
EeV scale dark matter: production mechanism and experimental signatures

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1901.XXXX, 1902.XXXX

Lucien HEURTIER



IBS, Daejeon, January 2019

The dark sector is still quite obscure...

What is Dark Matter?

(One unique particle? A set of particles? Primordial black holes?...)

How is it produced?

(Thermal Freeze Out, Non-thermal/Freeze-In, dynamical dark matter , out-of-equilibrium decay ...)

What is its mass?

(From micro-eV to EeV and further...)

The WIMP and beyond

- The WIMP miracle

$$m_{DM} \sim m_{EW} \quad \text{scale and} \quad \langle \sigma v \rangle \sim m_{EW}^{-2}$$

should have been detected already...

(According to [\[Slatyer, Beacom et al., 18'\]](#) the WIMP isn't slightly ruled out though)

- Unitarity bound on thermal scenarios $\rightarrow m_{DM} < 500 \text{ TeV}$
- Non-WIMP scenarios allow a wider range of masses
- How to evade the WIMP paradigm?
 - \longrightarrow Progressive thermal production ?
 - \longrightarrow Out-of-equilibrium decay ? (Reheating?)
 - \longrightarrow Hidden sector containing dark matter?

Can Dark Matter be (very) Heavy?

Two examples:

- WIMPzilla : Dark matter produced at the reheating time, in the right proportion.

→ Either under the form of an instantaneous decay at $T = T_{RH}$;

→ Or by an extended production during the reheating period $T_{RH} \lesssim T \lesssim T_{max}$.;

e.g., in [Dudas, Mambrini, Olive, 17'] An EeV gravitino is proposed as a DM candidate produced mainly during the reheating.

- The diluton [Hooper, 16'] is another option to generate thermally a DM relic density for DM particles as heavy as $m_{DM} \gtrsim PeV$

Can Dark Matter be (very) Heavy?

Two examples:

- WIMPzilla : Dark matter produced at the reheating temperature T_{RH} with the right proportion.

→ Either under the form of a thermal production as decay at $T = T_{RH}$;

→ Or by an external source during the reheating period $T_{RH} \lesssim T \lesssim T_m$.

e.g., in [Dine, 16, 17] An EeV gravitino is proposed as a DM candidate produced mainly during the reheating.

- The photon [Hooper, 16'] is another option to generate thermally a DM relic density for DM particles as heavy as $m_{DM} \gtrsim PeV$

If dark matter is heavy : The role played by the reheating might be of a crucial importance

Outline

Part I

- Motivations: Connection between inflation and dark matter phenomenology
- The inflaton portal to dark matter: General Idea
- Study of a toy model: The inflaton miracle
- Cosmological Constraints

Part II

- UHECR detection with neutrino detectors
- ANITA anomalous events
- A (very) heavy dark matter explanation
- Complementarity with IceCube detection

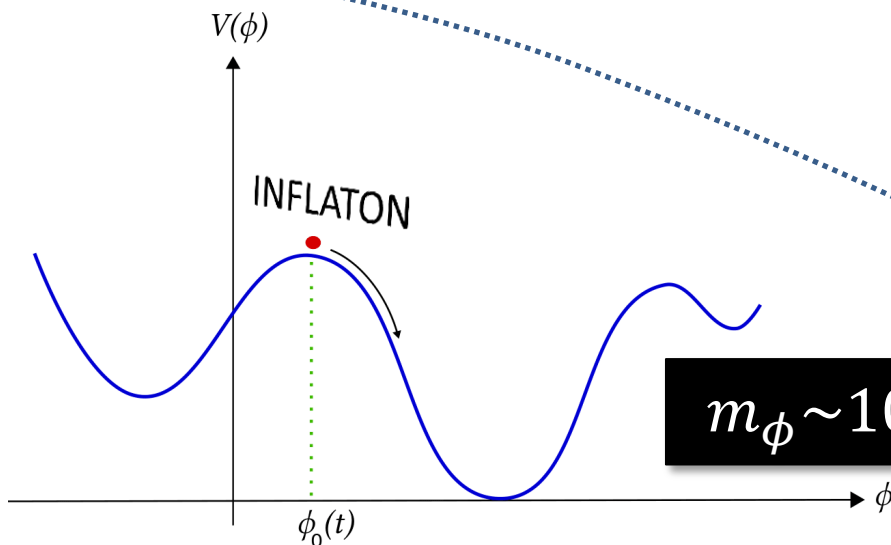
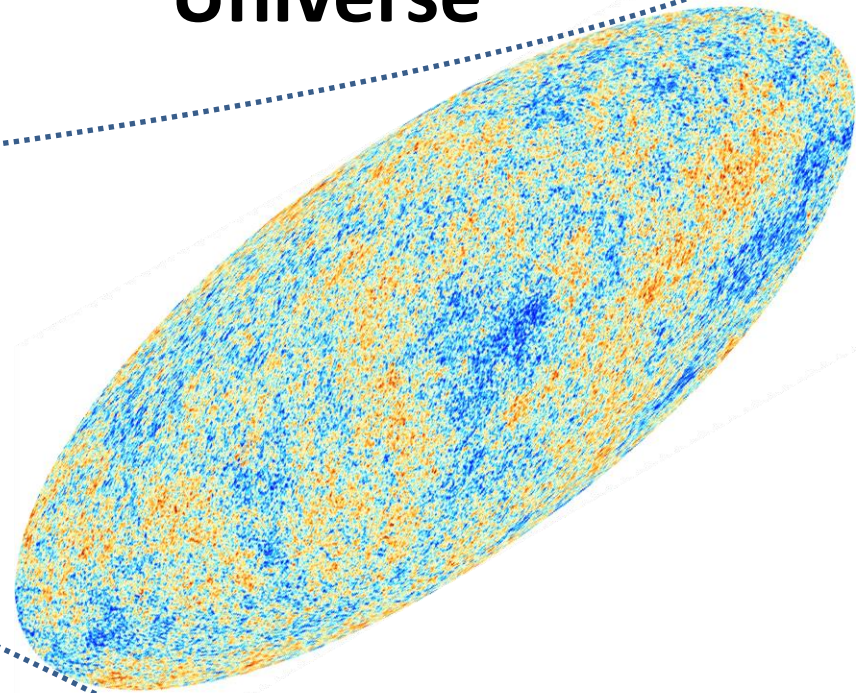
The importance of the inflaton couplings

Primordial
Universe

Homogeneous, flat
Universe

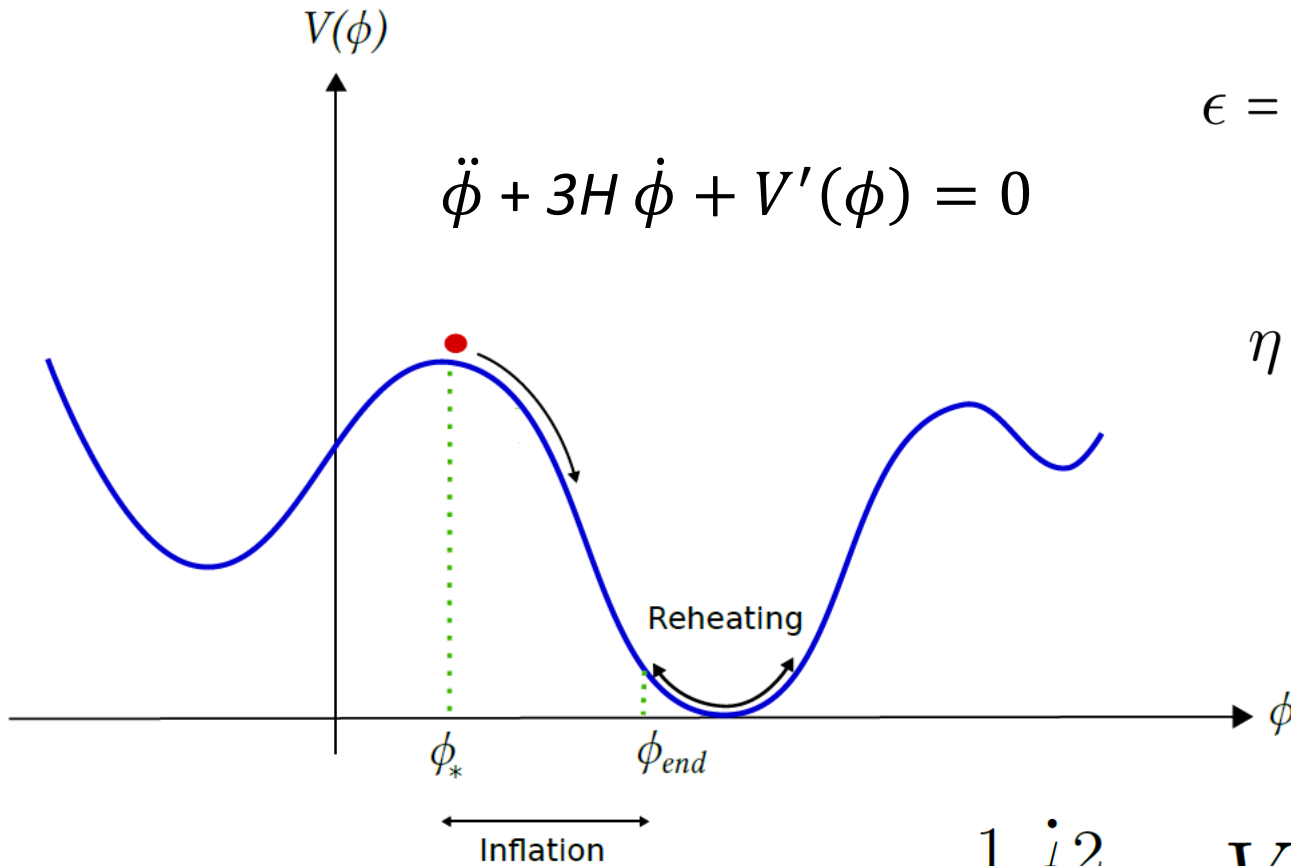


INFLATION



$$m_{\phi} \sim 10^{13} \text{ GeV (chaotic inflation)}$$

Single Field Inflation



$$\ddot{\phi} + 3H \dot{\phi} + V'(\phi) = 0$$

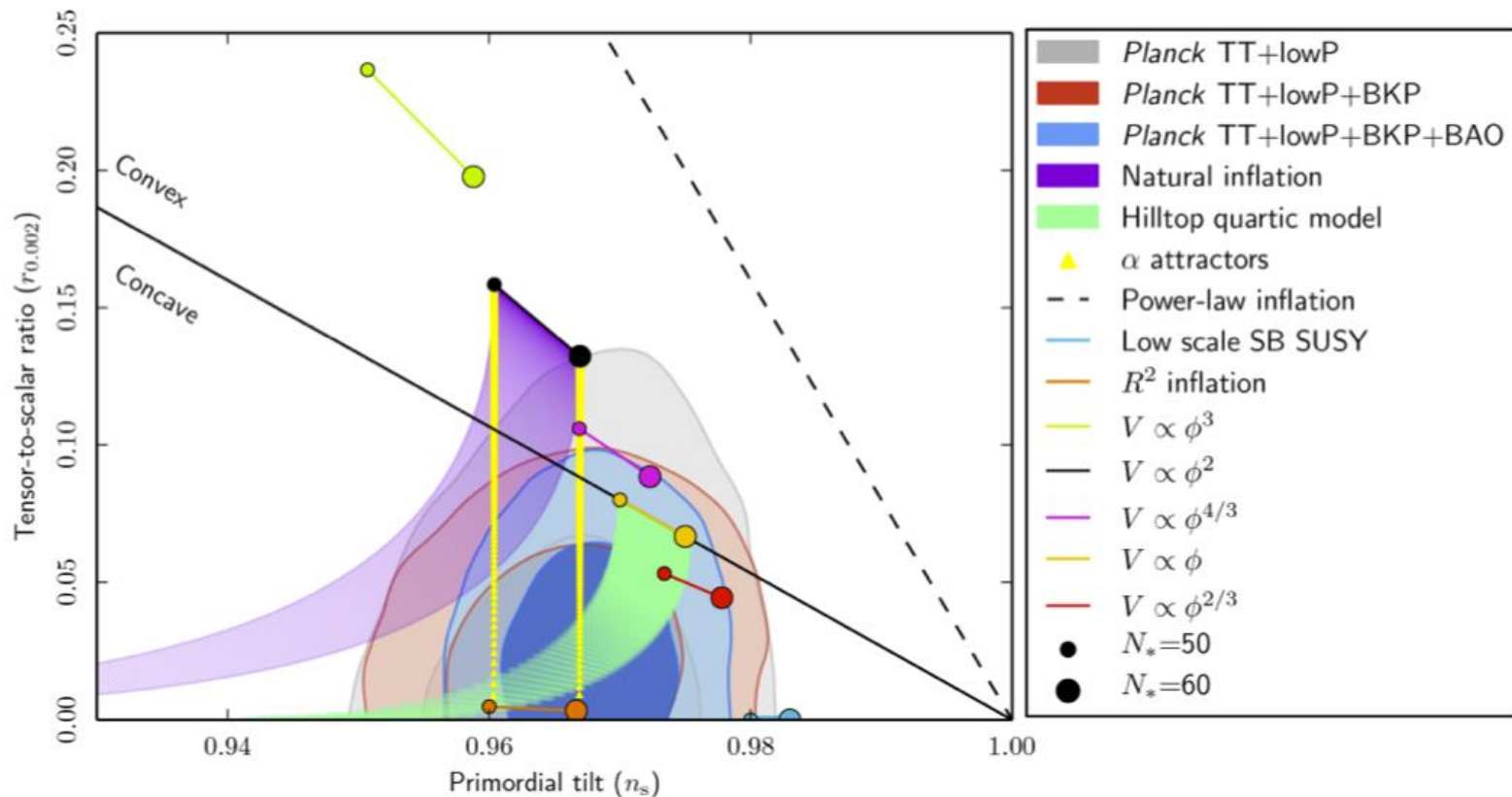
$$\epsilon = \frac{1}{2} \left(\frac{V'}{V} \right)^2 \ll 1$$

$$\eta = \left| \frac{V''}{V} \right| \ll 1$$

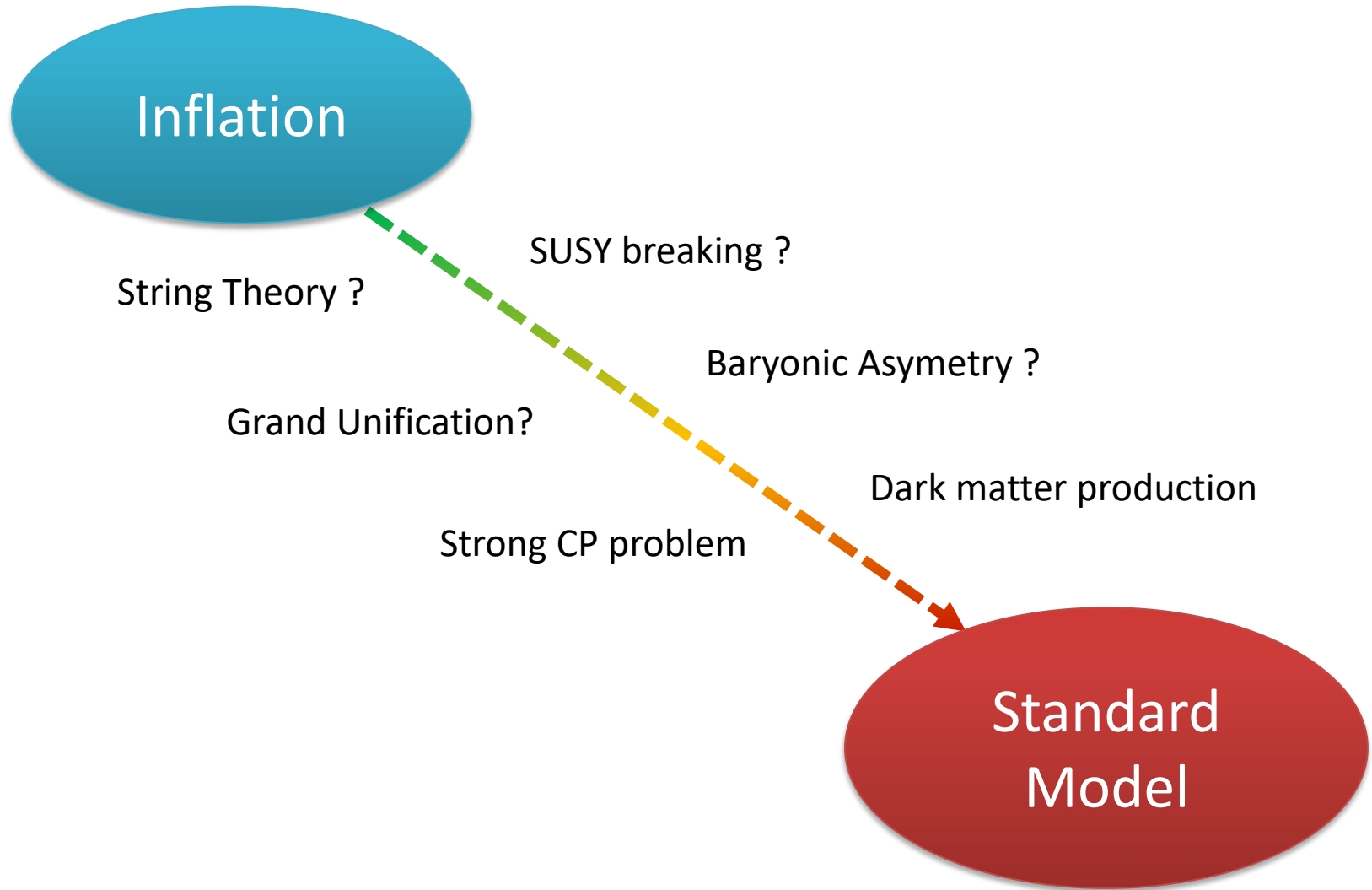
$$\omega_{\phi} = \frac{\frac{1}{2} \dot{\phi}^2 - V(\phi)}{\frac{1}{2} \dot{\phi}^2 + V(\phi)} \sim -1$$

Inflation constraints

- Tensor to scalar ratio : $r = 16 \varepsilon$
- Spectral index : $n_s = 1 - 6 \varepsilon + 2 \eta$



From Inflation to Low energy physics



From Inflation to Low energy physics

SUSY breaking **at the end of inflation**
perturbs the inflaton trajectory

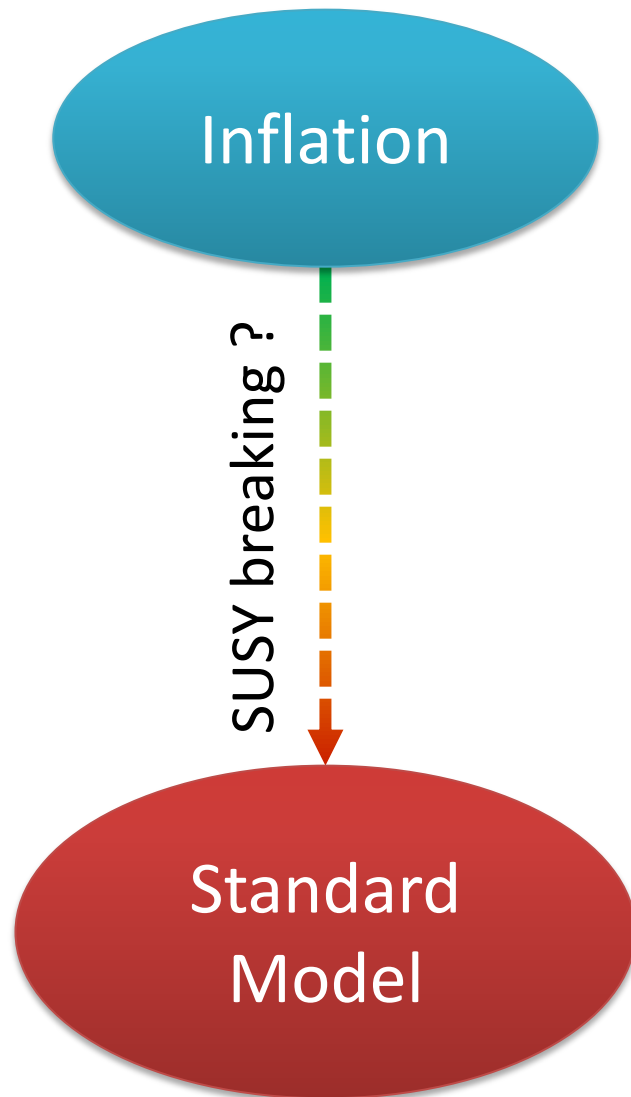
→ $f_{SUSY} > H$ can destabilize chaotic inflation for models with stabilizers ($W = mS\phi$)

[Buchmüller, Dudas, LH, Wieck, '14]

→ $f_{SUSY} \sim H$ can rescue models without stabilizers ($W = m\phi^2$)

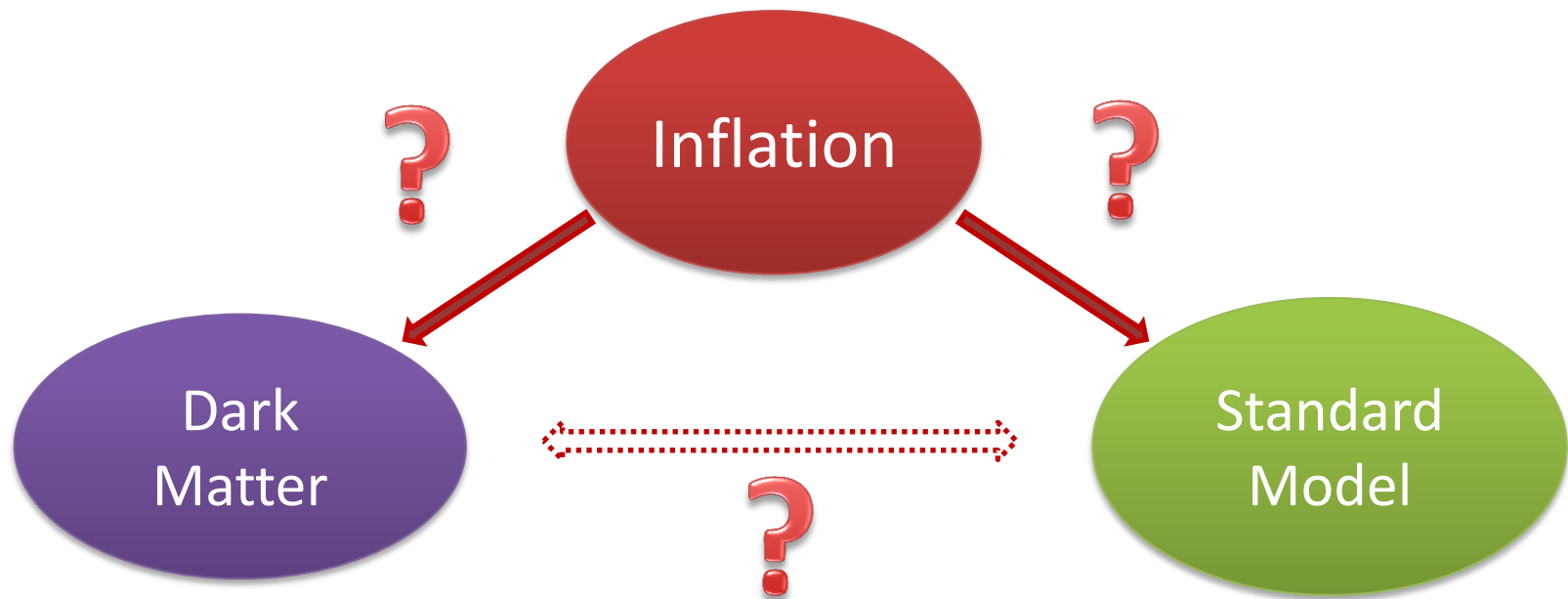
[Buchmüller, Dudas, LH, Westphal, Wieck, Winkler '15]

→ In a complete model including low energy SUSY breaking and inflation: constraints from many different perspectives... [Argurio, Dries, LH, Mariotti '16]



The importance of the inflaton couplings

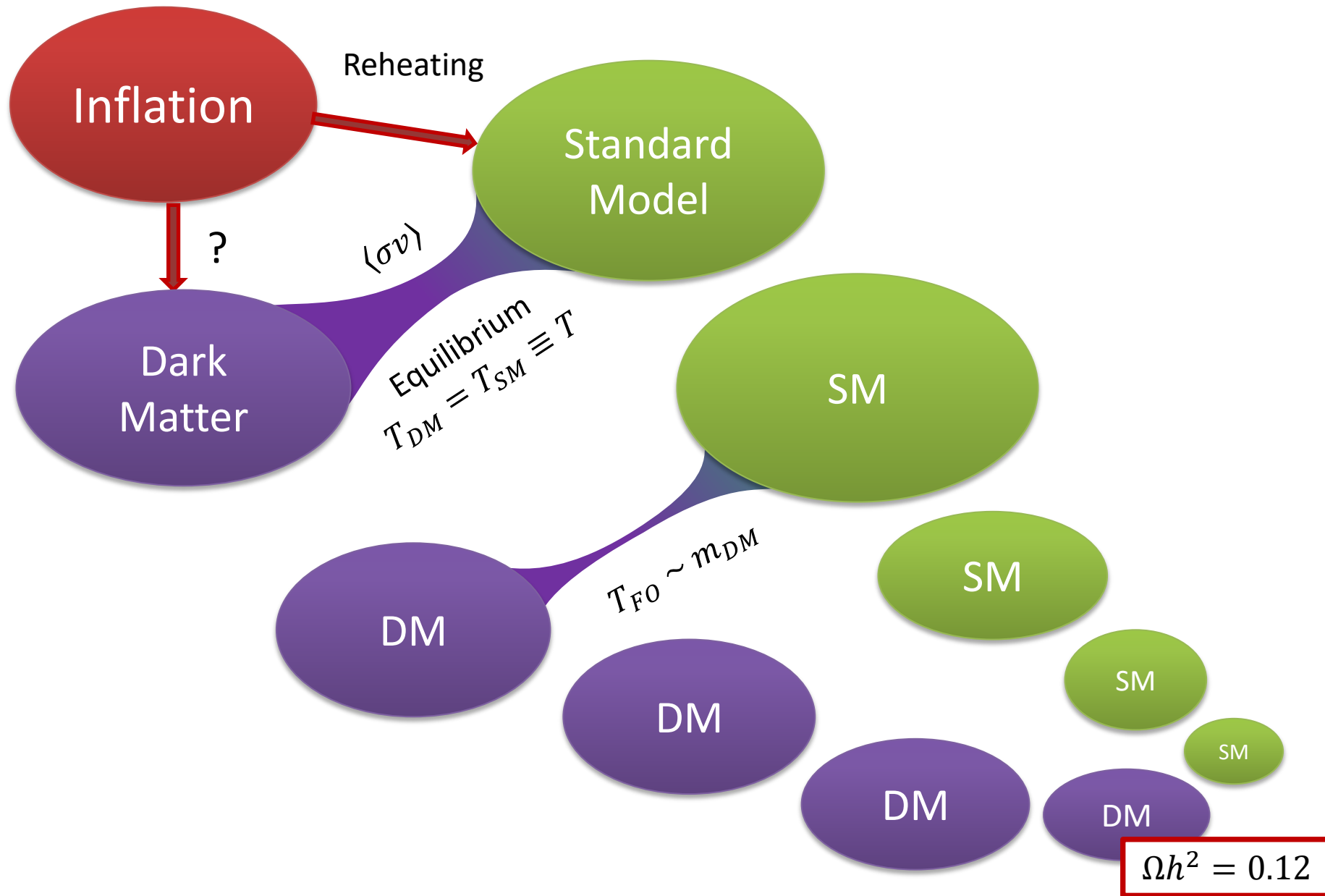
Inflaton Decay : Reheating



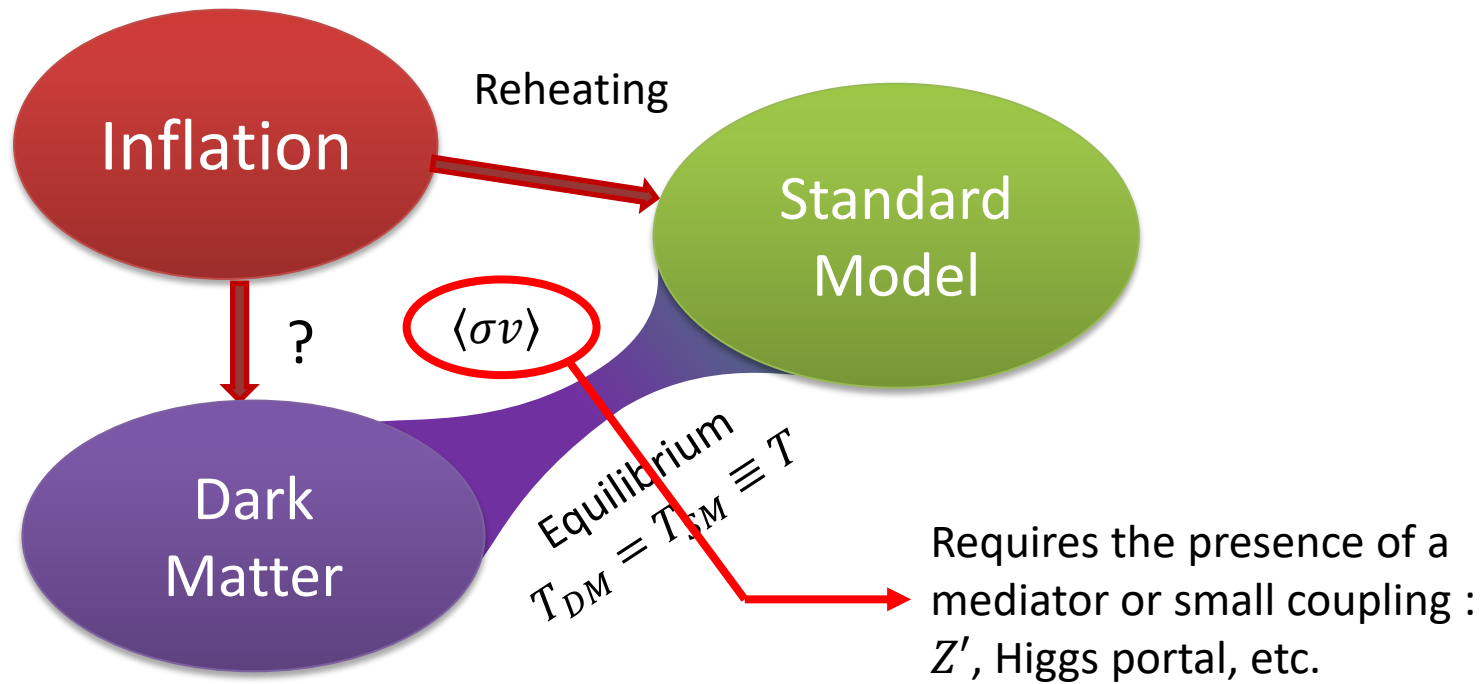
Being explicit about the reheating lagrangian fixes initial conditions for dark matter production ...

When is it relevant to DM production ?

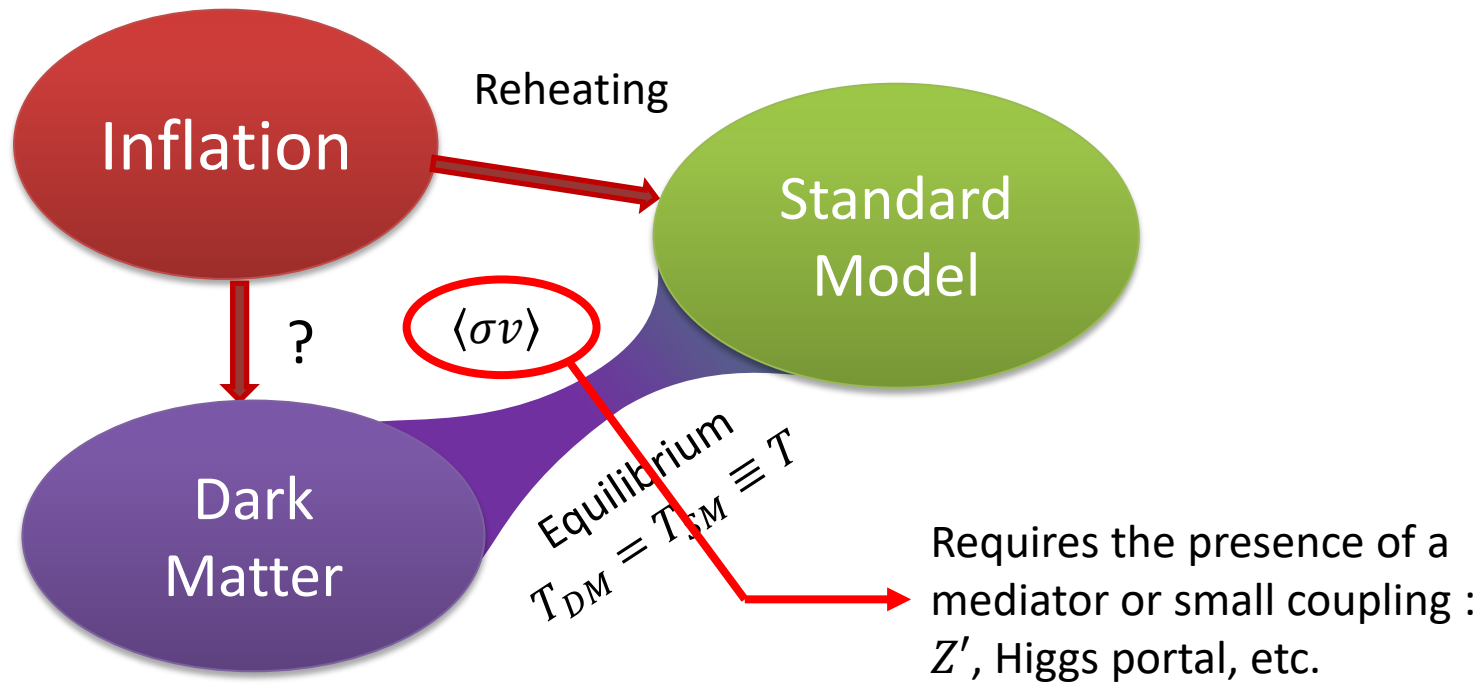
Thermal scenario of Dark matter production :



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Thermal scenario of Dark matter production :



**More and more disfavored
by direct detection
experiments...**



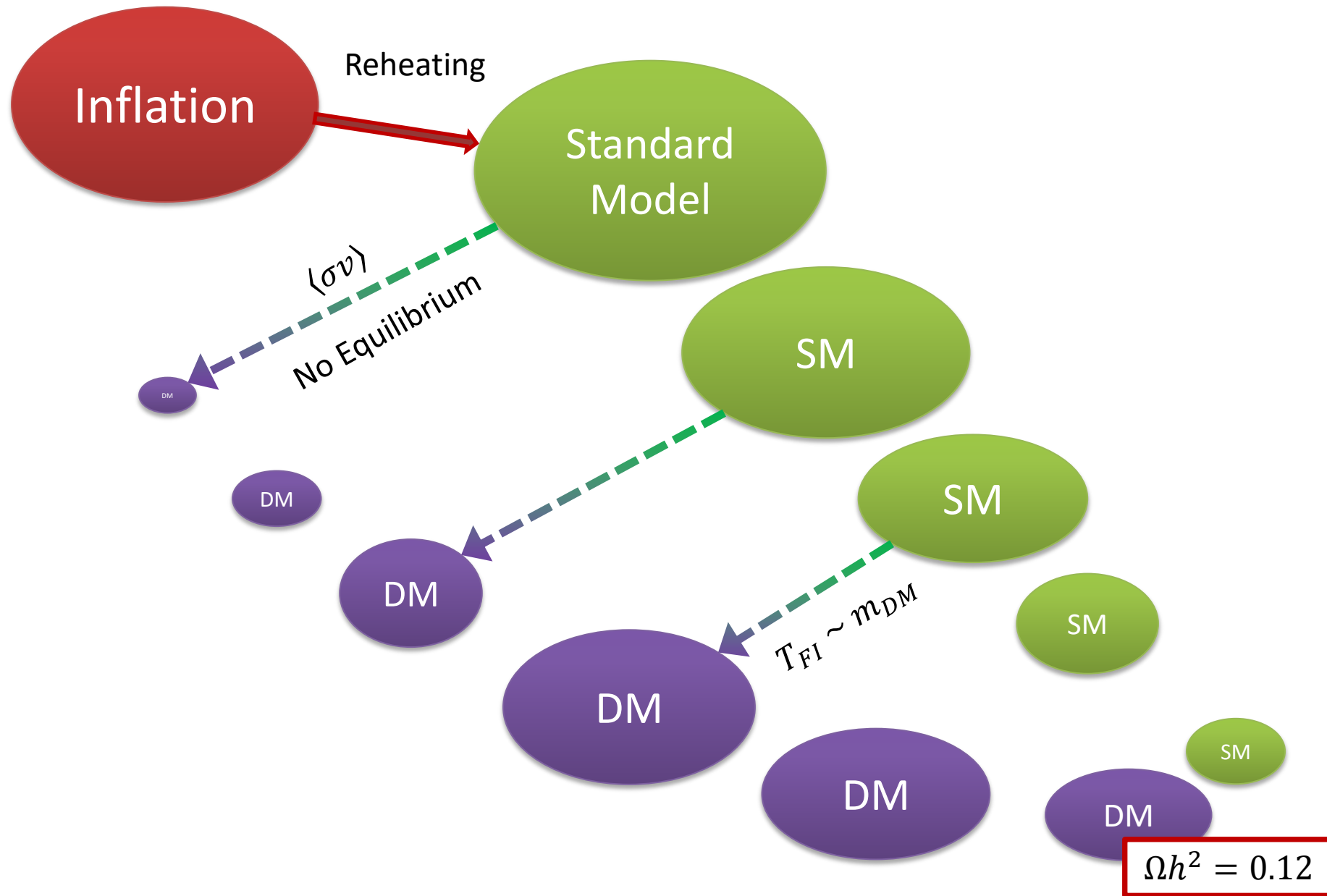
WIMP miracle:

$$\langle\sigma v\rangle \sim \langle\sigma v\rangle_{EW} \quad \text{and} \quad m_{DM} \sim \mathcal{O}(100) \text{ GeV}$$

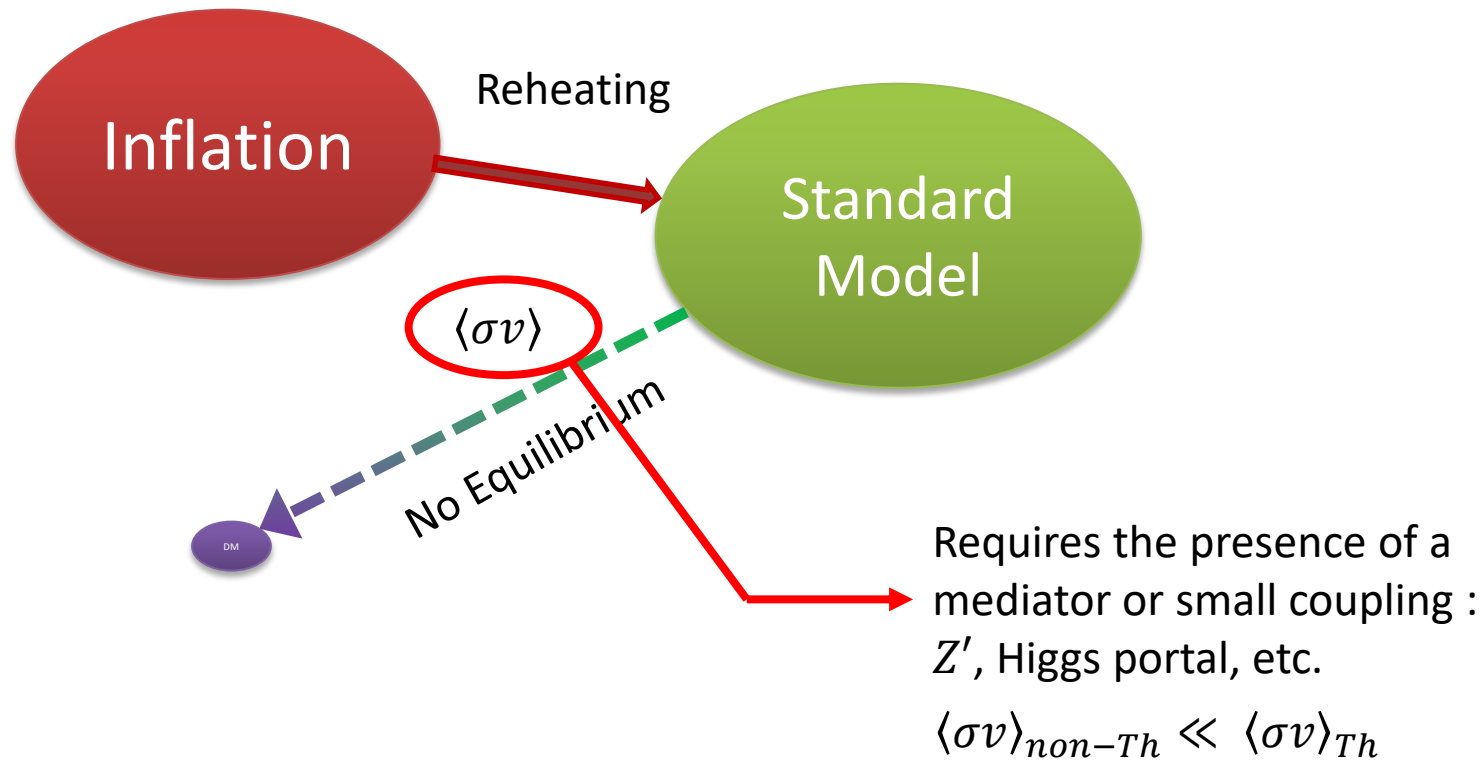


$$\Omega h^2 \sim 0.12$$

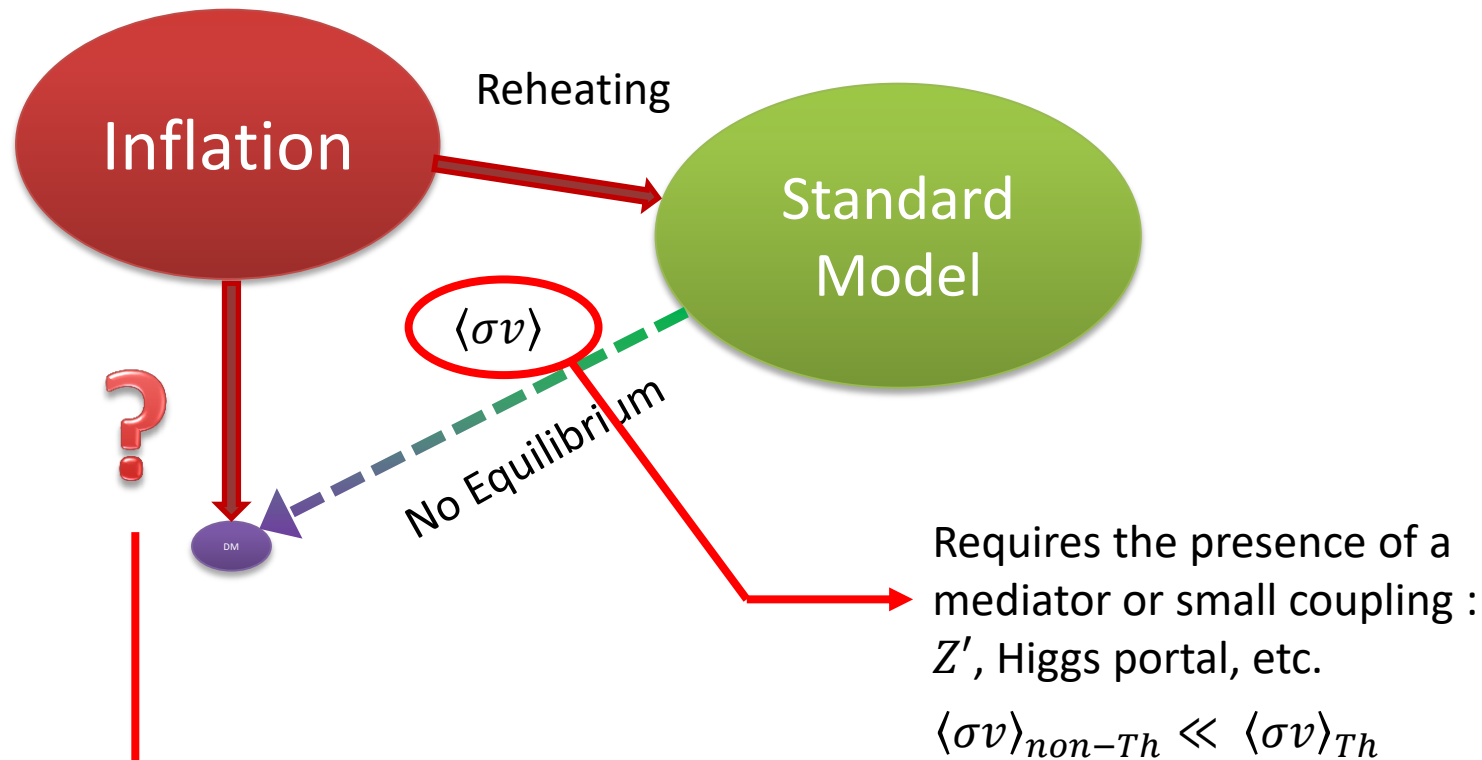
Non-Thermal scenario of Dark matter production :



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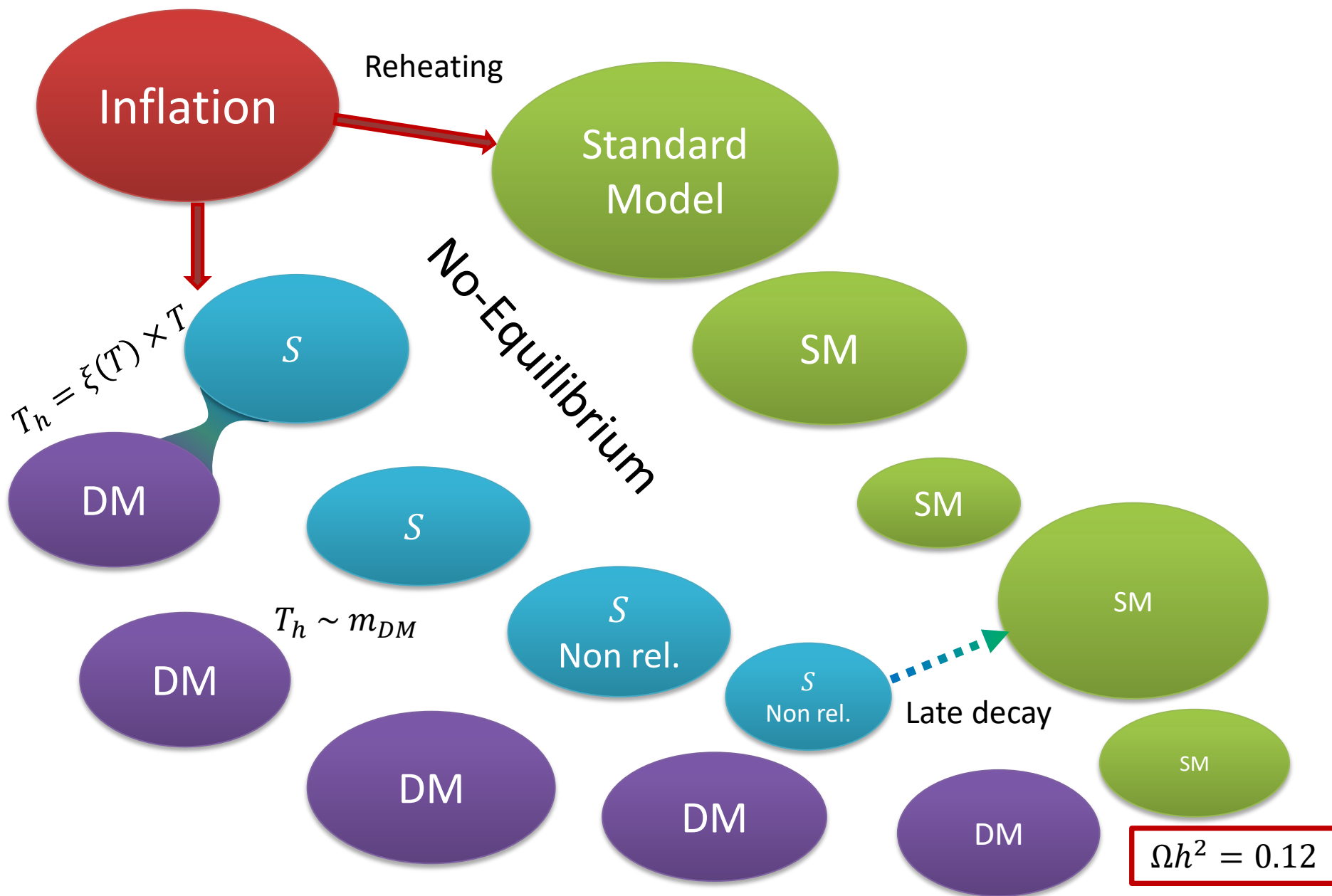
Non-Thermal scenario of Dark matter production :



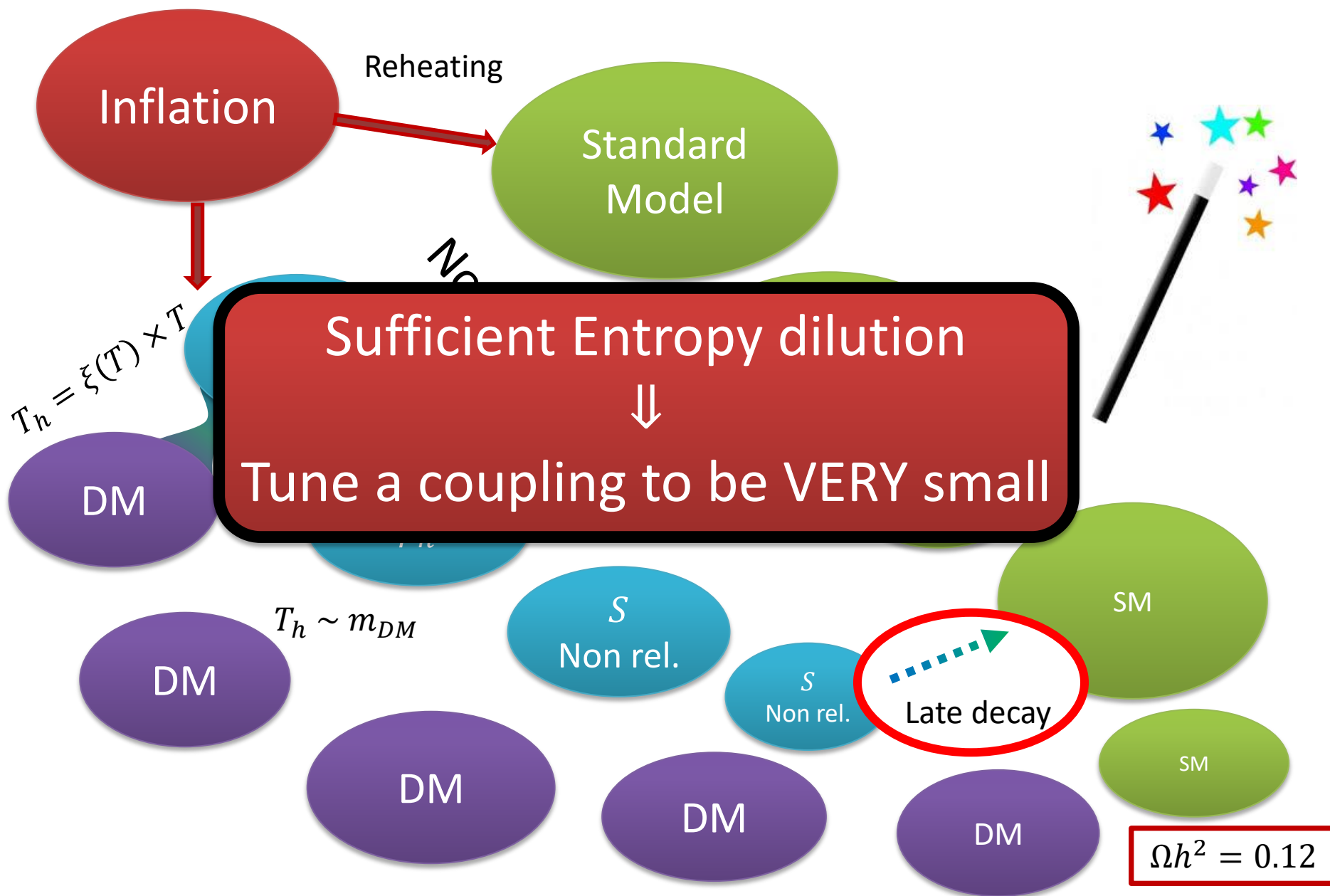
No reason a priori to suppress the production of DM through inflaton decay...

Such coupling HAS to be there at the loop level in many models of reheating ...
[Kaneta, Mambrini, Olive '19]

Decoupled Hidden sector [Hooper et al., '16]



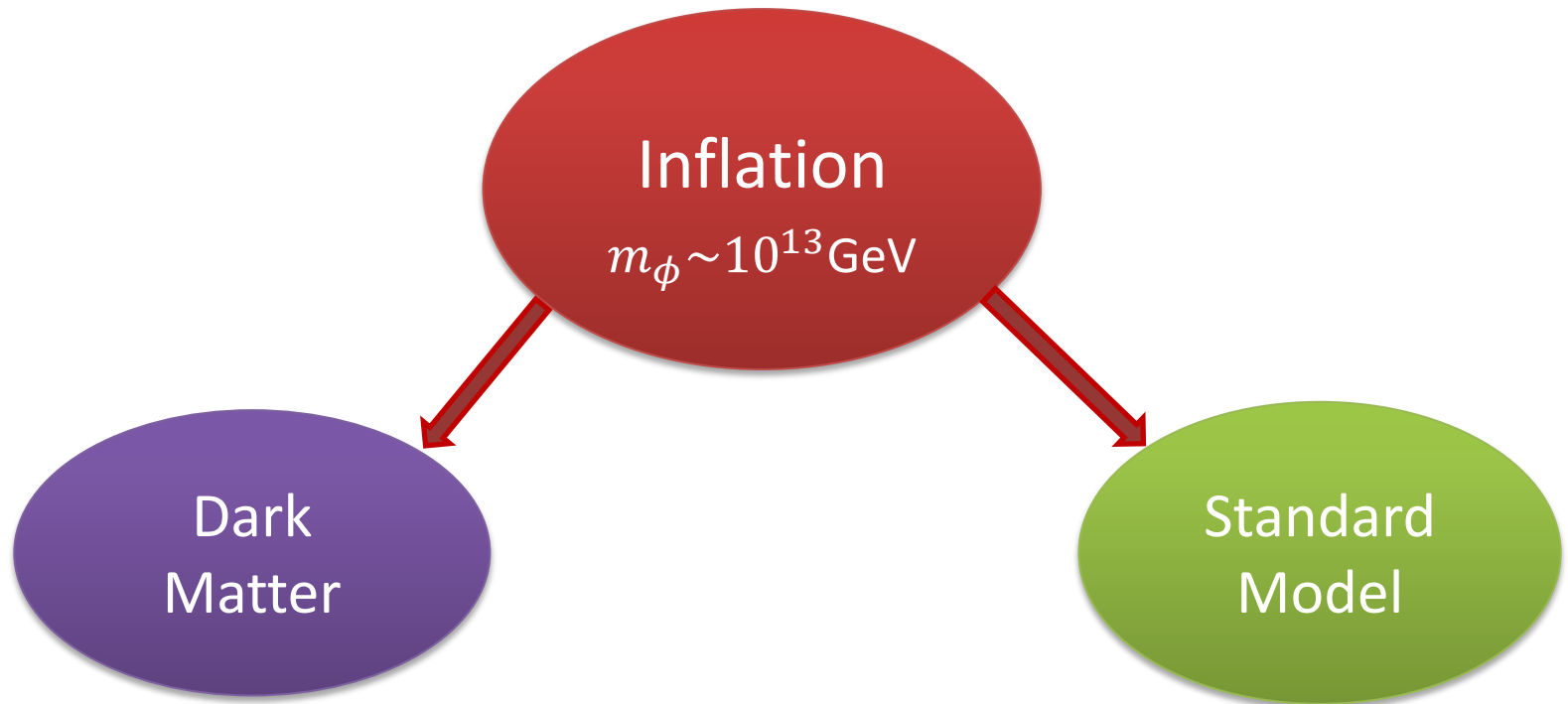
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Part I

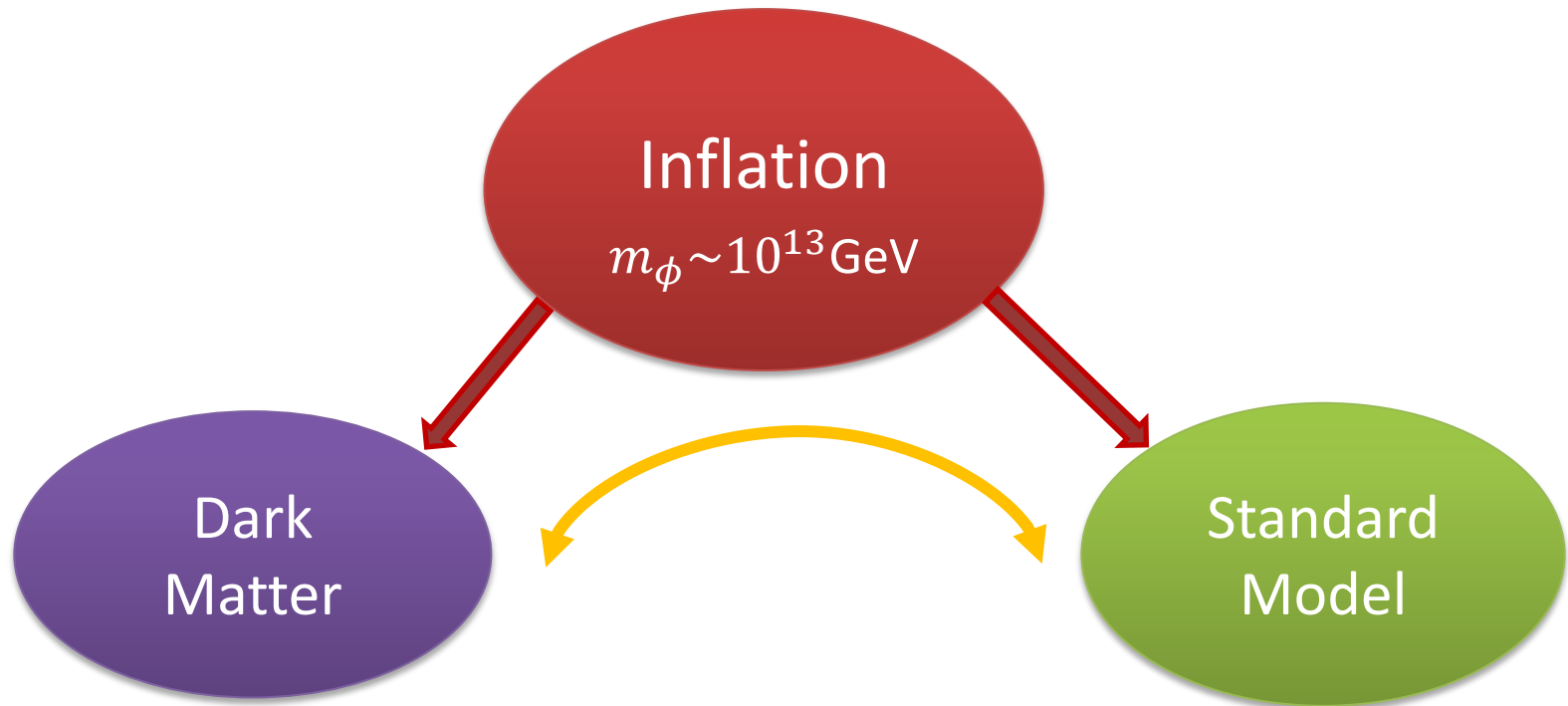
The Inflaton Portal to Dark Matter

The inflaton portal to DM



[Dev, Mazumdar, Qutub 13'], [Heurtier 17']

The inflaton portal to DM

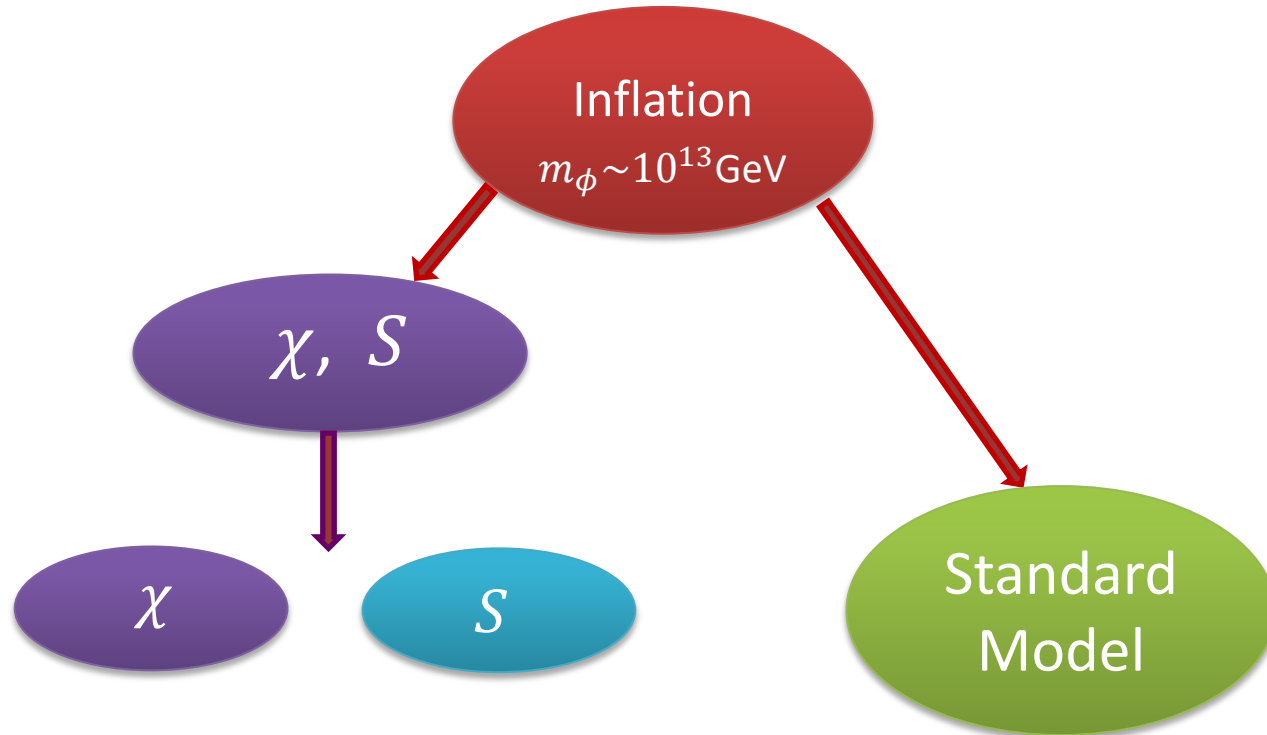


$m_\phi \sim 10^{13} \text{ GeV}$ → Annihilation cross section feeble
→ No possible thermal scenario

[Dev, Mazumdar, Qutub 13'], [Heurtier 17']

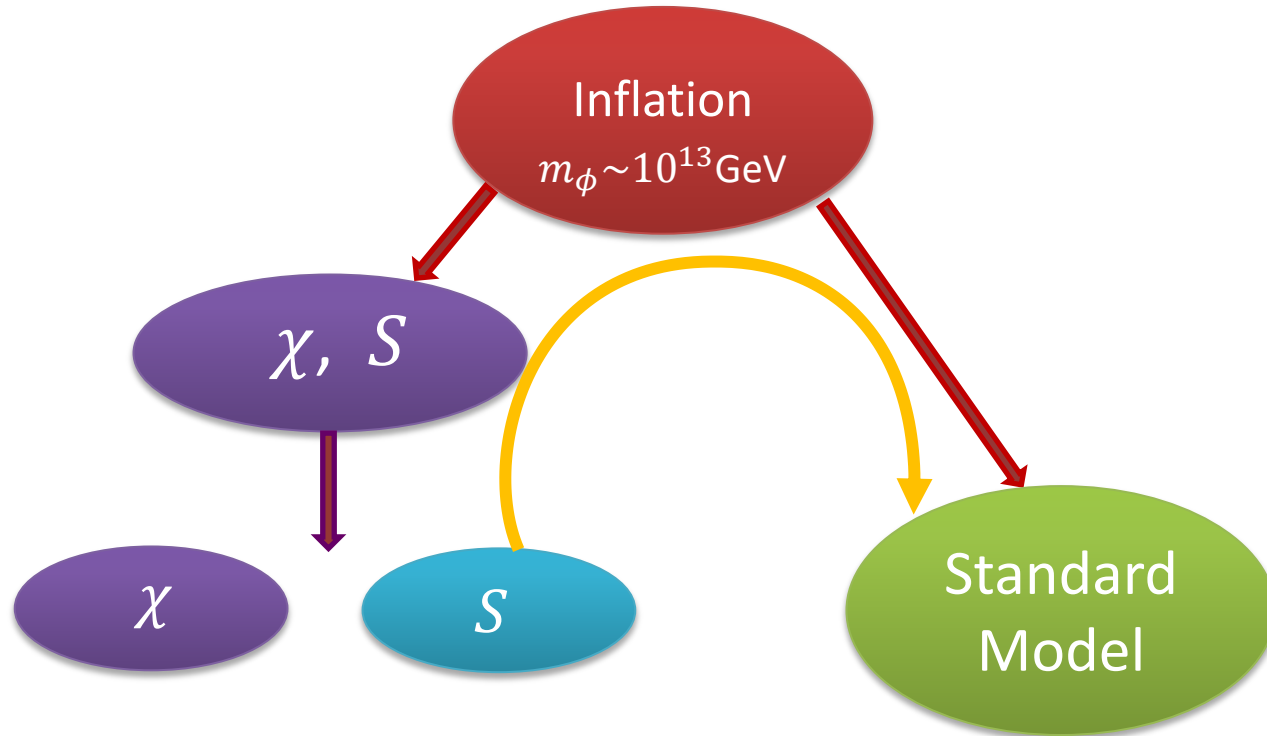
The inflaton portal to DM

Highly decoupled sectors?



The inflaton portal to DM

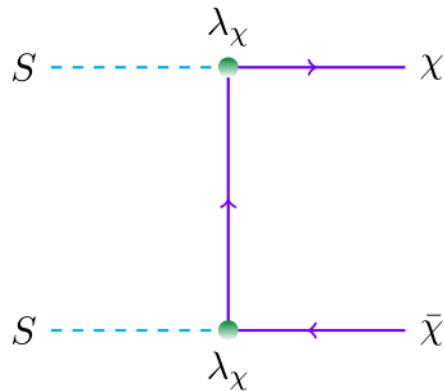
Highly decoupled sectors?



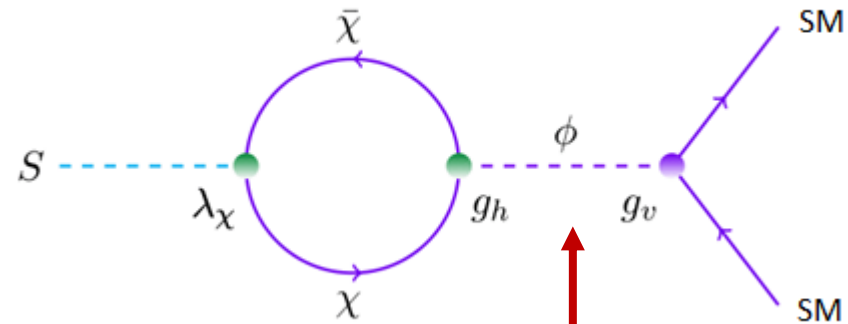
$m_\phi \sim 10^{13} \text{ GeV} \longrightarrow$ Late decay of the hidden sector

The Model

[F.Huang, L.H., coming soon]



Thermal decoupling of dark matter in the dark sector

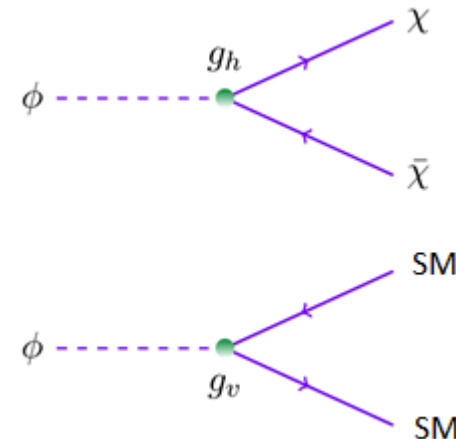


$$m_\phi = 10^{13} \text{ GeV}$$

Natural suppression of the hidden scalar decay width...

$$g_h/g_v \longrightarrow T_h/T_v \text{ after inflation}$$

$$\xi_{\text{inf}} \equiv \left(\frac{T_h}{T} \right)_{\text{inf}} = \left(\frac{g_{\text{inf}}^*}{g_{h,\text{inf}}^*} \right)^{1/4} \times \left(\frac{\rho_h}{\rho_v} \right)^{1/4}$$



Relic Density

[F.Huang, L.H., coming soon]

$$\Omega_\chi h^2 \longrightarrow \Omega_\chi h^2 \times \frac{S_i}{S_f}$$

Thermal decoupling
in the dark sector

Entropy Suppression

$$\Omega_\chi h^2 \approx 8.5 \times 10^{-11} \frac{m_\chi \sqrt{g_\star + g_\star^h \xi^4}}{T_f^h g_\star} \left(\frac{a + 3\xi b m_\chi / T_f^h}{\text{GeV}^{-2}} \right)^{-1}$$

$$\frac{S_f}{S_i} \approx 1.83 \langle g_\star^{1/3} \rangle^{3/4} \frac{m_S Y_S \tau_S^{1/2}}{m_p^{1/2}}$$

$$a \sim b \sim \alpha_\chi^2 / m_\chi^2$$

Inflaton suppressed decay rate

$$\Gamma_S \sim (g_h g_v)^2 \frac{m_S^5}{m_\phi^4}$$

Relic Density

[F.Huang, L.H., coming soon]

$$\Omega_\chi h^2 \longrightarrow \Omega_\chi h^2 \times \frac{S_i}{S_f}$$

Generic range of masses

$$\frac{\Omega_\chi h^2}{S_f/S_i} \approx 0.38 \left(\frac{10^{13} \text{GeV}}{m_\phi} \right)^2 \left(\frac{g_h g_v}{(0.1)^2} \right) \left(\frac{m_\chi}{10 \text{PeV}} \right)^{7/2} \left(\frac{m_\chi/m_S}{10} \right)^{-3/2} \left(\frac{0.1}{\alpha_\chi} \right)^2$$

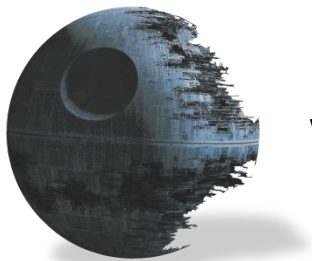
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WIMP miracle



**The Inflaton
miracle**



Relic Density

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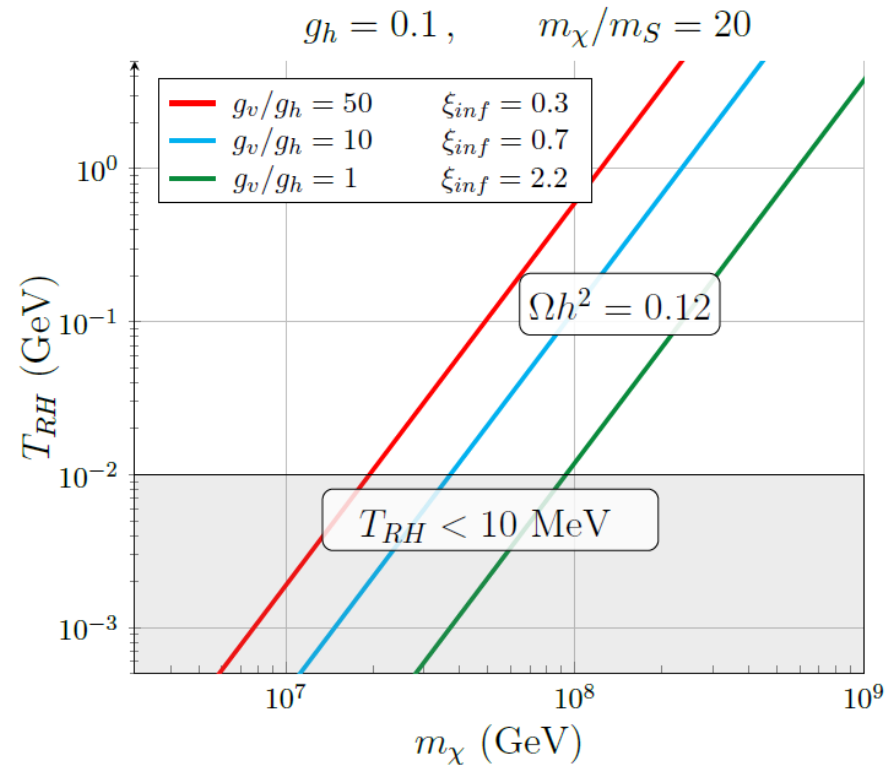
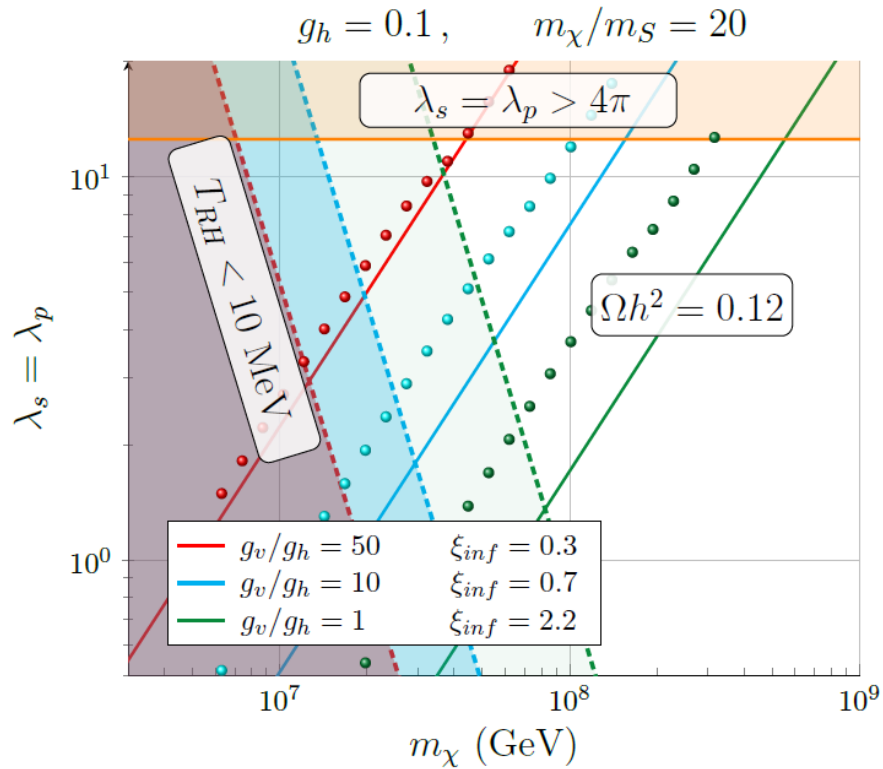
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$$T_{RH} \equiv \left(\frac{90}{8\pi^3 g_\star} \right)^{1/4} \sqrt{\frac{M_p}{\tau_s}} > 10 \text{ MeV}$$

$$\frac{\Omega_\chi h^2}{S_f/S_i} \approx 0.16 \times \xi_f^{-2} \left(\frac{T_R}{10 \text{MeV}} \right) \left(\frac{m_\chi}{50 \text{PeV}} \right) \left(\frac{m_\chi/m_S}{20} \right) \left(\frac{0.8}{\alpha_\chi} \right)^2$$

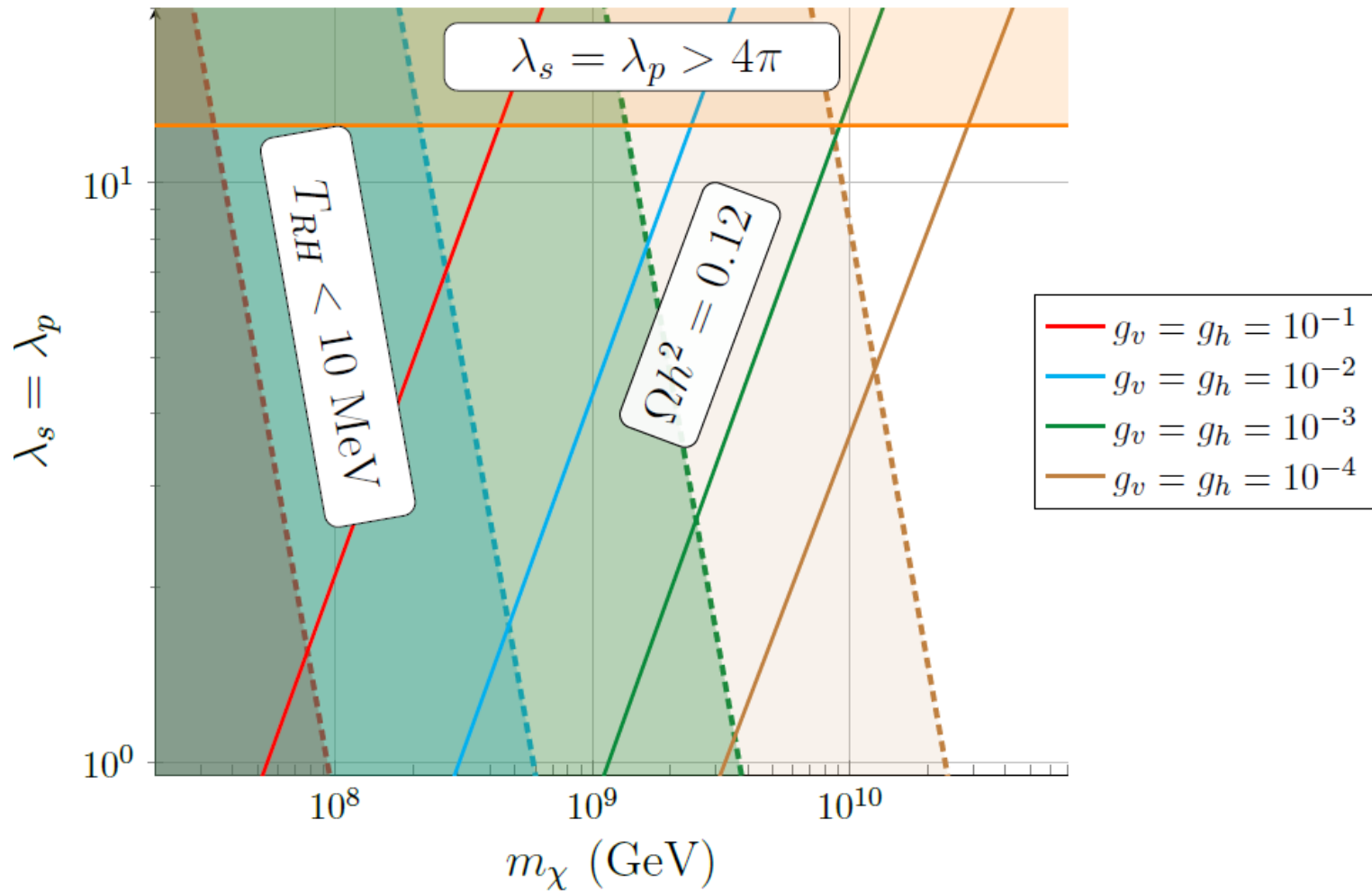
Relic Density

[F.Huang, L.H., coming soon]



Relic Density

[F.Huang, L.H., coming soon]



Experimental signatures ?

Dark matter features :

- 10 PeV – EeV dark matter
- Very feeble interaction with the standard model

 **No Direct Detection constraints**

- Significant annihilation into dark scalars
- Dark scalar lifetime $< 0.01\text{s}$

 **Indirect Detection ?** Same problem: Low DM number density, low flux of decay products, extremely hard to detect.

Vacuum mass versus inflation parameters

- In most minimal models, the inflaton mass in the vacuum is given by the inflation scenario

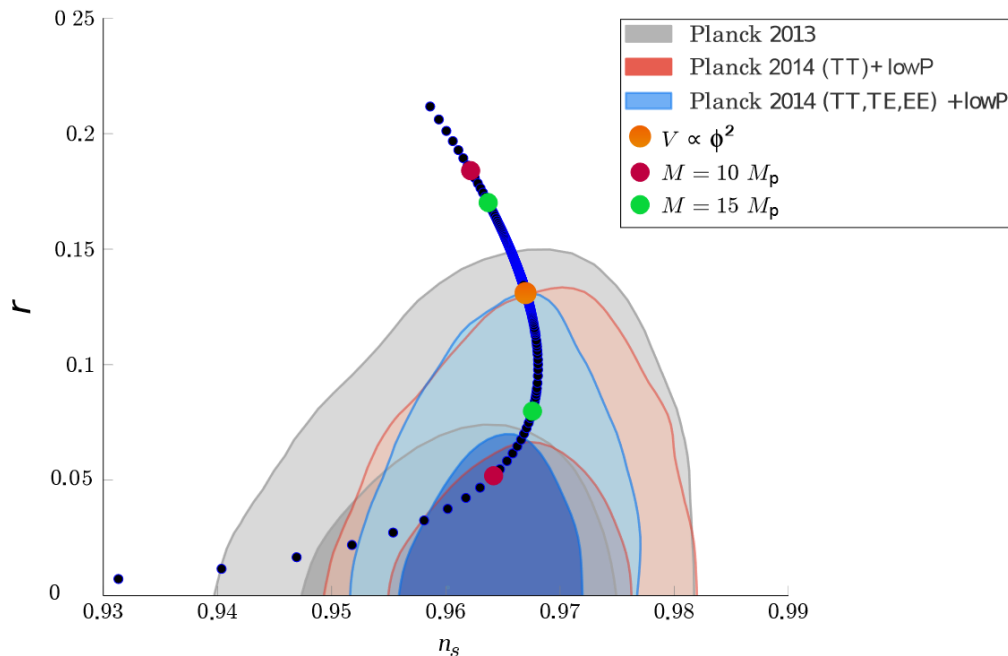
$$V(\phi) \longrightarrow m_{\phi}^2 = V''(\phi) \Big|_{\phi=0}$$

Chaotic inflation is the simplest option....

Vacuum mass versus inflation parameters

- Quadratic / quartic models

$$V(\phi) = \frac{\lambda}{4} (\phi^2 - v^2)^2$$



[Linde '07]

[LH, Moursy '15]

Vacuum mass versus inflation parameters

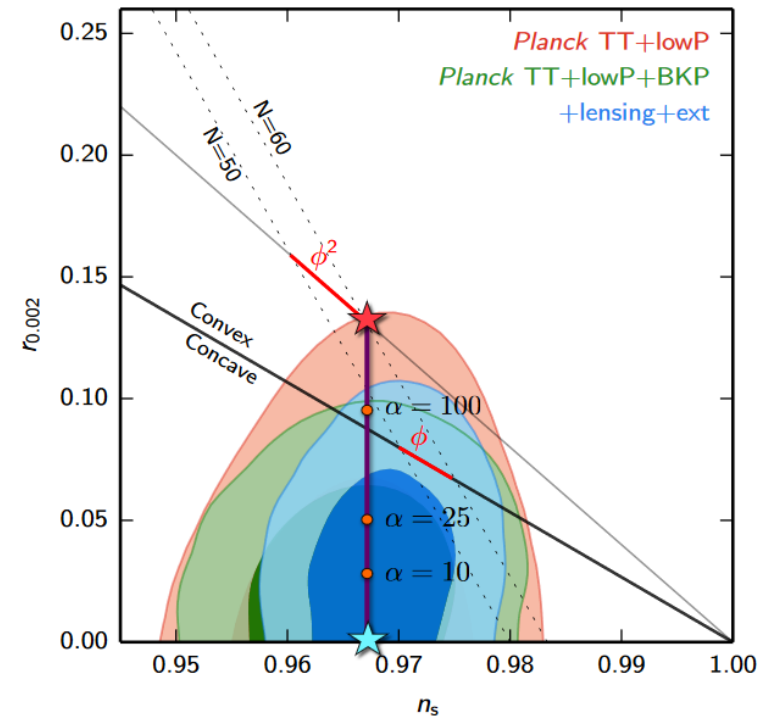
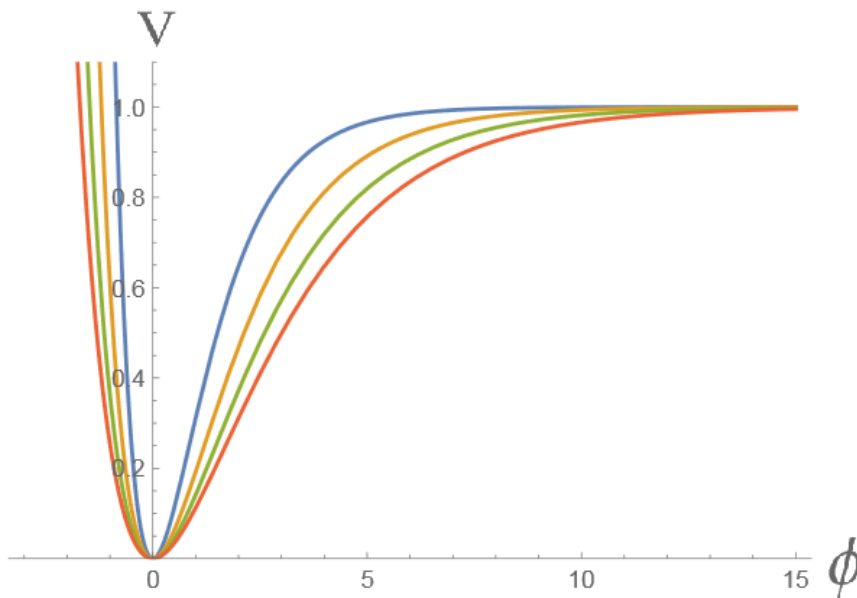
- α -attractors (E-model)

$$V(\phi) = V_0 \left(1 - e^{-\sqrt{\frac{2}{3\alpha}} \frac{\phi}{M_{Pl}}} \right)^2$$

$$m_\phi^2 = \left. \frac{\partial^2 V}{\partial \phi^2} \right|_{\phi=0} = \frac{4V_0}{3\alpha M_{Pl}^2} = \frac{m^2}{\alpha}$$

$$m \sim 3 \cdot 10^{13} \text{ GeV}$$

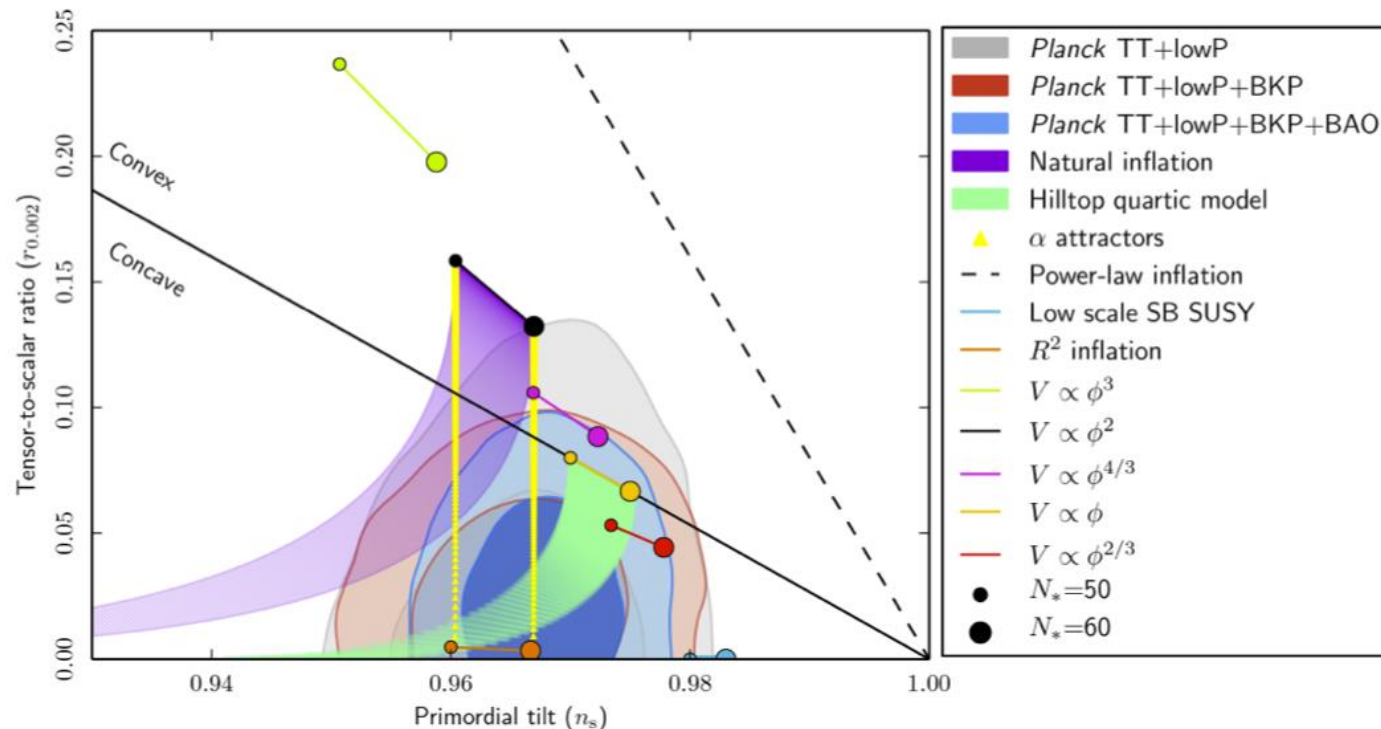
[Carrasco, Kallosh, Linde '15]



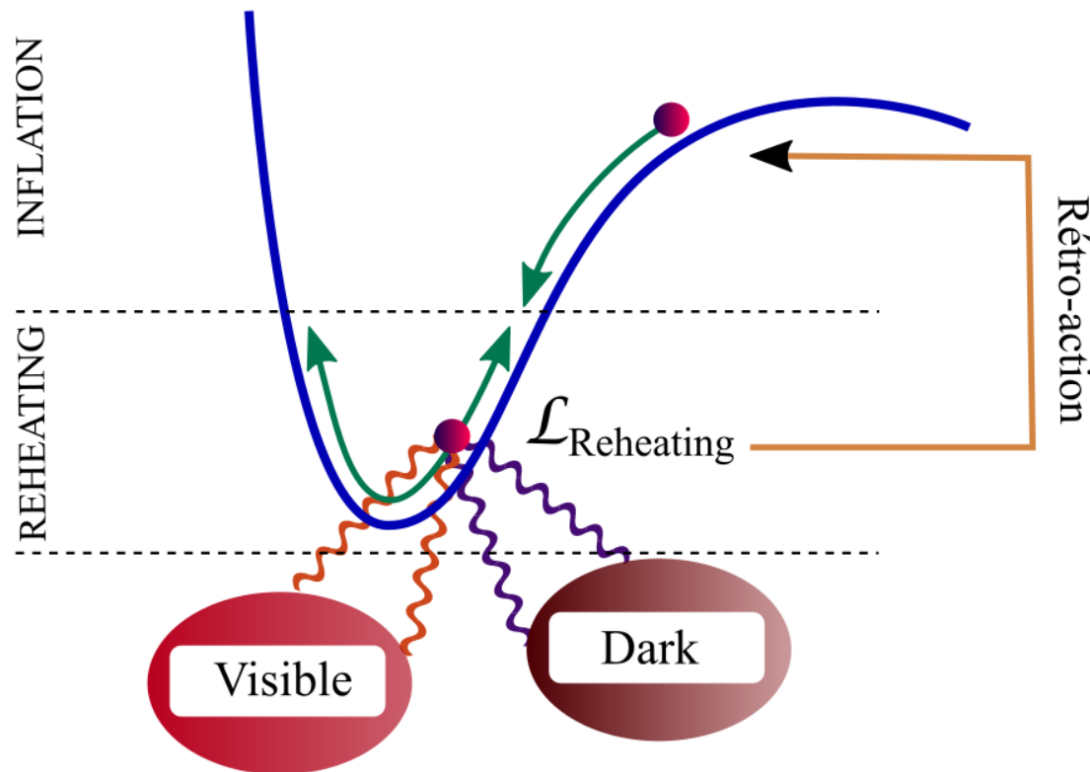
Vacuum mass versus inflation parameters

- Natural inflation

$$V(\phi) = \Lambda^4 \left[1 \pm \cos \left(a \frac{\phi}{f} \right) \right] \quad m_\phi \sim M_{GUT}^2 / M_{Pl} \sim 4 \times 10^{13} \text{ GeV}$$



What on the inflation side?



Loop corrections backreact on the inflationary trajectory

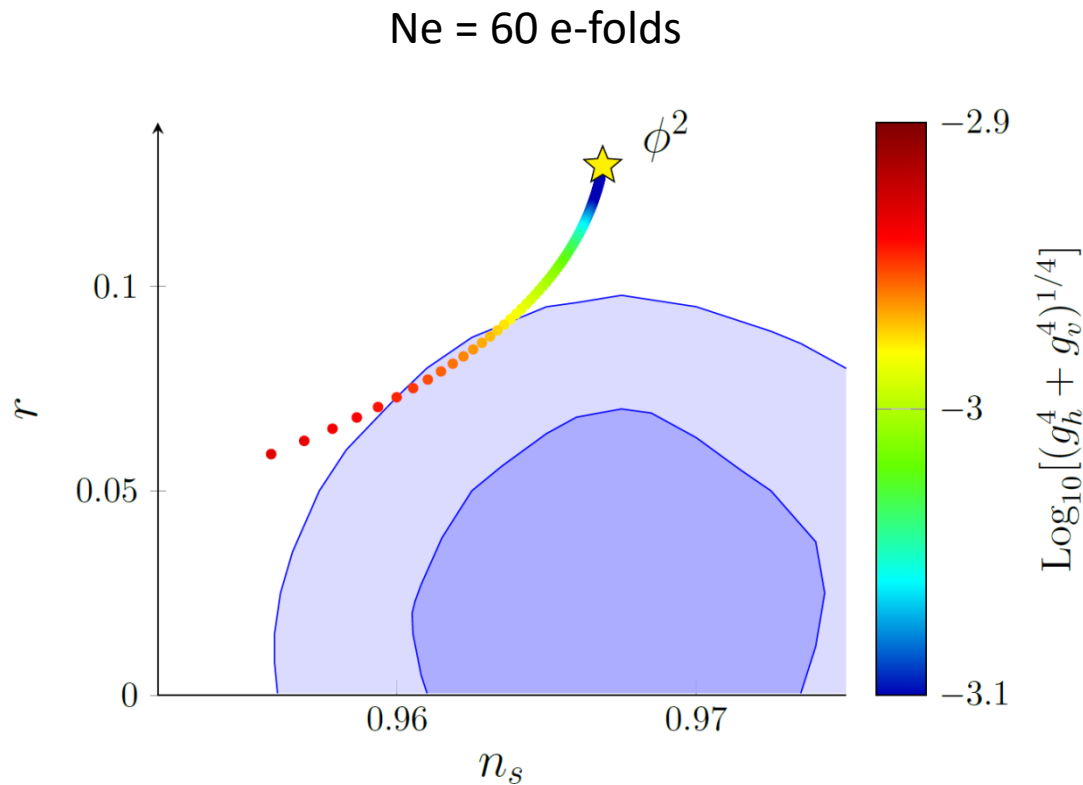


Destabilizes the inflaton at large field values



Modifies the predictions for inflation observables (r , n_s)

What on the inflation side?



Loop corrections backreact on the inflationary trajectory



Destabilizes the inflaton at large field values



Modifies the predictions for inflation observables (r , n_s)

What about the primordial universe?

After dark matter decouples, the light dark component becomes non-relativistic and dominates the energy density before decaying...

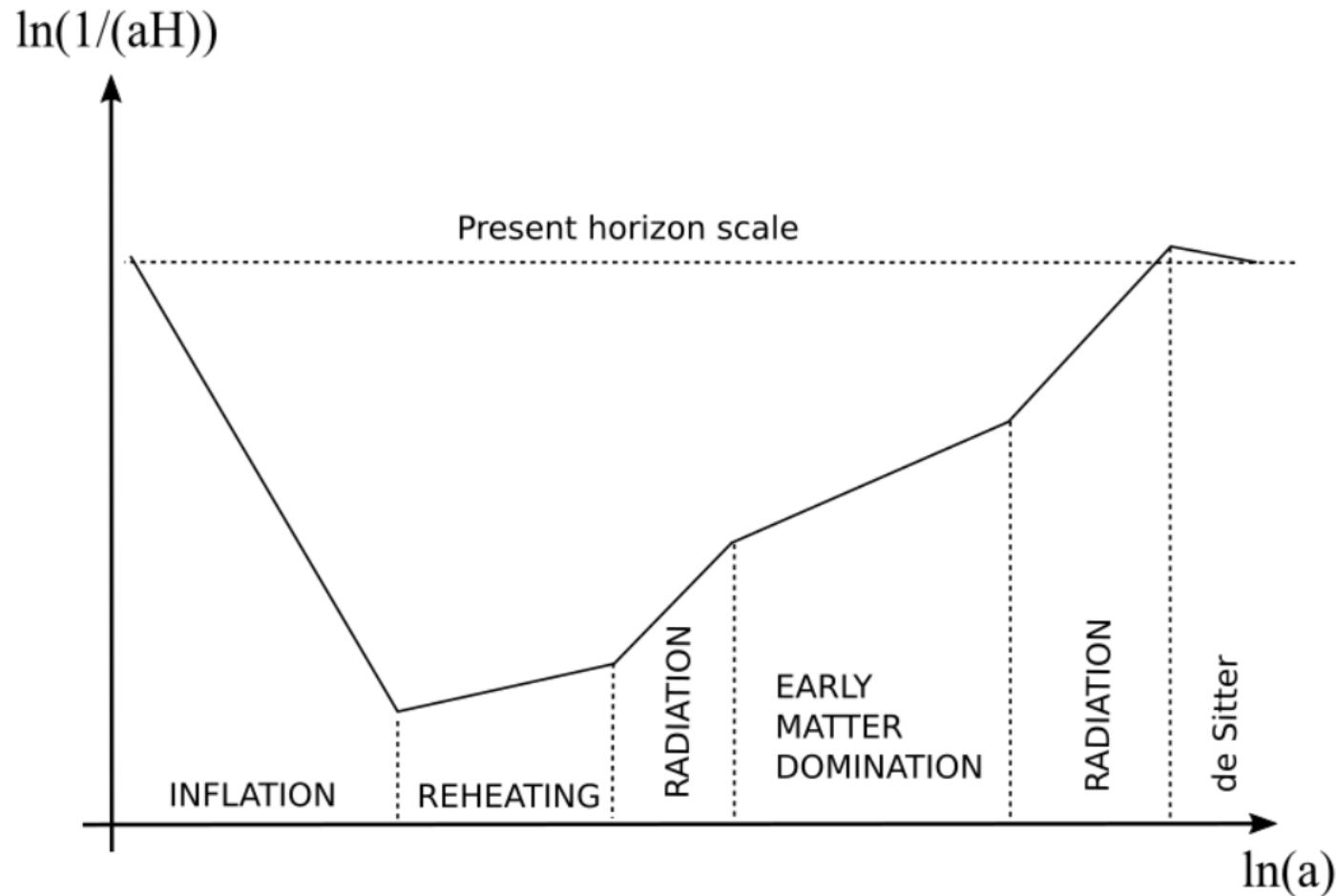


Period of early matter domination



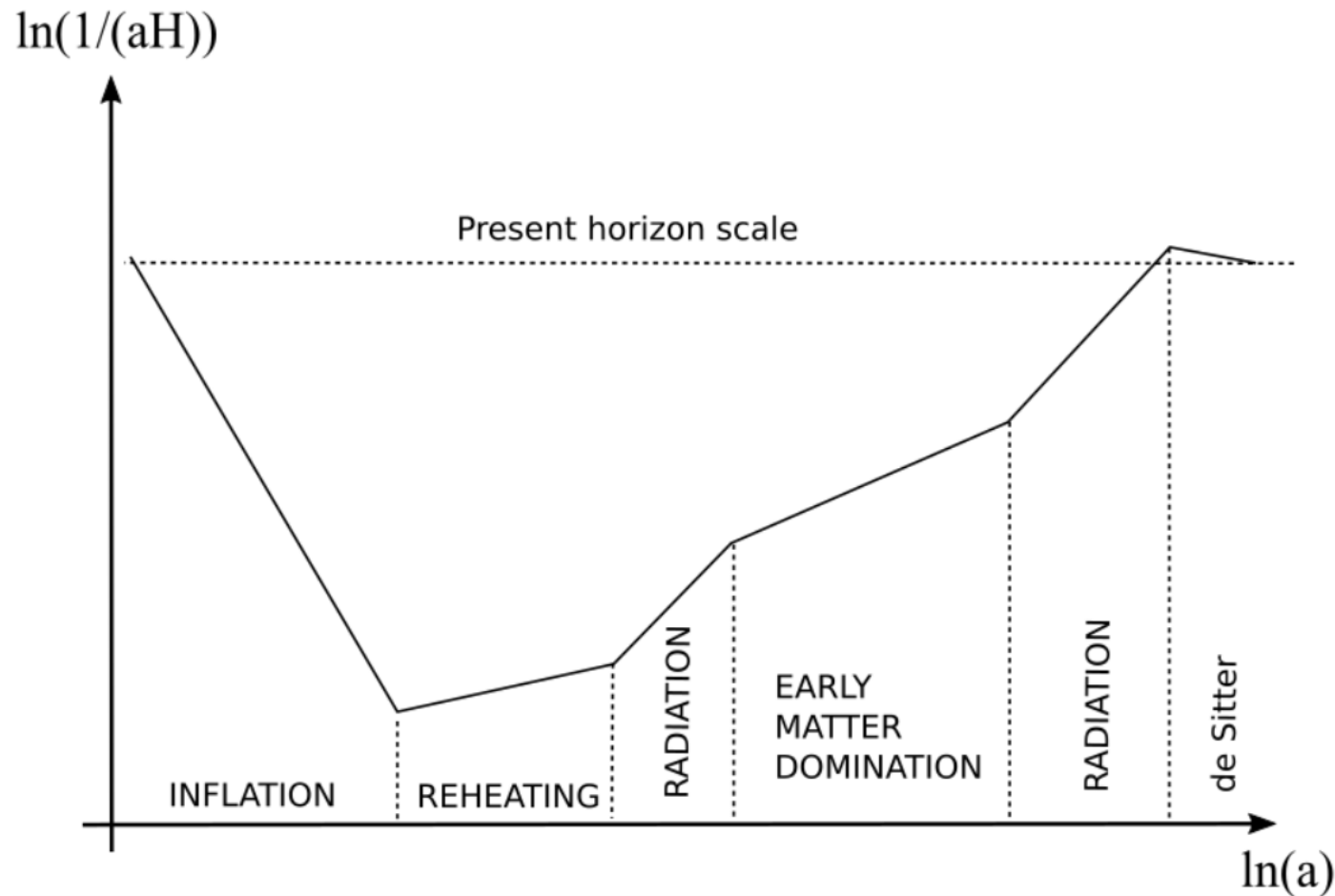
Modifies the number of e-folds

What about the primordial universe?



$$\frac{k}{a_0 H_0} = \frac{a_k H_k}{a_0 H_0} = e^{-N(k)} \frac{a_{\text{end}}}{a_{\text{reh}}} \frac{a_{\text{reh}}}{a_{\text{dom}}} \frac{a_{\text{dom}}}{a_{\text{eq}}} \frac{H_k}{H_{\text{eq}}} \frac{a_{\text{eq}} H_{\text{eq}}}{a_0 H_0}$$

What about the primordial universe?



Chaotic inflation $\rightarrow N < 50$ e-folds **EXCLUDED**

Other inflation scenarios : To appear soon !

Conclusion(I)

- Dark matter production usually requires fine tuning or the introduction of arbitrary mass scales
- We propose an ***inflaton portal*** to a highly decoupled dark sector
 - Reheating process explicitly present in the scenario
 - Natural choices of couplings lead to the correct relic abundance
- The model escapes any direct or indirect detection
- The presence of light fields and of an early matter domination period provides strong constraints on the inflationary sector...
- Chaotic inflation is ruled out in this framework

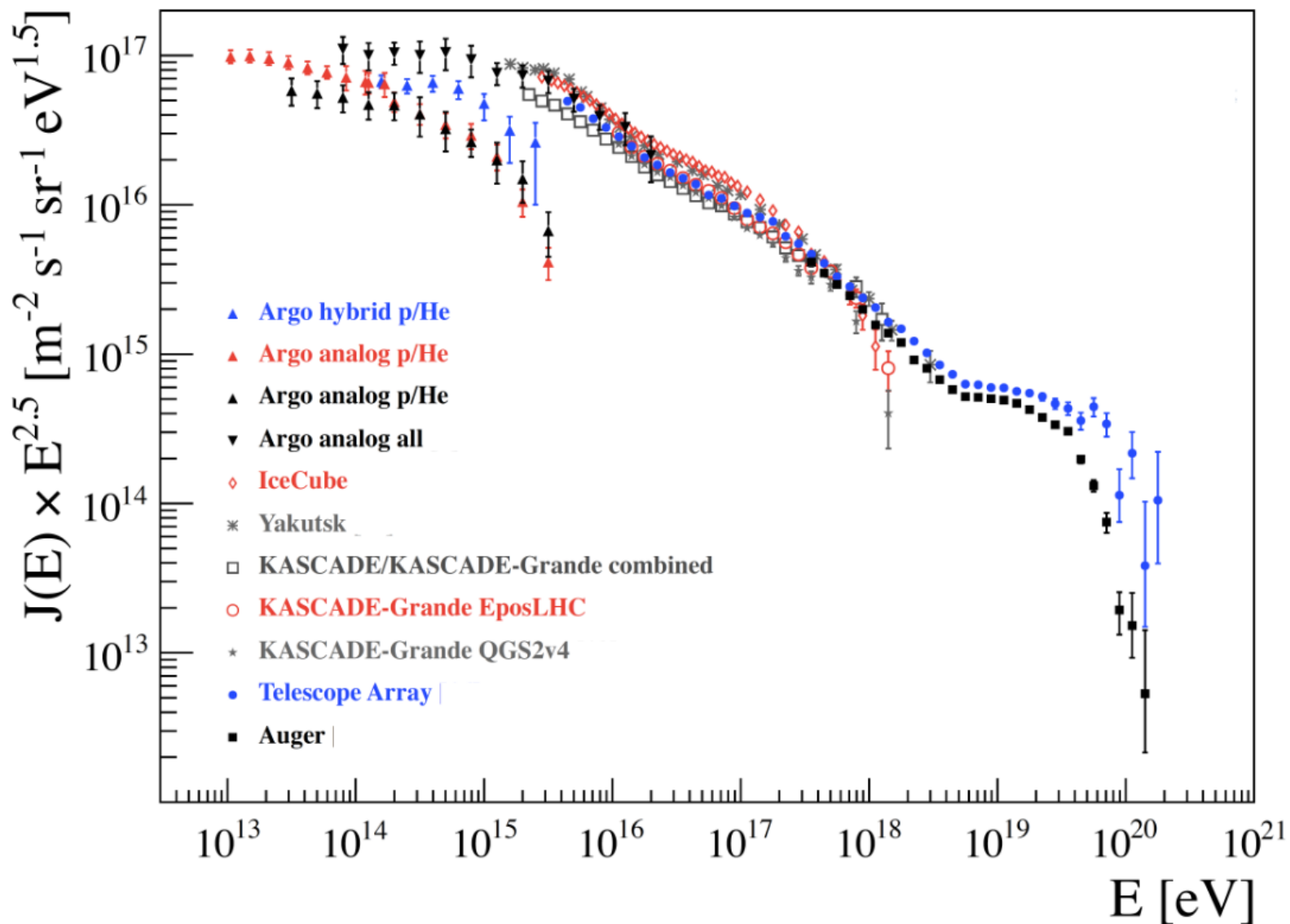
Part II

Signatures of EeV scale dark matter in the neutrino sector

What do we expect at ultra high energy?

Up to $E < 5 \times 10^{19} \text{ GeV}$, primary proton cosmic rays can propagate (also known as Greisen-Zatsepin-Kuzmin (GZK) limit)

Above such energy, primary cosmic rays scatter on the CMB to produce *cosmogenic* neutrinos which remain yet undetected...



Searching for Heavy Dark Matter in the Neutrino Sector

- Presence of decaying HDM in the galaxy
→ Ultra-High-Energy Cosmic Rays

At energies larger than $>\text{TeV}$ **photons scatter on the CMB...**

... but **neutrinos can propagate!**

→ Search for decay products in the neutrino sector !

Experimental Status

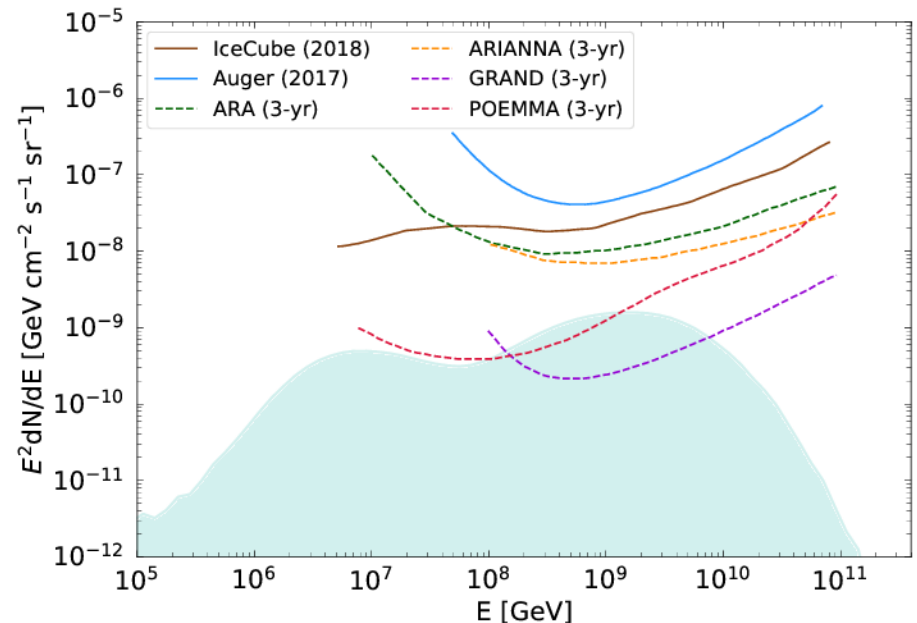
- IceCube, Auger and ANITA have been operating for several years

Signals:

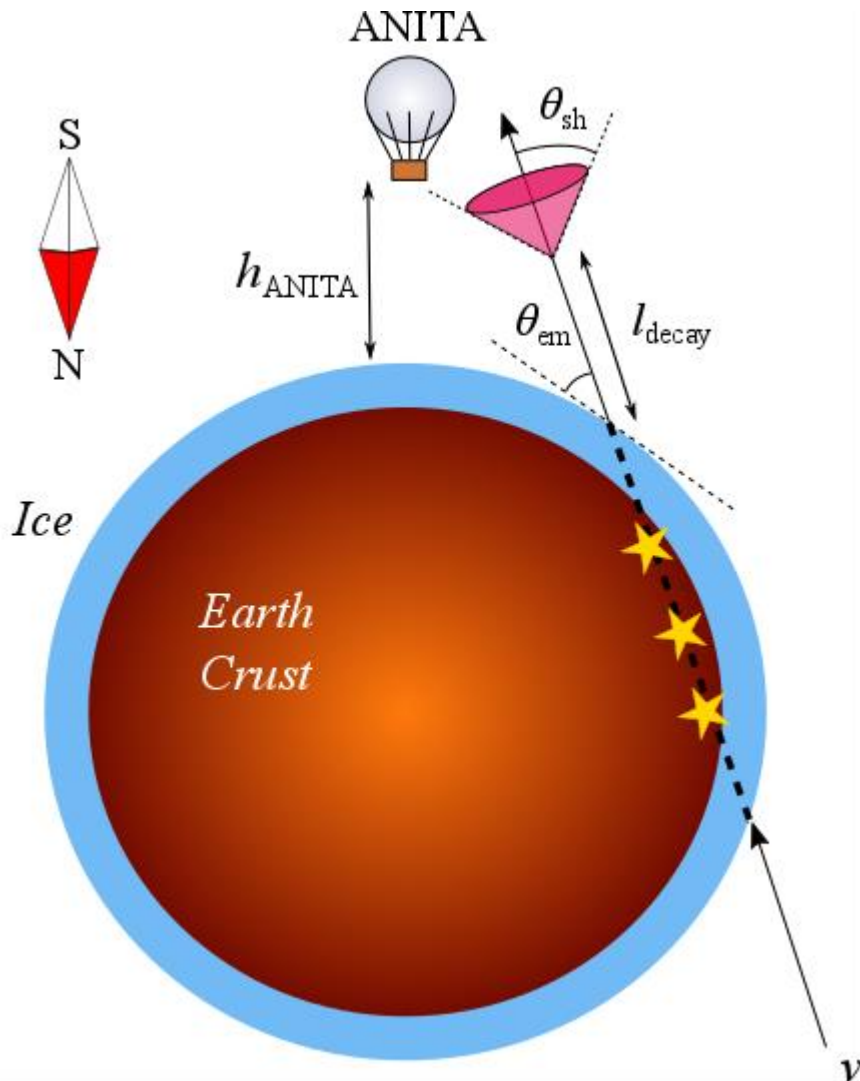
- IceCube : A few PeV events in ten years
- ANITA : 2 anomalous upgoing EeV events in 85 days

Future detectors:

POEMMA, GRANDE, ...



The ANITA experiment



Three types of interactions in the SM:

- Charged current: Neutrinos \rightarrow leptons
- Neutral current : Neutrinos \rightarrow Neutrinos
- Regeneration: Lepton \rightarrow Neutrino + sh.

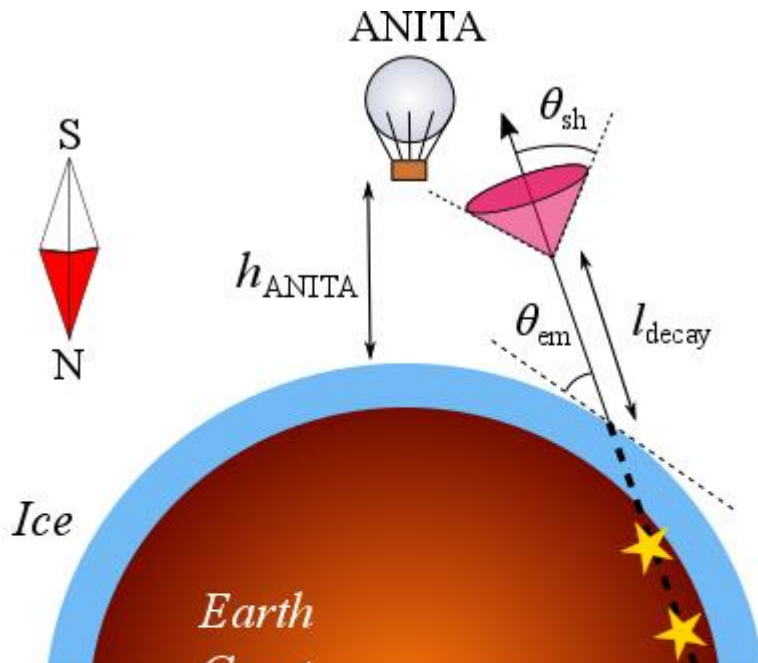
Propagation depends on the particle energy, local density, and particle interactions...

Two anomalous events with emergence angles 27° and 35° and energies $O(1)\text{EeV}$.

At such emergence angles, a SM neutrino cannot cross the Earth.

No astrophysical source identified.

The ANITA experiment



Three types of interactions in the SM:

- Charged current: Neutrinos \rightarrow leptons
- Neutral current : Neutrinos \rightarrow Neutrinos
- Regeneration: Lepton \rightarrow Neutrino + sh.

Propagation depends on the particle energy, local density, and particle interactions...

Property	AAE 061228	AAE 141220
Flight & Event	ANITA-I #3985267	ANITA-III #15717147
Date & Time (UTC)	2006-12-28 00:33:20	2014-12-20 08:33:22.5
Equatorial coordinates (J2000)	R.A. 282°14064, Dec. +20°33043	R.A. 50°78203, Dec. +38°65498
Energy ε_{cr}	0.6 ± 0.4 EeV	$0.56^{+0.30}_{-0.20}$ EeV
Zenith angle z'/z	$117.4 / 116.8 \pm 0.3$	$125.0 / 124.5 \pm 0.3$
Earth chord length ℓ	5740 ± 60 km	7210 ± 55 km
Mean interaction length for $\varepsilon_\nu = 1$ EeV	290 km	265 km
$p_{SM}(\varepsilon_\tau > 0.1 \text{ EeV})$ for $\varepsilon_\nu = 1$ EeV	4.4×10^{-7}	3.2×10^{-8}
$p_{SM}(z > z_{obs})$ for $\varepsilon_\nu = 1$ EeV, $\varepsilon_\tau > 0.1$ EeV	6.7×10^{-5}	3.8×10^{-6}
$n_\tau(1-10 \text{ PeV}) : n_\tau(10-100 \text{ PeV}) : n_\tau(> 0.1 \text{ EeV})$	34 : 35 : 1	270 : 120 : 1

The ANITA experiment

[Fox, Sigurdson, Murase *et al.*, Nov 18']

The SM can't explain ANITA events.

+

A re-analysis of IceCube data

~~Two massive events observed at 0.015 EeV~~

Property	IceCube-140611	IceCube-140109	IceCube-121205
EHE Northern Track ID	#27	#24	#20
Date & Time (UTC or MJD)	2014-06-11 04:54:24	56666.5	56266.6
Equatorial coordinates (J2000)	R.A. $110^{\circ}34 \pm 0^{\circ}22$, Dec. $+11^{\circ}42 \pm 0^{\circ}08$	R.A. $293^{\circ}29$, Dec. $+32^{\circ}82$	R.A. $169^{\circ}61$, Dec. $+28^{\circ}04$
Zenith angle z	$101^{\circ}42$	$122^{\circ}82$	$118^{\circ}04$
Earth chord length ℓ	2535 km	6910 km	5990 km
As tau: $\varepsilon_{\tau, \text{obs}}$ (median)	70 PeV	13 PeV	12 PeV
Mean interaction length for $\varepsilon_{\nu} = 1 \text{ EeV}$	340 km	270 km	285 km
$p_{\text{SM}}(\varepsilon_{\tau} > \varepsilon_{\tau, \text{obs}})$ for $\varepsilon_{\nu} = 1 \text{ EeV}$	2.2×10^{-4}	3.8×10^{-6}	1.0×10^{-5}
$p_{\text{SM}}(z > z_{\text{obs}})$ for $\varepsilon_{\nu} = 1 \text{ EeV}$, $\varepsilon_{\tau} > \varepsilon_{\tau, \text{obs}}$	5.0×10^{-3}	4.5×10^{-5}	1.8×10^{-4}

Question: Can both signals have a common origin?

A right-handed neutrino interpretation

Consider:

- A right-handed neutrino mixes with the tau neutrino
- A scalar dark-matter particle decays into sterile neutrinos in the galaxy

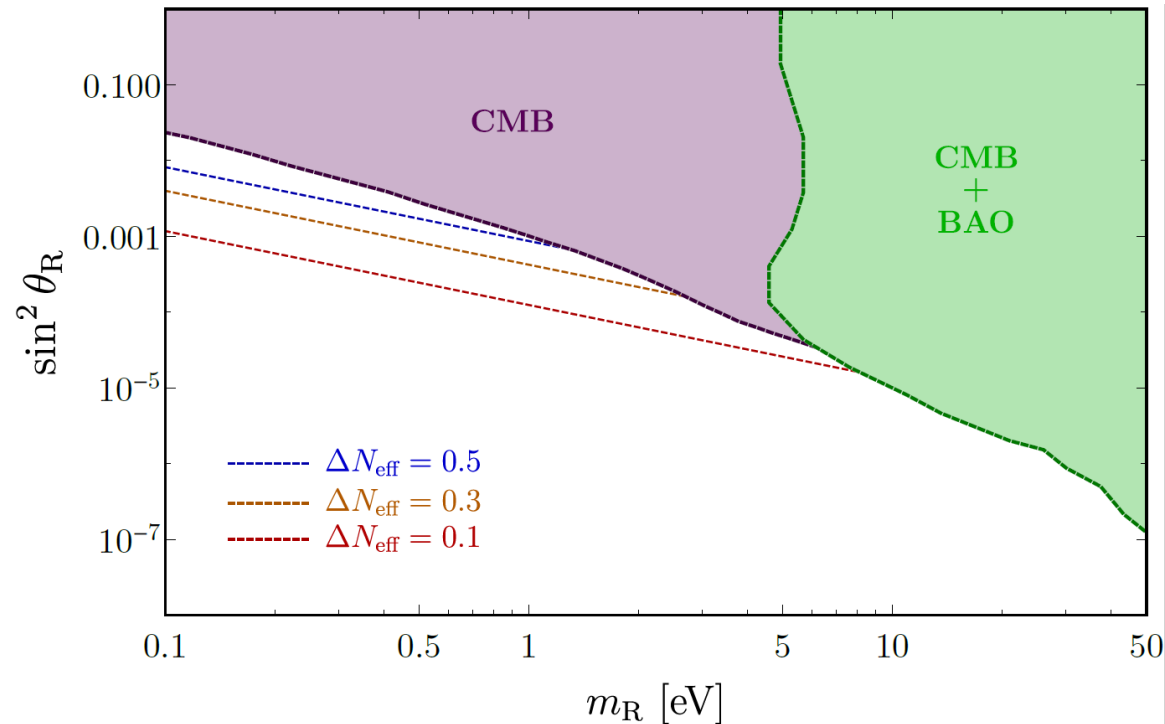
One needs:

→ RH neutrinos sufficiently long lived

→ To satisfy observational bounds (BBN, CMB, direct searches)

A right-handed neutrino interpretation

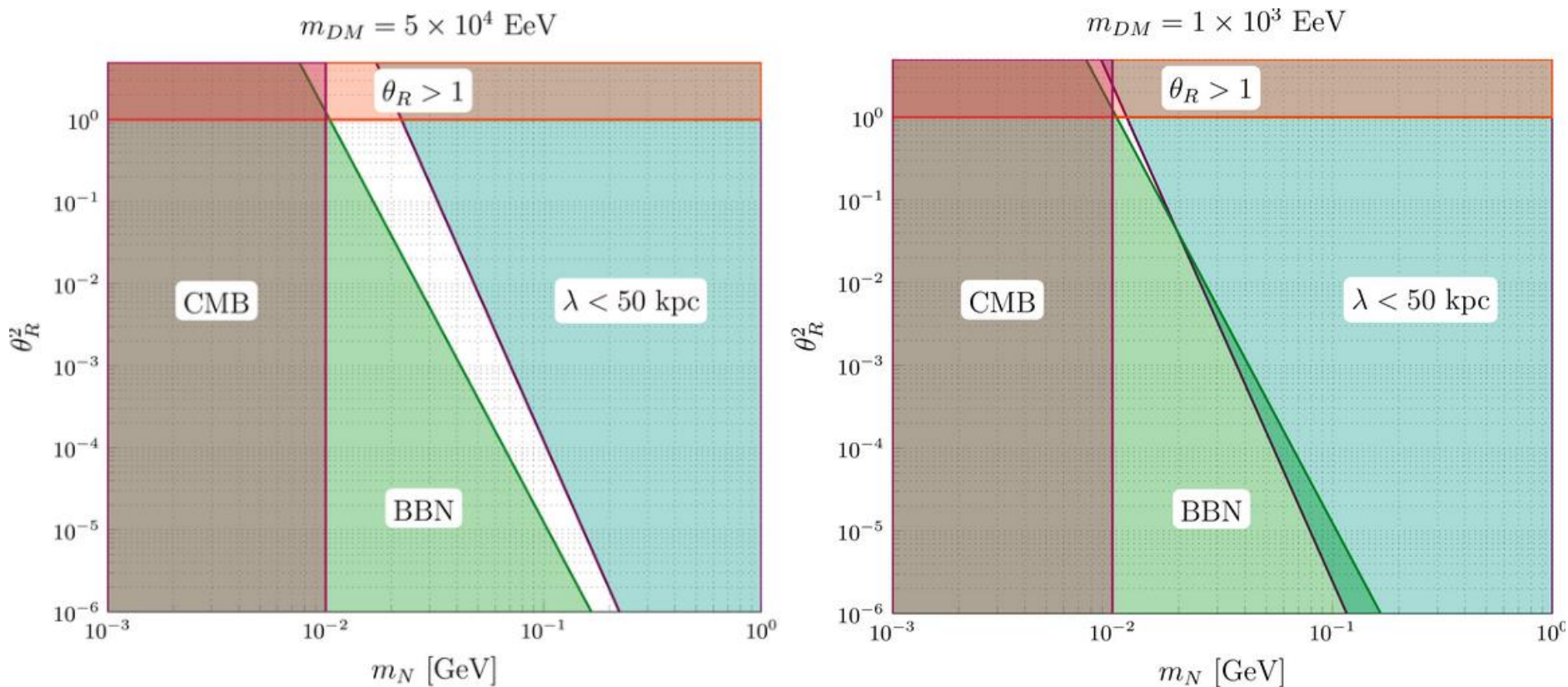
$$m_R < 10 \text{ eV} \text{ or } m_R \sim 0.1 \text{ GeV}$$



$$m_R < 10 \text{ eV} \longrightarrow m_{DM} \sim 10 \text{ EeV}$$

A right-handed neutrino interpretation

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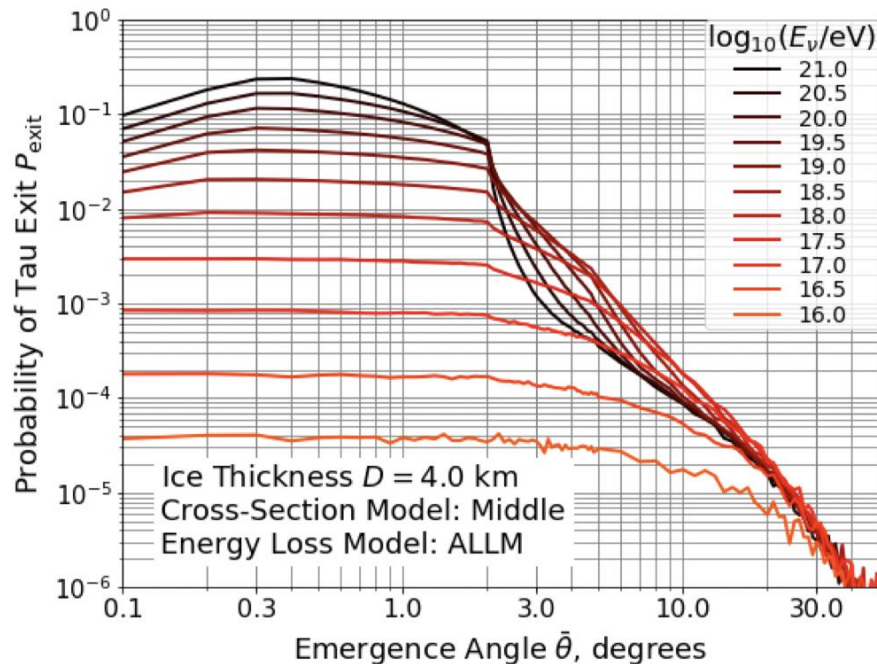


$$m_R \sim 0.1 \text{ GeV} \longrightarrow m_{DM} \gtrsim 10^3 \text{ EeV} = 10^{12} \text{ GeV}$$

Approaches the inflaton mass...

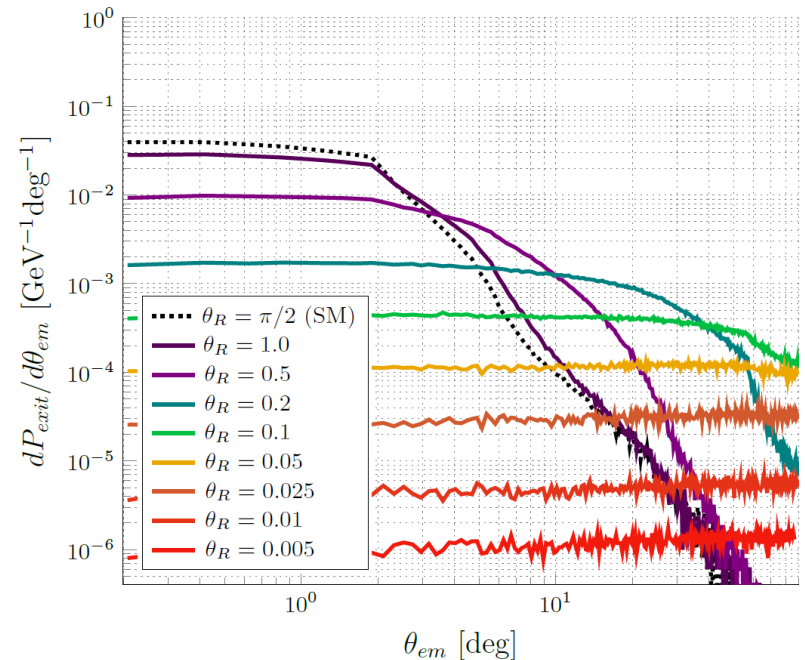
A right-handed neutrino interpretation

Propagation and conversion into tau's



SM

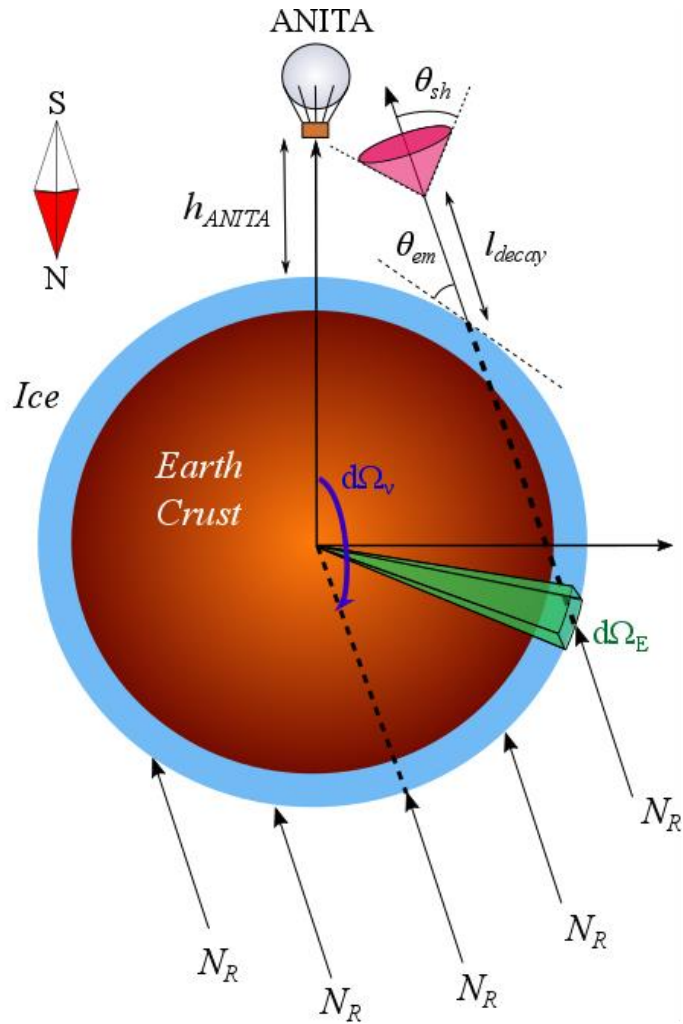
[Romero-Wolf *et al.*, Dec 17']



SM + RH neutrinos and mixing angle

[LH, Y. Mambrini, M. Pierre, To appear]

Effective Area Calculation



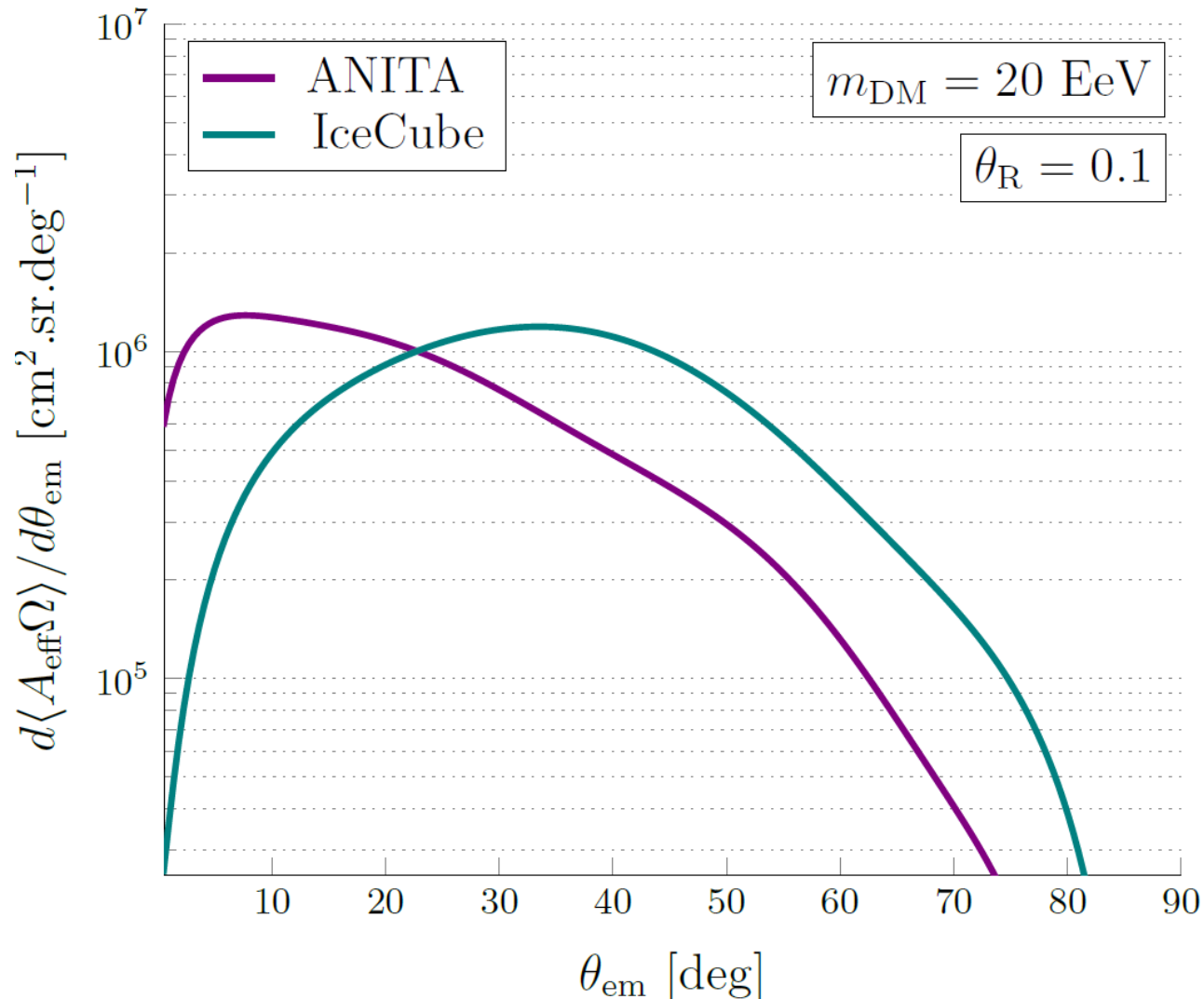
$$\begin{aligned} & \frac{d^2 A_{\text{eff}}}{dE_{\text{exit}} d\theta_{\text{em}}} (E_{\text{exit}}, \theta_{\text{em}} \mid E_{\text{N}}, \theta_{\text{N}}, \phi_{\text{N}}) \\ &= R_{\text{E}}^2 \int d\Omega_{\text{E}} \vec{n}_{\text{N}} \cdot \vec{n}_{\text{E}} \\ &\times \int dE_{\text{exit}} \frac{dP_{\text{exit}}}{dE_{\text{exit}}} (E_{\text{exit}}, \theta_{\text{em}} \mid E_{\text{N}}, \theta_{\text{N}}, \phi_{\text{N}}, \theta_{\text{E}}, \phi_{\text{E}}) \\ &\times \int \frac{dP_{\text{decay}}}{dl} (l \mid E_{\text{exit}}) \times P_{\text{det}}(\theta_{\text{sh}} | l, \theta_{\text{N}}, \phi_{\text{N}}, \theta_{\text{E}}, \phi_{\text{E}}) dl \end{aligned}$$

ANITA and IceCube detection

Simulate the detection probability of a shower
for IceCube and ANITA

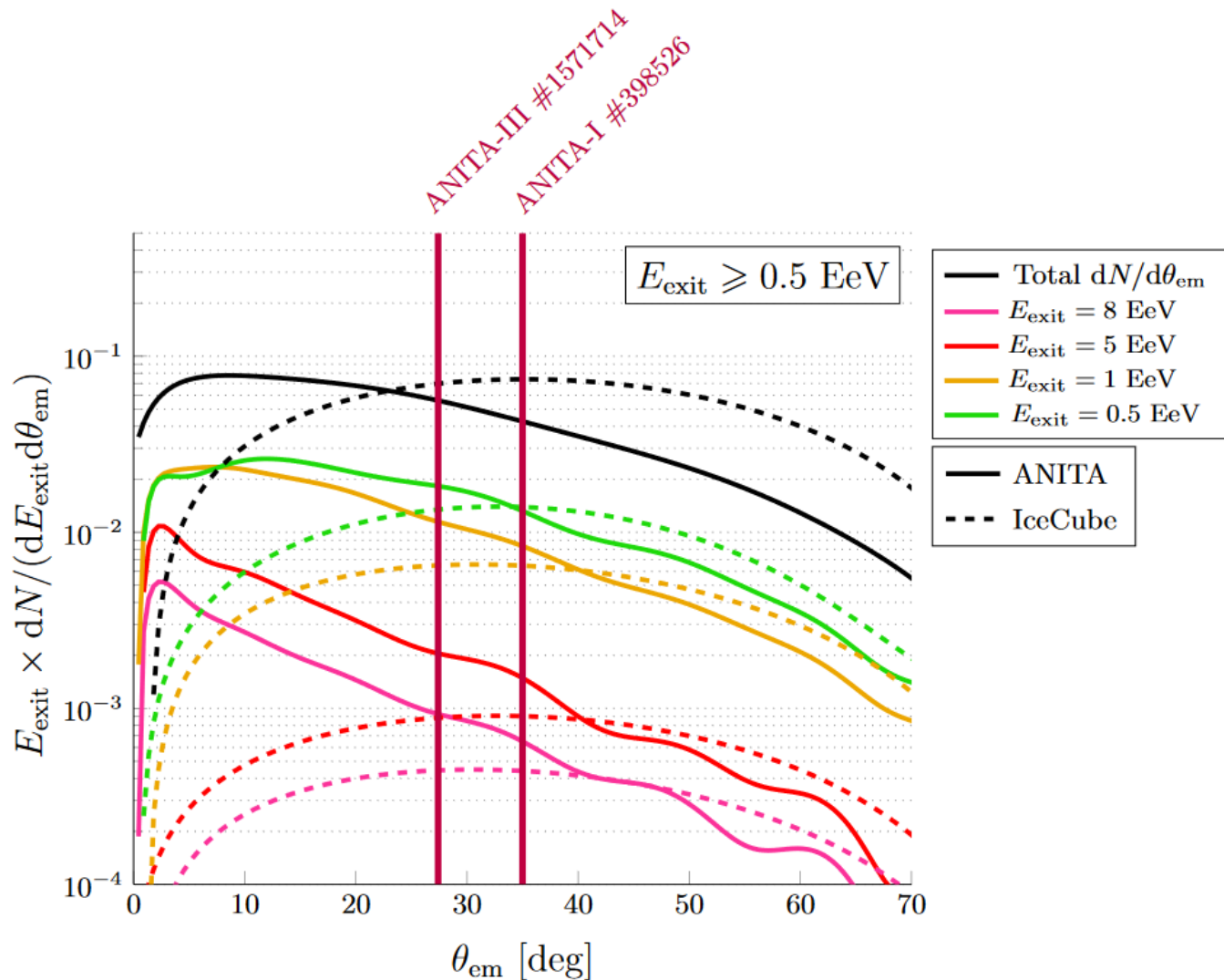
- IceCube : Fiducial volume of 1km under the Earth surface
- ANITA : shower must be produced in the low atmosphere, within a cone of around 1.5° away from the detector

ANITA and IceCube detection

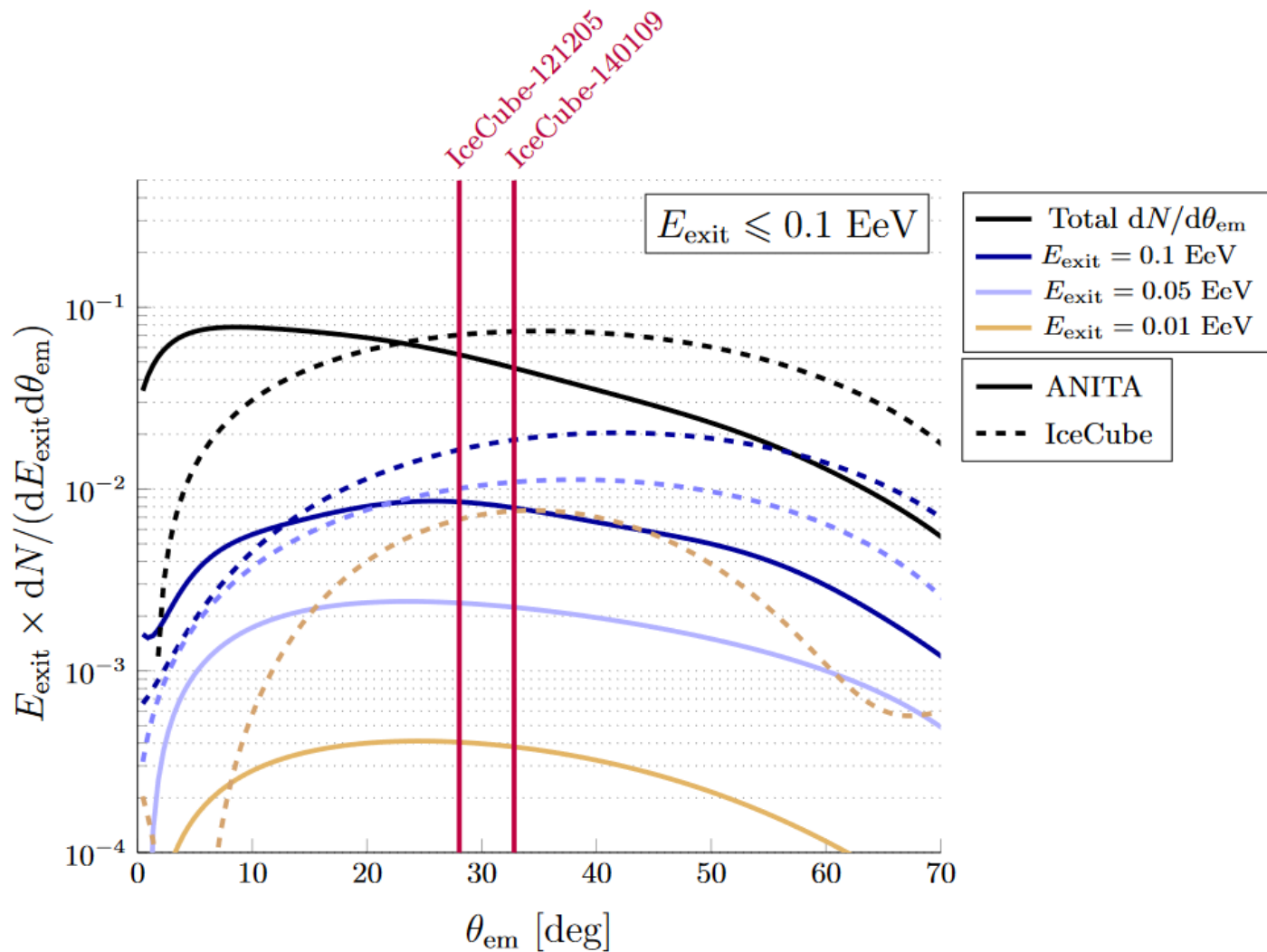


$N_{\text{IC}} = 2.35 \quad (2431 \text{ days})$
 $N_{\text{ANITA}} = 2.29 \quad (85 \text{ days})$

ANITA detection



IceCube detection



Conclusion (II)

- A dark matter of mass as large as $\gtrsim 10\text{EeV}$ decaying indirectly into neutrinos might lead to interesting signatures in the neutrino sector.
- The presence of a BSM sector modifies the way the neutrinos can propagate in the Earth
- A dark scalar decaying into right-handed neutrinos can explain the recent measurements of both IceCube and ANITA
- Given this model, a dark matter mass of mass $>10\text{ EeV}$ predicts a perfect complementarity between the two collaborations.

Thank you very much