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高エネルギー加速器研究機構



Primordial Black Holes in Matter Dominated Universe

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Why PBHs?

- We need $30 M_{\odot}$ BHs to explain LIGO/VIRGO GW events
- We do not know the origin of $30 M_{\odot}$ BHs to be astrophysical or cosmological.
- PBHs should be a good candidate for Cold Dark Matter (CDM), but we need to know the full cosmic history.
- Some inflation modes predict PBHs formed at small scales in the early Universe (before 1 sec)
- Scenarios have been constrained by **BBN**, **CMB anisotropies**, **lensing**, **gamma-ray**, and so on. *Carr, Kohri, Sendouda, J.Yokoyama (2010)*
- In future, scenarios can be investigated further by **PIXIE** (CMB μ -distortion), **SKA/Ominiscope** (21cm), **CTA** (gamma-ray), **DECIGO** (Gravitational Wave), ...

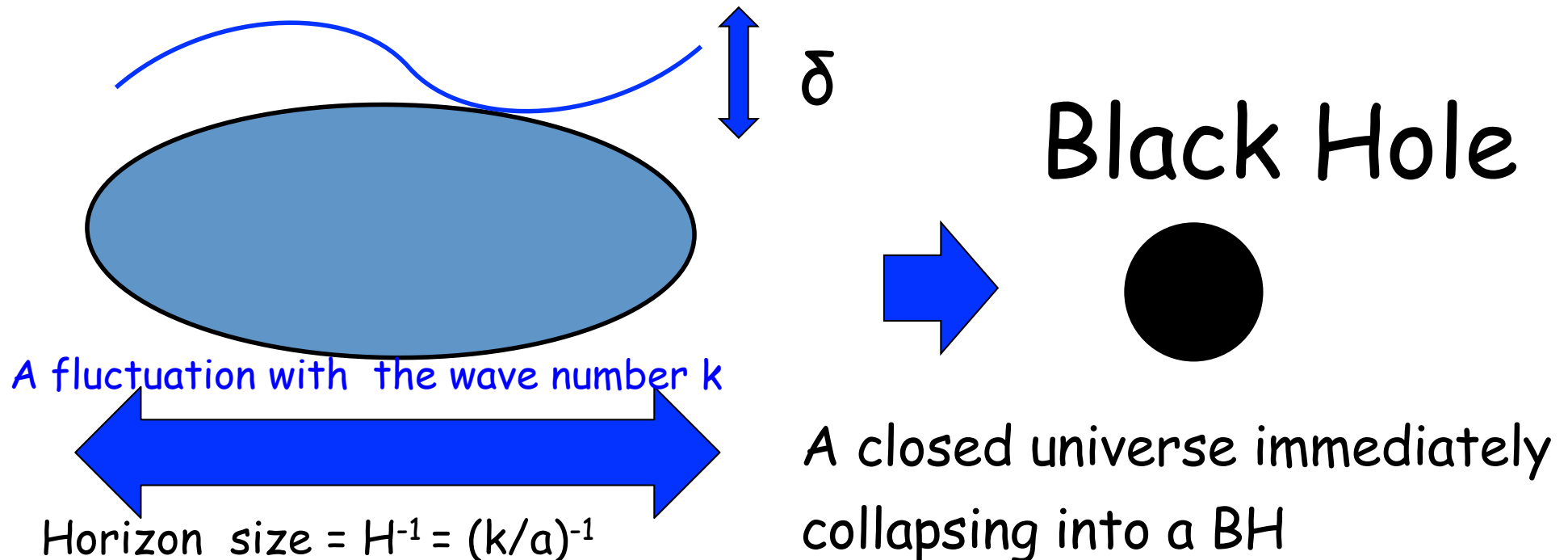
Conditions for a PBH formation in Radiation dominated (RD) Universe

Zel'dovich and Novikov (1967), Hawking (1971), Carr (1975)

Harada, Yoo and KK (2013)

- Gravity could be stronger than pressure

$$\delta > \delta_c \sim p / \rho \sim c_s^2 = w = 1/3$$



$P_\zeta(k)$ and PBH abundance $\beta(M)$

- Fraction of PBH to the total at its formation epoch with Gaussian fluctuation.

$$\beta(M) \equiv \frac{\rho_{\text{PBH}}(M)}{\rho_{\text{tot}}} = 2 \int_{\delta_{\text{th}}}^{\infty} d\delta \frac{1}{\sqrt{2\pi\sigma^2}} \exp\left(-\frac{\delta^2}{2\sigma^2}\right) = \text{erfc}\left(\frac{\delta_{\text{th}}}{\sqrt{2}\sigma}\right)$$

δ_{th} $\sim 1/3 - 0.5$

- Finally we have a relation between β and fluctuation σ (or β and Ω)

$$\beta(M) \sim \text{erfc}\left(\frac{\delta_{\text{th}}}{\sqrt{2}\sigma}\right) \simeq \sqrt{\frac{2}{\pi}} \frac{\sigma}{\delta_{\text{th}}} \exp\left(-\frac{\delta_{\text{th}}^2}{2\sigma^2}\right)$$

$$= 1.5 \times 10^{-18} \left(\frac{m_{\text{PBH}}}{10^{15} \text{ g}}\right)^{1/2} \left(\frac{\Omega_{\text{PBH}} h^2}{0.1}\right)$$

$\sim P_\zeta$

Typical quantities of PBHs in RD

- Mass (horizon mass = $\rho(t_{\text{form}}) H(t_{\text{form}})^{-3}$)

$$M_{\text{PBH}} \sim M_{\text{pl}}^2 t_{\text{form}} \sim \frac{M_{\text{pl}}^3}{T_{\text{form}}^2} \sim 10^{15} \text{ g} \left(\frac{T_{\text{form}}}{3 \times 10^8 \text{ GeV}} \right)^{-2} \sim 30 M_{\odot} \left(\frac{T_{\text{form}}}{40 \text{ MeV}} \right)^{-2}$$

- Lifetime

$$\tau_{\text{PBH}} \sim \frac{M_{\text{PBH}}^3}{M_{\text{pl}}^4} \sim 4 \times 10^{17} \text{ sec} \left(\frac{M_{\text{PBH}}}{10^{15} \text{ g}} \right)^3 \sim 3 \times 10^{68} \text{ yrs} \left(\frac{M_{\text{PBH}}}{30 M_{\odot}} \right)^3$$

- Hawking Temperature

$$T_{\text{PBH}} \sim \frac{M_{\text{pl}}^2}{M_{\text{PBH}}} \sim 0.1 \text{ MeV} \left(\frac{M_{\text{PBH}}}{10^{15} \text{ g}} \right)^{-1} \sim 3 \times 10^{-11} \text{ K} \left(\frac{M_{\text{PBH}}}{30 M_{\odot}} \right)^{-1}$$

- Wave number of horizon length

$$k = aH \sim 10^5 \text{ Mpc}^{-1} \left(\frac{M_{\text{PBH}}}{10^4 M_{\odot}} \right)^{-1/2} \sim 10^5 \text{ Mpc}^{-1} \left(\frac{T_{\text{form}}}{\text{MeV}} \right)^{+1}$$

- Fraction to CDM

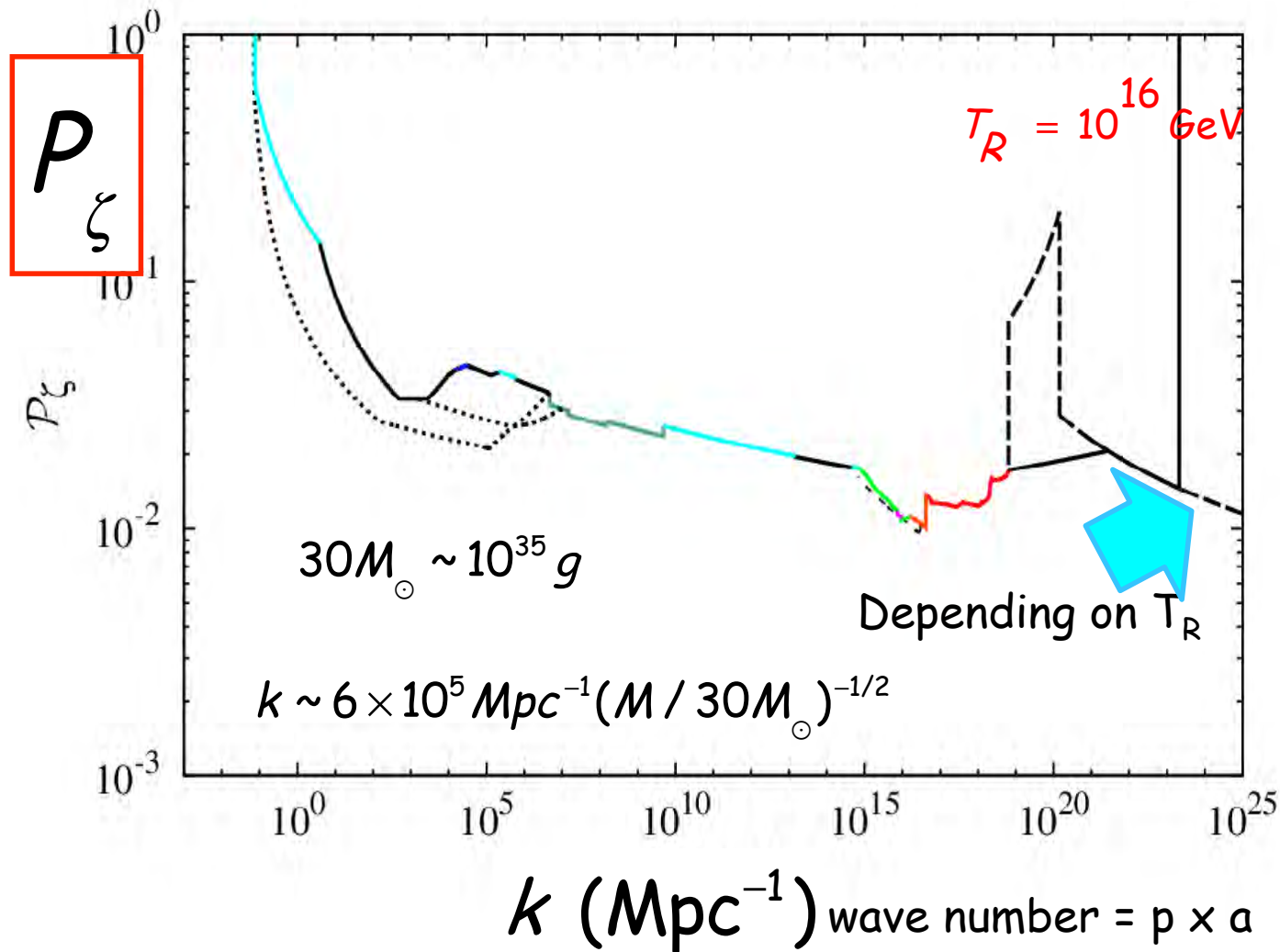
$$f_{\text{fraction}} \equiv \frac{\Omega_{\text{PBH}}}{\Omega_{\text{CDM}}} \sim \left(\frac{\beta}{10^{-18}} \right) \left(\frac{M_{\text{PBH}}}{10^{15} \text{ g}} \right)^{-1/2} \sim \left(\frac{\beta}{10^{-8}} \right) \left(\frac{M_{\text{PBH}}}{30 M_{\odot}} \right)^{-1/2} \sim 10^8 \left(\frac{M_{\text{PBH}}}{30 M_{\odot}} \right)^{-1/2} \sqrt{P_{\delta}} \exp \left[-\frac{1}{18 P_{\delta}} \right]$$

Upper bounds on P_ζ for PBH

Carr, KK, Sendouda, J.Yokoyama (2010)

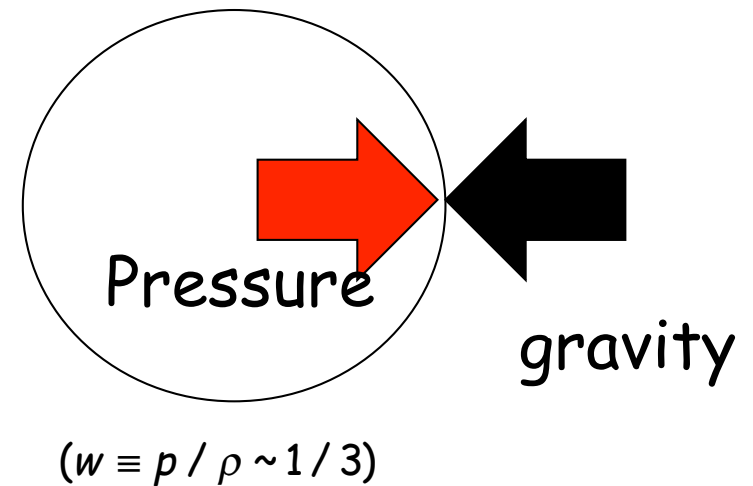
Alabidi, KK, Sasaki, Sendouda (2012)

Amplitude of curvature perturbation



Features of PBH formations in RD

- Perfectly spherical due to radiation pressure



- No evolutions of density perturbations
- Small angular momentum

PBH formation at the (early) matter dominated (MD) Universe

Polnarev and Khlopov (1982)

Harada, yoo, KK, Nakao, Jhingan (2016)

- **Pressure is zero**, which could induce an immediate collapse and producing more PBHs?
- **Density perturbations can evolve**, which produces non-spherical objects and cannot be included inside the Horizon radius. That means less PBHs can be produced?

Matter domination

- Three radius in Lagrangian coordinate q_i

$$r_1 = (a - \alpha b)q_1 \quad \text{Zel'dovich Approximation}$$

$$r_2 = (a - \beta b)q_2$$

$$r_3 = (a - \gamma b)q_3$$

- Eccentricity

$$e^2 = 1 - \left(\frac{r_2(t_c)}{r_3(t_c)} \right)^2 = 1 - \left(\frac{\alpha - \beta}{\alpha - \gamma} \right)^2$$

- Hoop

$$\mathcal{C} = 16 \left(1 - \frac{\gamma}{\alpha} \right) E \left(\sqrt{1 - \left(\frac{\alpha - \beta}{\alpha - \gamma} \right)^2} \right) r_f$$

- Hoop conjecture for PBH production

$$\mathcal{C} \lesssim 2\pi r_g$$

Abundance of PBHs formed in MD

- Probability distribution by peak statistics (BBKS)

Doroshkevich (1970)

$$\begin{aligned}
 & w(\alpha, \beta, \gamma) d\alpha d\beta d\gamma \\
 &= -\frac{27}{8\sqrt{5}\pi\sigma_3^6} \exp \left[-\frac{1}{10\sigma_3^2} (\alpha + \beta + \gamma)^2 - \frac{1}{4\sigma_3^2} \{(\alpha - \beta)^2 + (\beta - \gamma)^2 + (\gamma - \alpha)^2\} \right] \\
 & \cdot (\alpha - \beta)(\beta - \gamma)(\gamma - \alpha) d\alpha d\beta d\gamma.
 \end{aligned}$$

$\sigma_H = \sqrt{5}\sigma_3$

- Probability

$$\beta_0 = \int_0^\infty d\alpha \int_{-\infty}^\alpha d\beta \int_{-\infty}^\beta d\gamma \theta(1 - h(\alpha, \beta, \gamma)) w(\alpha, \beta, \gamma)$$

$$h(\alpha, \beta, \gamma) = \frac{2}{\pi} \frac{\alpha - \gamma}{\alpha^2} E \left(\sqrt{1 - \left(\frac{\alpha - \beta}{\alpha - \gamma} \right)^2} \right)$$

$$h(\alpha, \beta, \gamma) := \mathcal{C} / (2\pi r_g)$$

Effects by finite angular momentum

Harada, Yoo, KK, Nakao (2017)

- Probability distribution

$$a_* := L/(GM^2/c)$$
$$f_{\text{BH}(2)}(a_*) da_* \propto \frac{1}{a_*^{5/3}} \exp\left(-\frac{1}{2\sigma_H^{2/3}} \left(\frac{2}{5}\mathcal{I}\right)^{4/3} \frac{1}{a_*^{4/3}}\right) da_*$$

- Probability

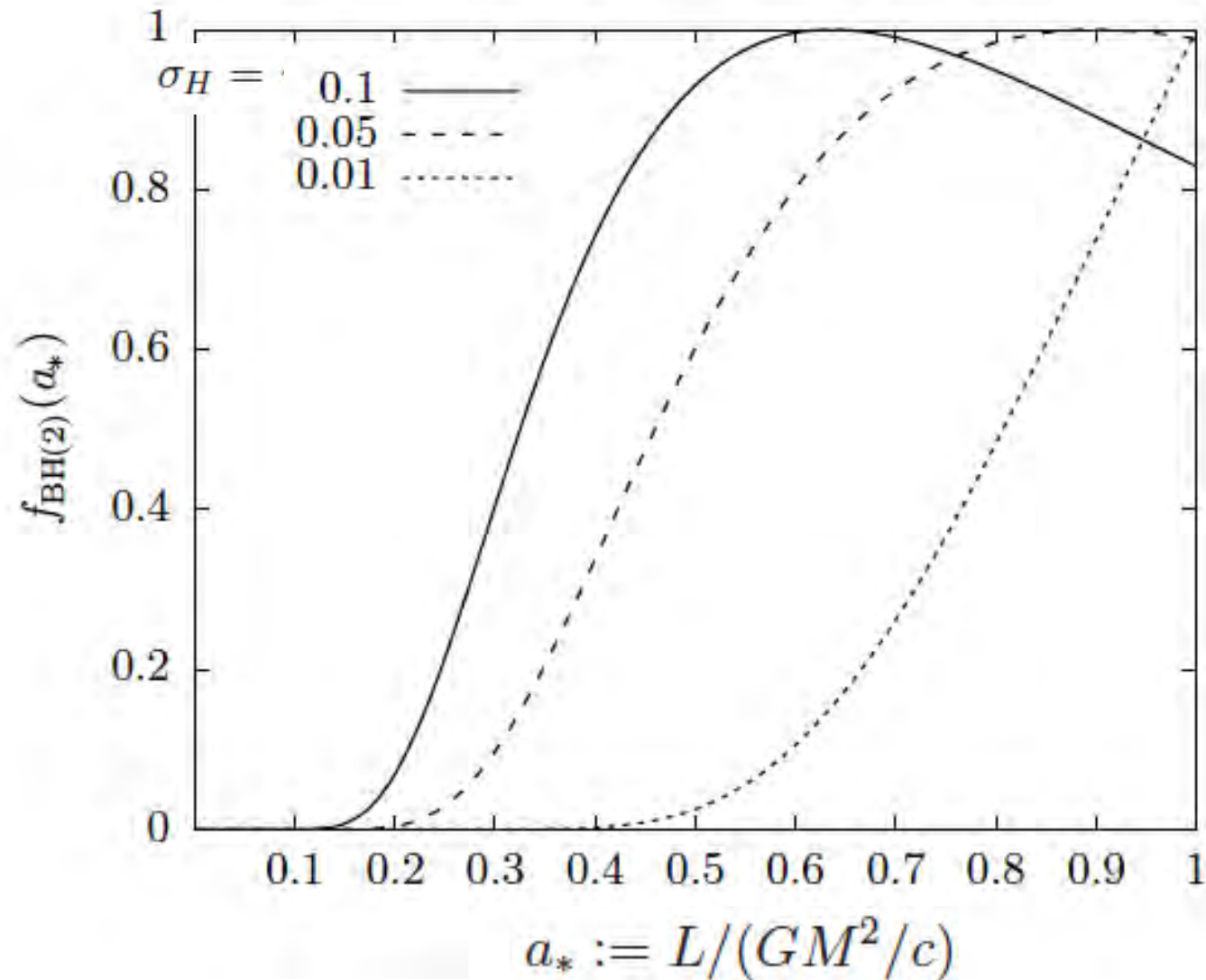
$$\beta_0 \simeq \int_0^\infty d\alpha \int_{-\infty}^\alpha d\beta \int_{-\infty}^\beta d\gamma \theta[\delta_H(\alpha, \beta, \gamma) - \delta_{\text{th}}] \theta[1 - h(\alpha, \beta, \gamma)] w(\alpha, \beta, \gamma)$$

$$\delta_{,H}(\alpha, \beta, \gamma) = \alpha + \beta + \gamma \quad \delta_{\text{th}} := \left(\frac{2}{5}\mathcal{I}\sigma_H\right)^{2/3}$$

Spin distribution

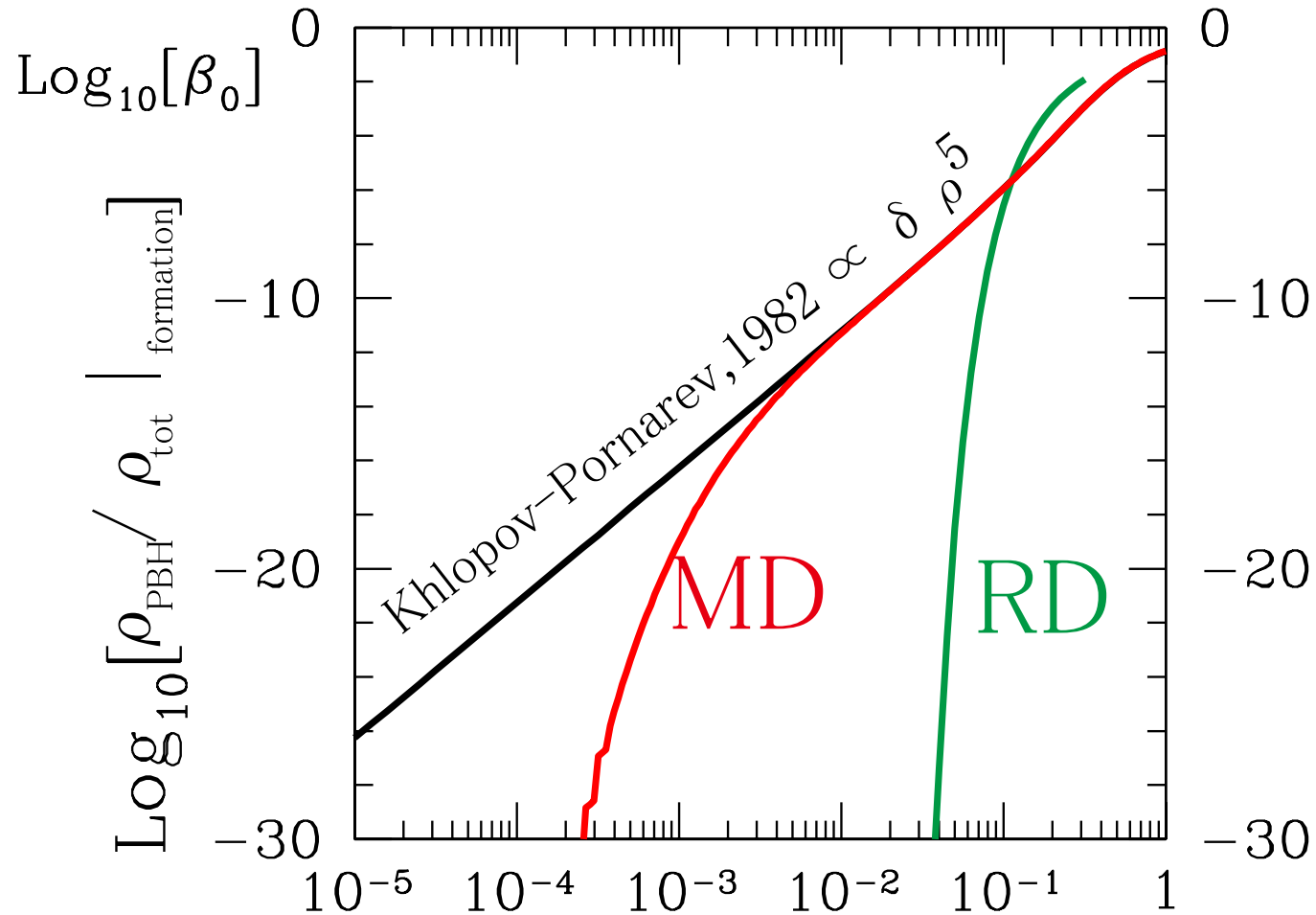
More highly-spinning halos cannot collapse into PBHs, which means that the PBHs produced tend to have high spins in MD

Harada, Yoo, KK, Nakao (2017)



Beta in matter-domination

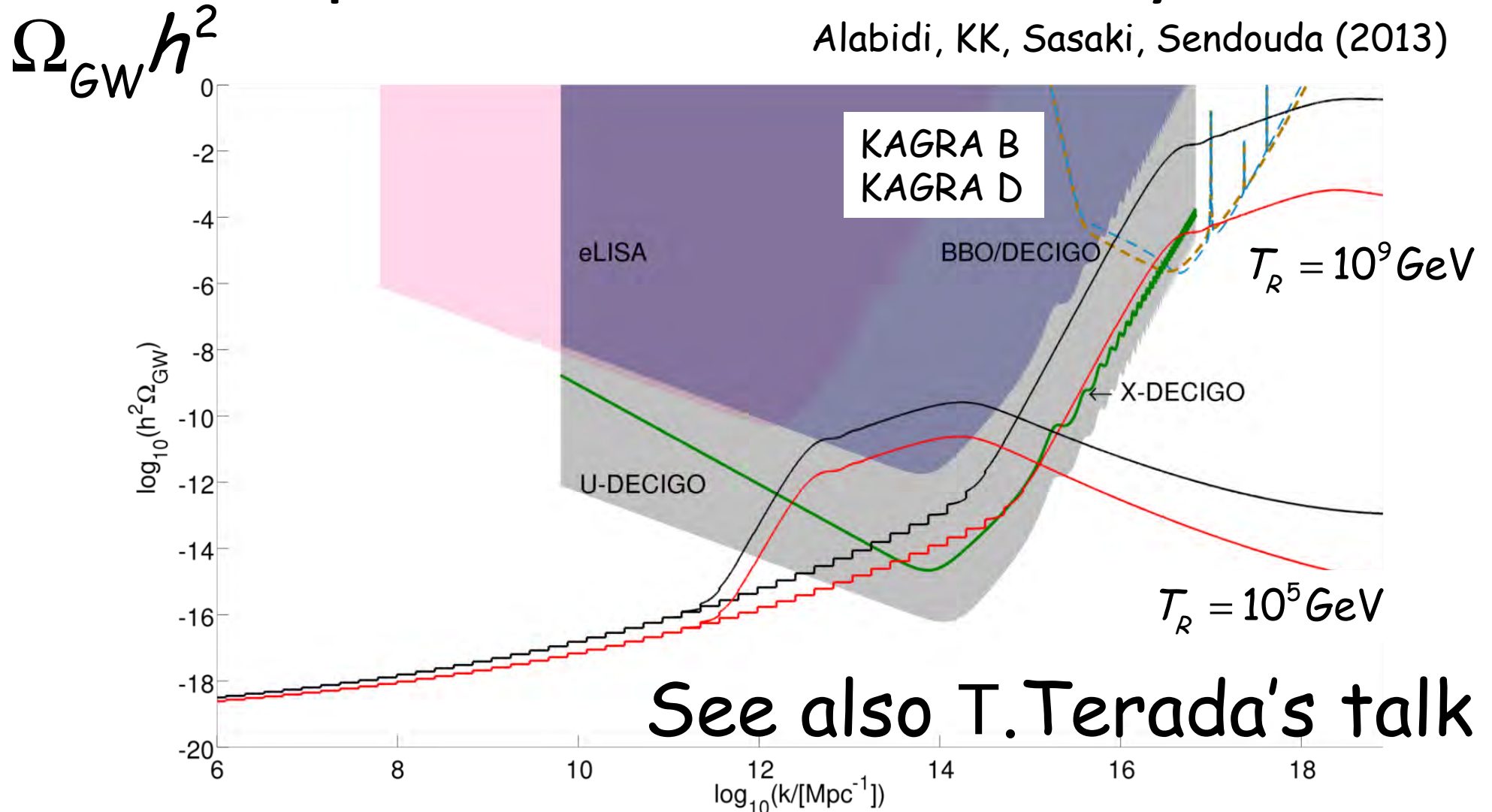
Harada, Yoo, KK, Nakao (2017)



$$\sigma_H = \sqrt{5}\sigma_3 = \delta \rho / \rho$$

In Hilltop models with low reheating temperature with the early MD

Alabidi, KK, Sasaki, Sendouda (2013)



See also T.Terada's talk

$k (Mpc^{-1})$

See also Nakayama, Saito, Suwa, Yokoyama (2008)

CMB bound on PBHs by disk-accretion in the late MD epoch

Poulin, Serpico, Calore, Clesse, KK (2017)

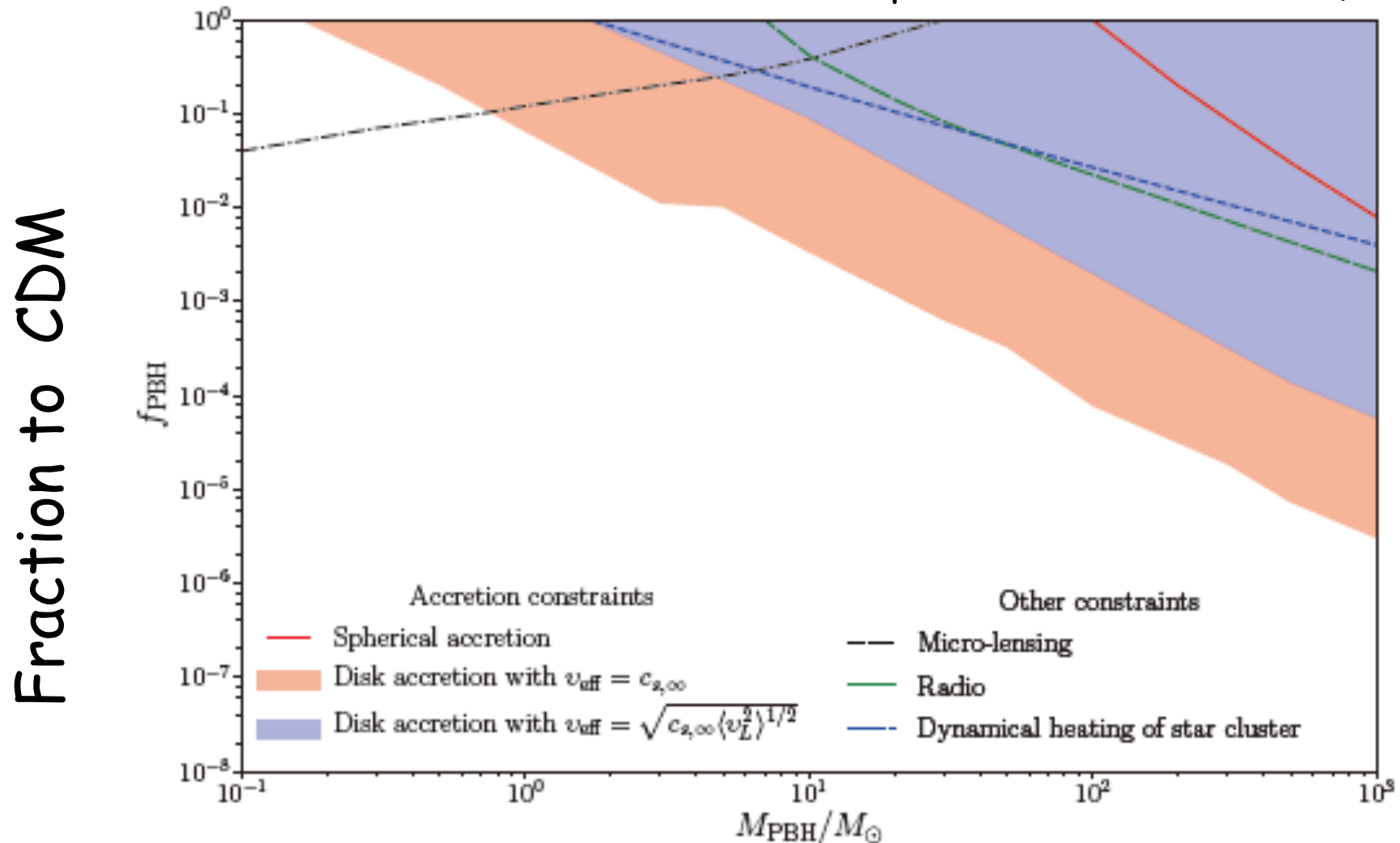
- A non-spherical accretion disk around a PBH caused by an angular momentum emits radiation

$$\dot{M}_{\text{HB}} \equiv 4\pi\lambda\rho_{\infty}v_{\text{eff}}r_{\text{HB}}^2 \equiv 4\pi\lambda\rho_{\infty}\frac{(GM)^2}{v_{\text{eff}}^3}$$
$$l \simeq \omega r_{\text{HB}}^2 \simeq \left(\frac{\delta\rho}{\rho} + \frac{\delta v}{v_{\text{eff}}}\right)v_{\text{eff}}r_{\text{HB}}$$

- CMB anisotropies are affected
- From observations, we can constrain the number density of PBHs.

CMB bound by disk-accretion in the latest MD epoch

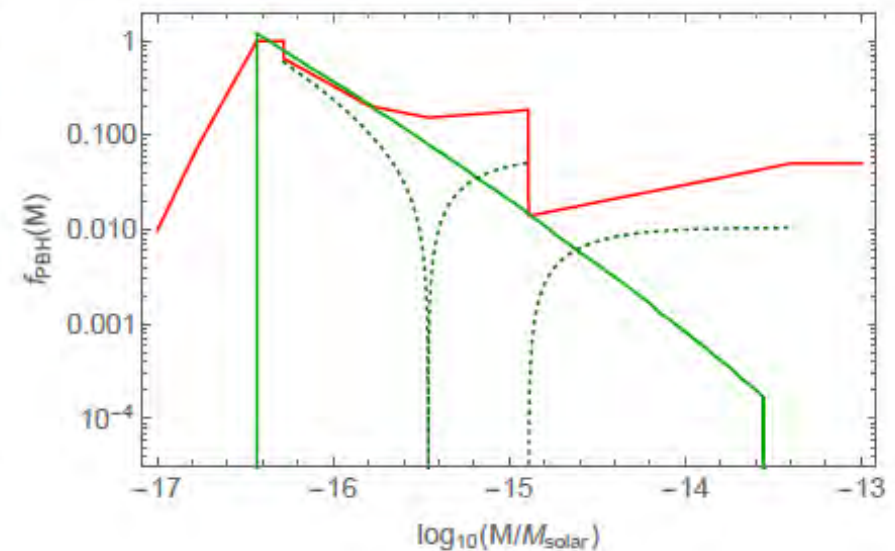
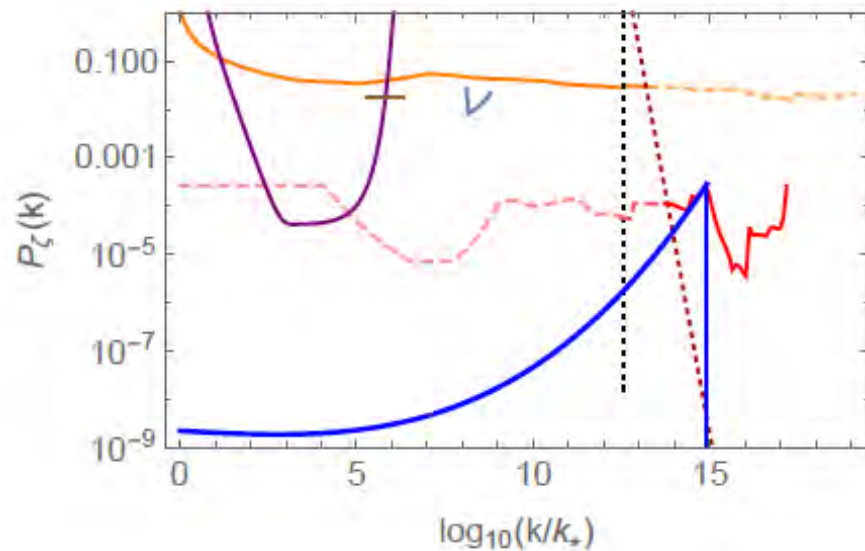
Poulin, Serpico, Calore, Clesse, KK (2017)



100 % Dark Matter by PBHs

KK and T.Terada, 2018

$$n_s = 0.96, \alpha_s = 0, \beta_s = 0.0019485.$$



black dotted line shows
 $T_R = 10^4 \text{ GeV}$.

Summary

- PBH can be formed at small scales even in both radiation and matter dominated epochs
- More PBHs are produced in MD
- We may detect gravitational wave signals secondarily-induced by large SCALAR fluctuations at small scales by e.g., [aLIGO](#), [KAGRA](#), [DECIGO](#) ...
- We will be able to distinguish a model from others by using future small-scale probes such as [PIXIE-like satellite](#) (CMB μ -distortion), [SKA/Ominiscope](#) (21cm, Pulsar timing), [CTA](#) (gamma-ray), [DECIGO](#) (GW)...