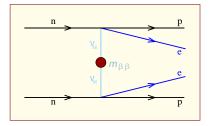
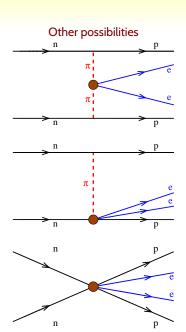


# Review of $Ov\beta\beta$ Decay

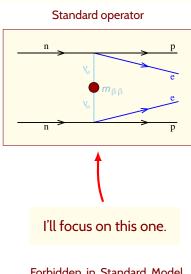
Standard operator



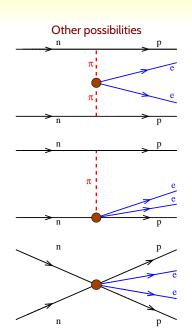


Forbidden in Standard Model. New physics inside blobs.

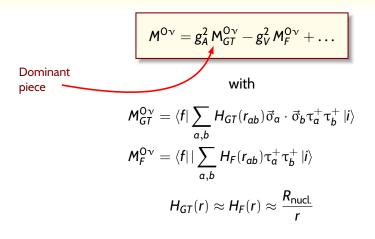
# Review of $Ov\beta\beta$ Decay



Forbidden in Standard Model. New physics inside blobs.



## Nuclear Matrix Element (Simplified)



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Dominant piece with 
$$M_{GT}^{O\nu} = g_A^2 M_{GT}^{O\nu} - g_V^2 M_F^{O\nu} + \dots$$

$$M_{GT}^{O\nu} = \langle f | \sum_{a,b} H_{GT}(r_{ab}) \vec{\sigma}_a \cdot \vec{\sigma}_b \tau_a^+ \tau_b^+ | i \rangle$$

$$M_F^{O\nu} = \langle f | | \sum_{a,b} H_F(r_{ab}) \tau_a^+ \tau_b^+ | i \rangle$$

$$H_{GT}(r) \approx H_F(r) \approx \frac{R_{\text{nucl}}}{r}$$

Also:

$$M_{2\nu} = g_A^2 \sum_m rac{\langle f | \sum_{\alpha} \vec{\sigma}_{\alpha} au_{\alpha}^+ | m 
angle \cdot \langle m | \sum_{b} \vec{\sigma}_{b} au_{b}^+ | i 
angle}{E_m - rac{E_f + E_i}{2}}$$

## **Nuclear Matrix Element (Simplified)**

Dominant piece with 
$$M^{OV} = g_A^2 M_{GT}^{OV} - g_V^2 M_F^{OV} + \dots$$
But the idea that there is a single " $g_A$  in medium" is too much of a simplification.
$$H_{GT}(r) \approx H_F(r) \approx \frac{R_{\text{nucl.}}}{r}$$

Also:

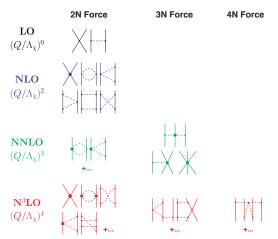
$$M_{2\nu} = g_A^2 \, \sum_m \, \frac{\langle f | \sum_a \, \vec{\sigma}_a \tau_a^+ \, | m \rangle \cdot \langle m | \sum_b \, \vec{\sigma}_b \tau_b^+ \, | i \rangle}{E_m - \frac{E_f + E_i}{2}} \label{eq:M2v}$$

#### Ab Initio Nuclear Structure

Often starts with chiral effective-field theory

Nucleons, pions sufficient below chiral-symmetry breaking scale. Expansion of operators in powers of  $Q/\Lambda_{\chi}$ .

 $Q=m_\pi$  or typical nucleon momentum.



All require many CPU-hours.

Quantum Monte Carlo in light nuclei: More or less exact solution of many-body Schrödinger equation.

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- Coupled-clusters ansatz:

$$|\Psi
angle = \exp\left(t_{ij}^1 a_i^\dagger a_j + t_{ijkl}^2 a_i^\dagger a_j^\dagger a_k a_l + \dots\right) | \text{Slater det.} 
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 In-medium similarity-renormalization group: Flow equations that gradually decouple low-lying states (will explain this).

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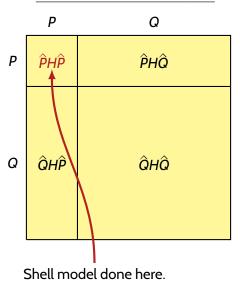
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**)** 

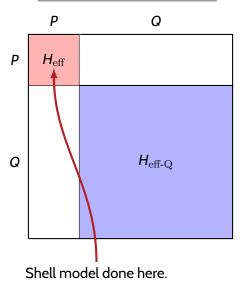
#### Partition of Full Hilbert Space



P = valence spaceQ = the rest

Task: Find unitary transformation to make H block-diagonal in P and Q, with  $H_{\rm eff}$  in P reproducing d most important eigenvalues.

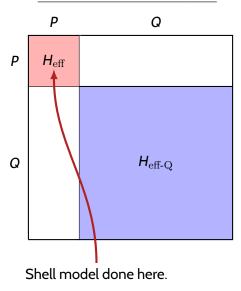
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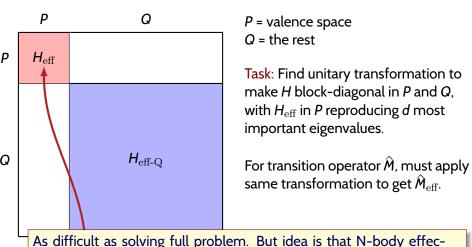


P = valence spaceQ = the rest

Task: Find unitary transformation to make H block-diagonal in P and Q, with  $H_{\rm eff}$  in P reproducing d most important eigenvalues.

For transition operator  $\hat{M}$ , must apply same transformation to get  $\hat{M}_{\rm eff}$ .



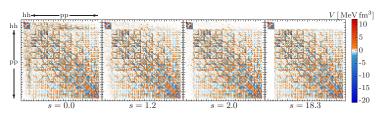


As difficult as solving full problem. But idea is that N-body effective operators beyond N >2 or 3 can be treated approximately.

## In-Medium Similarity Renormalization Group

One way to determine the transformation

Flow equation for effective Hamiltonian. Gradually decouples shell-model space.



Hergert et al.

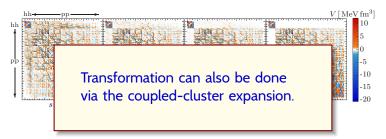
Trick is to keep all 1- and 2-body terms in *H* at each step *after normal ordering* (approximate treatment of 3-, 4-...terms.

If shell-model space contains just a single state, approach yields ground-state energy. If it is a typical valence space, result is effective interaction and operators.

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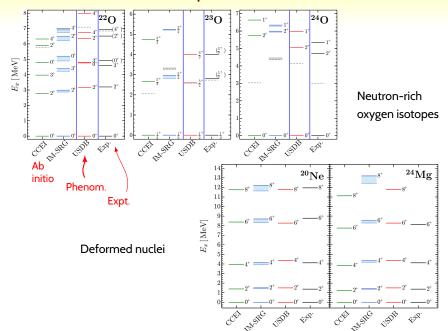
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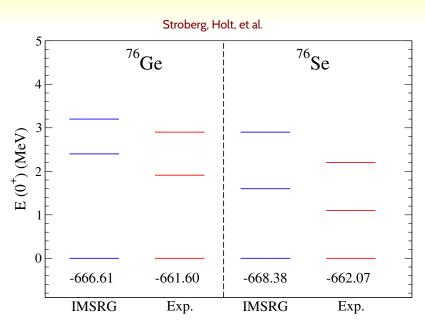
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#### Ab Initio Calculations of Spectra



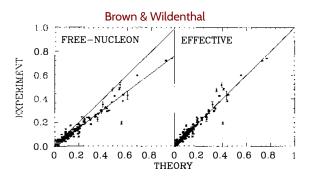
## Ab Initio <sup>76</sup>Ge and <sup>76</sup>Se



## Gamow-Teller β Decay

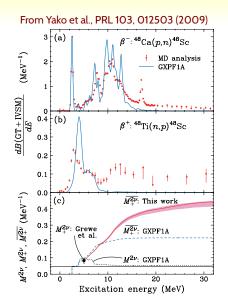
Leading order decay operator is  $\vec{\sigma}\tau_+$ .

50-Year-Old Problem: Effective  $g_A$  needed in all calculations of shell-model (or related) type.



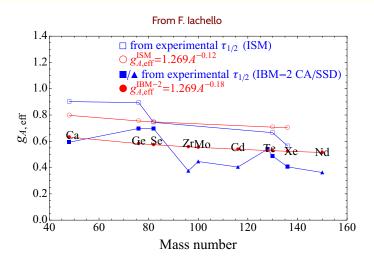
Many suggestions about the cause but, until recently, no consensus.

## Other Tests of $\vec{\sigma}\tau$ Strength Also Show Suppression



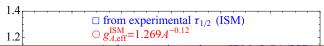
Only about 2/3 of theoretically expected strength observed.

## And $2\nu\beta\beta$ Decay...



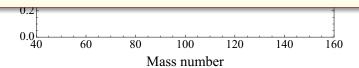
## And $2\nu\beta\beta$ Decay...





What explains all the over-prediction of matrix elements?

In ab initio calculation with chiral EFT, the answer must be a combination of many-body approximations and truncation of chiral expansion of current operator.



#### **Axial Weak Current in Chiral EFT**

 $\beta$  Decay (simplified) with electron lines omitted

#### Leading order:



Usual  $\beta$ -decay current. Finite-momentum corrections at next order.

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#### Higher order:

$$\begin{array}{c|c}
n/p & p & p \\
\hline
 & \pi & p & p \\
\hline
 & \Delta + \dots \\
\hline
 & \alpha_1, p & n & p \\
\hline
 & \alpha_1, p & n & p \\
\hline
 & \alpha_2, \alpha_3, \alpha_4$$

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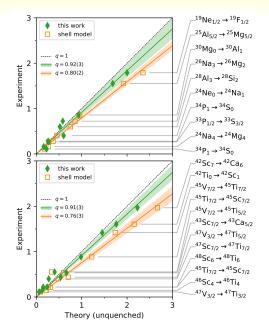
Higher order:

$$n/p$$
 $n/p$ 
 $n/p$ 

Coefficients same as in three-body interaction:

$$\begin{array}{c|cccc}
n/p & p & p \\
\hline
\pi & \pi \\
n/p & n & n \\
\hline
c_3, c_4
\end{array}$$

# Quenching in the sd and pf Shells



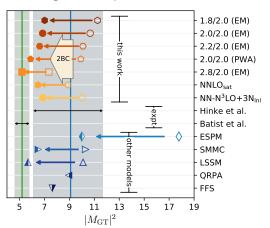
IMSRG calculation, Holt et al, preliminary

Shell model seems to include most correlations. Bulk of quenching comes from two-body current.

#### ...And in 100Sn

#### Coupled-Cluster Calculation of $\beta$ Decay

Hagen et al, unpublished

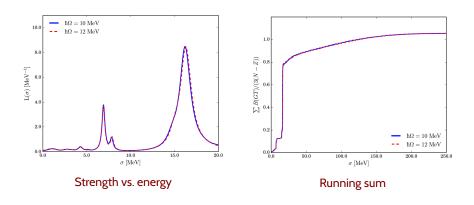


Again, good part of the quenching accounted for by two-body current.

Quenching increases with mass, at least up to Sn.

Spectator nucleons contribute coherently to two-body current.

# Gamow-Teller Strength in <sup>132</sup>Sn

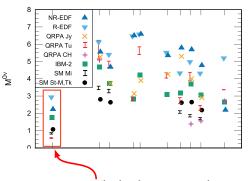


Almost 20% of strength above 30 MeV and 10% above 50 MeV.

## And $Ov\beta\beta$ Decay?

Preliminary results in 48 Ca

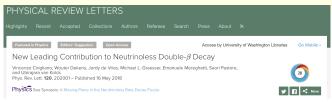
Two-body currents not yet included, but preliminary indications are that their effects are not as large as in  $\beta$  decay.



Coupled -clusters result is red bar at bottom.

But <sup>48</sup>Ca is not typical. <sup>76</sup>Ge coming soon.

## Small Fly in the Ointment



#### Usual light neutrino exchange:



must be supplemented, at same order in chiral EFT, by short-range operator (representing high-energy  $\nu$  exchange):

Coefficient of this term is unknown.

Looking for ways to fit to, e.g., pion double-charge exchange



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# That's all; thanks.

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