

A Particle Probing Thermodynamics in Rotating AdS Black Hole

Based on arXiv:1509.06691 "Cosmic Censorship of Rotating Anti-de Sitter Black Hole with a Probe"

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Motivation

- We test the validity of cosmic censorship in a 4-dimensional rotating AdS black hole through a particle absorption.
- The particle conserved quantities are redefined to satisfy the 1st and 2nd law of thermodynamics of the black hole.
- We can show the validity of cosmic censorship in the extremal black hole.

4-dimensional rotating AdS black hole

- The metric:

$$ds^2 = -\frac{\Delta_r}{\rho^2} \left(dt - \frac{a \sin^2 \theta}{\Xi} d\phi \right)^2 + \frac{\rho^2}{\Delta_r} dr^2 + \frac{\rho^2}{\Delta_\theta} d\theta^2 + \frac{\Delta_\theta \sin^2 \theta}{\rho^2} \left(a dt - \frac{r^2 + a^2}{\Xi} d\phi \right)^2,$$
$$\rho^2 = r^2 + a^2 \cos^2 \theta, \quad \Delta_r = (r^2 + a^2) \left(1 + \frac{r^2}{\ell^2} \right) - 2Mr, \quad \Delta_\theta = 1 - \frac{a^2}{\ell^2} \cos^2 \theta, \quad \Xi = 1 - \frac{a^2}{\ell^2}.$$

- AdS radius, mass parameter, spin parameter.
- Horizon area and Bekenstein-Hawking entropy:

$$\mathcal{A}_h = \frac{4\pi (r_h^2 + a^2)}{\Xi}, \quad S_{BH} = \frac{1}{4} \mathcal{A}_h = \frac{\pi (r_h^2 + a^2)}{\Xi},$$

The 1st law of thermodynamics in AdS BH

- The rotating velocity at the horizon: $\Omega_h = \frac{a\Xi}{r_h^2 + a^2}$.
- But, $dM_B \neq T_H dS_{BH} + \Omega_h dJ_B$
- The rotating velocity at $r \rightarrow \infty$: $\Omega_\infty = -\frac{a}{\ell^2}$.
- Redefine the rotating velocity at the horizon: $\Omega = \Omega_h - \Omega_\infty = \frac{a \left(1 + \frac{r_h^2}{\ell^2}\right)}{r_h^2 + a^2}$
- Under this redefinition, the 1st law of thermodynamics:

$$dM_B = T_H dS_{BH} + \Omega dJ_B$$

- The black hole mass and angular momentum: $M_B = \frac{M}{\Xi^2}$, $J_B = \frac{aM}{\Xi^2}$,
M. M. Caldarelli, G. Cognola, D. Klemm, Class. Quant. Grav. 17, 399 (2000)
G. W. Gibbons, M. J. Perry, C. N. Pope, Class. Quant. Grav. 22, 1503 (2005)

The particle equations of motions

- The black hole properties can be changed by a particle quantities.
- Using Hamilton-Jacobi method:

$$\mathcal{H} = \frac{1}{2}g^{\mu\nu} p_\mu p_\nu \quad S = \frac{1}{2}m^2\lambda - Et + L\phi + S_r(r) + S_\theta(\theta)$$

- Equations of motions with a separate variable:

$$p^r = \dot{r} = \frac{\Delta_r}{\rho^2} \sqrt{R(r)}, \quad p^\theta = \dot{\theta} = \frac{\Delta_\theta}{\rho^2} \sqrt{\Theta(\theta)}.$$

$$R(r) = \frac{\mathcal{K}}{\Delta_r} + \frac{1}{\Delta_r^2} \left(a\Xi L - (r^2 + a^2) E \right)^2 - \frac{m^2 r^2}{\Delta_r},$$

$$\Theta(\theta) = -\frac{\mathcal{K}}{\Delta_\theta} - \frac{1}{\Delta_\theta^2} \left(\Xi L \csc \theta - aE \sin \theta \right)^2 - \frac{m^2 a^2 \cos^2 \theta}{\Delta_\theta}.$$

The particle absorption

- The particle energy is obtained for given location and momenta:

$$\alpha E^2 + 2\beta E + \gamma = 0,$$

$$\alpha = \frac{(r^2 + a^2)^2}{\Delta_r} - \frac{a^2 \sin^2 \theta}{\Delta_\theta}, \quad \beta = -\frac{(r^2 + a^2)(aL\Xi)}{\Delta_r} + \frac{aL\Xi}{\Delta_\theta}, \quad \gamma = -\frac{(p^r)^2 \rho^4 - a^2 L^2 \Xi^2}{\Delta_r} - \frac{(p^\theta)^2 \rho^4 + L^2 \Xi^2 \csc^2 \theta}{\Delta_\theta} - m^2 \rho^2.$$

- Solution at the horizon:

$$E_h = \frac{a\Xi}{r_h^2 + a^2} L + \frac{\rho_h^2}{r_h^2 + a^2} |p^r|.$$

- However, this equation is not consistent with thermodynamic laws.

- This can be resolved by energy redefinition.

- The particle energy moving φ -plane at the boundary: $E_\infty = -\frac{a}{\ell^2} L < 0$

- To avoid negative energy:

$$E = E_h - E_\infty = \frac{a \left(1 + \frac{r_h^2}{\ell^2} \right)}{r_h^2 + a^2} L + \frac{\rho_h^2}{r_h^2 + a^2} |p^r|.$$

- The black hole infinitesimally changes by the particle:

$$E = \delta M_B, \quad L = \delta J_B,$$

- Under the changes, the black hole entropy always increases.

$$\delta S_{BH} = \frac{4\pi\rho_h^2}{\dot{D}} |p^r| \geq 0$$

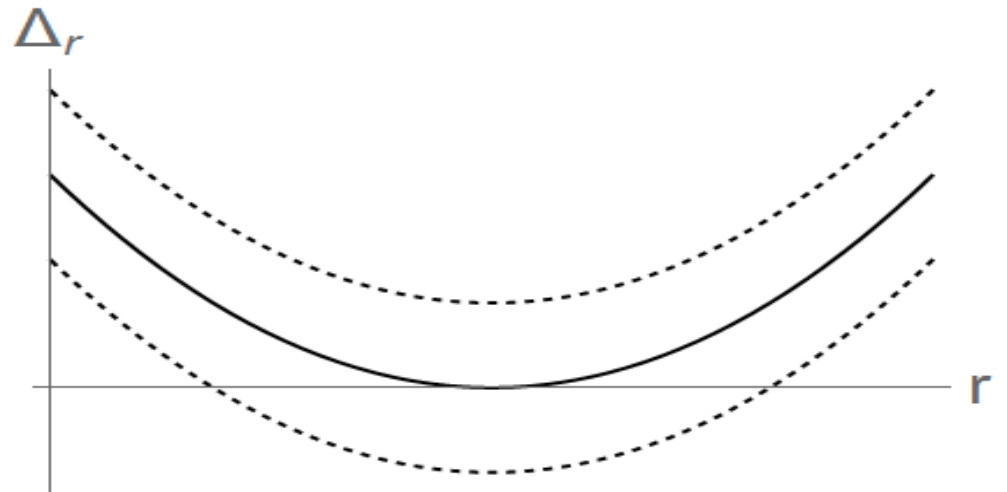
- This satisfies the 2nd law of thermodynamics.
- Using this relation, the redefined particle energy becomes:

$$\delta M_B = \Omega \delta J_B + T_H \delta S_{BH}$$

- This is the 1st law of thermodynamics.
- The particle is now consistent with the laws of thermodynamics.

Testing the validity of cosmic censorship

- We will investigate whether the extremal black hole can be overspun through particle absorption.
- If the horizon disappears, the black hole becomes a naked singularity, and cosmic censorship is invalid.
- The black hole horizon only depends on the function Δ_r .
- The extremal black hole is changed as $(M, J) \rightarrow (M + \delta M, J + \delta J)$



- The location of minimum point: $\delta r_e = -\frac{\dot{D}_M P_J + \dot{D}_J L}{\ddot{D}} - \frac{P_P}{\ddot{D}} |p^r|$, $\ddot{D} = \frac{\partial \dot{D}}{\partial r_h} = 2 \left(1 + \frac{r_h^2}{\ell^2}\right) + \frac{8r_h^2}{\ell^2} + \frac{2(r_h^2 + a^2)}{\ell^2}$,
 $\dot{D}_M = \frac{\partial \dot{D}}{\partial M_B} = -\frac{8a^2 \Xi}{\ell^2} - 2\Xi^2 - \frac{4a^2 r_h \Xi^2}{M \ell^2}$,
 $\dot{D}_J = \frac{\partial \dot{D}}{\partial J_B} = \frac{8a \Xi}{\ell^2} + \frac{4a r_h \Xi^2}{M \ell^2}$.
- The function Δ_r : $\Delta_r(r_h + \delta r_e) = D_M P_P |p^r|$, $D_M = \frac{\partial \Delta_h}{\partial M_B} = -\frac{8a^2 r_h \Xi}{\ell^2} - 2r_h \Xi^2 - \frac{2a^2 \left(1 + \frac{r_h^2}{\ell^2}\right) \Xi^2}{M}$,
 ≤ 0 $D_J = \frac{\partial \Delta_h}{\partial J_B} = -\frac{8a r_h \Xi}{\ell^2} + \frac{2a \left(1 + \frac{r_h^2}{\ell^2}\right) \Xi^2}{M}$.
- The black hole mass increases more than the angular momentum of the black hole, and the black hole becomes non-extremal one.
- Thus, cosmic censorship is valid under the particle absorption.

Summary

- We construct particle energy equation to satisfy the 1st and 2nd laws of thermodynamics using energy regularization.
- In the particle absorption, the extremal black hole mass increases more than the angular momentum of the black hole.
- It becomes non-extremal one, and the horizon still exists.
- Therefore, cosmic censorship is valid.

THANK YOU!