

Radiative Neutrino Mass via Fermion Kinetic Mixing in the $SU(3)_c \times SU(2)_L \times U(1)_Y \times U(1)_D$ Model

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Overview

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Introduction/Motivation

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- Generate neutrino masses
- Naturely accomodate DM
- Connect the existance of DM to non-zero neutrino masses

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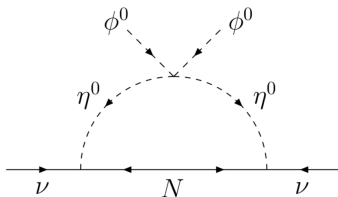
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- See-saw-1, Add SM singlet fermion, $m_\nu = -\frac{m_D^2}{M_N}$
- See-saw-2, add SM triplet scalar, $m_\nu = -\frac{f\mu\nu}{M^2}$
- See-saw-3, add SM triplet fermion



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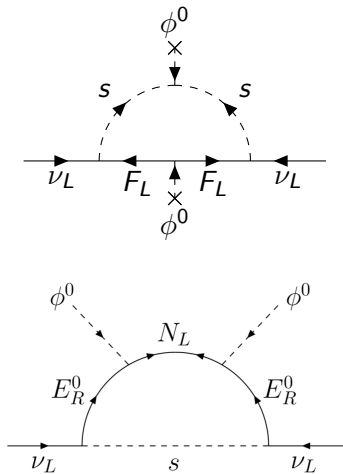
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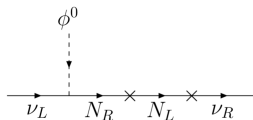
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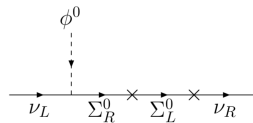


Brief review

- Insert a Dirac fermion singlet N which does not transform under \mathcal{S} , then break \mathcal{S} softly by the dimension-three $\bar{\nu}_R N_L$ term.

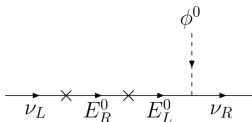


- Insert a Dirac fermion triplet $(\Sigma^+, \Sigma^0, \Sigma^-)$ which does not transform under \mathcal{S} , then break \mathcal{S} and $SU(2)_L \times U(1)$ together spontaneously to obtain the dimension-three $\bar{\nu}_R \Sigma_L^0$ term.

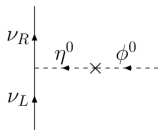


Brief review

- Insert a Dirac fermion doublet (E^0, E^-) which transforms as ν_R under \mathcal{S} , then break \mathcal{S} softly by the dimension-three $(\bar{E}^0\nu_L + E^+e^-)$ term.



- Insert a scalar doublet (η^+, η^0) which transforms as ν_R under \mathcal{S} , then break \mathcal{S} softly by the dimension-two $(\eta^- \phi^+ + \bar{\eta}^0 \phi^0)$ term.



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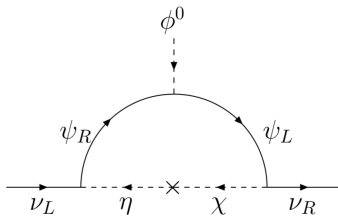
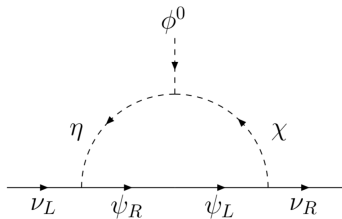
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Model particle content

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Field	$SU(3)_c$	$SU(2)_L$	$U(1)_Y$	$U(1)_D$	Flavors	copies
$Q \sim (u, d)^T$	3	2	$\frac{1}{6}$	0	3	1
u^c	$\bar{3}$	1	$-\frac{2}{3}$	0	3	1
d^c	$\bar{3}$	1	$\frac{1}{3}$	0	3	1
$L \sim (\nu, e)^T$	1	2	$-\frac{1}{2}$	0	3	1
e^c	1	1	1	0	3	1
$H \sim (H^+, H_0)^T$	1	2	$\frac{1}{2}$	0	1	1
A_L	1	1	0	3	3	1
C_L	1	1	0	1	3	5
N_L	1	1	0	-4	3	1
N_R^c	1	1	0	4	3	1
F_{2L}	1	1	0	-2	3	4
Ψ_L	1	1	0	$\frac{5}{2}$	3	1
Ψ_R^c	1	1	0	$-\frac{5}{2}$	3	1
$\eta_L \sim (\eta^0, \eta^{-1})^T$	1	2	$-\frac{1}{2}$	3	1	1
η_D	1	1	0	-1	1	1
ϕ	1	1	0	2	1	1
s_7	1	1	0	$\frac{7}{2}$	1	1
s_{11}	1	1	0	$-\frac{11}{2}$	1	1

Lagrangian

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$$-\mathcal{L}_{Yuk}^{SM} = \bar{u}_{Ra} Y_u^{ab} Q_{Lbi} H_j \epsilon^{ij} + \bar{d}_{Ra} Y_d^{ab} Q_{Lbi} H^{\dagger i} + \bar{l}_a Y_e^{ab} e_{Rb} H_i + \text{h.c.}$$

$$-\mathcal{L}_{m_F}^{New} = \bar{\Psi}_{La} m_{\Psi}^{ab} \Psi_{Rb} + \bar{N}_{La} m_N^{ab} N_{Rb} + \text{h.c.}$$

$$\begin{aligned} -\mathcal{L}_{Yuk}^{New} = & L_{ai} Y_L^{ab} A_{Lb} \eta_L^{\dagger i} + \Psi_{La} Y_A^{ab} A_{Lb} s_{11} + \Psi_{La} Y_C^{ab\alpha} C_{Lb\alpha} s_7^* \\ & + A_{La} Y_N^{ab} N_{Lb} \eta_D^* + C_{La\alpha} Y_{cc}^{ab\alpha\beta} C_{Lb\beta} \phi^* + A_{La} Y_{AF}^{ab\alpha} F_{2Lb\alpha} \eta_D \\ & + C_{La\alpha} Y_{CF}^{ab\alpha\beta} F_{2Lb\beta} \eta_D^* + \bar{F}_{2La\alpha} Y_{NF}^{ab\alpha} N_{Rb} \phi + \text{h.c.} \end{aligned}$$

$$V_0 = \sum_{\substack{H, \phi, \eta_L, \eta_D, \\ s_{11}, s_7 \in X}} \left((-1)^{q_x} m_x^2 |x|^2 + \frac{\lambda_x}{2} |x|^4 \right) + \sum_{\substack{H, \phi, \eta_L, \eta_D, s_{11}, \\ s_7 \in \{x < y\}}} \lambda_{xy} |x|^2 |y|^2$$

$$V_3 = \mu_D \eta_D^2 \phi + \mu_3 \phi s_{11} s_7 + \text{h.c.}$$

$$V_4 = \lambda_{H\eta\phi} H_i \eta_{Lj} \eta_D \phi^* \epsilon^{ij} + \lambda_{s\eta} \eta_D^2 s_7^* s_{11}^* + \text{h.c.}$$

Triangle Anomalies

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$$\sum_i Q_{Di} = 1 \times (3) + 5 \times (1) + 4 \times (-2) + 1 \times (-4) + 1 \times (4)$$

$$+ 1 \times \left(\frac{5}{2}\right) + 1 \times \left(-\frac{5}{2}\right) = 0$$

$$\sum_i Q_{Di}^3 = 1 \times (3)^3 + 5 \times (1)^3 + 4 \times (-2)^3 + 1 \times (-4)^3 + 1 \times (4)^3$$

$$+ 1 \times \left(\frac{5}{2}\right)^3 + 1 \times \left(-\frac{5}{2}\right)^3 = 0$$

Radiative neutrino mass generation

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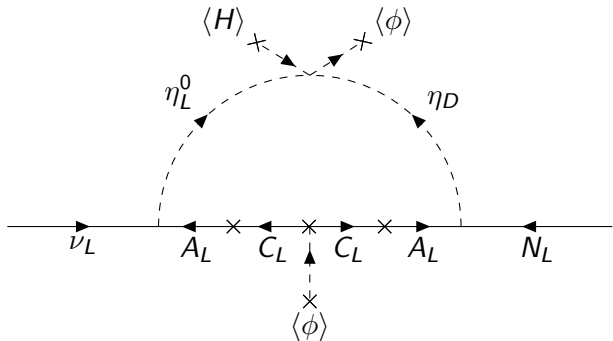
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$$m_\nu = \frac{1}{16\pi^2} Y_N m_{F_2} Y_L \frac{\epsilon}{1-\epsilon^2} F(m_{S_i}, m_{F_2}, \theta_{\eta_L, \eta_D})$$

Radiative Kinetic Mixing of Fermions

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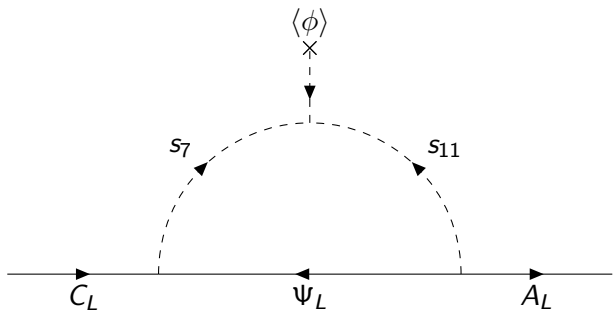
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$$\begin{aligned}\epsilon &= \frac{1}{16\pi^2} Y_A^* \gamma(m_{s_i}, m_\Psi, \theta_{s_7, s_{11}}) Y_C e^{-i2\Delta} \\ \begin{pmatrix} A_L \\ C_L \end{pmatrix} &= U(\pi/4, \Delta)^\dagger R_{\text{rescale}}^{-1} U(\alpha)^\dagger \begin{pmatrix} F_{1L} \\ F_{2L} \end{pmatrix} = \\ &\begin{pmatrix} 1 & -\frac{\epsilon}{\sqrt{1-\epsilon^2}} \\ 0 & \frac{\epsilon}{\sqrt{1-\epsilon^2}} \end{pmatrix} \begin{pmatrix} F_{1L} \\ F_{2L} \end{pmatrix} \\ \tan(2\alpha) &= \frac{-\sqrt{1-\epsilon^2}}{\epsilon}\end{aligned}$$

Hidden Sector Fermion Masses

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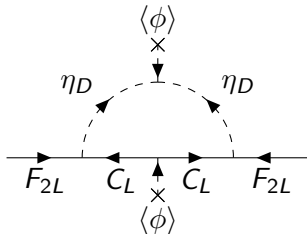
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- $m_{AD}^2 = \frac{g_d^2}{2} (2\nu_\phi)^2$
- 5×5 C_L mass matrix generated by $C_{La\alpha} Y_{cc}^{ab\alpha\beta} C_{Lb\beta} \phi^*$
- $\begin{pmatrix} 0 & m_\nu & 0 & 0 \\ m_\nu & 0 & m_N & 0 \\ 0 & m_N & 0 & Y_{NF} \nu_\phi \\ 0 & 0 & Y_{NF} \nu_\phi & 0 \end{pmatrix}$ in the $(\nu_L, N_L, N_R^c, F_{2Li})$ bases.



Mass of A dark fermion

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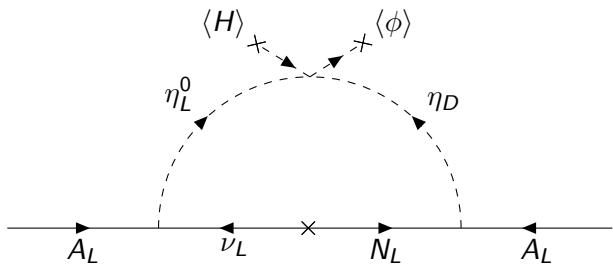
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Scalar Masses

$$\blacksquare m_{H\phi R} = \begin{pmatrix} \lambda_H \nu^2 & \lambda_{H\phi} \nu \nu_\phi \\ \lambda_{H\phi} \nu \nu_\phi & \lambda_\phi \nu_\phi^2 \end{pmatrix}$$

$$\blacksquare m_{\eta LD[R/I]} = \begin{pmatrix} m_{\eta L}^2 + \frac{\lambda_{HL} \nu^2}{2} + \frac{\lambda_{\phi L} \nu_\phi^2}{2} & \mp \frac{\lambda_{H\eta\phi} \nu \nu_\phi}{2} \\ \mp \frac{\lambda_{H\eta\phi} \nu \nu_\phi}{2} & m_{\eta D}^2 + \frac{\lambda_{HD} \nu^2}{2} + \frac{\lambda_{\phi D} \nu_\phi^2}{2} \pm \sqrt{2} \mu_D \nu_\phi \end{pmatrix}$$

$$\blacksquare m_{s[R/I]} = \begin{pmatrix} m_{s7}^2 + \lambda_{H7} \nu^2 + \lambda_{\phi 7} \nu_\phi^2 & \pm \frac{\mu_3 \nu_\phi}{\sqrt{2}} \\ \pm \frac{\mu_3 \nu_\phi}{\sqrt{2}} & m_{s11}^2 + \lambda_{H11} \nu^2 + \lambda_{\phi 11} \nu_\phi^2 \end{pmatrix}$$

Dark matter candidates

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- 2 dark separated sectors
- Q_D odd charged particles: A, C, η_L, η_D (equivalent to canonical Scotogenic case)
- Particles with fractional Q_D charges: Ψ, s_7, s_{11} (Needed for kinetic mixing), like dark sector within dark sector

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- generate naturally small neutrino masses
- naturally accommodate multilayer DM

Thank you!