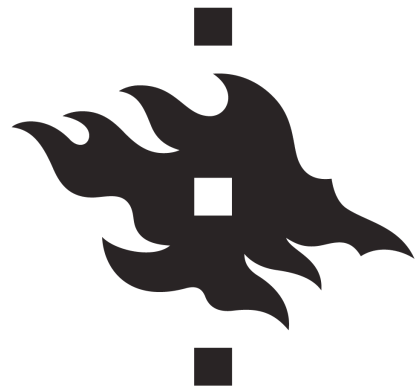




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Cosmic Microwave Background Constraints for Global Defects

JCAP 1707 (2017) 026
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COSMO-18
2018/08/28

In collaboration with: M. Hindmars, J. Lizarraga, J. Urrestilla

OBJECTIVE:

Obtain first CMB power spectra predictions and constraints for Global Defects

1. Perform *Field Theory Simulations* for global defects
2. Compute the *UnEqual Time Correlators*
3. Model Radiation-Matter *Transition*
4. Obtain *Power Spectra* and *Constraints*

Cosmic Defects

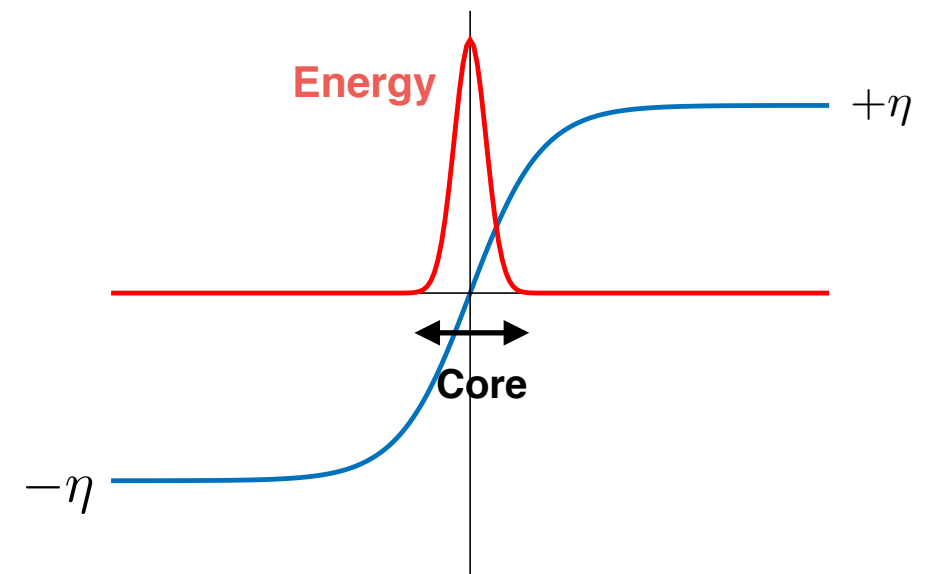
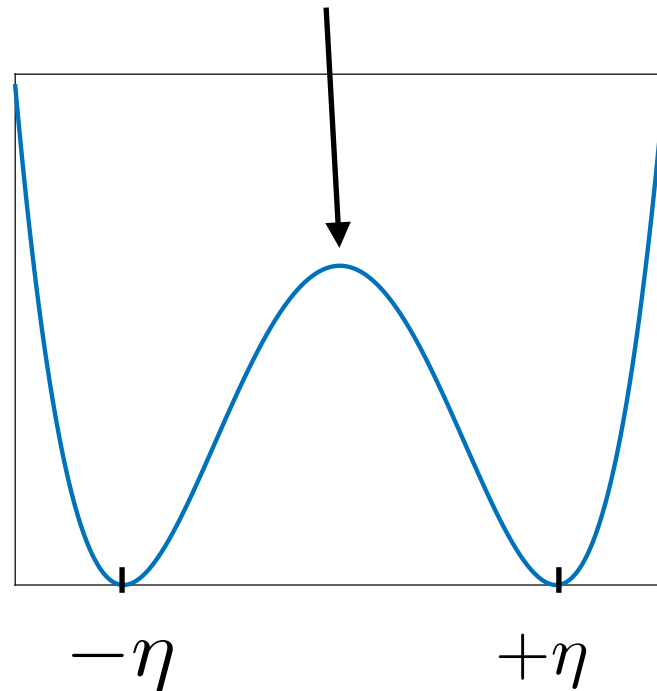
+ Kink (1+1):

$$\mathcal{L} = (\partial_t \phi)^2 - (\partial_x \phi)^2 - V(\phi) \quad \phi(t, x) \in \mathbb{R}$$

$$V(\phi) = \frac{\lambda}{2}(\phi^2 - \eta^2)^2 \xrightarrow{\phi \rightarrow 0} \frac{\lambda\eta^4}{2}$$

$$\phi(t, -\infty) = -\eta$$

$$\phi(t, +\infty) = +\eta$$



Defect Types:

$\pi_0(\mathcal{M}) \neq 1 \rightarrow$ Domain Wall

$\pi_1(\mathcal{M}) \neq 1 \rightarrow$ Cosmic String

$\pi_2(\mathcal{M}) \neq 1 \rightarrow$ Monopole

$\pi_3(\mathcal{M}) \neq 1 \rightarrow$ Texture

Global Defect Model

+ 3+1

$$\mathcal{L} = \frac{1}{2} \partial_\mu \Phi^i \partial^\mu \Phi^i - \frac{1}{4} \lambda (|\Phi|^2 - \eta^2)^2 \quad \begin{array}{l} \Phi^i \in \mathbb{R} \\ i = 1, \dots, N \end{array}$$

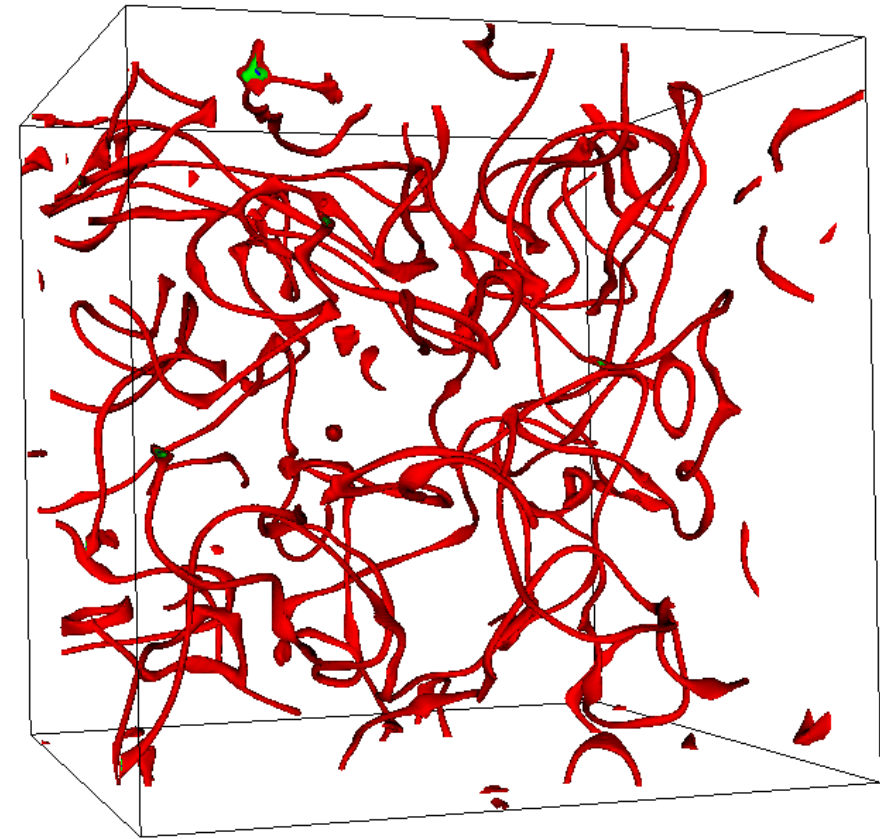
$$\phi(\mathbf{x}) \rightarrow e^{i\alpha} \phi(x) \mapsto \text{Global}$$

Our Simulations	$N = 2$	$\longrightarrow$	Strings
	$N = 3$	$\longrightarrow$	Monopoles
	$N = 4$	$\longrightarrow$	Textures
	$N = 8, 12, 20$		
	Large-N	$\longrightarrow$	Analytic Solution

- Presence of Goldstone bosons $\rightarrow$ energy not confined
- Long range interactions

Field Theory Simulations

- Field Theory Simulations
- Discrete e.o.m.
- Evolve e.o.m. in N^3 lattices
- Expanding background: Radiation, Matter and Transition
- Scaling is indispensable $\xi(\tau) \propto \tau$:
 - In Radiation and Matter eras ✓
 - In Rad-Mat transition ✗
- Extraction of UETCs



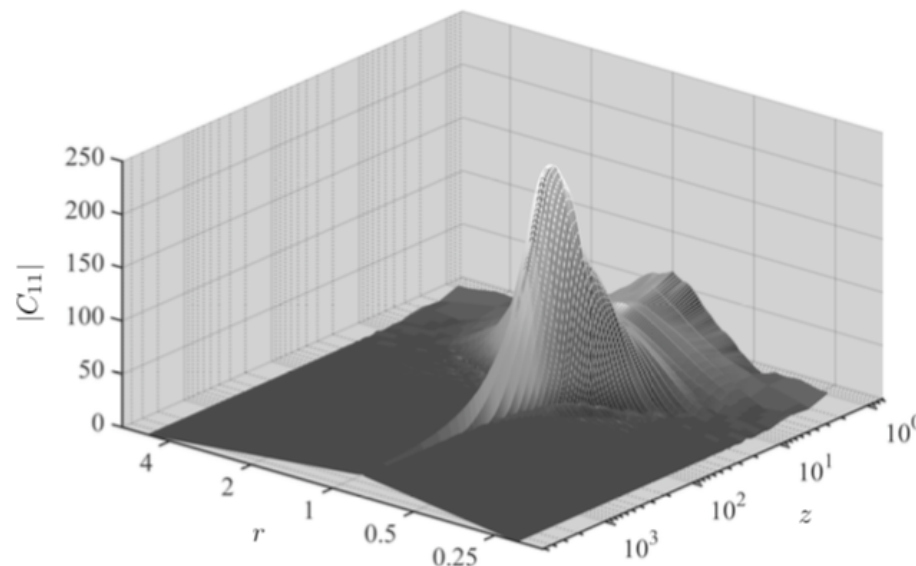
UETC method

- Cosmic defects perturb background
- Creating CMB Temperature and Polarization anisotropies

Power Spectra can be obtained using:

Unequal **T**ime **C**orrelators of the energy-momentum tensor

$$U_{\lambda\kappa\mu\nu}(\mathbf{k}, \tau, \tau') = \langle T_{\lambda\kappa}(\mathbf{k}, \tau) T_{\mu\nu}^*(\mathbf{k}, \tau') \rangle$$



UETC method

$$U_{\lambda\kappa\mu\nu}(\mathbf{k}, \tau, \tau') = \langle T_{\lambda\kappa}(\mathbf{k}, \tau) T_{\mu\nu}^*(\mathbf{k}, \tau') \rangle$$

Scaling + Symmetry

$$C(k\tau, k\tau')$$

Decompose into
coherent functions

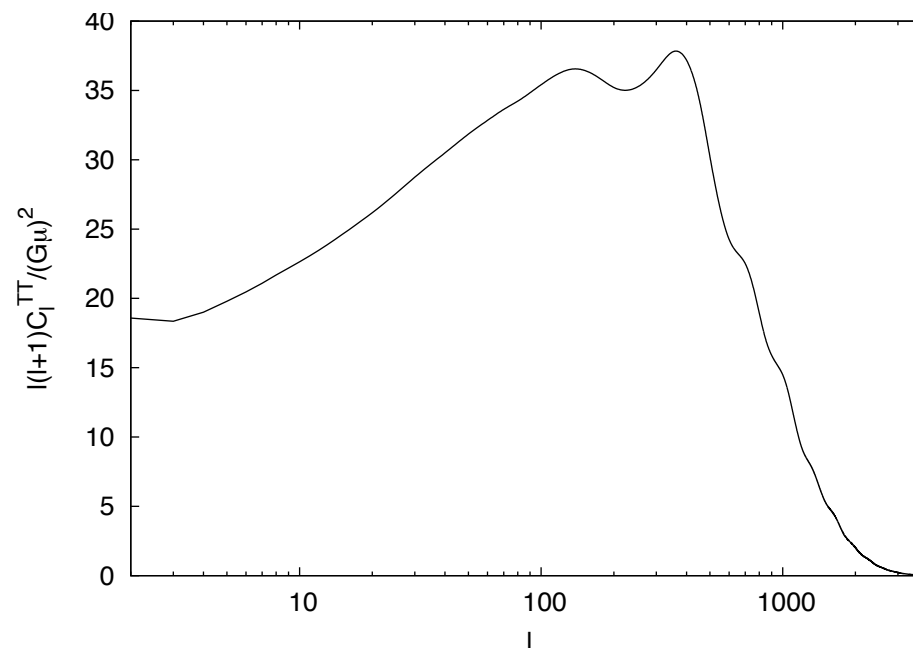
$$\begin{matrix} C_{\phi\phi} & C_{\psi\phi} \\ C_{\psi\psi} & C_{vv} & C_{tt} \end{matrix}$$

$$\sqrt{\lambda_n} c_n(k\tau)$$

Pen, Spergel and Turok
PRD49 692-729

CMBeasy, CAMB
CLASS

Einstein-Boltzmann
solver



Transition

- Important information in cosmological transition
- Scaling is broken
- We can model transition with fixed-k interpolation method *Daverio et al. 2015*

$$C(k, \tau, \tau') = f \left(\frac{\sqrt{\tau\tau'}}{\tau_{eq}} \right) C^R(k\tau, k\tau') + \left(1 - f \left(\frac{\sqrt{\tau\tau'}}{\tau_{eq}} \right) \right) C^M(k\tau, k\tau')$$

$$f_{ab}(k, \tau) = \frac{E_{ab}^{RM}(k, \tau) - \bar{E}_{ab}^M(k\tau)}{\bar{E}_{ab}^R(k\tau) - \bar{E}_{ab}^M(k\tau)}$$

Equal Time Correlator: $E(k, \tau) = C(k\tau, k\tau)$

Simulations at transition
to compute:

$$E^{RM}(k, \tau)$$

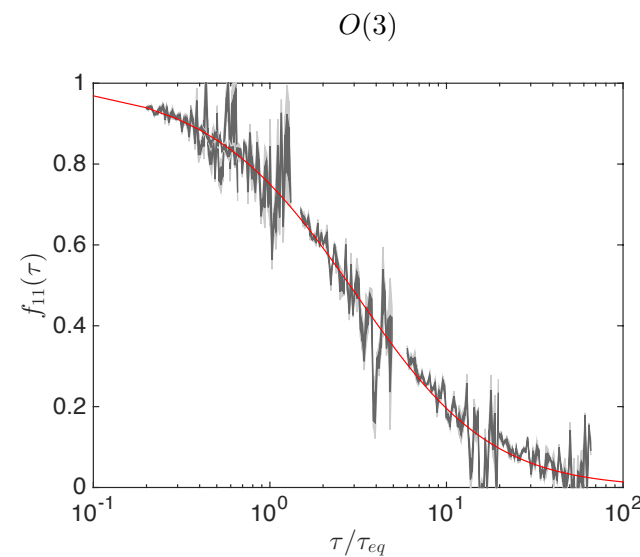
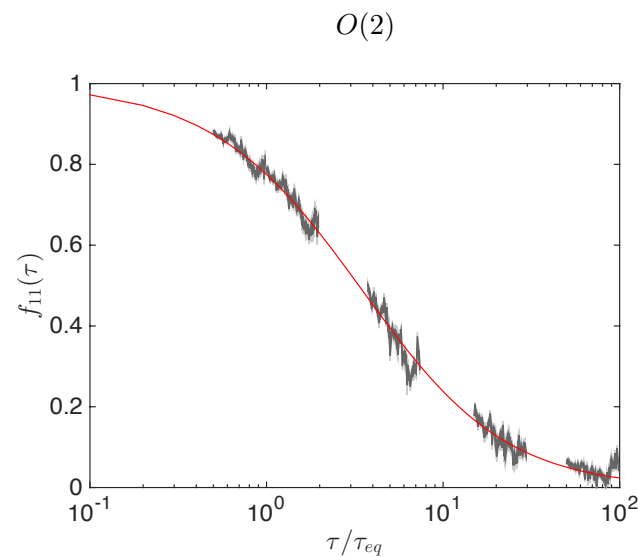
$$f(\tau) = \left(1 + \gamma \frac{\tau}{\tau_{eq}} \right)^\kappa$$

Source functions (for every k): $\sqrt{\lambda_n} c_n(\tau/\tau_{eq})$

Transition

O(2) and O(3)

$$f(\tau) = \left(1 + \gamma \frac{\tau}{\tau_{\text{eq}}}\right)^{\kappa}$$



	γ	κ
O(2)	0.26 ± 0.03	-1.15 ± 0.02
O(3)	0.23 ± 0.05	-1.4 ± 0.2
AH	0.25	-1

- Independent of k
- Slow transition
- Same for all correlators

N > 3

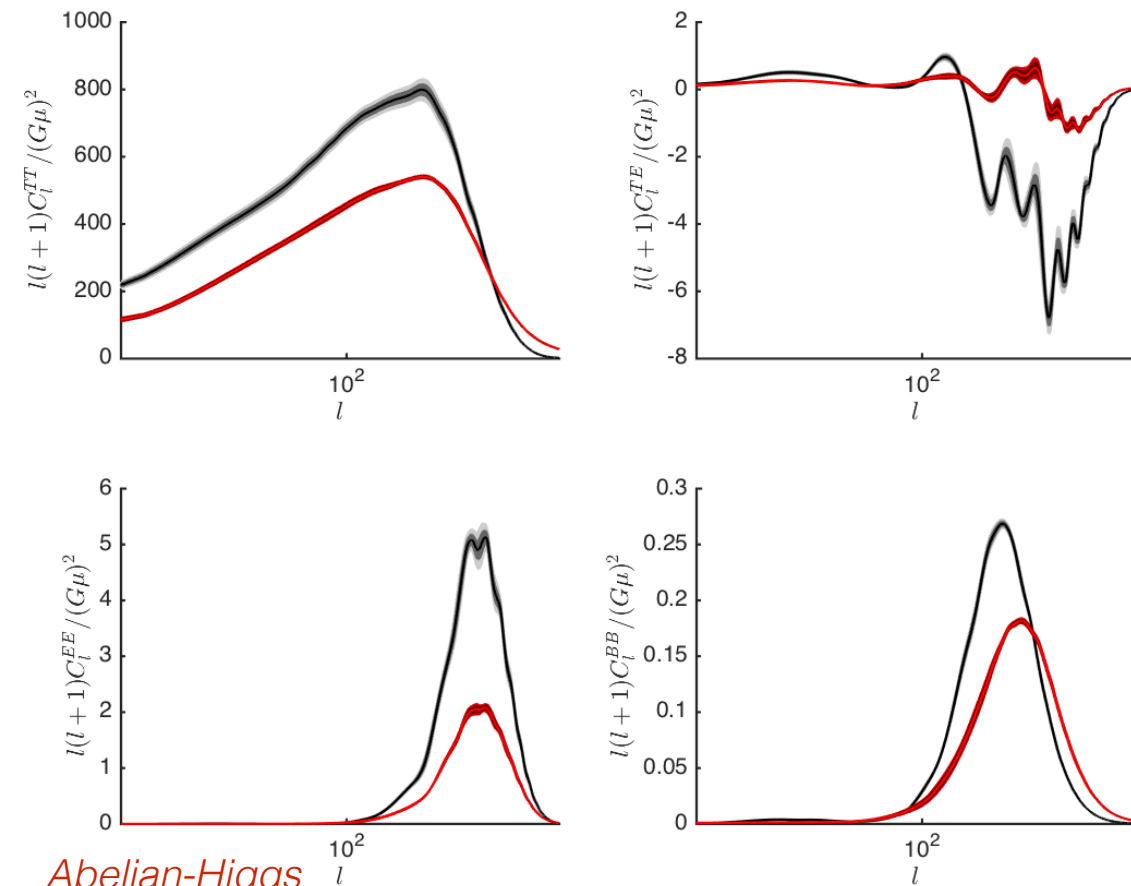
- Interpolation function depends on k
- Currently working on new approach to describe transition

(Work in progress)

Power spectra

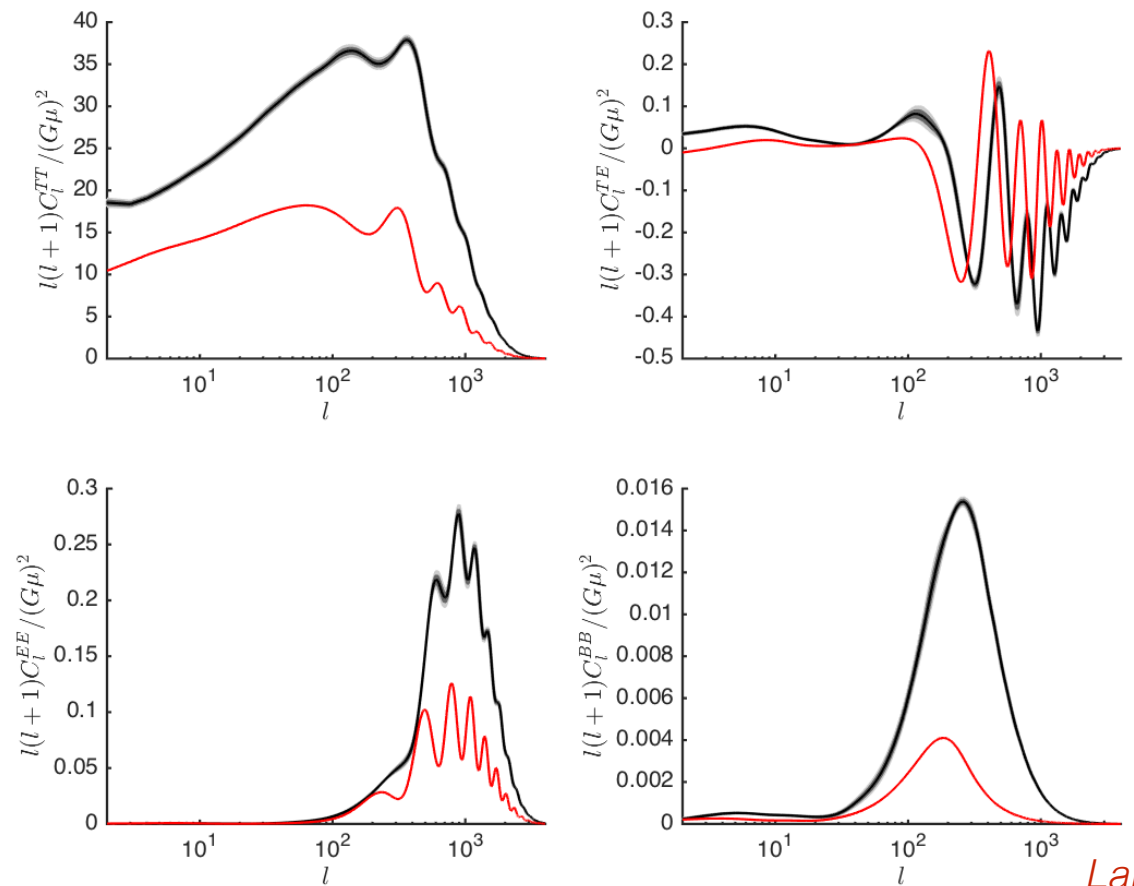
O(2)

O(3)



Abelian-Higgs

Lizarraga et al. 2016



Large-N

Fenu et al. 2013

- O(2) and O(3) relatively similar although O(2) has bigger amplitude
- O(2) and AH relatively similar although O(2) has bigger amplitude
- O(2) and Large-N are not similar
- O(3) and Large-N are more similar:
 - The amplitudes are different
 - Details are different

Bounds

Dataset	<i>Planck</i> 2015 CMB		
Defect	O(2)	O(3)	
Model	$\mathcal{P}\mathcal{L} + G\mu$	$\mathcal{P}\mathcal{L} + G\mu$	$\mathcal{P}\mathcal{L}$
f_{10}	< 0.017	< 0.024	—
$10^{12}(G\mu)^2$	< 0.031	< 0.73	—
$-\ln \mathcal{L}_{\max}$	6472	6470	6472

95% upper limits obtained using Monte Carlo analysis

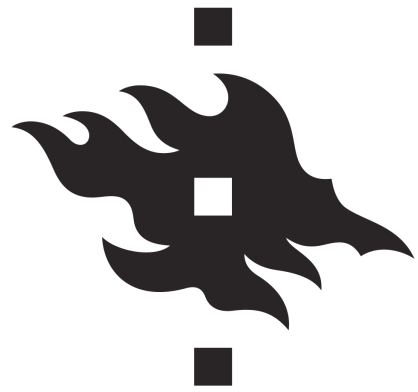
- Models are consistent with no defects

$$G\mu \xrightarrow[\text{Figueroa et al. 2013}]{} \Omega_{GW} \xrightarrow{\text{O(2) and O(3)}} \Omega_{GW} \lesssim 2 \times 10^{-15}$$

- The gravitational wave background below the LISA sensitivity window

Summary

- **Field theory simulations** of global defects $O(N)$
 - Global Strings
 - Global Monopoles
 - Work in progress with higher N 's
- **UnEqual Time Correlation** extraction
- **Transition** (Rad-Mat)
 - k independent transition in $N=2,3$
 - New approach needed for $N>3$ (*Work in progress*)
- **First CMB PS predictions and constraints**
- GWs spectrum: below LISA sensitivity



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