

Shell-crossing structure of cold dark matter with Lagrangian perturbation theory

Shohei Saga (Kyoto University)

arXiv:1805.08787 (S.S, Atsushi Taruya, Stéphane Colombi)

COSMO18, Institute for Basic Science, Daejeon, 08/27-08/31

Contents

1. Introduction

- Single/Multi-stream flow
- Vlasov solver

2. Lagrangian Perturbation Theory

- Recursion relation
- Initial condition

3. Results

- Simulation
- LPT

4. Summary

1. Introduction

Λ CDM model

Cold Dark Matter:

- Initially **negligible velocity dispersion**
 - single-stream flow = fluid approximation is valid.
- Breakdown of single-stream flow at **shell-crossing**
 - singular density field
 - after the shell-crossing, cold dark matters undergo multi-stream flow

$\delta \lesssim 1$: Quasi non-linear regime $\leftrightarrow$ single-stream flow

$\delta \gtrsim 1$: Non-linear regime $\leftrightarrow$ multi-stream flow

On small scales, **multi-stream flow** is important.

1.1 How much do we know multi-stream flow?

Simulation:

ColDICE: 6D Vlasov-Poisson solver

T.Sousbie and S.Colombi [1509.07720]

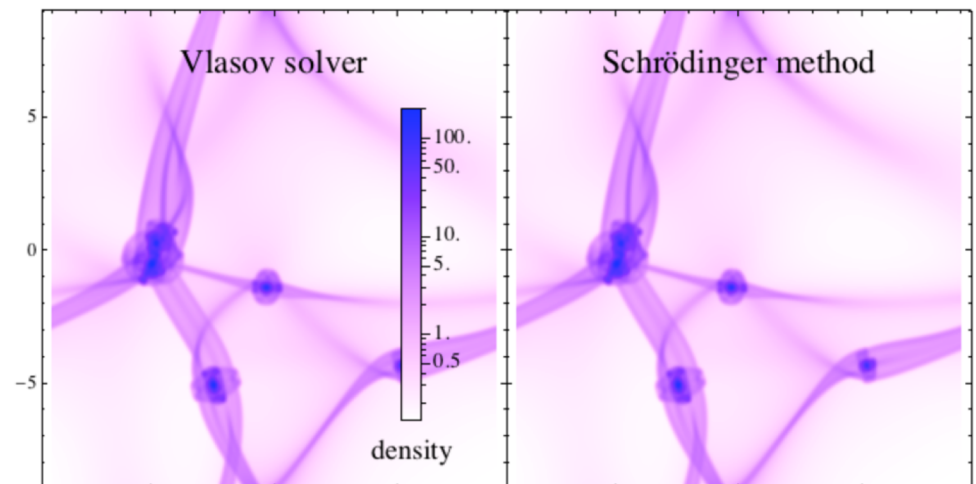
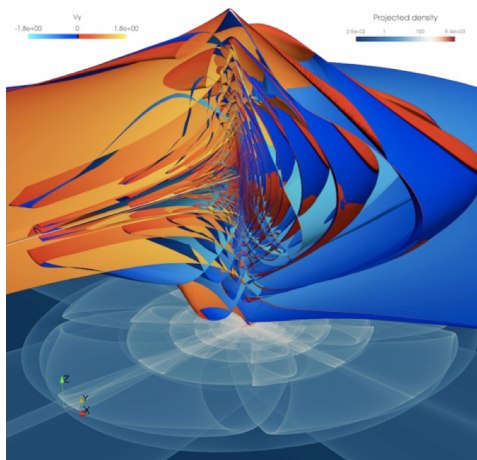
ScM: 4D Schrödinger-Poisson solver

Solve the Schrödinger-Poisson equation: wave function ψ

$$i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \psi + mV(\mathbf{r})\psi \quad \nabla^2 V = 4\pi G \psi \psi^*$$

$$f \sim |\psi|^2 + \boxed{\hbar \text{ dependent term}}$$

M.Kopp, K.Vattis, and C.Skordis [1711.00140]



1.2 Why do we need Vlasov-Poisson solver?

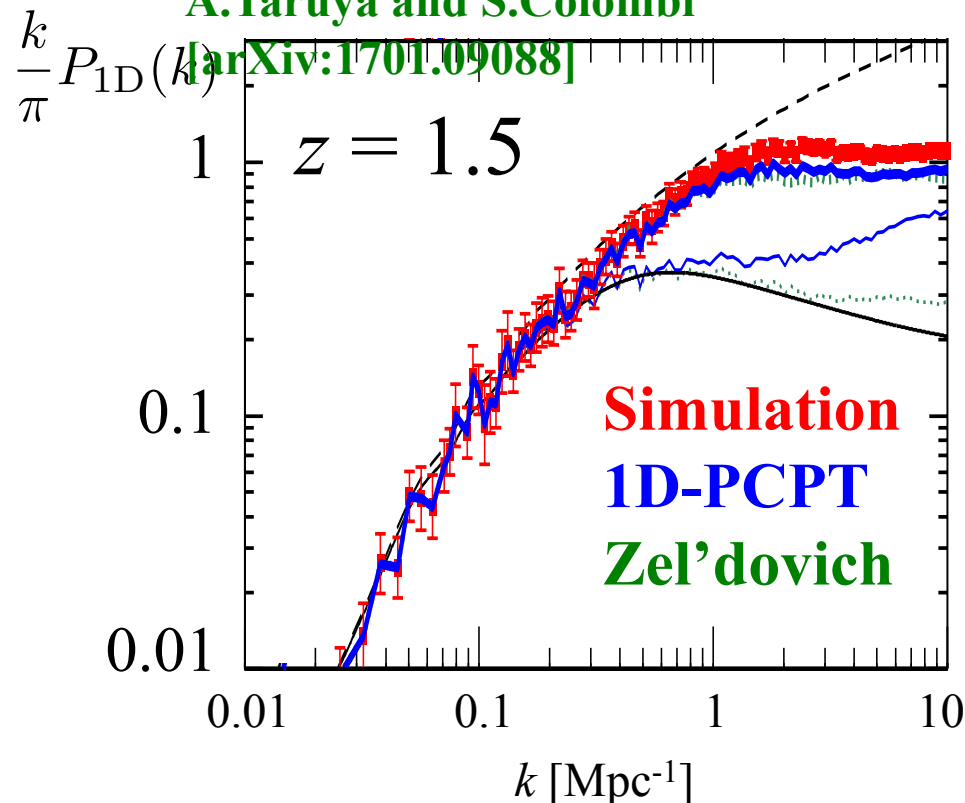
<https://vlasix.org/>

Solving in phase-space:

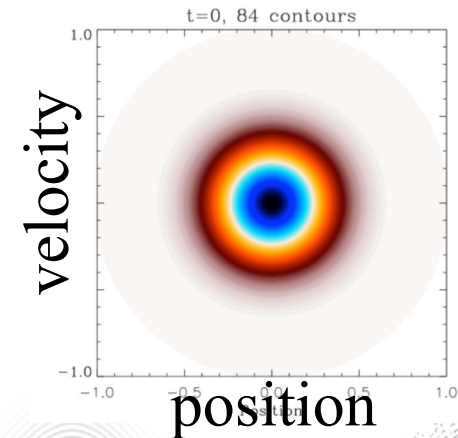
- ✓ including shell-crossing process
- ✓ without shot noise
- ✓ available in the sparse region

A.Taruya and S.Colombi

[arXiv:1701.09088]



1D collapse



Vlasov

N-body

S.Colombi et al. [1504.07337]

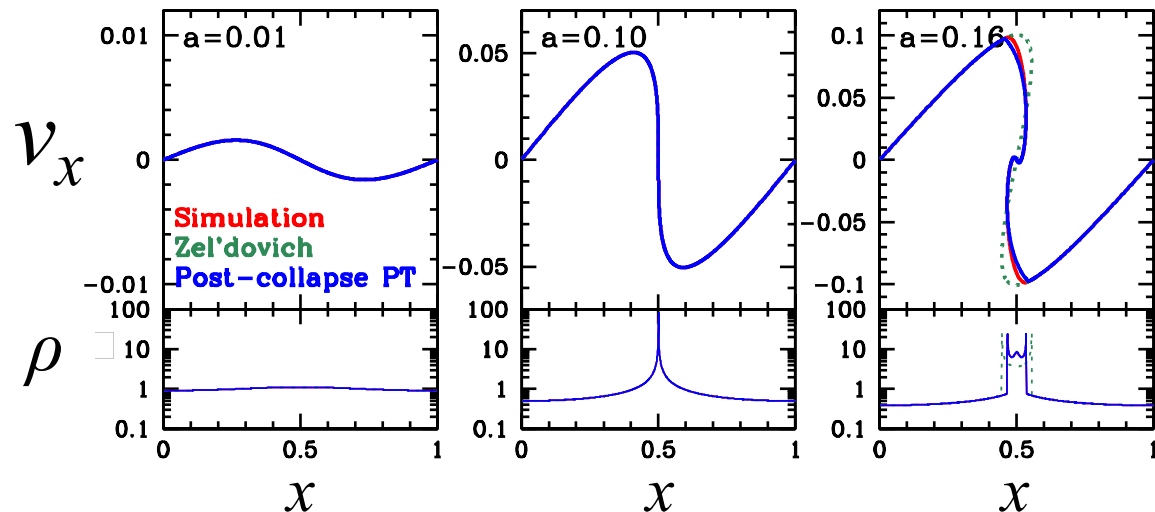
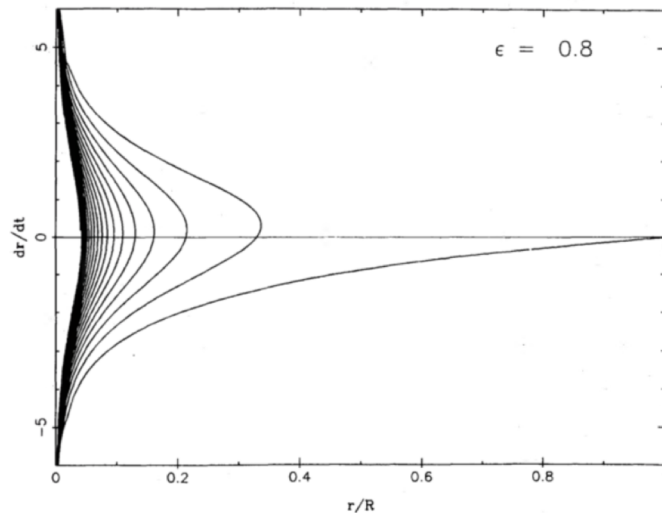
1.3 How much do we know multi-stream flow?

Theory:

- self-similar solution with a certain symmetry
e.g., Planar (1D), Cylindrical (2D), Spherical (3D)

J.A.Fillmore and P.Goldreich [ApJ,281(1984)]

- Post-collapse perturbation theory (1D)



1.4 Our work

(1D) Post-collapse perturbation theory (1D-PCPT)

Pre-collapse

Zel'dovich solution

$$x(q, t) = q + \psi(q)D_+(t)$$



Post-collapse

1D-PCPT

based on **Zel'dovich solution**
(= post-collapse dynamics)

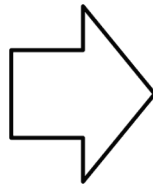
A.Taruya and S.Colombi [arXiv:1701.09088]

(2D/3D) Post-collapse perturbation theory (2D/3D-PCPT)

Pre-collapse

Perturbation

$$x(q, t) = \text{This work}$$



Post-collapse

2D/3D-PCPT

based on **?**
(= post-collapse dynamics)

2 Lagrangian Perturbation Theory

$$\mathbf{x}(\mathbf{q}, t) = \mathbf{q} + \Psi(\mathbf{q}, t)$$

Pressureless, Newtonian gravity:

Longitudinal part: $\nabla_x \cdot [\ddot{\Psi} + 2H\dot{\Psi}] = -4\pi G\bar{\rho}_m\delta$

Transverse part: $\nabla_x \times [\ddot{\Psi} + 2H\dot{\Psi}] = 0$

Nonlinear equation for the displacement field Ψ

T.Matsubara [arXiv:1505.01481]

Longitudinal part: $\left(\hat{\mathcal{T}} - 4\pi G\bar{\rho}_m\right) \Psi_{i,i} = -\epsilon_{ijk}\epsilon_{ipq}\Psi_{j,p}\left(\hat{\mathcal{T}} - 2\pi G\bar{\rho}\right)\Psi_{k,q} \quad \underline{\hspace{1cm}} \quad \text{2D}$

$$\hat{\mathcal{T}} \equiv \frac{\partial^2}{\partial t^2} + 2H\frac{\partial}{\partial t} \quad -\frac{1}{2}\epsilon_{ijk}\epsilon_{pqr}\Psi_{i,p}\Psi_{j,q}\left(\hat{\mathcal{T}} - \frac{4\pi G}{3}\bar{\rho}\right)\Psi_{k,r} \quad \underline{\hspace{1cm}} \quad \text{3D}$$

Transverse part: $\hat{\mathcal{T}}\epsilon_{ijk}\Psi_{j,k} = -\epsilon_{ijk}\Psi_{a,j}\hat{\mathcal{T}}\Psi_{a,k} \quad \underline{\hspace{1cm}} \quad \text{2D, 3D}$

$$\Psi_i = \nabla^{-2} [\nabla_i \Psi_{a,a} + \epsilon_{iab} \nabla_a (\epsilon_{bcd} \Psi_{c,d})]$$

2.1 How to solve the recursion relations?

We need to perform time- and spatial-integration.

- **Time-part**

Fastest-growing mode

$$\Psi(\mathbf{q}, \eta) = \sum_{n=1}^{\infty} D_+^n \Psi^{(n)}(\mathbf{q}) = \sum_{n=1}^{\infty} e^{n\eta} \Psi^{(n)}(\mathbf{q})$$

- **Spatial-part**

Specific initial condition (three crossed initial sine waves)

$$\Psi^{(1)}(\mathbf{q}, t) = D_+(t) \begin{pmatrix} \epsilon_x \sin q_x \\ \epsilon_y \sin q_y \\ \epsilon_z \sin q_z \end{pmatrix}$$

We finally obtain LPT solutions

up to 10th order.

3 Results

CoLDICE: Vlasov-Poisson solver

T.Sousbie and S.Colombi [1509.07720]

<https://vlasix.org/>

$$\frac{\partial f(t, \mathbf{x}, \mathbf{u})}{\partial t} + \mathbf{u} \cdot \nabla f(t, \mathbf{x}, \mathbf{u}) - \nabla \phi \cdot \nabla f(t, \mathbf{x}, \mathbf{u}) = 0$$

$$\nabla^2 \phi = 4\pi G \int d^3 \mathbf{u} f(t, \mathbf{x}, \mathbf{u})$$

Initial condition:

$$\Psi^{(1)}(\mathbf{q}, t) = D_+(t) \begin{pmatrix} \epsilon_x \sin q_x \\ \epsilon_y \sin q_y \\ \epsilon_z \sin q_z \end{pmatrix}$$

Q1D

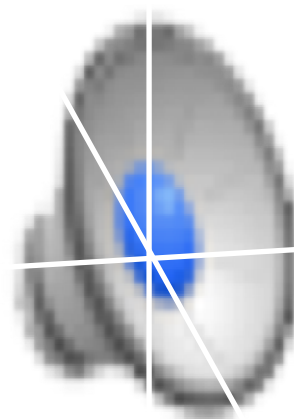
$$(\epsilon_x, \epsilon_y, \epsilon_z) = (-24, -4, -3)$$

ASY

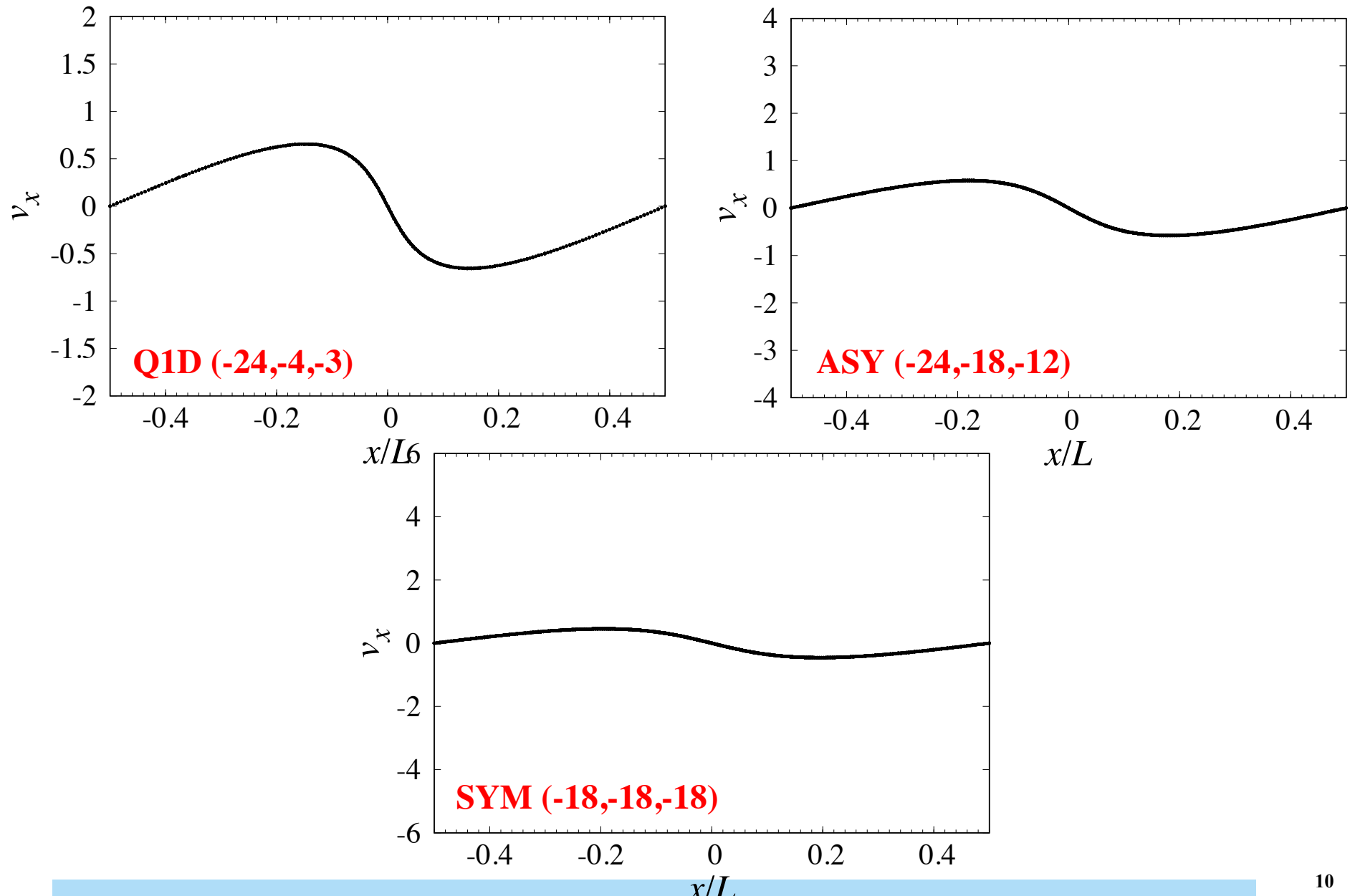
$$(\epsilon_x, \epsilon_y, \epsilon_z) = (-24, -18, -12)$$

SYM

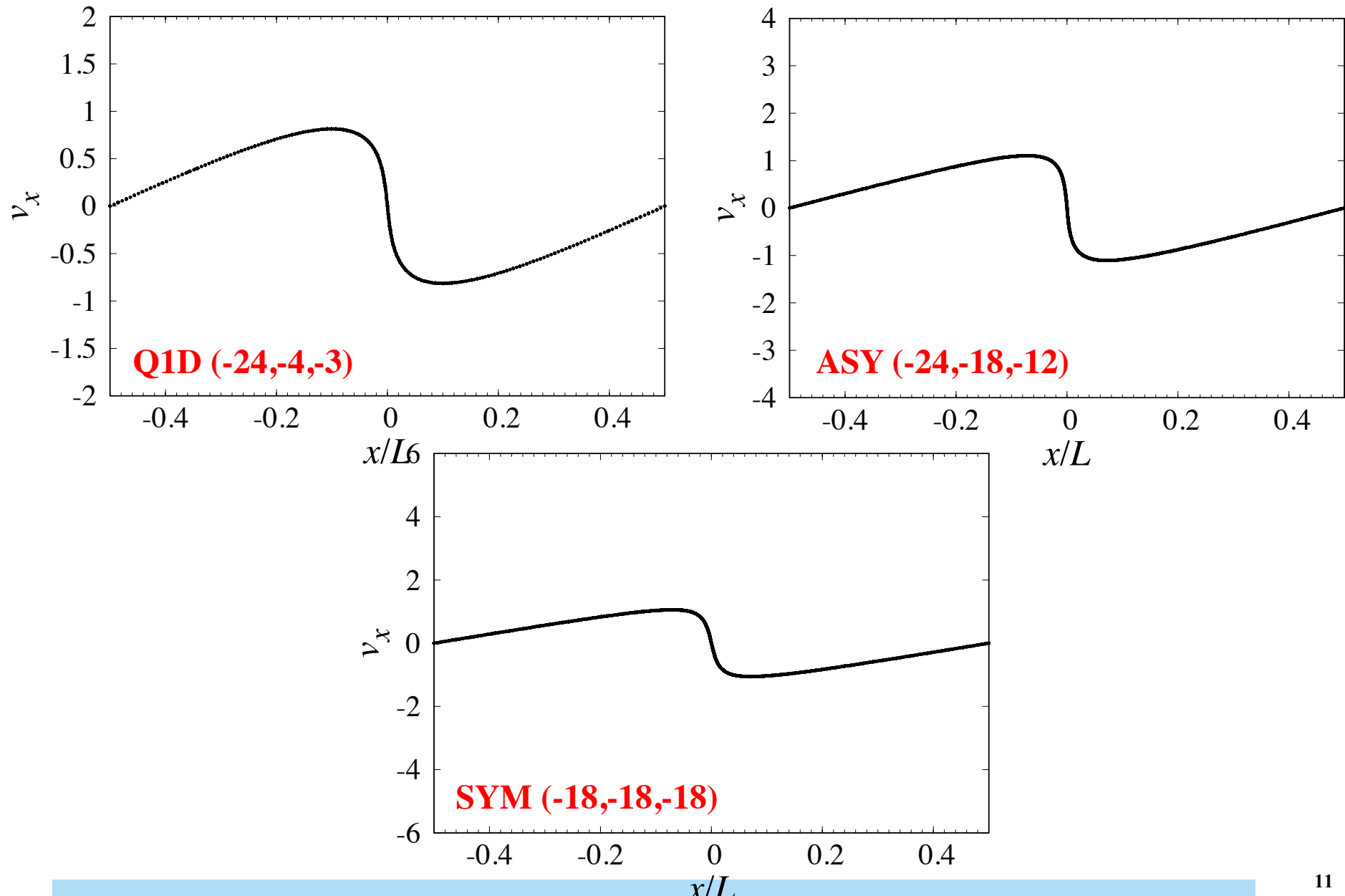
$$(\epsilon_x, \epsilon_y, \epsilon_z) = (-18, -18, -18)$$



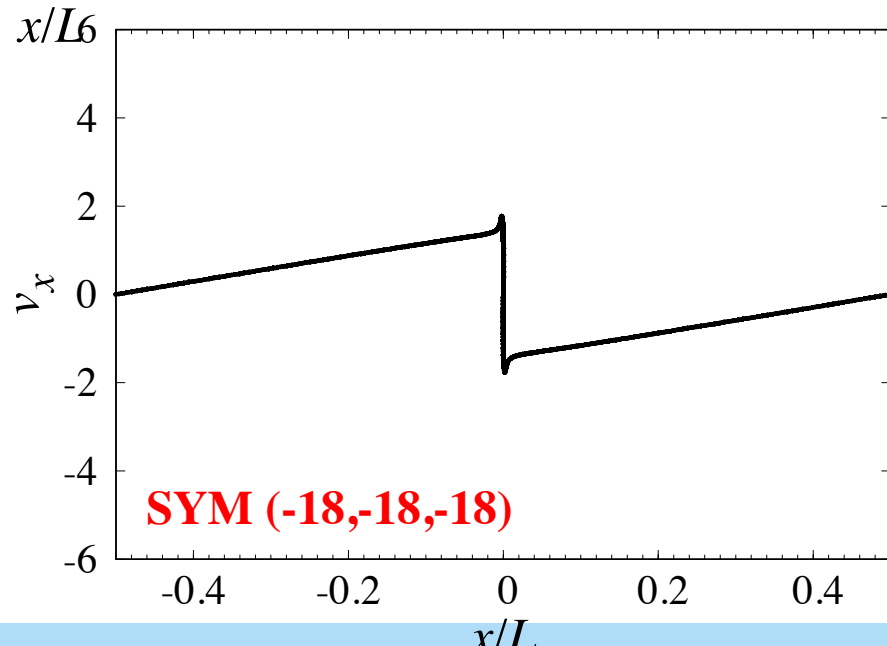
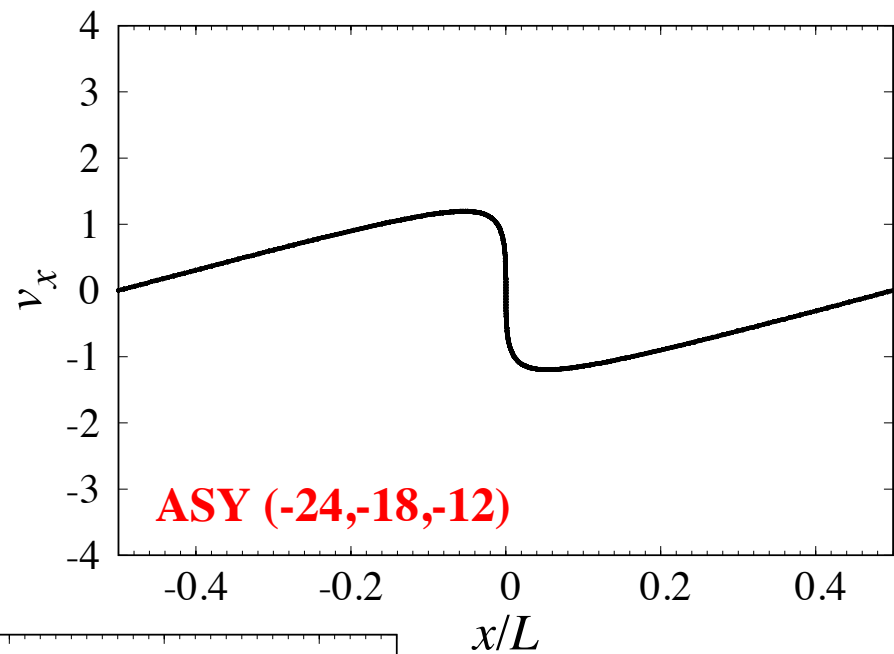
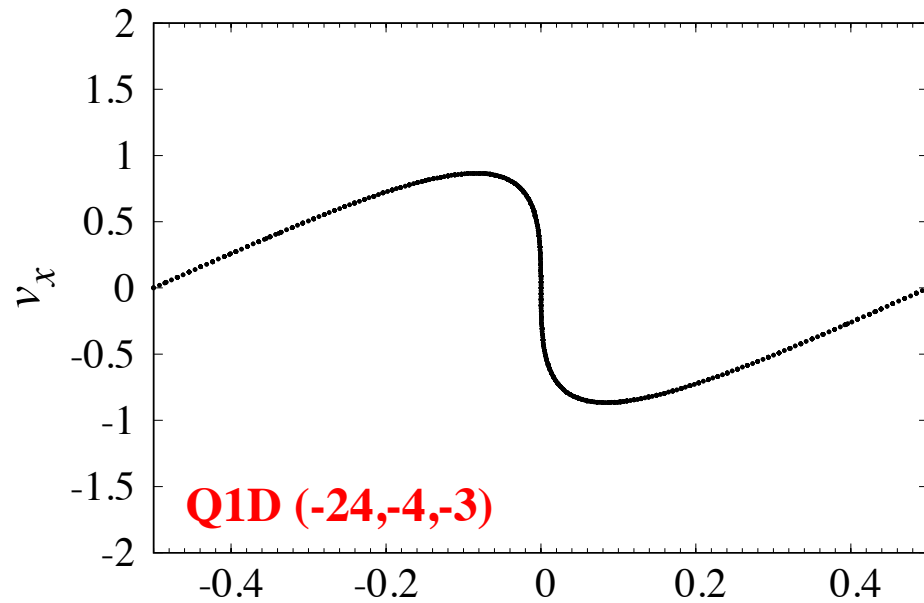
3.1 Vlasov-Poisson simulation (Pre-collapse)



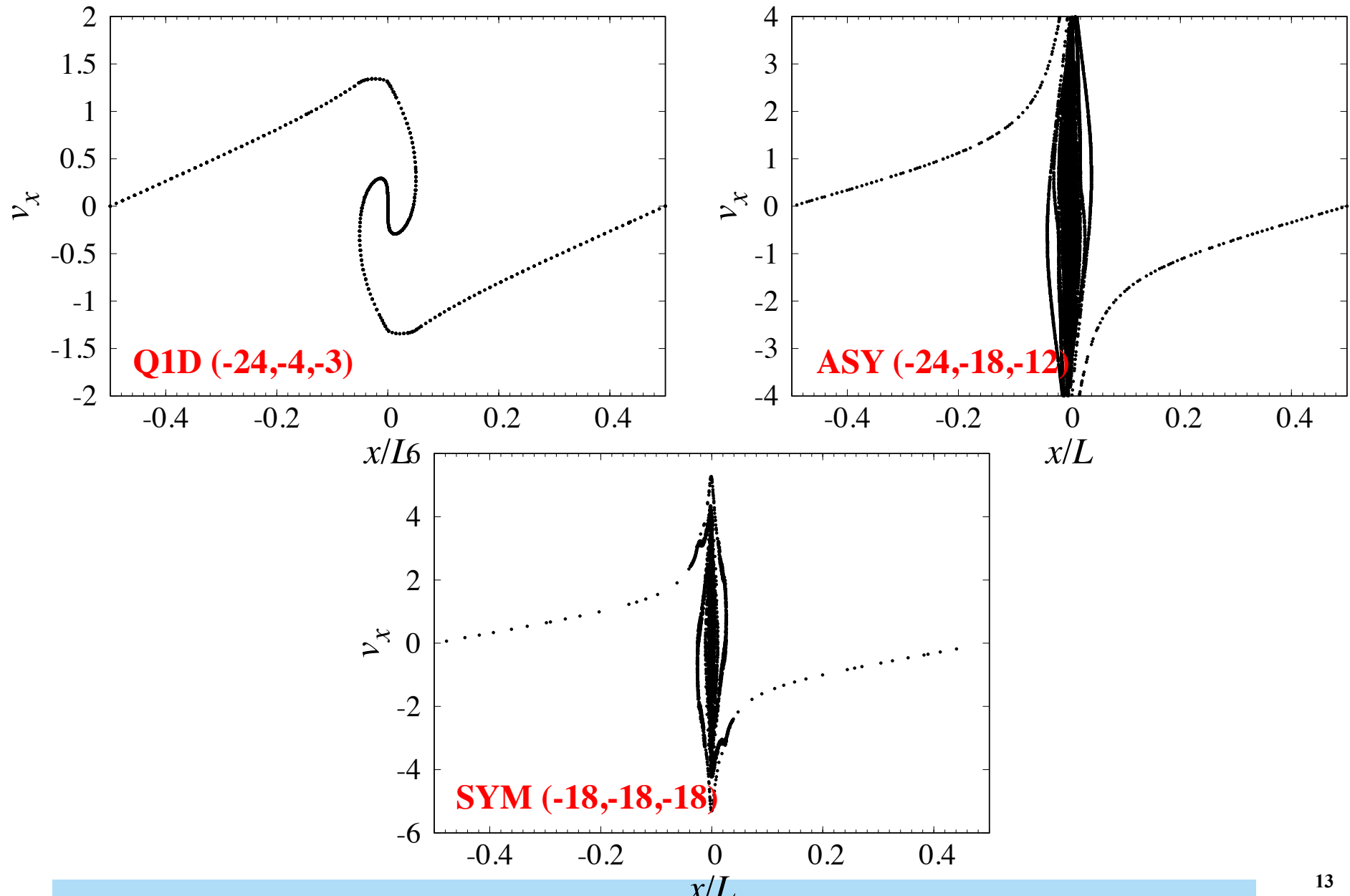
3.1 Vlasov-Poisson simulation (Pre-collapse)



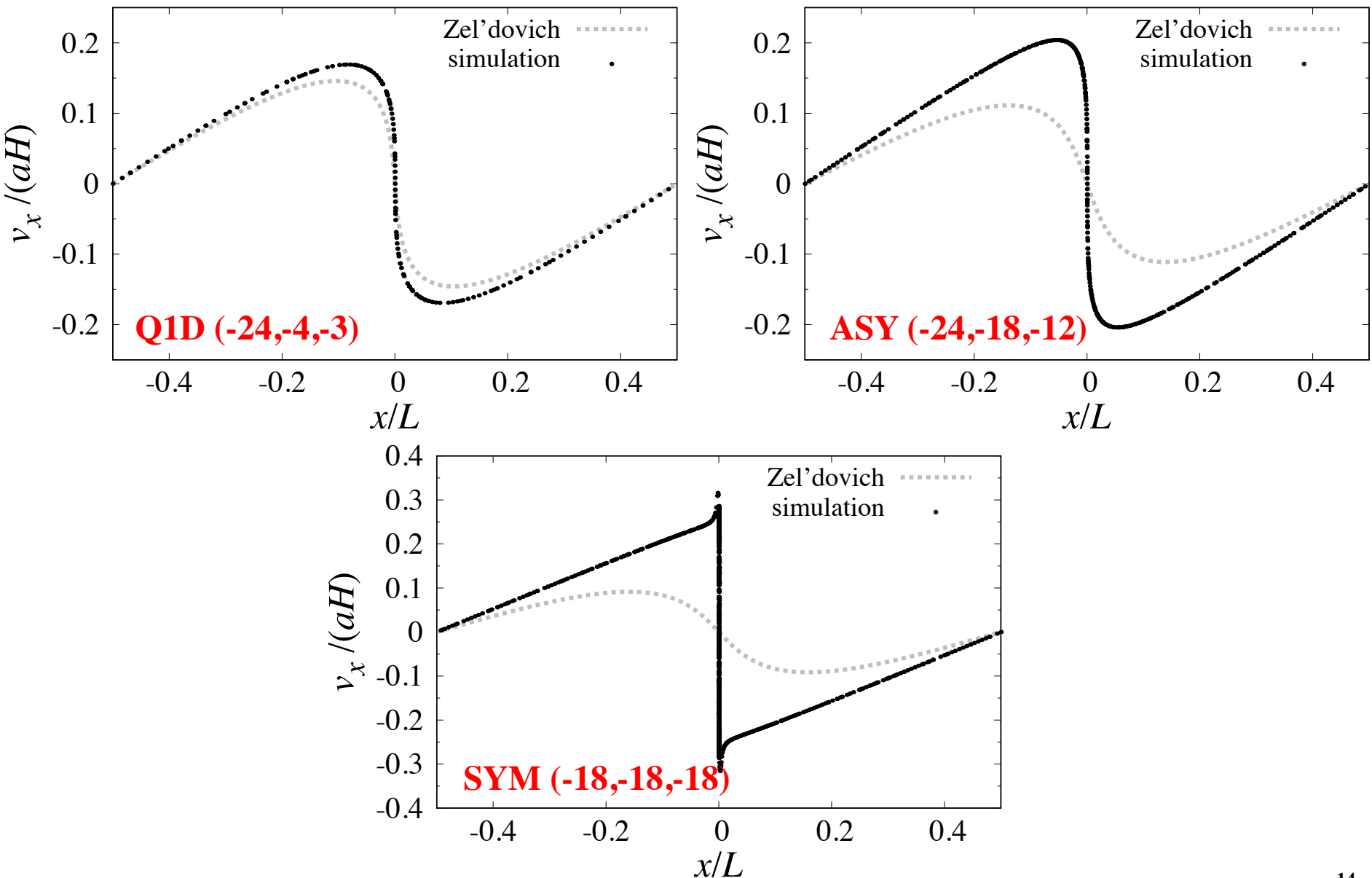
3.1 Vlasov-Poisson simulation (Shell-crossing)



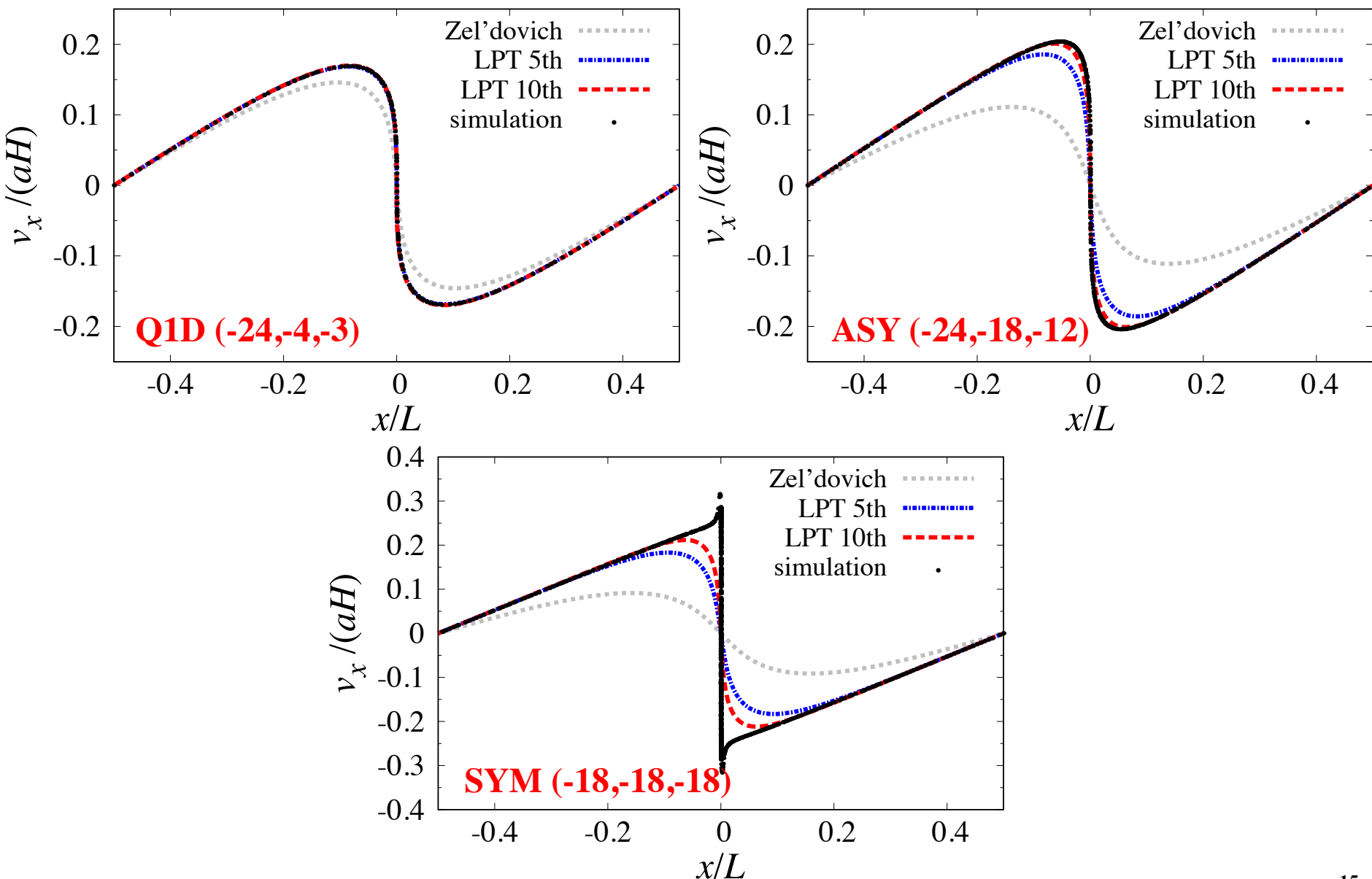
3.1 Vlasov-Poisson simulation (Post-collapse)



3.2 Zel'dovich approximation at shell-crossing



3.3 10th LPT at shell-crossing



3.4 Fitting function and Extrapolation

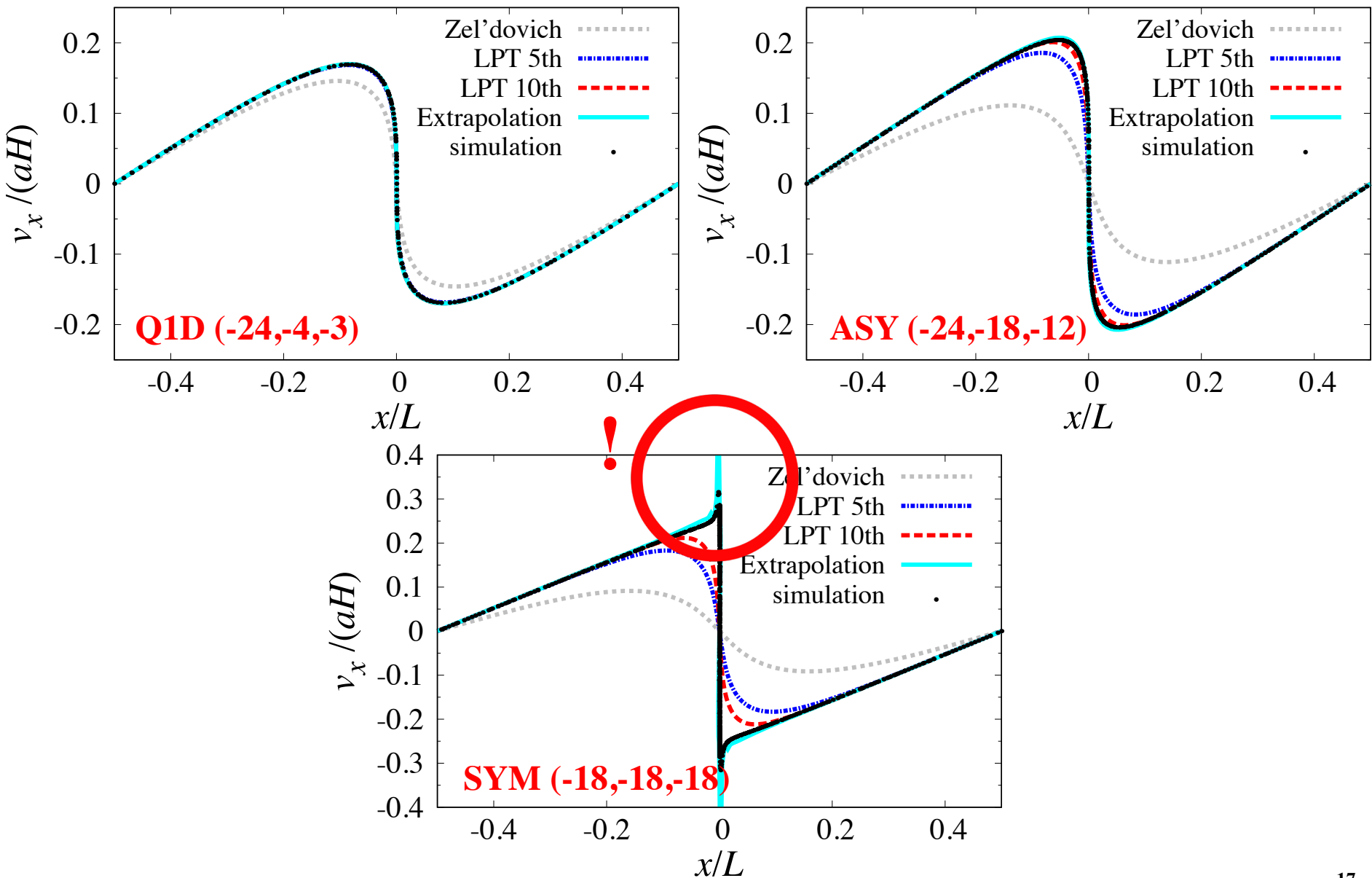
We found the fitting formula $x(\mathbf{q}, n)$ as a function of order n .

$$\begin{aligned} x_i(\mathbf{q}, t, n) &= q_i + \sum_{a=1}^n \Psi_i^{(a)}(\mathbf{q}, t) \\ &= a_i(\mathbf{q}, t) + \frac{1}{b_i(\mathbf{q}, t) + c_i(\mathbf{q}, t) \exp [d_i(\mathbf{q}, t) n^{e_i(\mathbf{q}, t)}]} \end{aligned}$$

(in an empirical way)

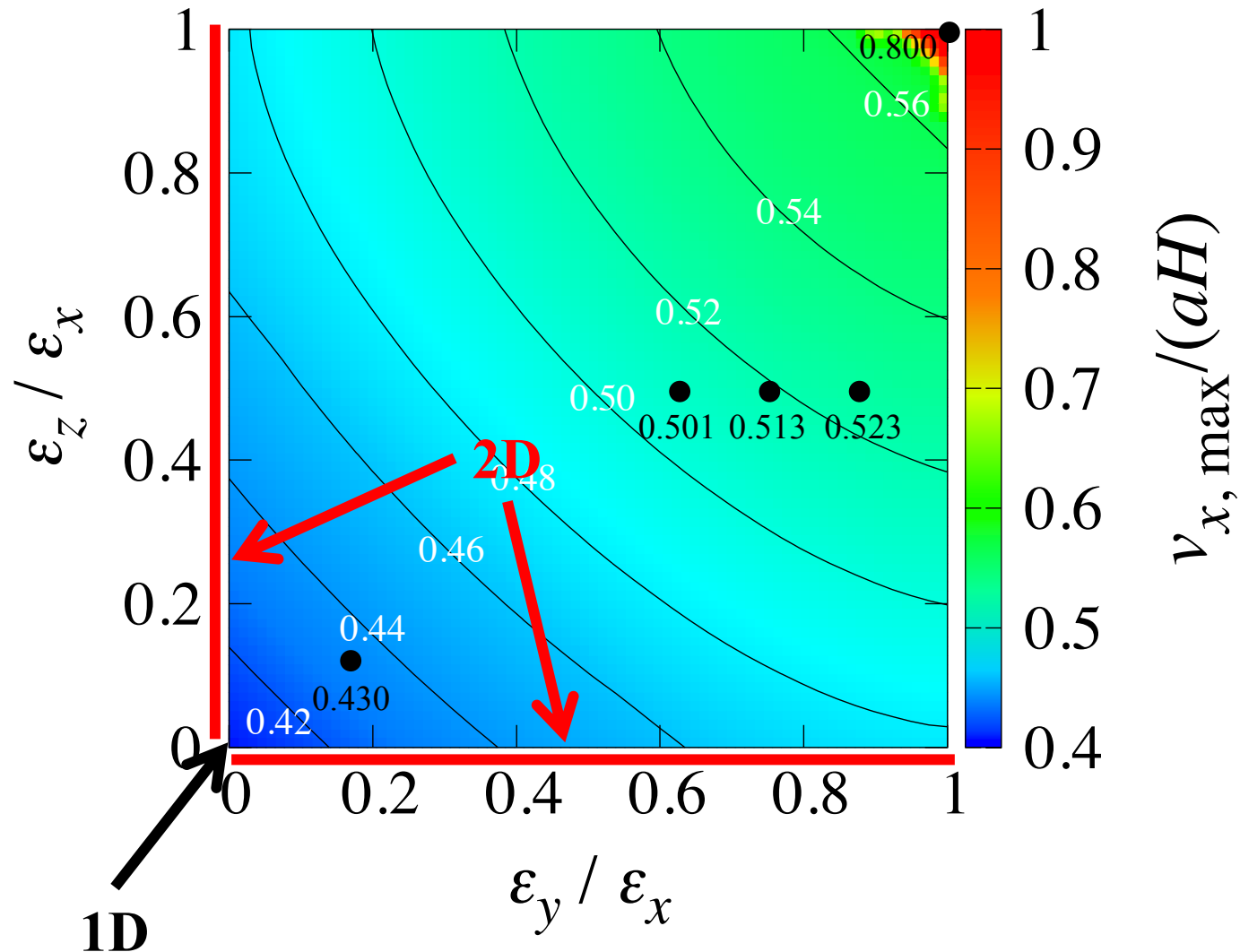
$n \rightarrow \infty$: “Effective” infinite-order LPT

3.5 Extrapolation at Shell-crossing



3.6 Exploring parameter space@Shell-crossing

Focusing on maximum velocity as a function of $(\epsilon_y/\epsilon_x, \epsilon_z/\epsilon_x)$.



4. Summary

- We consider primordial dark matter halos seeded by three crossed initial sine waves.
- We perform **10th order LPT and beyond**.
- Except for rare cases such as the triaxial symmetric configuration, **collapse is generally expected to produce a planar singularity**.

1D $\rightarrow$ 2D: No transitions

2D $\rightarrow$ 3D: Spiky feature appears!