Aligned natural inflation: Oscillations in primordial power spectrum

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Suppose we achieve *super-Planckian axion scale* through *alignment mechanism*

$$V(\phi) = \Lambda_1^4 \left[1 - \cos \frac{\phi}{f_{\text{eff}}} \right] + \Lambda_2^4 \left[1 - \cos \left(\frac{\phi}{f} + \delta \right) \right]$$

$$\Lambda_2 \ll \Lambda_1$$

What would be the *observational consequences* of this model?

 $f \ll f_{\text{eff}}$

Natural Inflation

$$V(\phi) = \Lambda^4 \left[1 - \cos \frac{\phi}{f} \right]$$

The amplitude can determined by observed temperature fluctuation

The axion scale is required to be super-Planckian

Axion scale

Especially in string theory,

$$f \simeq \frac{g^2}{8\pi^2} M_P \sim 10^{16} \text{GeV}$$

Axion scale is normally *sub-Planckian* in weak coupling limit

Proposals

Let's introduce extra axions!

Axion alignment

Kim, Nilles, Peloso '04

N-flation

Dimopolous '05

Other variations

Choi, HK, Yun '14

Higaki, Takahashi '14

Tye, Wong '14

Ben-Dayan, Pedro, Westphal '14

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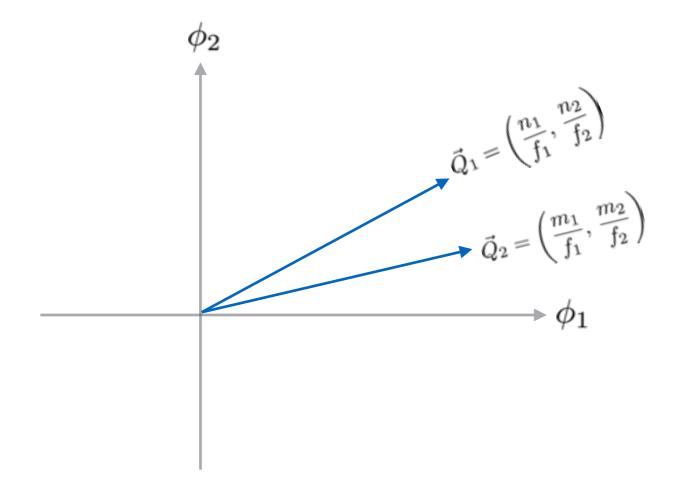
Tye, Wong '14

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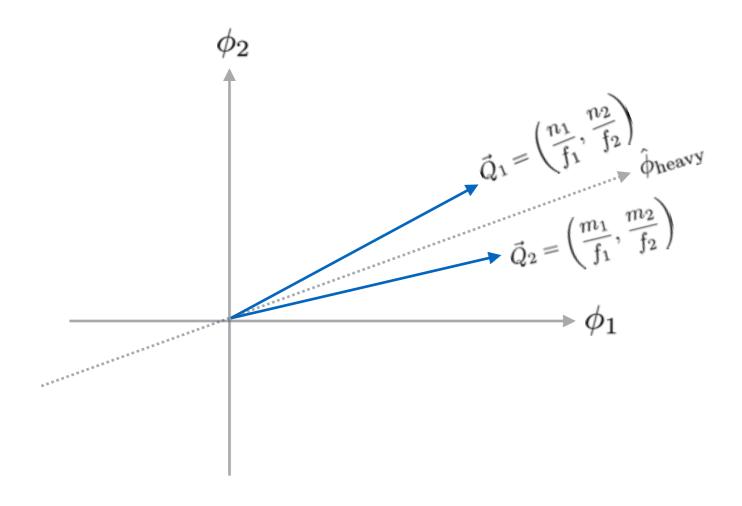
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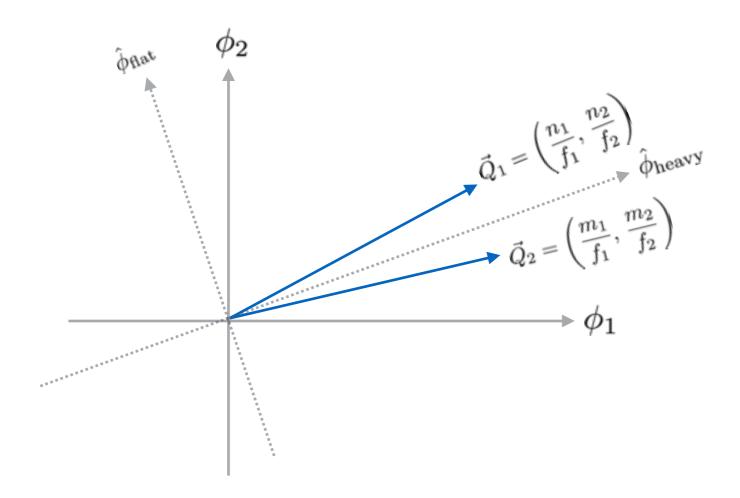
$$V = \Lambda^4 e^{-S_1} \left[1 - \cos \left(\frac{n_1}{f_1} \phi_1 + \frac{n_2}{f_2} \phi_2 \right) \right] + \Lambda^4 e^{-S_2} \left[1 - \cos \left(\frac{m_1}{f_1} \phi_1 + \frac{m_2}{f_2} \phi_2 \right) \right]$$



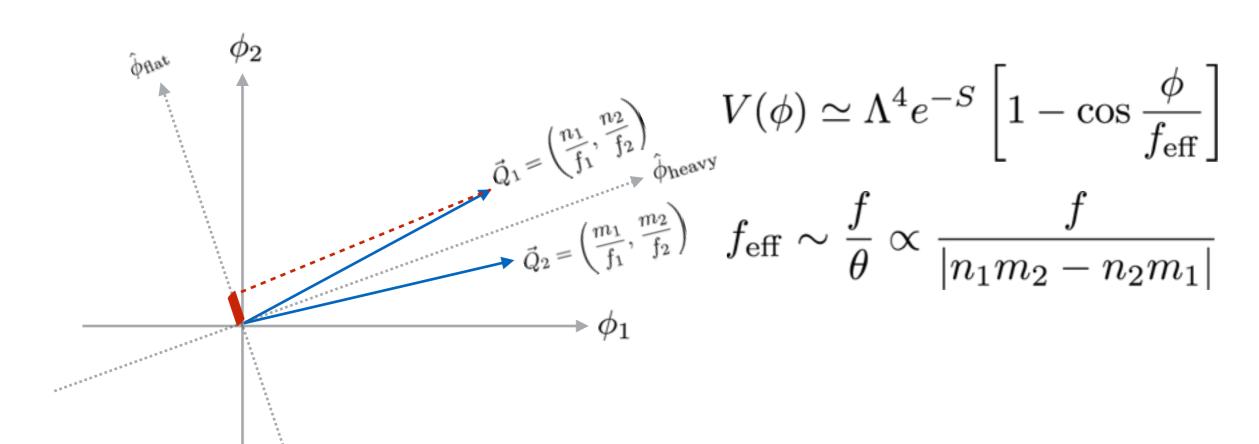
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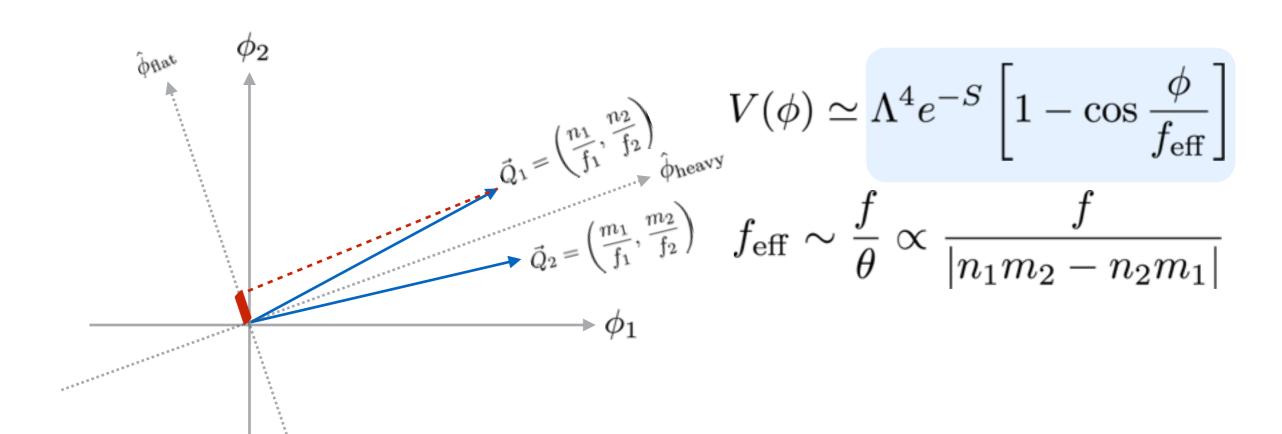
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Consider a theory of *U(1)* gauge boson + gravity

$$q \ge \frac{m}{M_P}$$

Consider a theory of **axion** + **gravity**

$$\frac{1}{f} \ge \frac{S_{\text{inst}}}{M_P} \quad \Leftrightarrow \quad f \le \frac{M_P}{S_{\text{inst}}}$$

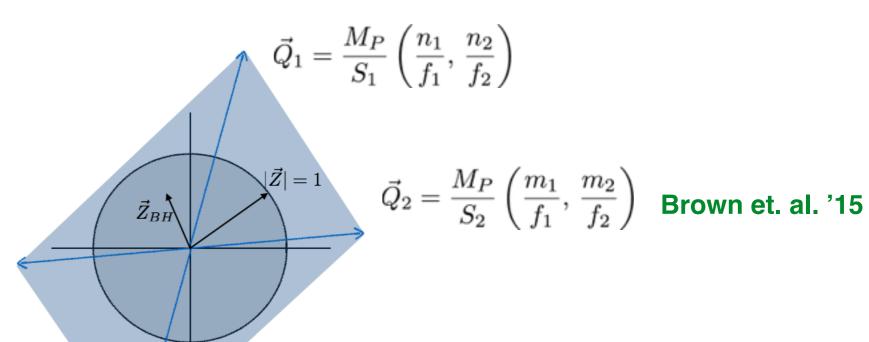
axion scale is bounded from above!

How about **N axions?**

Convex Hull condition

Cheung et. al. '14

$$V = \Lambda^4 e^{-S_1} \left[1 - \cos \left(\frac{n_1}{f_1} \phi_1 + \frac{n_2}{f_2} \phi_2 \right) \right] + \Lambda^4 e^{-S_2} \left[1 - \cos \left(\frac{m_1}{f_1} \phi_1 + \frac{m_2}{f_2} \phi_2 \right) \right]$$



$$f_{\rm eff} < M_P$$

Rudelius '15 Montero et. al. '15

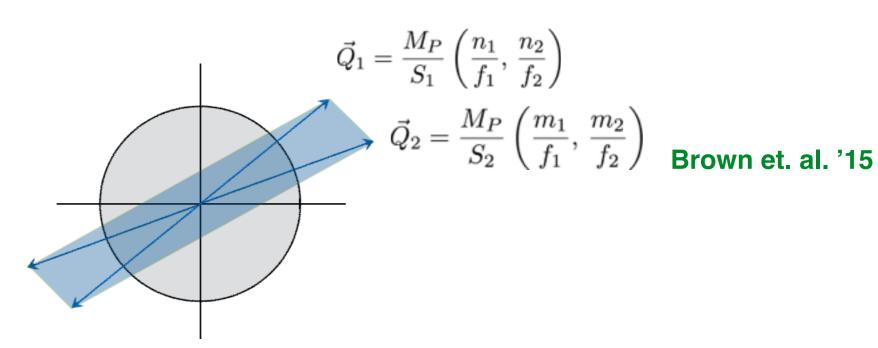
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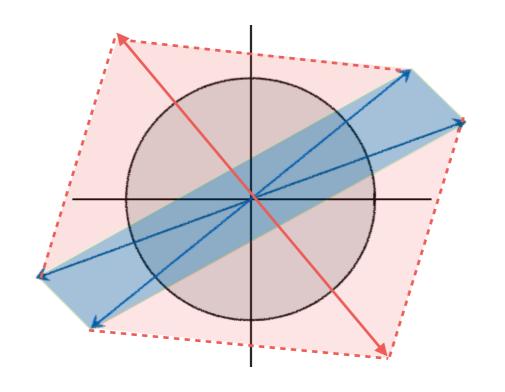
Alignment case

$$f_{\rm eff} < M_P$$

Rudelius '15 Montero et. al. '15

...

One way of circumventing this is to add extra instanton!



$$f \ll f_{\text{eff}}$$

 $S \ll S_{\text{large}}$

$$V(\phi) = \Lambda^4 e^{-S} \left[1 - \cos \frac{\phi}{f_{\text{eff}}} \right] + \Lambda^4 e^{-S_{\text{large}}} \left[1 - \cos \left(\frac{\phi}{f} + \delta \right) \right]$$

Back to the original question

If the modulation is natural consequence of alignment mechanism,

$$V(\phi) = \Lambda_1^4 \left[1 - \cos \frac{\phi}{f_{\text{eff}}} \right] + \Lambda_2^4 \left[1 - \cos \left(\frac{\phi}{f} + \delta \right) \right]$$
$$f \ll f_{\text{eff}}$$

How do they change the shape of primordial power spectrum?

How could they be constrained by the recent observational data?

Former study: Axion monodromy

Axion monodromy potenial with cosine modulation

$$V(\phi) = \mu^{3}\phi + \Lambda^{4}\cos\left(\frac{\phi}{f} + \delta\right)$$

induces oscillations in primordial power spectrum

$$\mathcal{P}_{\mathcal{R}}(k) = A_{\mathcal{R}*} \left(\frac{k}{k_*}\right)^{n_s - 1} \left[1 + \delta n_s \cos\left(\frac{\phi_k}{f} + \alpha\right) \right]$$

Natural inflation with modulations

Natural inflation with cosine modulation

$$V(\phi) = \Lambda_1^4 \left[1 - \cos \frac{\phi}{f_{\text{eff}}} \right] + \Lambda_2^4 \left[1 - \cos \left(\frac{\phi}{f} + \delta \right) \right]$$

also induces oscillations in primordial power spectrum

$$\mathcal{P}_{\mathcal{R}}(k) = A_{\mathcal{R}*} \left(\frac{k}{k_*}\right)^{n_s - 1} \left[1 + \delta n_s \cos\left(\frac{\phi_k}{f} + \alpha\right) \right]$$

Natural inflation with modulations

Natural inflation with cosine modulation

$$\delta n_s = \left(\frac{\Lambda_2^4 f_{\text{eff}}}{\Lambda_1^4 f}\right) \frac{3\sqrt{2\pi\gamma \coth\frac{\pi}{2\gamma}}}{\sin(\phi_*/f_{\text{eff}})\sqrt{(1+\frac{3}{2}\gamma^2/f_{\text{eff}}^2)^2 + (3\gamma)^2}} \quad \gamma = f_{\text{eff}} f \tan\frac{\phi_*}{2f_{\text{eff}}}$$

also induces oscillations

$$\mathcal{P}_{\mathcal{R}}(k) = A_{\mathcal{R}*} \left(\frac{k}{k_*}\right)^{n_s - 1} \left[1 + \delta n_s \cos\left(\frac{\phi_k}{f} + \alpha\right) \right]$$

Oscillations in power spectrum

For ordinary slow-roll inflation, power spectrum

$$\mathcal{P}_{\mathcal{R}}(k) = A_* \left(\frac{k}{k_*}\right)^{(n_s - 1) + \frac{1}{2} \frac{dn_s}{d \ln k} \ln(k/k_*) + \frac{1}{6} \frac{d^2 n_s}{d \ln k^2} \ln^2(k/k_*) + \cdots}$$

is described by a set of parameters

$$A_*$$
 n_s $\frac{dn_s}{d\ln k}$ $\frac{d^2n_s}{d\ln k^2}$ \cdots

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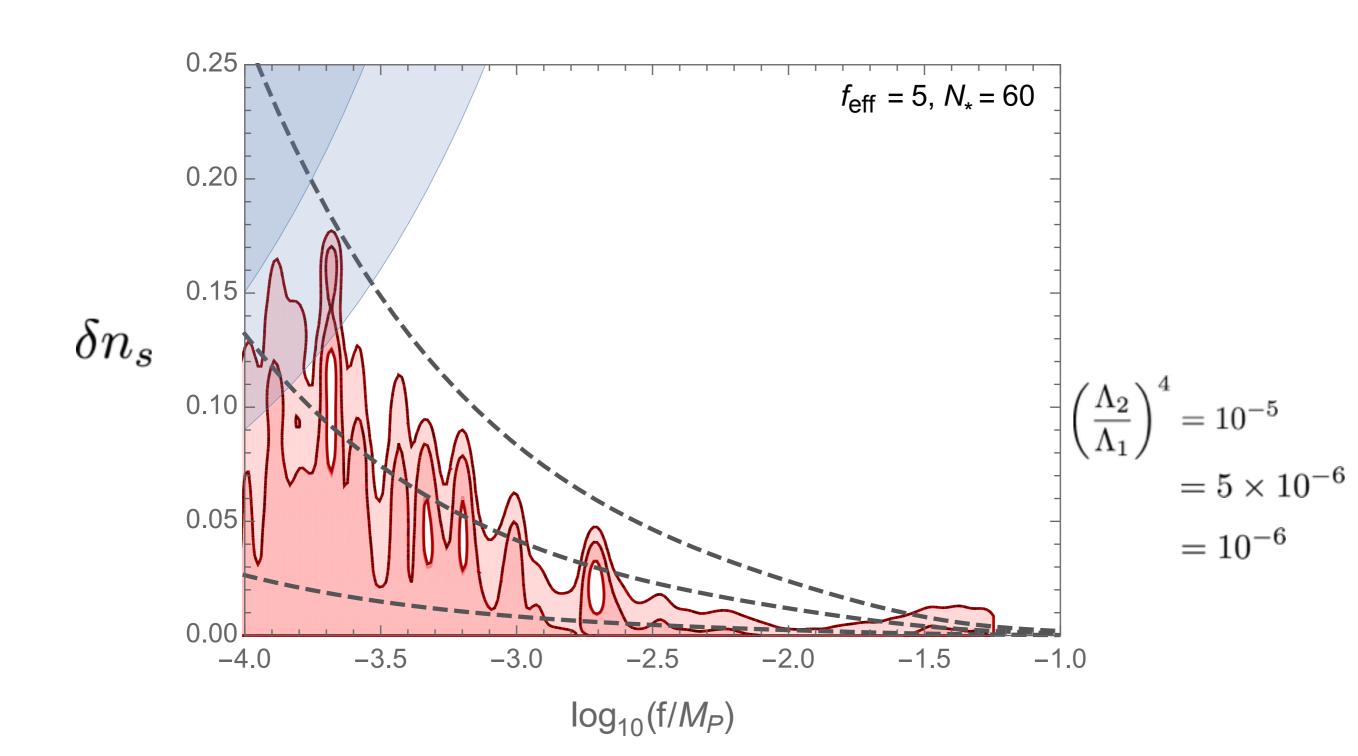
In this scenario, power spectrum looks like

$$\mathcal{P}_{\mathcal{R}}(k) = A_{\mathcal{R}*} \left(\frac{k}{k_*}\right)^{n_s - 1} \left[1 + \frac{\delta n_s}{f} \cos\left(\frac{\phi_k}{f} + \alpha\right) \right]$$

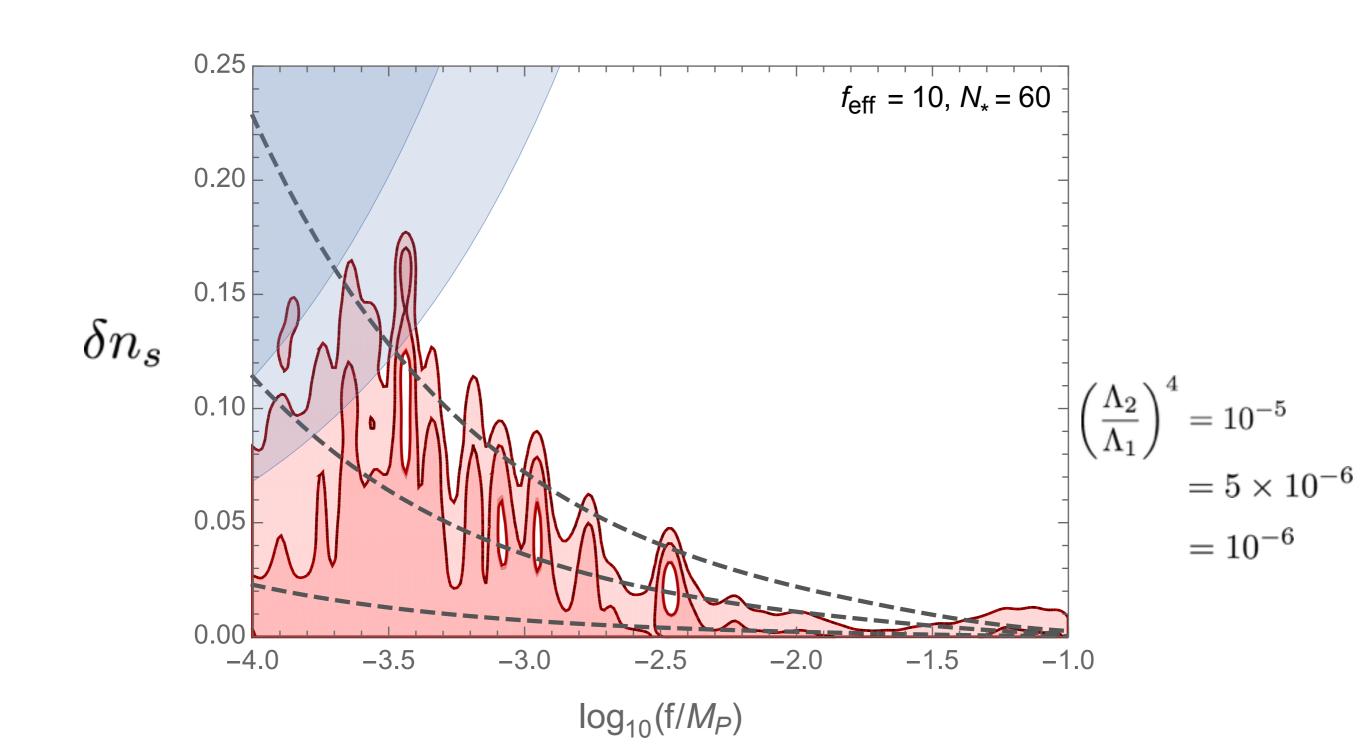
the standard template $(A_{\mathcal{R}*}, n_s)$

does not capture characteristic features of this scenario.

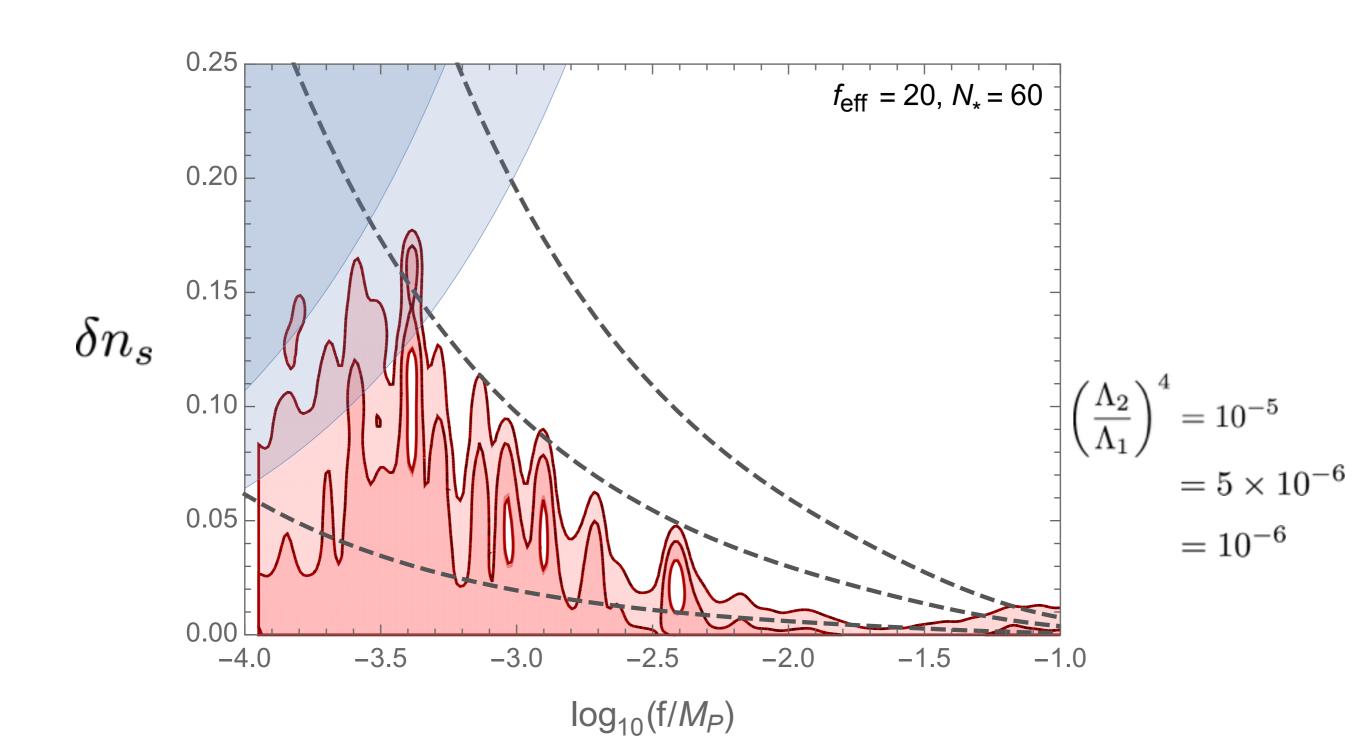
Constraints on oscillations



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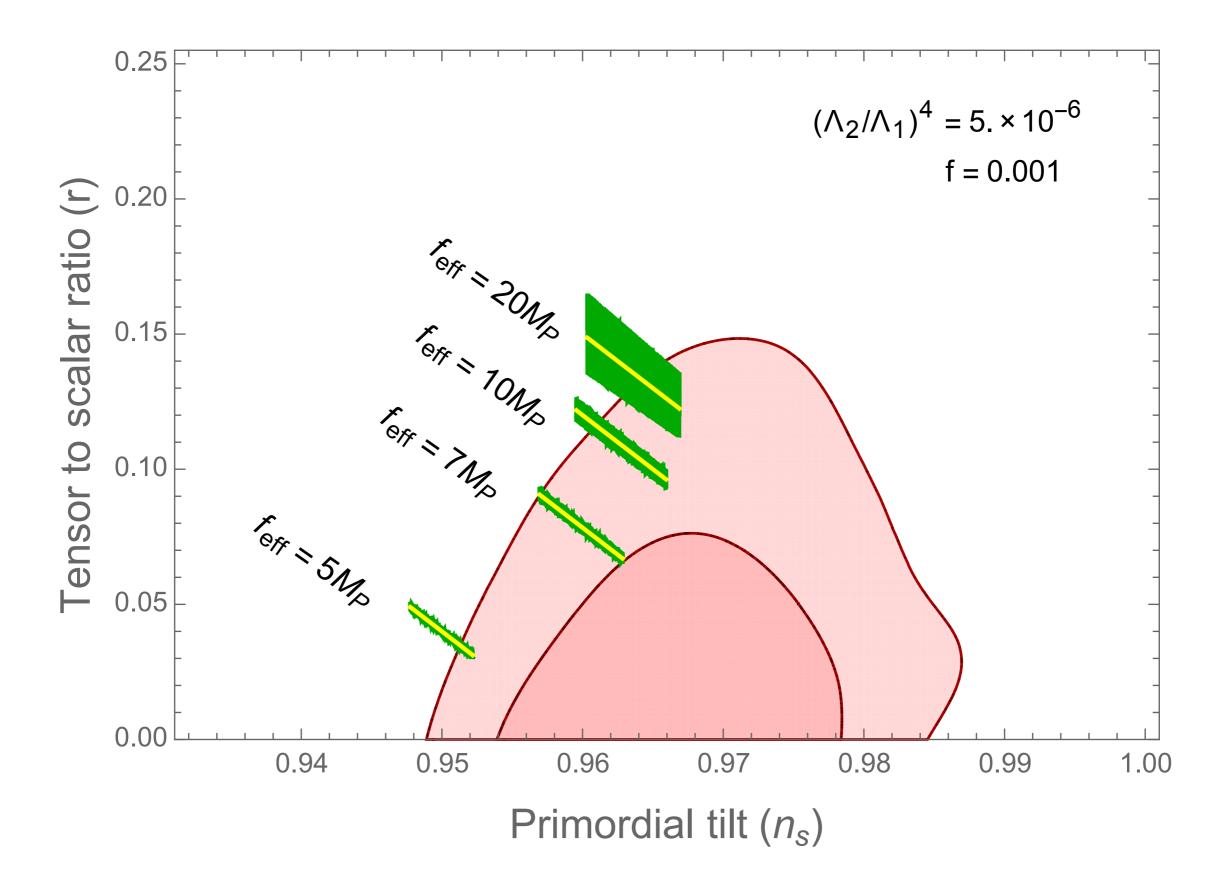


tensor-to-scalar ratio?

tensor-to-scalar ratio can change, but only by

$$r = \frac{\mathcal{P}_t}{\mathcal{P}_{\mathcal{R}}} = \frac{r_0}{1 + \delta n_s \cos(\phi_k/f)}$$

as the amplitude of oscillation is constrained by observation, only ~10% change is possible



Summary

$$V(\phi) = \Lambda_1^4 \left[1 - \cos \frac{\phi}{f_{\text{eff}}} \right] + \Lambda_2^4 \left[1 - \cos \left(\frac{\phi}{f} + \delta \right) \right]$$

Observational signature?

Oscillations!

Any contraints?

Amplitude of oscillation leads to $(\Lambda_2/\Lambda_1)^4 \lesssim 5 \times 10^{-6}$

tensor-to-scalar ratio can change by $\sim 10\%$