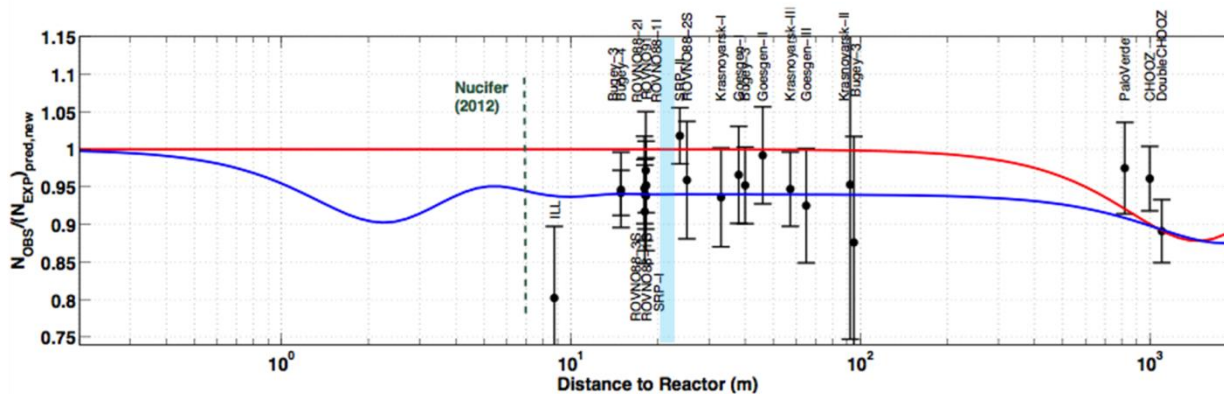


Remarks on Analysis and Systematics?

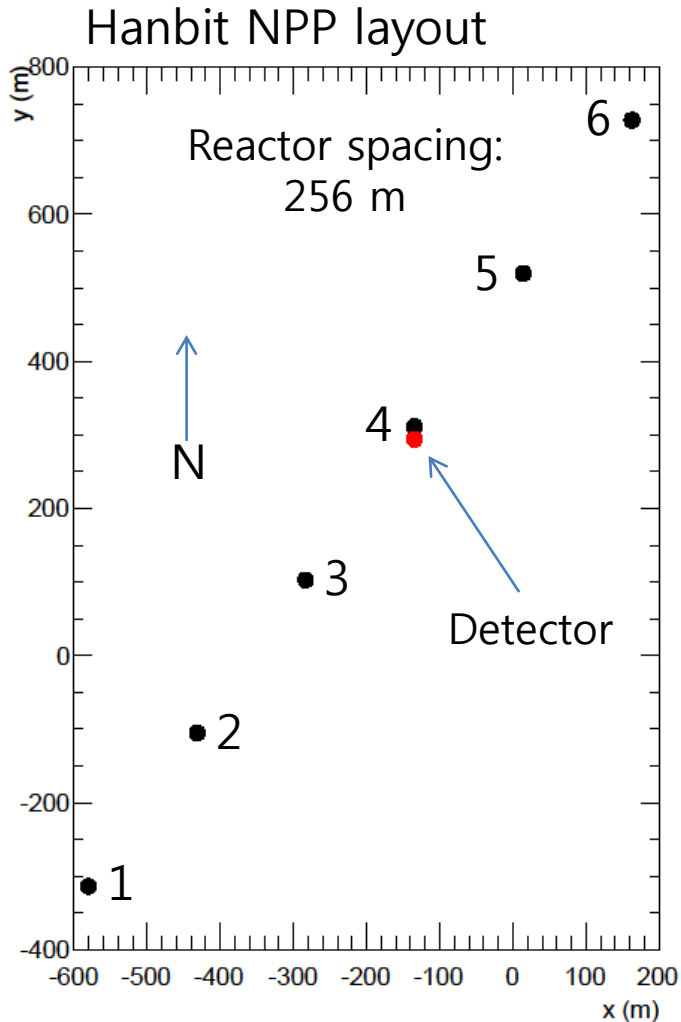
Hyunsoo Kim
Chonbuk National Univ.

- Because of using a single detector, the uncertainty in the neutrino flux (counts) degrades the sensitivity.
- The shape is of the concern.

$$\text{Prob}(\bar{\nu}_e \rightarrow \bar{\nu}_e) \approx 1 - \sin^2(2\theta_{ee}) \sin^2\left(\frac{1.27 \Delta m_{41}^2 L}{E_{\bar{\nu}_e}}\right)$$



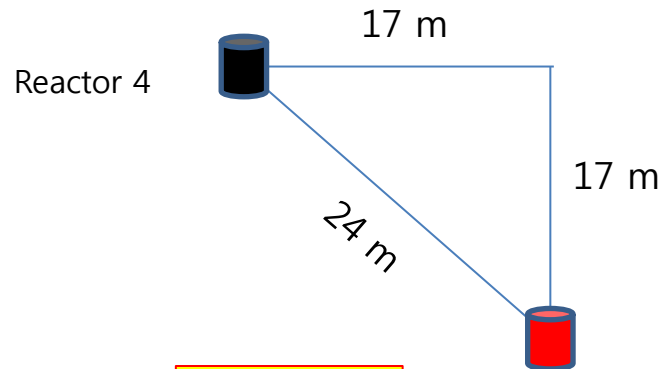
Calculation Setup



Reactor

Fuel assembly :
r = 1.56 m, h = 3.18 m

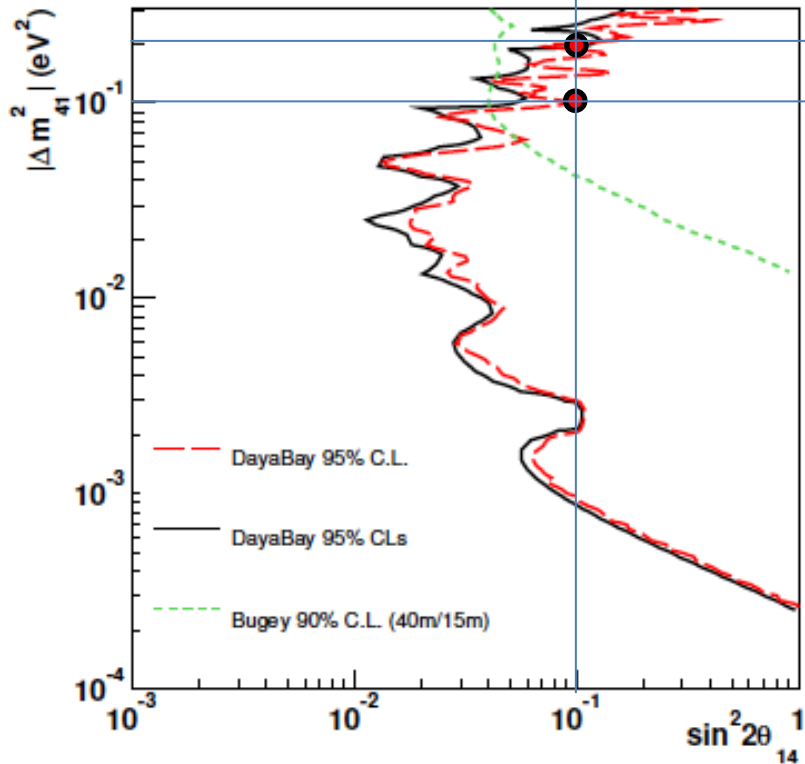
$P_{th} = 2.73$ GW
 ^{235}U : 0.6977
 ^{238}U : 0.0765
 ^{239}Pu : 0.1912
 ^{241}Pu : 0.0346



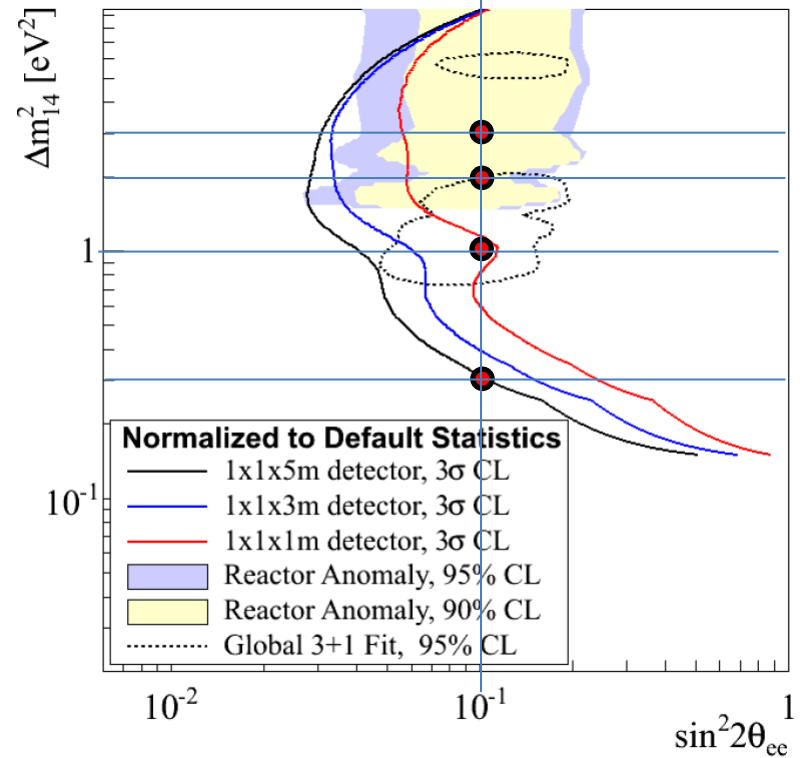
Detector

Active volume :
r = 0.5 m, h = 1.2 m
Energy resolution $\frac{\delta E}{E} = \frac{0.06}{\sqrt{E}}$
Acceptance: 60%
No boundary effects

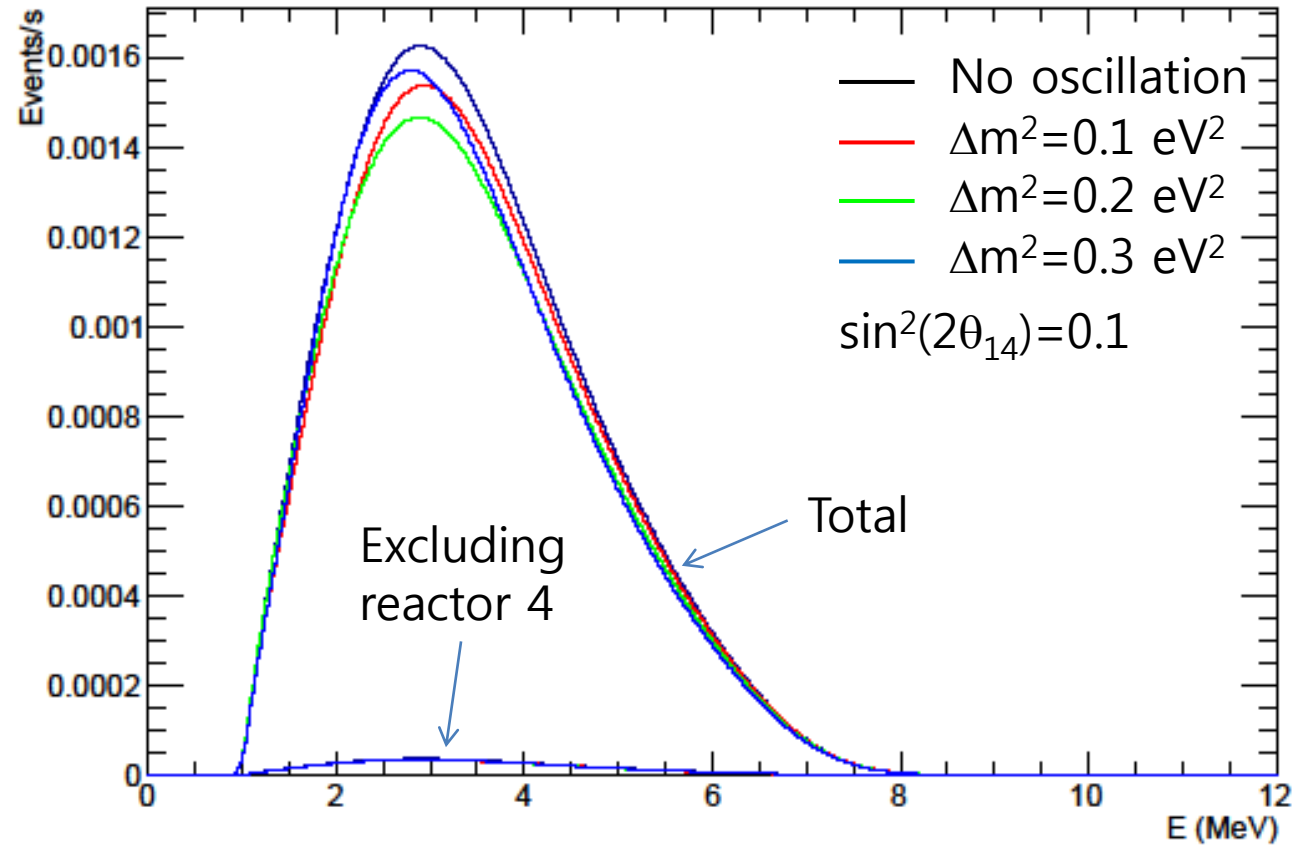
Daya-Bay 95% CL limits



arXiv:1212.2182v1 (feasibility study)

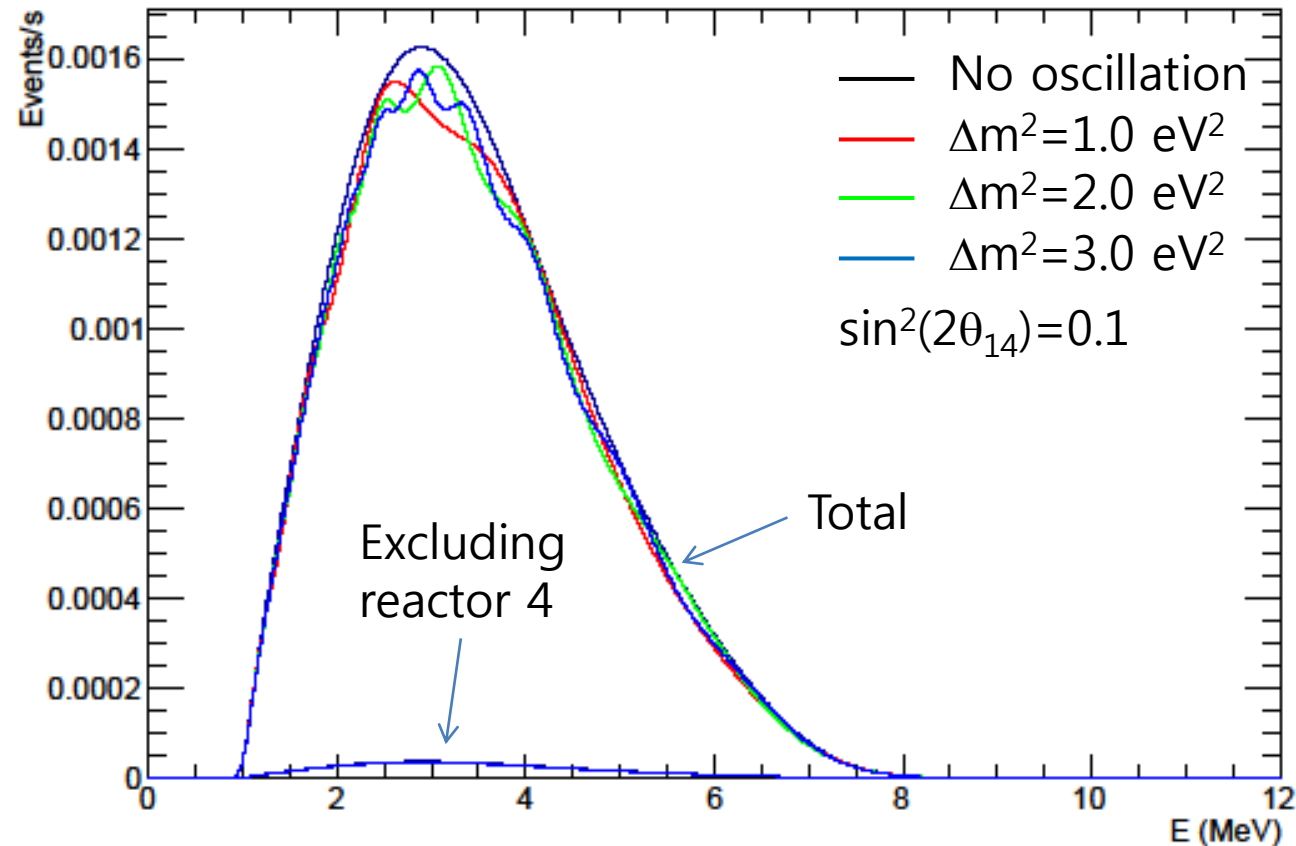


Energy Spectra



Spectra shifted
Low sensitivity if large
energy scale uncertainty
and large flux uncertainty.

Energy Spectra



Spectra distorted
Low sensitivity if large
energy scale uncertainty
and large flux uncertainty.

χ^2 Fit

From "Experimental Parameters for a Reactor Antineutrino Experiment at Very Short Baselines" (arXiv:1212.2182v1)

$$\chi^2 = \sum_{i,j} \frac{\left[\overset{\text{Data in the } i^{\text{th}} \text{ energy bin}}{M_{ij}} - (\alpha + \alpha_e^i + \alpha_r^j) \overset{\text{background in the } i^{\text{th}} \text{ energy bin}}{T_{ij}} - (1 + \alpha_b) B_{ij} \right]^2}{\underset{\text{signal in the } i^{\text{th}} \text{ energy bin}}{T_{ij}} + (\sigma_{b2b} B_{ij})^2} + \frac{\alpha^2}{\sigma^2} + \sum_j \left(\frac{\alpha_r^j}{\sigma_r} \right)^2 + \sum_i \left(\frac{\alpha_e^i}{\sigma_e^i} \right)^2 + \frac{\alpha_b^2}{\sigma_b^2}$$

Absolute normalisation
Bin-by-bin normalisation
background normalisation

Minimise with respect to nuisance parameters $\{\alpha, \alpha_e^i, \alpha_b\}$

χ^2 Fit Parameters

$$\chi^2 = \sum_{i,j} \frac{[M_{ij} - (\alpha + \alpha_e^i + \alpha_r^j)T_{ij} - (1 + \alpha_b)B_{ij}]^2}{T_{ij} + (\sigma_{b2b}B_{ij})^2} + \frac{\alpha^2}{\sigma^2} + \sum_j \left(\frac{\alpha_r^j}{\sigma_r}\right)^2 + \sum_i \left(\frac{\alpha_e^i}{\sigma_e}\right)^2 + \frac{\alpha_b^2}{\sigma_b^2}$$

$\frac{\alpha^2}{\sigma^2}$ In our case, floating α , i.e. σ =some large number...

$\left(\frac{\alpha_e^i}{\sigma_e}\right)^2$ Bin-by-bin uncertainties: very simple minded. not sure we can account for the correlations between systematics this way.

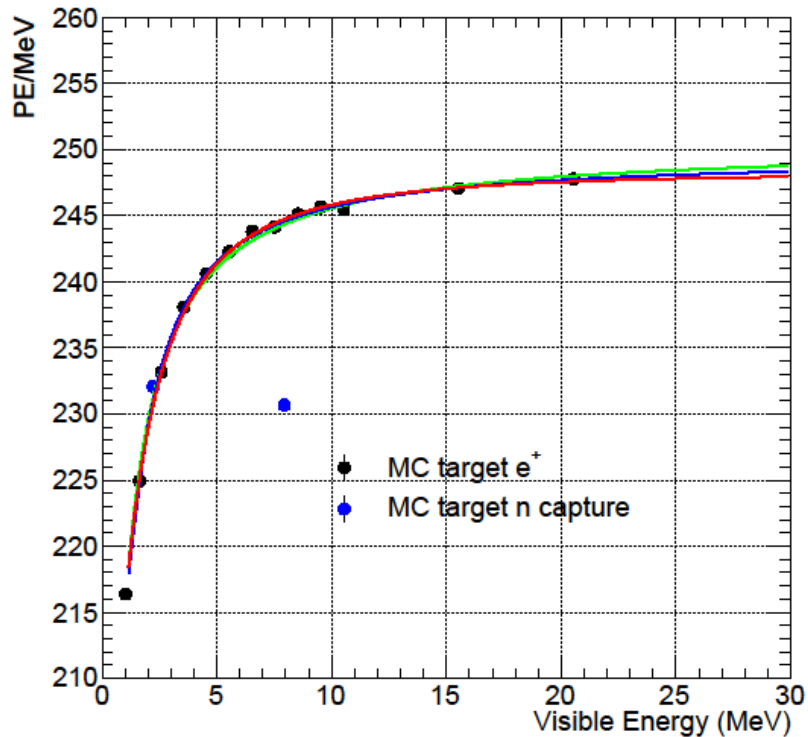
For example, energy scale uncertainty is correlated and contribution goes in to T_i as

$$\varepsilon_i \left. \frac{\partial T_i}{\partial \varepsilon_i} \right|_{\varepsilon_i=0} \quad \text{where } E' = (1 + \varepsilon)E$$

$\frac{\alpha_b^2}{\sigma_b^2}$ Completely bin-by-bin correlated background uncertainty. But not so simple in reality.

Energy Scale

positrons energy deposition: MC (RENO)



Need various calibration sources.
Photon and positron responses are different.

We need to tune the MC as well.

χ^2 Fit Parameters

χ^2 function once used at RENO

Spectral Fit

"current" χ^2 function

$$\chi^2 = \min_{\{\omega\}} \left\{ \sum_{d=N,F} \left[\sum_{i=1}^{N_b} \left(\frac{O_i^d - N_i^d}{U_i^d} \right)^2 + (\beta^d)^2 + \left(\frac{\delta^d}{\sigma_d^d} \right)^2 + \left(\frac{\epsilon^d}{\sigma_{cml}^d} \right)^2 \right] + \left(\frac{\alpha}{\sigma_a} \right)^2 + \sum_{i=1}^{N_b} \sum_{j=1}^{n_{iso}} \left(\frac{\gamma_{ij}}{\sigma_{shape_{ij}}} \right)^2 + \sum_{r=1}^{N_c} \left(\frac{\xi_r}{\sigma_{cfl}} \right)^2 \right\} + F_i^d$$

5 MeV component

where

$$N_i^d = (1 + a + e^d) \sum_{r=1}^{N_c} (1 + f_r) T_{ir}^d + g^d M_i^d + b^d \cdot \sigma(B_i^d)$$

$$T_i^d \simeq T_i^d(g^d = 0) + g^d M_i^d, \quad M_i^d = \left. \frac{dT_i^d}{dg^d} \right|_{g^d=0}, \quad T_{ir}^d = \sum_{j=1}^{n_{iso}} (1 + \gamma_{ij}) T_{ir,j}^d$$

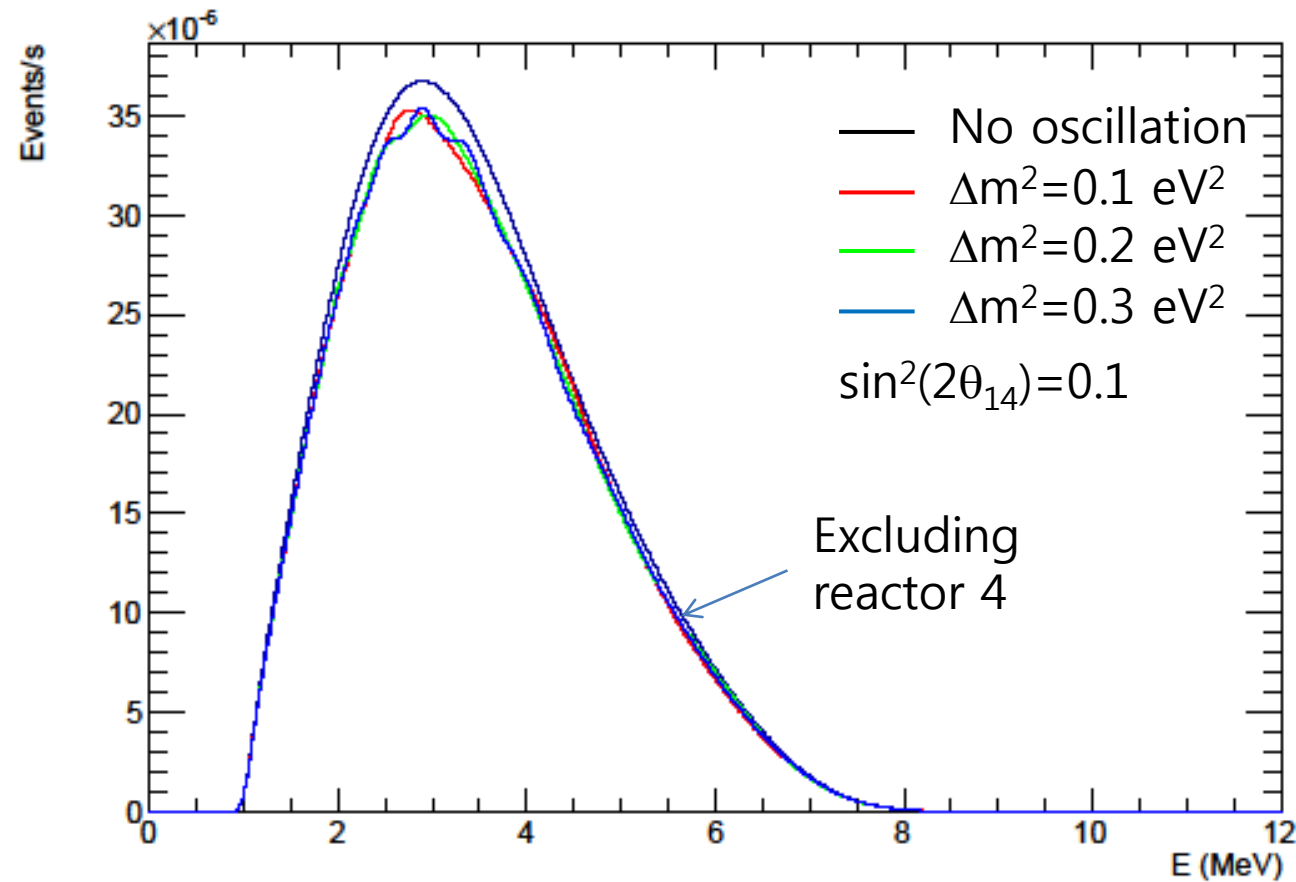
and

$$U_i^d = \sqrt{O_i^d + B_{ij}^d + (b^d \cdot \sigma(B_i^d))} \quad (?)$$

Summary

- There's a long and arduous way ahead.

Energy Spectra



Spectra distorted

Energy Spectra

