



Dark Matter Searches at the CMS

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On behalf of the CMS collaboration

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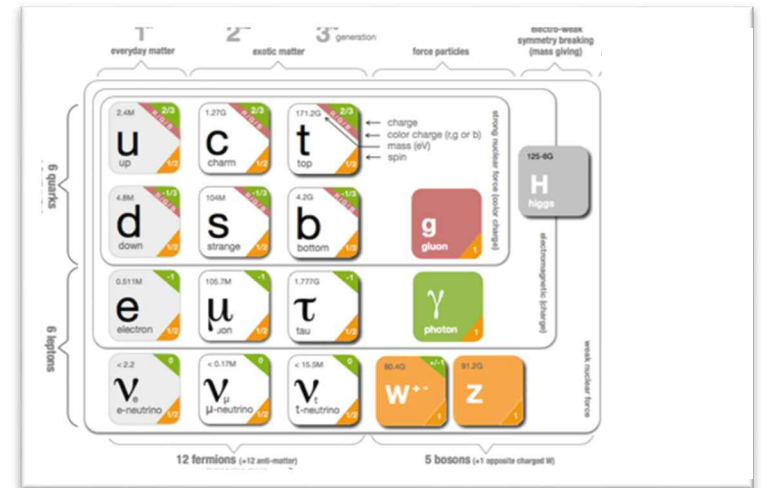
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**** SUSY searches are not included in this talk due to the time limit.**

Introduction

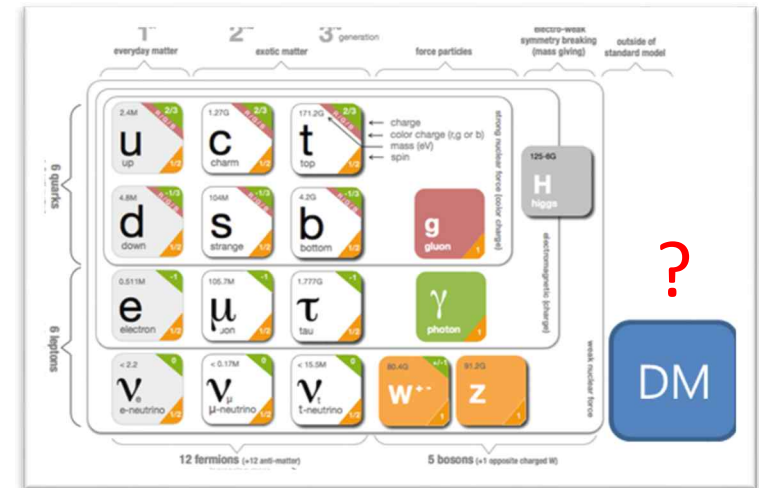
- The discovery of the Higgs particle implies the Standard Model(SM) describes nature very well.
- However, the incompleteness of the SM may hint the bigger quantum symmetrized universe.
- The Higgs discovery naturally leads to attempt to understand whether new phenomena beyond the SM exist or not.
 - The beyond SM like SUSY, Extra Dimensions and many extended SMs.



- Moreover, the dark matter(DM) is one of the biggest mysteries in understanding of universe after the discovery of Higgs.
- The discovery of a dark matter is very important since the existence of dark matter has been certified by astronomic observations.
- The LHC at CERN with new energy has a great chance to discover a dark matter or probe unknown nature.
- Searches of a dark matter (and unknown phenomena) at the LHC would be one of the most exciting moments enhancing mankind knowledge in the Universe !

Introduction

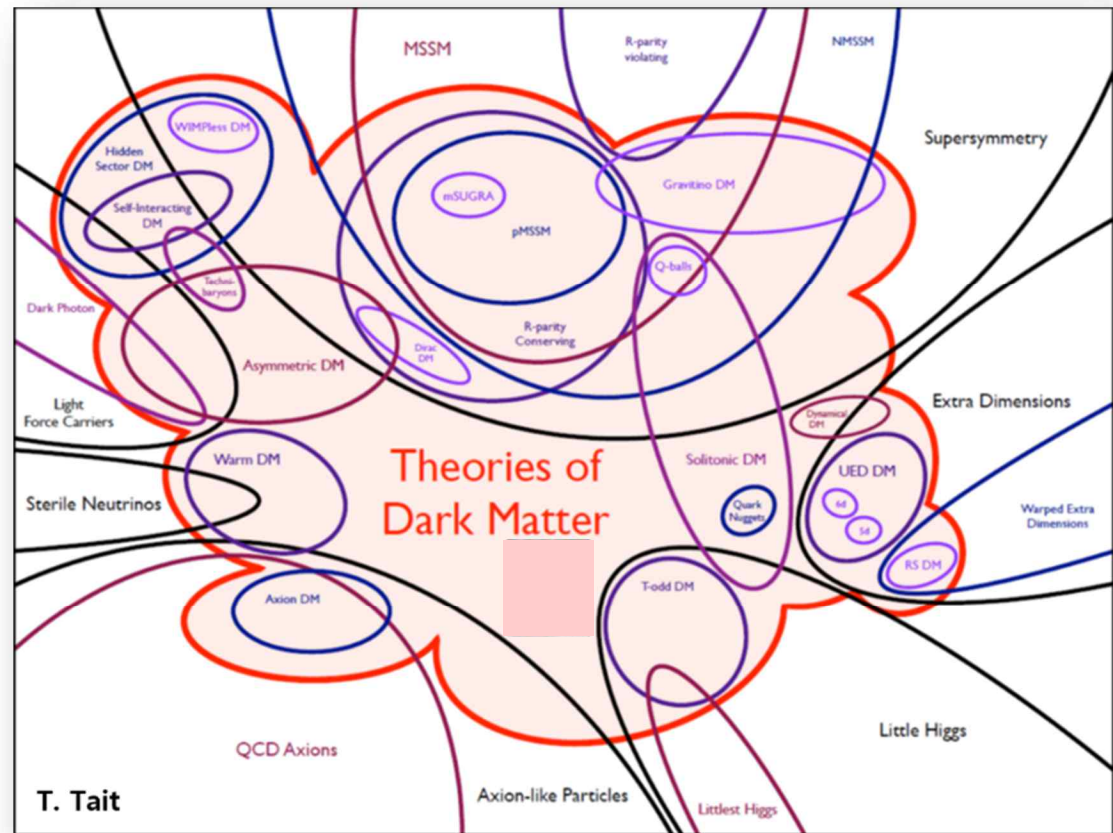
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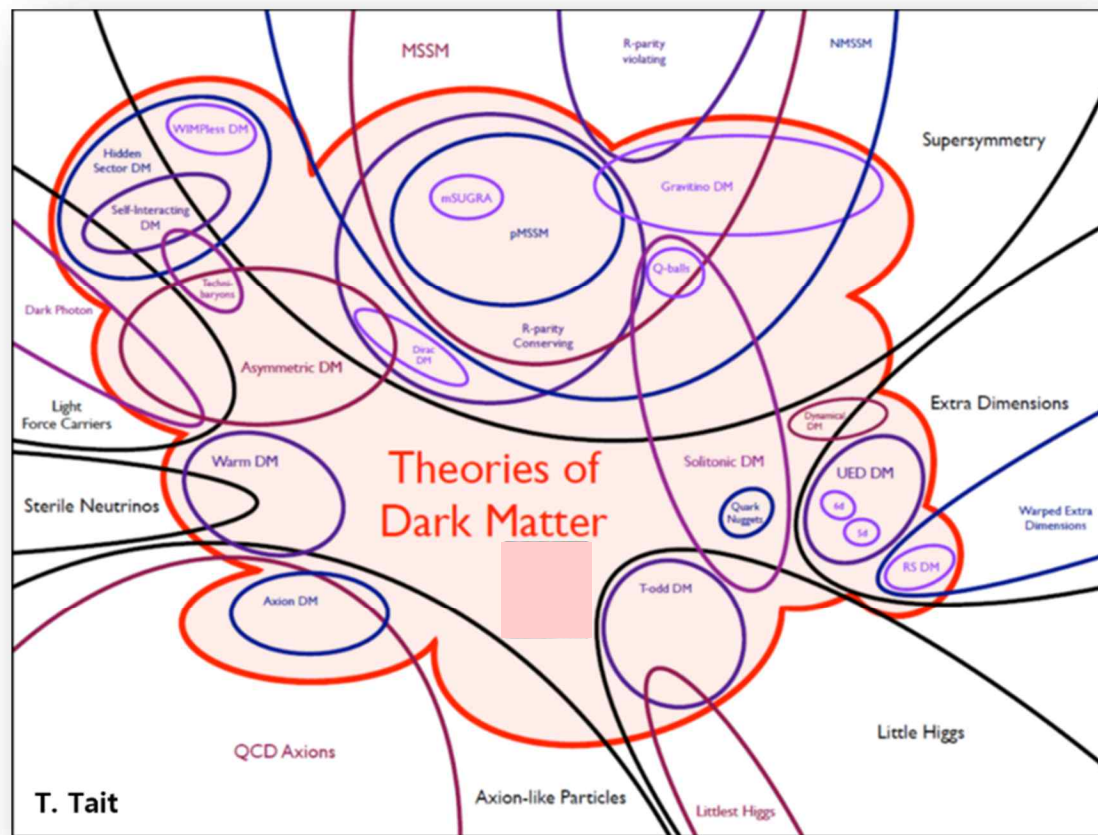
Dark matter ?

- WIMPs: Weakly Interacting Massive Particles ?
- Super-WIMPs ?
- Axions ?
- Sterile neutrinos ?
- FIMPs : Feebly IMPs ?



Dark matter ?

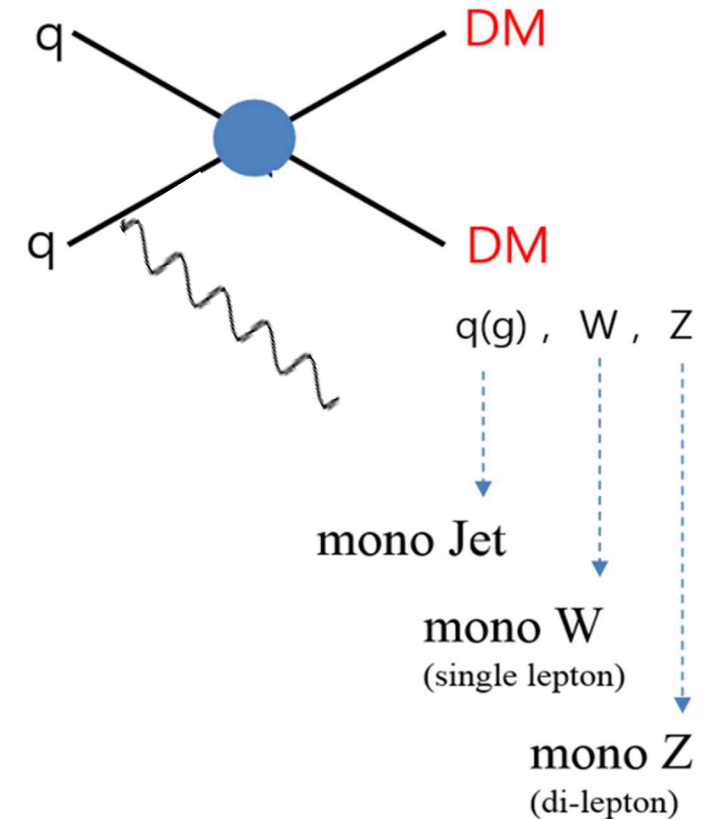
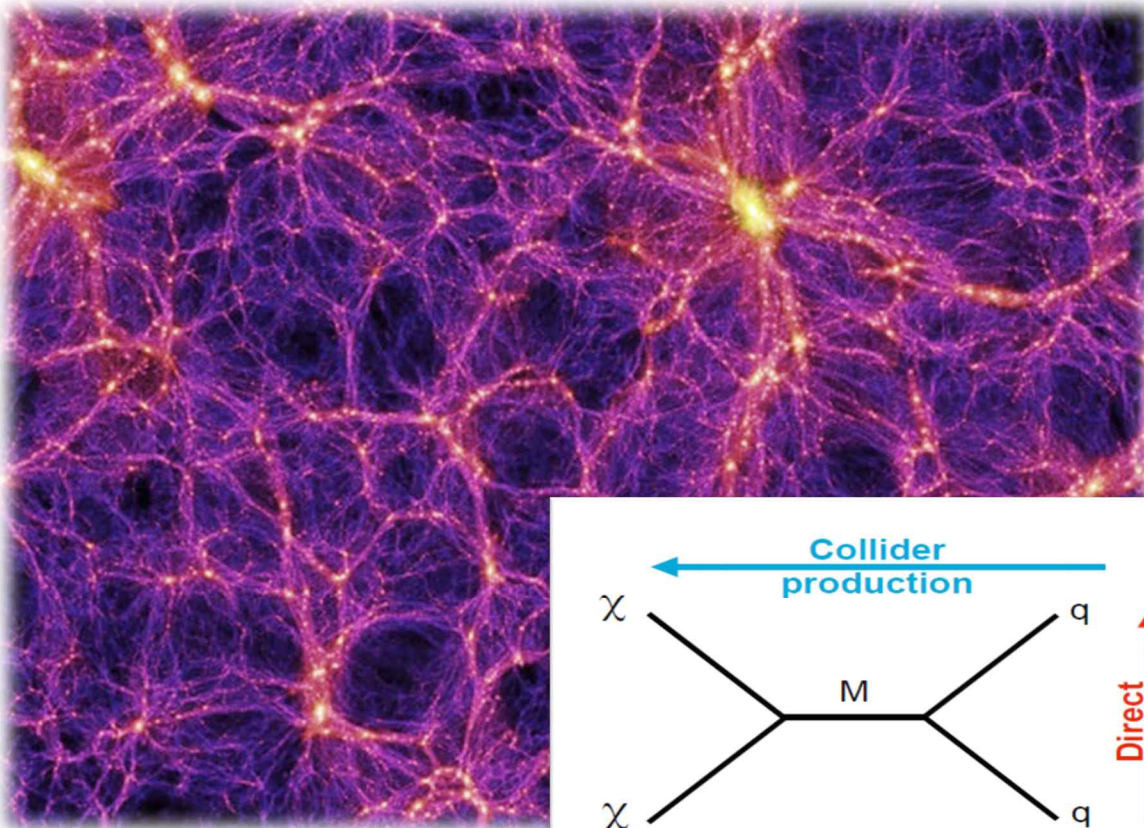
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→ Something invisible !!

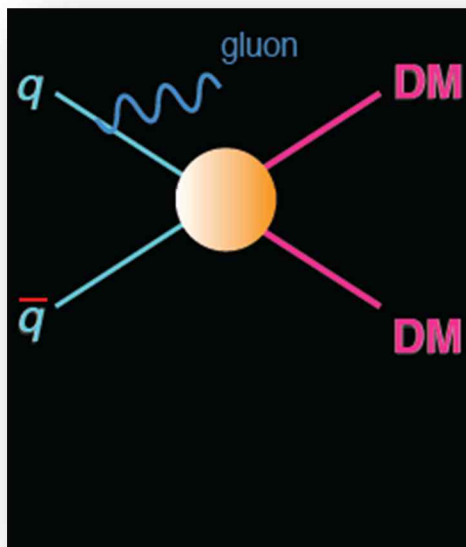
Dark Matter Production at colliders

- There are many models from the theoretical point of view.
- However the final states are simple from the experimental point of view.
 - Events characterized by ISR recoil of X
 - mono X + missing energy in collider experiments



Effective Field Theory for dark matter

- Effective Field Theory(EFT) interaction between DM and SM particles is a contact interaction.
- Consider operator at the new physics scale Λ .



eff. operator $\mathcal{O}_S = \frac{1}{\Lambda^2}(\bar{\chi}\chi)(\bar{q}q)$

parton-level differential cross section:

$$\frac{d^2\hat{\sigma}_{\text{eff}}}{dp_T d\eta} = \frac{\alpha_s}{36\pi^2} \frac{1}{p_T} \frac{1}{\Lambda^4} \frac{[Q_{\text{tr}}^2 - 4m_{\text{DM}}^2]^{3/2} \left[1 + \frac{Q_{\text{tr}}^4}{(x_1 x_2 s)^2}\right]}{Q_{\text{tr}}}$$

matching: $\frac{1}{\Lambda^2} = \frac{g_\chi g_q}{M^2}$

The Results applied to Run 1 analysis at the LHC

Advantage : Limited number of degrees of freedom scale of interaction

Disadvantage : Only applicable at low momentum transfer

DM Production in EFT

General Interaction Lagrangian for DM in EFT frame can be

$$\mathcal{L} = \sum_q \left\{ \frac{m_q}{\Lambda_{D1}^3} \bar{q}q \bar{\chi}\chi + \frac{1}{\Lambda_{D8}^2} \bar{q}\gamma^\mu \gamma_5 q \bar{\chi}\gamma_\mu \gamma_5 \chi + \frac{1}{\Lambda_{D5}^2} \bar{q}\gamma^\mu q \bar{\chi}\gamma_\mu \chi + \frac{1}{\Lambda_{D9}^2} \bar{q}\sigma^{\mu\nu} q \bar{\chi}\sigma_{\mu\nu} \chi \right\}$$

D1 : Scalar, D8 : Axial-Vector, D5 : Vector and D9 : Tensor

χ : Dirac Fermion Dark Matter

Λ : interaction scale(or suppression scale)

- Mediate DM + q/g interaction under $M > M_\chi$
 - M_χ : Dark Matter mass, M : Mediator mass
- Suppression scale by Λ (or M_*) $\sim M/(g_\chi g_q)^{1/2}$
- $M_\chi < M/2$ required for four-m/m conservation
- $g_\chi g_q < (4\pi)^2$ required for perturbation
 $\rightarrow \Lambda > M_\chi / 2\pi$

Dirac Fermions

Name	Operator	Coefficient
D1	$\bar{\chi}\chi\bar{q}q$	m_q/M_*^3
D2	$\bar{\chi}\gamma^5\chi\bar{q}q$	im_q/M_*^3
D3	$\bar{\chi}\chi\bar{q}\gamma^5q$	im_q/M_*^3
D4	$\bar{\chi}\gamma^5\chi\bar{q}\gamma^5q$	m_q/M_*^3
D5	$\bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu q$	$1/M_*^2$
D6	$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu q$	$1/M_*^2$
D7	$\bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu\gamma^5q$	$1/M_*^2$
D8	$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu\gamma^5q$	$1/M_*^2$
D9	$\bar{\chi}\sigma^{\mu\nu}\chi\bar{q}\sigma_{\mu\nu}q$	$1/M_*^2$
D10	$\bar{\chi}\sigma_{\mu\nu}\gamma^5\chi\bar{q}\sigma_{\alpha\beta}q$	i/M_*^2
D11	$\bar{\chi}\chi G_{\mu\nu}G^{\mu\nu}$	$\alpha_s/4M_*^3$
D12	$\bar{\chi}\gamma^5\chi G_{\mu\nu}G^{\mu\nu}$	$i\alpha_s/4M_*^3$
D13	$\bar{\chi}\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$i\alpha_s/4M_*^3$
D14	$\bar{\chi}\gamma^5\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$\alpha_s/4M_*^3$

scalar

vector

axial-vector

tensor

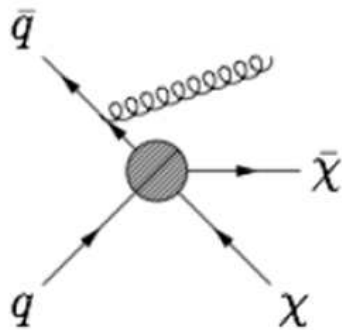
scalar

Spin-independent

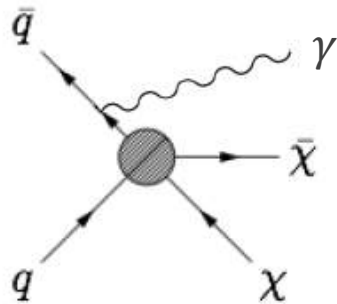
Spin dependent

J. Goodman et. al, Phys. Rev. D 82, 116010 (2010)

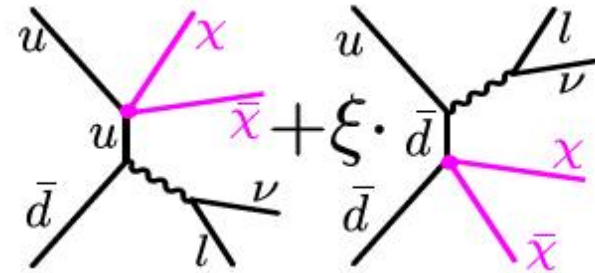
Channels for Dark matter at the LHC



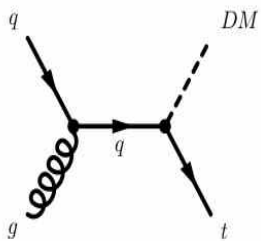
Mono-jet



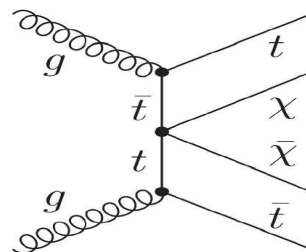
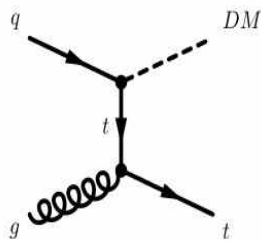
Mono-photon



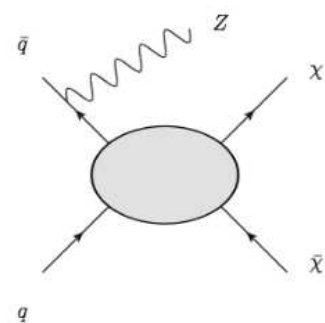
Mono-W



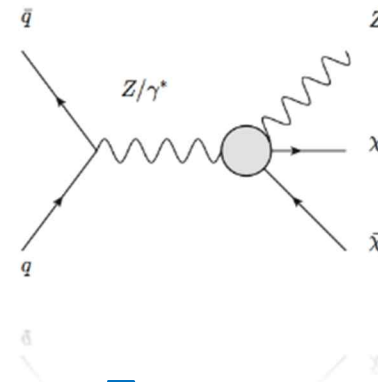
Mono-top



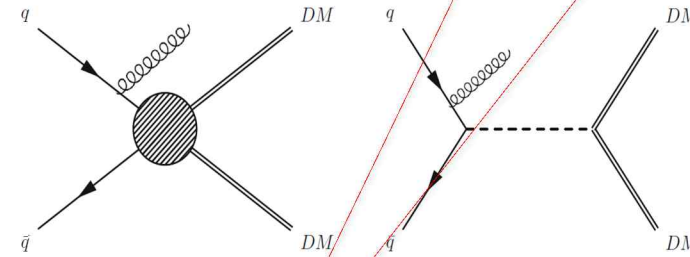
Top pair



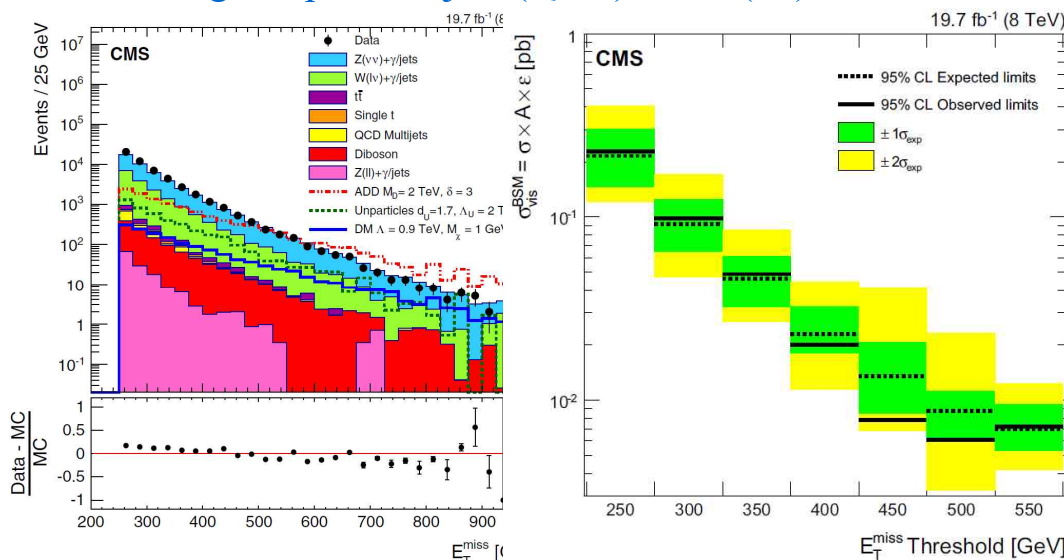
Mono-Z



- A Jet from ISR
- one energetic jet and large missing energy (MET)
 - $Z(\rightarrow \nu\nu) + \text{jet}$, irreducible background and $W + \text{jets}(\text{lost lepton})$: main backgrounds(62 and 36%)
 - QCD jet misidentified, $t\bar{t}$ etc.
- Selection \rightarrow very large MET !
 - $p_T > 110 \text{ GeV}$, $|\eta| < 2.4$
 - second jet are allowed , if $\Delta\phi(\text{jet1}, \text{jet2}) < 2.5$
 - $\text{MET} > 250, 300, 350, 400, 450, 500, 550 \text{ GeV}$
 - Veto events with isolated leptons.
- Estimates of Standard Model contributions
 - Data-driven : $Z(\nu\nu) + \text{jets}$ (from $Z \rightarrow \mu\mu$)
 - Data-driven : $W + \text{jets}$ (from $W \rightarrow \mu\nu$)
 - single-top, multi-jets (QCD), $t\bar{t}$, $Z(\ell\ell)$

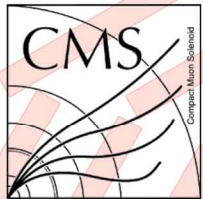


$E_T^{\text{miss}} \text{ (GeV)} \rightarrow$	>250	>300	>350	>400	>450	>500	>550
$Z(\nu\nu) + \text{jets}$	32100 \pm 1600	12700 \pm 720	5450 \pm 360	2740 \pm 220	1460 \pm 140	747 \pm 96	362 \pm 64
$W + \text{jets}$	17600 \pm 900	6060 \pm 320	2380 \pm 130	1030 \pm 65	501 \pm 36	249 \pm 22	123 \pm 13
$(t\bar{t})$	446 \pm 220	167 \pm 84	69 \pm 35	31 \pm 16	15 \pm 7.7	6.6 \pm 3.3	2.8 \pm 1.4
$Z(\ell\ell) + \text{jets}$	139 \pm 70	44 \pm 22	18 \pm 9.0	8.9 \pm 4.4	5.2 \pm 2.6	2.3 \pm 1.2	1.0 \pm 0.5
Single t	155 \pm 77	53 \pm 26	18 \pm 9.1	6.1 \pm 3.1	0.9 \pm 0.4	–	–
QCD multijets	443 \pm 270	94 \pm 57	29 \pm 18	4.9 \pm 3.0	2.0 \pm 1.2	1.0 \pm 0.6	0.5 \pm 0.3
Diboson	980 \pm 490	440 \pm 220	220 \pm 110	118 \pm 59	65 \pm 33	36 \pm 18	20 \pm 10
Total SM	51800 \pm 2000	19600 \pm 830	8190 \pm 400	3930 \pm 230	2050 \pm 150	1040 \pm 100	509 \pm 66
Data	52200	19800	8320	3830	1830	934	519
Exp. upper limit $+1\sigma$	5940	2470	1200	639	410	221	187
Exp. upper limit -1σ	2870	1270	638	357	168	123	104
Exp. upper limit	4250	1800	910	452	266	173	137
Obs. upper limit	4510	1940	961	397	154	120	142



- Consistent with the SM backgrounds
- Set the model-independent observed 95% CL upper limits on the visible x-section for non-SM production according to MET cuts

Typical Mono-jet Event



CMS Experiment at LHC, CERN
Data recorded: Fri Oct 5 20:41:32 2012 CEST
Run/Event: 204553 / 26729384
Lumi section: 31

Jet 0,
et = 921.98
eta = -0.463
phi = 2.508

MET 0,
pt = 913.68
eta = 0.000
phi = -0.657

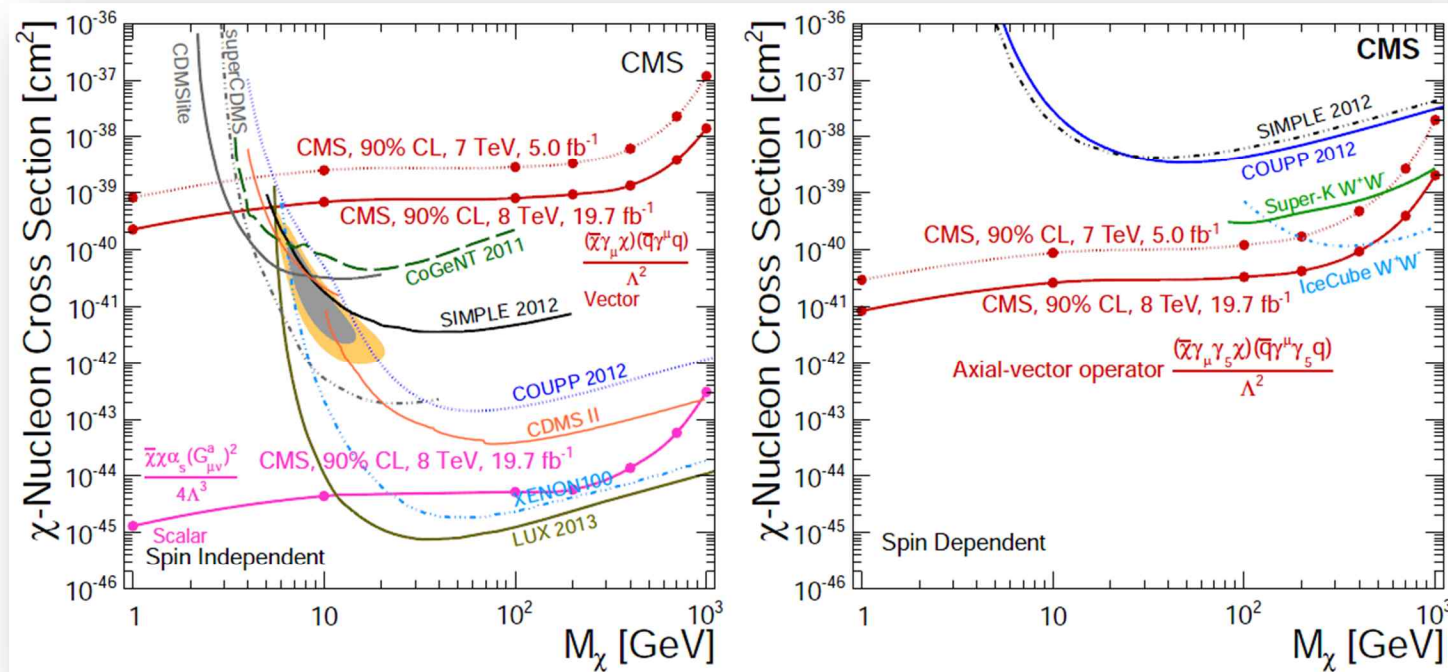


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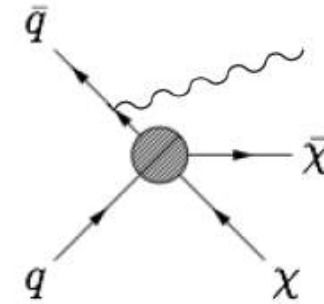
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Mono-jet (DM mass vs x-section)

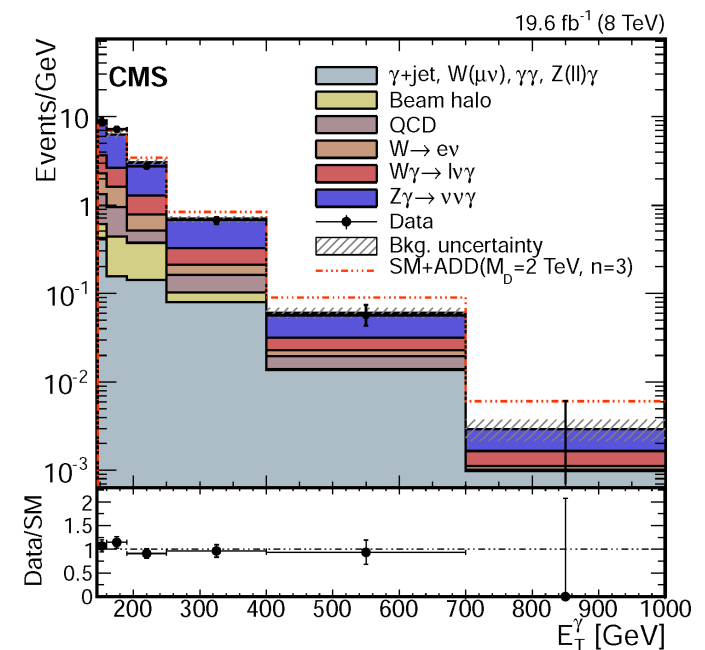
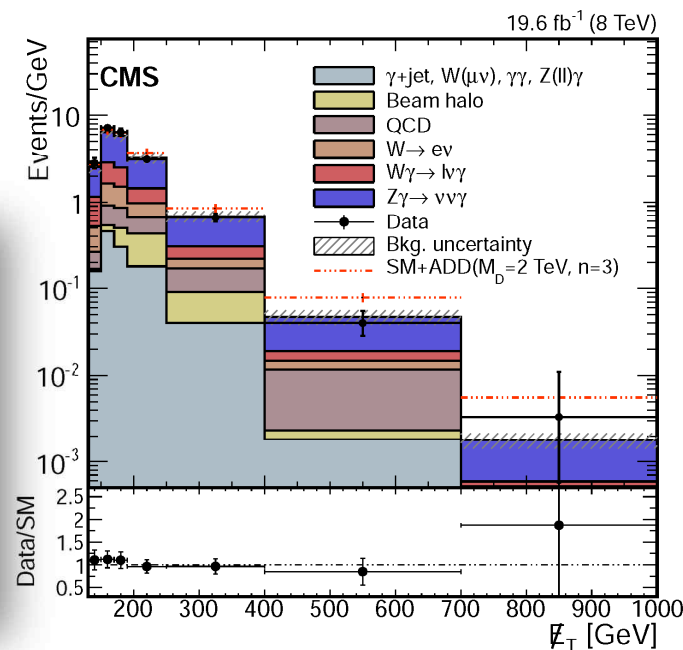


- The observed limit on the x-section depends on M_χ and nature of its interaction with SM particles.
- The limits on the effective contact interaction scale as a function of M_χ translated into a limit on the dark matter-nucleon scattering cross section in the framework of EFT.
- Shown for the vector & scalar case(SI) and axial vector(SD).
- Better results shown in Spin Dependent than direct detection experiments $\rightarrow \sim 1000$ times better
- Sensitive to low M_χ for Spin Independent case.

- Dark matter particle can be produced in process of photon radiated by one of incoming quarks.
- Signature : 1 photon + 2 dark matters
 - One isolated photon
 - Very large MET (> 140 GeV) and photon $p_T(> 145$ GeV)
 - $Z\gamma \rightarrow \nu\bar{\nu}\gamma$ as irreducible background, $W\gamma$, W , γ +jets, QCD multijet, diphotons etc..



Process	Estimate
$Z(\rightarrow \nu\bar{\nu}) + \gamma$	345 ± 43
$W(\rightarrow \ell\nu) + \gamma$	103 ± 21
$W \rightarrow e\nu$	60 ± 6
jet $\rightarrow \gamma$ MisID	45 ± 14
Beam halo	25 ± 6
Others	36 ± 3
Total background	614 ± 63
Data	630

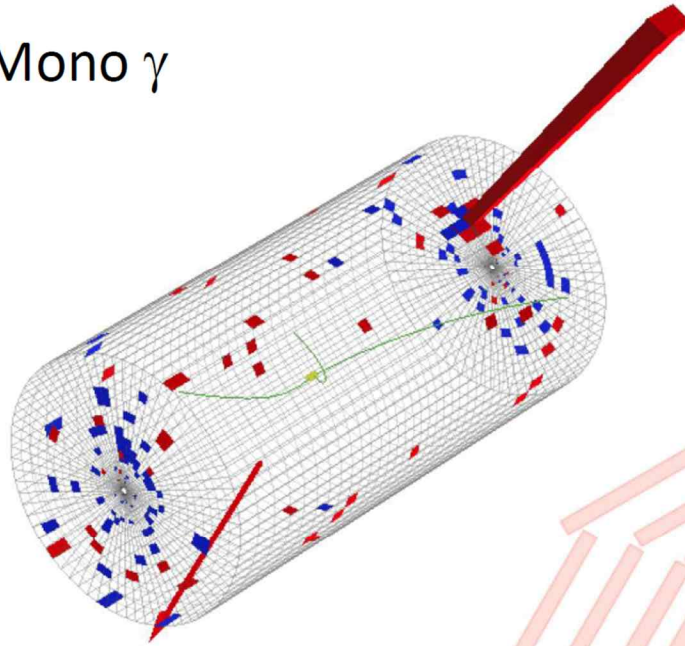


A good agreement with expectations from SM background.

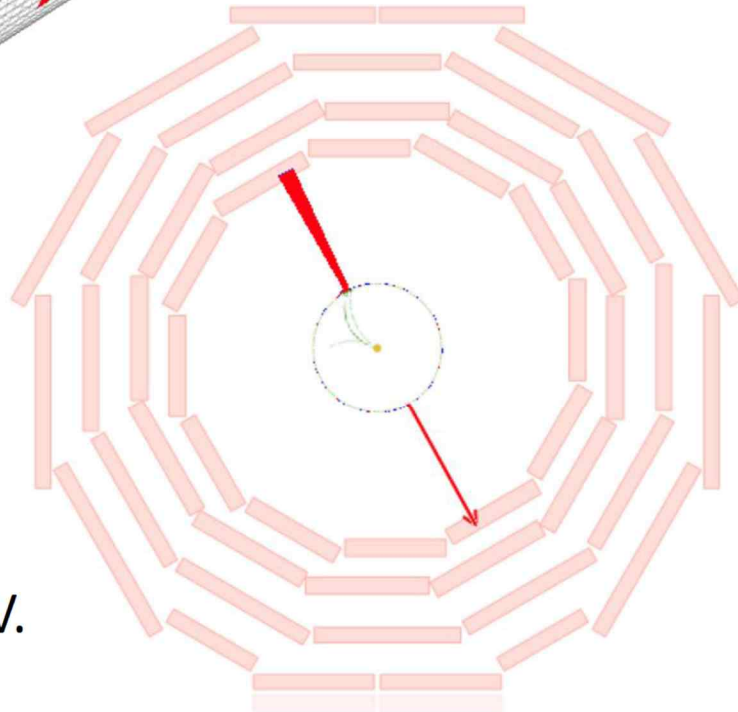
Mono-photon event



Mono γ



CMS Experiment at LHC, CERN
Data recorded: Sun Apr 24 22:57:52 2011 CDT
Run/Event: 163374 / 314736281
Lumi section: 604

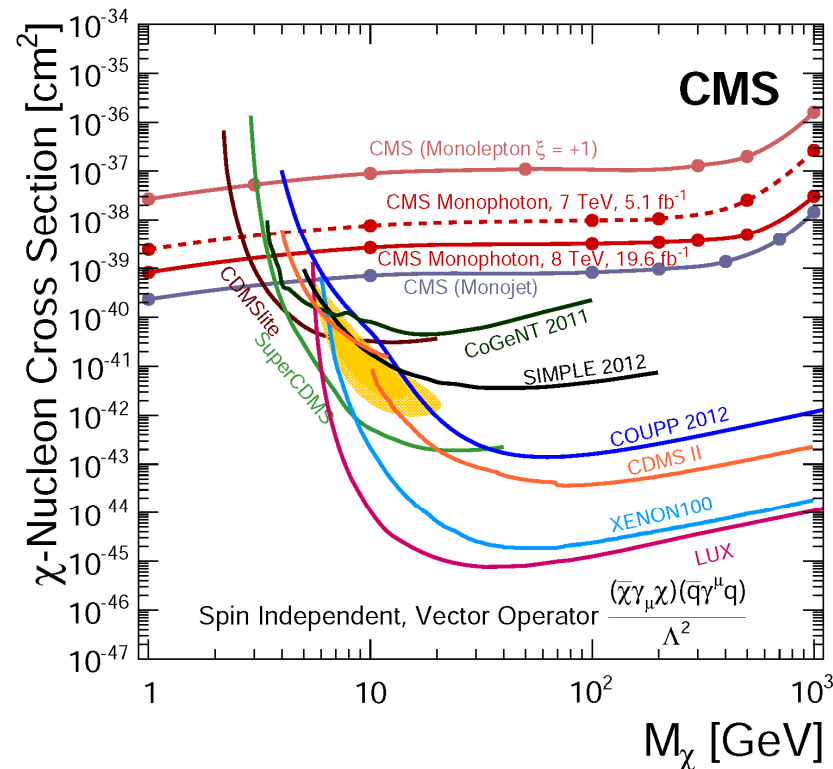
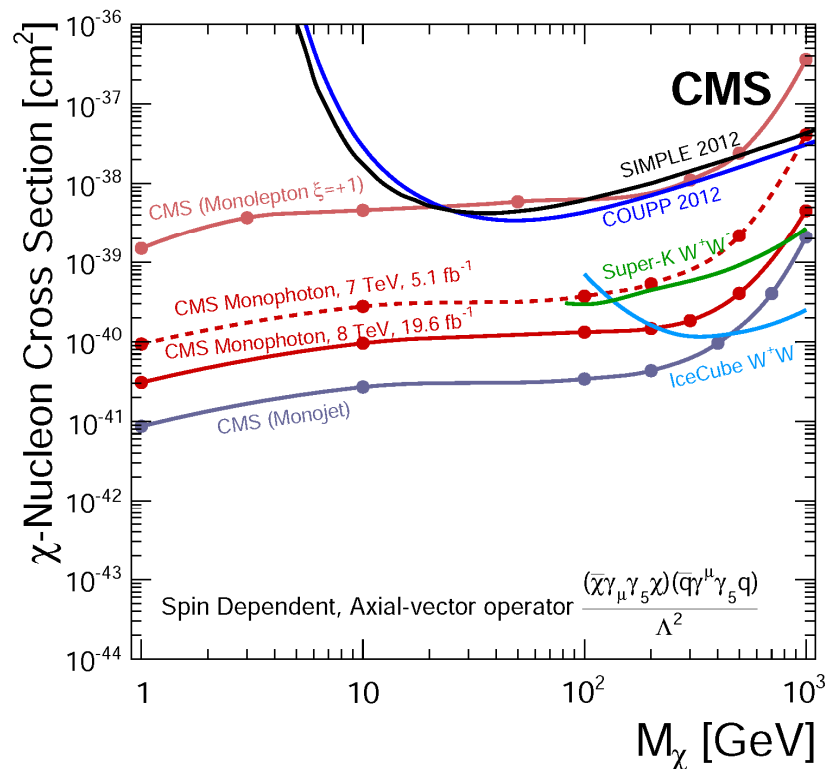


$p_T^\gamma = 384 \text{ GeV}$, $\text{MET} = 407 \text{ GeV}$.

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Mono-photon results

- The 90% CL upper limits on the χ -nucleon cross section as a function of the DM mass (M_χ)
- Results also shown for CMS mono-jet and mono-lepton signatures
 - ζ is the interference parameter
- Also shown several direct detection experiments for comparison
 - The yellow contours show the 68% and 95% CL contours respectively for a possible signal from CDMS

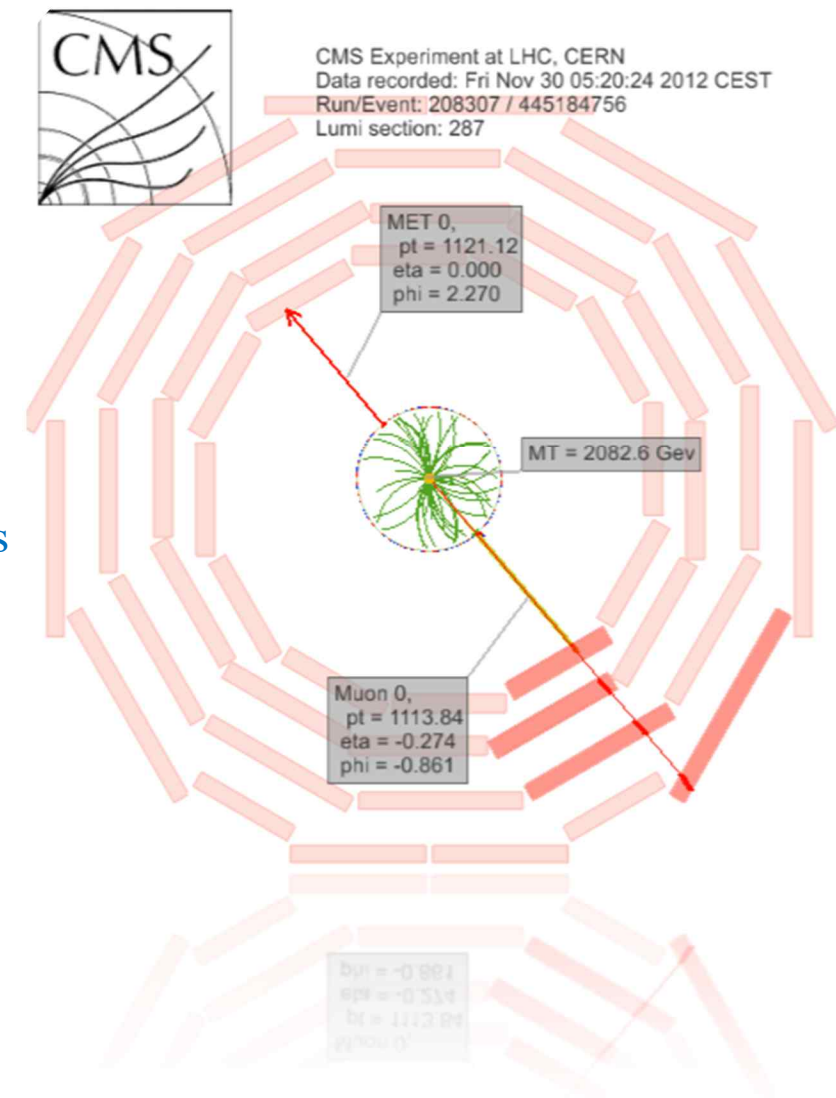


Mono-W(to lepton)

CMS-EXO-12-060

PHYSICAL REVIEW D **91**, 092005 (2015)

- Experimental point of view :
 - Easier to trigger (mono-lepton)
 - Clean, well simulated background
 - Low systematic uncertainties from the detector
- Theoretical point of view :
 - Higher production cross section than mono-Z
 - Well simulated, Quark-sensitive interference effects
- Plenty reinterpretations to the model
 - Sequential Standard Model W'
 - Interference W - W'
 - Mass Limit $M(W') > 3.35$ TeV
 - Split-UED model
 - 2-dim. limits are set in $(1/R, \mu)$ space
 - Helicity-Non-Conserving Contact Interaction (HNC-CI)
 - Limits are set on Λ (binding energy)

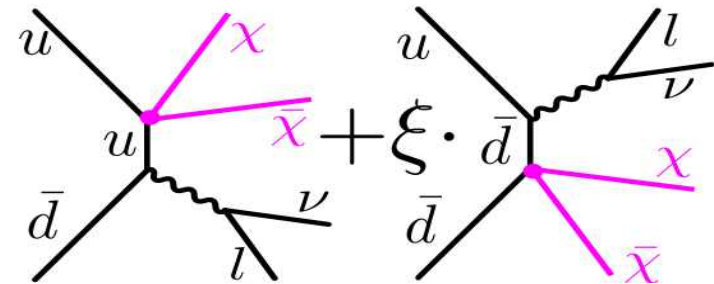


Huge Interference effects

- The DM particles either couple to the up- or the down-type quark with their relative coupling strength parameterized by the factor ξ .

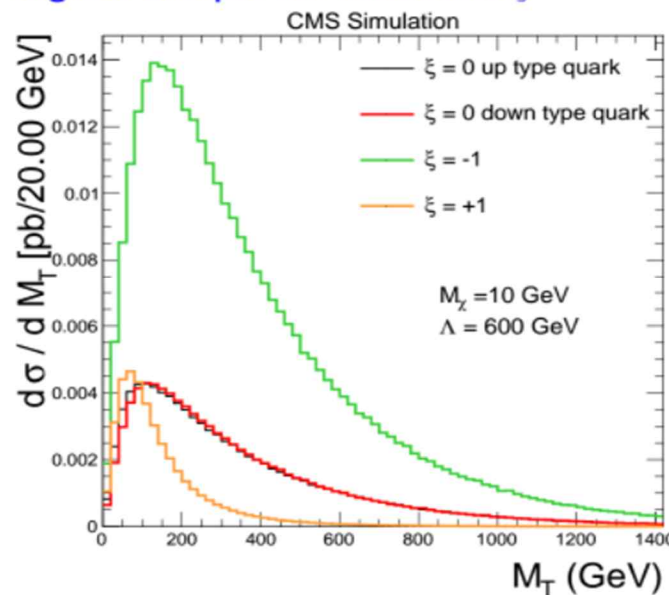
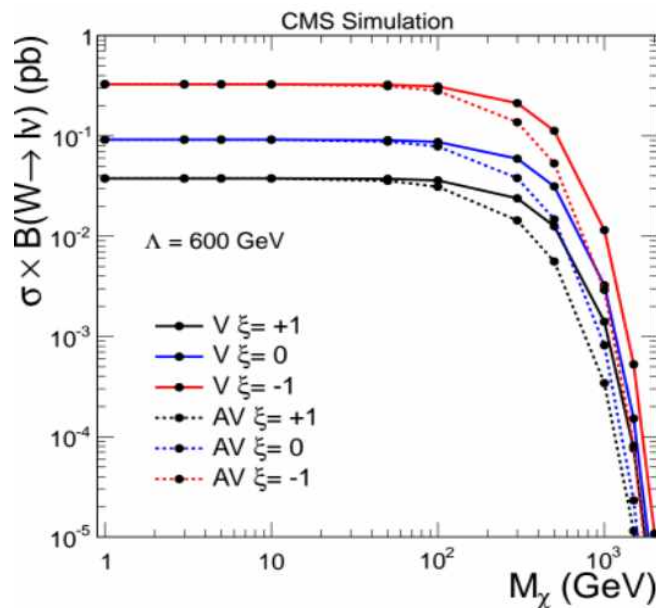
$$V: \frac{1}{\Lambda^2} \bar{\chi} \gamma^\mu \chi (\bar{u} \gamma^\mu u + \xi \bar{d} \gamma^\mu d)$$

$$AV: \frac{1}{\Lambda^2} \bar{\chi} \gamma^\mu \gamma^5 \chi (\bar{u} \gamma^\mu \gamma^5 u + \xi \bar{d} \gamma^\mu \gamma^5 d)$$



The interesting values are $\xi = -1, 0, +1$

signal shape for different ξ

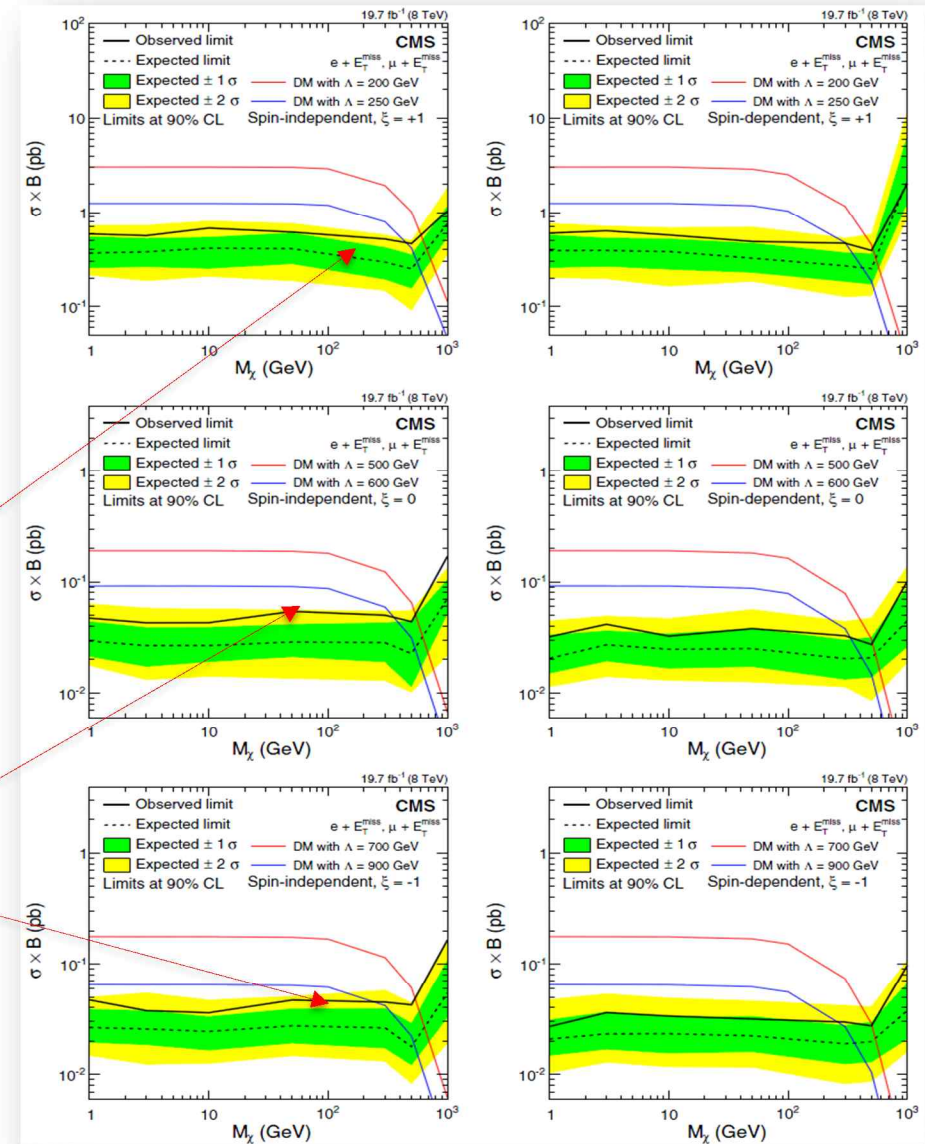


- Strong dependence on ξ
- factor $\xi = +1$ (destructive)
- 0 (u/d no differ-)
- 1 (constructive)
- The interference changes the total cross-section and steepens the M_T spectrum

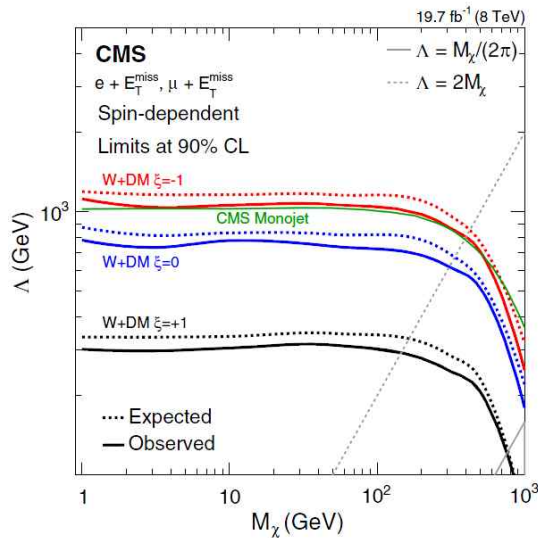
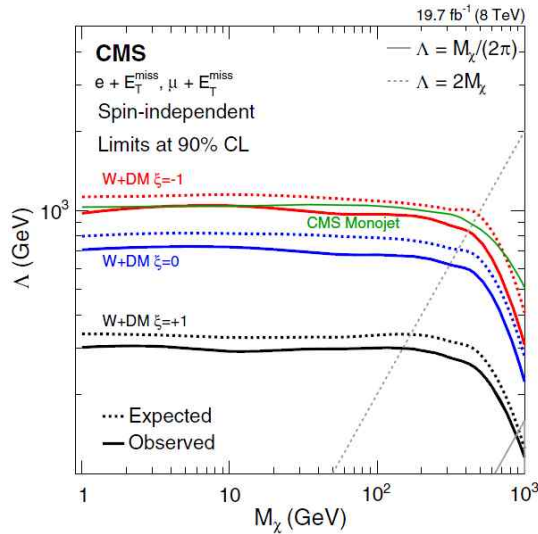
One can expect the results have a strong dependence on the interference.

X-sections w(w/o) interference

- The excluded cross section is flat as a function of M_χ ,
 → because the signal kinematics do not change appreciably for different M_χ
- The coupling does not have a large effect on the excluded cross section
- The different interference scenarios have a visible influence on the limit
 → for $\xi = +1$, a cross section > 0.6 pb excluded
 → for $\xi = 0$ and $\xi = -1$, the x-section limit is 0.05 pb
- For high M_χ the phase space to produce two heavy particles and a W-boson is small
 → the signal cross section is reduced



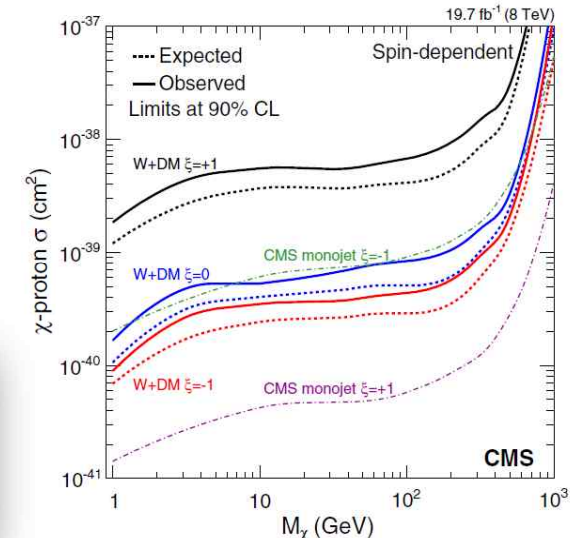
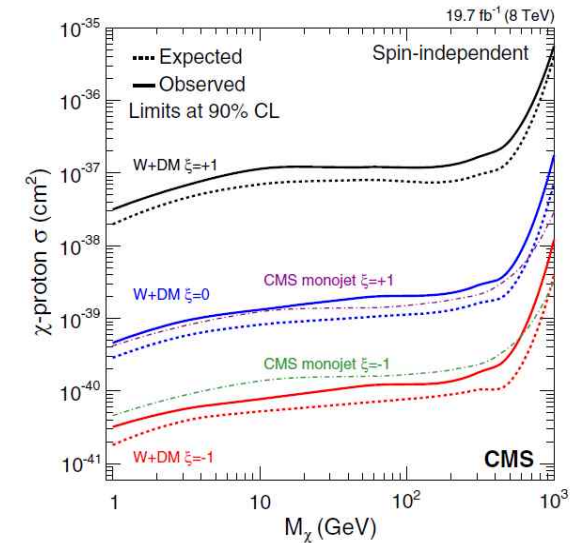
Λ and σ_χ vs M_χ



- Cross section upper limits transformed into lower limits on the effective scale parameter Λ as a function of M_χ
- For lower masses a constant Λ exclusion obtained for $M_\chi \leq 100$ GeV
 - $\Lambda < 300$ GeV for $\xi=1$, $\Lambda < 700$ GeV for $\xi=0$, and $\Lambda < 1000$ GeV for $\xi=-1$.
- The difference between vector-like and axial-vector-like couplings is small for low M_χ for all three values of ξ , but a difference is observed in the high- M_χ region, > 100 GeV
- The χ -proton cross section upper limits at 90% for $M_\chi = 10$ GeV as an example



ξ	Vector coupling (cm ²)	Axial-vector coupling (cm ²)
-1	4×10^{-41}	1×10^{-40}
0	6×10^{-40}	2×10^{-40}
+1	3×10^{-38}	2×10^{-39}

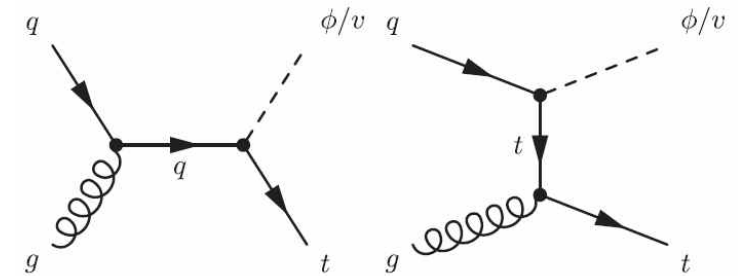


Mono-lepton results are better than mono-jet for the construction case.

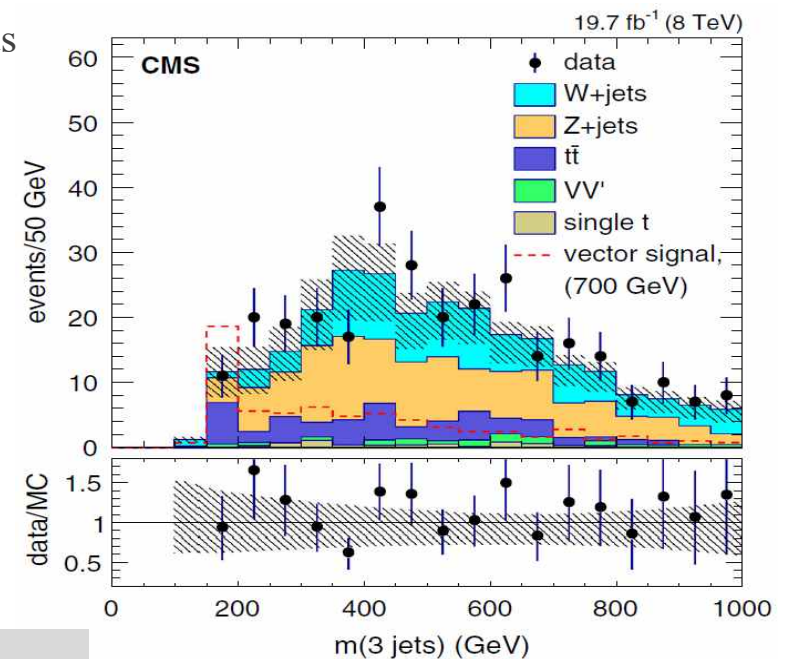
Mono-top

PRL 114, 101801 (2015)

- DM particle(either scalar or vector) could be produced in association with a top quark.
 - Could be invisible particle produced top quark by quark flavor changing
 - R-parity violating SUSY or FC interaction mediated by invisible state
 - Both with branching fraction to unity assumed
- A single top decays a bottom quark and two light quarks
 - Only jets with $p_T \geq 35$ GeV and $|\eta| < 2.4$
 - The two highest- p_T (leading) jets > 60 GeV, third highest $p_T > 40$ GeV
 - The inv. mass of the 3 jets < 250 GeV and rejected additional jets with a $p_T > 35$ GeV
 - requiring one of the three jets to be identified as a candidate jet from a b quark
- 3 jets with MET, Z + jets(73%), W + jets(13%) and $t\bar{t}$ as backgrounds



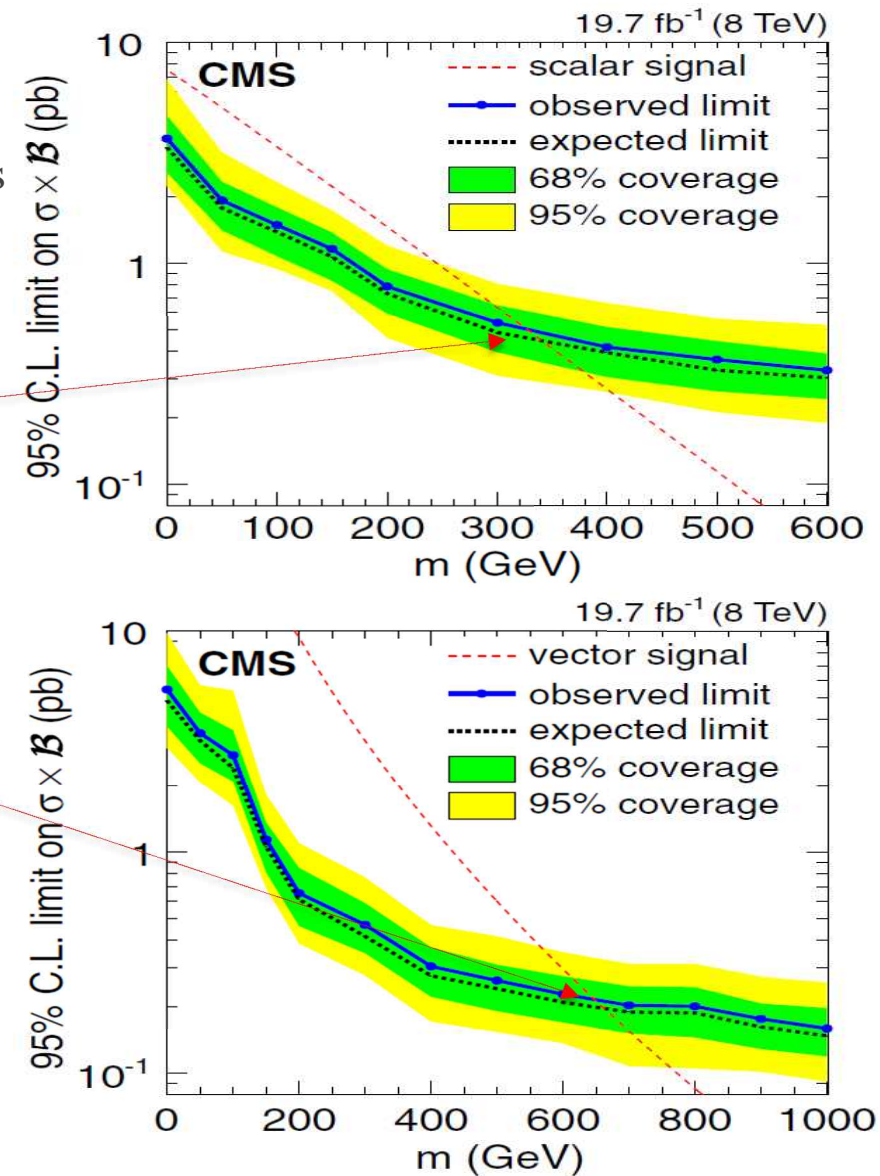
	No b tag	One b tag
$t\bar{t}$	$6 \pm 0 \pm 5$	$12 \pm 0 \pm 12$
W + jets	$18 \pm 9 \pm 7$	$3 \pm 1 \pm 2$
Z + jets	$103 \pm 33 \pm 9$	$11 \pm 10 \pm 1$
Single top	$2 \pm 1 \pm 1$	$1 \pm 1 \pm 1$
VV'	$5 \pm 0 \pm 0$	$0 \pm 0 \pm 0$
Multijet	$6(\pm 39)$	$1(\pm 9)$
Total background	140 ± 36	28 ± 16
Signal	2 ± 6	3 ± 11
Data	143	30



No excess is observed above the background expectation

Mono-top Results

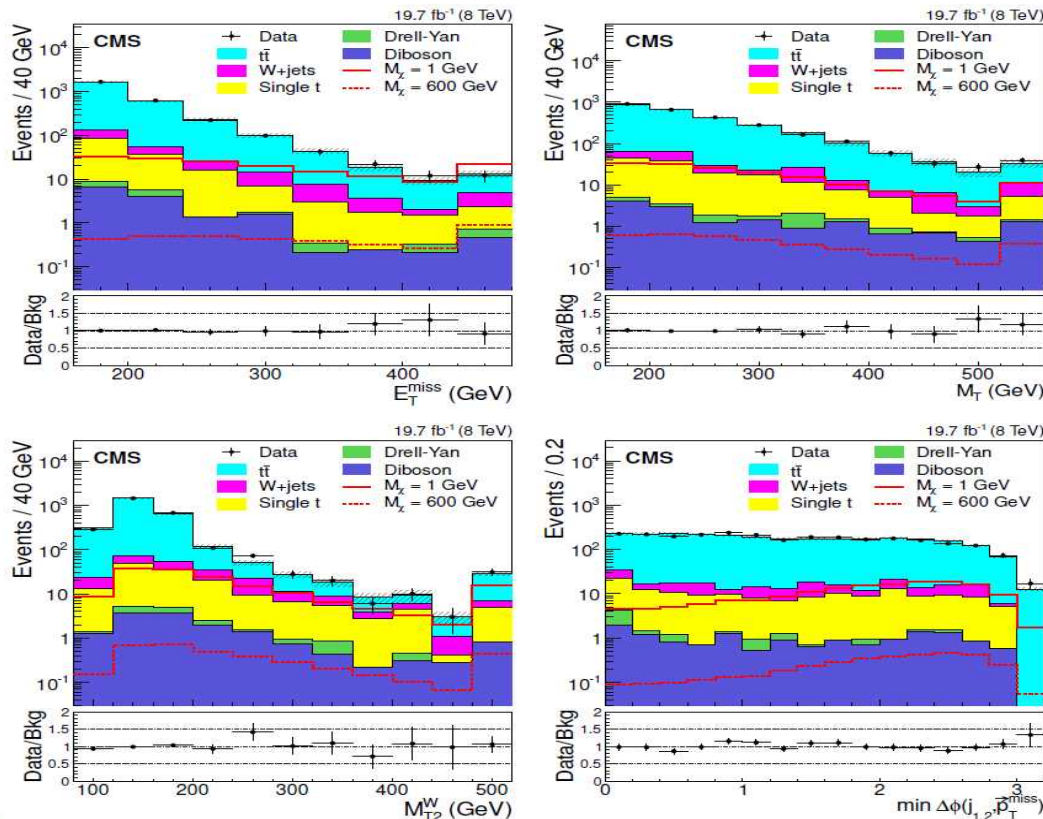
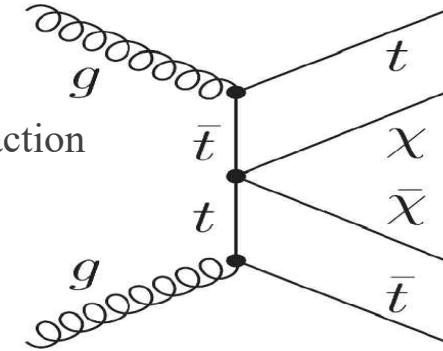
- The 95% C.L. expected and observed limits on product of production cross section of mono-top as function of mass of invisible bosonic state. (scalar and vector fields)
- The scalar and vector particles, with masses below 330 and 650 GeV, respectively, are excluded at 95% confidence level.
- The substantial Extension of previous limit on the search from the CDF
- Complementary with the ATLAS results obtained with the leptonic top decay channel.
- This results can be applied to the non-thermal DM model as well.



Top-quark pairs associated

JHEP06(2015)121

- Assume only one new Dirac fermion related to DM by EFT Lagrangian
 - Assume coupling strength proportional to the mass of quark.
 - suppressed to light quarks → dominant to 3rd generation quarks
- Consider DM production in association with top quark pair and only scalar interaction
- Requiring the presence of one lepton, multiple jets and large MET
- Backgrounds : $t\bar{t}$ +jets, $t\bar{t}$ + γ /W/Z, W+jets, single top, dibosons and DY

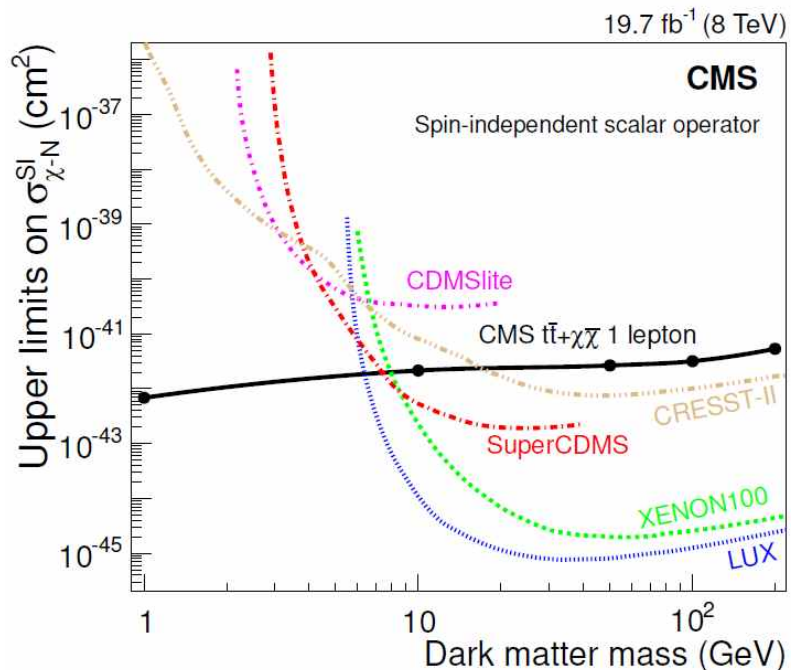
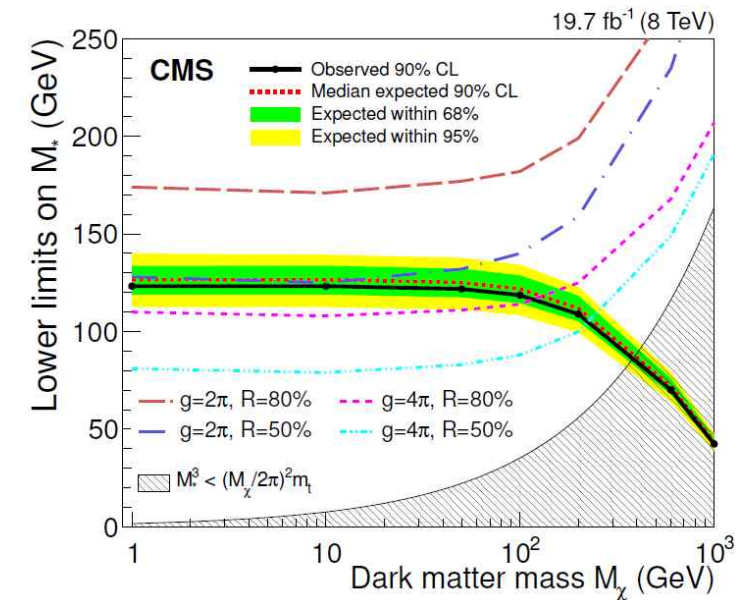


Source	Yield ($\pm\text{stat} \pm\text{syst}$)
$t\bar{t}$	$8.2 \pm 0.6 \pm 1.9$
W	$5.2 \pm 1.8 \pm 2.1$
Single top	$2.3 \pm 1.1 \pm 1.1$
Diboson	$0.5 \pm 0.2 \pm 0.2$
Drell-Yan	$0.3 \pm 0.3 \pm 0.1$
Total Bkg	$16.4 \pm 2.2 \pm 2.9$
Data	18

The data are in good agreement with expectations from SM background.

Top quark-pair : Results

- Observed exclusion limits in the plane of (M_χ, M_*) at a 90% CL.
- The background-only expectations are represented by dashed line (68% and 95% CL bands).
- A lower bound of the validity of the EFT indicated by the upper edge of the hatched area.
- The four curves corresponding to different g and R values
 - represent the lower bound on M_* for different (g, R) set
 - indicate further restrictions on the applicability of EFT



- The 90% CL upper limits on the DM-nucleon spin-independent scattering cross section as a function of M_χ for the scalar operator.
- Also shown are 90% CL limits from various direct DM search experiments

Better result in $M_\chi < \sim 5 - 10$ GeV region in comparison with DD experiments

Mono-Z(to Dilepton)

CMS PAS EXO-12-054

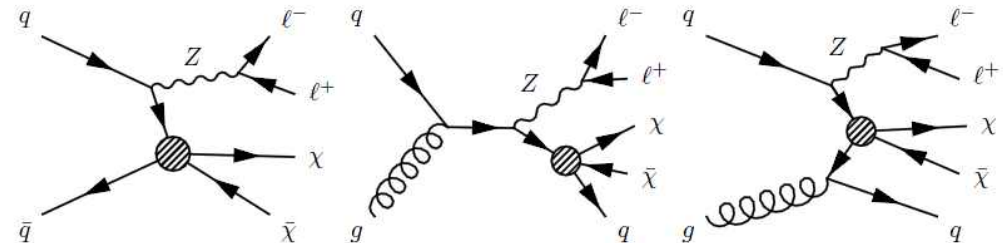
- Dark matter production : Z boson is radiated(ISR) or interacts directly with DMs

- Signature : 2 leptons + 2 dark matters

- 2 isolated leptons (electrons or muons)

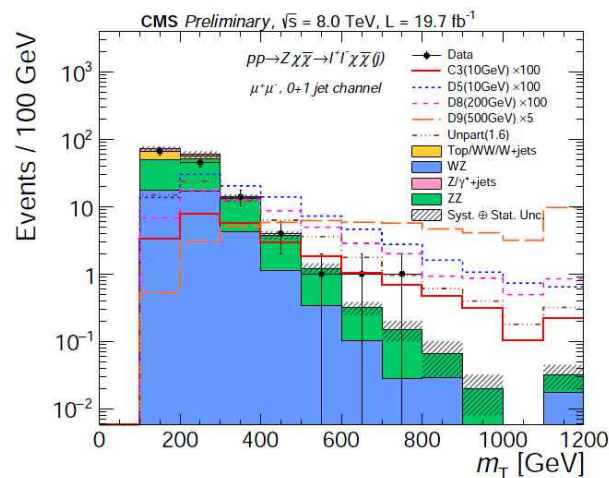
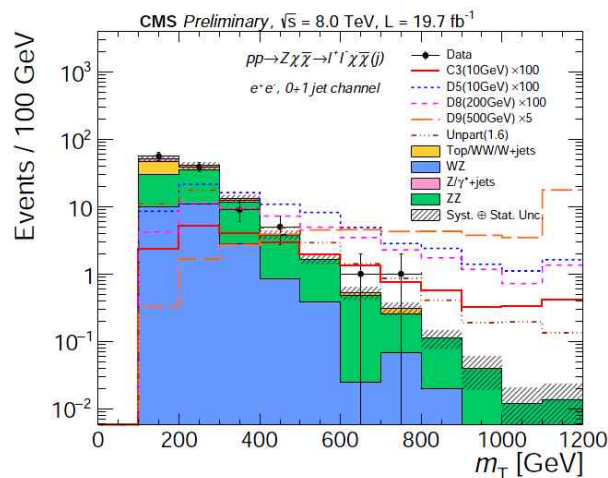
- Mass(l_l) in Z boson mass window

- Large missing transverse energy



- ZZ/WZ main backgrounds, $t\bar{t}$, tW , WW etc. as well.

- Applying to vector (D5), axial-vector (D8) and tensor (D9) coupling to the SM quarks for Dirac fermion DM and vector (C3) coupling for complex scalar DM.

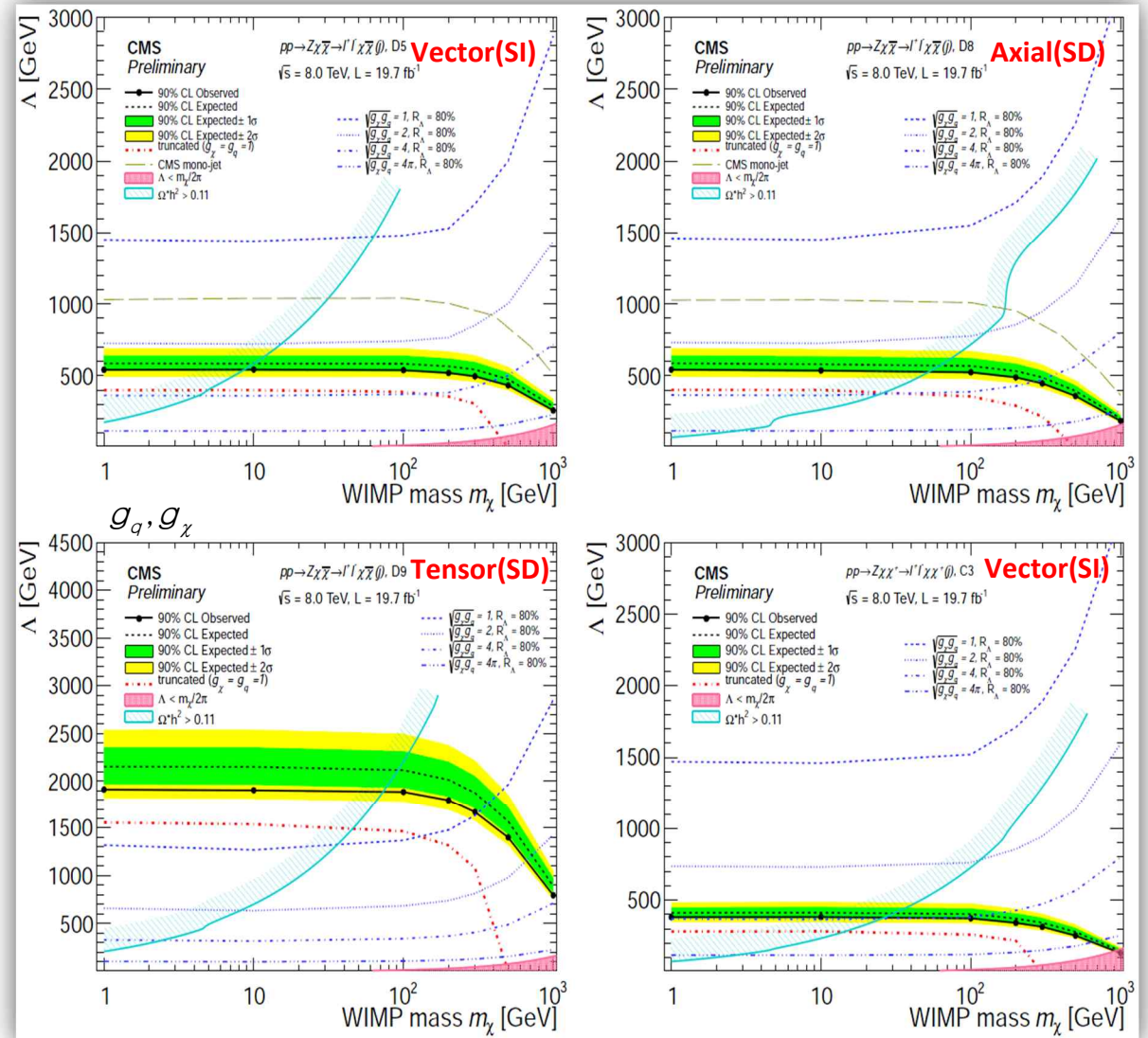


Process	e^+e^-	$\mu^+\mu^-$
C3(10GeV)	$0.20 \pm 0.004 \pm 0.02$	$0.24 \pm 0.005 \pm 0.02$
D5(10GeV)	$0.79 \pm 0.02 \pm 0.09$	$0.97 \pm 0.02 \pm 0.09$
D8(200GeV)	$0.48 \pm 0.01 \pm 0.06$	$0.59 \pm 0.01 \pm 0.05$
D9(500GeV)	$10.24 \pm 0.12 \pm 1.59$	$10.80 \pm 0.13 \pm 0.98$
Unparticle(1.6)	$48.96 \pm 0.88 \pm 3.26$	$65.80 \pm 1.07 \pm 4.38$
$Z/\gamma^* \rightarrow \ell^+\ell^-$	$8.18 \pm 1.93 \pm 0.82$	$8.59 \pm 3.01 \pm 1.02$
$WZ \rightarrow 3\ell\nu$	$25.08 \pm 0.53 \pm 2.84$	$40.70 \pm 0.69 \pm 4.50$
$ZZ \rightarrow 2\ell 2\nu$	$58.81 \pm 0.65 \pm 10.30$	$78.68 \pm 0.79 \pm 13.83$
$Top/WW/Z \rightarrow \tau^+\tau^-$	$18.74 \pm 3.39 \pm 3.27$	$22.92 \pm 2.29 \pm 3.44$
$W + jets$	$1.84 \pm 0.61 \pm 0.27$	
Total bkg.	$113 \pm 4 \pm 13$	$151 \pm 4 \pm 18$
Data	111	133

The data agree with SM background prediction and no excess is observed.

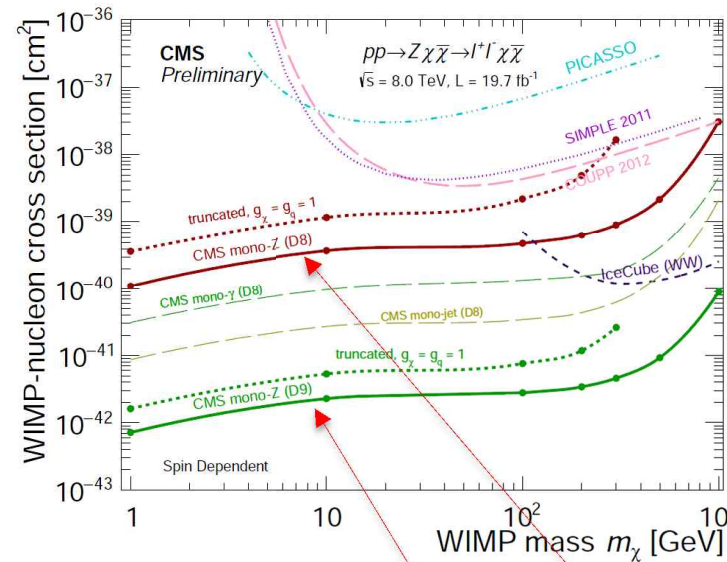
Dilepton - Λ vs M_χ

- 90% expected and observed CL Limits set on (M_χ, Λ) plane
- Appropriate choices of couplings and masses made for valid EFT
 - Consider massive mediator(M) couples to DM and quarks.(g_q and g_χ)
 - Can be $\Lambda \sim M/\sqrt{g_q g_\chi}$
- Lower bound set to $\Lambda > M_\chi/2\pi$ with condition $\sqrt{g_q g_\chi} < 4\pi$
- Realistically R_Λ to be 80% with various values as well as truncated limits
 - R_Λ more realistic min. constraint $Q_{tr} < M$ applied \rightarrow fraction of total events w/ the constraint of Q less than cutoff scale
 - Truncated : removed events w/ $Q_{tr} > M$ at gen. level \rightarrow this limit curve pattern coincides with the observed one in shape
- $\Omega h^2 = 0.11$ measured by WMAP
- In overall, tune very much within EFT
- Poor results than CMS mono-jet for D5 and D8



Dilepton – Nucleon xsection vs M_χ

- The 90% CL upper limits on the DM-nucleon cross section as a function of the DM particle mass
- Spin-dependent(D8 and D9) and Spin-Independent(C3 and D5) limits
- Comparisons with various direct detection experiments
- CMS mono-jet, mono-photon and Higgs portal DM results as well.
- In general, dominant contributions to SD against direct detection experiments
- In the region $M_\chi < 5$ to 10 GeV, LHC is better than DD experiments for the case of SI

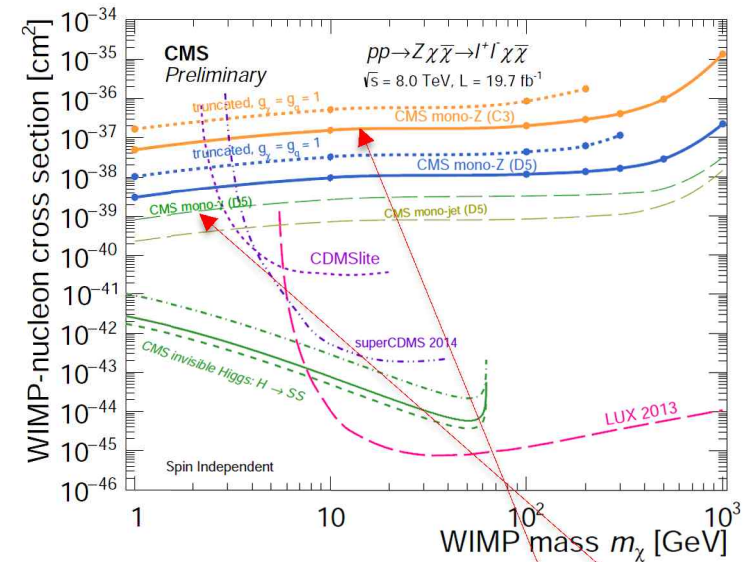


D8

M_χ (GeV)	Expected		Expected -1σ		Expected $+1\sigma$		Observed	
	Λ (GeV)	$\sigma_{\chi-N}$ (cm²)	Λ (GeV)	$\sigma_{\chi-N}$ (cm²)	Λ (GeV)	$\sigma_{\chi-N}$ (cm²)	Λ (GeV)	$\sigma_{\chi-N}$ (cm²)
1	586	8.05×10^{-41}	641	5.63×10^{-41}	535	1.16×10^{-40}	542	1.10×10^{-40}
10	579	2.65×10^{-40}	633	1.85×10^{-40}	528	3.82×10^{-40}	536	3.62×10^{-40}
100	568	3.37×10^{-40}	621	2.36×10^{-40}	518	4.86×10^{-40}	523	4.66×10^{-40}
200	532	4.42×10^{-40}	581	3.09×10^{-40}	485	6.38×10^{-40}	488	6.22×10^{-40}
300	488	6.24×10^{-40}	534	4.37×10^{-40}	445	9.01×10^{-40}	446	8.96×10^{-40}
500	394	1.48×10^{-39}	430	1.04×10^{-39}	359	2.14×10^{-39}	358	2.15×10^{-39}
1000	203	2.10×10^{-38}	222	1.47×10^{-38}	185	3.03×10^{-38}	184	3.07×10^{-38}

D9

M_χ (GeV)	Expected		Expected -1σ		Expected $+1\sigma$		Observed	
	Λ (GeV)	$\sigma_{\chi-N}$ (cm²)	Λ (GeV)	$\sigma_{\chi-N}$ (cm²)	Λ (GeV)	$\sigma_{\chi-N}$ (cm²)	Λ (GeV)	$\sigma_{\chi-N}$ (cm²)
1	2151	4.44×10^{-43}	2352	3.10×10^{-43}	1962	6.40×10^{-43}	1909	7.14×10^{-43}
10	2148	1.40×10^{-42}	2350	9.78×10^{-43}	1960	2.02×10^{-42}	1902	2.28×10^{-42}
100	2113	1.75×10^{-42}	2311	1.23×10^{-42}	1928	2.53×10^{-42}	1884	2.78×10^{-42}
200	2011	2.16×10^{-42}	2199	1.51×10^{-42}	1835	3.12×10^{-42}	1795	3.40×10^{-42}
300	1873	2.88×10^{-42}	2048	2.01×10^{-42}	1709	4.16×10^{-42}	1668	4.58×10^{-42}
500	1570	5.84×10^{-42}	1717	4.09×10^{-42}	1433	8.43×10^{-42}	1401	9.23×10^{-42}
1000	892	5.63×10^{-41}	975	3.94×10^{-41}	813	8.13×10^{-41}	795	8.90×10^{-41}



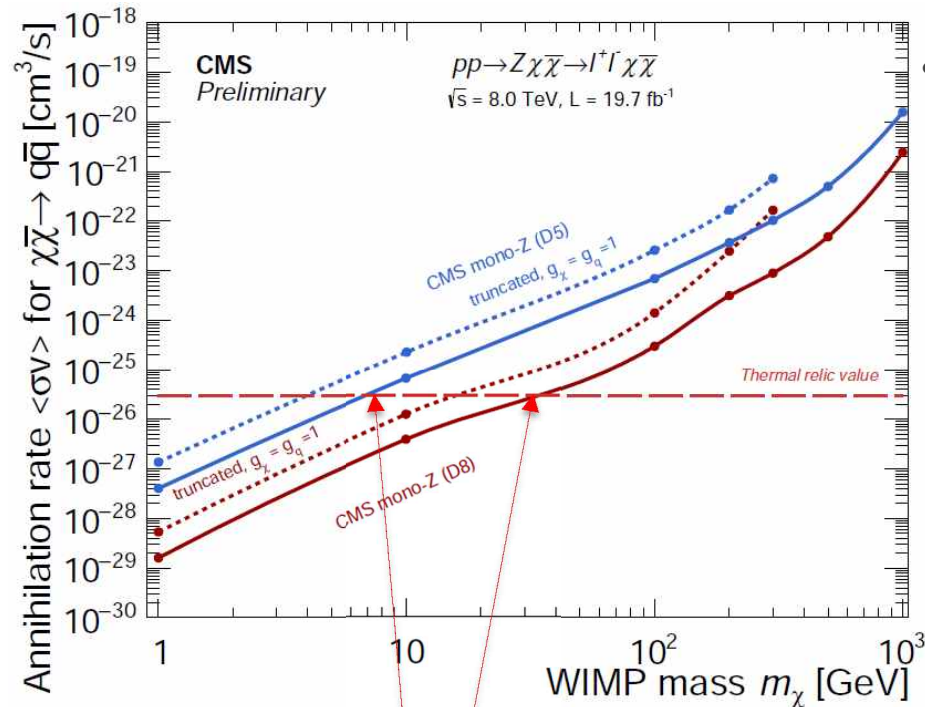
D5

M_χ (GeV)	Expected		Expected -1σ		Expected $+1\sigma$		Observed	
	Λ (GeV)	$\sigma_{\chi-N}$ (cm²)	Λ (GeV)	$\sigma_{\chi-N}$ (cm²)	Λ (GeV)	$\sigma_{\chi-N}$ (cm²)	Λ (GeV)	$\sigma_{\chi-N}$ (cm²)
1	585	2.23×10^{-39}	640	1.56×10^{-39}	534	3.22×10^{-39}	541	3.05×10^{-39}
10	584	7.08×10^{-39}	638	4.95×10^{-39}	532	1.02×10^{-38}	542	9.49×10^{-39}
100	583	8.34×10^{-39}	638	5.83×10^{-39}	532	1.20×10^{-38}	538	1.15×10^{-38}
200	565	9.55×10^{-39}	618	6.68×10^{-39}	515	1.38×10^{-38}	519	1.34×10^{-38}
300	542	1.13×10^{-38}	593	7.88×10^{-39}	495	1.63×10^{-38}	495	1.62×10^{-38}
500	474	1.94×10^{-38}	519	1.35×10^{-38}	433	2.79×10^{-38}	432	2.81×10^{-38}
1000	286	1.46×10^{-37}	313	1.02×10^{-37}	261	2.11×10^{-37}	258	2.21×10^{-37}

C3

M_χ (GeV)	Expected		Expected -1σ		Expected $+1\sigma$		Observed	
	Λ (GeV)	$\sigma_{\chi-N}$ (cm²)	Λ (GeV)	$\sigma_{\chi-N}$ (cm²)	Λ (GeV)	$\sigma_{\chi-N}$ (cm²)	Λ (GeV)	$\sigma_{\chi-N}$ (cm²)
1	411	3.66×10^{-38}	450	2.56×10^{-38}	375	5.28×10^{-38}	382	4.93×10^{-38}
10	412	1.14×10^{-37}	451	7.96×10^{-38}	376	1.64×10^{-37}	382	1.54×10^{-37}
100	403	1.47×10^{-37}	440	1.02×10^{-37}	367	2.11×10^{-37}	373	1.99×10^{-37}
200	371	2.05×10^{-37}	406	1.43×10^{-37}	339	2.95×10^{-37}	341	2.87×10^{-37}
300	342	2.84×10^{-37}	374	1.99×10^{-37}	312	4.10×10^{-37}	314	4.03×10^{-37}
500	278	6.53×10^{-37}	304	4.56×10^{-37}	254	9.42×10^{-37}	253	9.62×10^{-37}
1000	145	8.85×10^{-36}	159	6.18×10^{-36}	132	1.28×10^{-35}	131	1.34×10^{-35}

DM Annihilation Rate limit



- The limits from D5(vector) and D8(axial-vector) translated into upper limits on the dark matter annihilation rate.

- A 100% branching fraction of DM annihilating to quarks is assumed.
- The corresponding truncated limits for D5 and D8 with $\text{sqrt}(g_q g_\chi) = 1$ are also presented.
- The value required for DM to make up the relic abundance labeled “Thermal relic value” is shown as a red dashed line.

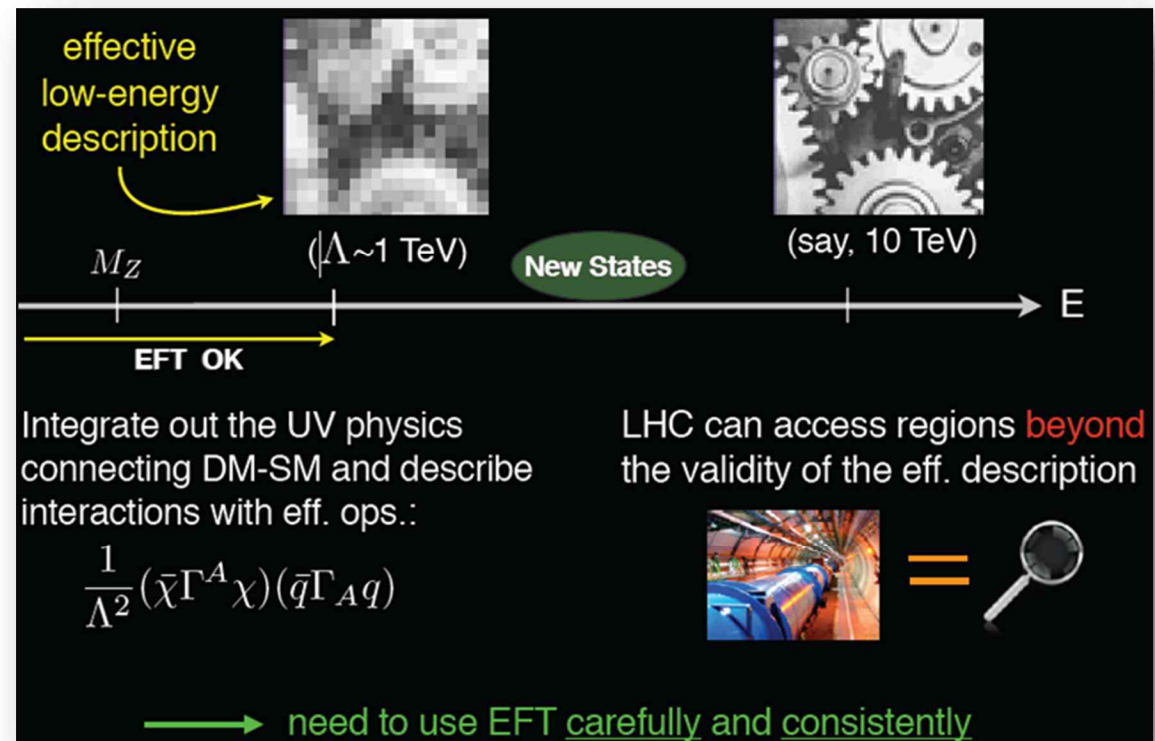
- The annihilation rate of Dirac fermion DM ruled out at 90% CL.
 - $m_\chi < 7 \text{ GeV}$ for vector coupling
 - $m_\chi < 32 \text{ GeV}$ in case of axial-vector coupling.
- Indicate same conclusions in comparison with FNAL-LAT experiment.

Run 1 Summary

- Searches for DM at collider have evolved since the Higgs discovery.
 - The most recent one : Mono-Z(to Dilepton) analysis contains the most tuned EFT
 - Indicate no-longer valid EFT → Requiring new DM model rather than EFT at new energy
- Qualitative picture on DM searches by the LHC is complementary with Direct Detection experiments regardless of decay channels by the LHC.
 - Dominant contribution in Spin-dependent by the LHC
 - Dominant in spin-independent by the Direct Detection experiments but the LHC is better in the region of DM mass less than $\sim 5 - 10$ GeV region.
- Need more reliable model rather than EFT at the new energy LHC era since momentum transfer is very high.

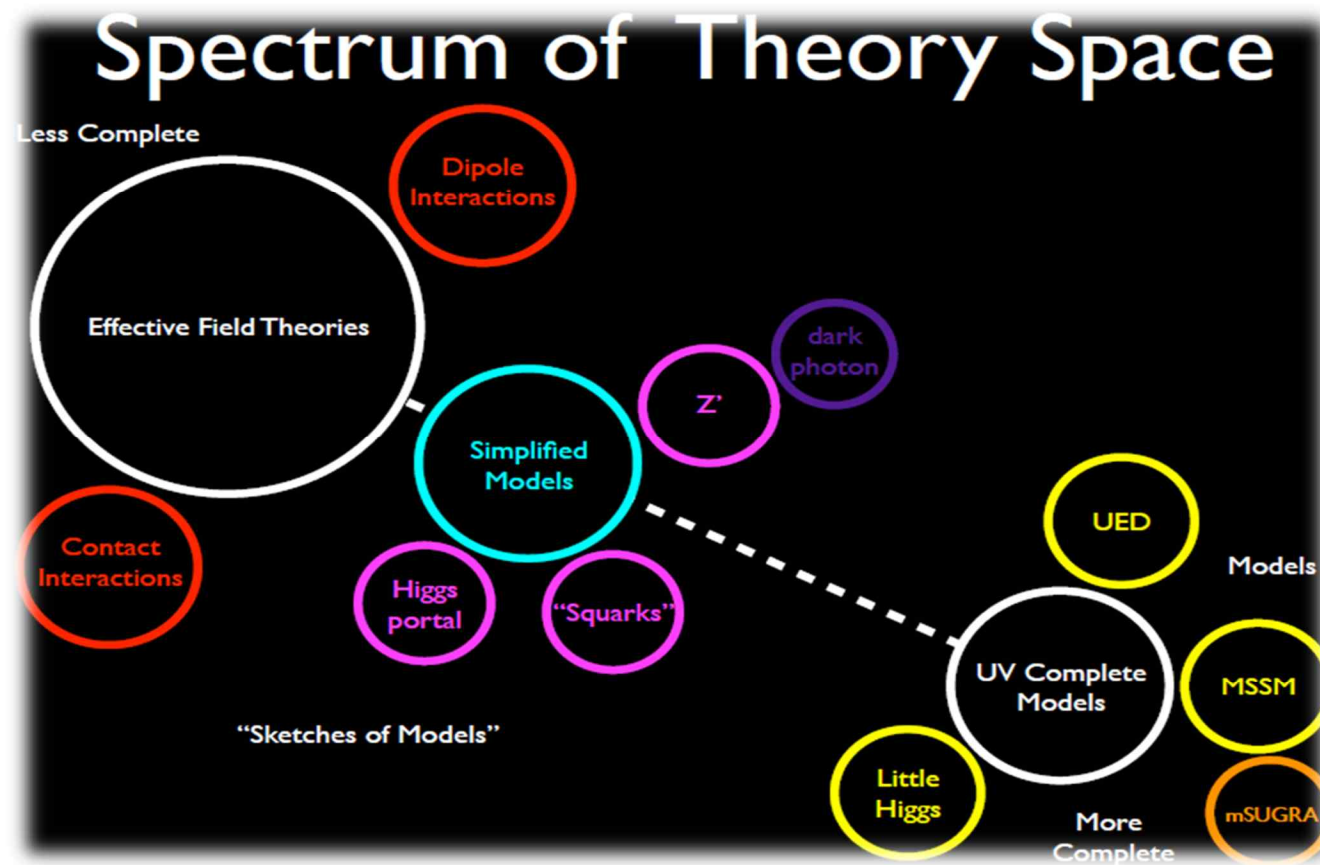
EFT applicable to Run 2 ?

- Since the start-up of the LHC, generic searches for DM production at colliders, and their comparison with direct-detection experiments such as LUX, have become a focal point for both the experimental and theoretical particle and astro-particle communities.
- A comparison of DM searches at collider and DD experiments using the EFT approach does not provide an accurate description of the complementarity of the two search strategies .
- How to go beyond the problematic comparison of DM searches in the EFT framework ??
- An alternative to the EFT interpretation is **the Simplified Models**.
- Simplified models are able to capture properly the relevant kinematic properties of collider searches with only a few free parameters.

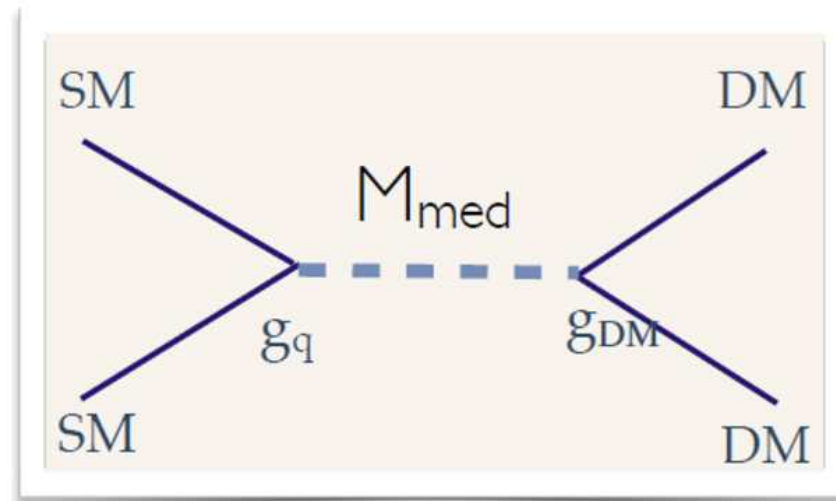


Minimal Simplified DM Model

- The model where only the particles and parameters most relevant to the search are included.
- The Minimal Simplified DM(MSDM) model can be understood as a phenomenological sketch of a complete mode.



MSDM Model



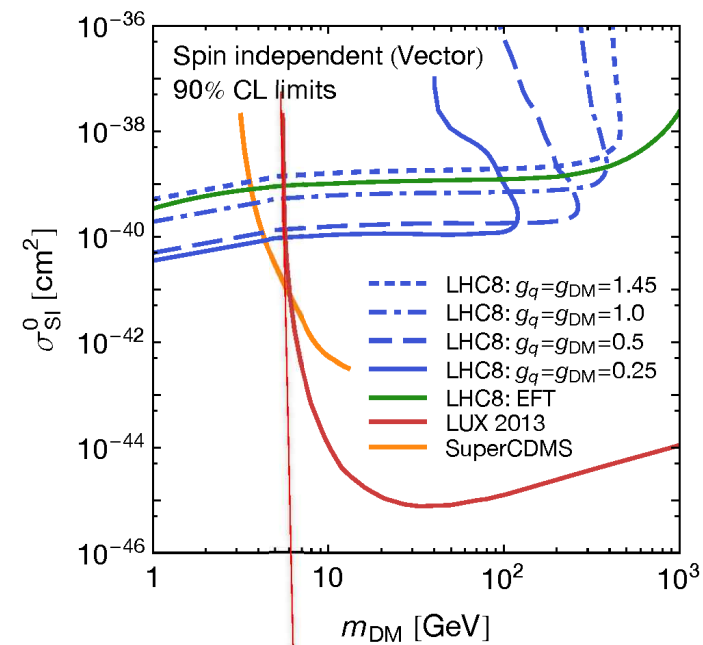
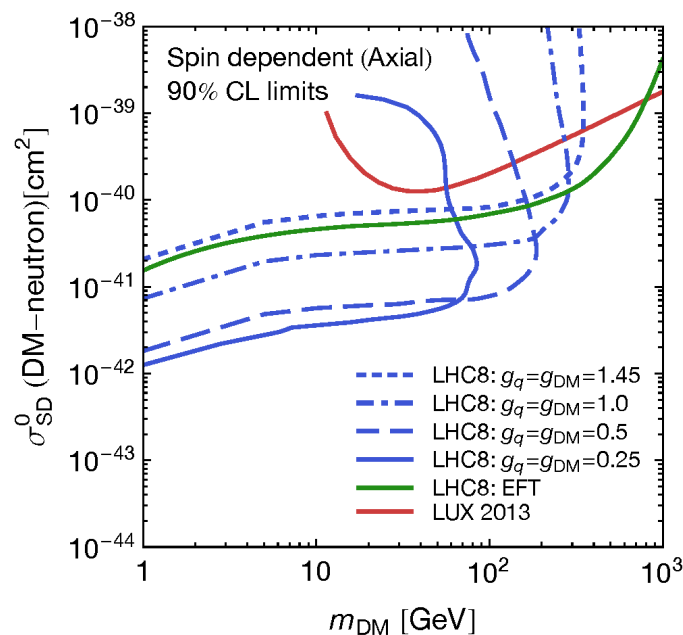
Define simplified model with (minimum) 4 parameters		DM		Consider comprehensive set of diagrams for mediator	
Mediator mass (M_{med})	DM mass (M_{DM})	Dirac fermion	Scalar - real	Vector	Axial-vector
g_q	g_{DM}	Majorana fermion	Scalar - complex	Scalar	Pseudoscalar

$$g_q = g_{\text{DM}} = [0.25, 0.5, 1.0, 1.45] \quad M_{\text{DM}} \text{ vs } M_{\text{med}} \text{ assuming } g_q = g_{\text{DM}}$$

S. Malik et. al, arXiv:1409.4075 (2014) : White Report

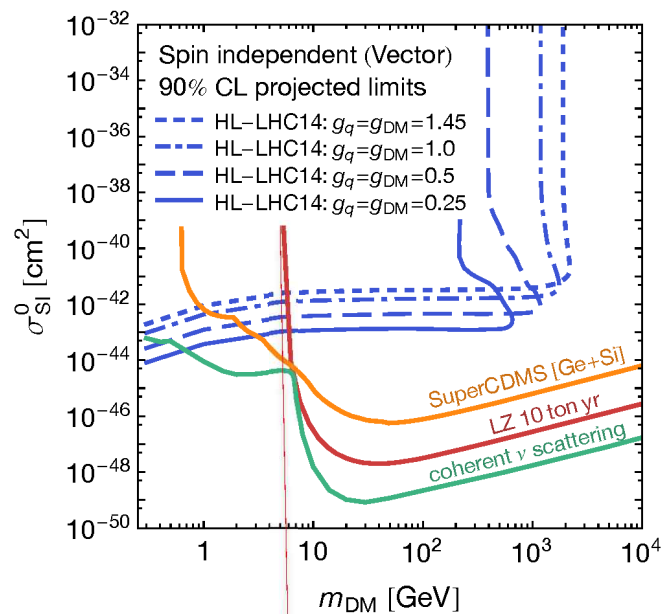
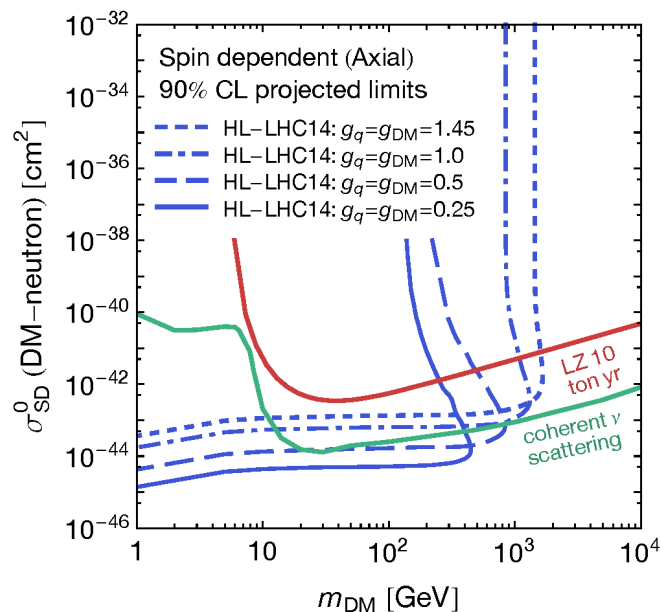
Dark Matter Searches Results(Run 1)

- The most recent comparison of the LHC and underground experiments
 - LUX : Current world best limit among underground experiments(350 kg Xe)
 - CMS : Mono-jet channel at 8 TeV
 - The LHC results are much better in the case of Spin dependent(Axial)
 - But the LHC is only sensitive in low DM mass(<5 GeV) region in the case of Spin-independent



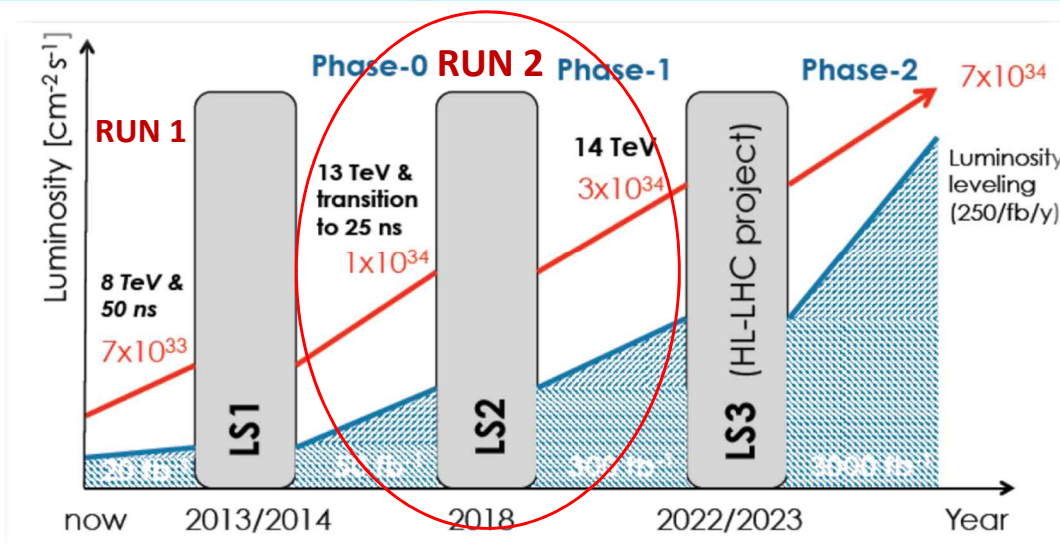
Prospect of Run 2

- Expectations with 14 TeV LHC
 - Comparison with future underground experiments
 - LZ(LUX-ZEPLIN : 7 ton planned) and SuperCDMS
 - Show much better results(10 -100 times more in x-section) with the LHC for spin-dependent case
 - LHC results sensitive in very low dark mass region for spin-independent case



Collider Dark Matter Search is complementary with underground experiments

The LHC Run 2 Original Plan

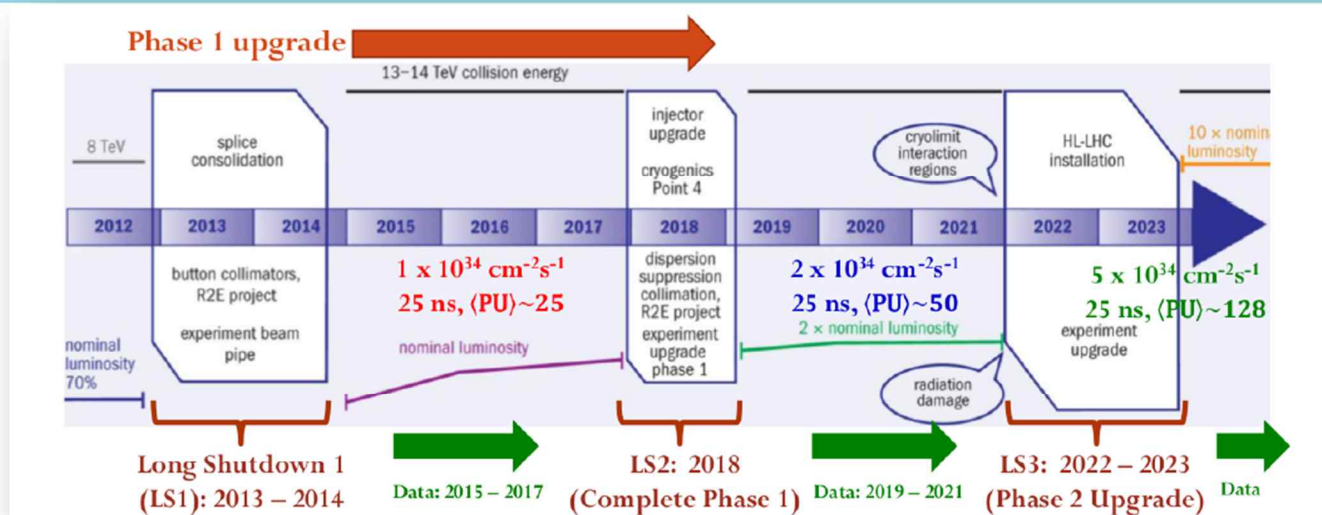


The most important period in Particle Physics Field

Period	Activity	Label	Peak \mathcal{L}	Int \mathcal{L} (fb ⁻¹)	\sqrt{s} (TeV)	Bunch (ns)
2013-2014	Install Phase 0	LS1				
2015-2017	Running		10^{34}	~100	~14	25
2018	Install Phase I	LS2				
2019-2021	Running		$\sim 2 \times 10^{34}$	~300	~14	25
2022-2023	Install Phase II	LS3				
2024-2036	Running		$\sim 5 \times 10^{34}$	~3000	~14	25

The Run 2 is main while Run 1 was like introductory even Higgs had been found.

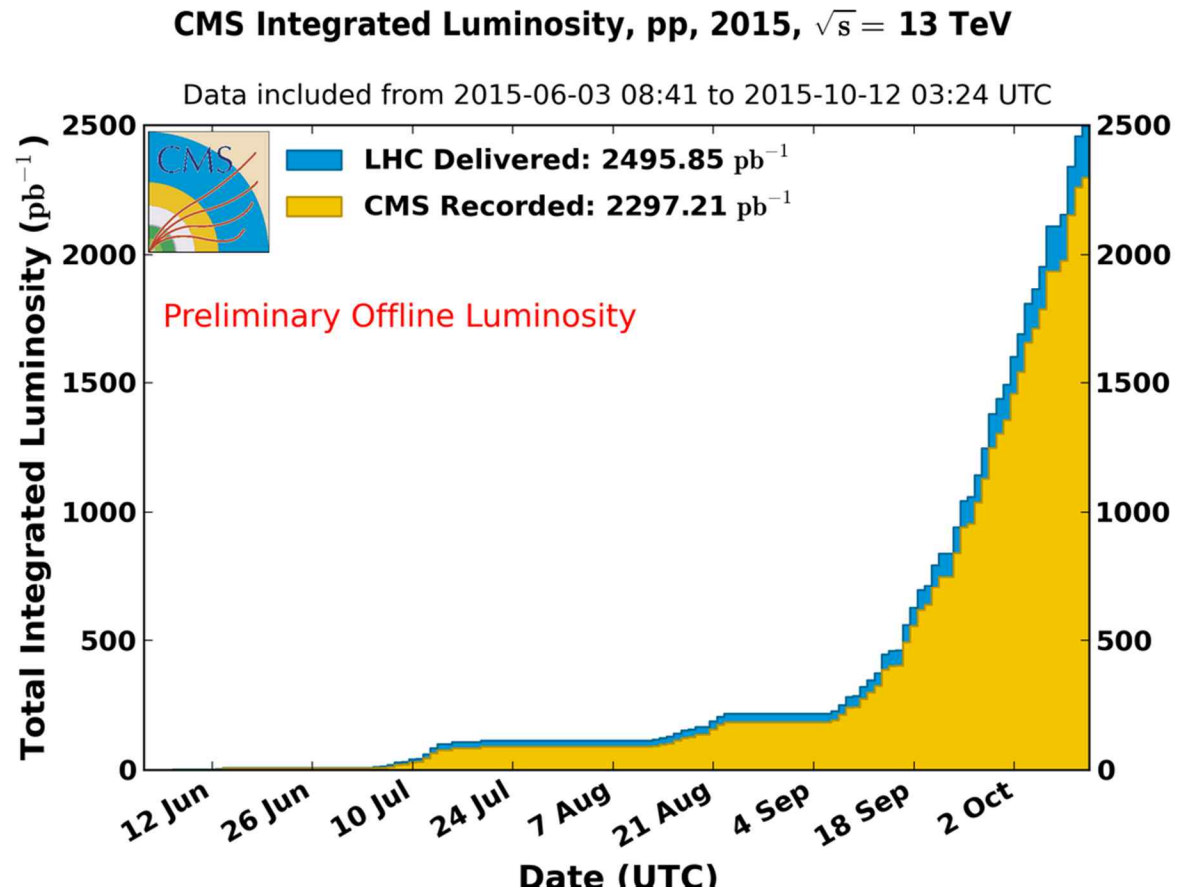
Physics Milestone(Run 2 & 3)



- Run 2 (Phase 1: 2015-2021) : 13 and 14 TeV, $\sim 2 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$, 300 fb⁻¹
 - Doubling of new particle mass reach : SUSY and(/or) other unknown phenomena
 - Potential searches of dark matters
 - Higgs Properties : mass, width, couplings and even self-coupling
 - Looking for anomalies through vector-vector scattering
 - Measurements of rare SM processes
- Run 3(Phase 2: 2024-2036) : 14 TeV, $\sim 5 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$, 3000 fb⁻¹
 - Characterization of new particle discovered at run 2, or
 - Extension of new particle mass reach : SUSY and(/or) Exotic discovery
 - Precision measurements of Higgs properties : couplings (< 5%), self-coupling ($\sim 20\%$) etc.
 - Precision Measurements of rare SM process

Current Status

- As of Oct 12, 2015, $\sim 2.3 \text{ fb}^{-1}$ recorded.
- Some fraction of data with no magnetic field on ($\sim 20\%$)
- Data delivered exponentially (similar situation with year 2011 run)



- First results were presented at EPS-HEP conference in July 2015.
- First paper already submitted to PLB for charged Hadron Production.
- All the important physics topics are underway of analysis mode.

Summary

- The LHC experiments are compatible and complementary with the underground experiments.
- The LHC is very strong with spin-dependent DM case and low mass($\sim < 5$ to 10 GeV) DM region in spin-dependent case.
- There had been a development of Minimal Simplified DM model in order to compare with direct detection experiments in parallel.
- We have many interesting results on DM searches at the LHC from Run 1, but have not yet found a signal.
- Dark Matter searches at the LHC are hot topics at the LHC Run 2.
- Hopefully the DM will be unveiled at the LHC Run 2 period.

Summary

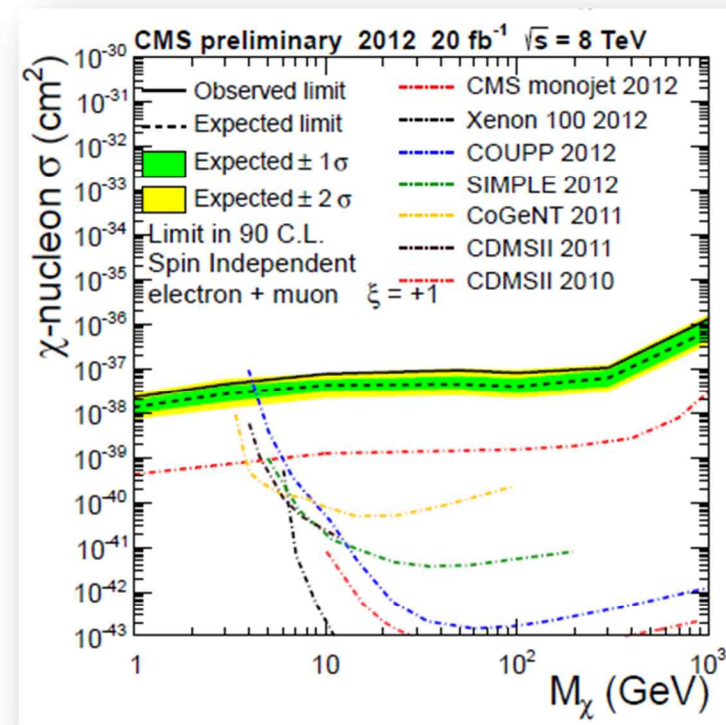
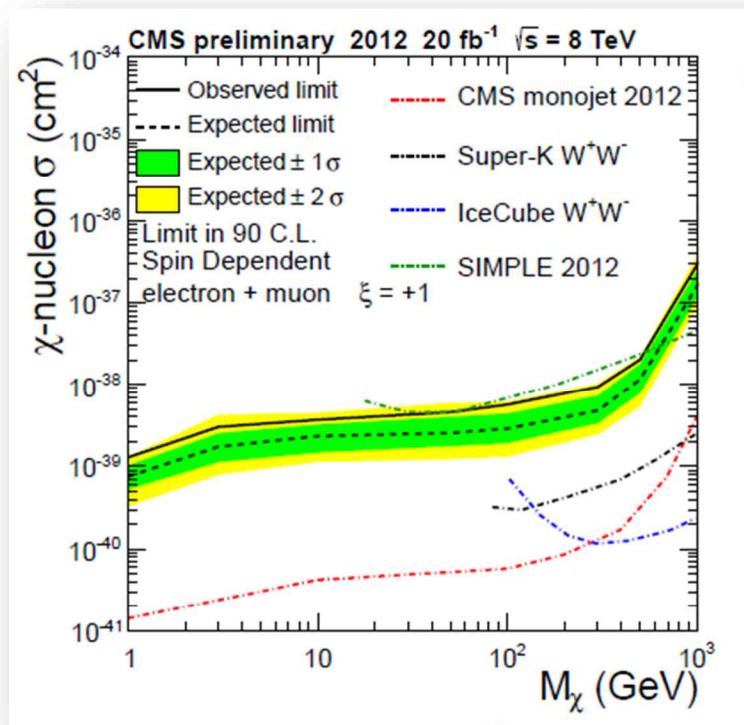
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- Dark Matter searches at the LHC are hot topics at the LHC Run 2.
- Hopefully the DM will be unveiled at the LHC Run 2 period.

‘ The truth is only an agreement between a cognitive spirit and Things.
Well, Nature never surrenders against any traditions by human being.’

BACKUP

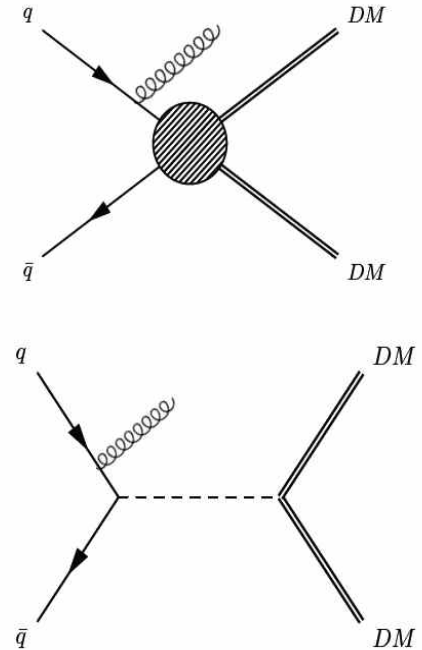
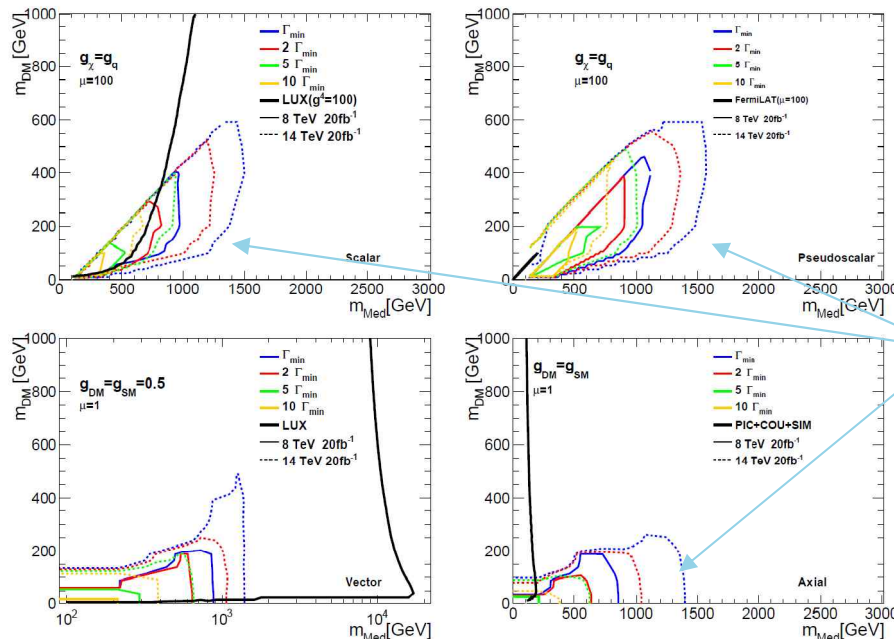
Comparisons with the Direct Detection Expts

- Direct detection much more sensitive on SI(V) coupling for $M_\chi > 5\text{GeV}$
- Strong in SD(axial) mode :
- very similar with mono-jet case but better for the case constructive interference



Mono-jet + missing energy

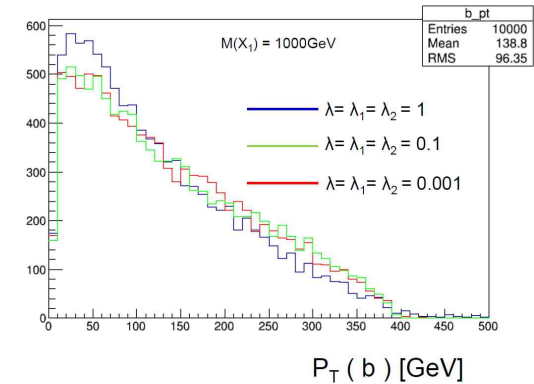
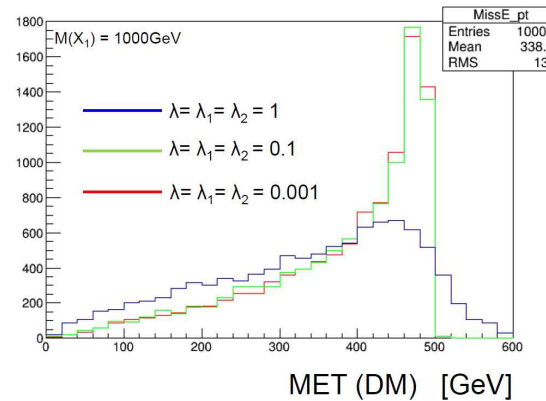
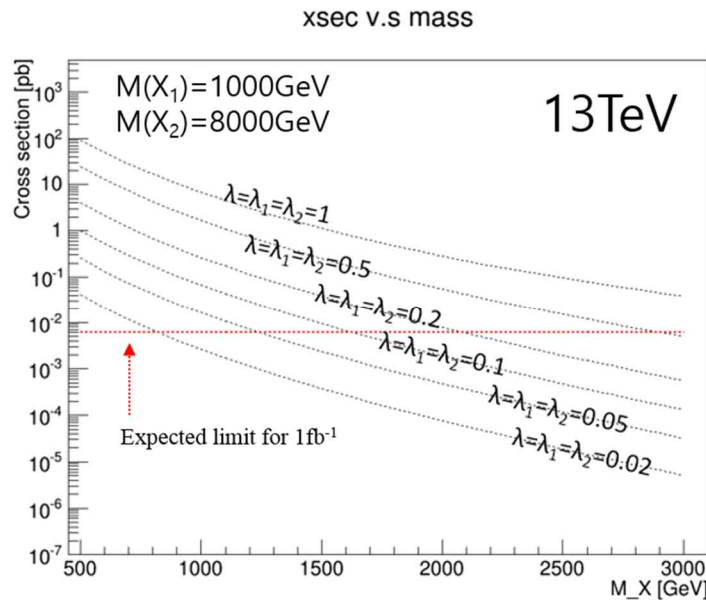
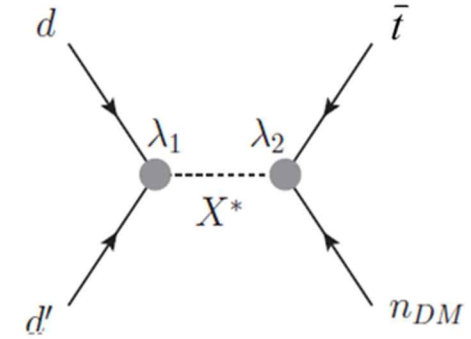
- Signature
 - One light jet(gluon initiated) with large missing energy from undetected particle
- Dark matter production scenarios
 - Contact interaction
 - Mediator is heavier than the typical energy transfer at the LHC
 - The framework of effective field theory is used
 - mediator can be scalar, pseudo-scalar, vector or axial vector field.
 - Exchange of mediator
 - Mediator is light enough to be produced at the LHC



- LHC 14TeV Run can probe much wider region than any other dark matter experiments !

Mono-top Studies at new energy

- Mono-top + missing energy channel
 - Top quark + dark matter production via resonance of scalar particle is considered.
 - Production cross section of pp collision at 13 TeV are calculated.
 - Kinematical properties are studied to setup the trigger and selection scheme for data analysis.



- Confirmed the mono-top channel can be probed for the dark matter
- To be analyzed with the CMS new energy data