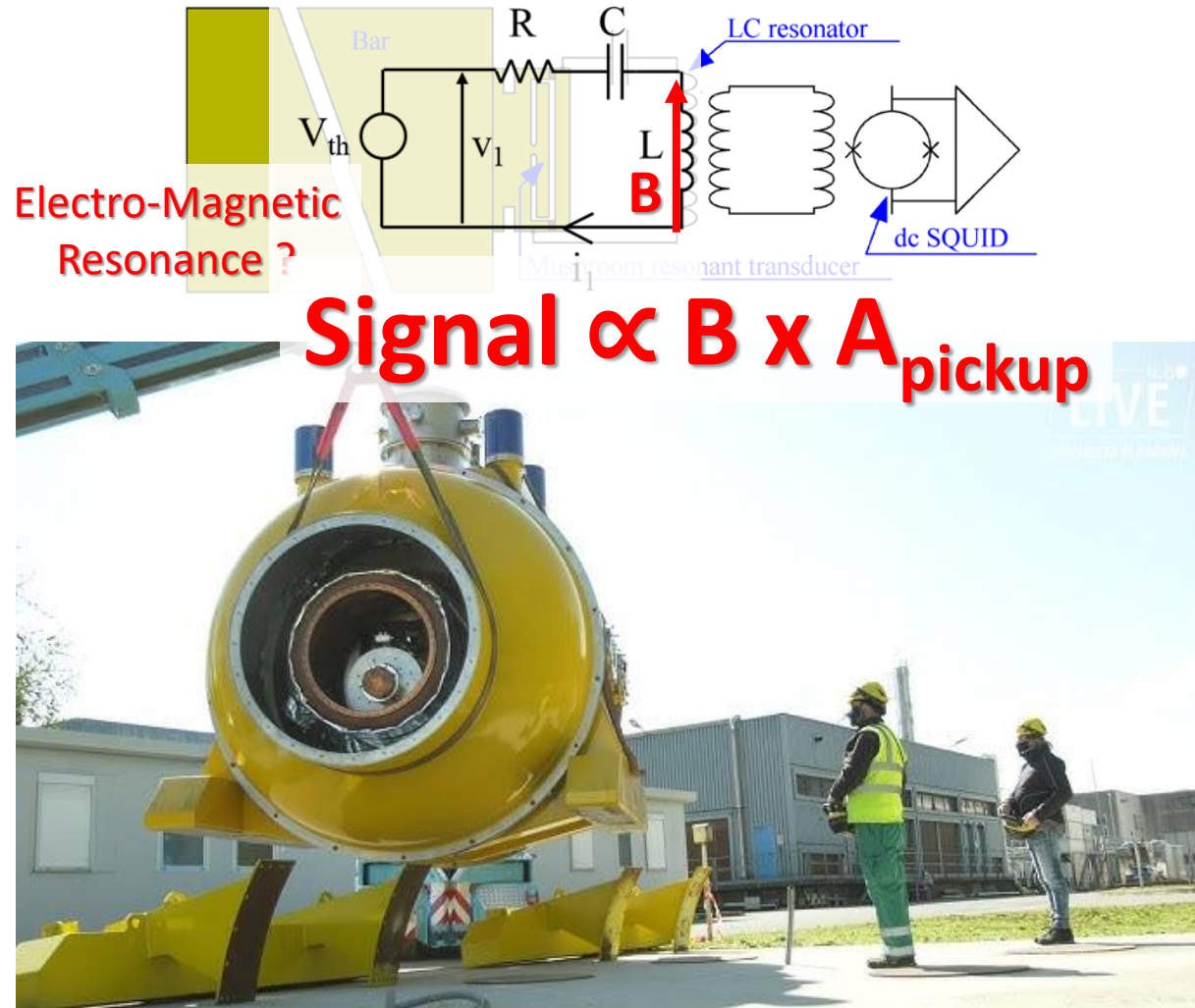


Pierluigi Fortini, Enrico Montanari,
Antonello Ortolan, Gerhard Schäfer,
Gravitational Wave Interaction with Normal
and Superconducting Circuits, Annals of
Physics, Volume 248, Issue 1, (1996)



SQUID Readout

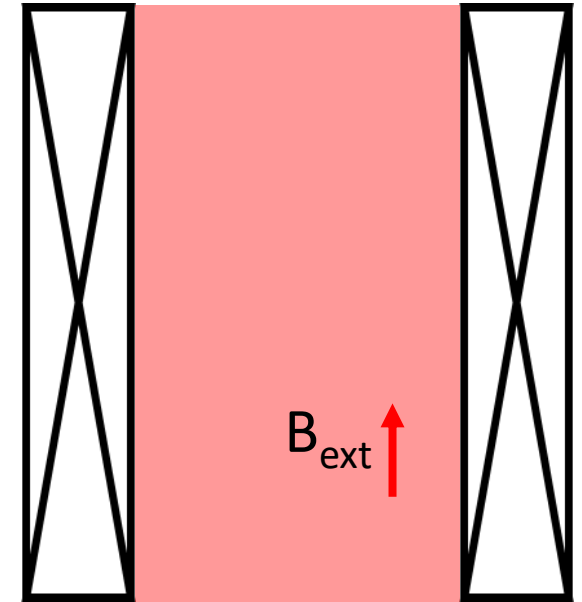
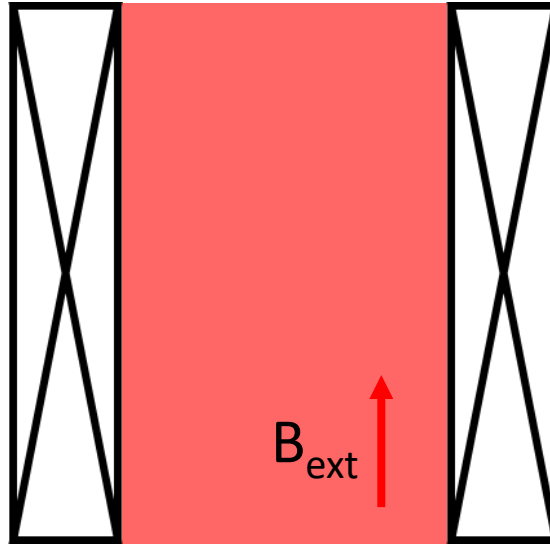
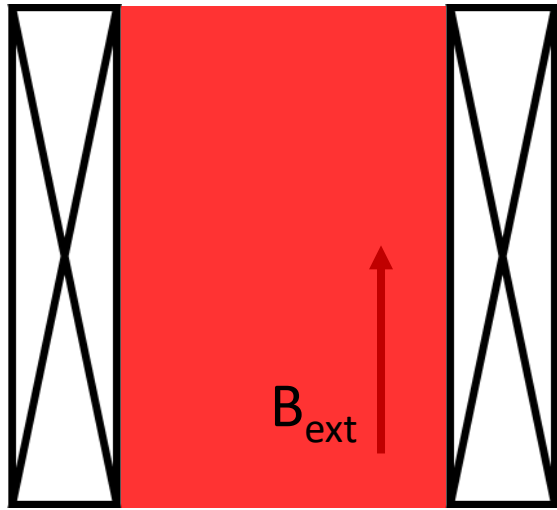
Magnetic Field Detector with SQUID Readout



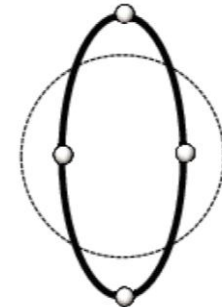
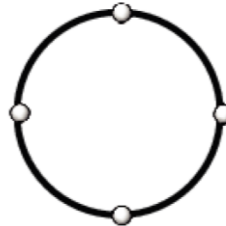
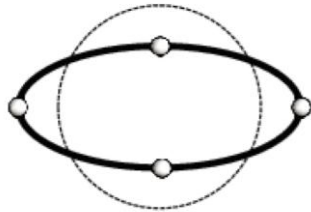
SQUID Readout

Magnetic Field Detector with SQUID Readout

$$\lambda_{GW}(\sim 10 \text{ km}) \gg \lambda_{Det}(\sim 1 \text{ m})$$



h^+



SQUID Readout

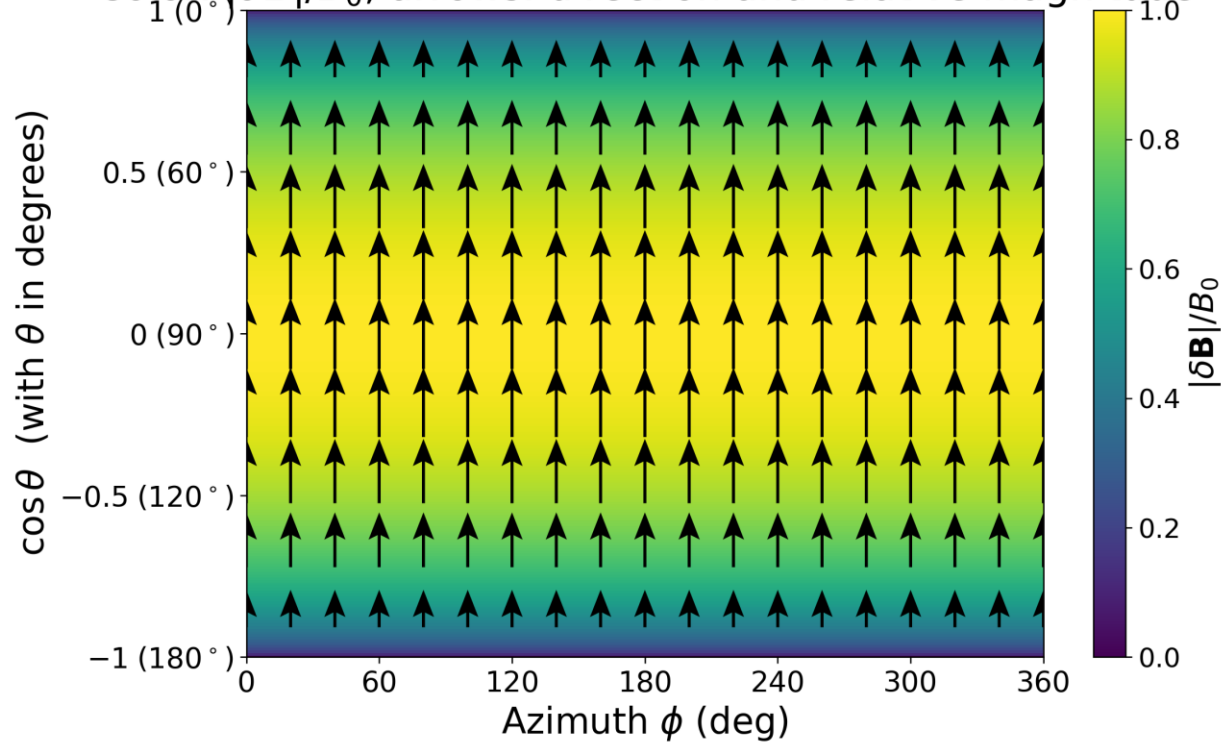
Magnetic Field Detector with SQUID Readout

$$\lambda_{GW}(\sim 10 \text{ km}) \gg \lambda_{Det}(\sim 1 \text{ m})$$

$$\delta B_z = B_0 h_{GW} \sin^2 \theta \cos 2\psi$$

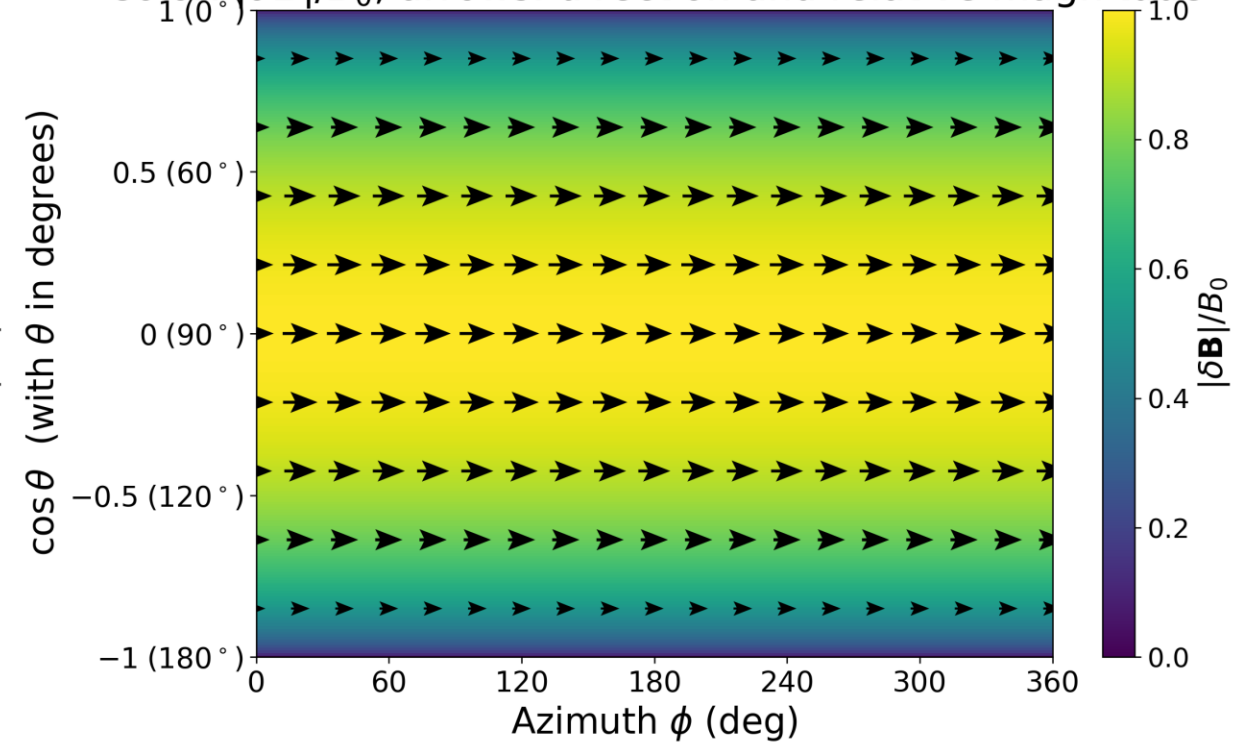
Plus polarization antenna pattern

Color: $|\delta \mathbf{B}|/B_0$, arrows: direction and relative magnitude



Cross polarization antenna pattern

Color: $|\delta \mathbf{B}|/B_0$, arrows: direction and relative magnitude

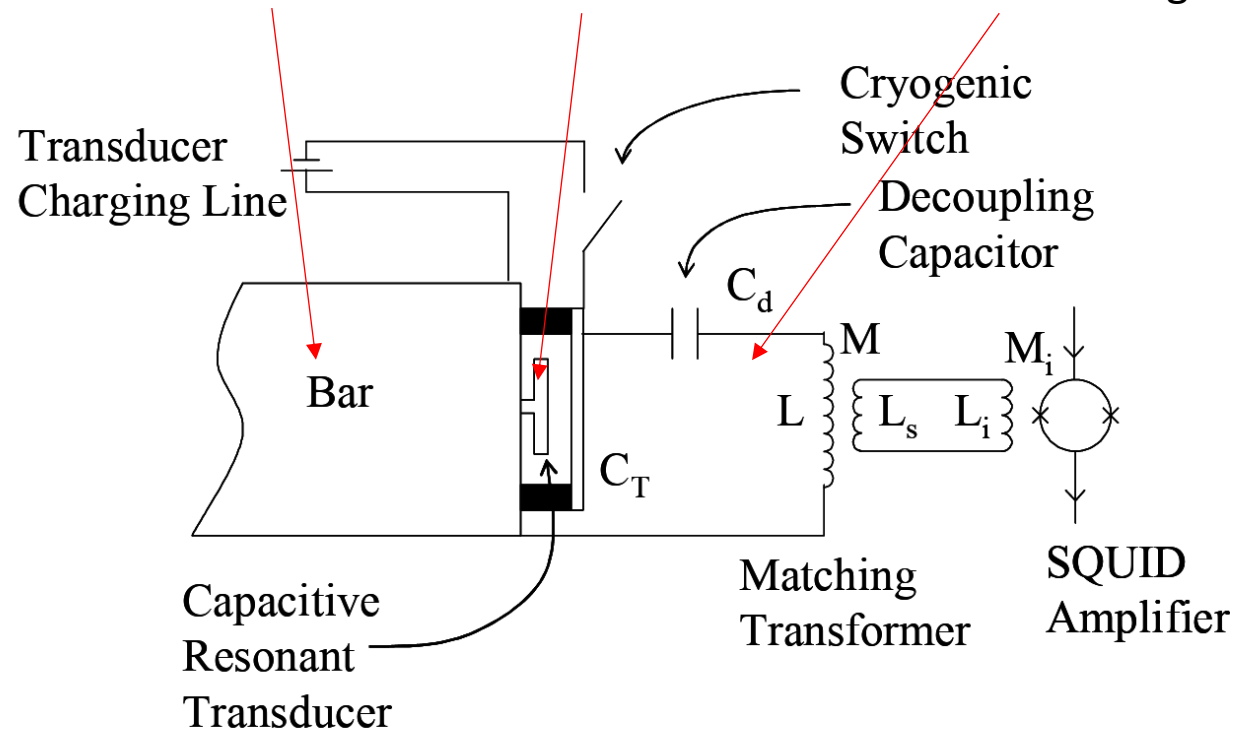


SQUID Readout

Motivation: Bar Detector with SQUID Readout

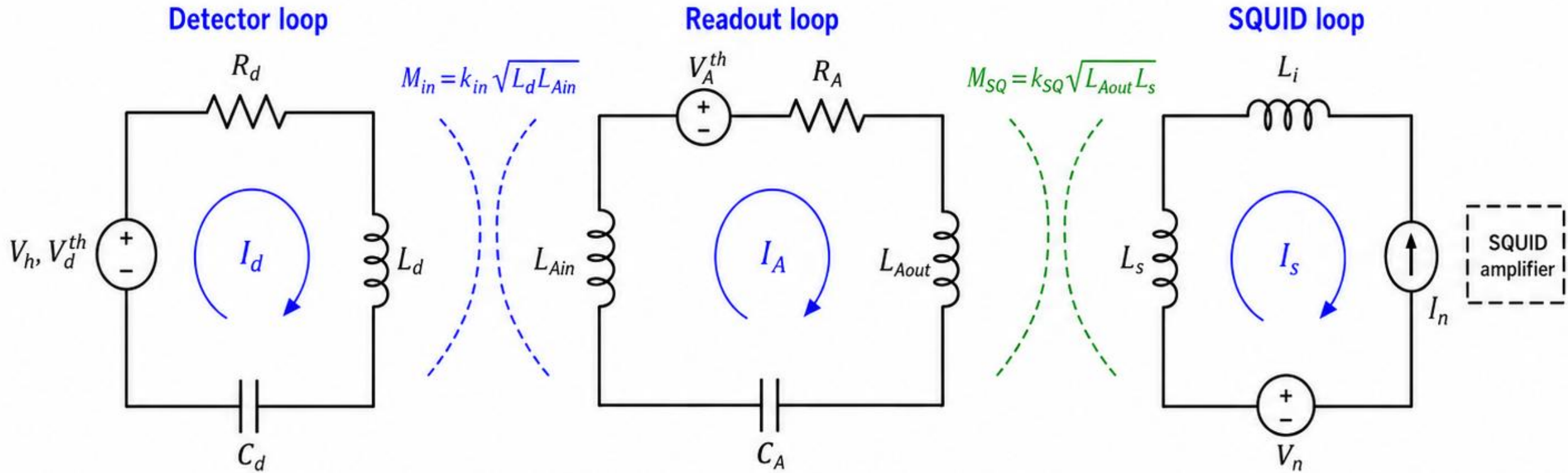
L. Baggio, M. Bignotto, M. Bonaldi, M. Cerdonio, L. Conti, **P. Falferi**, N. Liguori, A. Marin, R. Mezzena, **A. Ortolan**, S. Poggi, G. A. Prodi, F. Salemi, G. Soranzo, L. Taffarelli, G. Vedovato, A. Vinante, S. Vitale, and J. P. Zendri, “3-mode detection for widening the bandwidth of resonant gravitational wave detectors,” *Physical Review Letters* **94**, 241101 (2005).

Mechanical Resonant of Bar + Transducer Resonance + Electro Magnetic Mode



SQUID Readout

SQUID Readout for 100 Hz Bandwidth Detection



SQUID Readout

SQUID Readout for 100 Hz Bandwidth Detection

Signal source:

$$V_h = i\omega NABh$$

Detector impedance:

$$Z_d(\omega) = R_d + i\omega L_d + \frac{1}{i\omega C_d}$$

Readout impedance:

$$Z_A(\omega) = R_A + i\omega (L_{Ain} + L_{Aout}) + \frac{1}{i\omega C_A}$$

SQUID-loop impedance:

$$Z_s(\omega) = R_s + i\omega (L_s + L_i)$$

Measured SQUID flux:

$$\Phi_{meas} = L_i (I_s + I_n)$$

Loop-current vector:

$$I = (I_d, I_A, I_s)^T$$

Loop equations:

$$\begin{pmatrix} Z_d & i\omega M_{in} & 0 \\ i\omega M_{in} & Z_A & i\omega M_{SQ} \\ 0 & i\omega M_{SQ} & Z_s \end{pmatrix} \begin{pmatrix} I_d \\ I_A \\ I_s \end{pmatrix} = \begin{pmatrix} V_h + V_d^{th} \\ V_A^{th} \\ V_n \end{pmatrix}$$

Strain-noise model:

$$\mathcal{D}_3(\omega) = Z_s(\omega) [Z_d(\omega)Z_A(\omega) + \omega^2 M_{in}^2] + \omega^2 M_{SQ}^2 Z_d(\omega)$$

$$S_{hh} = S_d + S_A + S_{ba} + S_{imp}$$

$$S_d = \frac{4k_B T_d R_d}{\omega^2 (NAB)^2}$$

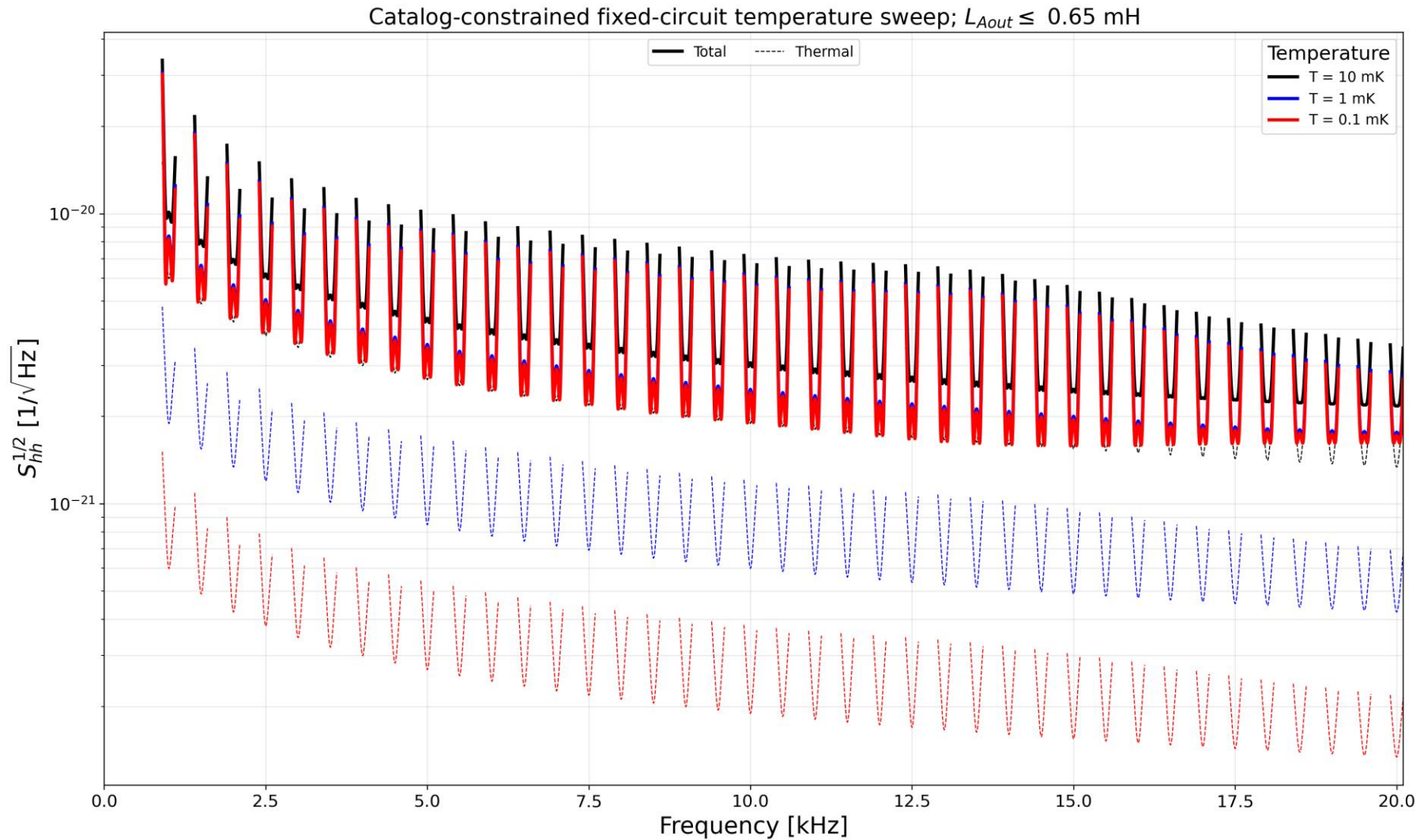
$$S_A = \frac{|Z_d|^2 4k_B T_A R_A}{\omega^4 (NAB)^2 M_{in}^2}$$

$$S_{ba} = \frac{|Z_d Z_A + \omega^2 M_{in}^2|^2 S_{V_n}}{\omega^6 (NAB)^2 M_{in}^2 M_{SQ}^2}$$

$$S_{imp} = \frac{|\mathcal{D}_3|^2 S_{I_n}}{\omega^6 (NAB)^2 M_{in}^2 M_{SQ}^2}$$

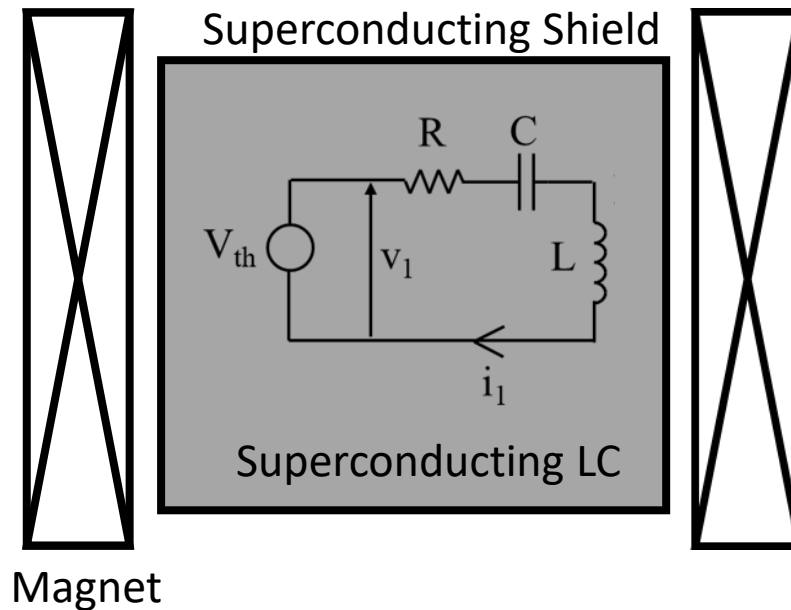
SQUID Readout

Strain Sensitivity



SQUID Readout

Superconducting LC Resonator



$Q = 500,000$ at 1.5 kHz

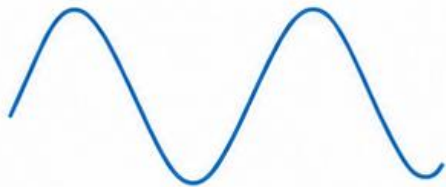
SQUID Readout

Piezoaxionic Effect

A. Arvanitaki, A. Madden, K. Van Tilburg, The Piezoaxionic Effect, arXiv:2112.11466

1. Oscillating axion dark matter

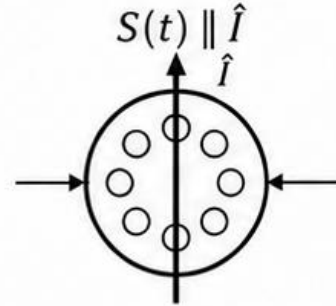
$$\theta_a(t) = \theta_0 \cos(\omega_a t)$$



$$\omega_a \simeq m_a$$

A spatially coherent QCD CP angle oscillates in time.

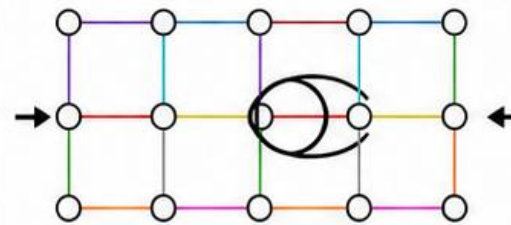
2. Schiff moment of a polarized nucleus



The axion-gluon coupling induces an oscillating P,T-odd nuclear charge distribution. The nuclear-spin polarization direction is fixed.

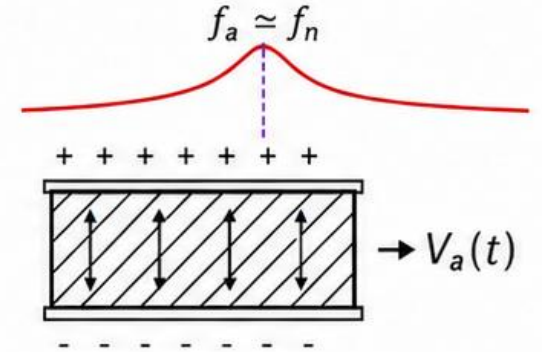
3. Alternating stress in a piezoelectric crystal

$$T_n^{(a)}(t) = -\xi_{nk} \hat{I}_k \theta_a(t)$$



In a non-centrosymmetric electronic structure, the equilibrium bonding shifts slightly. This produces alternating compressive and tensile stress throughout the crystal.

4. Acoustic resonance and piezoelectric voltage



Near resonance, the strain is amplified and converted to voltage by the piezoelectric effect.

Macroscopic constitutive equations of the crystal

$$T = c^D S - h^T D - \xi \hat{I} \theta_a$$

$$E = -h S + \beta^S D - \zeta \hat{I} \theta_a$$

Piezoaxionic channel: the ξ term drives stress $\rightarrow$ acoustic resonance $\rightarrow$ the piezoelectric h term generates an electrical signal.

Electroaxionic auxiliary channel: the ζ term creates a direct electric field, but it does not receive mechanical resonant enhancement.

Signal frequency set by the axion mass

$$f_{\text{sig}} = \frac{\omega_a}{2\pi} \simeq \frac{m_a c^2}{h}$$

$$f_{\text{sig}} \simeq 241.8 \text{ MHz} \times \left(\frac{m_a}{\mu\text{eV}} \right)$$

The mechanical resonance is tuned to this axion-induced signal frequency.

For $10^{-11} \text{ eV} \lesssim m_a \lesssim 10^{-7} \text{ eV}$, the corresponding signal frequency is about 2.4 kHz to 24 MHz.

Summary

Status and Prospects of QUAX-PD/LNL

- QUAX-LNL is searching for axion dark matter near 10 GHz, with the 2025–2026 data analysis currently in progress.
- Future upgrades include a wider-tuning, higher- Q cavity, a 14 T magnet, and an optimized TWPA to improve sensitivity and scan speed.
- The SMPD program aims to go beyond the standard quantum limit through single-microwave-photon counting.
- SQUID-based readout activities extend QUAX searches toward kHz-frequency axion and gravitational-wave signals.

Summary

Status and Prospects of QUAX-PD/LNL

QUAX Padova / Legnaro



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QUAX Frascati

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QUAX Trento

P. Falferi

QUAX Salerno

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Grazie per l'ascolto. 😊

Thank you for listening!