



The 2nd RISP Intensive Program on
“Rare Isotope Physics”

July 7, 2017

Nuclear Astrophysics Lecture #1

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Outline

- Lecture #1: Introduction of Nuclear Astrophysics
 - Historical background
 - Theoretical/Experimental considerations
 - Nuclear reactions in the Sun
 - Experiments using stable beams

- Lecture #2: Nuclear Astrophysics Experiments using Radioactive Ion Beams I
 - Why RIB?
 - Hydrogen burnings, HCNO, rp-process, r-process
 - Experiments at RIKEN & CNS
 - Prospect for the Korea RI Accelerator

Nuclear Astrophysics

- Addresses some of the most compelling questions in nature
 - The origins of the elements
 - How does the sun shine for so many years?
 - What is the total density of matter in the universe?
 - How did the stars, galaxies evolve?
- Require a considerable amount of nuclear physics information as input
 - Mostly based on theoretical models or extrapolations
- More complete and precise nuclear physics measurements are needed

Special: New Learning Series on Genetics, page 70

Complexity—the Science of Surprise | Your Inner Savant

Discover

FEBRUARY 2002

DISCOVER.COM

The
11
Greatest
Unanswered
Questions
of Physics

No.
9
What Is Gravity?

Question 3

How were the elements from iron to uranium made ?

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EAST LANS
Libraries

Based on National Academy of Science Report

[Committee for the Physics of the Universe (CPU)]



~90 Years of Nuclear Astrophysics

A Story of Success

1928 Gamow-Factor

George Gamow

1931 Stellar Structure & Theory of White Dwarfs

Subramanyan Chandrasekhar

1938 CNO Cycle - C. F. von Weizsacker

CNO Cycle, pp Chain - Hans Bethe

1957 Nucleosynthesis of Elements in Stars

The $0+ 2^{\text{nd}}$ Excited State in ^{12}C

Measured at 7.65 MeV in 1957

Margaret & Geoffrey Burbidge, William

Fowler and Fred Hoyle = B^2FH

1967 Nobel Prize: Hans Bethe

Reactions in the Sun: pp chain

Solar Neutrino Problem: $^7\text{Be}(\text{p},\text{g})$

1983 Nobel Prize: William Fowler

Subramanyan Chandrasekhar

2002 Nobel Prize for Neutrino Detection

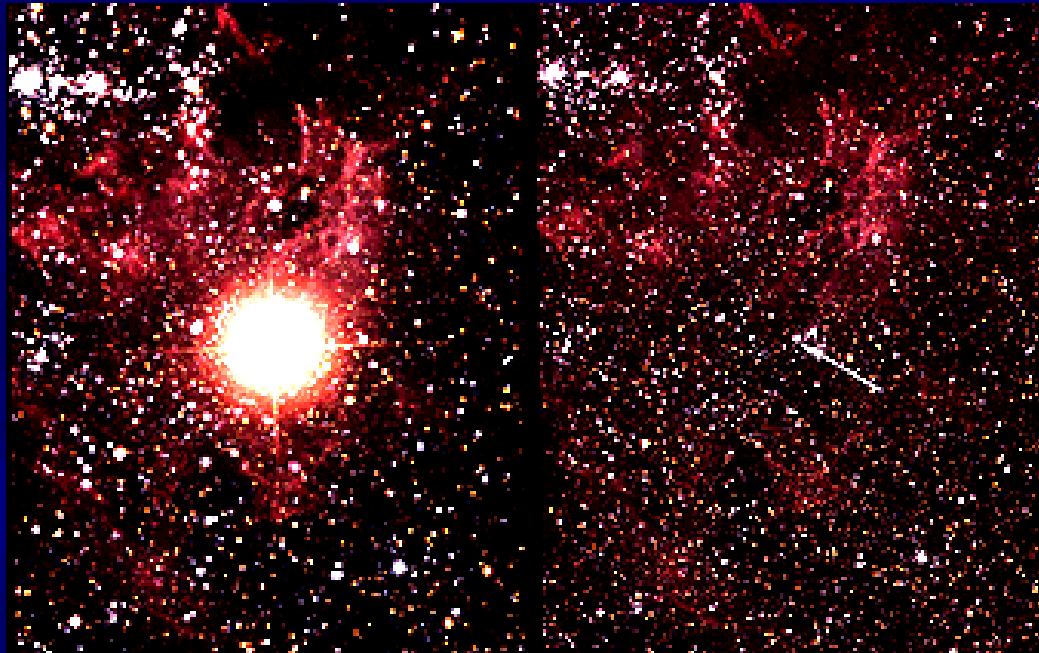
Raymond Davis & Masatoshi Koshiba

X-ray Astronomy - Riccardo Giacconi

2015 Nobel Prize for Neutrino oscillation

Takaaki Kajita & Arthur McDonald

Experimental Nuclear Astrophysics



Nuclear reactions in stars

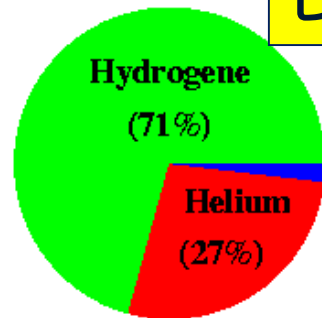
- ➡ **produce energy**
- ➡ **generate the elements**

Experimental Nuclear Astrophysics

- We try to observe nuclear reaction processes from
 - Heat from stars
 - probes only surface
 - Abundances of elements
 - Neutrino's from stars
 - probes interior of star
 - Lab studies of reaction cross-sections
 - Experimental nuclear physics

태양계의 원소비율

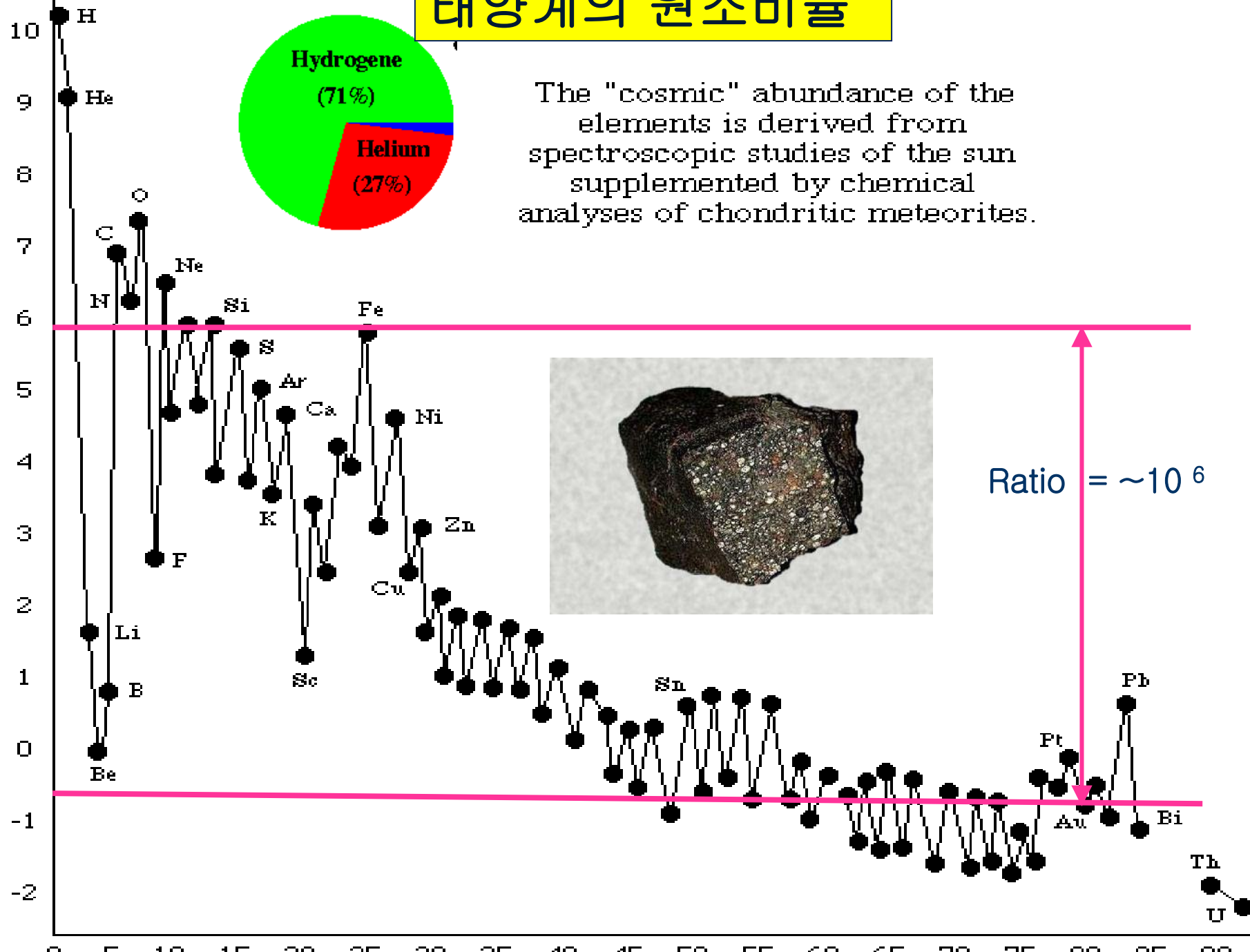
RELATIVE
ABUNDANCE
LOG
SCALE



The "cosmic" abundance of the elements is derived from spectroscopic studies of the sun supplemented by chemical analyses of chondritic meteorites.



Ratio = $\sim 10^6$



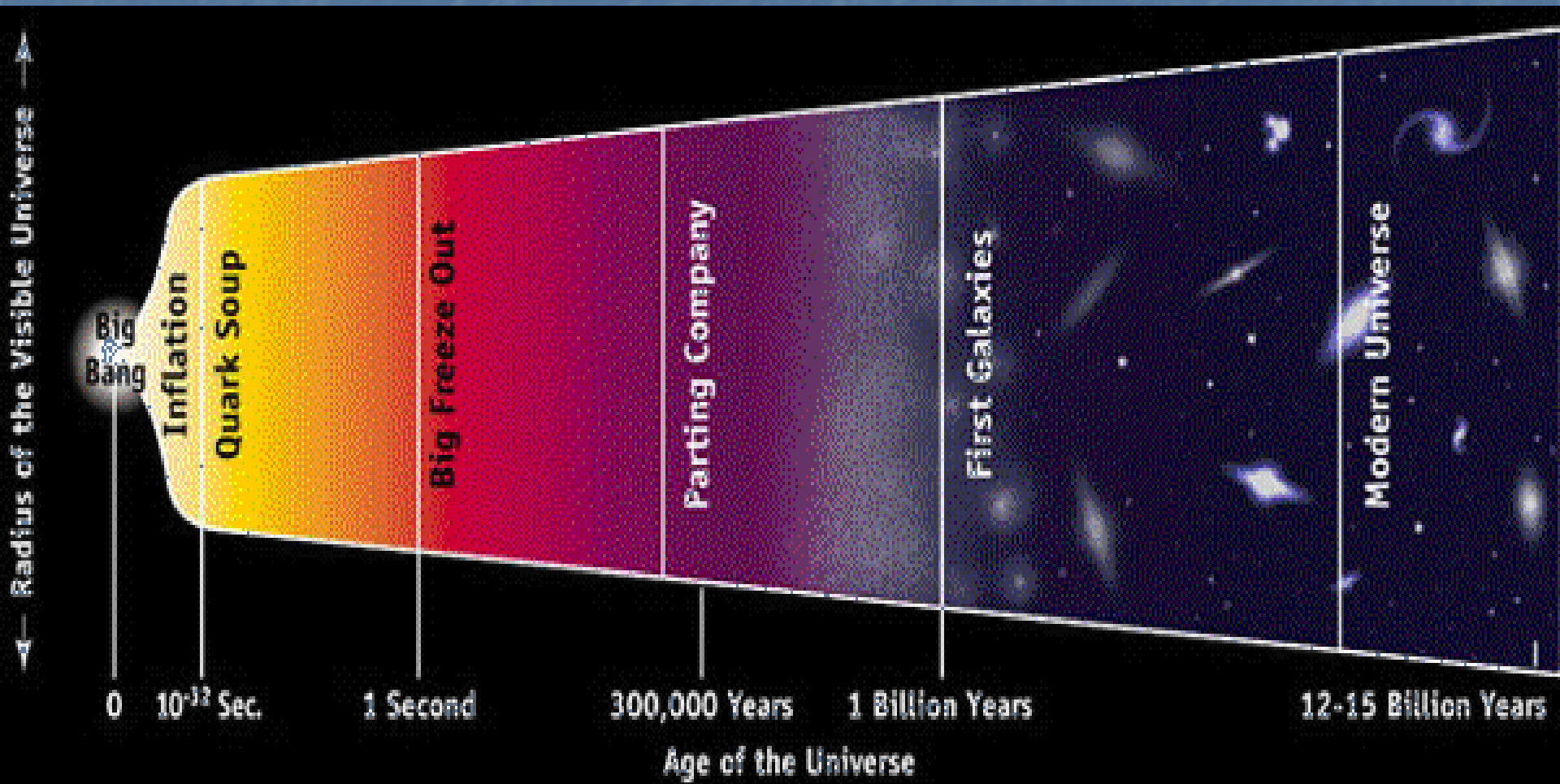
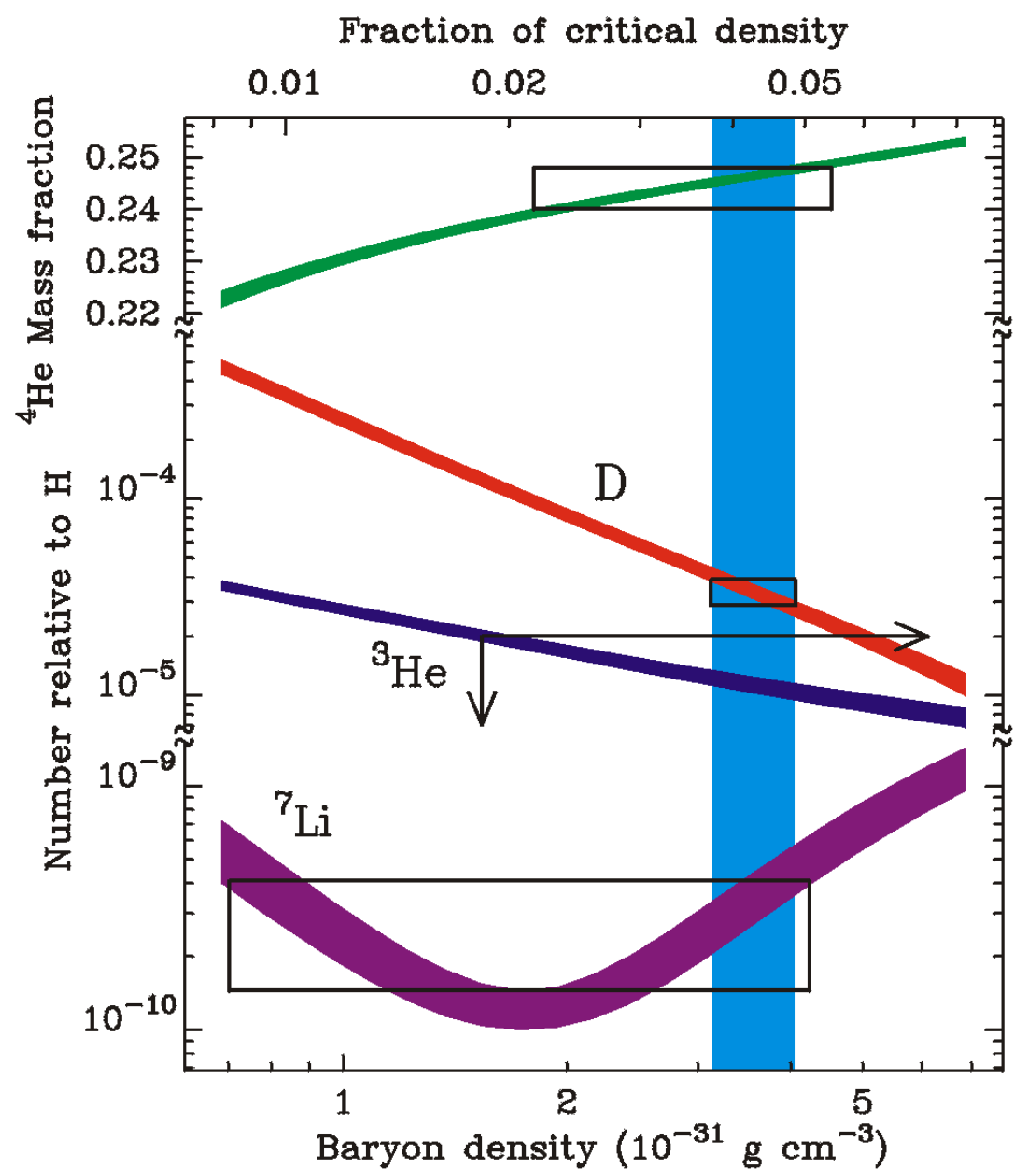


Figure taken from APS poster by Burles, Nollett & Turner



Question : What is the ratio of neutron over proton, n/p?

Nucleosynthesis in Cosmos

293

2,771

3,064

NNDC (BNL, 2000)

Stable
Observed Unstable

s process

rp process

Stellar evolution

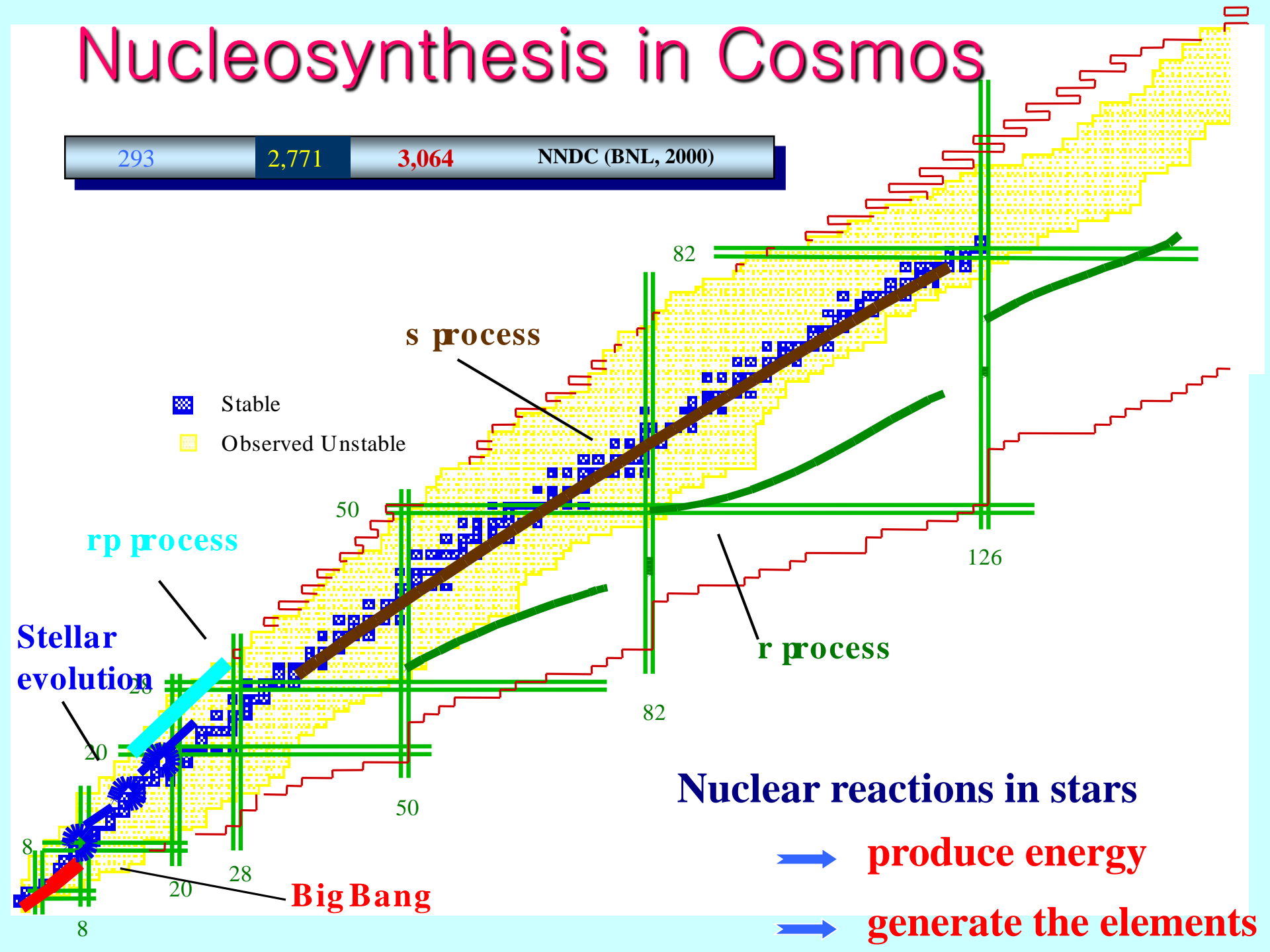
r process

Nuclear reactions in stars

→ produce energy

→ generate the elements

Big Bang

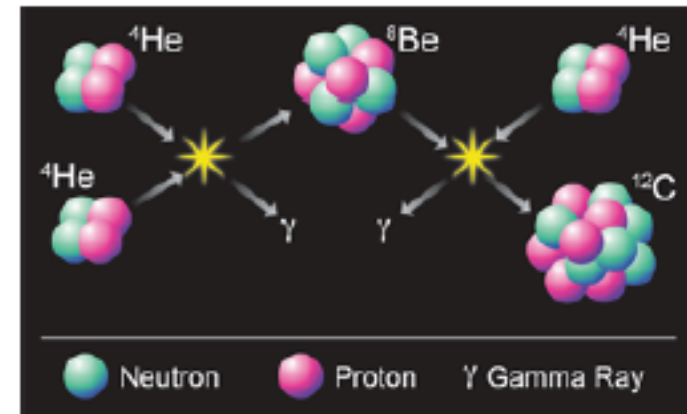
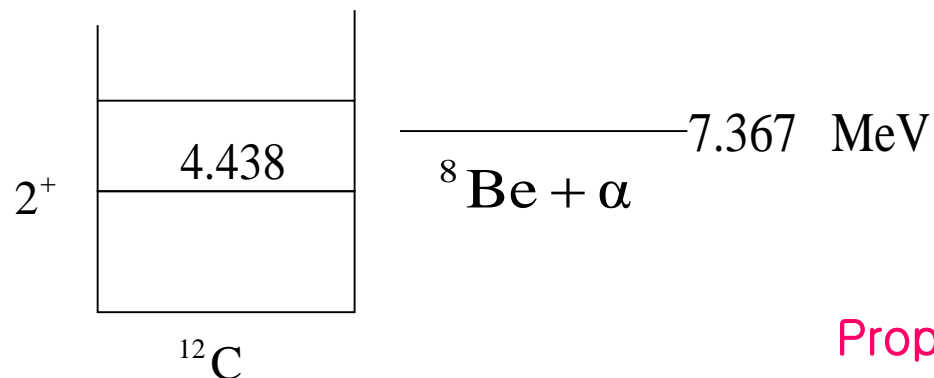
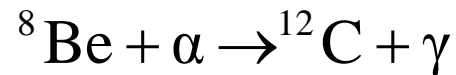


H, ^4He , ^3He , Li - Big Bang



< Hoyle in 1953 >

$3\alpha \rightarrow ^{12}\text{C}$ Is insufficient to explain the observed abundance



Proposed O^+ at 7.68 MeV in 1953

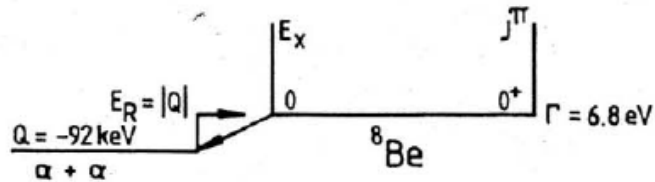
Measured at 7.65 MeV in 1957

→ removed the major roadblock for the theory that elements are made in stars

→ Nobel Prize in Physics 1983 for Willy Fowler

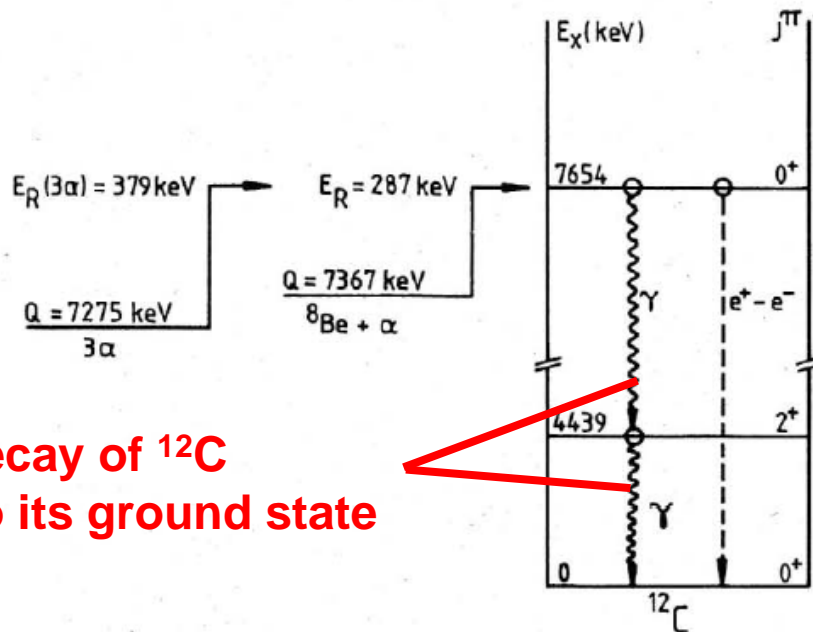
Third step completes the reaction:

FIRST STEP: $\alpha + \alpha \rightleftharpoons {}^8\text{Be}$

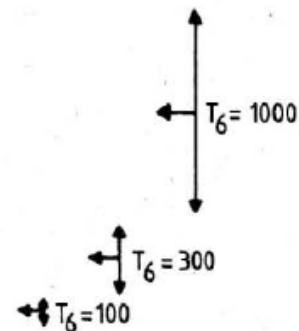


Note: ${}^8\text{Be}$ ground state is a 92 keV resonance for the $\alpha + \alpha$ reaction

SECOND STEP: ${}^8\text{Be} (\alpha, \gamma) {}^{12}\text{C}$



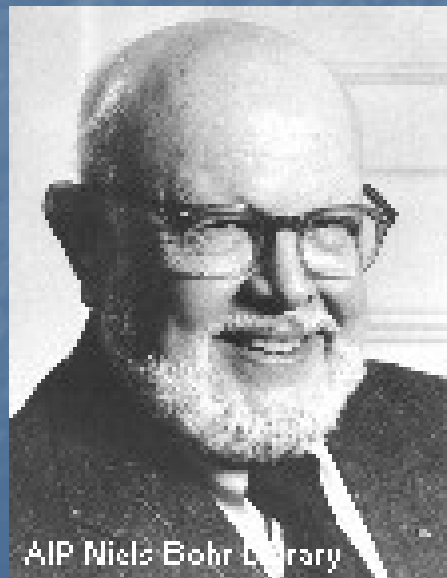
**γ decay of ${}^{12}\text{C}$
into its ground state**



Note:
 $\Gamma_\alpha / \Gamma_\gamma > 10^3$
so γ -decay is very rare !

“ It is a remarkable fact that humans, on the basis of experiments and measurements carried out in the lab, are able to understand the universe in the early stages of its evolution, even during the first three minutes of its existence.”

Fowler (Nobel prize 1983)



AIP Niels Bohr Library

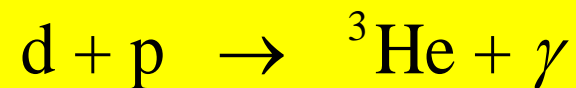
Nucleosynthetic reactions are typically dominated by Coulomb barriers

$$E_B = \frac{Z_1 Z_2 e^2}{R} = \frac{1.44 Z_1 Z_2}{R(fm)} \text{ MeV}$$

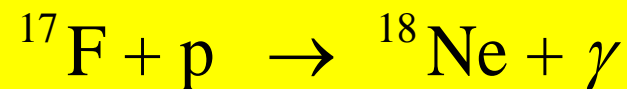
$$E_T \approx kT = 8.62 \times 10^{-8} T \text{ keV}$$

$$T \approx 10^8 \text{ K}$$

$$E_T \approx 10 \text{ keV}$$



$$E_B \approx 400 \text{ keV}$$



$$E_B \approx 2.52 \text{ MeV}$$



$$E_B \approx 4.00 \text{ MeV}$$

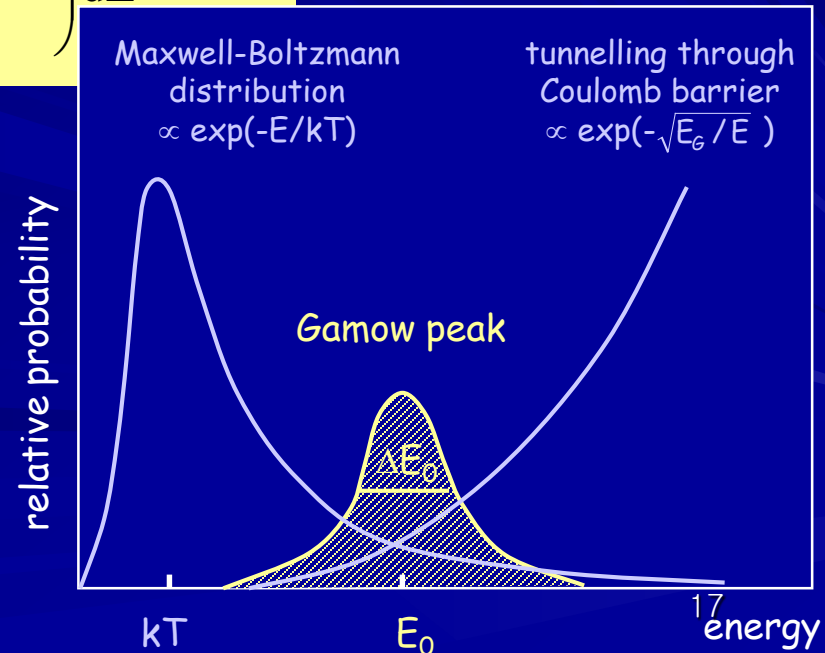
Thermonuclear reactions in stars

$$S(E) \equiv \sigma(E)E \exp\left(\frac{2\pi Z_1 Z_2 e^2}{\hbar v}\right)$$

$$\lambda = \langle \sigma v \rangle = \int_0^\infty \sigma(E) v(E) \Psi(E) dE$$

$$= \int_0^\infty \frac{S(E)}{E} \exp(-bE^{-1/2}) \sqrt{\frac{2E}{\mu}} \frac{2}{\sqrt{\pi}} \frac{E}{kT} \exp\left(-\frac{E}{kT}\right) \frac{dE}{(kTE)^{1/2}}$$

$$= \left(\frac{8}{\mu\pi}\right)^{1/2} \frac{1}{(kT)^{3/2}} \int_0^\infty \frac{S(E)}{E} \exp\left(-\frac{E}{kT} - bE^{-1/2}\right) dE$$



With above definition of cross section:

$$\langle \sigma v \rangle_{12} = \left(\frac{8}{\pi \mu_{12}} \right)^{1/2} \frac{1}{(kT)^{3/2}} \int_0^{\infty} S(E) \exp \left[-\frac{E}{kT} - \frac{b}{E^{1/2}} \right] dE$$

\downarrow varies smoothly with energy $\underbrace{\exp \left[-\frac{E}{kT} - \frac{b}{E^{1/2}} \right]}_{f(E)}$ governs energy dependence

MAXIMUM reaction rate:

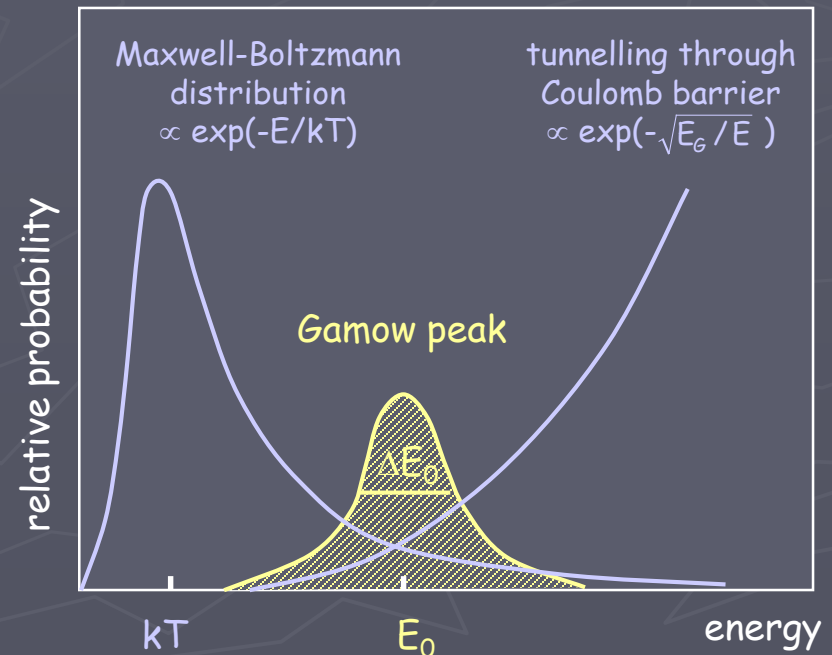
$$\frac{df(E)}{dE} = 0 \quad \Rightarrow \quad E_0 = \left(\frac{bkT}{2} \right)^{2/3}$$

$$\Delta E_0 < E_0$$

only small energy range contributes to reaction rate



OK to set $S(E) \sim S(E_0) = \text{const.}$



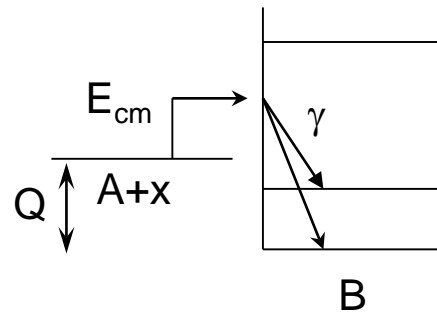
I. direct process

one-step process

direct transition into a bound state

example:

radiative capture $A(x,\gamma)B$



$$\sigma_{\gamma} \propto \left| \langle B | H_{\gamma} | A + x \rangle \right|^2$$

H_{γ} = electromagnetic operator describing the transition

- reaction cross section proportional to single matrix element
- can occur at all projectile energies
- smooth energy dependence of cross section

other direct processes: stripping, pickup, charge exchange, Coulomb excitation

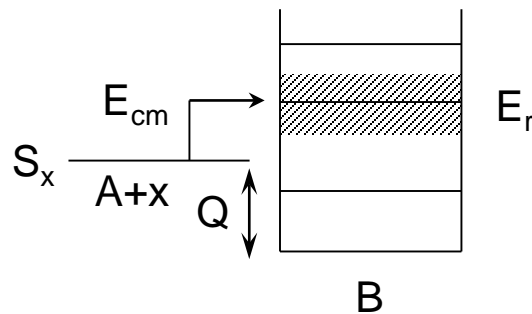
II. resonant process

two-step process

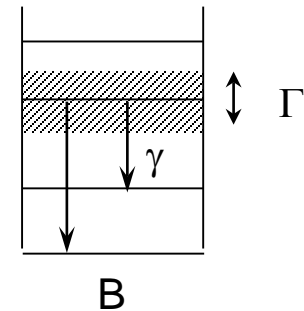
example:

resonant radiative capture $A(x,\gamma)B$

1. Compound nucleus formation
(in an unbound state)



2. Compound nucleus decay
(to lower excited states)



$$\sigma_{\gamma} \propto \underbrace{\left| \langle E_f | H_{\gamma} | E_r \rangle \right|^2}_{\text{compound decay probability} \propto \Gamma_{\gamma}} \underbrace{\left| \langle E_r | H_B | A + x \rangle \right|^2}_{\text{compound formation probability} \propto \Gamma_x}$$

- reaction cross section proportional to two matrix elements
- only occurs at energies $E_{cm} \sim E_r - Q$
- strong energy dependence of cross section

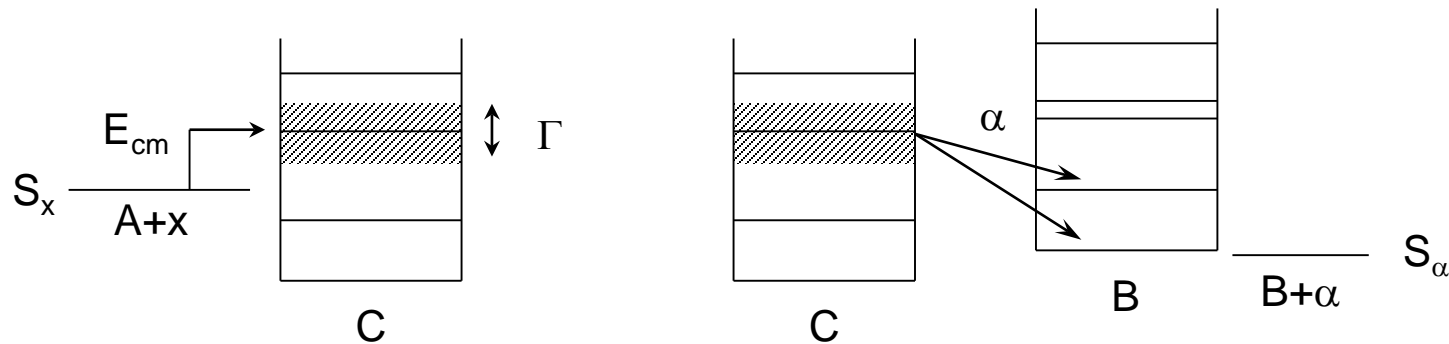
N. B. energy in entrance channel ($Q + E_{cm}$) has to match excitation energy E_r of resonant state, however all excited states have a width \Rightarrow there is always some cross section through tails

example:

resonant reaction $A(x,\alpha)B$

1. Compound nucleus formation
(in an unbound state)

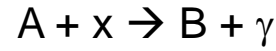
2. Compound nucleus decay
(by particle emission)



$$\sigma_{\gamma} \propto \underbrace{\left| \langle B + \alpha | H_{\alpha} | E_r \rangle \right|^2}_{\text{compound decay probability} \propto \Gamma_{\alpha}} \underbrace{\left| \langle E_r | H_x | A + x \rangle \right|^2}_{\text{compound formation probability} \propto \Gamma_x}$$

N. B. energy in entrance channel ($S_x + E_{cm}$) has to match excitation energy E_r of resonant state, however all excited states have a width \Rightarrow there is always some cross section through tails

example: direct capture



$$\sigma = \pi \lambda_x^2 \left| \langle B | H | x + A \rangle \right|^2 P_\ell(E)$$

“geometrical factor”
de Broglie wavelength
of projectile

$$\lambda = \frac{h}{p} = \frac{h}{\sqrt{2mE}}$$

matrix element
contains nuclear
properties of interaction

penetrability/transmission
probability for projectile to
reach target for interaction
depends on projectile’s angular
momentum ℓ and energy E

$$\sigma = \frac{1}{E} \cdot P_\ell(E) \cdot S(E)$$

$$\sigma = (\text{strong energy dependence}) \times (\text{weak energy dependence})$$

S(E) = astrophysical factor contains nuclear physics of reaction
+ can be easily: graphed, fitted, extrapolated (if needed)

$$\text{need expression for } P_\ell(E)$$

factors affecting transmission probability:

- Coulomb barrier (for charged particles only)
- centrifugal barrier (both for neutrons and charged particles)

for a single isolated resonance:

resonant cross section given by Breit-Wigner expression

$$\sigma(E) = \pi \hat{\lambda}^2 \frac{2J+1}{(2J_1+1)(2J_T+1)} \frac{\Gamma_1 \Gamma_2}{(E - E_r)^2 + (\Gamma/2)^2}$$

for reaction: $1 + T \rightarrow C \rightarrow F + 2$

geometrical factor

$\propto 1/E$

spin factor ω

J = spin of CN's state

J_1 = spin of projectile

J_T = spin of target

strongly energy-dependent term

Γ_1 = partial width for decay via emission of particle 1
= probability of compound formation via entrance channel

Γ_2 = partial width for decay via emission of particle 2
= probability of compound decay via exit channel

Γ = total width of compound's excited state
 $= \Gamma_1 + \Gamma_2 + \Gamma_\gamma + \dots$

E_r = resonance energy

what about penetrability considerations? \Rightarrow look for energy dependence in partial widths!

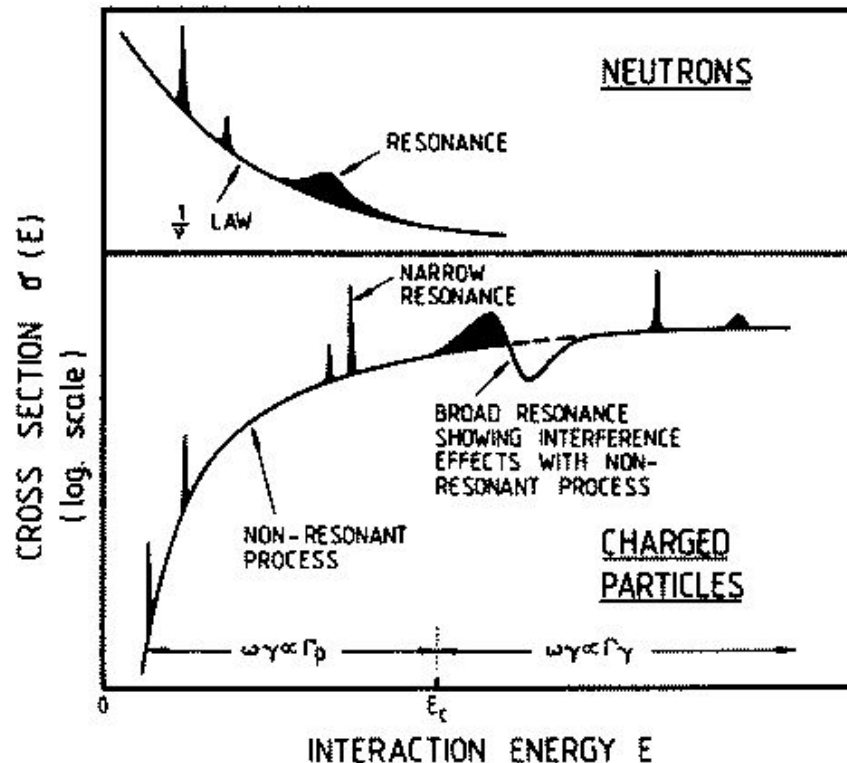
partial widths are NOT constant but energy dependent!

stellar reaction rate of nuclear reaction determined by the sum of contributions due to

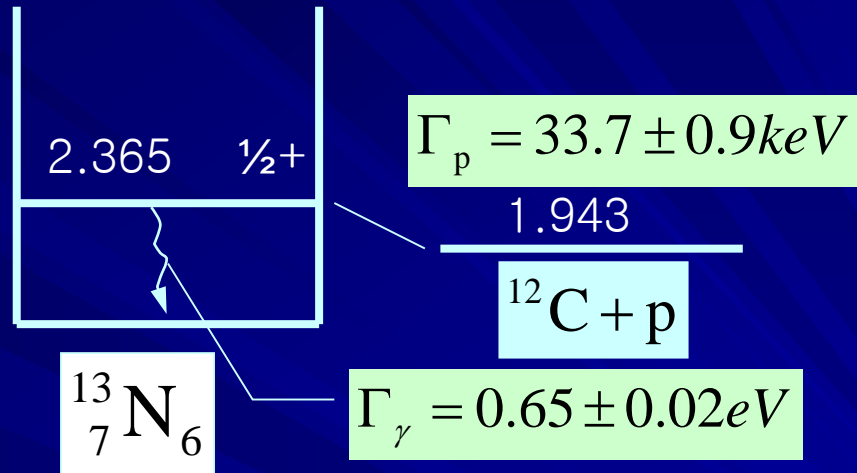
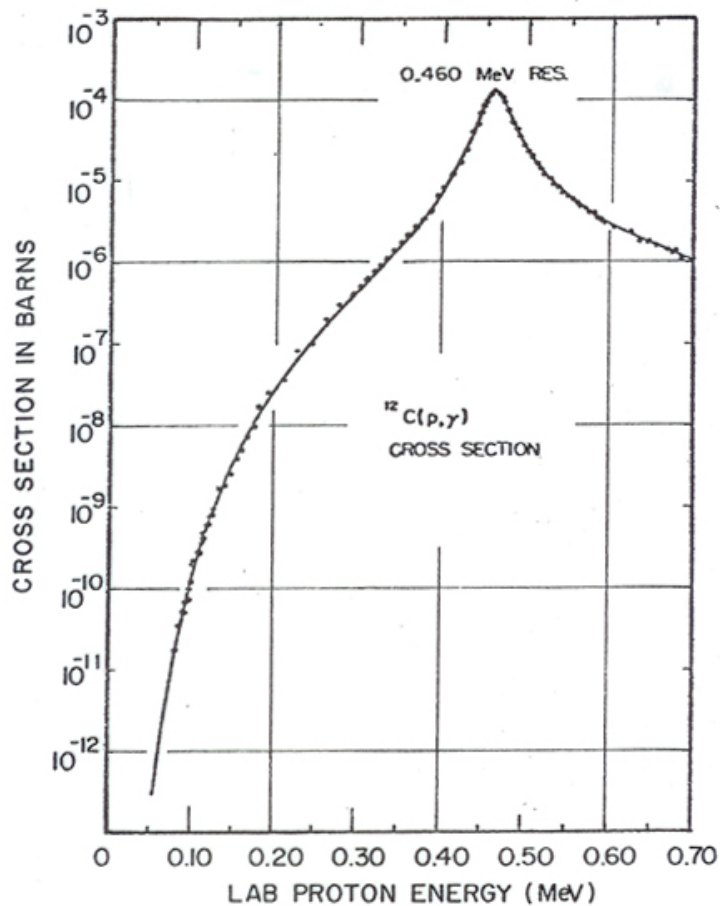
- direct transitions to the various bound states
- all narrow resonances in the relevant energy window
- broad resonances (tails) e.g. from higher lying resonances
- any interference term

total rate

$$\langle \sigma v \rangle = \sum_i \langle \sigma v \rangle_{\text{DCi}} + \sum_i \langle \sigma v \rangle_{\text{Ri}} + \langle \sigma v \rangle_{\text{tails}} + \langle \sigma v \rangle_{\text{interference}}$$



Rolfs & Rodney
Cauldrons in the Cosmos, 1988

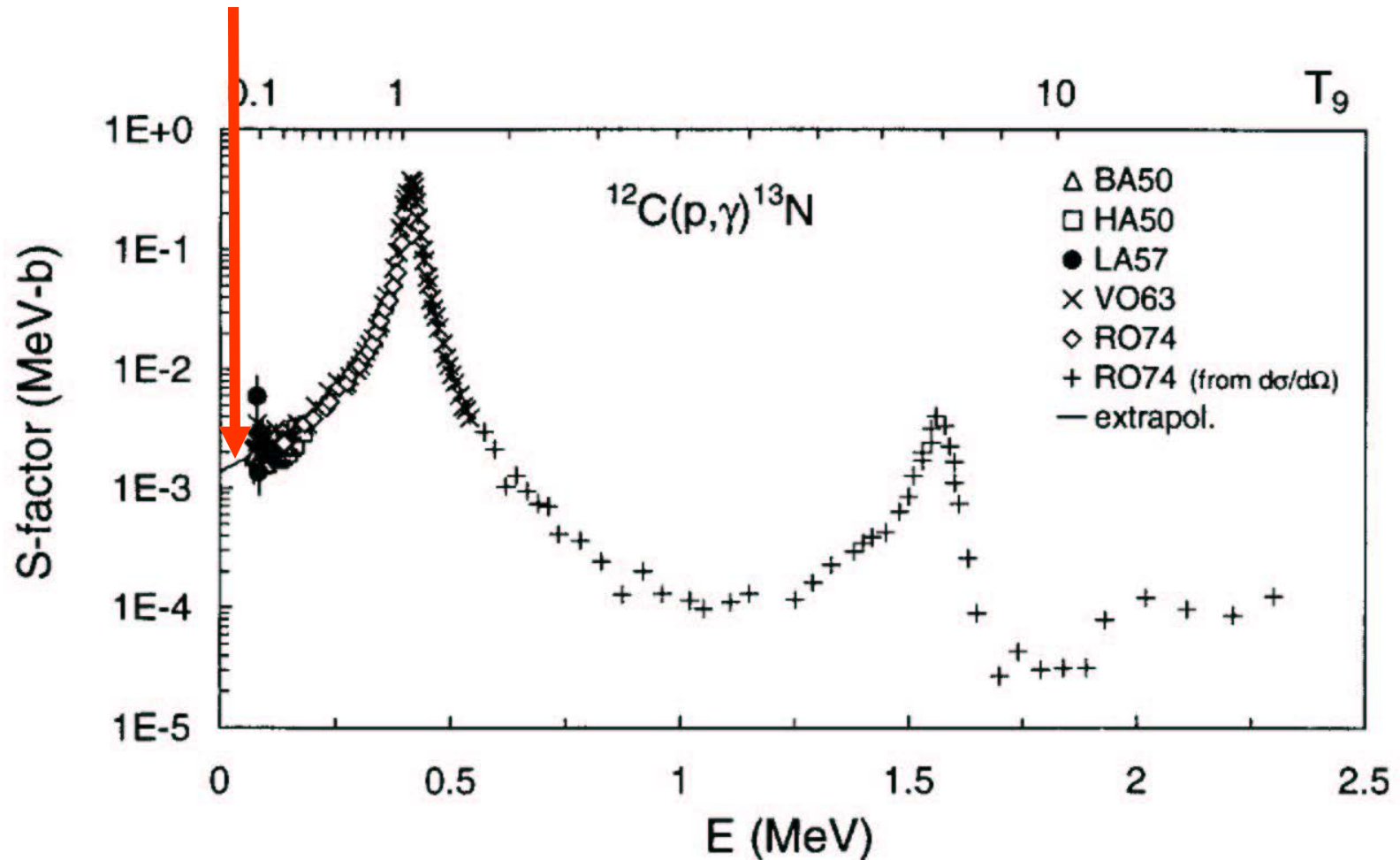


Breit-Wigner

$$\sigma \propto \frac{(2J+1)}{(2j_p+1)(2j_T+1)} \frac{\Gamma_p \Gamma_\gamma}{(E-E_\gamma)^2 + (\Gamma/2)^2}$$

S-Factor:

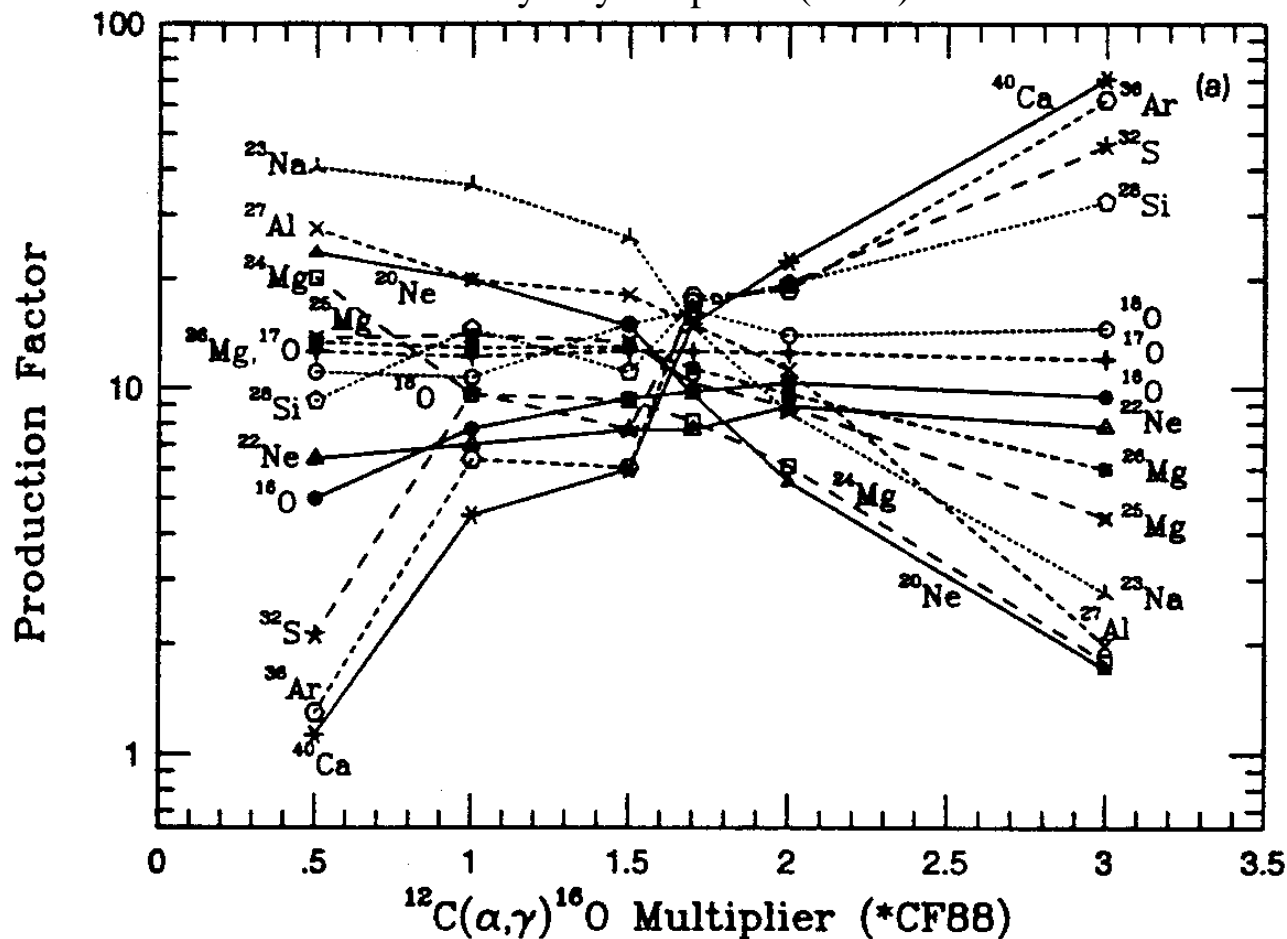
Need rate
about here



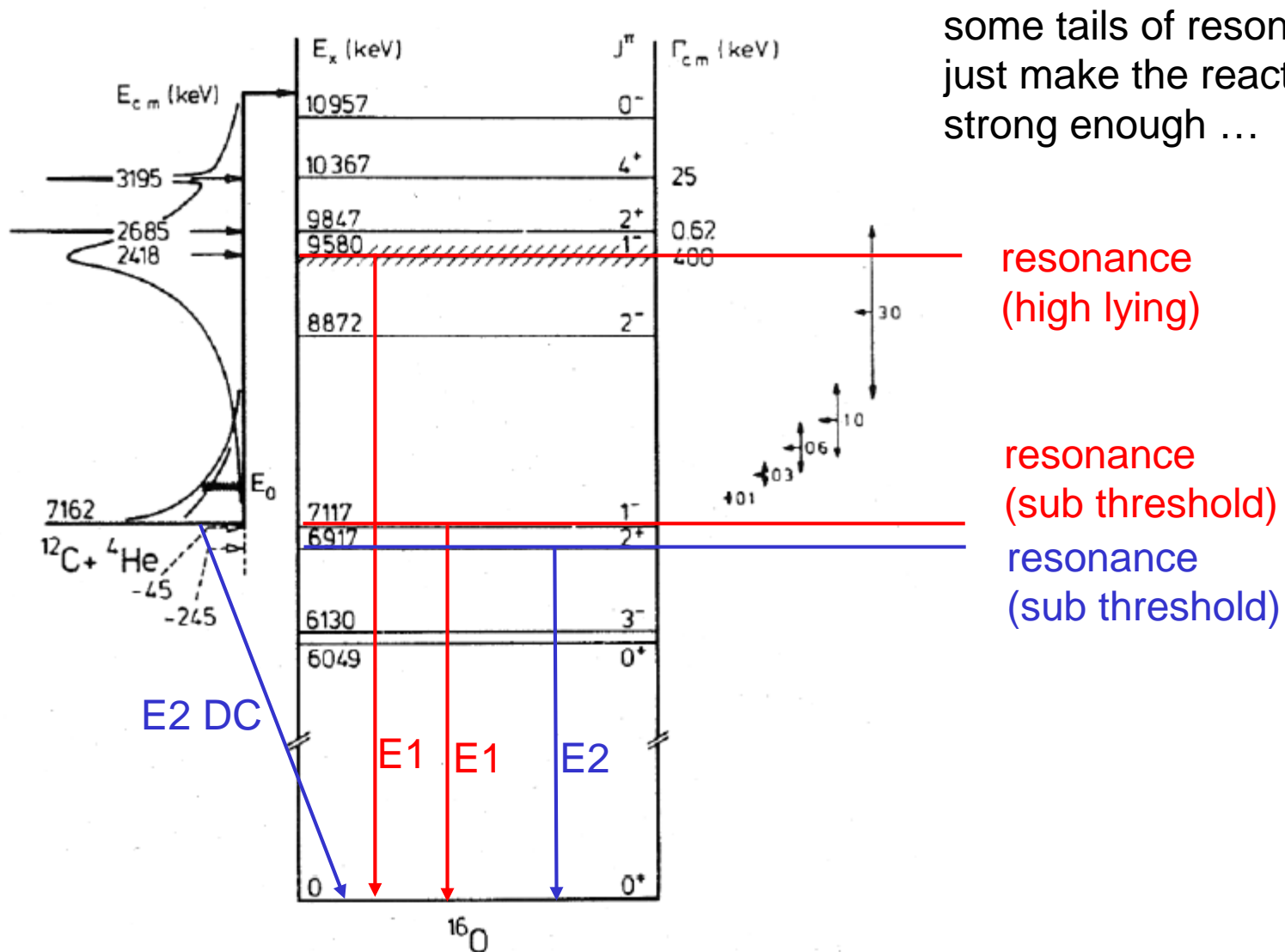
From the **NACRE compilation** of charged particle induced reaction rates on stable nuclei from H to Si (Angulo et al. Nucl. Phys. A 656 (1999) 3)

Massive star nucleosynthesis model as a function of $^{12}\text{C}(\alpha,\gamma)$ rate

Weaver and Woosley Phys Rep 227 (1993) 65



- This demonstrates the sensitivity
- One could deduce a preference for a total S(300) of ~120-220 (But of course we cannot be sure that the astrophysical model is right)



- complications:
- very low cross section makes direct measurement impossible
 - subthreshold resonances cannot be measured at resonance energy
 - Interference between the E1 and the E2 components

Therefore:

Uncertainty in the $^{12}\text{C}(\alpha,\gamma)$ rate is the single most important nuclear physics uncertainty in astrophysics

- Affects:
- C/O ratio \rightarrow further stellar evolution (C-burning or O-burning ?)
 - iron (and other) core sizes (outcome of SN explosion)
 - Nucleosynthesis (see next slide)

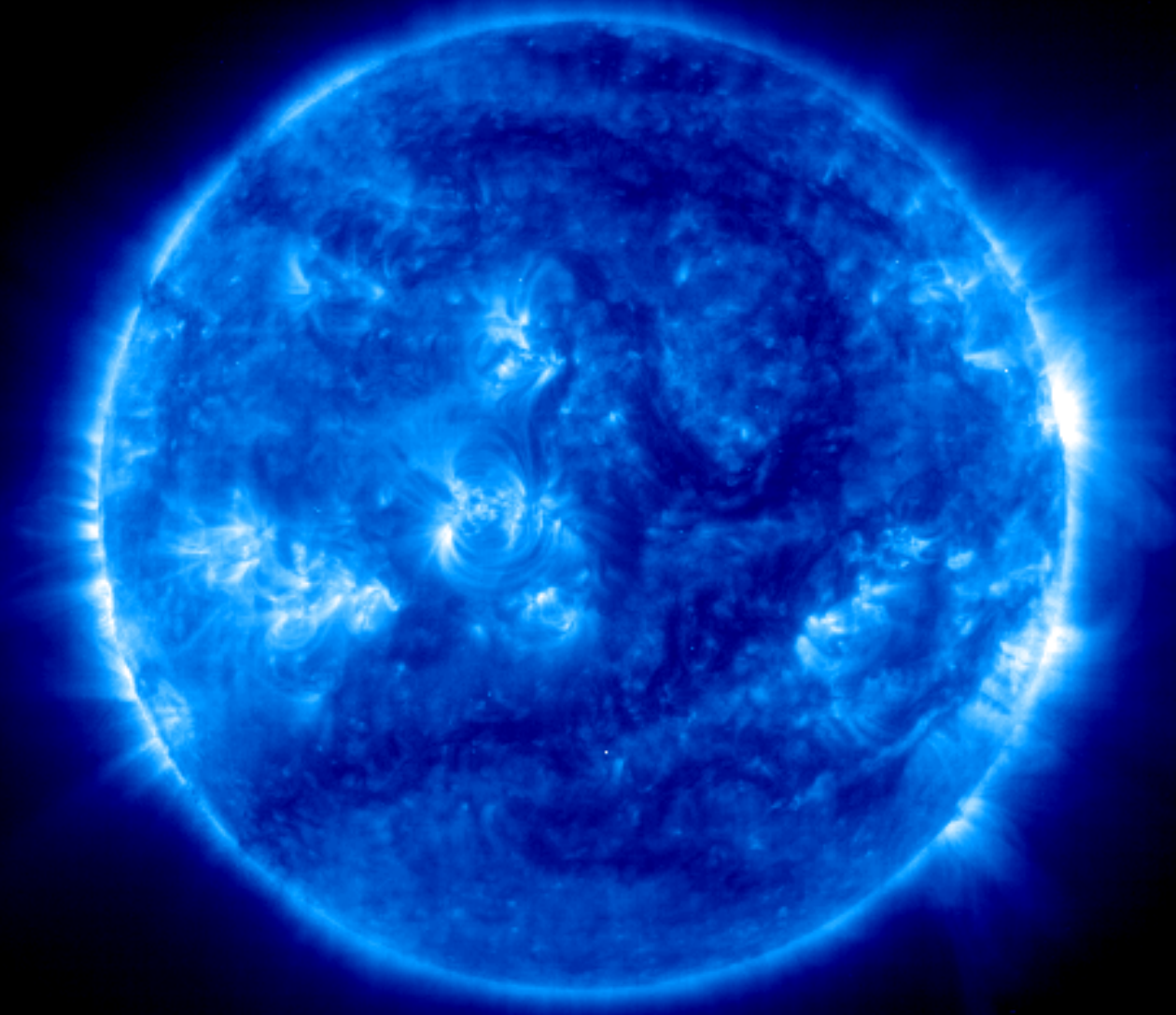
Some current results for S(300 keV):

$S_{E2} = 53 \pm 13$ -18 keV b (Tischhauser et al. PRL88(2002)2501)

$S_{E1} = 79 \pm 21$ -21 keV b (Azuma et al. PRC50 (1994) 1194)

But others range among groups larger !

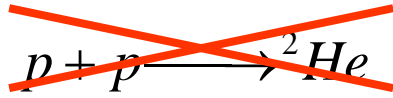
Nuclear Reactions in the Sun



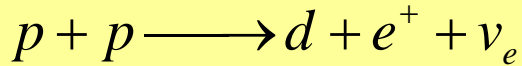
SOHO, 171A Fe emission line

The ppl chain

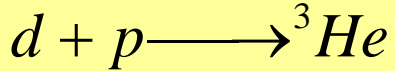
Step 1:



${}^2\text{He}$ is unstable



Step 2:



d abundance is too low

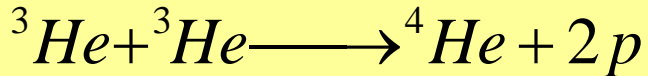
Step 3:



${}^4\text{Li}$ is unstable



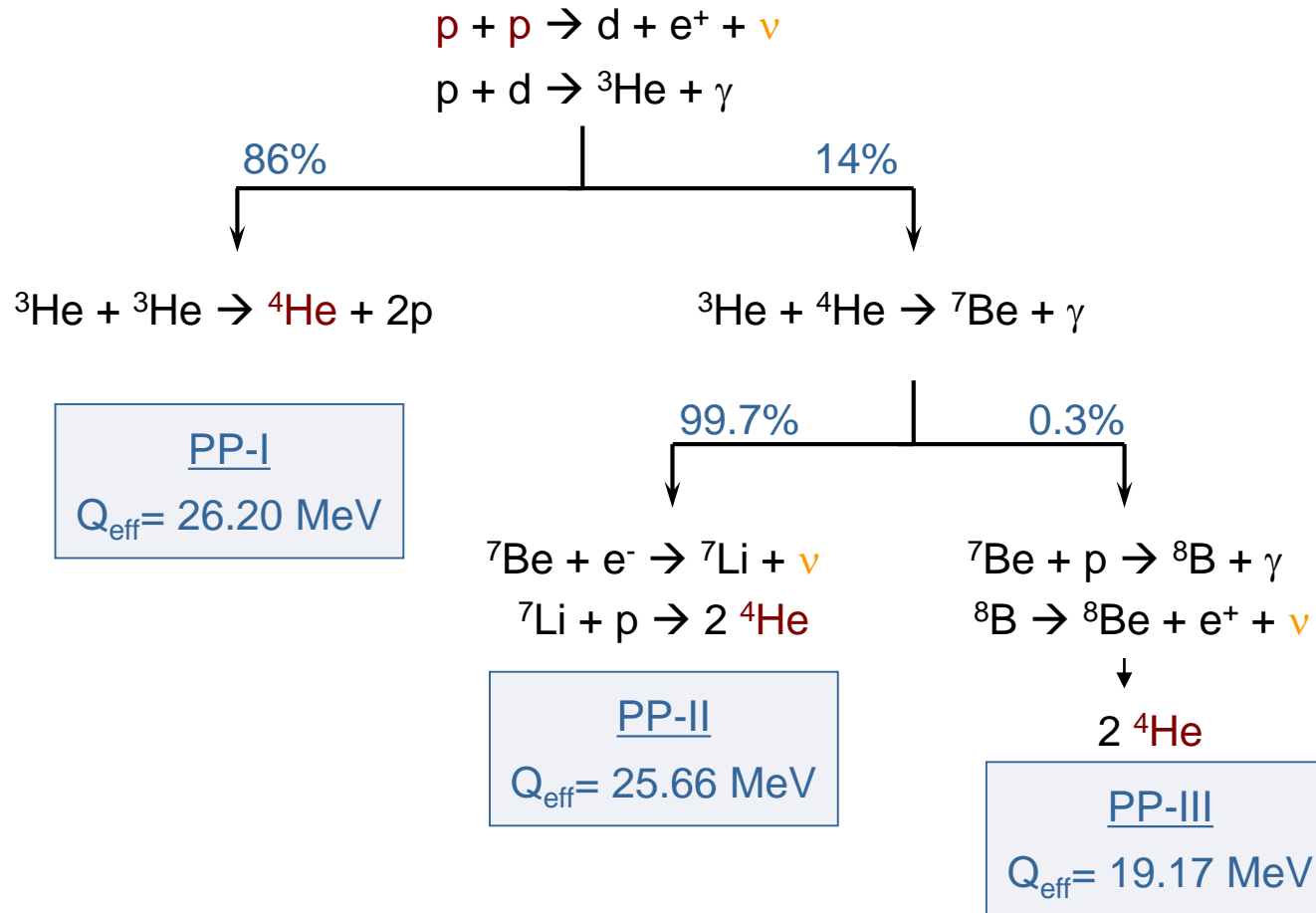
d abundance is too low



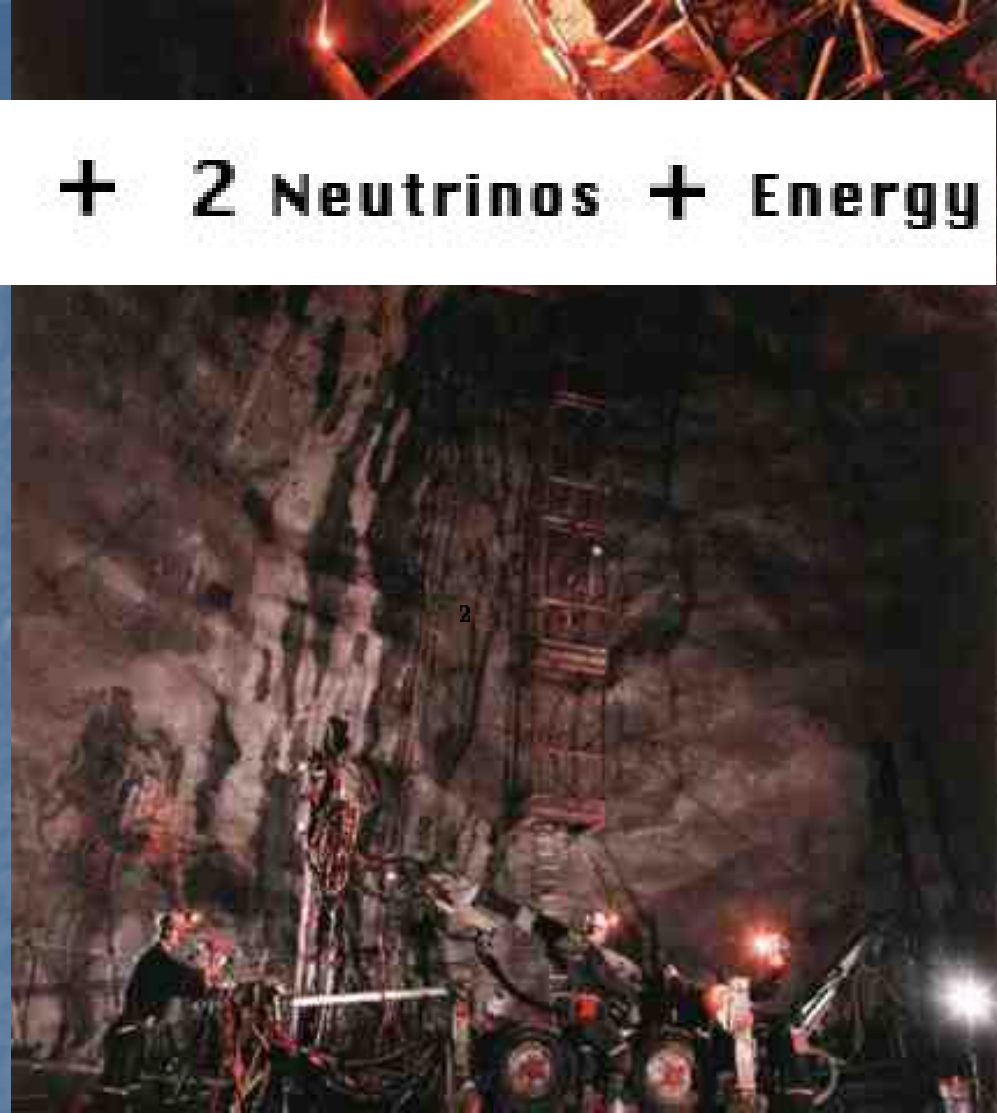
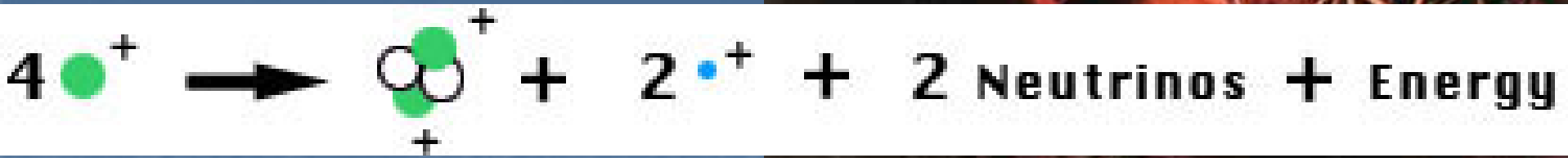
d+d not going because Y_d is small as d+p leads to rapid destruction

${}^3\text{He} + {}^3\text{He}$ goes because $Y_{{}^3\text{He}}$ gets large as there is no other rapid destruction

proton-proton chain

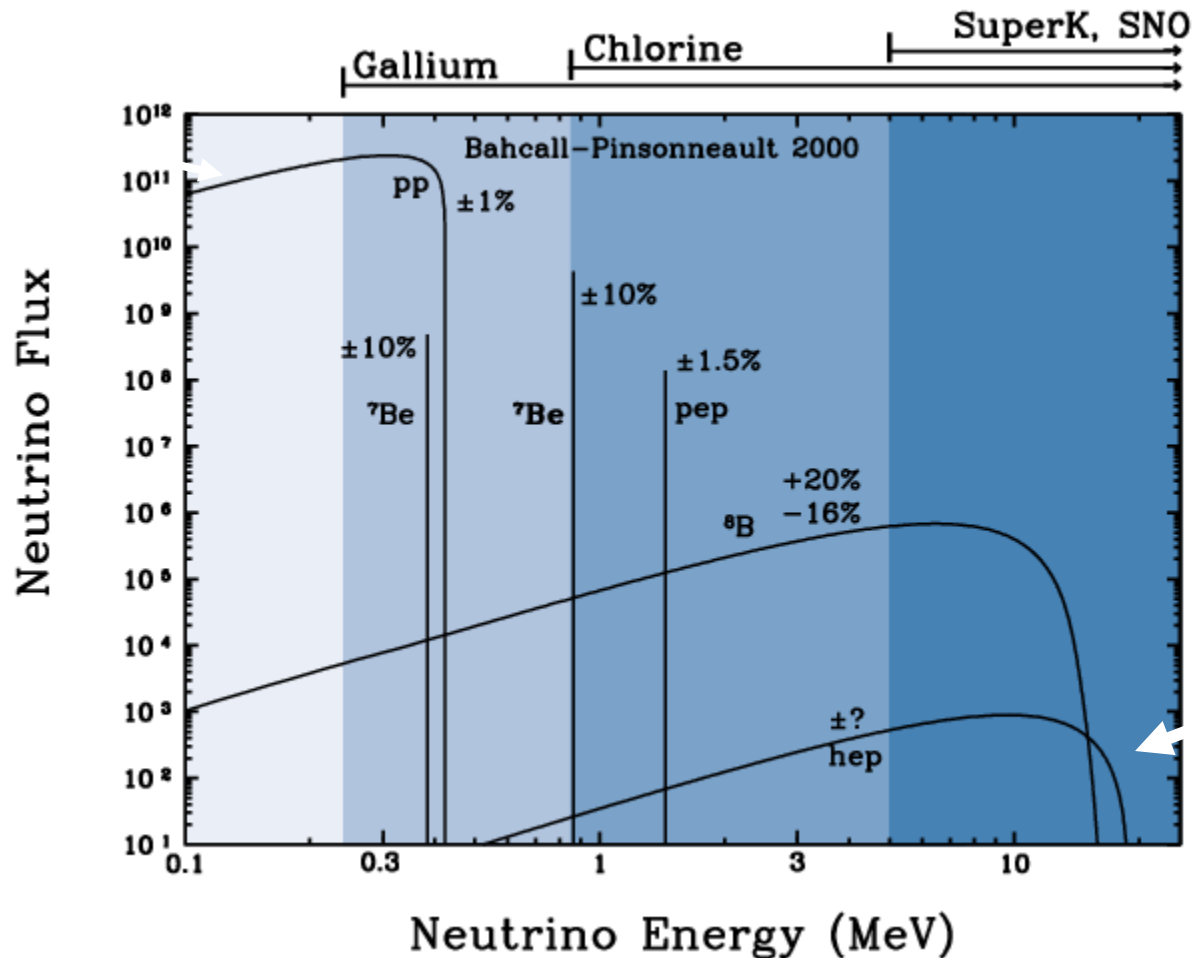


net result: $4p \rightarrow {}^4\text{He} + 2e^+ + 2\nu + Q_{\text{eff}}$



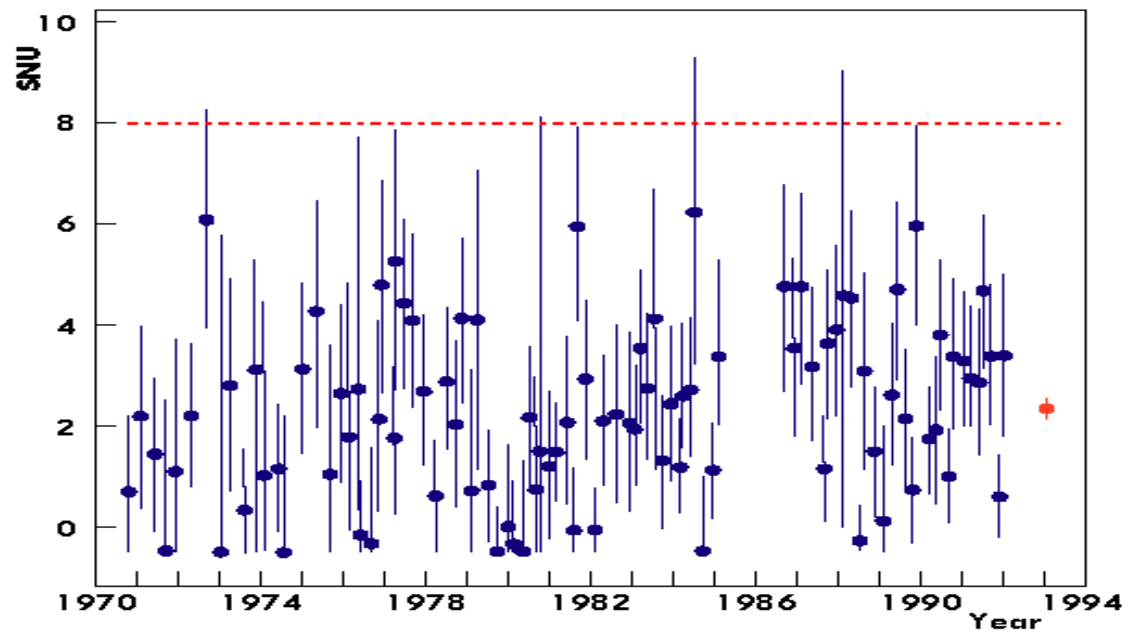
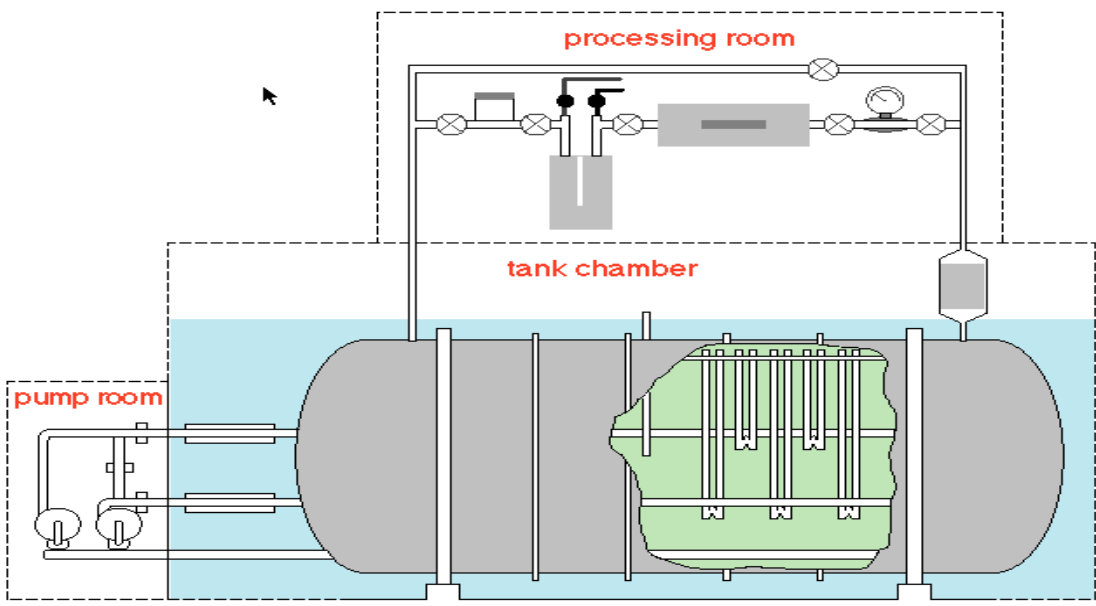
Question : solar constant $0.033 \text{ cal/sec/cm}^2$
 What is the number of protons consumed in the sun so far?

Solar Neutrino Spectrum



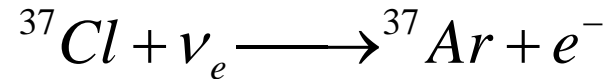
pp

${}^{\text{hep}}$



First experimental detection of solar neutrinos:

- **1964** John Bahcall and Ray Davis have the idea to detect solar neutrinos using the reaction:



- **1967 Homestake experiment starts taking data**
 - 100,000 Gallons of cleaning fluid in a tank 4850 feet underground
 - ^{37}Ar extracted chemically every few months (single atoms !)
and decay counted in counting station (35 days half-life)
 - event rate: ~1 neutrino capture per day !
- **1968 First results: only 34% of predicted neutrino flux !**

solar neutrino problem is born - for next 20 years no other detector !

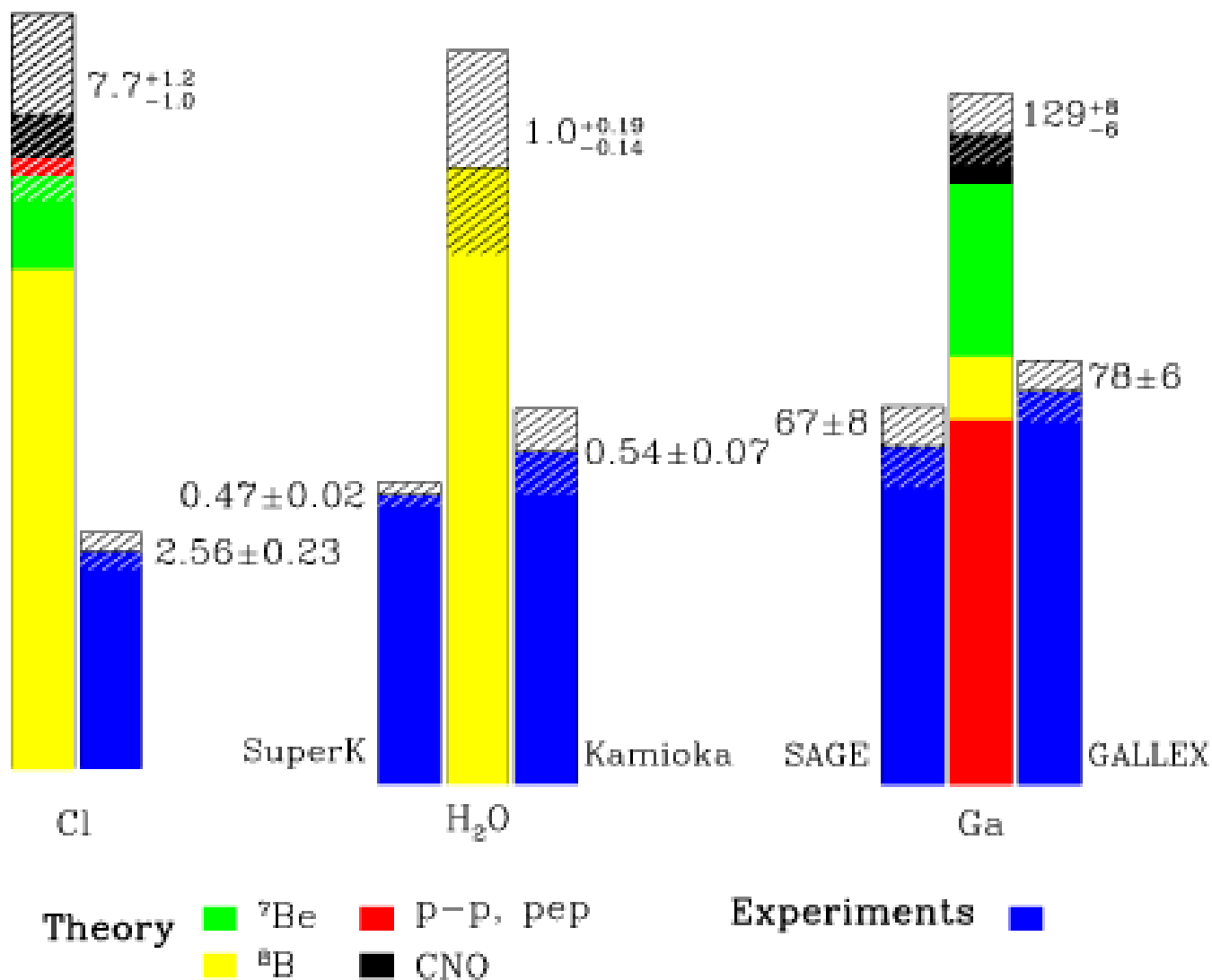
Neutrino production in solar core ~ T^{25}

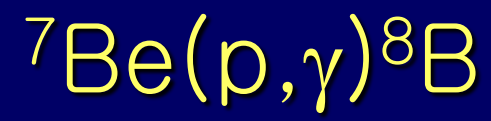
- **nuclear energy source of sun directly and unambiguously confirmed**
- **solar models precise enough so that deficit points to serious problem**



Total Rates: Standard Model vs. Experiment

Bahcall-Pinsonneault 98





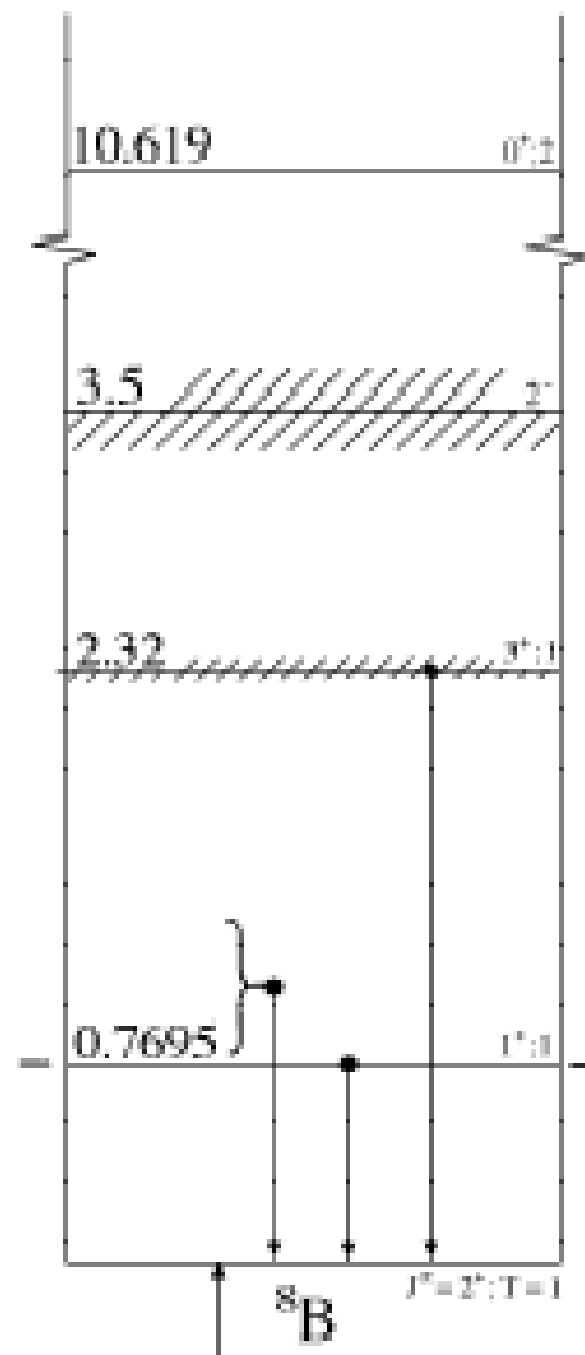
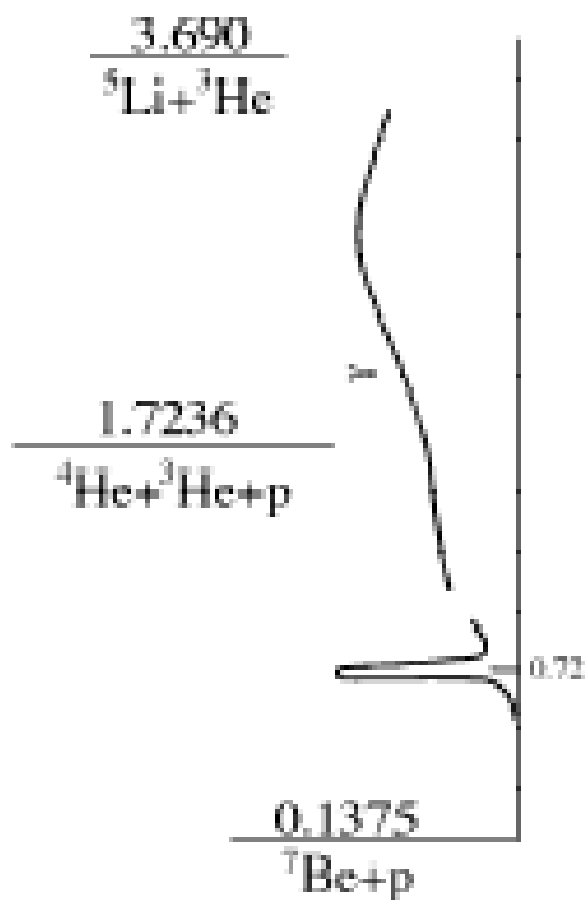
Levels and γ -ray branchings:

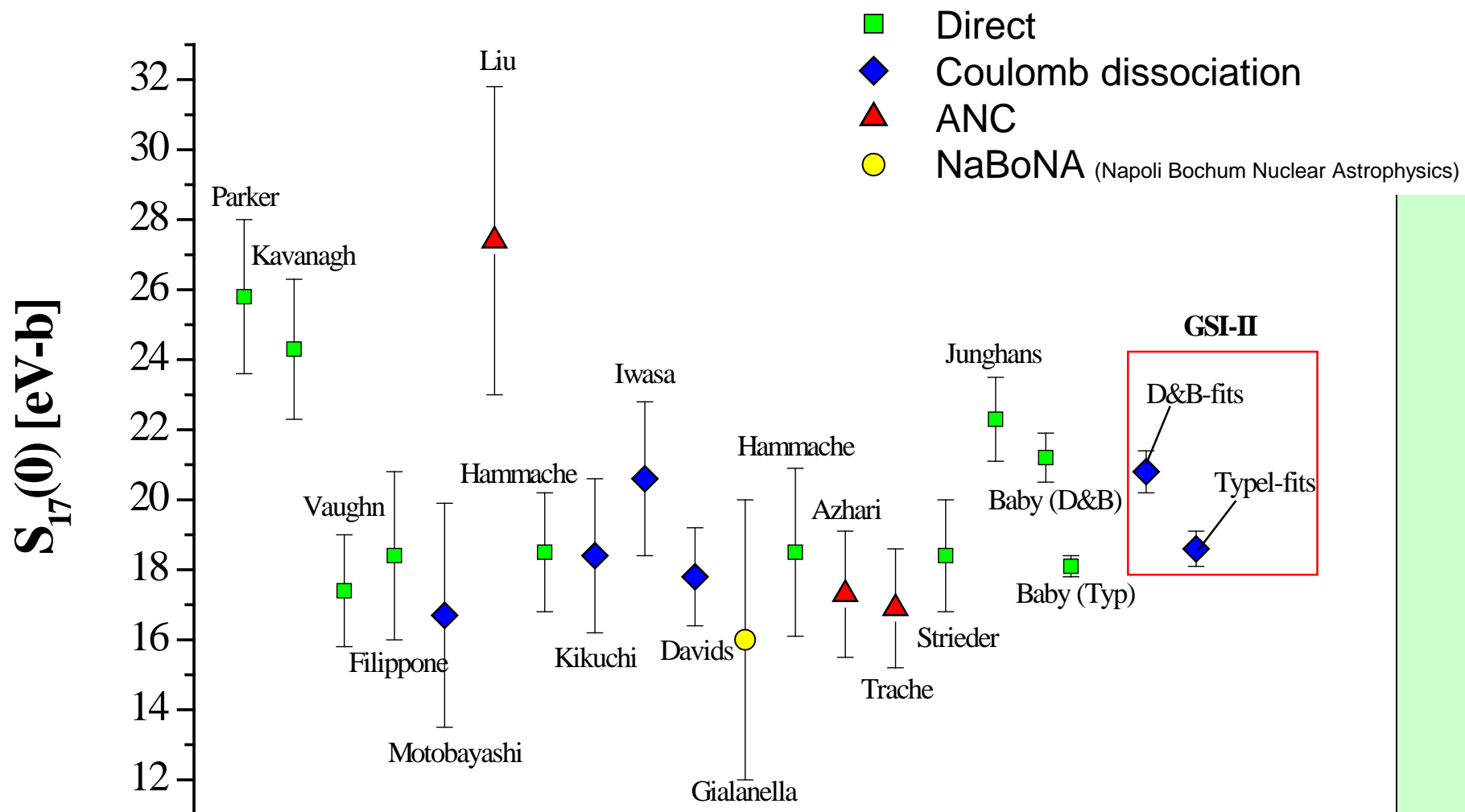
0, 2^+ , 770.3 ms, [ABCDEFGH], T=1,
 %EC+% β^+ =100, %EC2 α =100,
 $\mu=1.03553$

774.6, $\Gamma=37.5$ keV, [ABCEGH]
 γ_{0774} M1

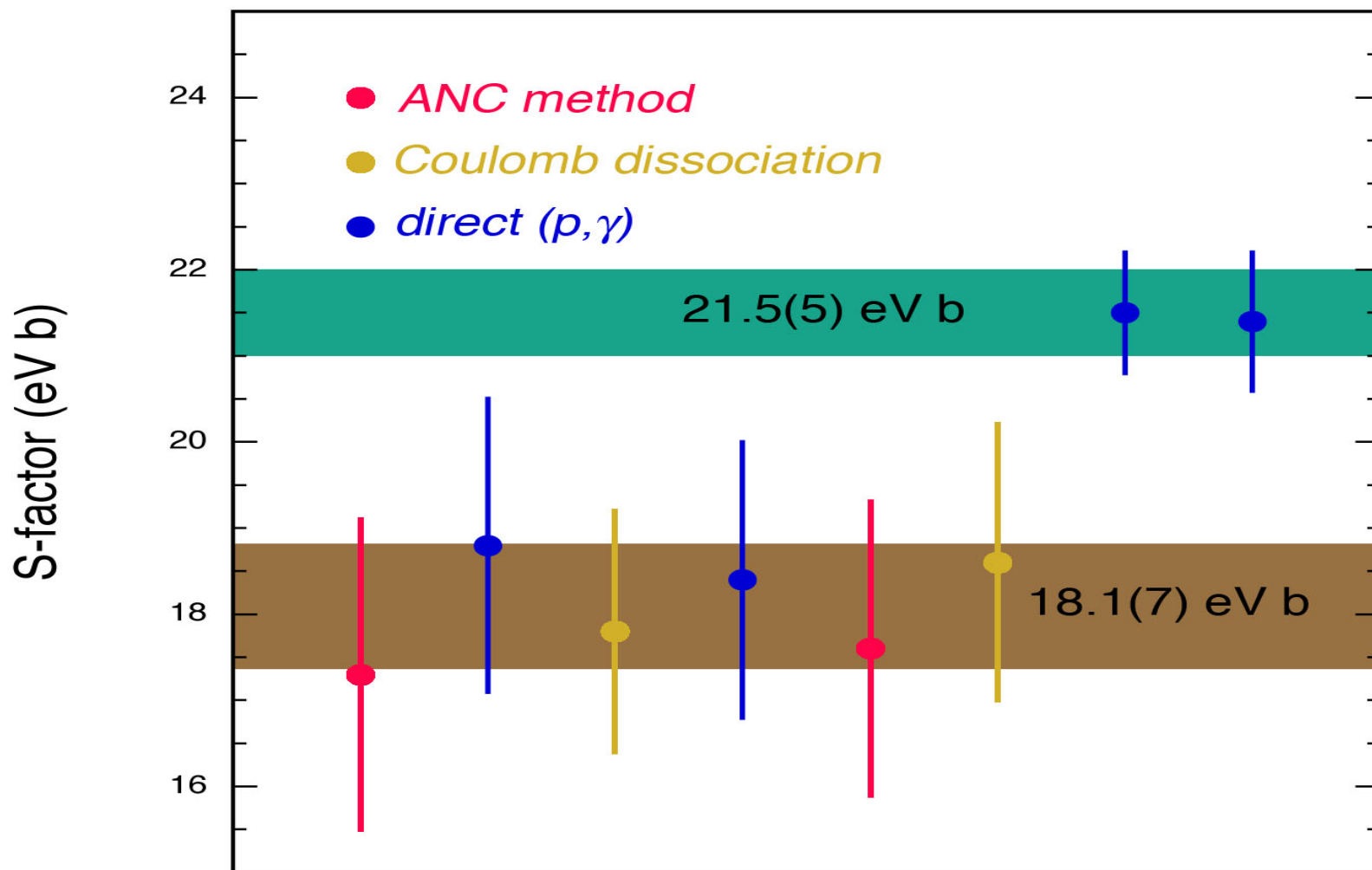
2320.30, 3^+ , $\Gamma=350.40$ keV, [CGH], T=1

10619.9, 0^+ , $\Gamma<60$ keV, [H], T=2





${}^7\text{Be}(p,\gamma){}^8\text{B}$, 2001 - present



Solar fusion cross sections II: the pp chain and CNO cycles

E. G. Adelberger, A. García, R. G. Hamish Robertson, and K. A. Snover

*Department of Physics and Center for Experimental Nuclear Physics and Astrophysics,
University of Washington, Seattle, WA 98195 USA*

A. B. Balantekin, K. Heeger, and M. J. Ramsey-Musolf

Department of Physics, University of Wisconsin, Madison, WI 53706 USA

D. Bemmerer and A. Junghans

Forschungszentrum Dresden-Rossendorf, D-01314 Dresden, Germany

C. A. Bertulani

Department of Physics and Astronomy, Texas A&M University, Commerce, TX 75429 USA

J.-W. Chen

Department of Physics and Center for Theoretical Sciences, National Taiwan University, Taipei 10617, Taiwan

H. Costantini and P. Prati

Università di Genova and INFN Sezione di Genova, Genova, Italy

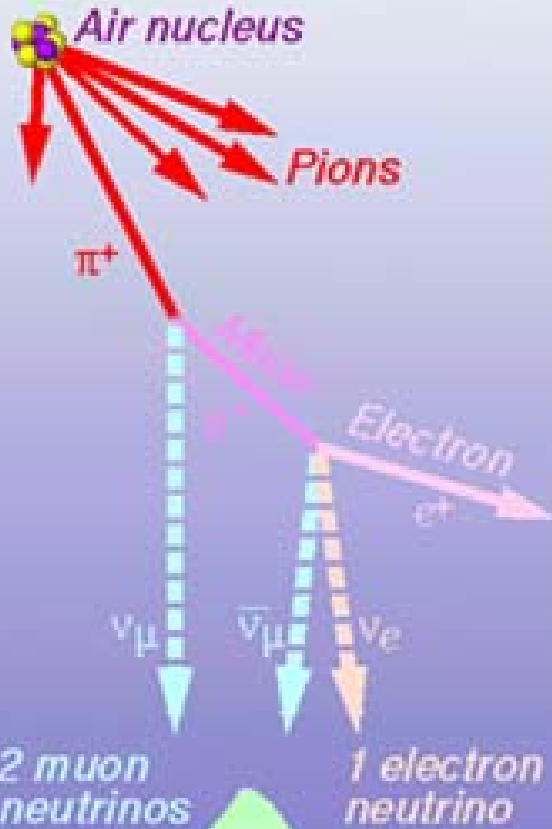
M. Couder, E. Überseder, and M. Wiescher

Department of Physics and JINA, University of Notre Dame, Notre Dame, IN 46556 USA

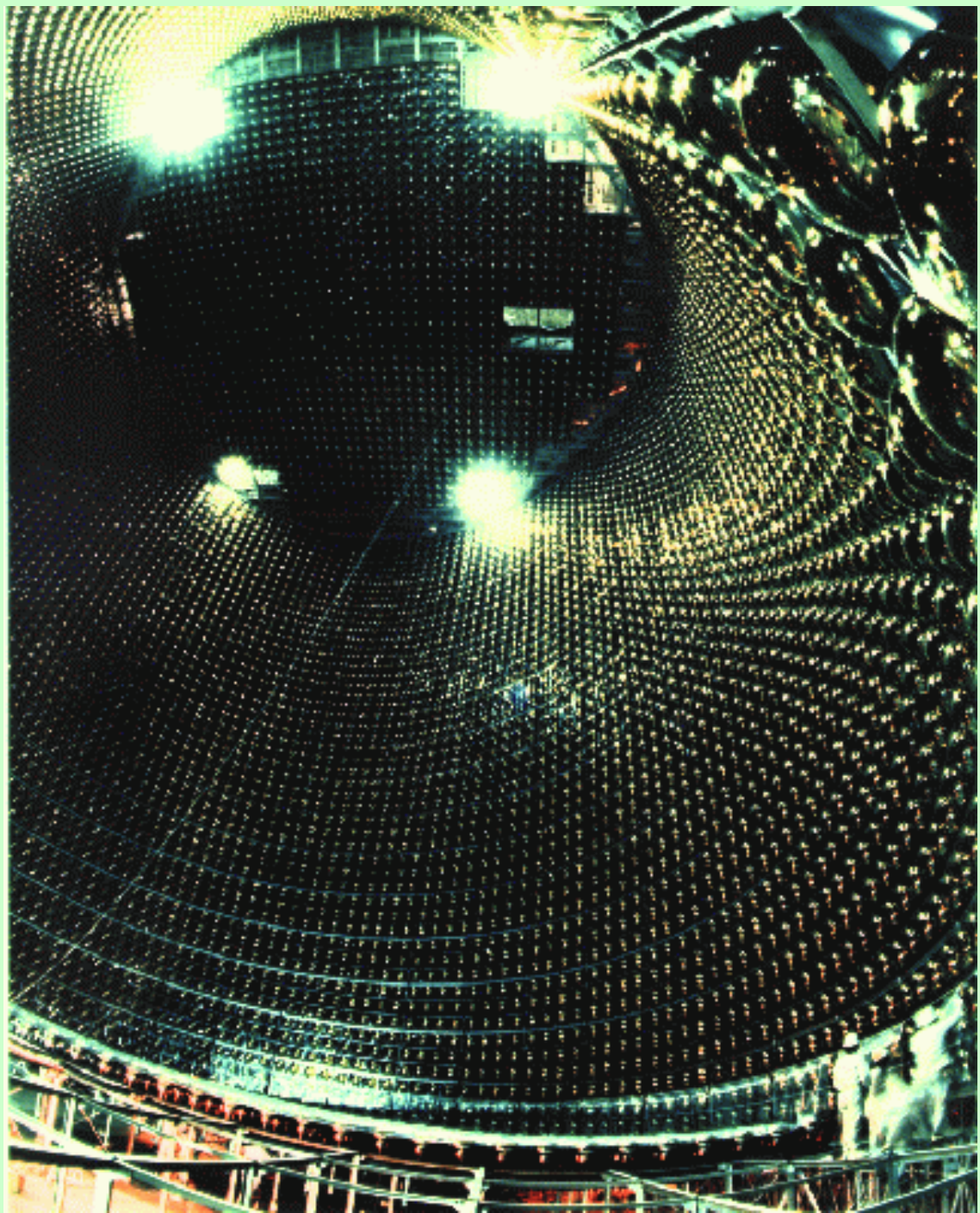
R. Cyburt

JINA and National Superconducting Cyclotron Laboratory, Michigan State University, East Lansing, MI 48824 USA

Cosmic Ray



Super-K
Detector

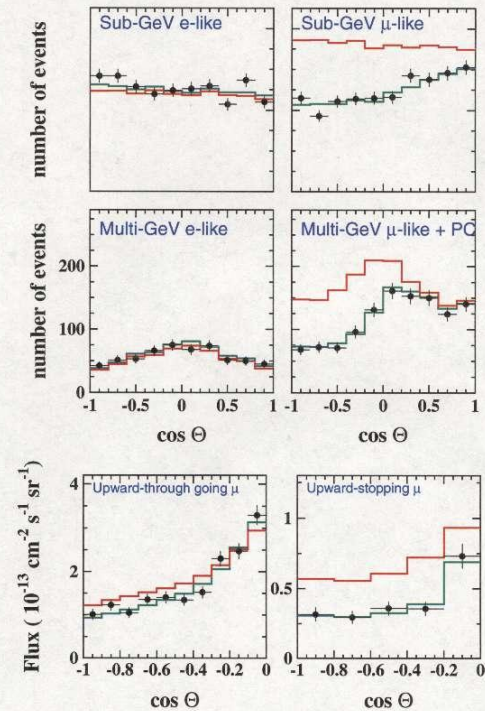


Atmospheric neutrinos

Super-Kamiokande



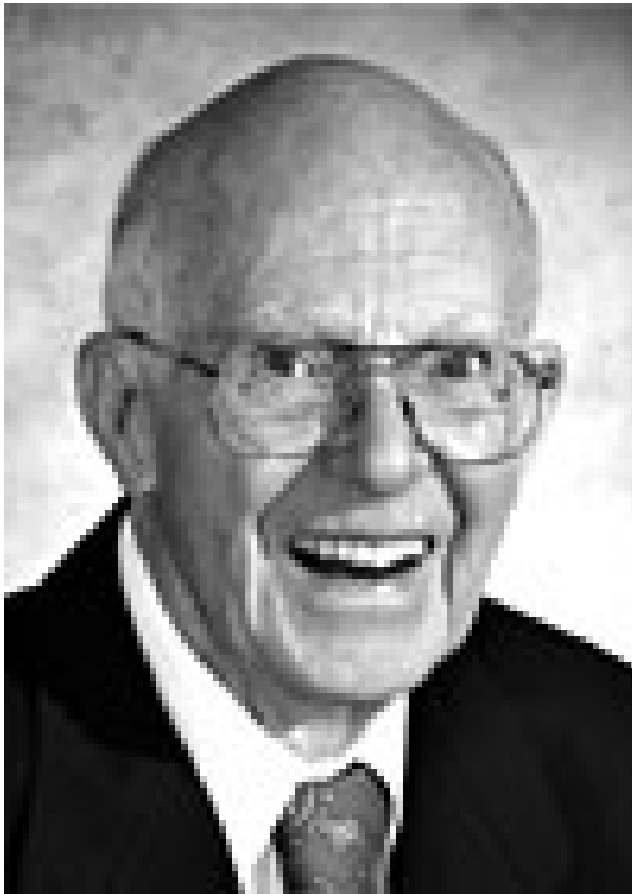
Zenith distributions with combined fit



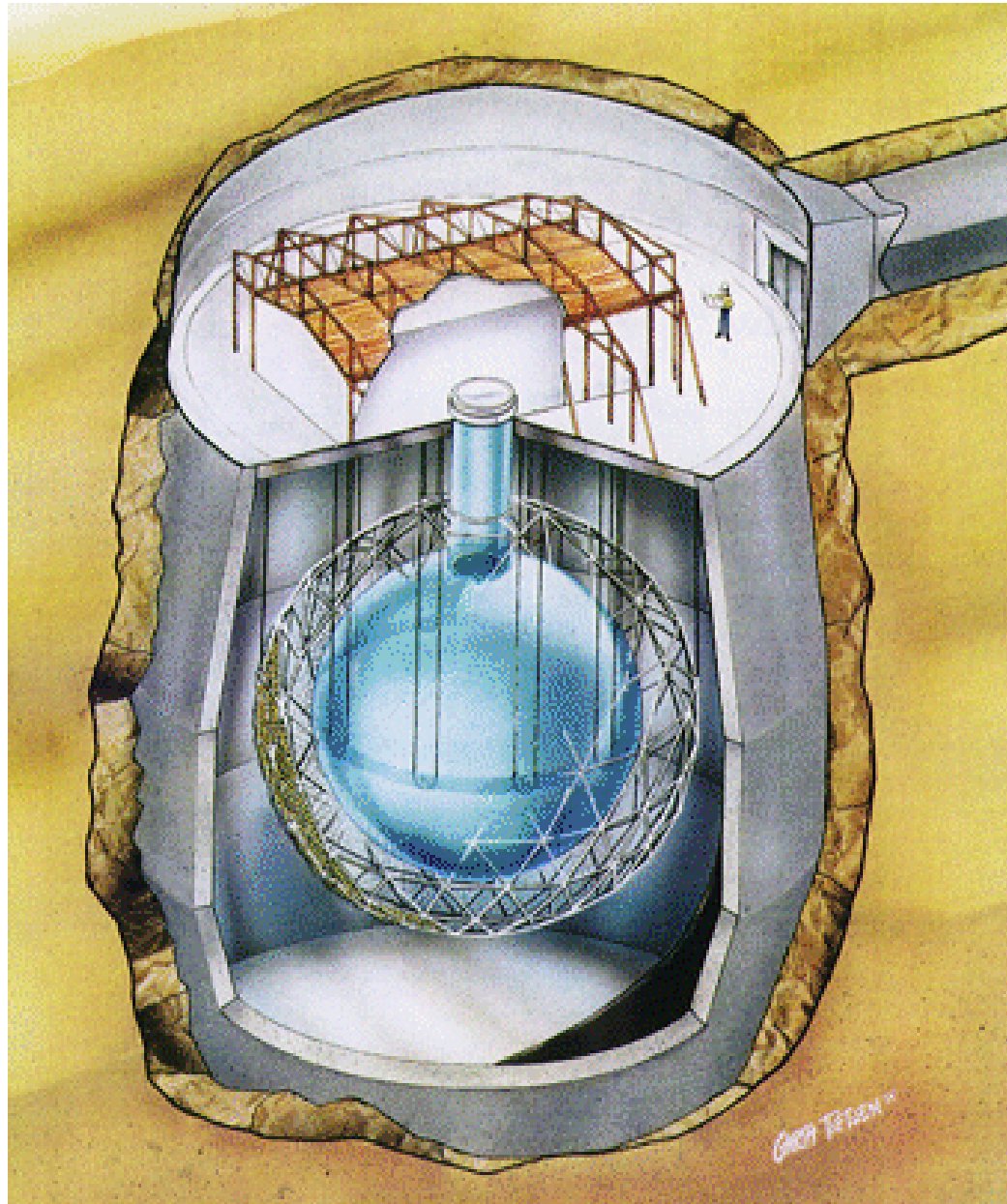


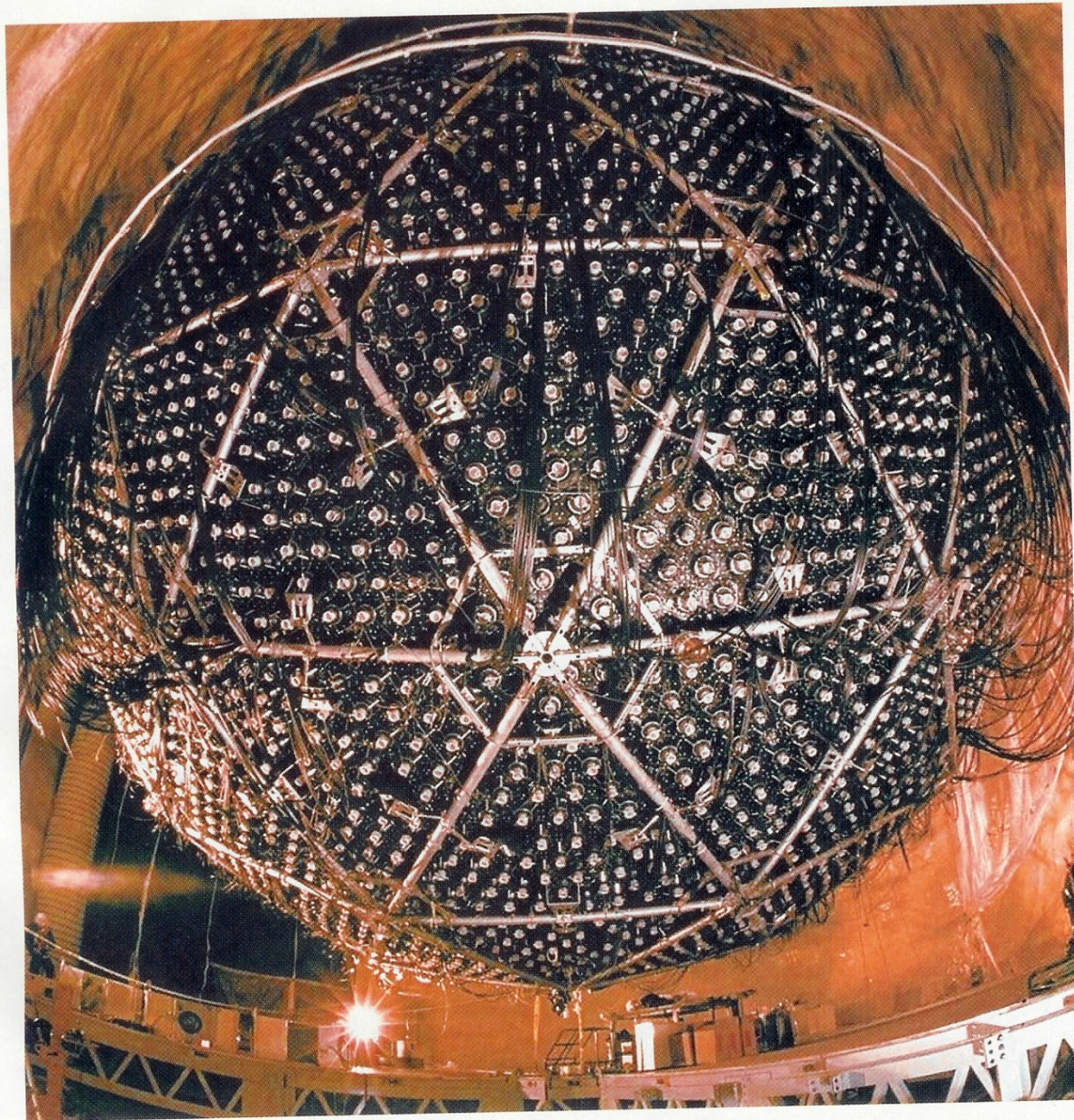
The Nobel Prize in Physics 2002

"for pioneering contributions to astrophysics, in particular for the detection of cosmic neutrinos"



Sudbury Neutrino Observatory

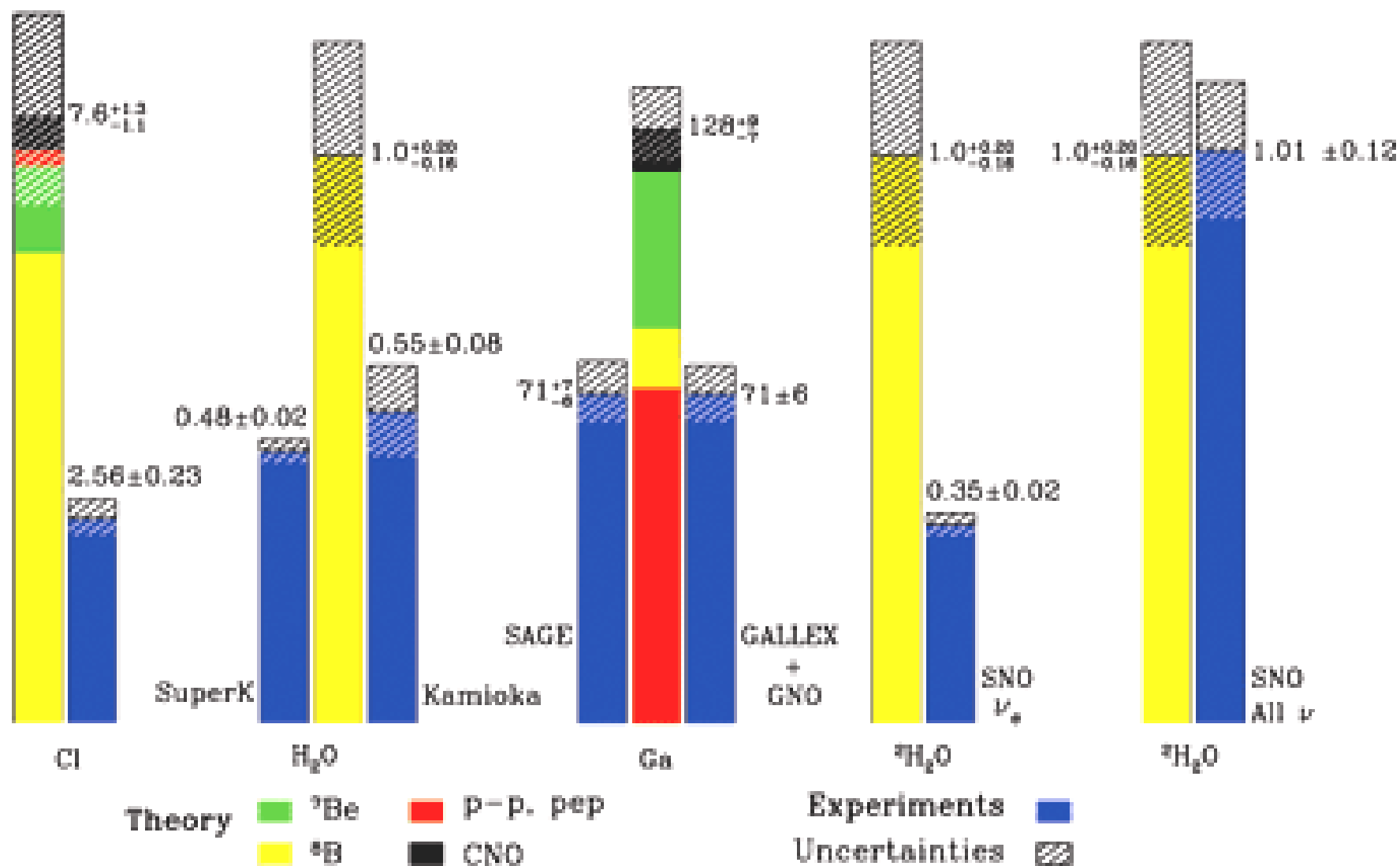




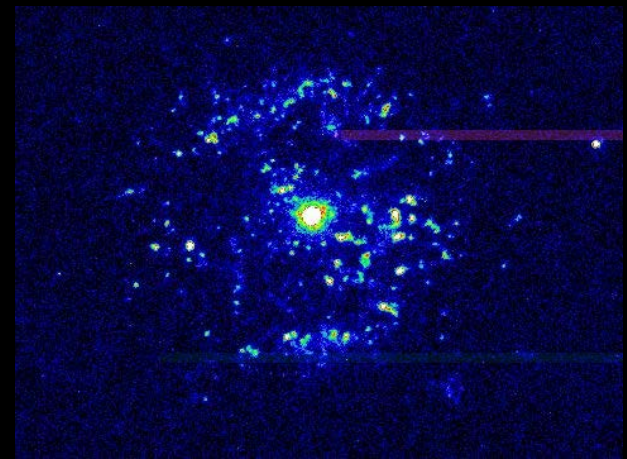
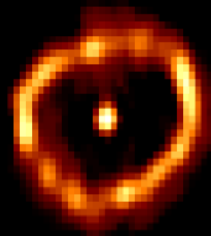
PHOTOMULTIPLIER TUBES—more than 9,500 of them—on a geodesic sphere 18 meters in diameter act as the eyes of the Sudbury Neutrino Observatory. The tubes surround and monitor a 12-meter-diameter acrylic sphere that contains 1,000 tons of heavy water. Each tube can detect a single photon

of light. The entire assembly is suspended in ordinary water. All the materials that make up the detector must be extraordinarily free of natural traces of radioactive elements to avoid overwhelming the tubes with false solar neutrino counts.

Total Rates: Standard Model vs. Experiment Bahcall-Pinsonneault 2000



Nova observations

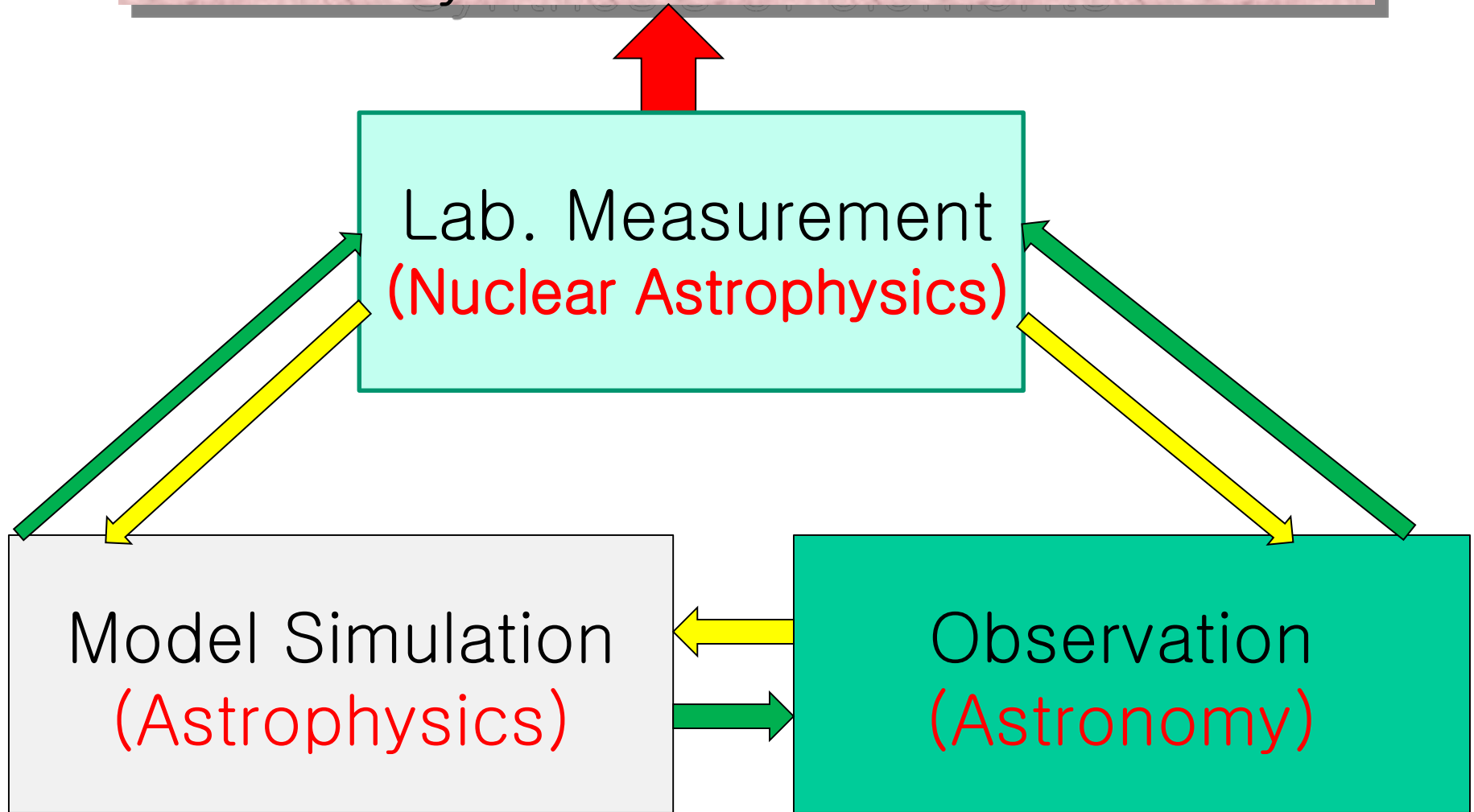


Evolution of the stars
Synthesis of elements

Lab. Measurement
(Nuclear Astrophysics)

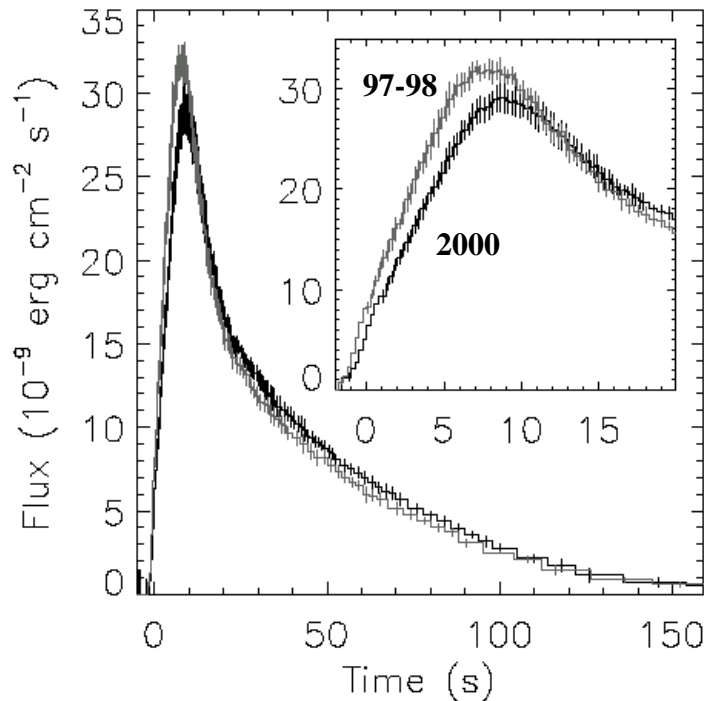
Model Simulation
(Astrophysics)

Observation
(Astronomy)



Another example: observation – model – nuclear astrophysics

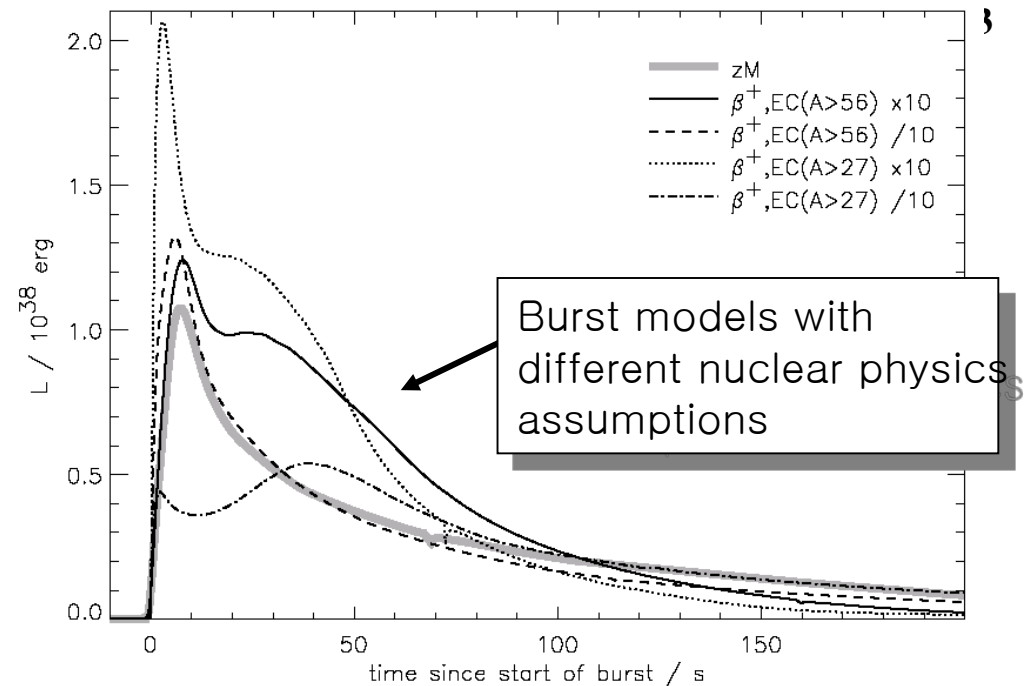
Precision X-ray observations (NASA's RXTE)



→ GS 1826-24 burst shape changes !

(Galloway 2003 astro/ph 0308122)

Uncertain models due to nuclear physics



Woosley et al. 2003 astro/ph 0307425



Need much more precise nuclear data to make full use of high quality observational data

The origin of galactic ^{26}Al

- Long-standing mystery in nuclear astrophysics:

provides important clues to the chemical evolution
of our galaxy

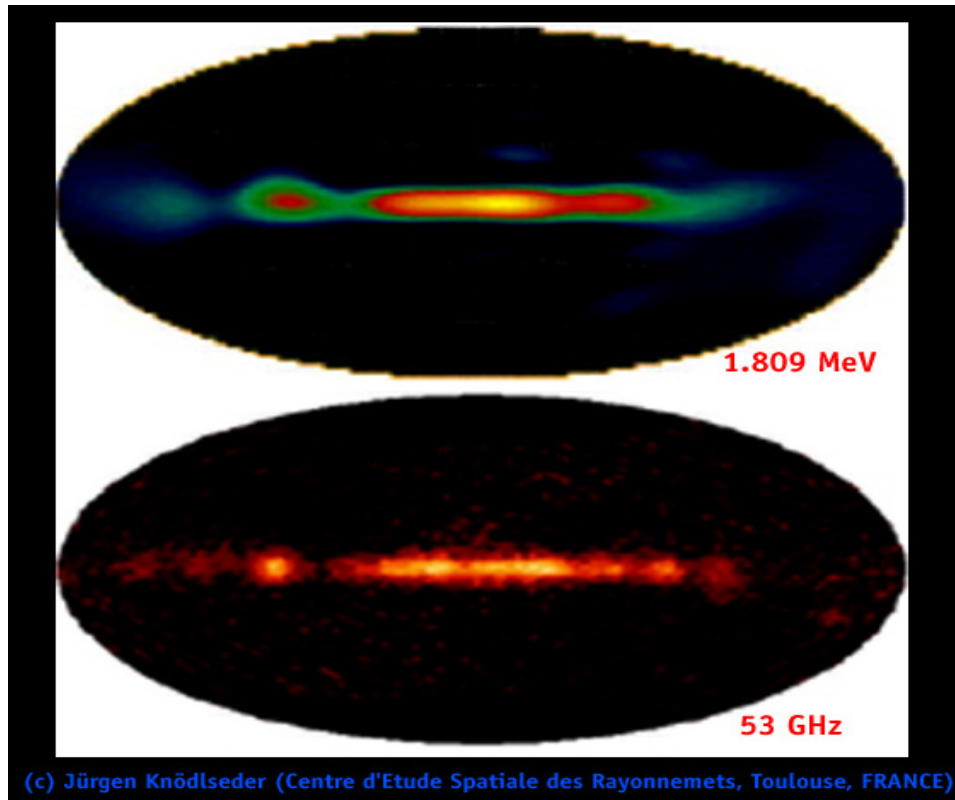
- ^{26}Al (ground state) decay:

$T_{1/2} = 0.7$ million years

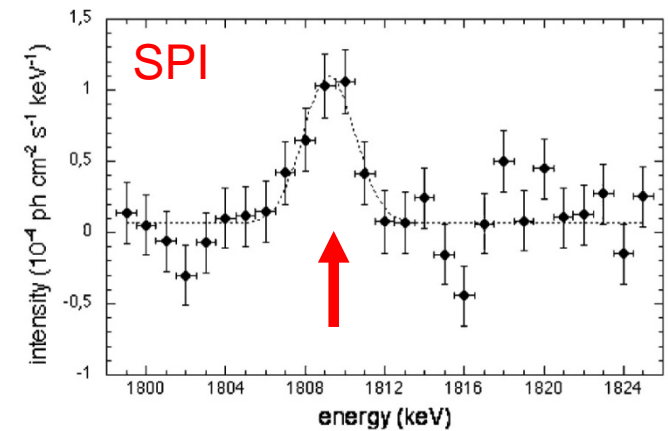
1.809 MeV γ -ray

- Contexts: [Gamma ray astronomy](#) and [presolar meteoritics](#)

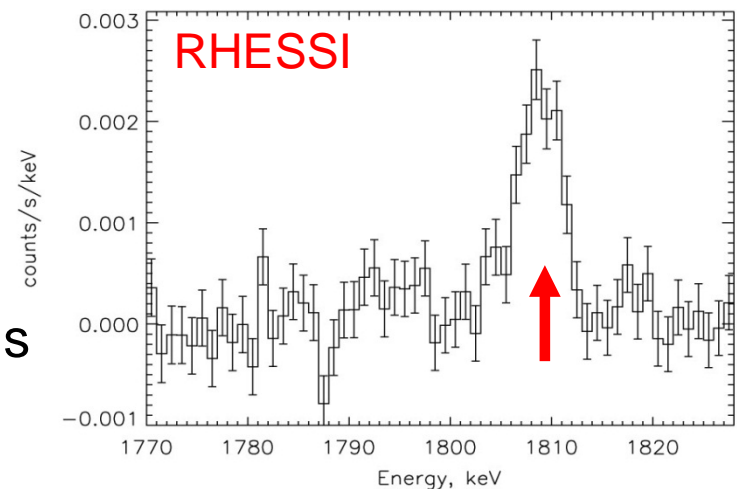
Galactic ^{26}Al and gamma ray astronomy



- Evidence for recent stellar nucleosynthesis
 - ~ 3 solar masses of ^{26}Al
 - $^{26}\text{Al} / ^{27}\text{Al} \sim 10^{-7}$
- (R. Diehl et al., Nature 2006)



J. Knödlseider, New Ast. Rev. (2004)

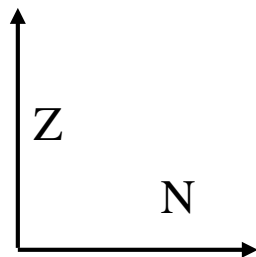


D.M. Smith, Ap J Lett. (2003)

From Alan Chen @ McMaster

Stellar sources of galactic ^{26}Al ?

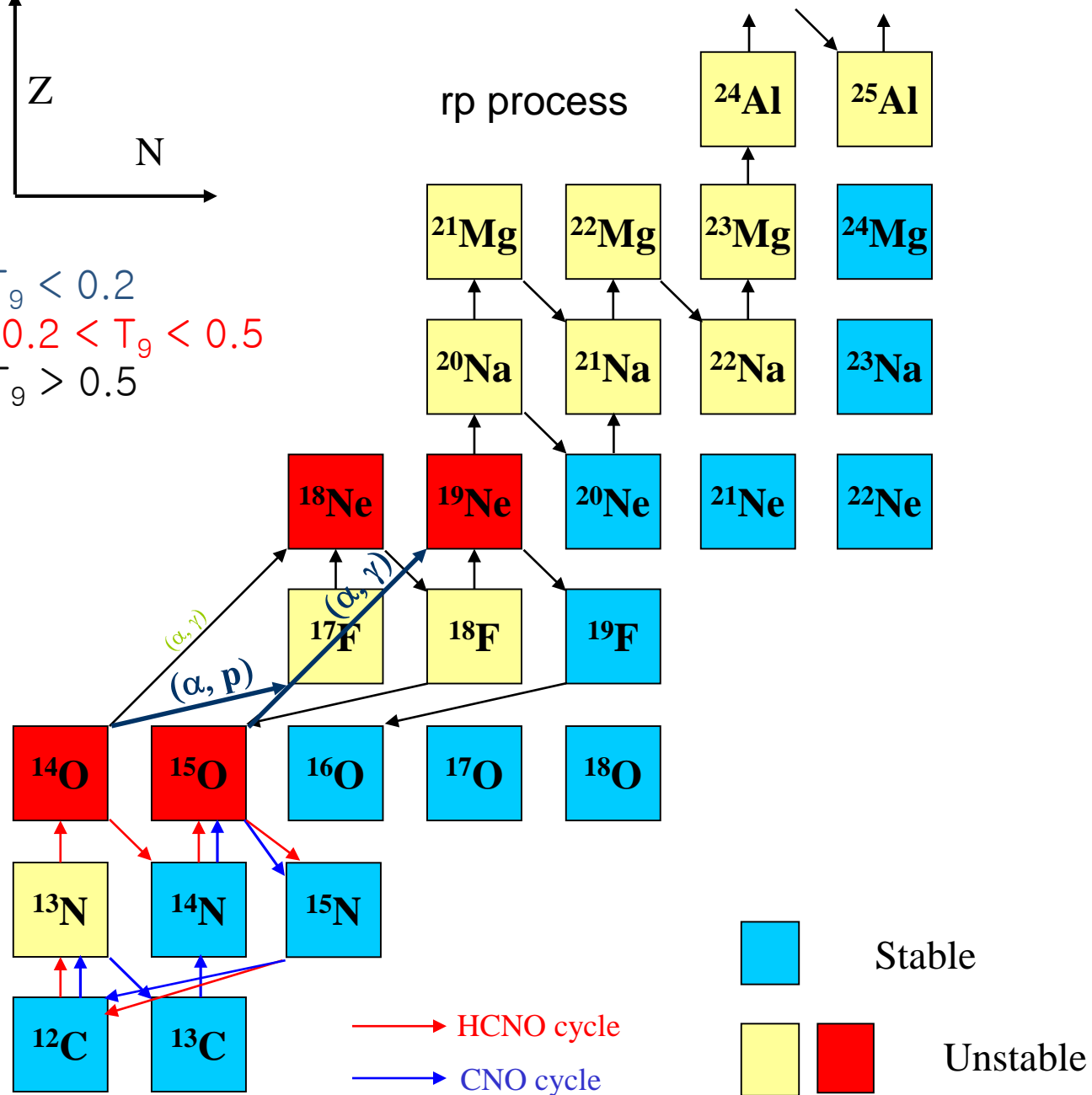
- Stellar production sites?
 - Gamma ray map favors massive stars
 - Core collapse supernovae
 - Wolf-Rayet stars
 - Contributions from other sources can be significant
 - Nova explosions
 - AGB stars
- **Novae:** ^{26}Al production in explosive hydrogen burning
- Important reactions: $^{26}\text{Al}(p,\gamma)^{27}\text{Si}$, $^{25}\text{Al}(p,\gamma)^{26}\text{Si}$



CNO cycle : $T_9 < 0.2$

HCNO cycle: $0.2 < T_9 < 0.5$

rp process : $T_9 > 0.5$



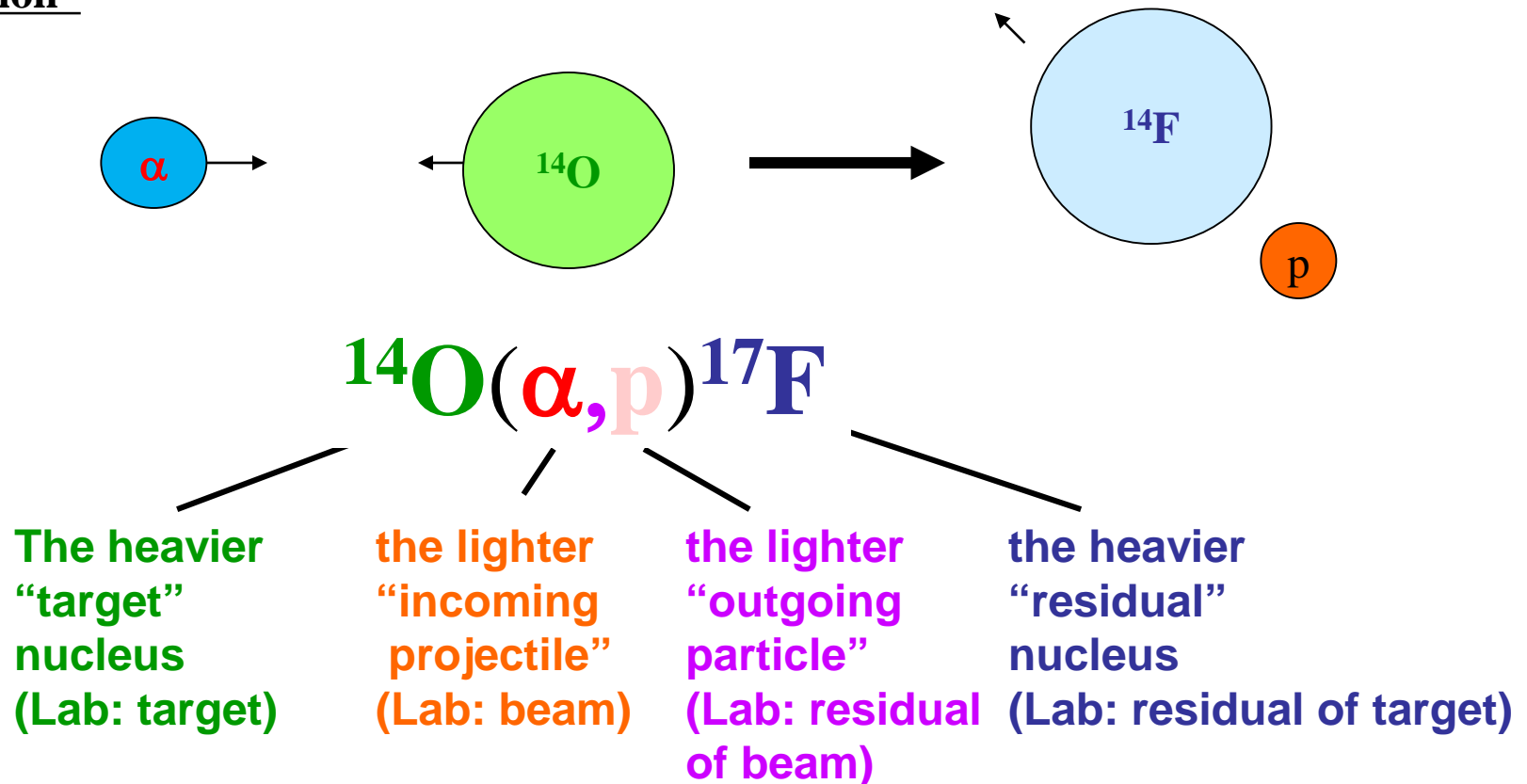
Studies on ^{18}Ne using stable beams and targets

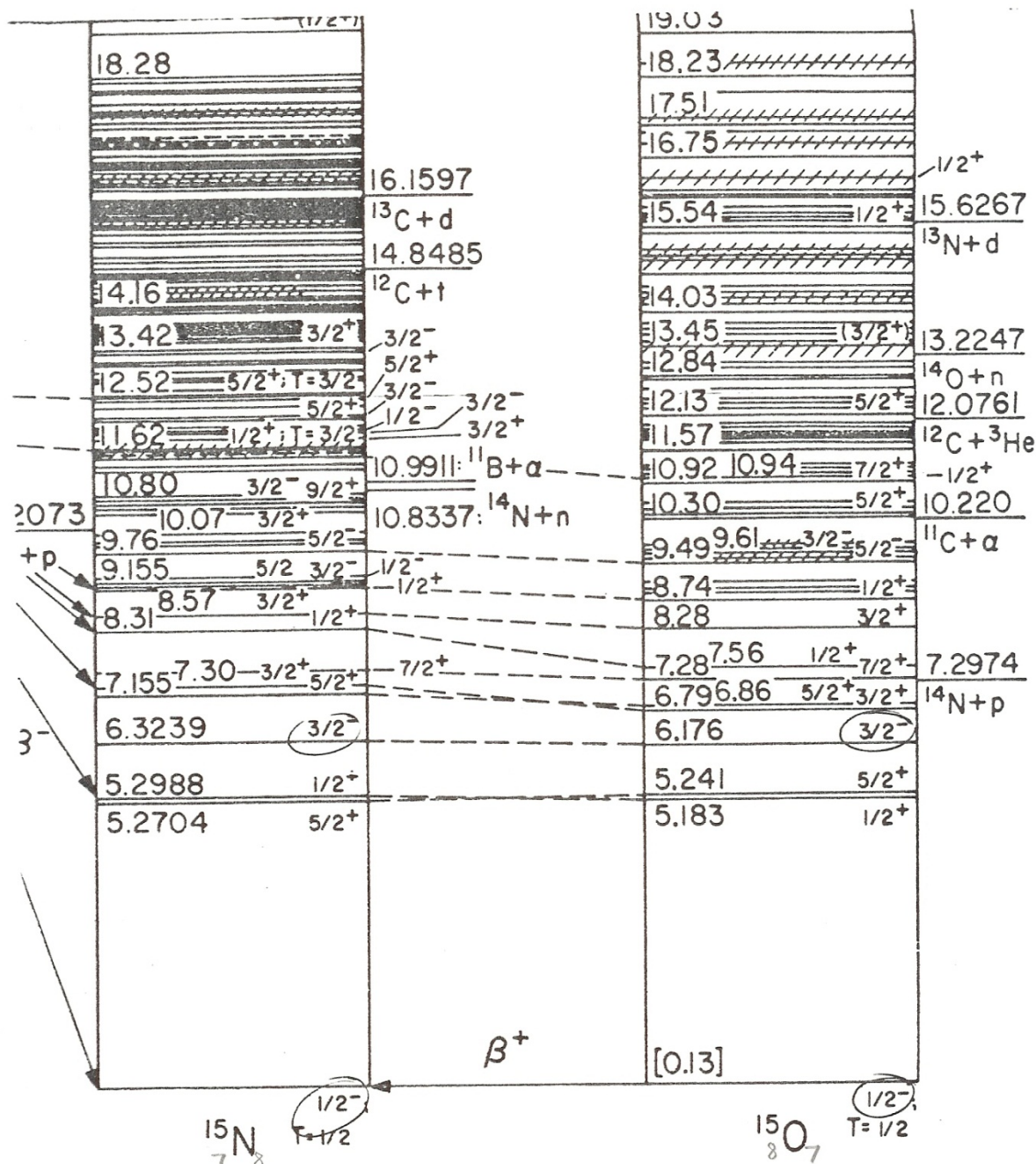
- $^{19}\text{F}(\text{p},\text{t})^{18}\text{Ne}$
@ Princeton, IUCF, INS(CNS)
- $^{16}\text{O}({}^3\text{He},\text{n})^{18}\text{Ne}$
@Univ. of Washington
- $^{12}\text{C}({}^{12}\text{C}, {}^6\text{He})^{18}\text{Ne}$
@Yale Univ.

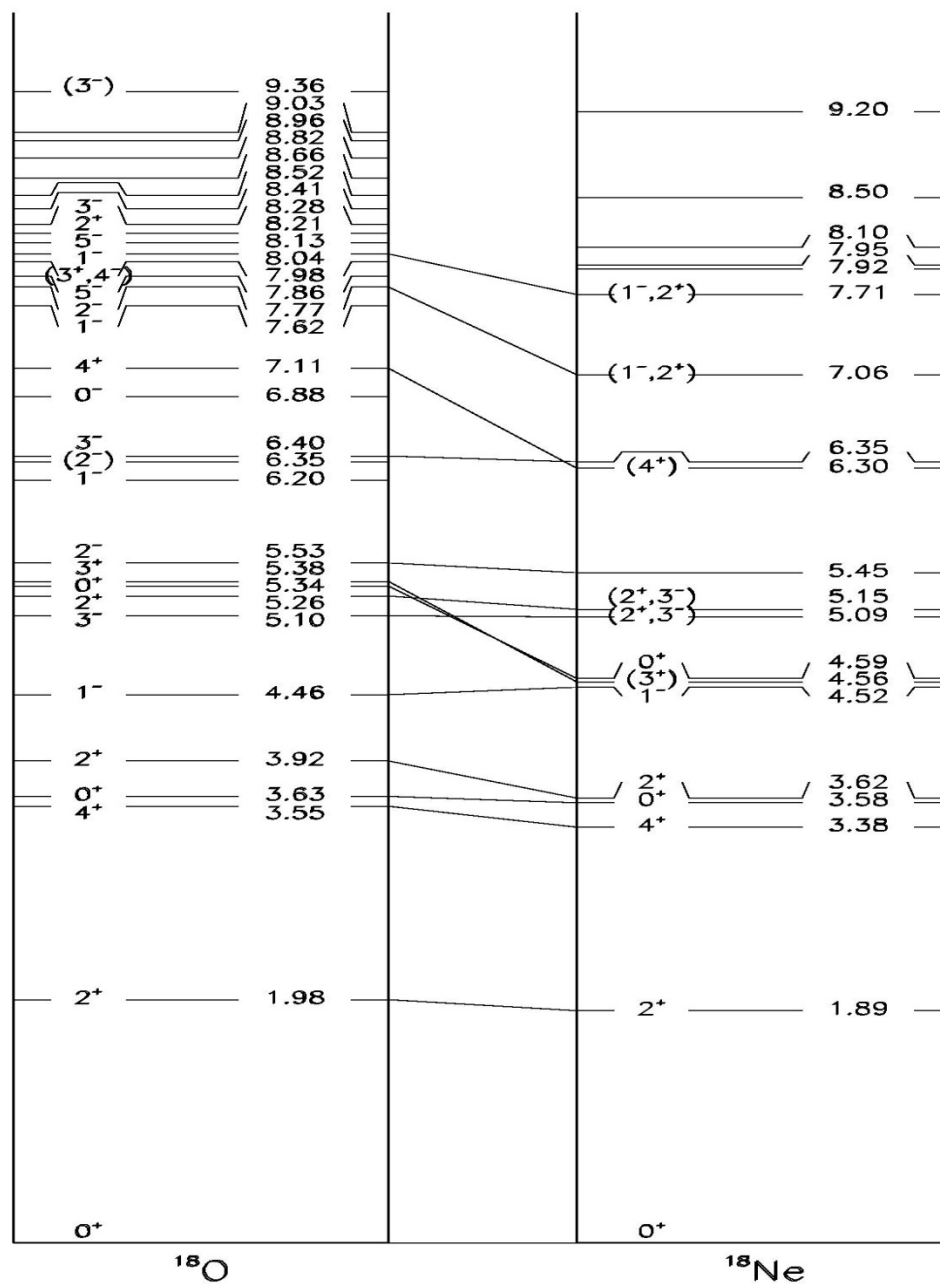
Nuclear Reactions



Notation

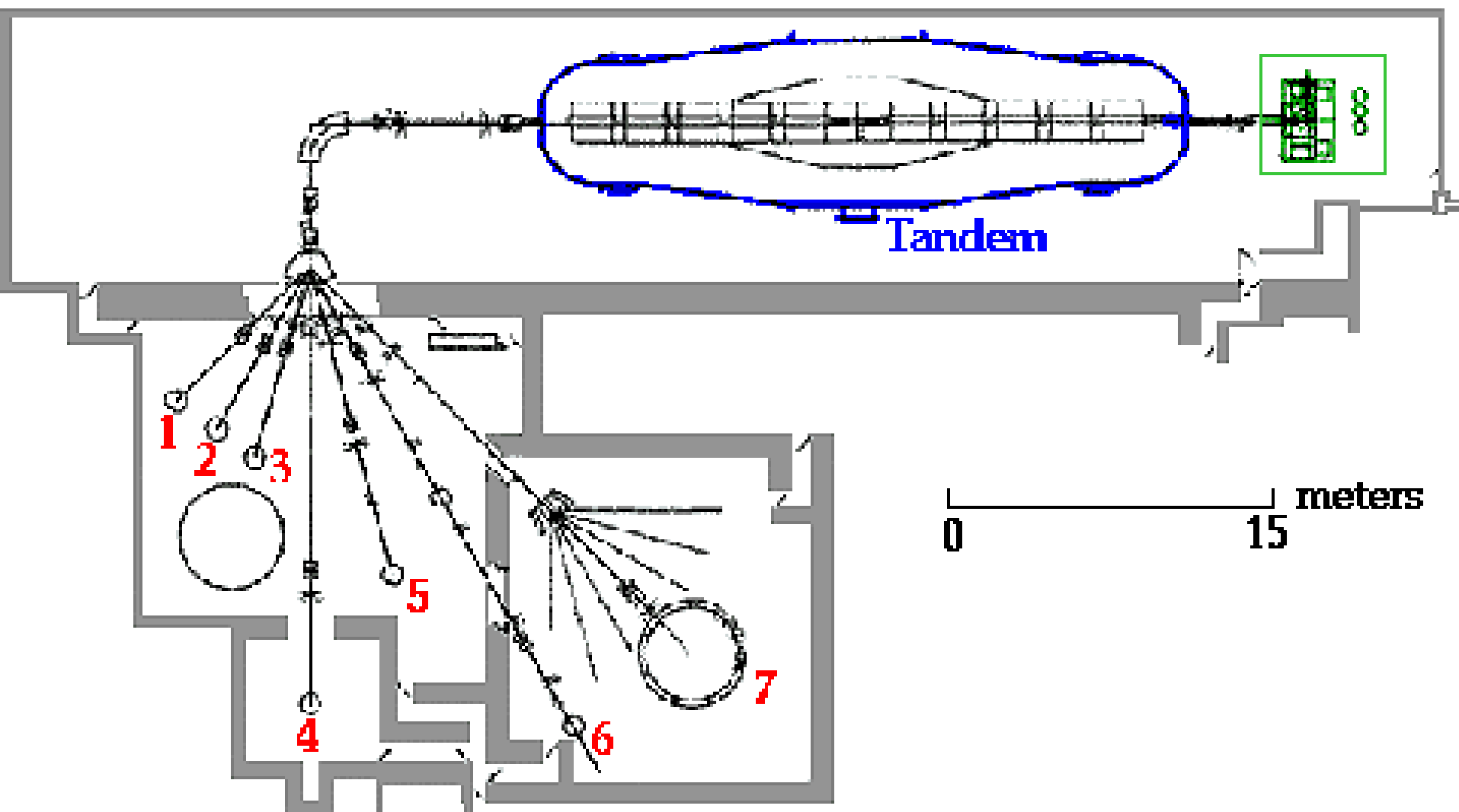


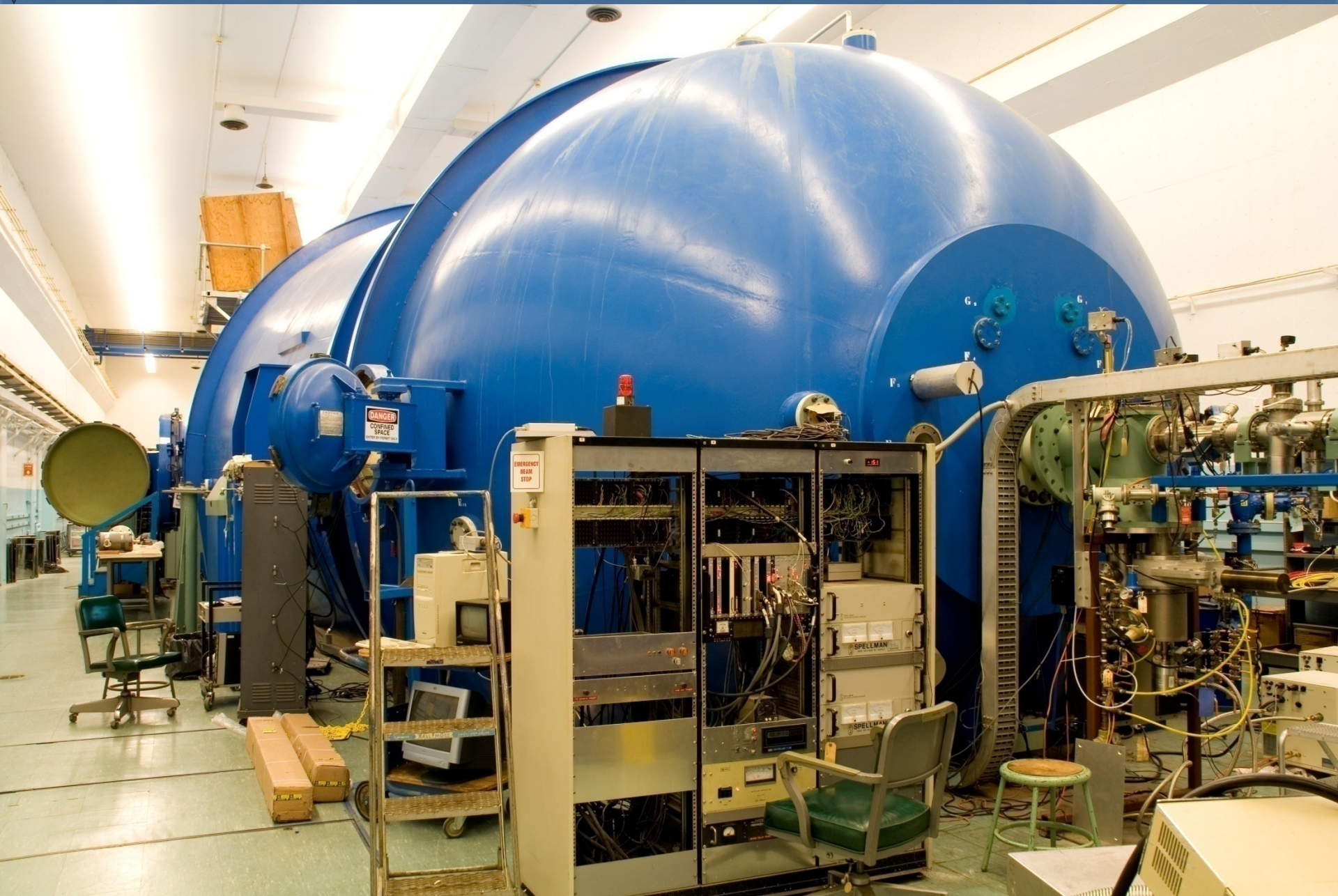


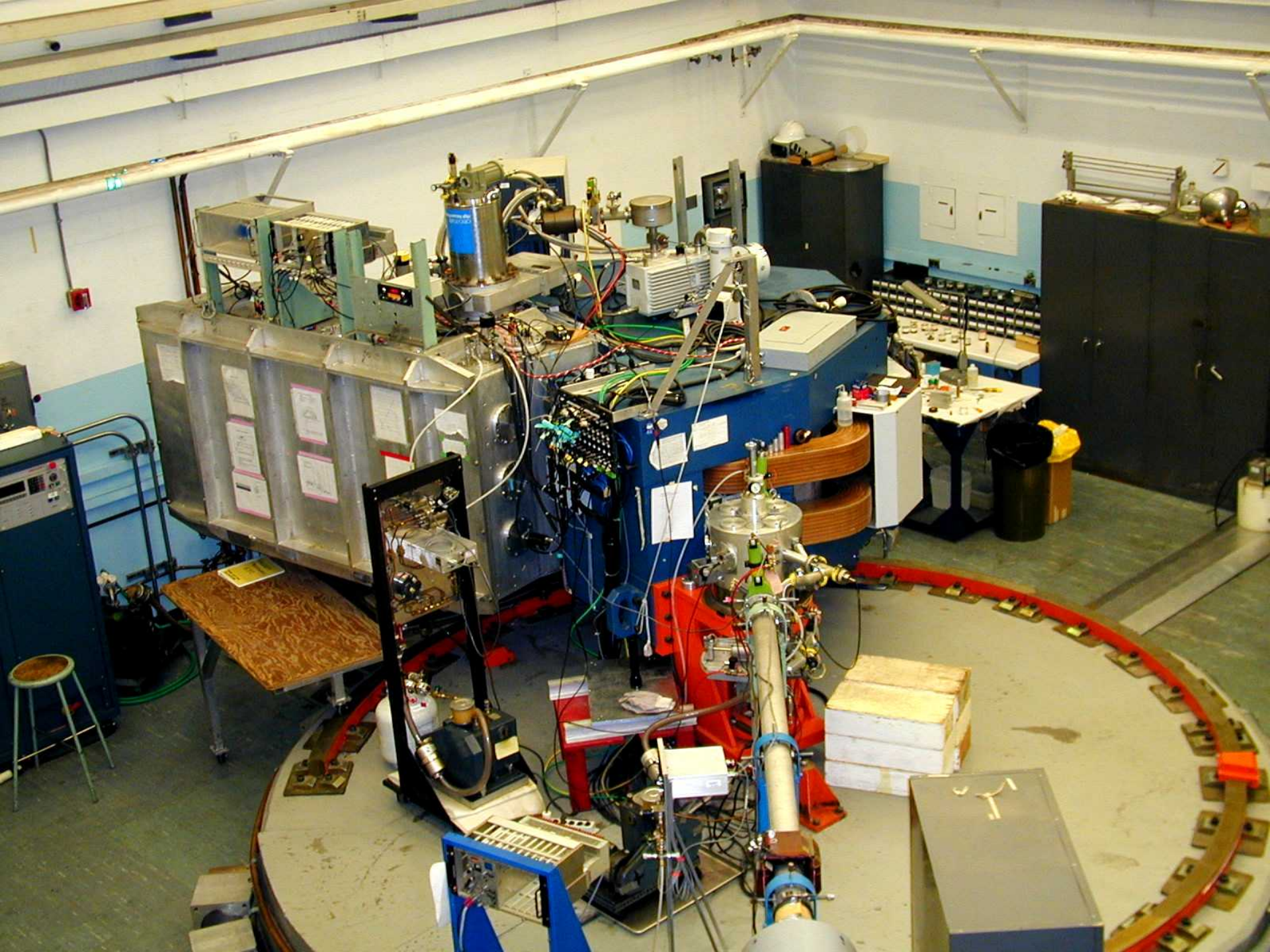


$$\frac{5.114}{^{14}\text{O} + \alpha}$$

$$\frac{3.922}{^{17}\text{F} + \text{p}}$$







Performance Characteristics of the Yale Split-Pole Spectrograph

Solid Angle: $160 \text{ mrad} \times 80 \text{ mrad} = 12.8 \text{ msr}$

Orbit radii at full solid angle: 51.1 cm to 92.0 cm

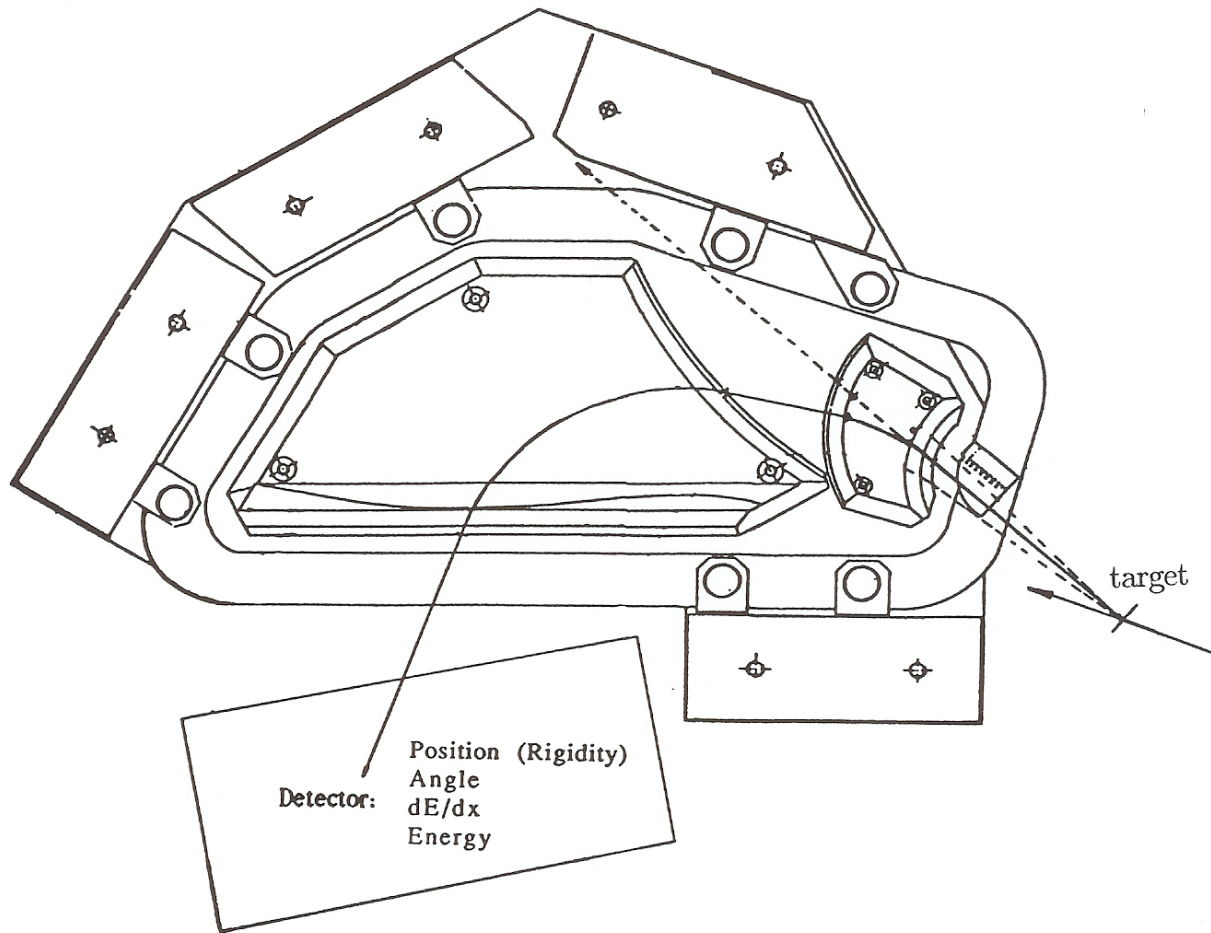
First order resolution for 1mm target spot:

$$\Delta p/p = 1/4290 \text{ for } \rho = 92 \text{ cm}$$

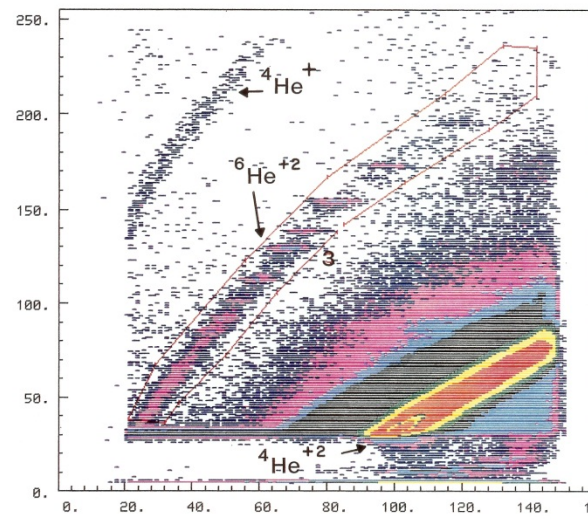
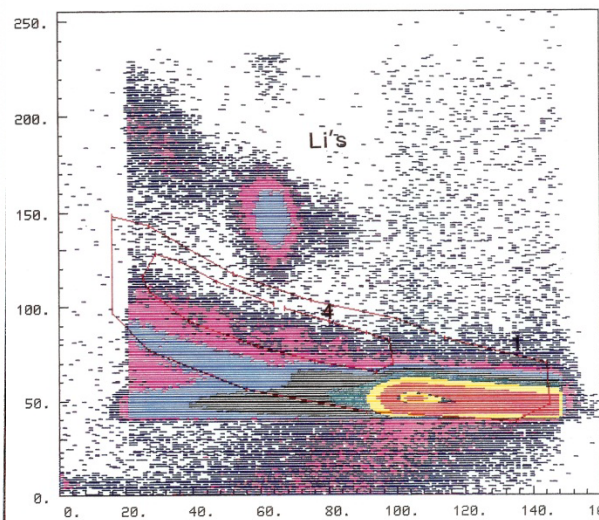
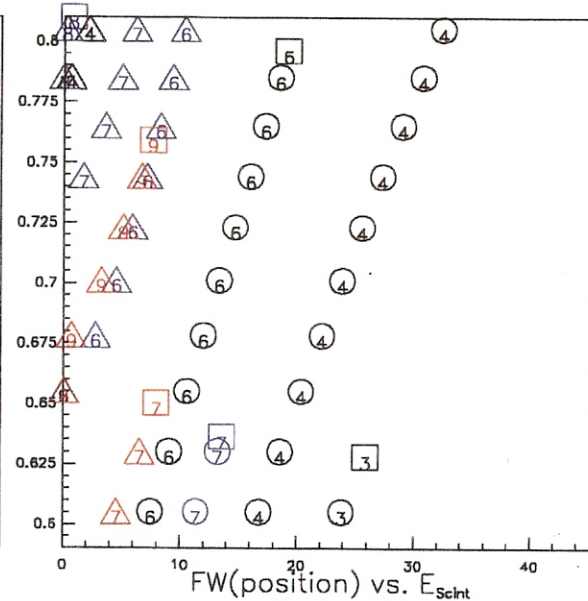
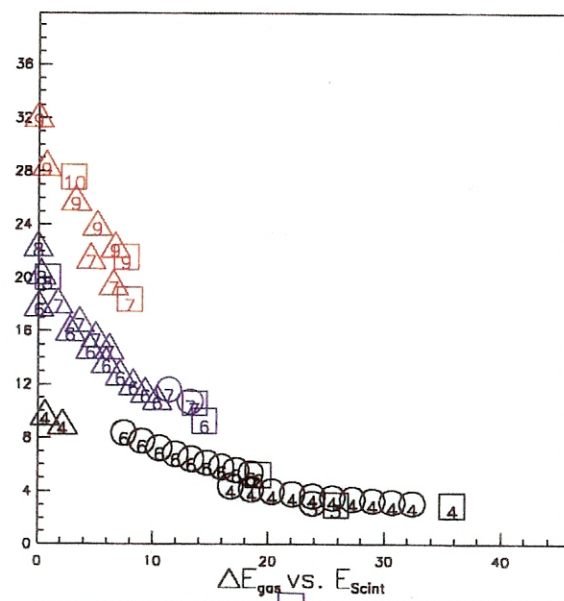
The momentum range: $P_{\text{max}}/P_{\text{min}} = 1.80$

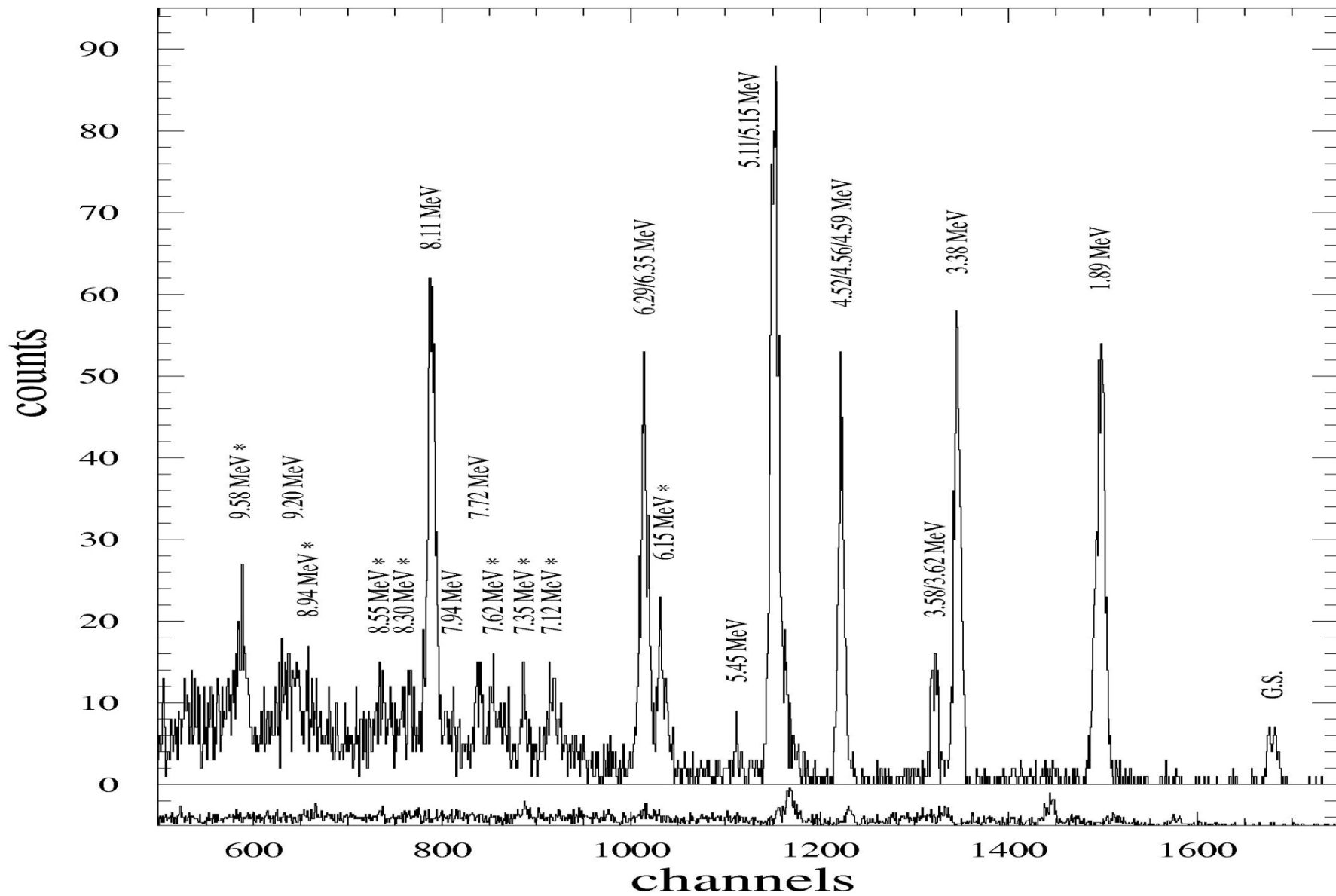
Maximum field strength: $B = 16.3 \text{ kG}$

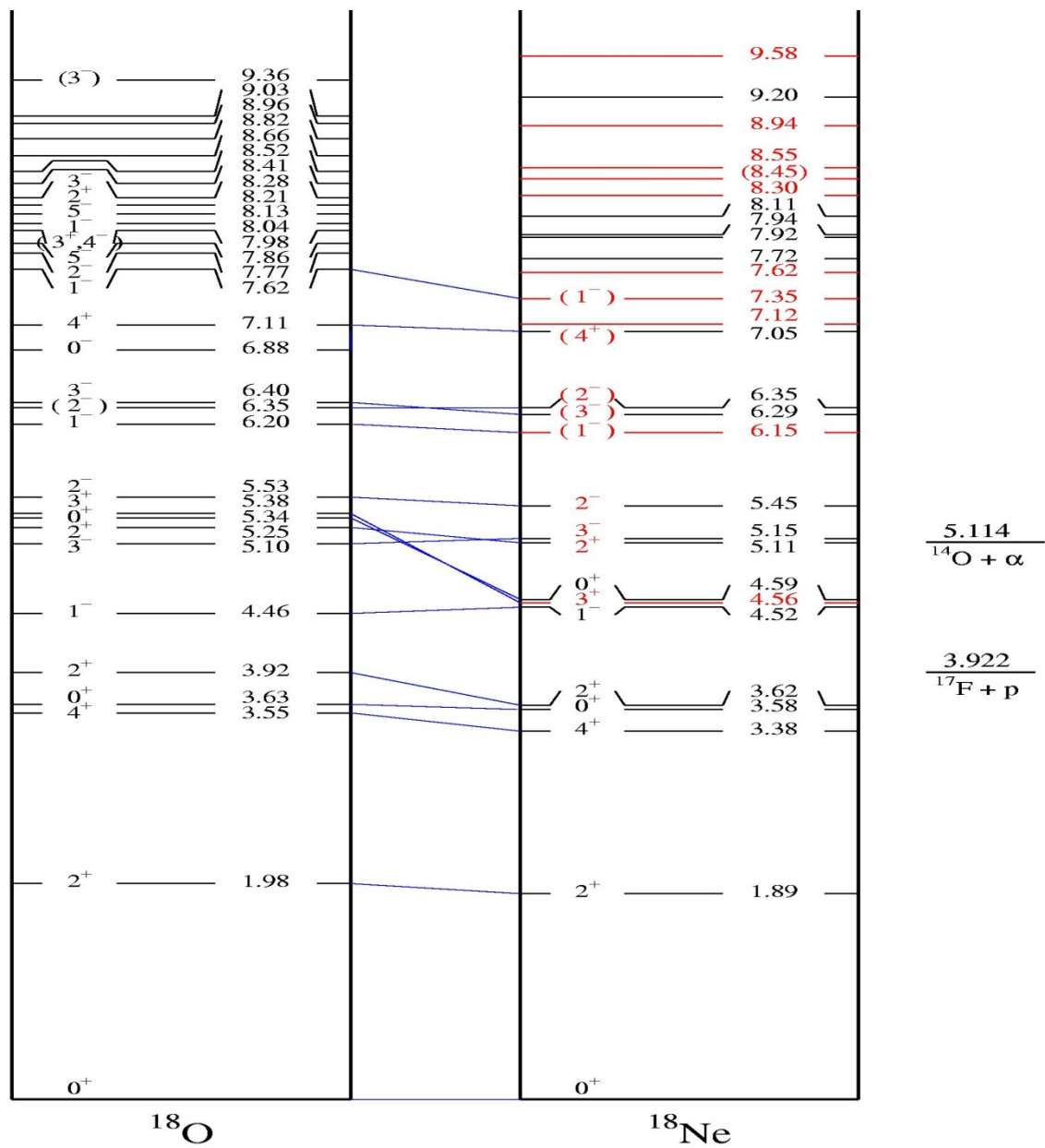
Magnifications: $M_X = 0.39, \quad M_Y = 2.9$



12C + 12C, Ebeam=80MeV, 10 Degrees, 14.8 kG







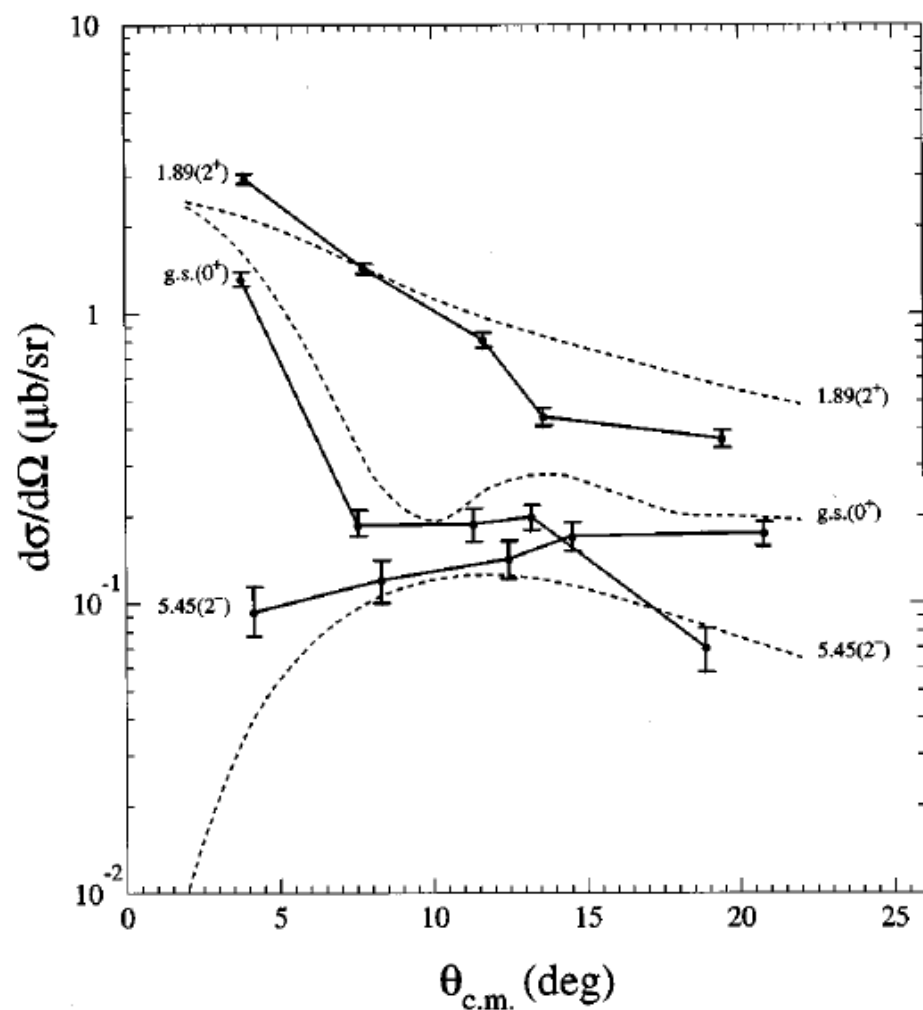


FIG. 8. Absolute Hauser-Feshbach statistical-model calculations compared to experimental angular distributions for ^{18}Ne states populated in the $^{12}\text{C}(^{12}\text{C}, ^6\text{He})^{18}\text{Ne}$ reaction at $E_{\text{lab}}=80.0$ MeV. The solid lines indicate experimental values; the dashed lines are model calculations.

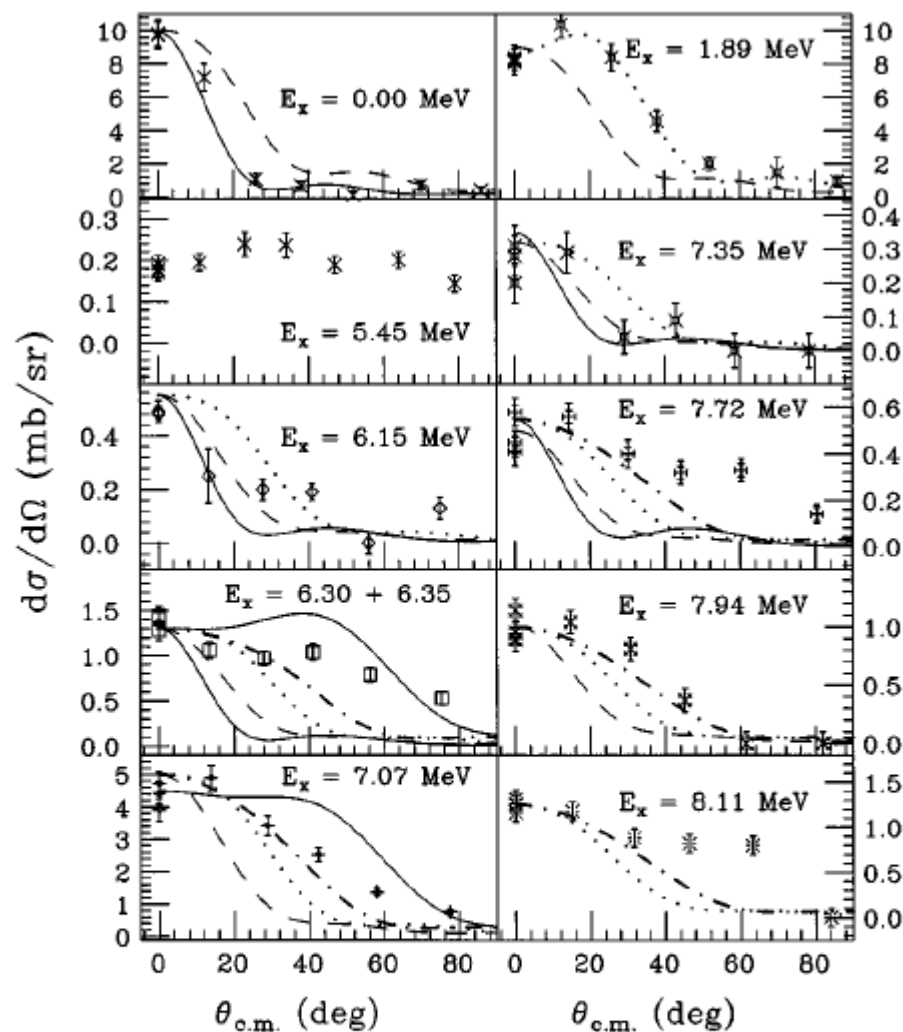


FIG. 5. $^{16}\text{O}(^3\text{He},n)$ angular distributions taken at $E_{^3\text{He}}=14.5$ MeV. The lines show DWBA calculations for different values of the transferred orbital angular momentum; solid line, $L=0$; dashed line, $L=1$; dotted line, $L=2$; dot-dashed line, $L=3$; solid line, $L=4$.

Books

- C.E. Rolfs and W.S. Rodney *Cauldrons in the Cosmos*
- The University of Chicago Press, 1988 (...the “Bible”)
- D.D. Clayton *Principles of stellar evolution and nucleosynthesis*
- The University of Chicago Press, 1983
- W.D. Arnett and J.W. Truran *Nucleosynthesis*
- The University of Chicago Press, 1968
- J. Audouze and S. Vauclair *An introduction to Nuclear Astrophysics*
- D. Reidel Publishing Company, Dordrecht, 1980
- E. Böhm–Vitense *Introduction to Stellar Astrophysics, vol. 3*
- Cambridge University Press, 1992