

Reconstructing the inflationary landscape

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Based on 1804.07315 & 1806.05202

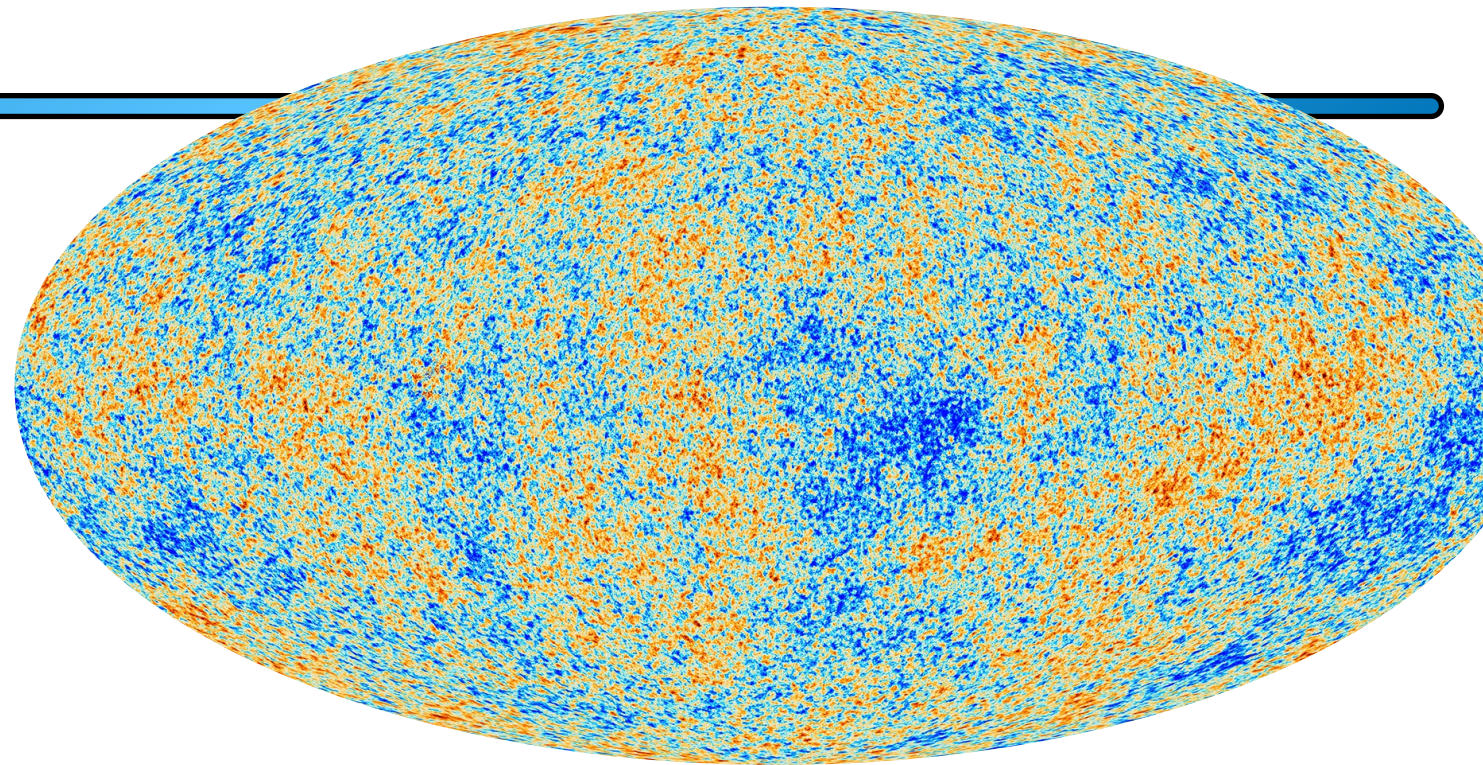
Work with Xingang Chen (Harvard)
Walter Riquelme, Bruno Scheihing & Spyros Sypsas (UCH)

COSMO18



Preamble

Inflation:



- **An early stage of accelerated expansion**
- **It explains the origin of the primordial fluctuations that seeded the CMB**
- **The simplest models of inflation predict a scale invariant and nearly Gaussian distribution of primordial fluctuations**

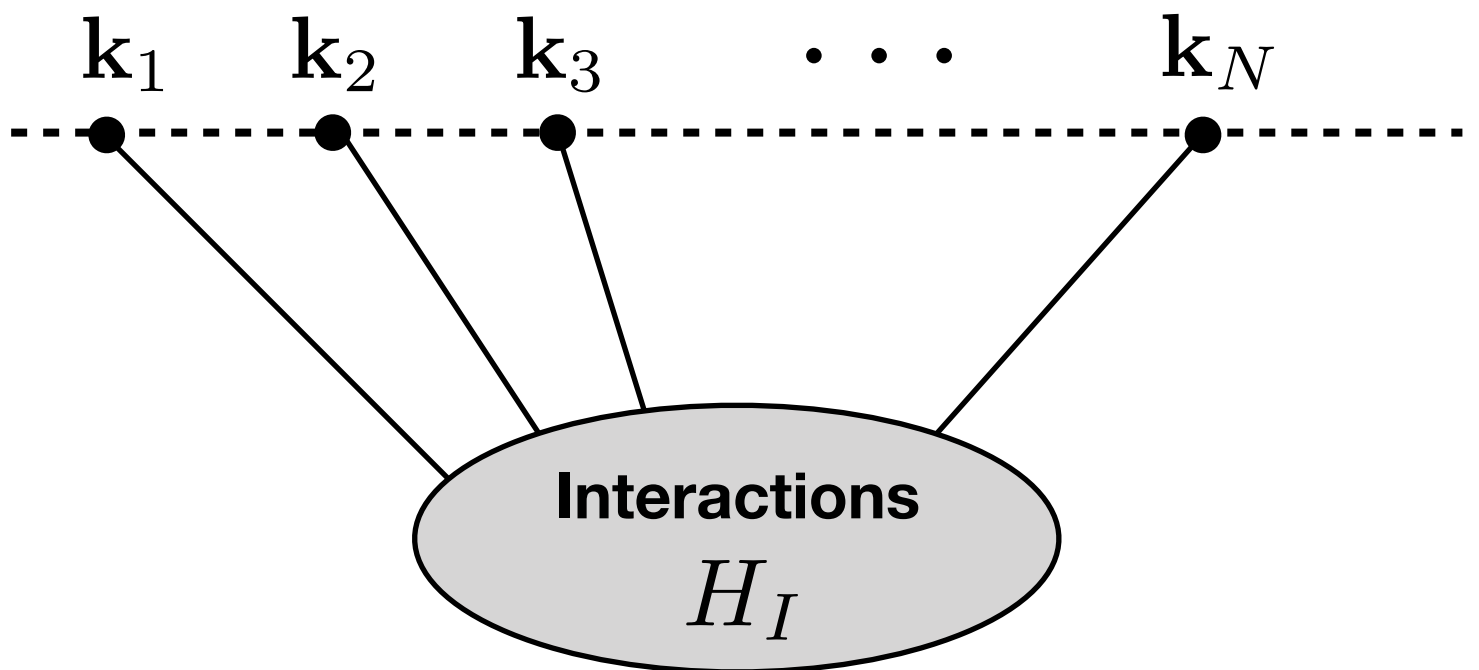
Inflation and primordial fluctuations 1

In inflation we are interested in predicting correlation functions:

$$ds^2 = -dt^2 + a^2(t)e^{2\zeta(t,\mathbf{x})}d\mathbf{x}^2 \qquad \zeta = -\frac{H}{\dot{\phi}_0}\delta\phi$$

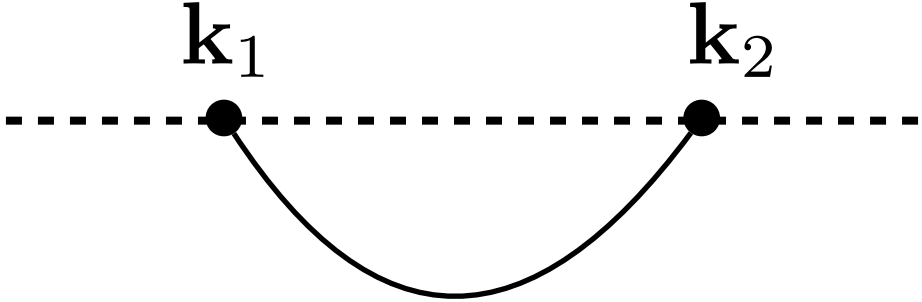
To compute them we use the in-in formalism:

$$H_I(t) = \mathcal{O}(\zeta_I^3) + \mathcal{O}(\zeta_I^4) + \dots$$

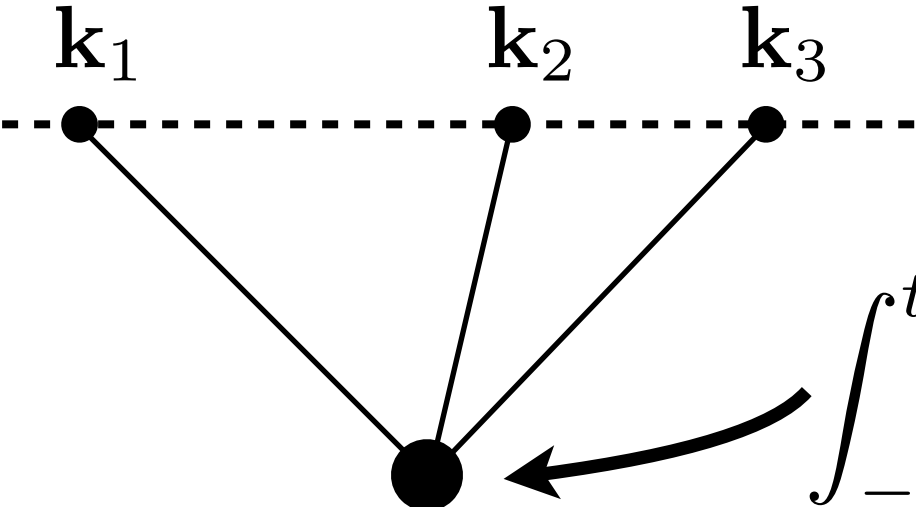
$$\langle \zeta_{\mathbf{k}_1}(t) \cdots \zeta_{\mathbf{k}_N}(t) \rangle =$$


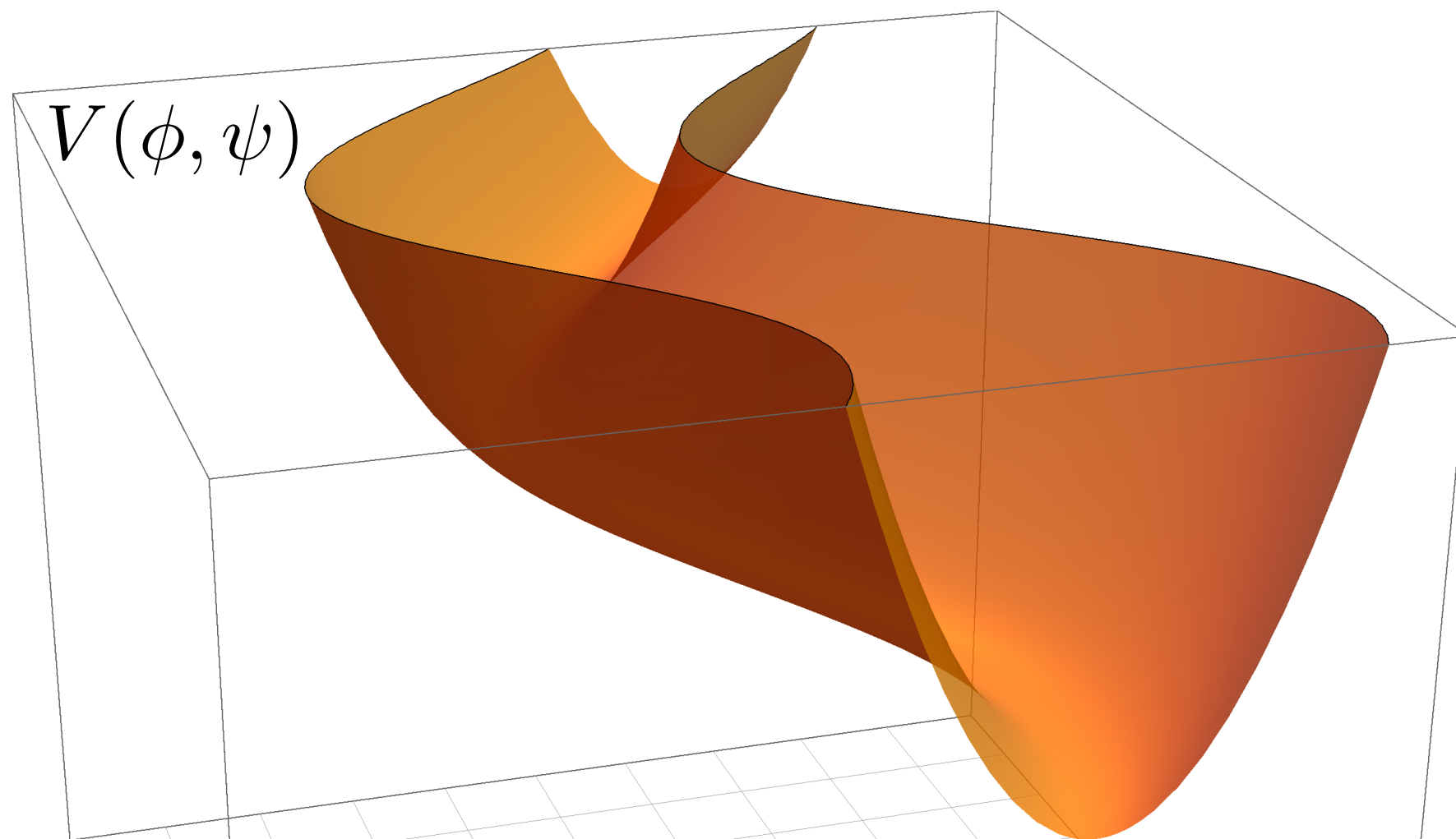
Inflation and primordial fluctuations 2

Power spectrum

$$\langle \zeta_{\mathbf{k}_1} \zeta_{\mathbf{k}_2} \rangle =$$

$$= (2\pi)^3 \delta(\mathbf{k}_1 + \mathbf{k}_2) \frac{2\pi^2}{k^3} \mathcal{P}_\zeta(k)$$

Bi-spectrum

$$\langle \zeta_{\mathbf{k}_1} \zeta_{\mathbf{k}_2} \zeta_{\mathbf{k}_3} \rangle =$$

$$\int_{-\infty}^t dt' H_I^{(3)}(t') = \mathcal{O}(\zeta_I^3)$$
$$= (2\pi)^3 \delta(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3) B(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3)$$



$$\mathcal{L} = \epsilon \left(\dot{\zeta} - \alpha \psi \right)^2 - \frac{\epsilon}{a^2} (\nabla \zeta)^2 + \frac{1}{2} \dot{\psi}^2 - \frac{1}{a^2} (\nabla \psi)^2 - \frac{1}{2} \mu^2 \psi^2 \\ + \epsilon \left(\dot{\zeta} - \alpha \psi \right)^3 + g \psi^3 + \dots$$

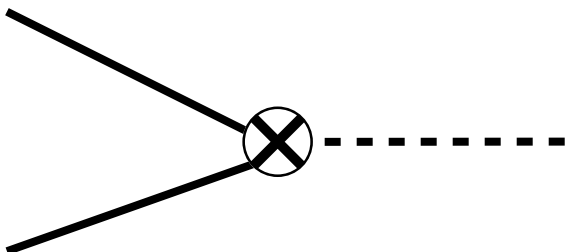
$$\mathcal{L} = \epsilon \left(\dot{\zeta} - \alpha \psi \right)^2 - \frac{\epsilon}{a^2} (\nabla \zeta)^2 + \frac{1}{2} \dot{\psi}^2 - \frac{1}{a^2} (\nabla \psi)^2 - \frac{1}{2} \mu^2 \psi^2$$

$$+ \epsilon \left(\dot{\zeta} - \alpha \psi \right)^3 + g \psi^3 + \dots$$

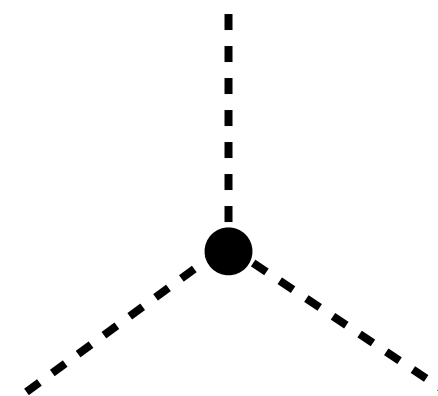
$\sim \alpha \dot{\zeta} \psi :$



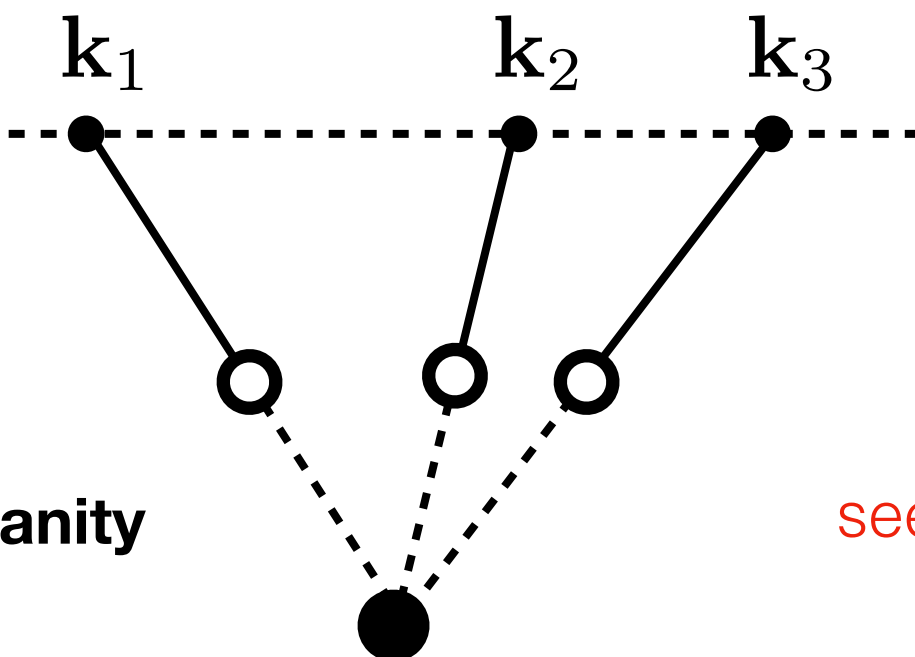
$\sim \alpha \dot{\zeta}^2 \psi :$



$g \psi^3 :$



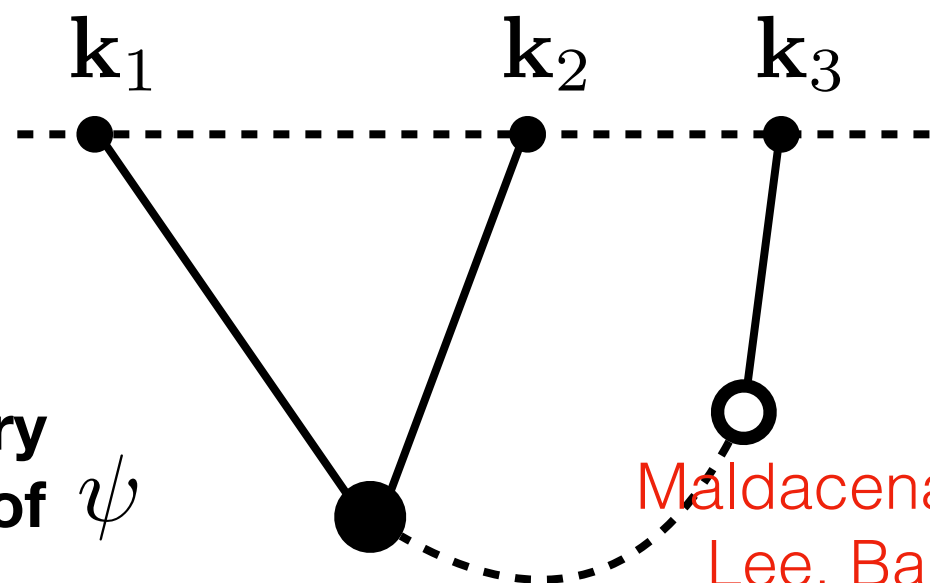
Quasi-single field:

$$\langle \zeta_{\mathbf{k}_1} \zeta_{\mathbf{k}_2} \zeta_{\mathbf{k}_3} \rangle =$$


Enhances local non Gaussianity

Chen & Wang (2012)
see also Assassi et al. (2013)

Cosmological colliders:

$$\langle \zeta_{\mathbf{k}_1} \zeta_{\mathbf{k}_2} \zeta_{\mathbf{k}_3} \rangle =$$


Bispectrum with oscillatory
patterns set by the mass of ψ

Maldacena & Arkani-Hamed (2016)
Lee, Baumann & Pimentel (2016)

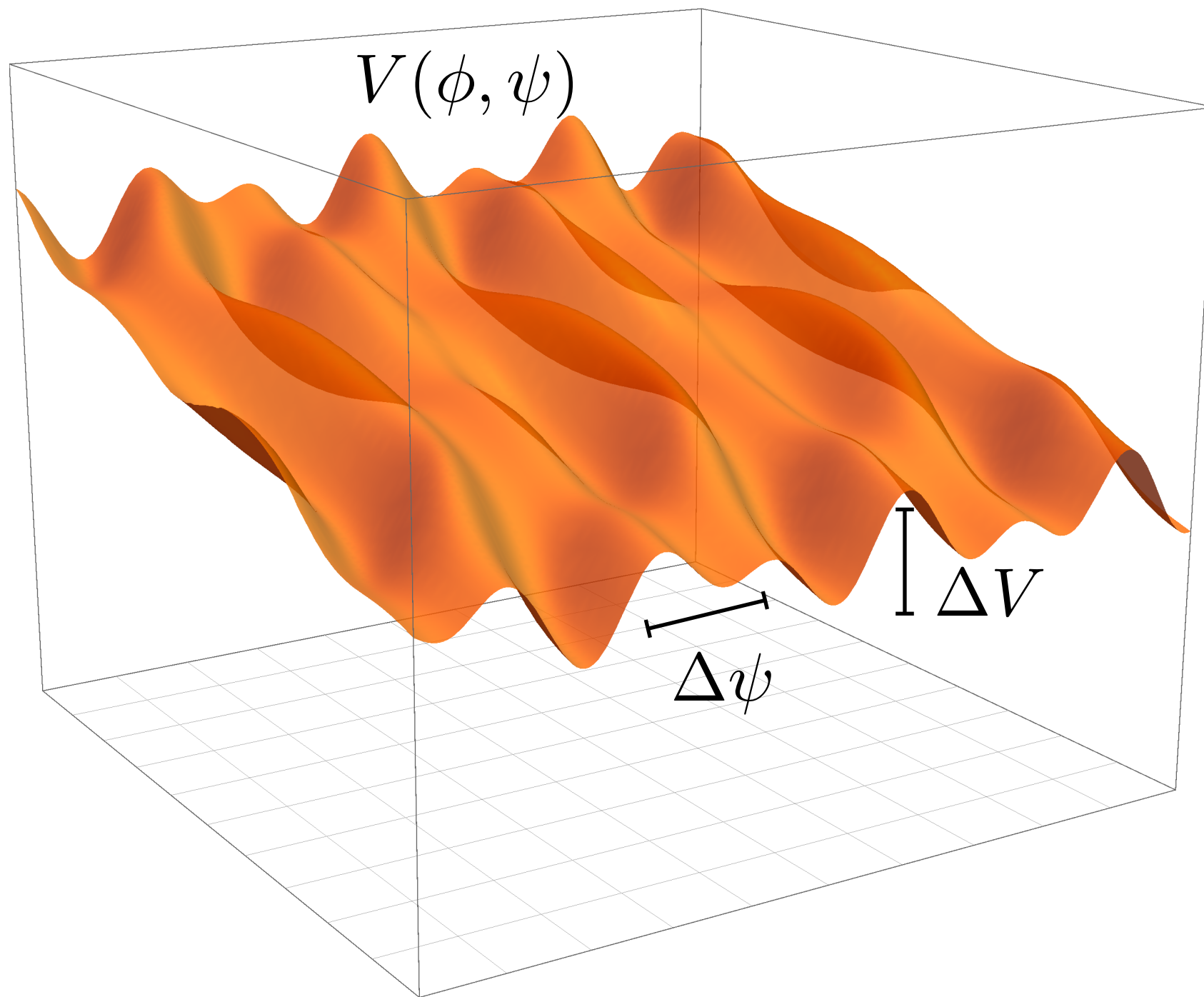
$$\mathcal{L} = \epsilon \left(\dot{\zeta} - \alpha \psi \right)^2 - \frac{\epsilon}{a^2} (\nabla \zeta)^2 + \frac{1}{2} \dot{\psi}^2 - \frac{1}{a^2} (\nabla \psi)^2 - \frac{1}{2} \mu^2 \psi^2 \\ + g \psi^3 + \dots$$

What about a more general situation:

$$\mathcal{L} = \epsilon \left(\dot{\zeta} - \alpha \psi \right)^2 - \frac{\epsilon}{a^2} (\nabla \zeta)^2 + \frac{1}{2} \dot{\psi}^2 - \frac{1}{a^2} (\nabla \psi)^2 - \frac{1}{2} \mu^2 \psi^2 \\ + \Delta V(\psi) + \dots$$

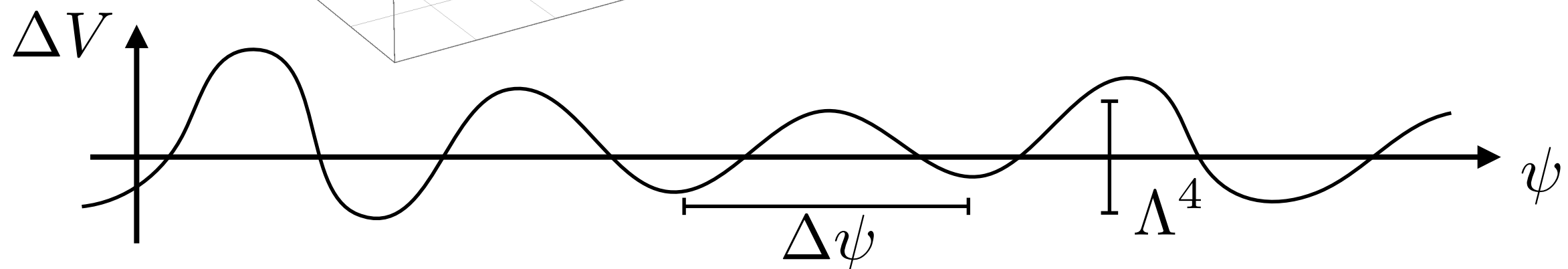
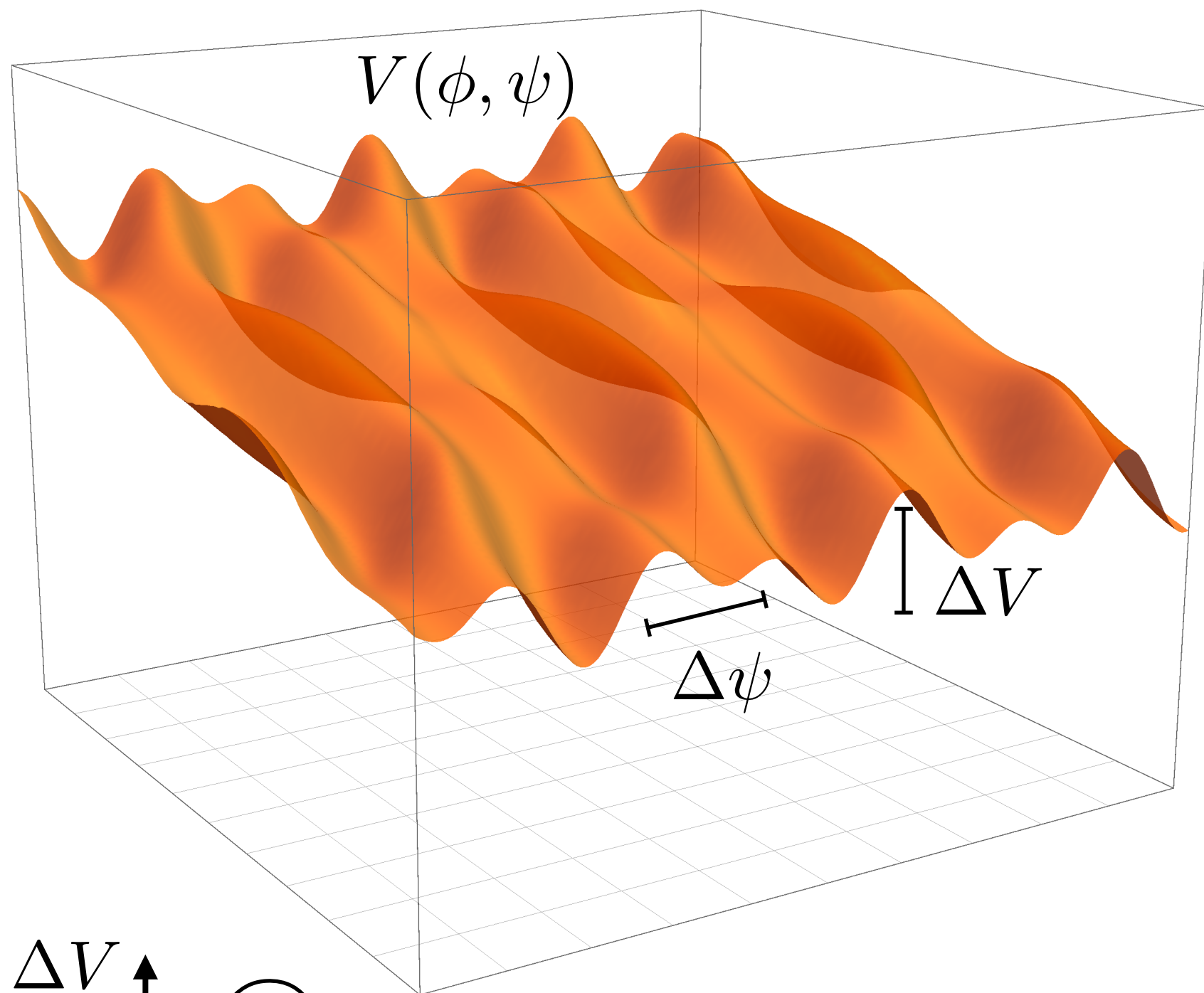
✿ Tomographic non-Gaussianity

7

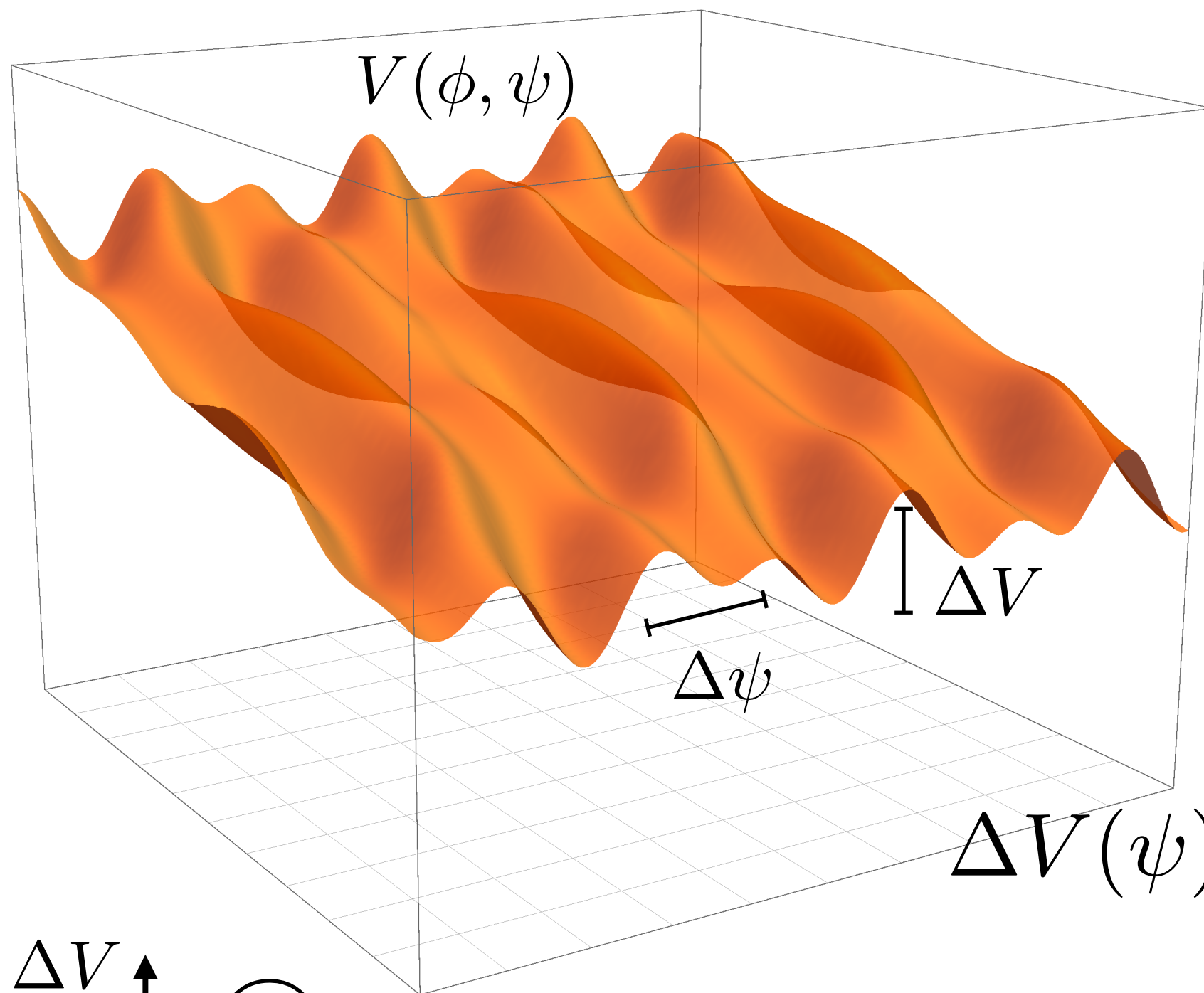


✿ Tomographic non-Gaussianity

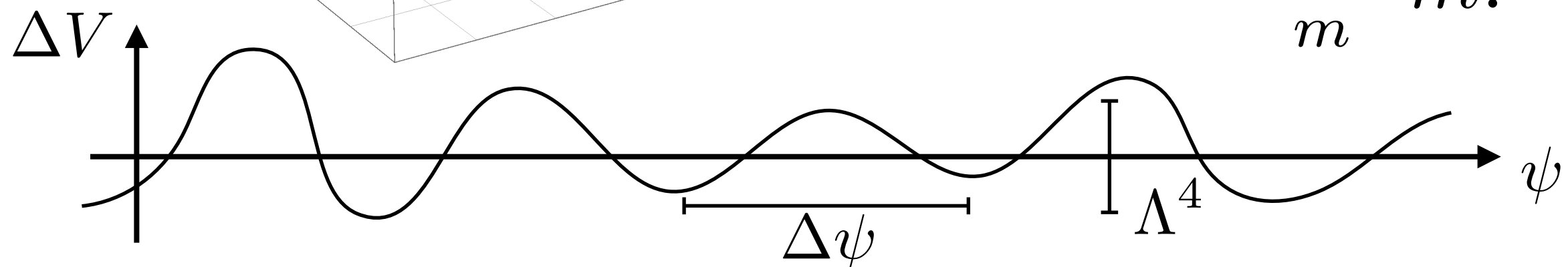
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Tomographic non-Gaussianity 7



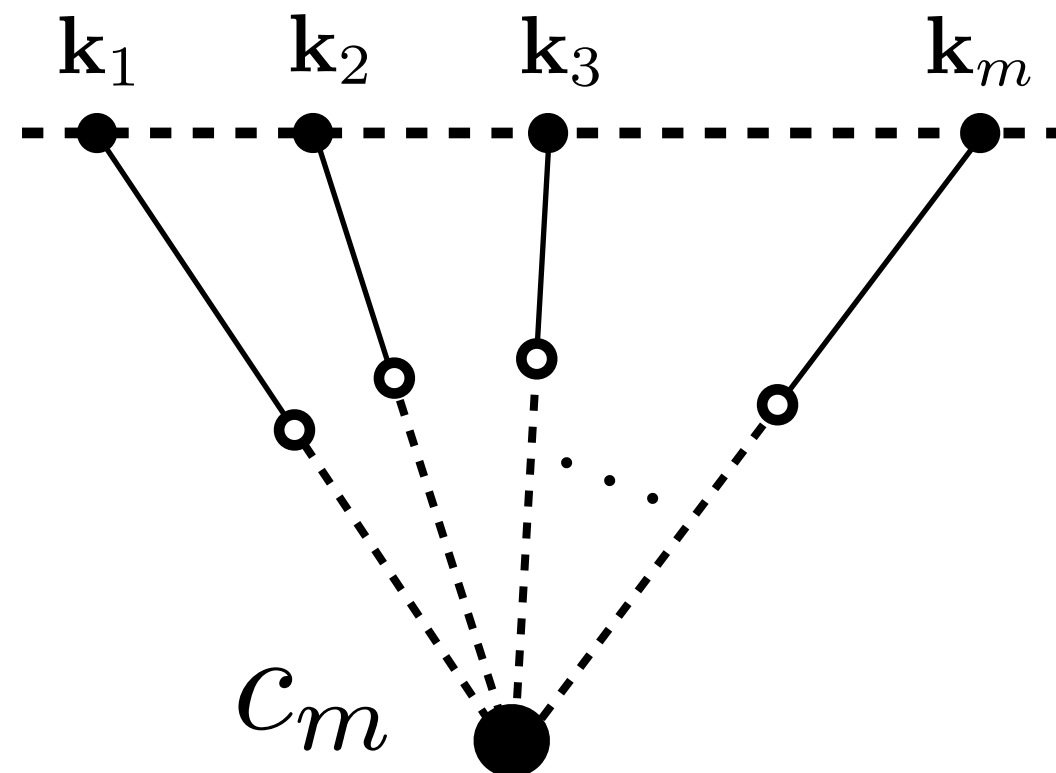
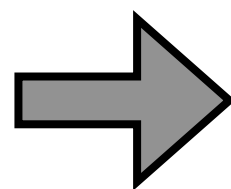
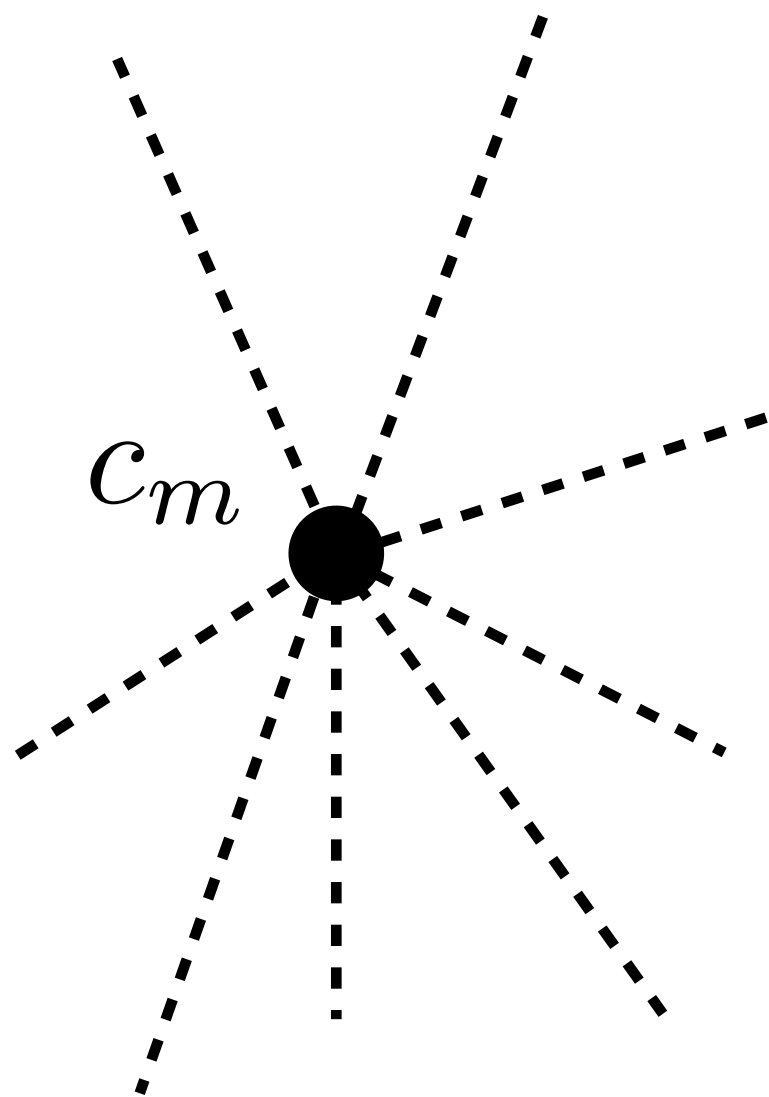
$$\Delta V(\psi) = \sum_m \frac{c_m}{m!} \psi^m$$



✿ n-point correlation functions

8

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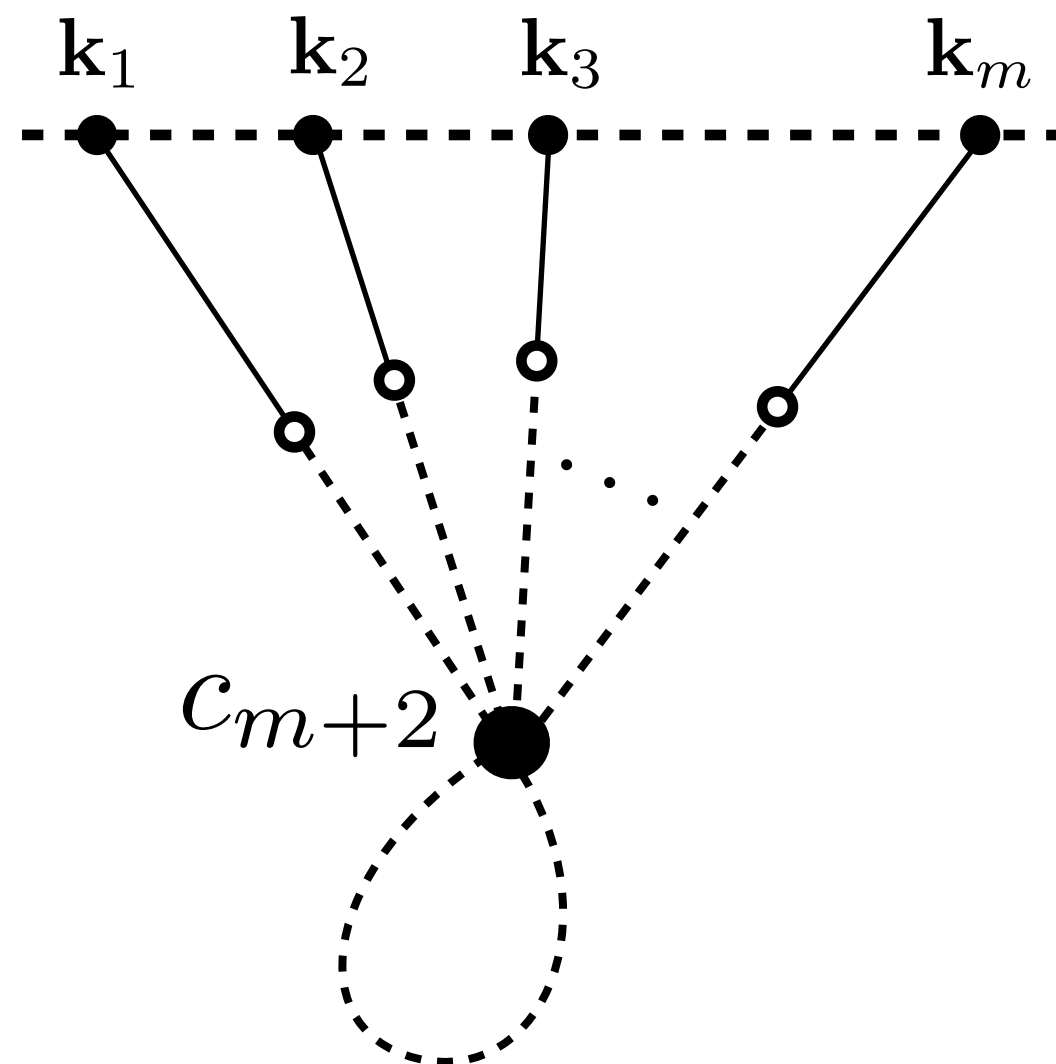
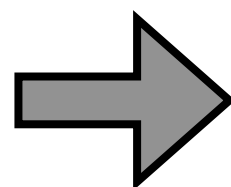
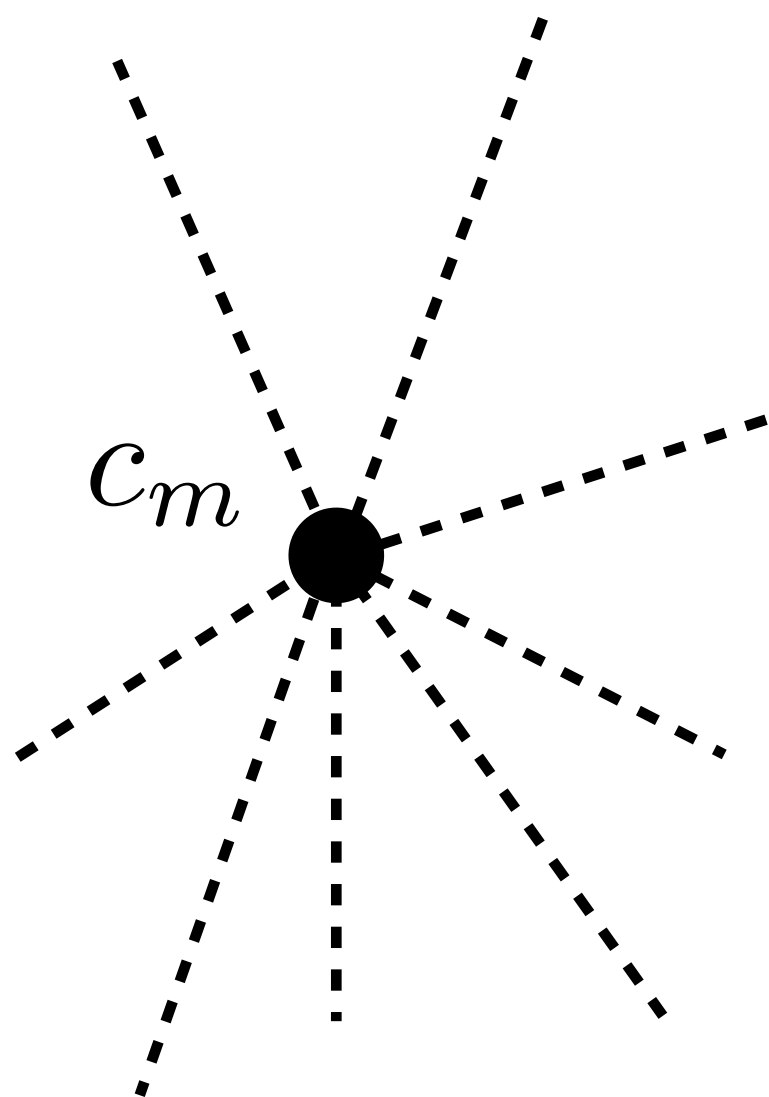


Recall Spyros Sypsas' talk

✿ n-point correlation functions

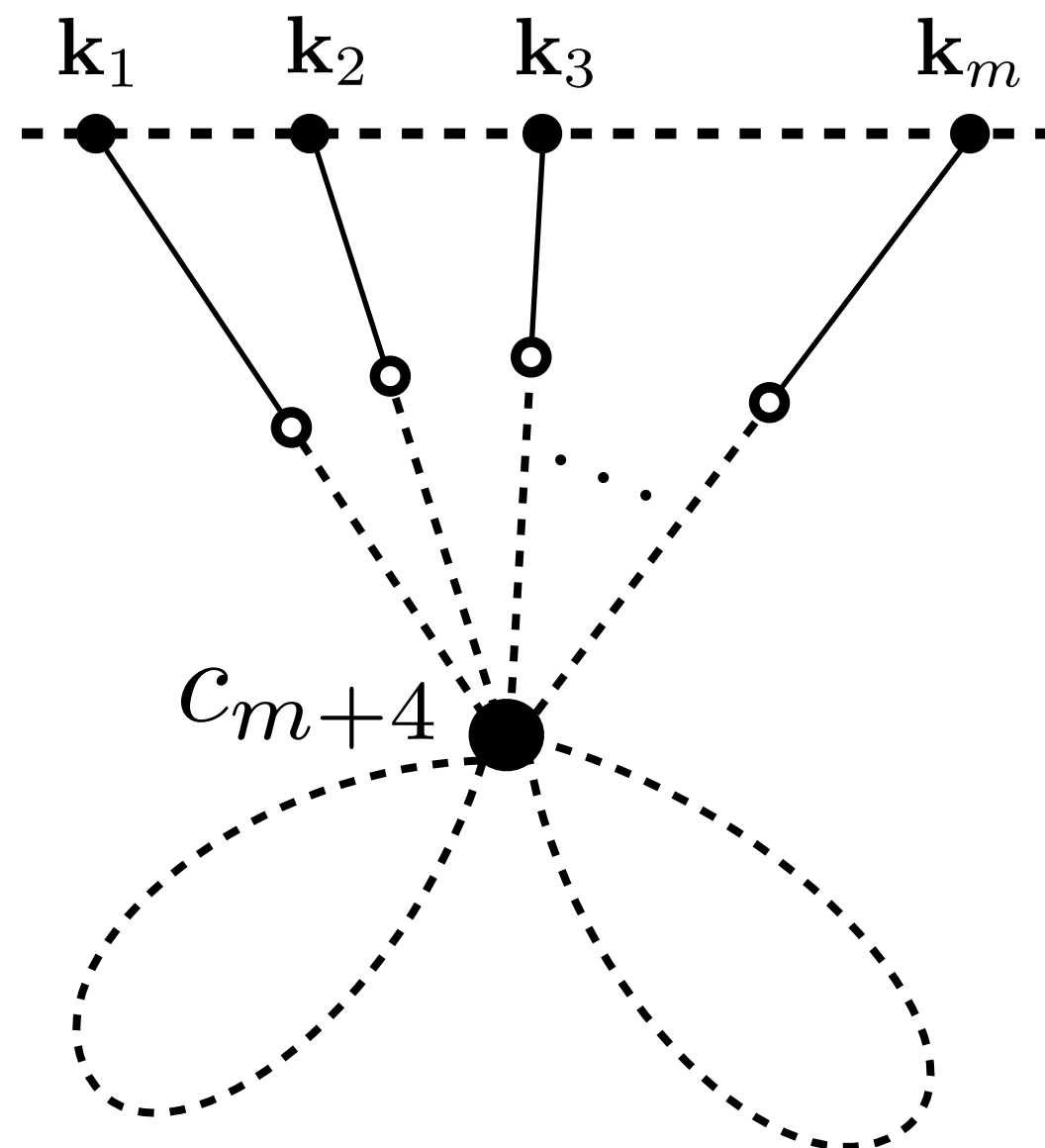
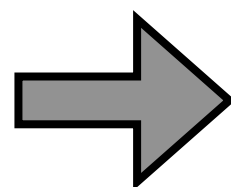
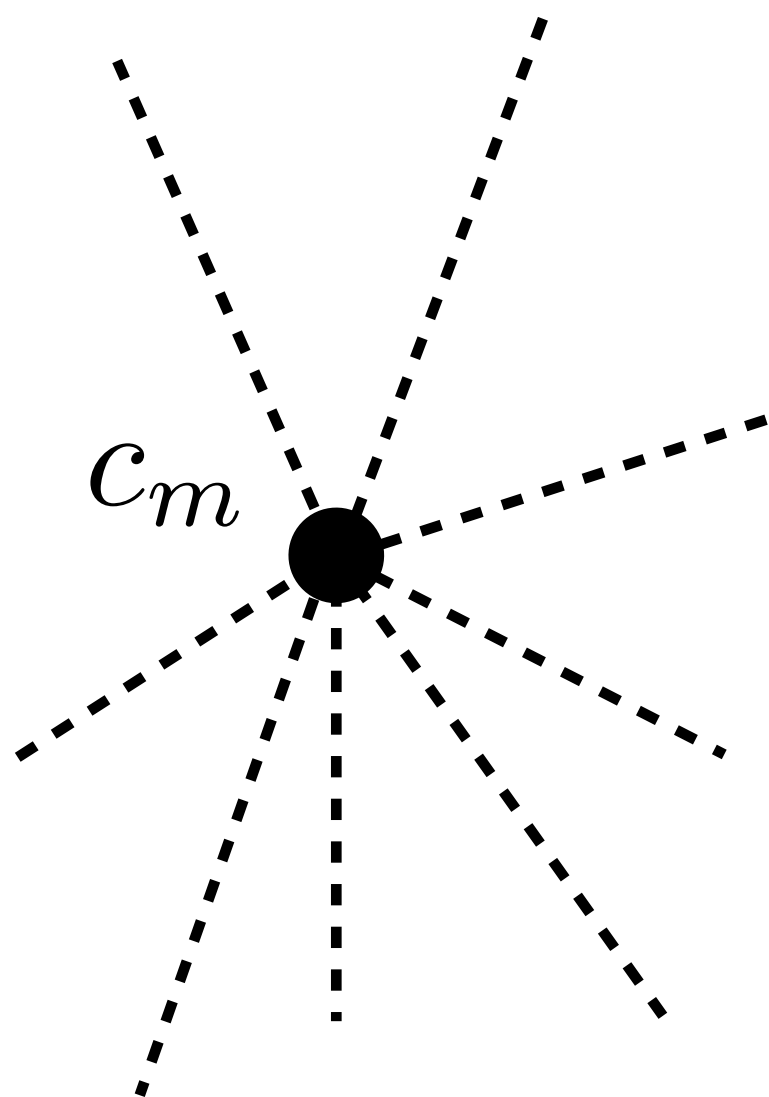
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$$\Delta V(\psi) = \sum_m \frac{c_m}{m!} \psi^m$$



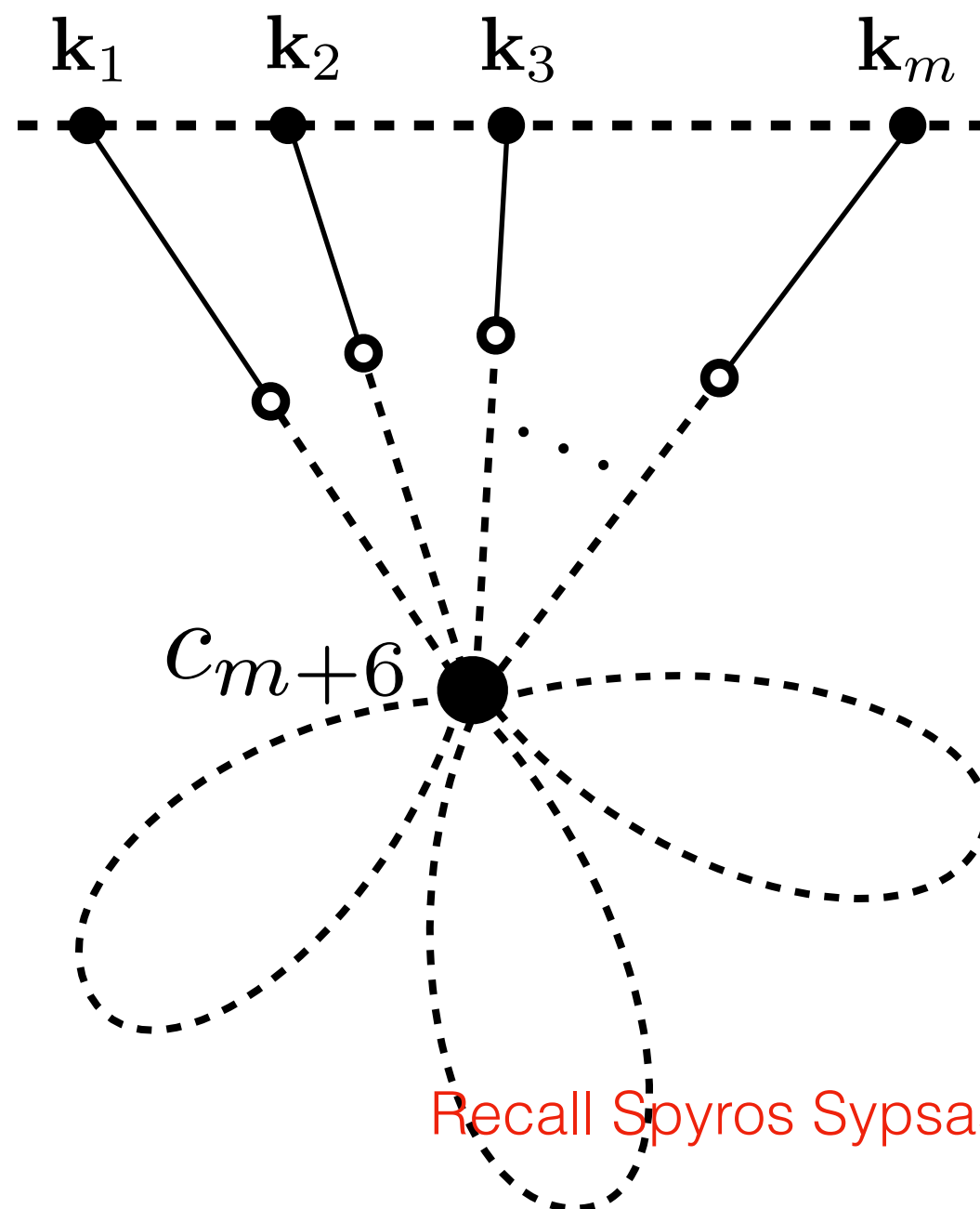
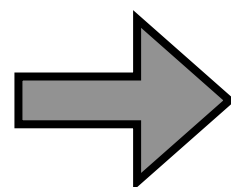
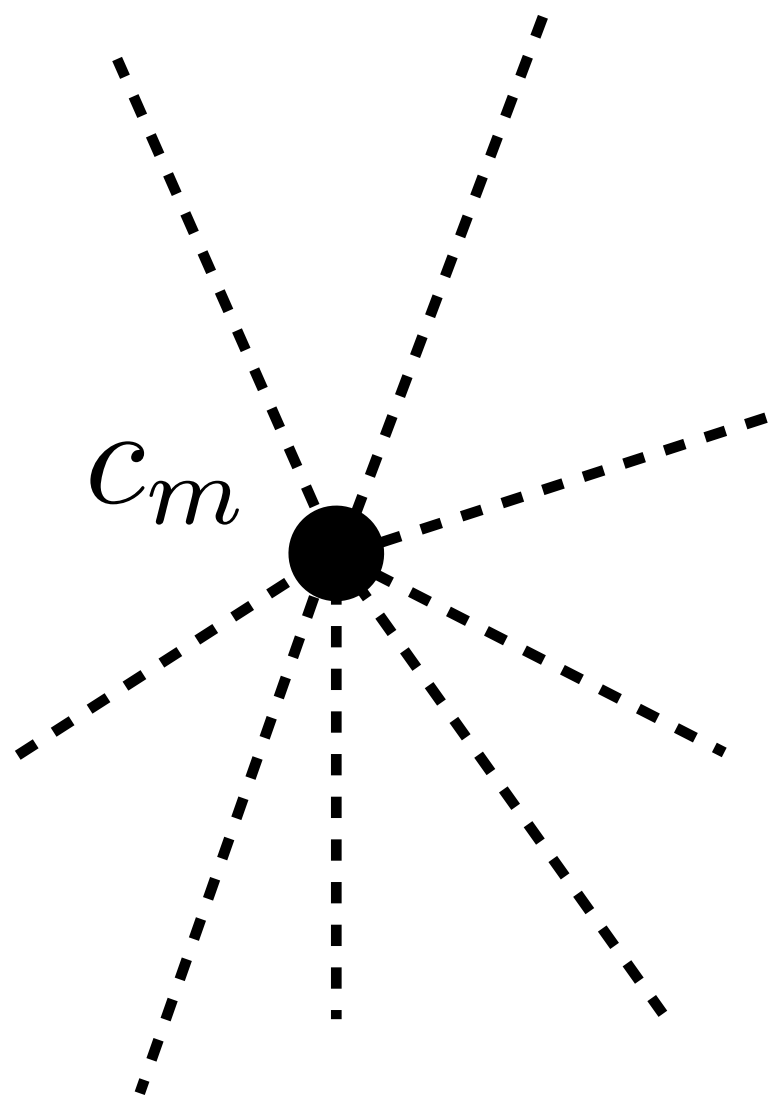
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$$\Delta V(\psi) = \sum_m \frac{c_m}{m!} \psi^m$$



Recall Spyros Sypsas' talk

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Recall Spyros Sypsas' talk

n-point correlation functions

One finds:

$$\langle \zeta_{\mathbf{k}_1}(t) \cdots \zeta_{\mathbf{k}_N}(t) \rangle = (2\pi)^3 h_n \delta^3\left(\sum \mathbf{k}\right) \frac{k_1^3 + \cdots + k_N^3}{k_1^3 \cdots k_N^3}$$

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$$h_n = \frac{1}{n} \left(\frac{\alpha H \Delta N}{2\sigma_L} \right)^n \frac{\Delta N}{3H^4} \int d\psi \frac{e^{-\frac{\psi^2}{2\sigma_L^2}}}{\sqrt{2\pi}\sigma_L} \text{He}_n(\psi/\sigma_L) \\ \times \left(\sigma_L^2 \frac{\partial^2}{\partial \psi^2} - \psi \frac{\partial}{\partial \psi} \right) \Delta V_{\text{ren}}(\psi),$$

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The reconstructed PDF is found to be

$$\rho(\zeta) = \frac{1}{\sqrt{2\pi}\sigma_\zeta} e^{-\frac{\zeta^2}{2\sigma_\zeta^2}} [1 + \Delta(\zeta)]$$

$$\begin{aligned} \Delta(\zeta) \equiv & \int_0^\infty \frac{dx}{x} \mathcal{K}(x) \int_{-\infty}^\infty d\bar{\zeta} \frac{\exp\left[-\frac{(\bar{\zeta} - \zeta(x))^2}{2\sigma_\zeta^2(x)}\right]}{\sqrt{2\pi}\sigma_\zeta(x)} \\ & \times \frac{\Delta N}{3H^4} \left(\sigma_\zeta^2 \frac{\partial^2}{\partial \bar{\zeta}^2} - \bar{\zeta} \frac{\partial}{\partial \bar{\zeta}} \right) \Delta V_{\text{ren}}(\psi_{\bar{\zeta}}) \end{aligned}$$

It looks complicated, but in fact it has some simple characteristics

The reconstructed PDF is found to be

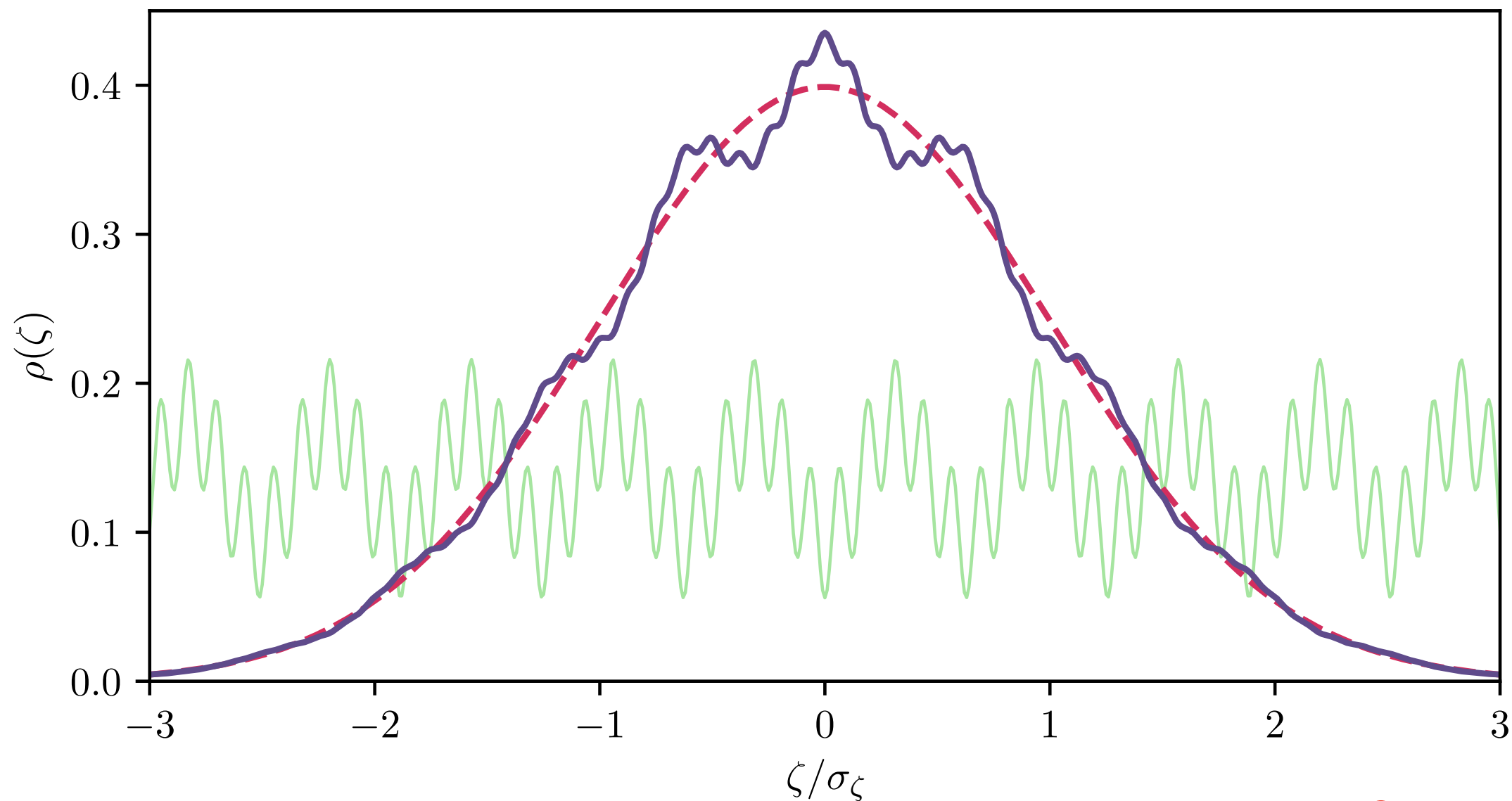
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It looks complicated, but in fact it has some simple characteristics

Example:

$$\Delta V(\psi) \propto [1 - \cos(\psi/f_1)] + [1 - \cos(\psi/f_2)]$$



Recall Spyros Sypsas' talk

We can compute: N-point functions at the end of inflation

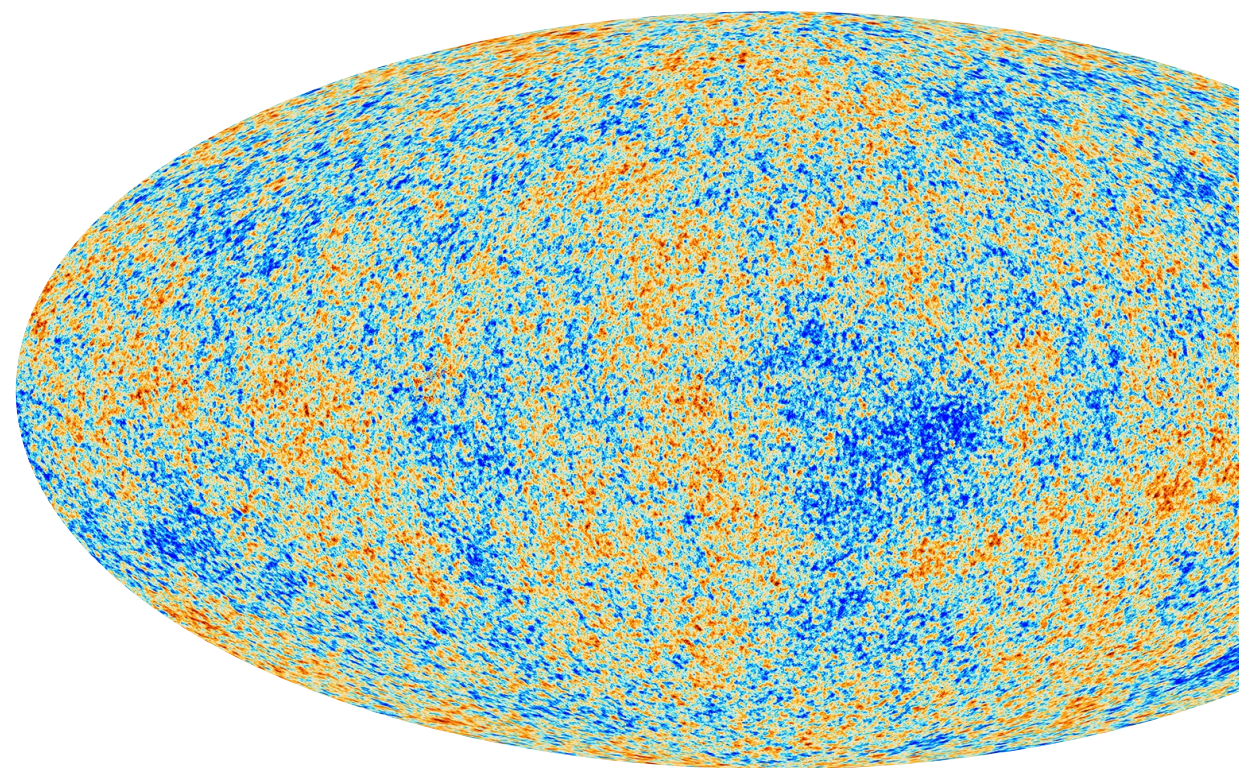
$$\langle \zeta(\mathbf{x}_1, t) \zeta(\mathbf{x}_2, t) \cdots \zeta(\mathbf{x}_N, t) \rangle_{t=t_{\text{end}}}$$

On Earth we observe: One realization of temperature fluctuations

$$\Theta(\hat{n}) \equiv \frac{\Delta T(\hat{n})}{T_0} = \frac{T(\hat{n}) - T_0}{T_0}$$

$$\Theta_{\mathbf{k}}(\hat{n}) = \mathcal{T}(k, \mu) \zeta_{\mathbf{k}}(t_{\text{end}})$$

$$\mu \equiv \mathbf{k} \cdot \hat{n}$$



We may assess non-Gaussianity in a different way: Let us introduce a probability distribution function (PDF)

$$\rho(\zeta) \longrightarrow \rho(\Theta)$$

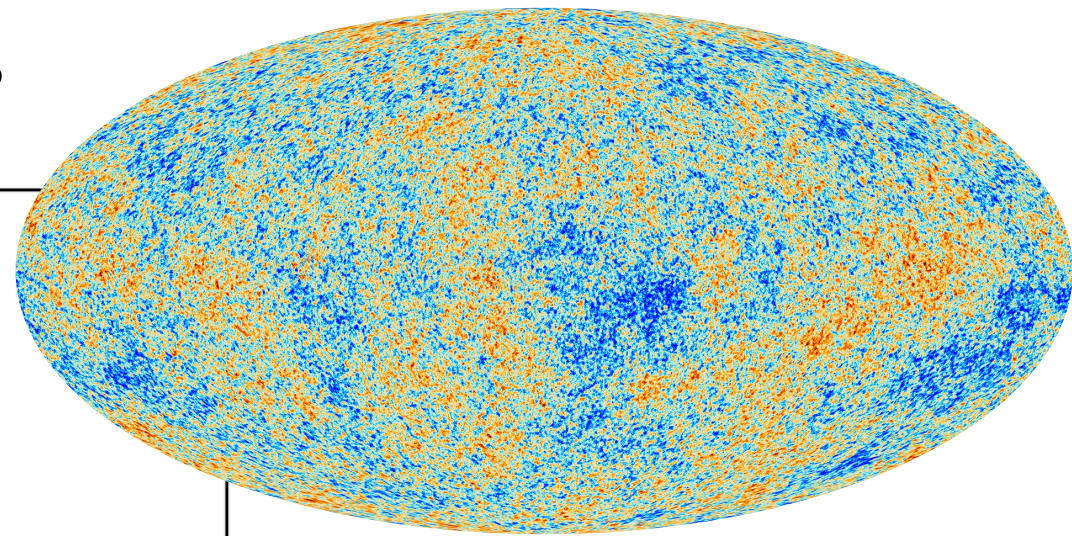
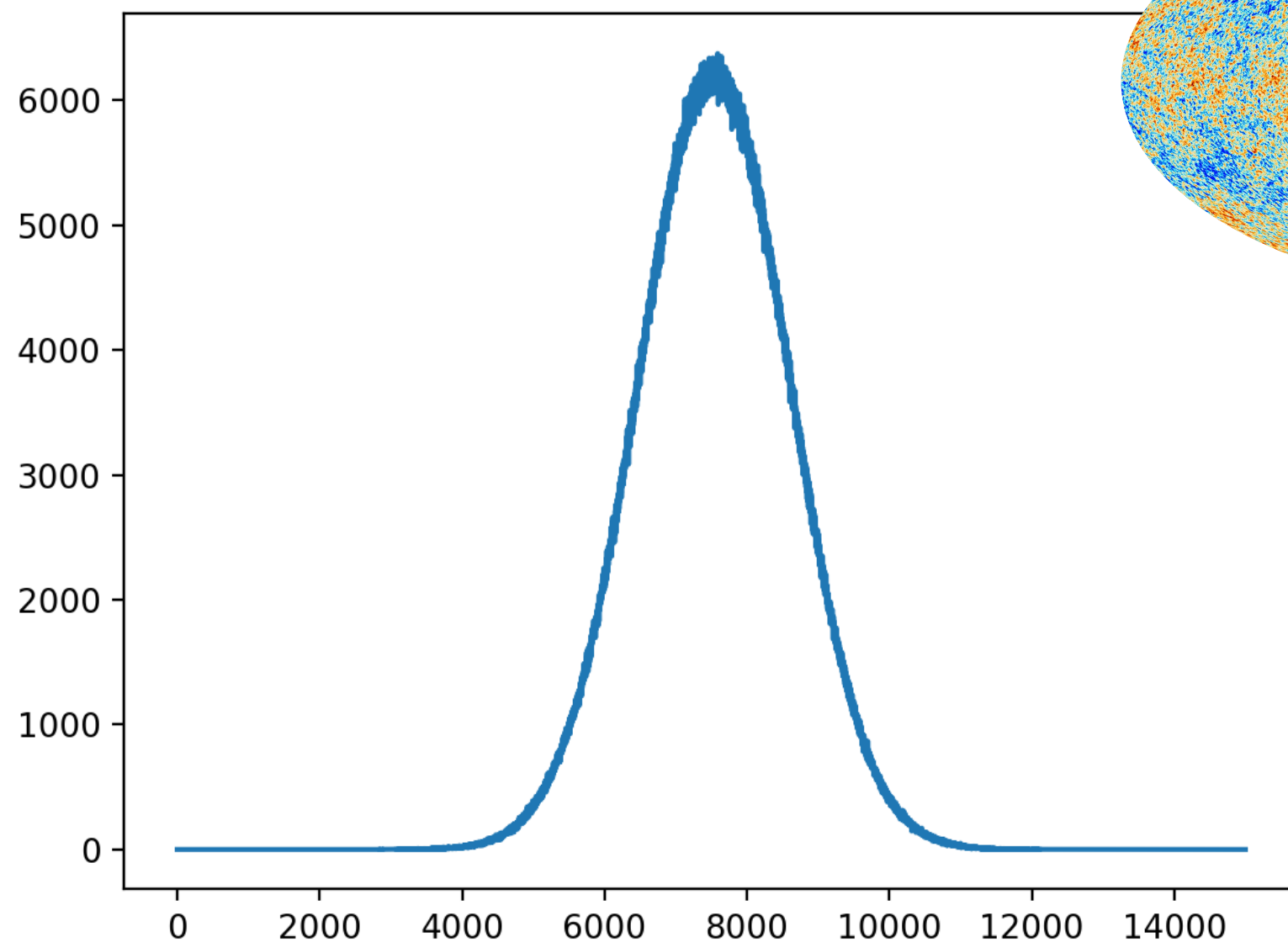
Such that $\langle \Theta^n \rangle = \int d\Theta \rho(\Theta) \Theta^n$

The variance of $\rho(\Theta)$ is found to be given by:

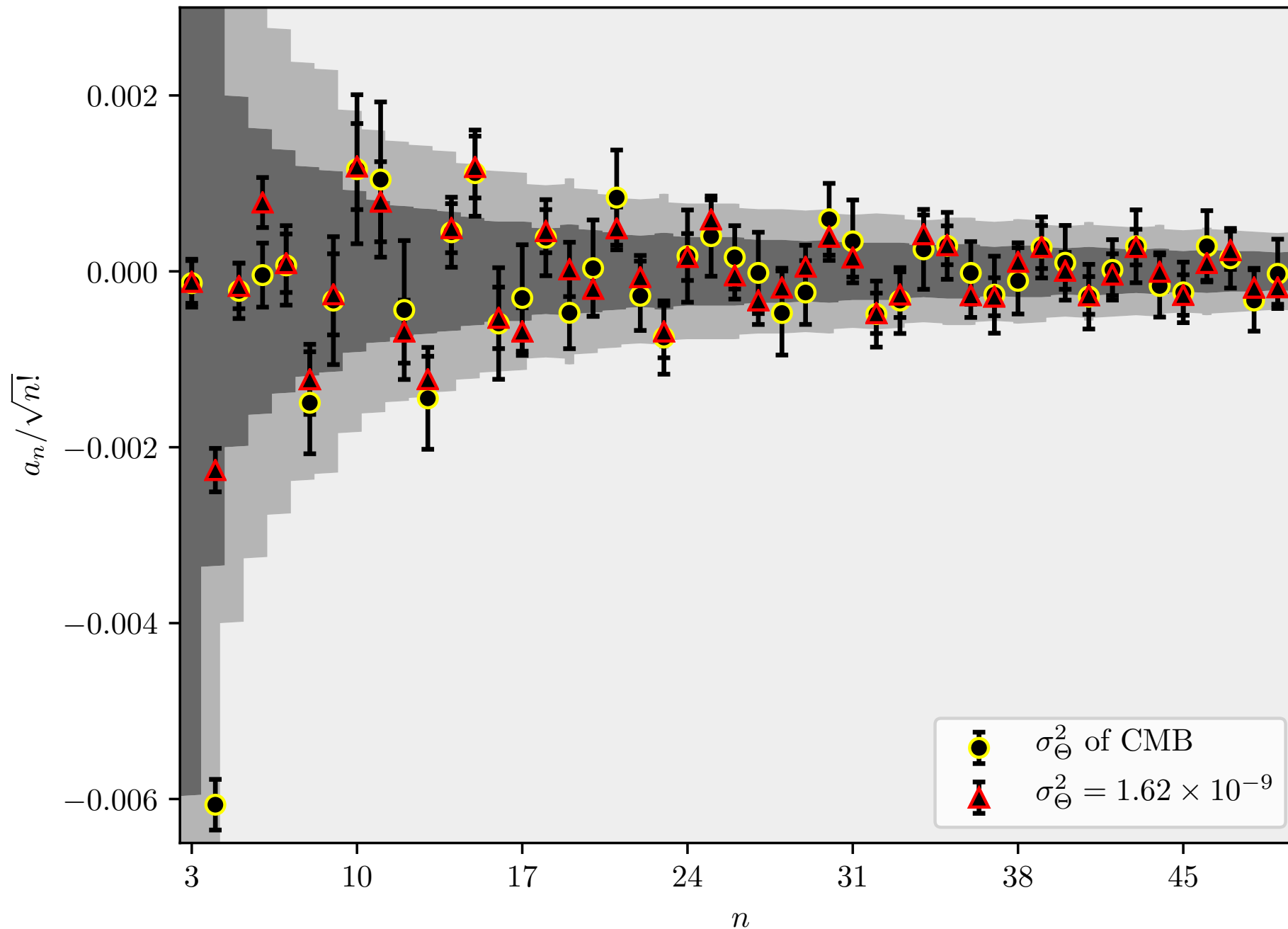
$$\sigma_{\Theta}^2 \equiv \langle \Theta(\hat{n}) \Theta(\hat{n}) \rangle = \int d\Theta \rho(\Theta) \Theta^2$$

$$\sigma_{\Theta}^2 = \frac{1}{4\pi} \sum_{\ell} (2\ell + 1) C_{\ell}$$

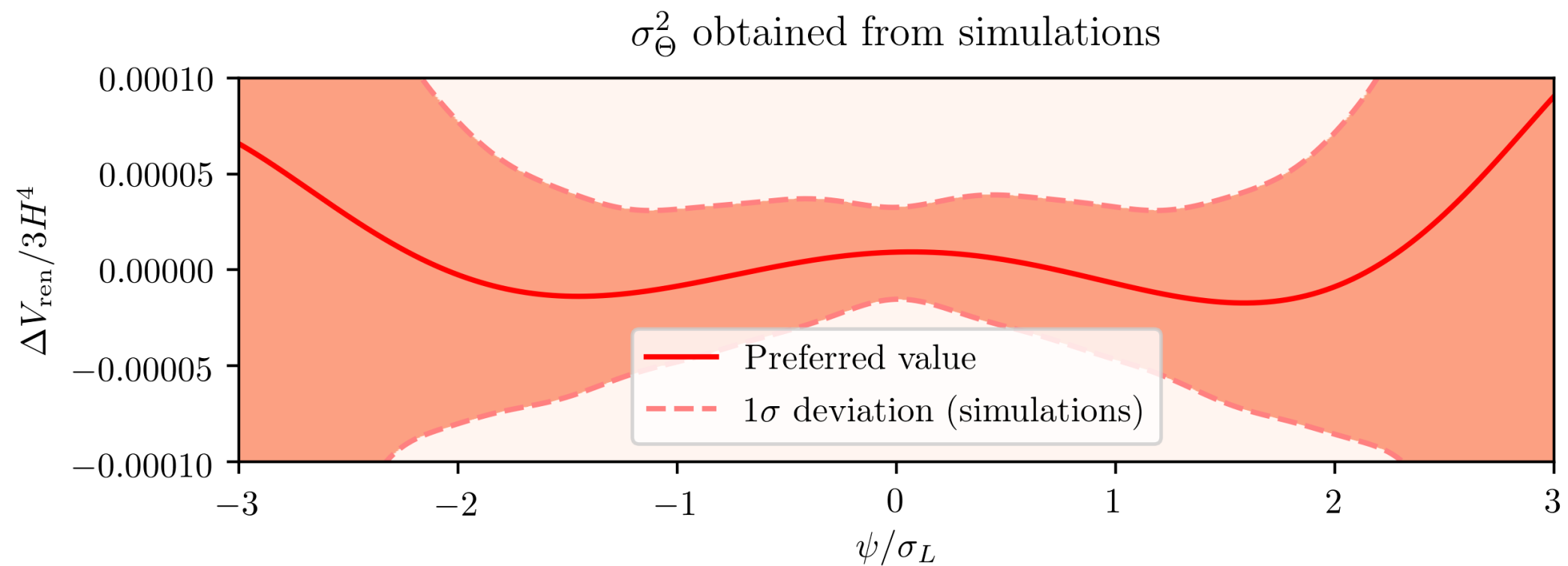
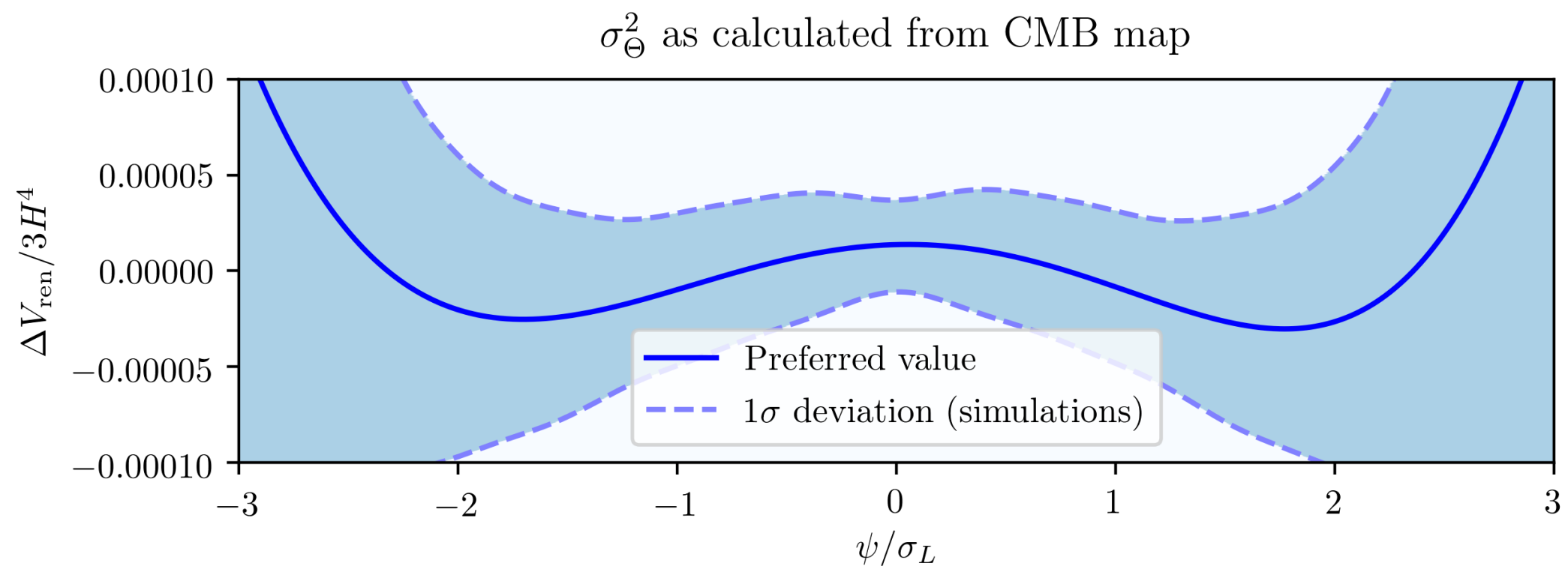
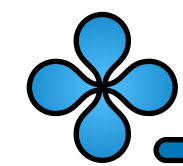
The reconstructed PDF from CMB data is



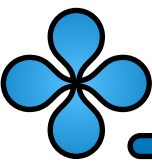
$$a_n \equiv \int d\Theta \, \rho(\Theta) H e_n(\Theta / \sigma_\Theta)$$



$$a_n \equiv \int d\Theta \, \rho(\Theta) H e_n(\Theta / \sigma_{\Theta}) \longrightarrow \Delta V(\psi)$$



- We have unveiled a novel class of NG that requires a description with many n-point functions
- Possibly, an effect only achievable with multi-field inflation
- Although not significant, current Planck data contains some level of tomographic NG
- Future CMB and LSS surveys will shed some light



Conclusions

Thanks!