

Confronting Grand Unification with LFV Dark Matter and LHC data

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*Colaboration with M. Canonni, J. Ellis, S. Lola, Ruiz de Austri
JCAP 1603 (2016)
and Q. Shafi, JHEP 1810 (2018) 062 Front.in Phys. 6 (2018) and
work in progress.*

OUTLINE

- GUT's and SUSY.

- SU(5), Flipped-SU(5) and SU(4)xSU(2)xSU(2)

- *Neutralino relic density and DM detection.*

- *Charged LFV in SUSY with a see-saw T1,*

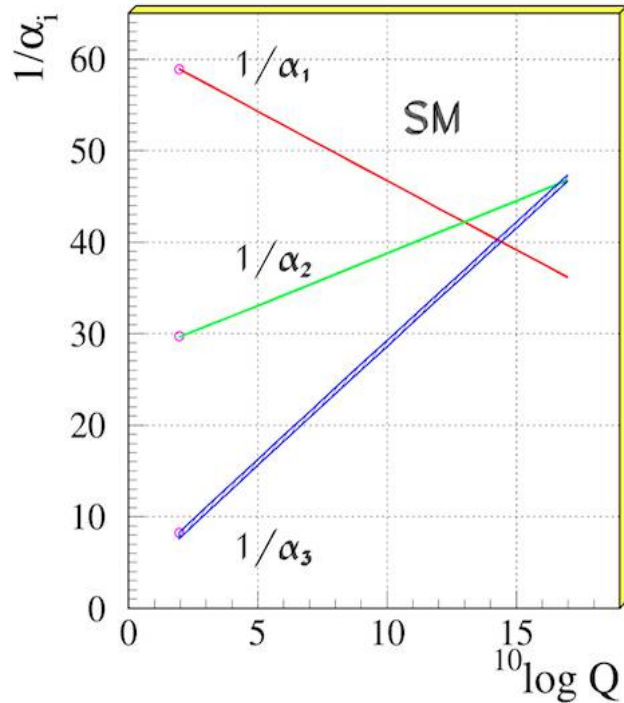
- SUSY Masses, LHC searches and rare lepton decays.*

Conclusions.

SUSY extension of SM

- Divergence cancellations
- Gauge unification
- Particles in the range of coming accelerators.
- DM candidate.
- Small but sizeable contribution to SM processes...

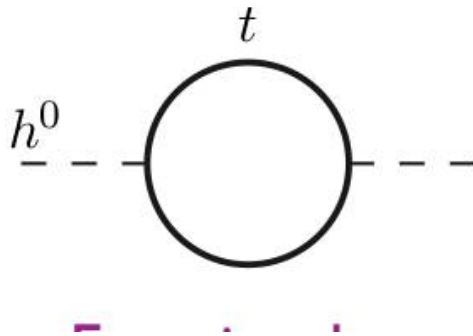
Unification & Hierarchy Problem



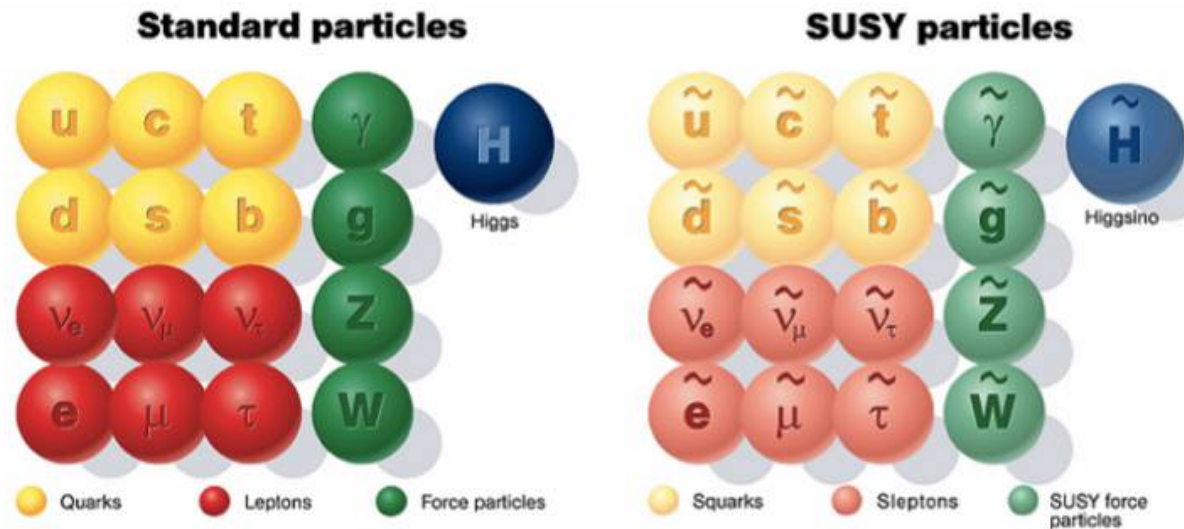
Almost unification at $\sim 10^{14}$ GeV

$$\mu \frac{d\alpha_i(\mu)}{d\mu} = -\frac{1}{2\pi} \left[b_i + \frac{1}{4\pi} \sum b_{ij} \alpha_j(\mu) \right] \alpha_i^2(\mu)$$

$$b_i = (0, -22/3, -11) + N_F(4/3, 4/3, 4/3) + N_H(1/10, 1/6, 0)$$



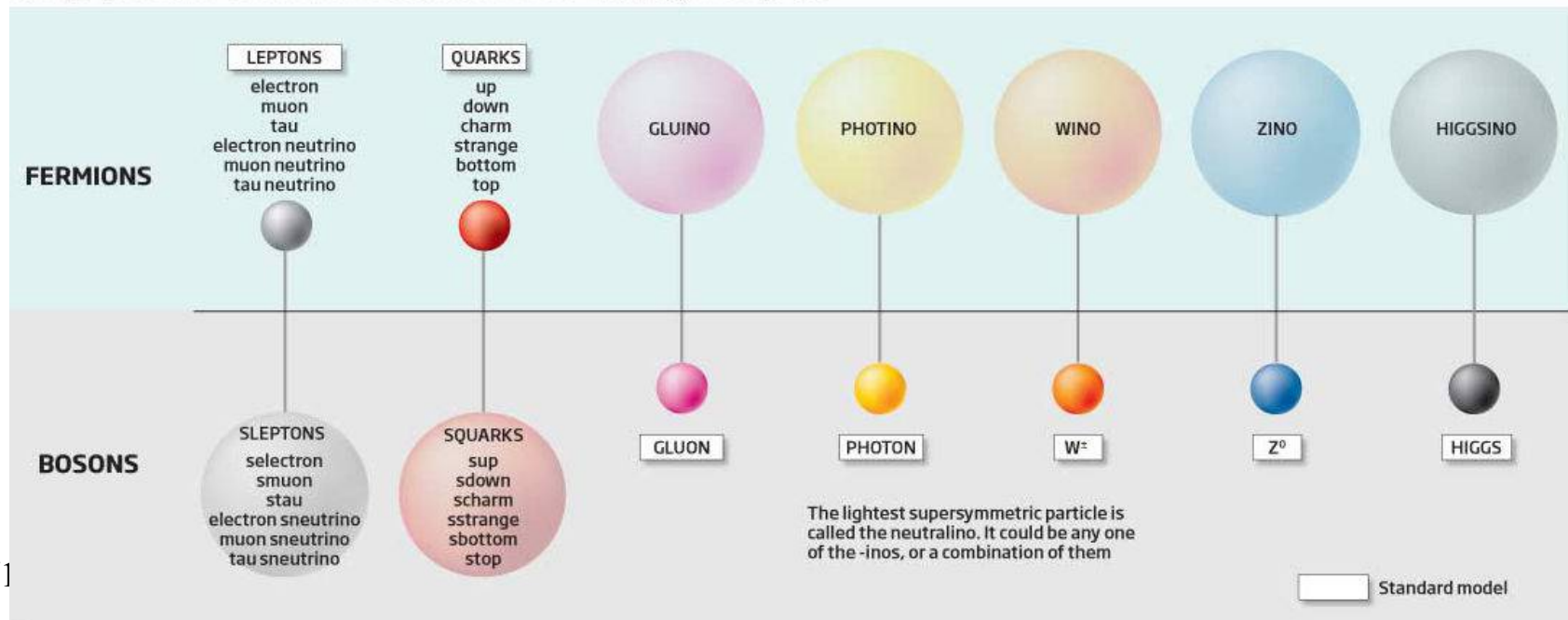
Divergent contribution to Higgs mass $\uparrow$ with $(m_{\text{scale}})^2$



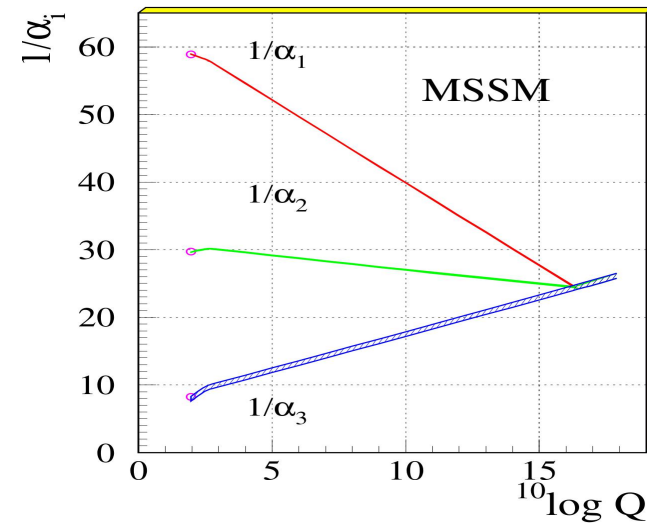
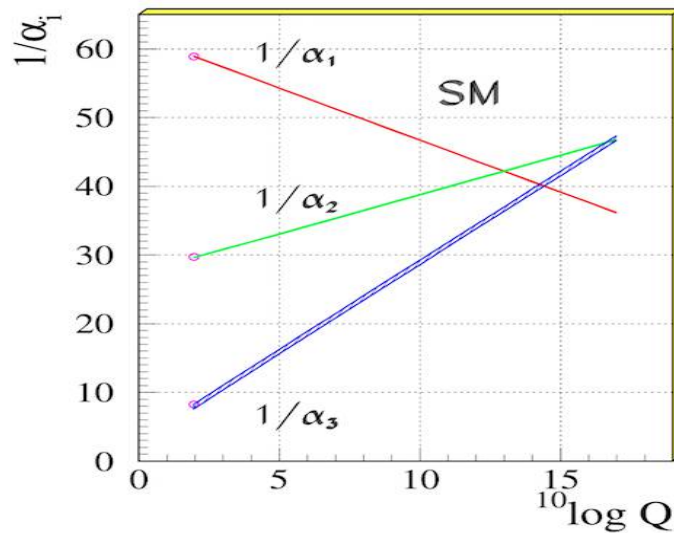
Particle zoo

©NewScientist

Particles are divided into two families called bosons and fermions. Among them are groups known as leptons, quarks and force-carrying particles like the photon. Supersymmetry doubles the number of particles, giving each fermion a massive boson as a super-partner and vice versa. The LHC is expected to find the first supersymmetric particle.

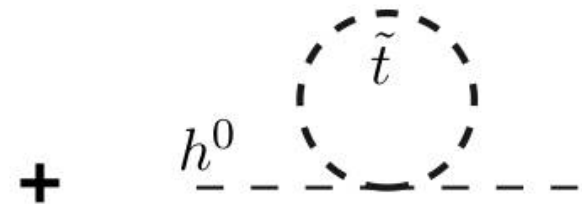
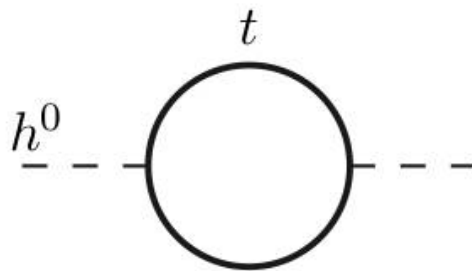


Unification & Hierarchy Problem

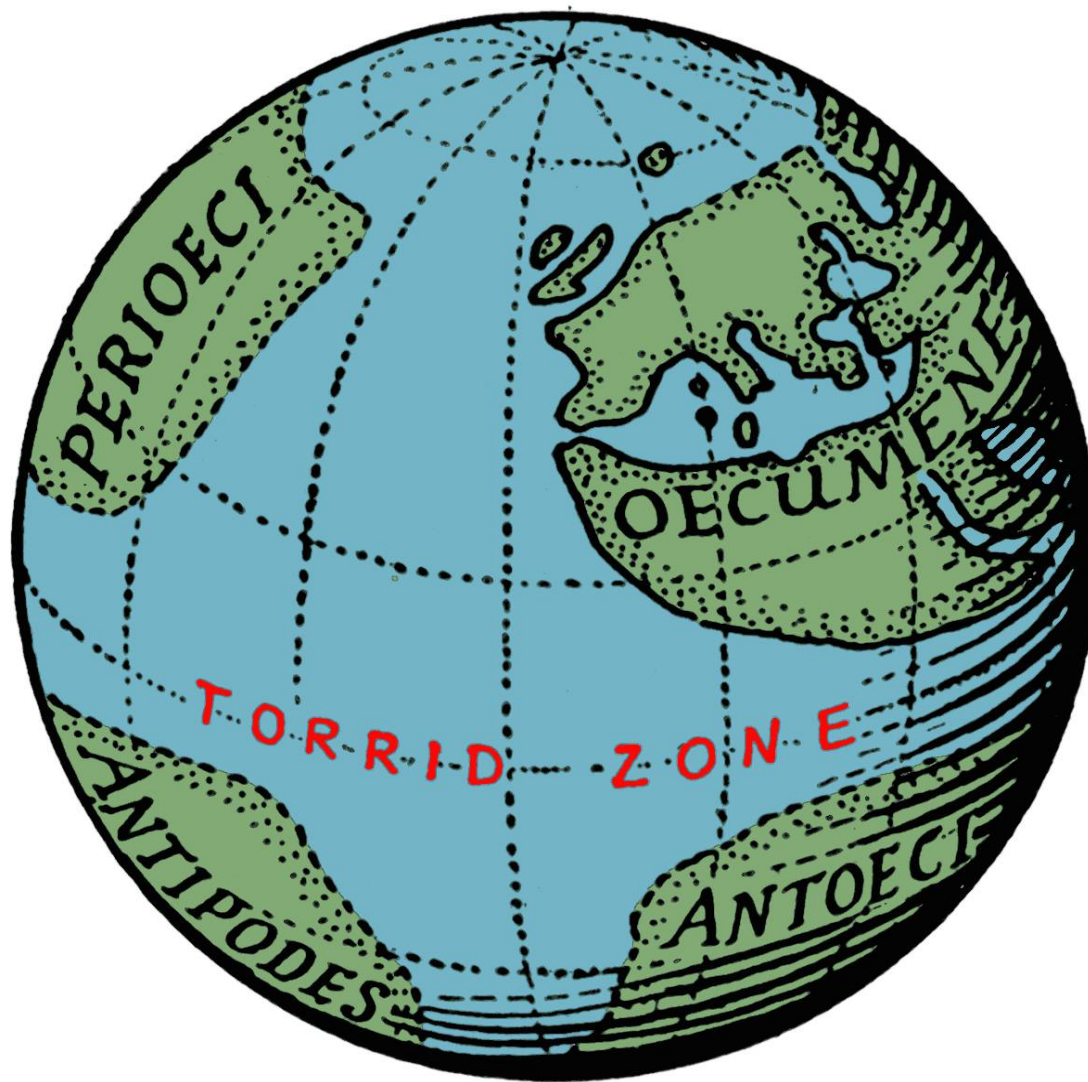


Unification at $\sim 10^{16}$ GeV

$$b_i = (0, -6, -9) + N_F(2, 2, 2) + N_H(3/10, 1/2, 0)$$



Cancelation of quadratic divergences



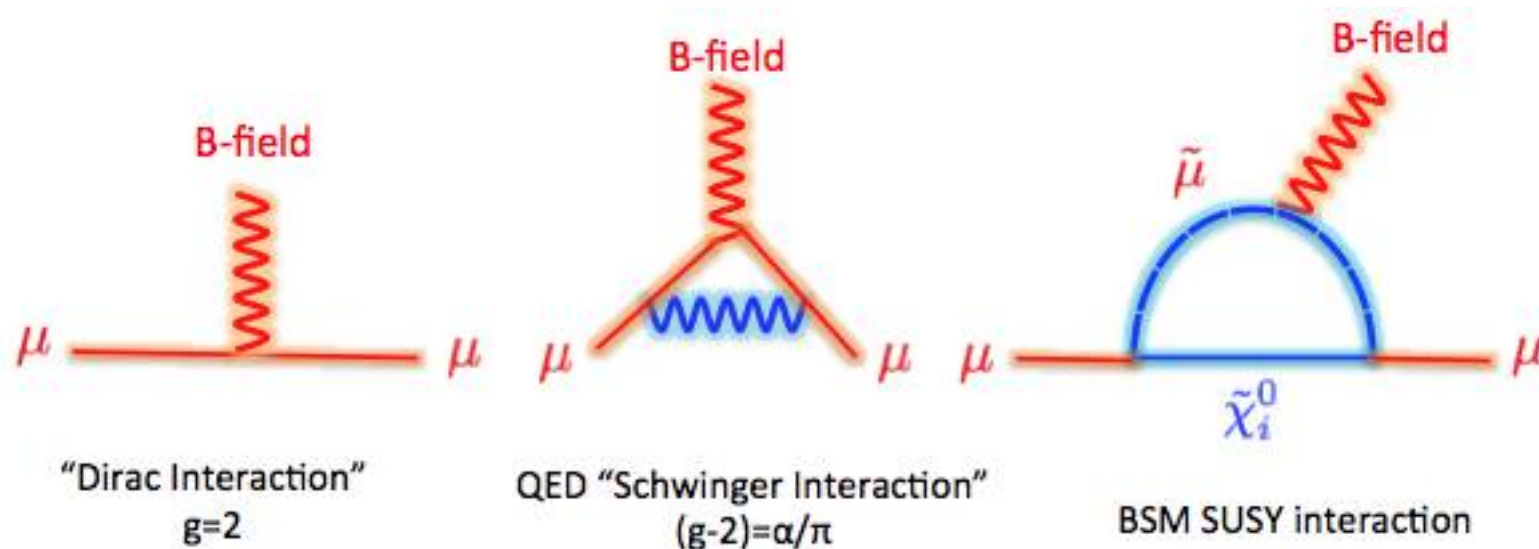
Antipodes theory:

Popular debate in the Middle Ages.

Probably referred in all exploration trip proposals until it was eliminated by direct observation.



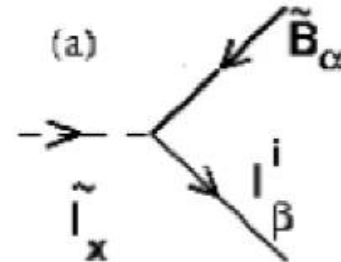
One loop contribution to SM proc.



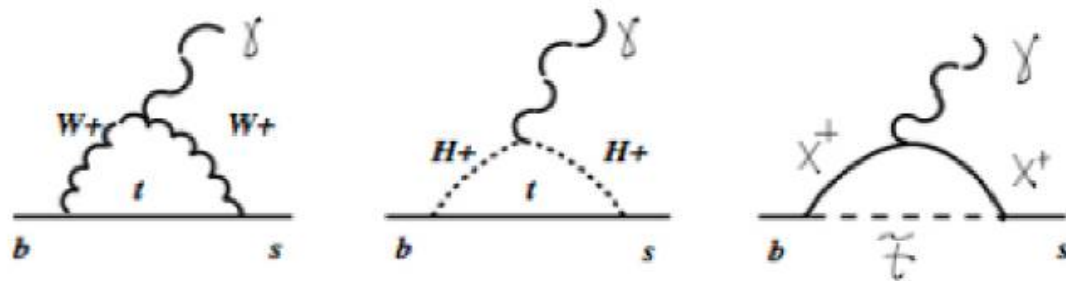
SM vs Experiment Discrepancies
anomalous magnetic dipole moment ($g_\mu - 2$)

SUSY FLAVOR

R-parity warranties that SUSY particles only appear in pairs:



therefore SM model phenomenology is only modified at *loops level*:



The present average given by the

$$BR(b \rightarrow s\gamma) = (3.32 \pm 0.15) \times 10^{-4}$$

The SM prediction: $BR(b \rightarrow s\gamma) = (3.36 \pm 0.23) \times 10^{-4}$

Neutrino Oscillations

$$\mathcal{L}_{\text{CC}} = - \frac{g}{\sqrt{2}} \sum_{l=e,\mu,\tau} \bar{l}_L(x) \gamma_\alpha \nu_{lL}(x) W^{\alpha\dagger}(x) + h.c.,$$

$$\nu_{lL}(x) = \sum_{j=1}^3 U_{lj} \nu_{jL}(x),$$

EW interaction
Mixes Lepton
Flavor.

**There is a matrix:
PNMS=U
Equivalent to CKM
in the Quark sector.**

$$P(\nu_{l(l')} \rightarrow \nu_{l'(l)}) \cong P(\bar{\nu}_{l(l')} \rightarrow \bar{\nu}_{l'(l)}) \cong \delta_{ll'} - 2|U_{ln}|^2 \left[\delta_{ll'} - |U_{l'n}|^2 \right] \\ (1 - \cos \frac{\Delta m_{n1}^2}{2p} L).$$

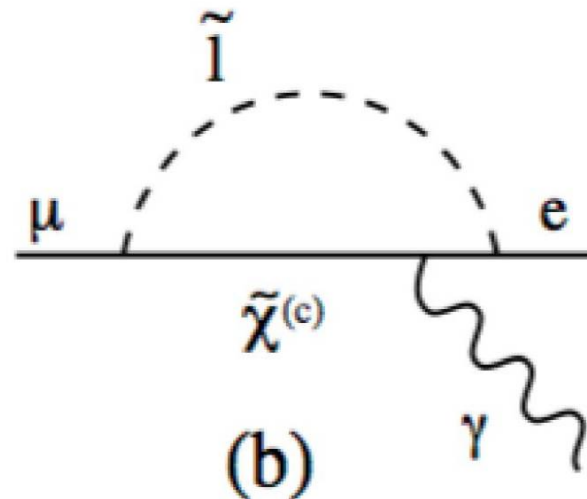
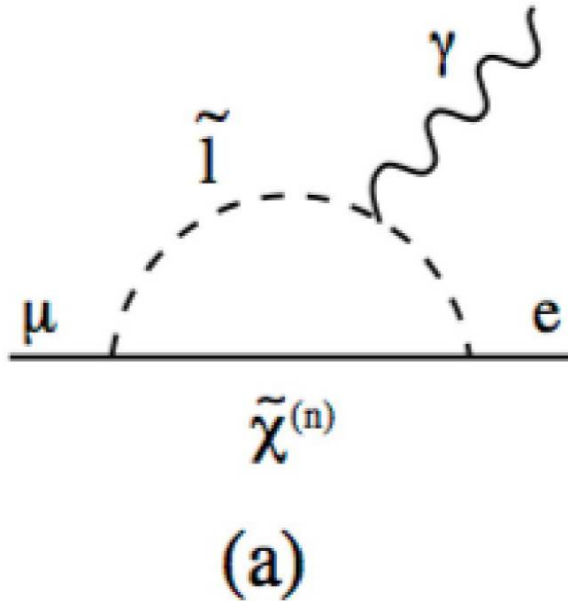
$$U = \begin{bmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{bmatrix} \\ \times \text{diag}(1, e^{i\frac{\alpha_{21}}{2}}, e^{i\frac{\alpha_{31}}{2}}).$$

Atmospheric problem	Solar problem
$\Delta m_{atm}^2 = (2.6^{+0.4}_{-0.7}) \times 10^{-3} \text{ eV}^2$ $\sin^2 2\theta_{atm} > 0.90$	$\Delta m_{sol}^2 = (8.1^{+0.5}_{-0.5}) \times 10^{-5} \text{ eV}^2$ $\sin^2 2\theta_{sol} = (0.86^{+0.05}_{-0.06})$

$$0.06 < \sin^2 2\theta_{13} < 0.13$$

Reactor data (RENO, Daya Bay).

In SUSY flavor mixing lepton-slepton vertices can induce LFV diagrams:



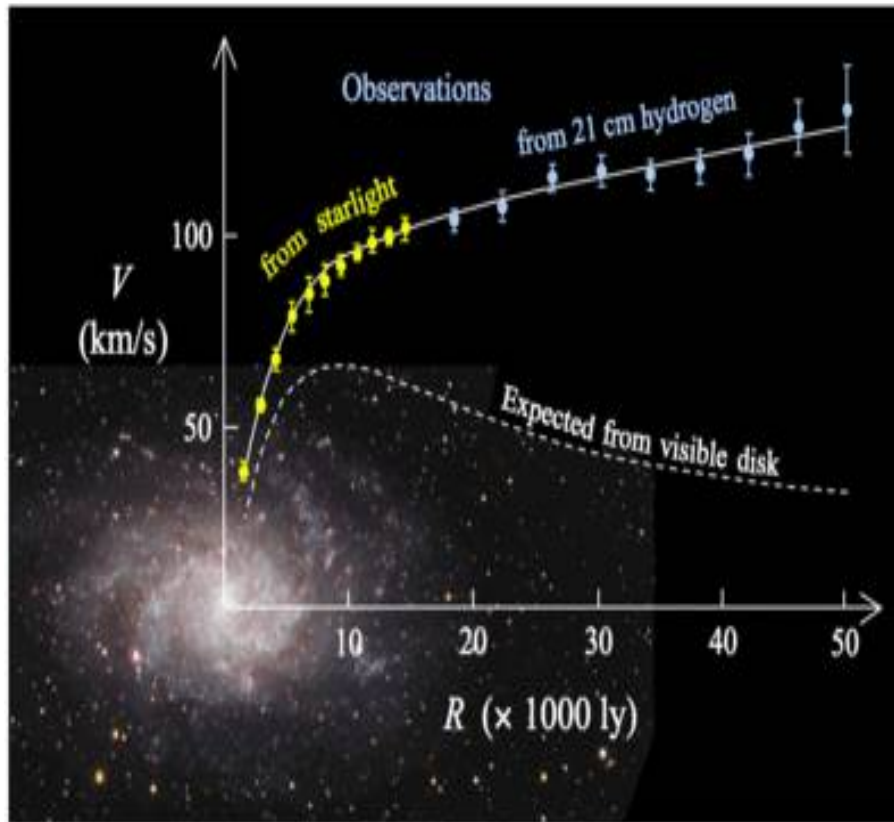
Lepton-slepton flavor mixing is very constrained by the experimental limits:

$$BR(\mu \rightarrow e \gamma) < 5.7 \times 10^{-13}$$

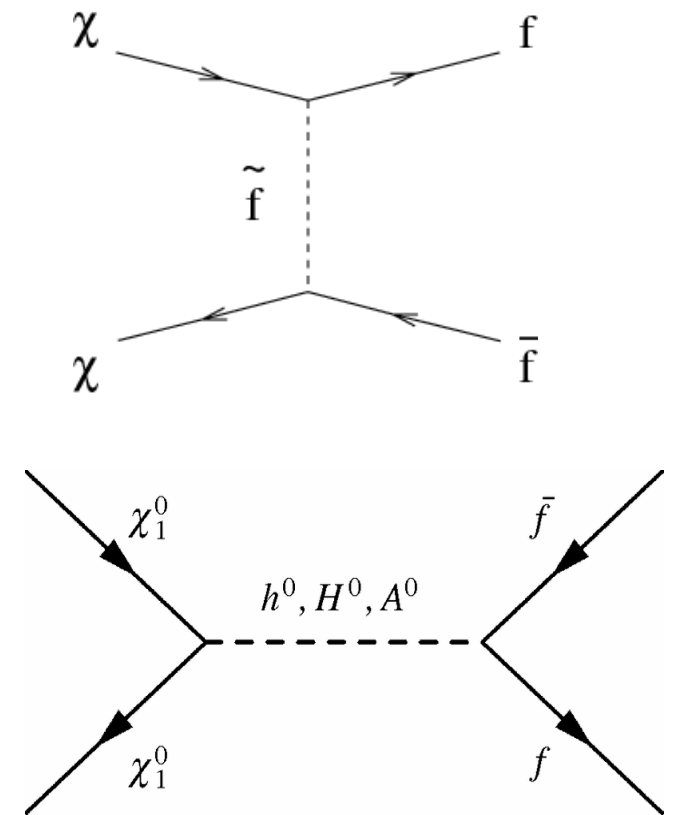
$$BR(\tau \rightarrow \mu \gamma) < 4.4 \times 10^{-8}$$

$$BR(\tau \rightarrow e \gamma) < 3.3 \times 10^{-8}$$

Dark Matter problem



Rotation curve of spiral galaxy M 33 (yellow and blue points with error bars), and a predicted one from distribution of the visible matter (white line). The discrepancy between the two curves can be accounted for by adding a dark matter halo surrounding the galaxy



$$\Omega h^2 \sim \frac{3 \times 10^{-27} \text{ cm}^3 \text{ s}^{-1}}{\langle \sigma_{\text{ann}} v_{\text{Mol}} \rangle}$$

Soft SUSY Breaking Terms

The soft SUSY breaking masses

$$\begin{aligned}
 -\mathcal{L}_{\text{soft}} = & -\frac{1}{2} \left(M_3 \lambda_g^a \lambda_g^a + M_2 \lambda_{\tilde{W}}^i \lambda_{\tilde{W}}^i + M_1 \lambda_{\tilde{B}} \lambda_{\tilde{B}} + \text{h.c.} \right) \\
 & + M_L^2 \tilde{L}^\dagger \tilde{L} + M_Q^2 \tilde{Q}^\dagger \tilde{Q} + M_U^2 \tilde{U}^* \tilde{U} + M_D^2 \tilde{D}^* \tilde{D} + M_E^2 \tilde{E}^* \tilde{E} + \\
 & m_{H_d}^2 \tilde{H}_d^\dagger \tilde{H}_d + m_{H_u}^2 H_u^\dagger H_u - \left(B\mu \tilde{H}_d^T H_u + \text{h.c.} \right) \\
 & + \left(y_\ell A_\ell H_d^\dagger \tilde{L} \tilde{E} + y_d A_d H_d^\dagger \tilde{Q} \tilde{D} - y_u A_u H_u^T \tilde{Q} \tilde{U} + \text{h.c.} \right),
 \end{aligned}$$

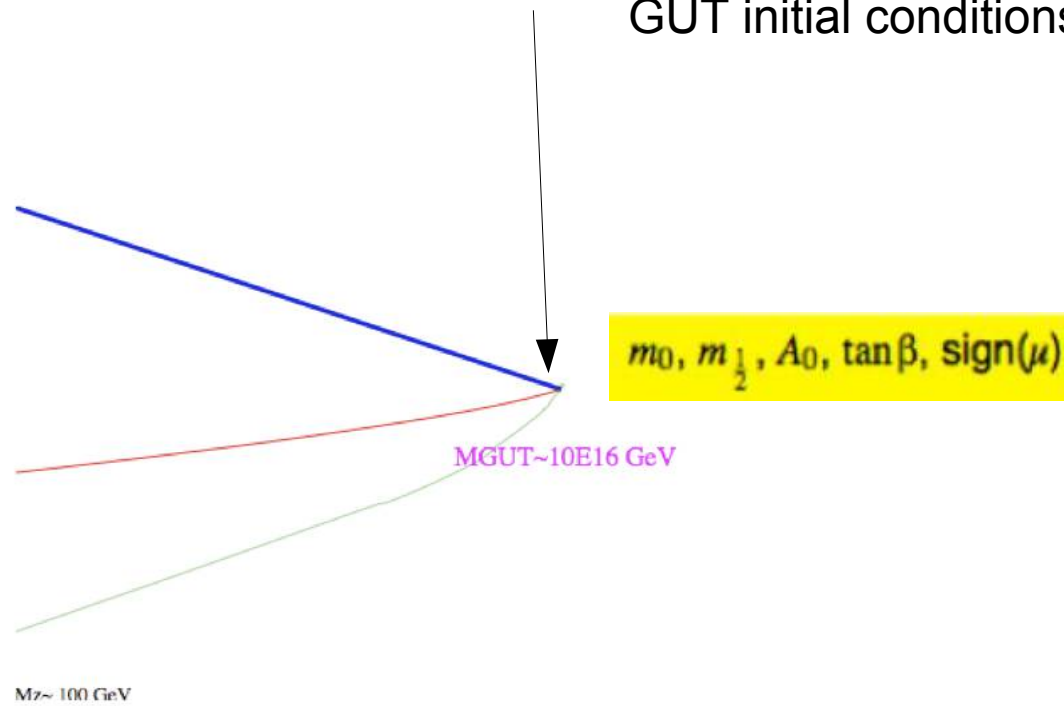
Inspired from supergravity assume universal soft breaking, $\mathcal{L}_{\text{soft}}$:

$$\sum_{f,H} m_0^2 \tilde{f} \tilde{f} + \sum_{\lambda} m_{\frac{1}{2}} \lambda \lambda + \sum_f A_0 Y_f \tilde{f} \tilde{F} H_f + B\mu H_u H_d$$

$$m_0, m_{\frac{1}{2}}, A_0, \tan\beta, \text{sign}(\mu)$$

μ and A_0 can be complex, however their phases constraint to be $< 0,2$ rad by the bounds on the fermion EDM.

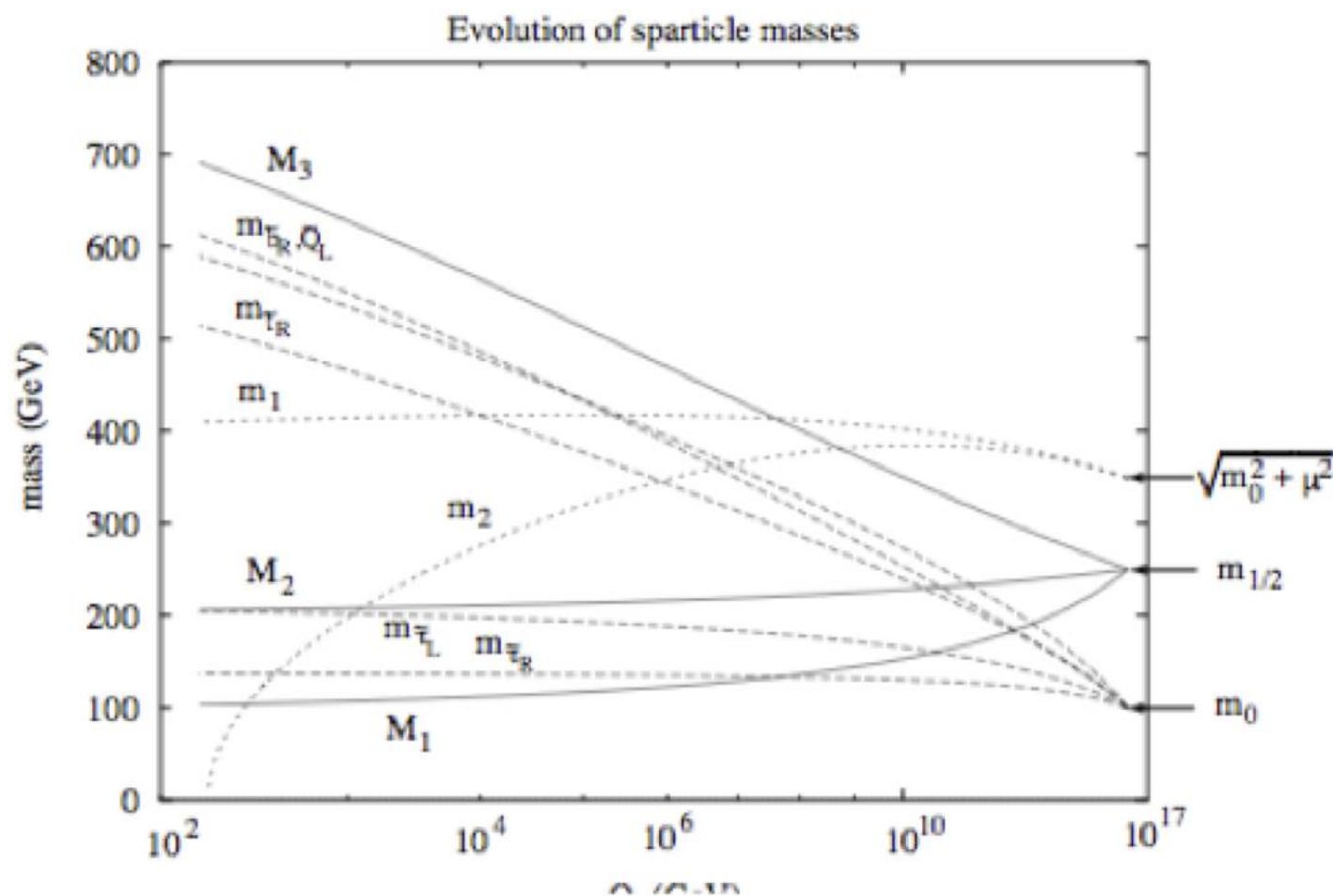
GUT initial conditions



SUSY spectrum

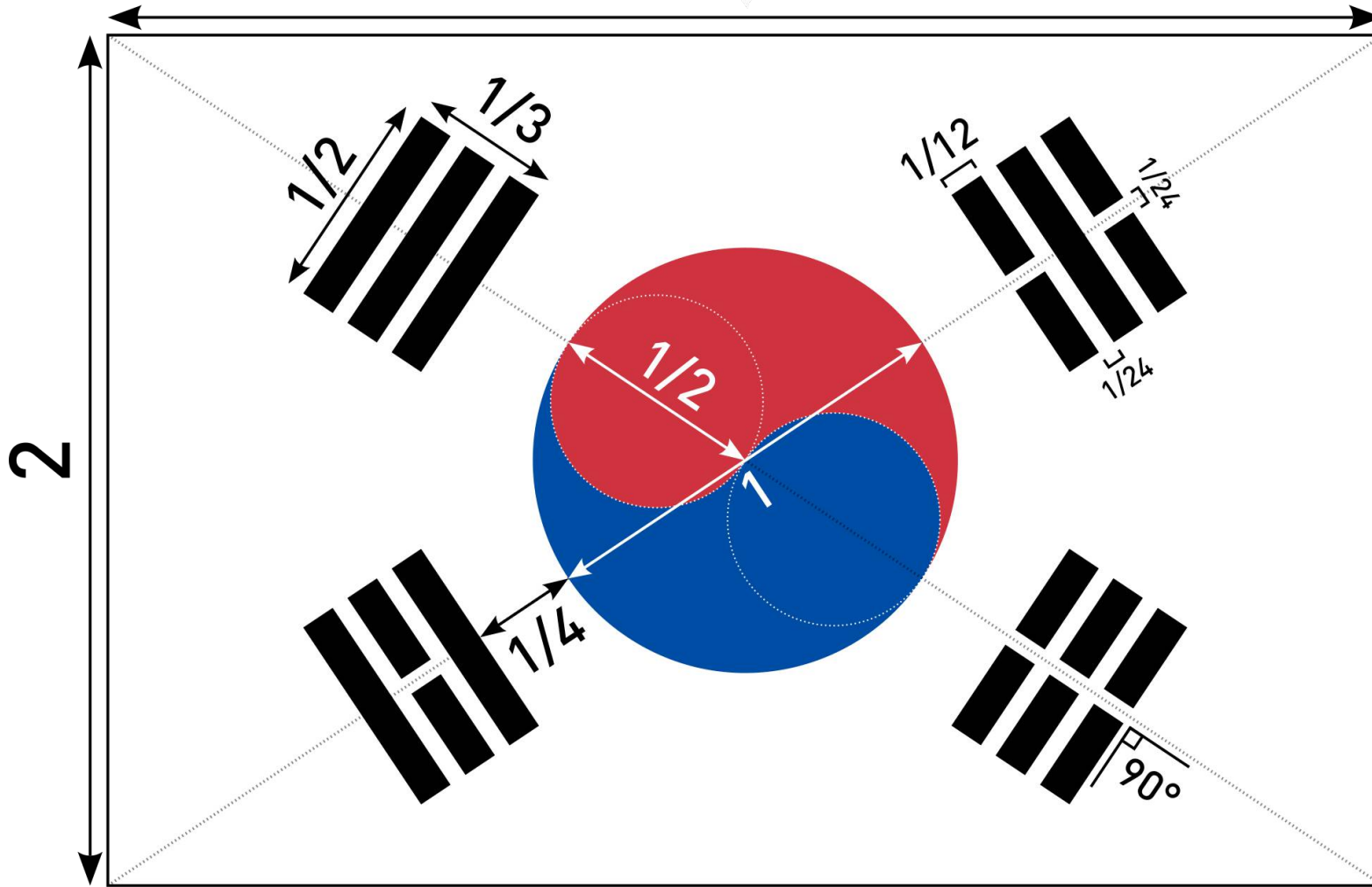
CMSSM, mSUGRA. Parametros de masa universales:

$m_0, M_{1/2}, A_0, \mu_0, \alpha_G, M_{GUT}, \tan\beta$.



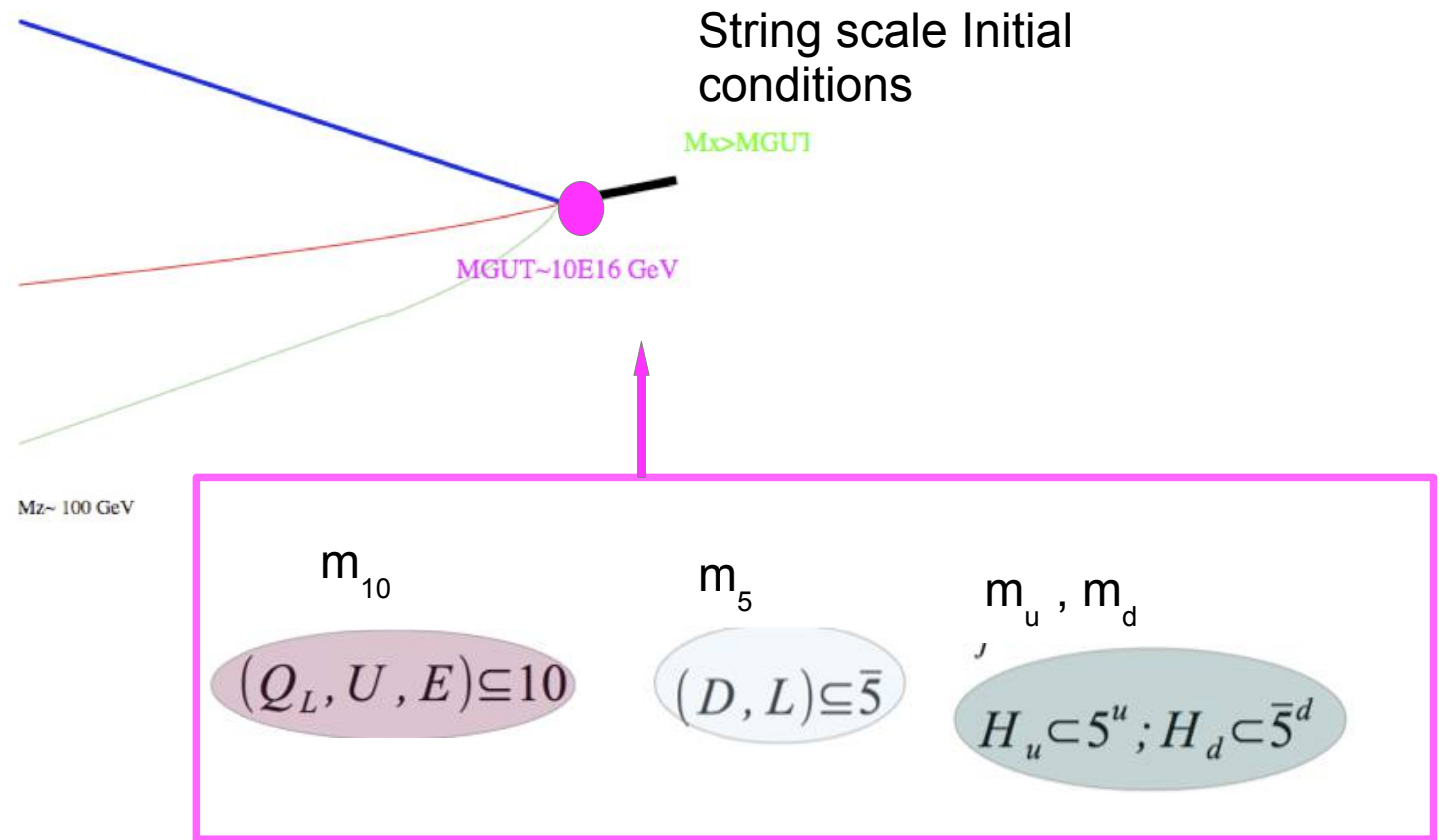
Symmetry

3



$SU(3) \times SU(2) \times U(1) \leftarrow GUT$

$$W_{SU(5)} = Y_u^{ij} 10_i 10_j 5^u + Y_d^{ij} 10_i \bar{5}_j \bar{5}^d$$



Non Universal scenarios

CMSSM choice:

- m_0 *Universal soft masses.*
- $m_{1/2}$ *Universal gaugino masses.*
- A_0 *Universal Trilinear terms.*

Representation-dependent choice

$$m_r = x_r m_0$$

$$A_r = Y_r A_0, \quad A_0 = a_0 m_0$$

Non Universal SU(5)

$$W_{SU(5)} = Y_u^{ij} 10_i 10_j 5^u + Y_d^{ij} 10_i \bar{5}_j \bar{5}^d$$

$$(Q_L, U, E) \subseteq 10$$

$$H_u \subset 5^u; H_d \subset \bar{5}^d$$

$$(D, L) \subseteq \bar{5}$$

The soft terms are taken at GUT as:

$$m_{10} = m_0;$$

$$m_5 = x_5 \cdot m_{10};$$

$$m_u = x_u \cdot m_{10};$$

$$m_d = x_d \cdot m_{10}.$$

$$A_{10,5} = a_0 \cdot m_0,$$

Okada, Shafi, Raza

Phys.Rev. D90 (2014)

Flipped SU(5)

$$\begin{array}{lcl}
 \text{SU(5)} & (Q, u^c, e^c)_i \in \mathbf{10}_i, (L, d^c)_i \in \bar{\mathbf{5}}_i, \nu_i^c \in \mathbf{1}_i. & \\
 & \downarrow \quad \quad \quad \uparrow \quad \quad \quad \swarrow & \\
 \text{Flipped SU(5)} & (Q, d^c, \nu^c)_i \in \mathbf{10}_i, (L, u^c)_i \in \bar{\mathbf{5}}_i, e_i^c \in \mathbf{1}_i. & \\
 & m_{10} = m_0, \quad m_5 = x_5 \cdot m_{10} \quad m_R = x_R \cdot m_{10} & \\
 & m_{H_u} = x_u \cdot m_{10} \quad m_{H_d} = x_d \cdot m_{10} ; &
 \end{array}$$

PATI-SALAM Unification

$$G_{PS} \equiv SU(4) \times SU(2)_L \times SU(2)_R$$

$$4_c 2_L 2_R$$

MATTER FIELDS

$$F_r \quad \begin{pmatrix} d_r & -u_r \\ e_r & -\nu_r \end{pmatrix} \quad (4, 2, 1)$$

$$F_r^c \quad \begin{pmatrix} u_r^c \\ d_r^c \end{pmatrix}, \begin{pmatrix} \nu_r^c \\ e_r^c \end{pmatrix} \quad (\bar{4}, 1, 2)$$

$$\langle \tilde{\nu}_H^c \rangle = \langle \tilde{\nu}_H^c \rangle \sim M$$

Non universal Higgs Mass terms due to D- terms.

$$m_{H_{u,d}}^2 = m_{10}^2 \mp 2M_D^2$$

HIGGS FIELDS

$$H^c \quad \begin{pmatrix} u_H^c \\ d_H^c \end{pmatrix}, \begin{pmatrix} \nu_H^c \\ e_H^c \end{pmatrix} \quad (\bar{4}, 1, 2)$$

$$\bar{H}^c \quad \begin{pmatrix} \bar{u}_H^c & \bar{d}_H^c \\ \bar{\nu}_H^c & \bar{e}_H^c \end{pmatrix} \quad (4, 1, 2)$$

$$h \quad \begin{pmatrix} h_2^+ & h_1^0 \\ h_2^0 & h_1^- \end{pmatrix} \quad (1, 2, 2)$$

$$G_{PS} \rightarrow SU(3)_C \otimes SU(2)_L \otimes U(1)_Y$$

$$M_1 = \frac{3}{5}M_2 + \frac{2}{5}M_3.$$

Condition for gaugino masses.

PS(4-2-1) *LR Asymmetry*

MATTER FIELDS

$$F_r \quad \begin{pmatrix} d_r & -u_r \\ e_r & -\nu_r \end{pmatrix} \quad (\mathbf{4}, \mathbf{2}, \mathbf{1})$$

$$F_r^c \quad \begin{pmatrix} u_r^c \\ d_r^c \end{pmatrix}, \quad \begin{pmatrix} \nu_r^c \\ e_r^c \end{pmatrix} \quad (\bar{\mathbf{4}}, \mathbf{1}, \mathbf{2})$$

$$\frac{m_L}{m_R}$$

New Parameter

$$x_{LR} = \frac{m_L}{m_R},$$

- SUSY SEARCH: SuperBayeS, MultiNest
- RGE's: SoftSusy
- Relic Density: MicroOMEGAs
- Direct DM detection: DarkSUSY
- Super Iso: $\delta a_{\mu}^{\text{SUSY}}$
- SusyBSG: B-Physics.

Likelihood function:

$$\begin{aligned}\ln \mathcal{L}_{\text{Joint}} = & \ln \mathcal{L}_{\text{EW}} + \ln \mathcal{L}_{\text{B}} + \ln \mathcal{L}_{\Omega_{\chi} h^2} \\ & + \ln \mathcal{L}_{\text{LUX}} + \ln \mathcal{L}_{\text{Higgs}} + \ln \mathcal{L}_{\text{SUSY}} + \ln \mathcal{L}_{\text{g-2}},\end{aligned}$$

Higgsino χ_1^0 :

$$h_f \equiv |N_{13}|^2 + |N_{14}|^2, \quad h_f > 0.1, \quad |m_A - 2m_\chi| > 0.1 m_\chi.$$

A/H resonances:

$$|m_A - 2m_\chi| \leq 0.1 m_\chi.$$

$\tilde{\tau}$ coannihilations:

$$h_f < 0.1, \quad (m_{\tilde{\tau}_1} - m_\chi) \leq 0.1 m_\chi$$

$\tilde{\tau} - \tilde{\nu}_\tau$ coannihilations:

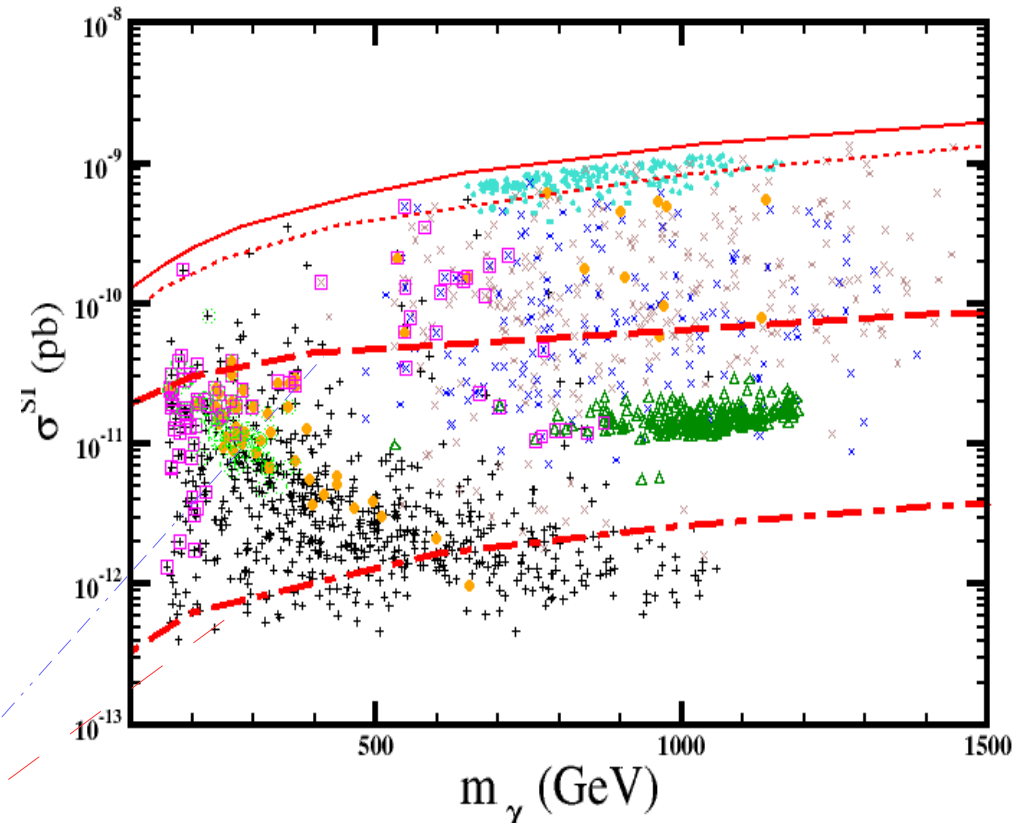
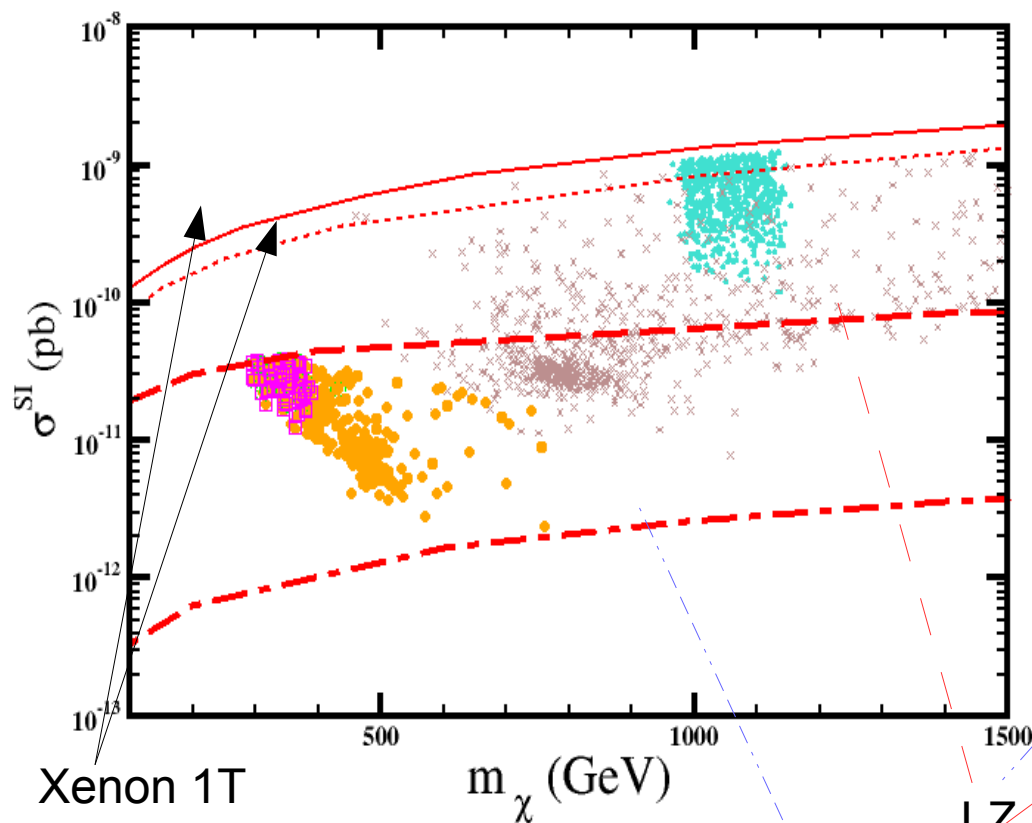
$$h_f < 0.1, \quad (m_{\tilde{\tau}_1} - m_\chi) \leq 0.1 m_\chi, \quad (m_{\tilde{\mu}_1} - m_\chi) \leq 0.1 m_\chi.$$

$\tilde{t}_1$ coannihilations:

$$h_f < 0.15, \quad (m_{\tilde{t}_1} - m_\chi) \leq 0.1 m_\chi.$$

Flipped-SU(5)

PS(4-2-1)



-higgsino DM



A/H Resonances



Coannihilations:



$\tilde{\tau}$

$+$

$\tilde{\chi}^+$



$\tilde{t}$



$\tilde{g}$

$\tilde{\tau} - \tilde{\nu}$



$\tilde{b}$



LHC excluded.



good (g-2)

MSSM extended by seesaw mechanism

- The superpotential for MSSM-Seesaw I can be written as

$$W = W_{\text{MSSM}} + Y_{\nu}^{ij} \epsilon_{\alpha\beta} H_2^{\alpha} N_i^c L_j^{\beta} + \frac{1}{2} M_N^{ij} N_i^c N_j^c, \quad (5)$$

- The full set of soft SUSY-breaking terms is given by,

$$\begin{aligned} -\mathcal{L}_{\text{soft,SI}} = & -\mathcal{L}_{\text{soft}} + (m_{\tilde{\nu}}^2)_j^i \tilde{\nu}_{Ri}^* \tilde{\nu}_R^j + \left(\frac{1}{2} B_{\nu}^{ij} M_N^{ij} \tilde{\nu}_{Ri}^* \tilde{\nu}_{Rj}^* \right. \\ & \left. + A_{\nu}^{ij} h_2 \tilde{\nu}_{Ri}^* \tilde{l}_{Lj} + \text{h.c.} \right), \end{aligned} \quad (6)$$

With W_{MSSM} and $\mathcal{L}_{\text{soft}}$ given in Eq. 3 and 4 respectively.

$$\mathcal{M} = \begin{pmatrix} 0 & m_\nu^D \\ m_\nu^{D^T} & M_R \end{pmatrix} \quad \text{“See-Saw” explanation for tiny masses.}$$

• The physical masses are:

1. $m_1 \equiv m_{light} \simeq \frac{(m_\nu^D)^2}{M_R}$
2. $m_2 \simeq M_R$

• For $(m_\nu^D)_{33} \approx (200 \text{ GeV})$ ($\lambda_\nu \approx \lambda_t$) and $M_{N_3} \approx O(10^{14} \text{ GeV})$, $m_{eff} \approx 0.05 \text{ eV}$

$$W = W_{\text{MSSM}} + \frac{1}{2} (Y_\nu L H_2)^T M_N^{-1} (Y_\nu L H_2).$$

$$m_{\text{eff}} = -\frac{1}{2} v_u^2 Y_\nu \cdot M_N^{-1} \cdot Y_\nu^T, \quad m_\nu^\delta = U^T m_{\text{eff}} U$$

Slepton flavor mixings

$$(m_{\tilde{L}}^2)_{ij} \sim \frac{1}{16\pi^2} (6m_0^2 + 2A_0^2) (Y_\nu^\dagger Y_\nu)_{ij} \log \left(\frac{M_{\text{GUT}}}{M_N} \right)$$

$$(m_{\tilde{e}}^2)_{ij} \sim 0$$

$$(A_l)_{ij} \sim \frac{3}{8\pi^2} A_0 Y_{li} (Y_\nu^\dagger Y_\nu)_{ij} \log \left(\frac{M_{\text{GUT}}}{M_N} \right)$$

Orthogonal
matrix

$$Y_\nu = \frac{\sqrt{2}}{v_u} \sqrt{M_R^\delta} R \sqrt{m_\nu^\delta} U^\dagger$$

Casas + Ibarra

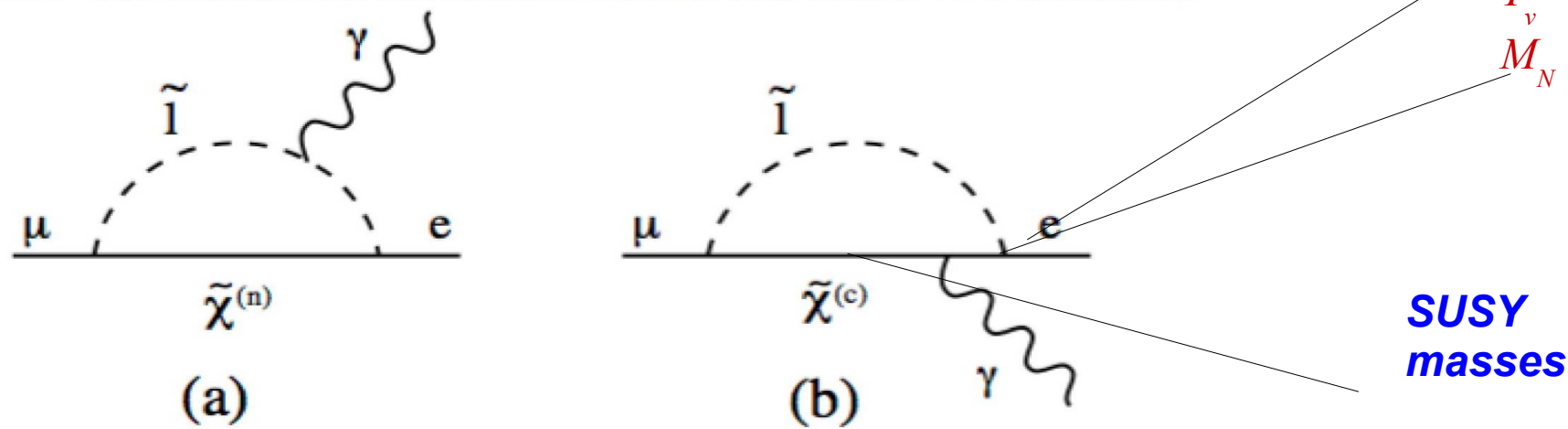
Diagonal Universal
1E14 GeV

Order 1

$$Y_\nu^\dagger Y_\nu = \frac{2}{v_u^2} M_R U m_\nu^\delta U^\dagger$$

Limit case of
degenerate MR

In SUSY flavor mixing lepton-slepton vertices can induce LFV diagrams:



Lepton-slepton flavor mixing is very constrained by the experimental limits:

$$BR(\mu \rightarrow e \gamma) < 4.2 \times 10^{-13}$$

$$BR(\tau \rightarrow \mu \gamma) < 4.4 \times 10^{-8}$$

$$BR(\tau \rightarrow e \gamma) < 3.3 \times 10^{-8}$$

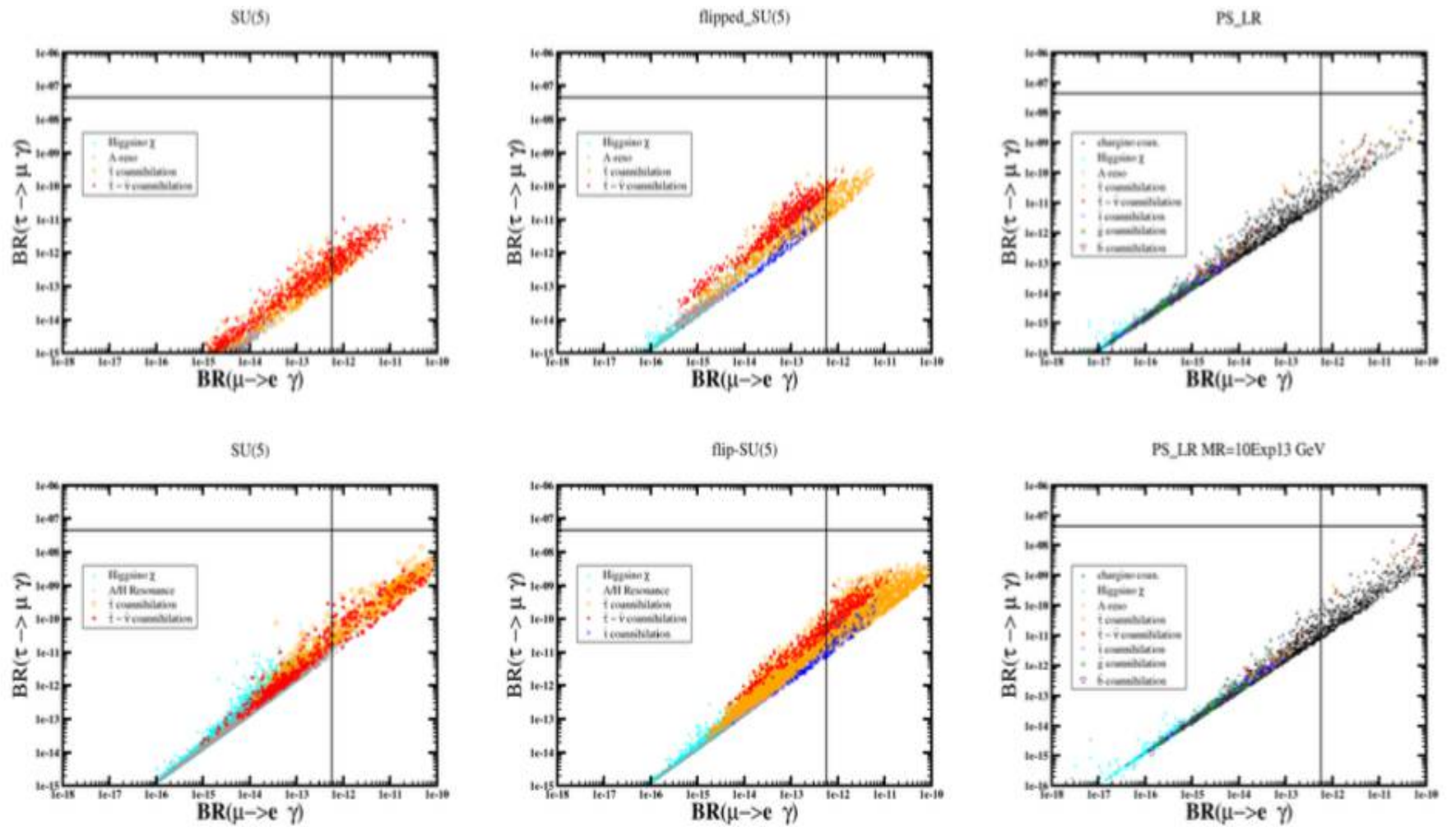
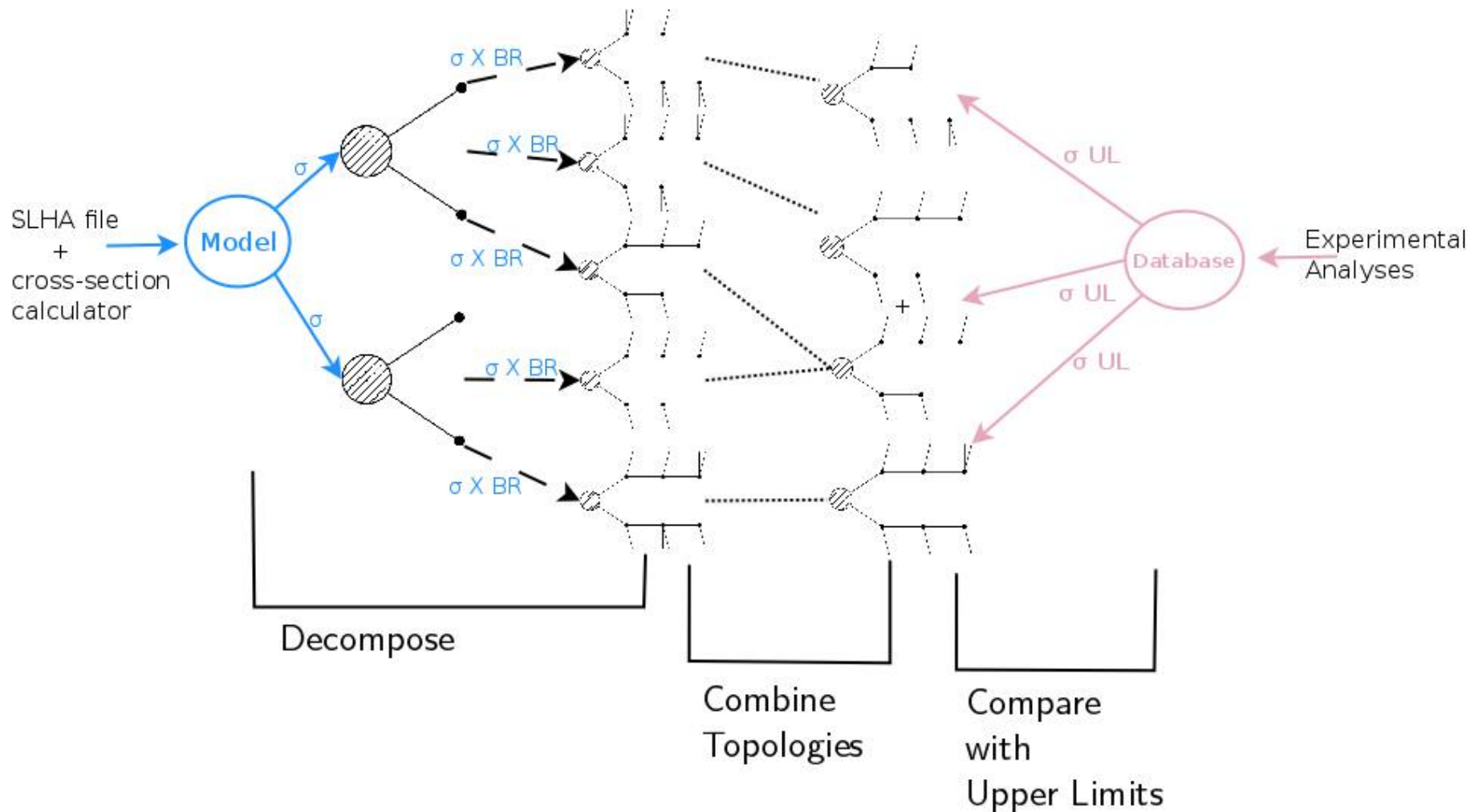


Figure 1: Prediction for $\text{BR}(\tau \rightarrow \mu\gamma)$ vs $\text{BR}(\mu \rightarrow e\gamma)$ for the SU(5), the f-SU(5) and the PS cases. The upper (lower) panels use $M_N = 2.5 \cdot 10^{12}$ ($M_N = 10^{13}$ GeV). The symbols correspond to the classes of models representing the DM scenarios described in the text: orange circles $\tilde{t} - \chi$ coannihilations ...



Analyzed points:



Excluded

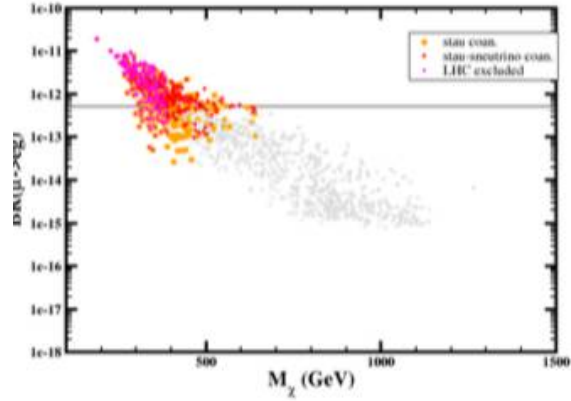


Allowed

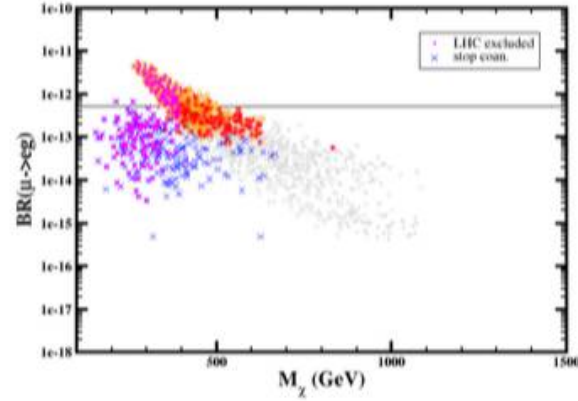


Not analyzed.

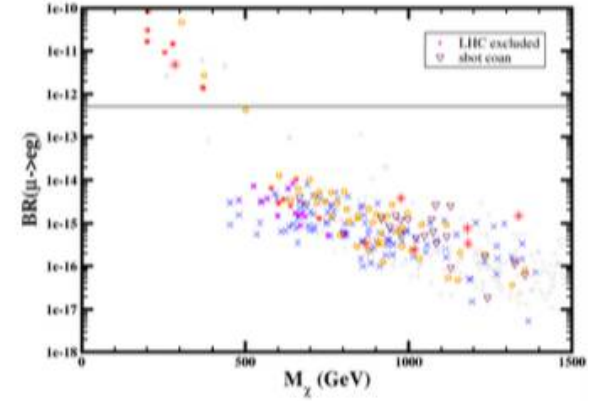
CRR MR 2.5 E12 SU5



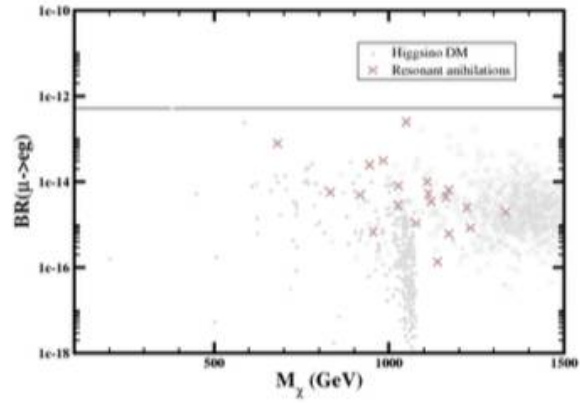
CRR MR 2.5 E12 FSU5



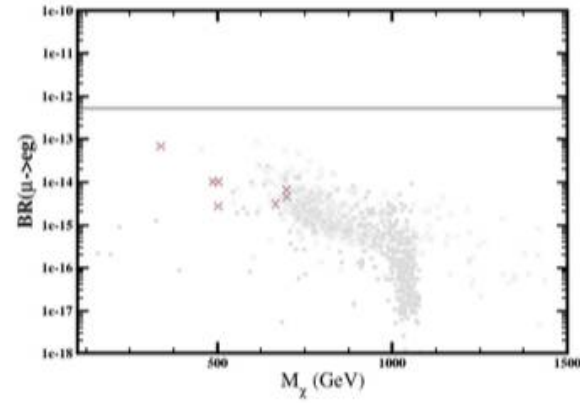
4-2-2 CSTOP MR 2.5 E12



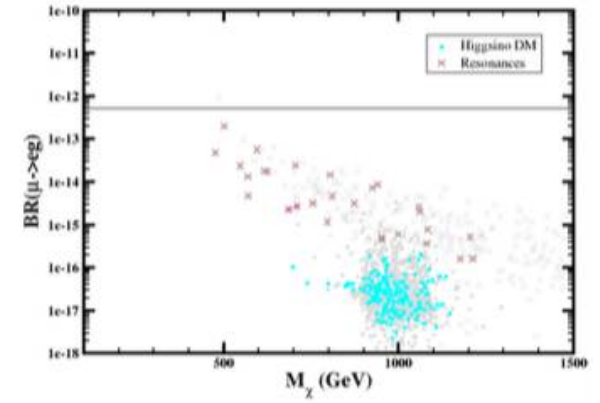
Higgsino-DM MR 2.5 E12 SU(5)

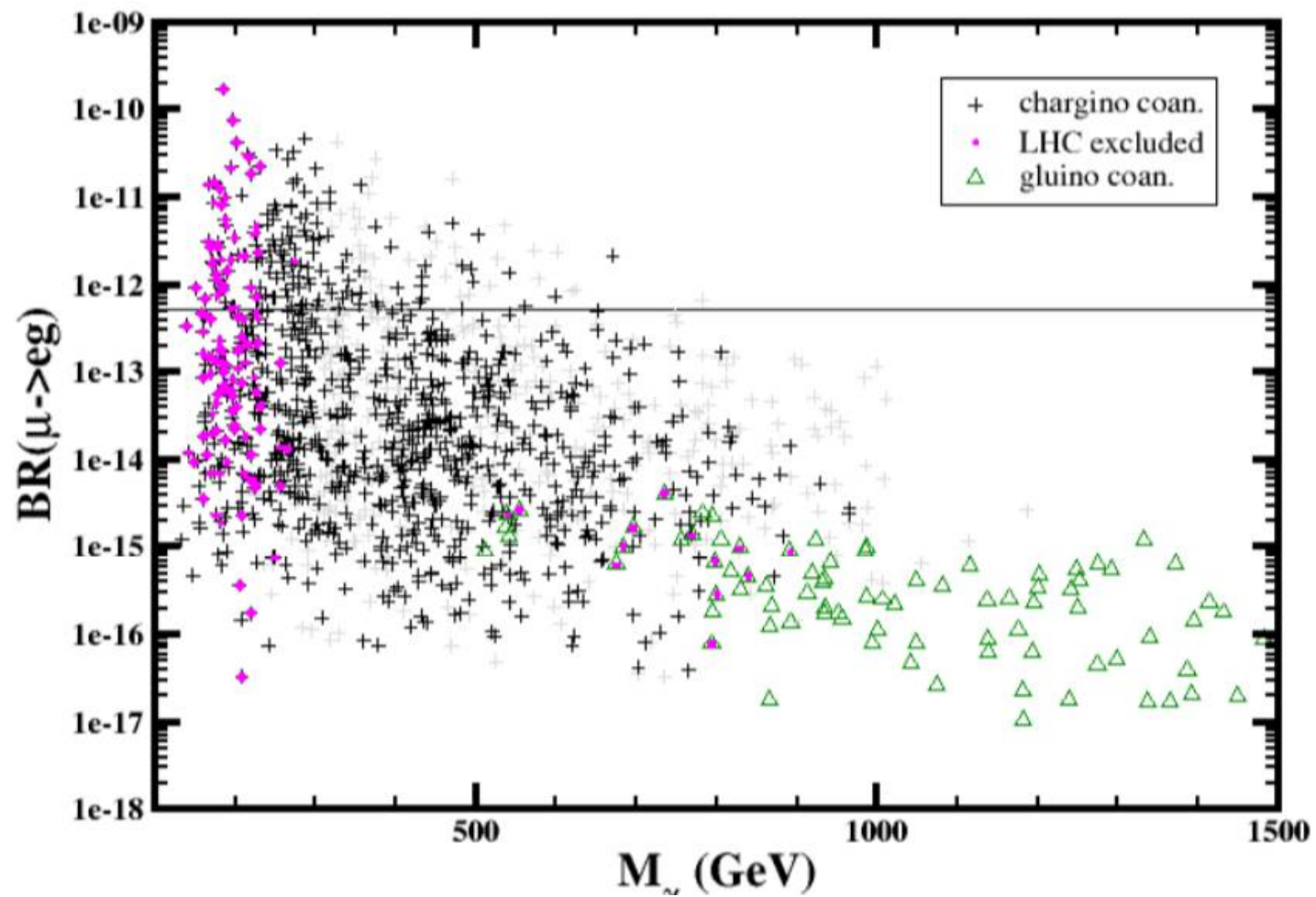


Higgsino DM MR=2.5 E12 FSU5



CRR MR 2.5 E12 PS



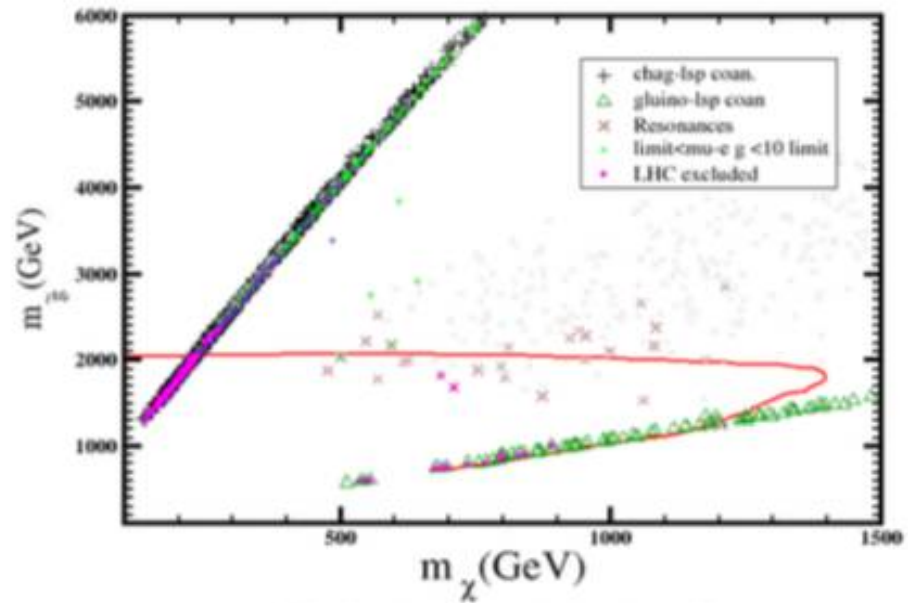
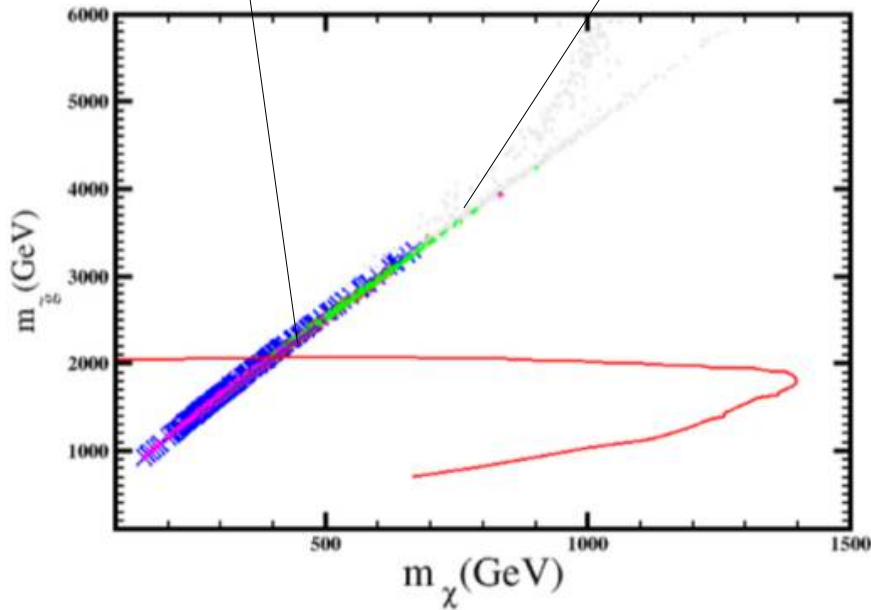


LFV excluded (indigo)

LFV prospect (green)

FSU5 all areas

PS Gauino and resonance areas



A/H Resonances



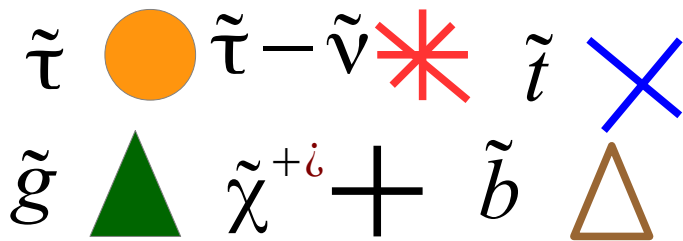
Higgsino DM



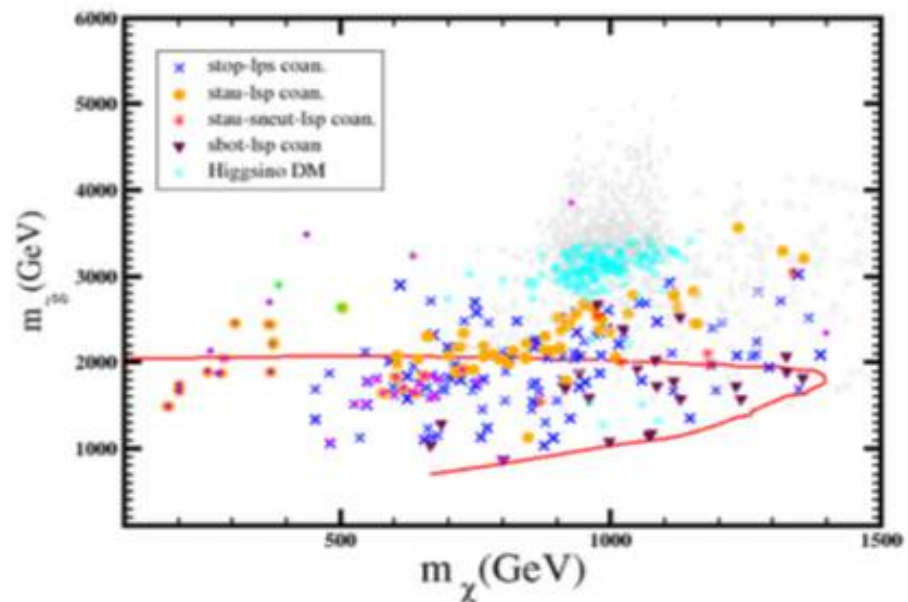
LHC excluded



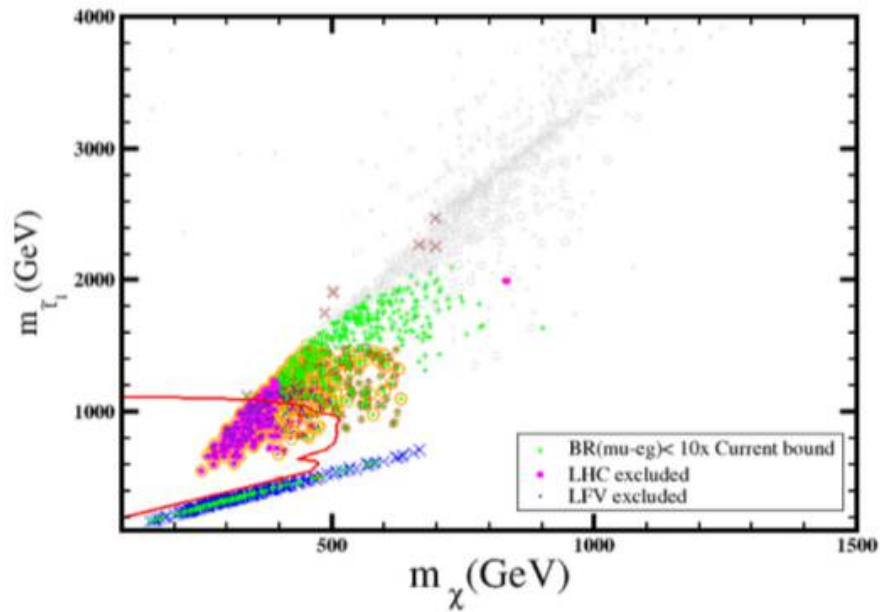
Coanihilations:



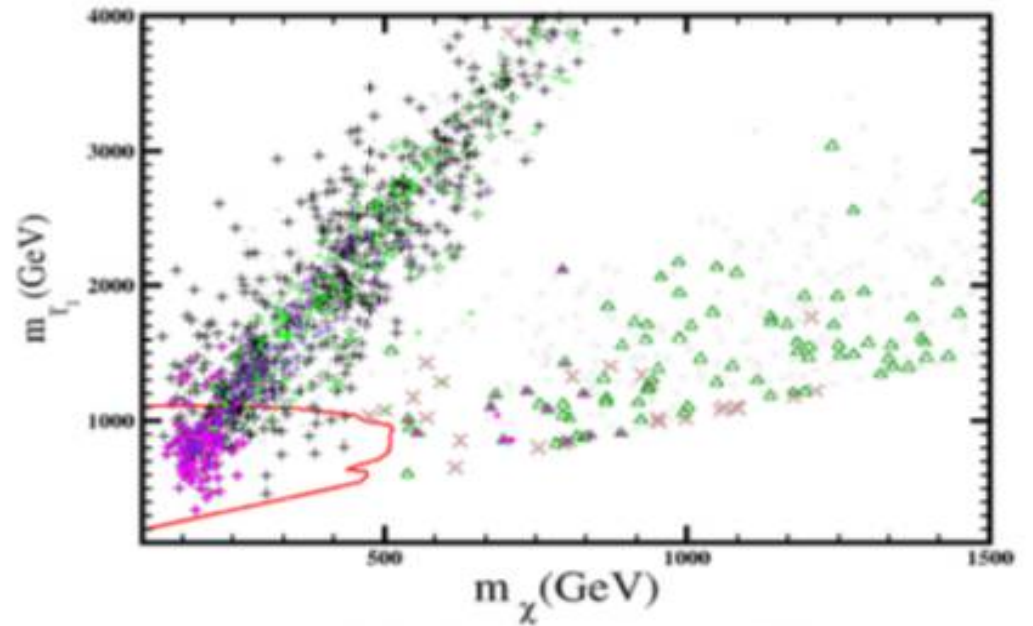
PS sfermions coan +Higgsino DM



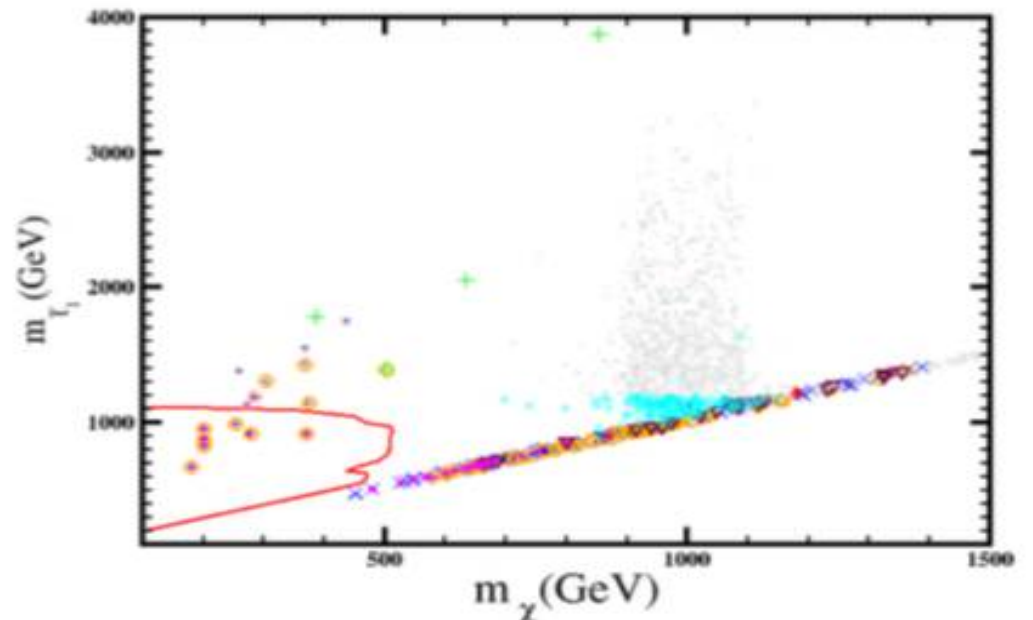
FSU5 All regions



PS-Gauginos+resonances



PS sfer-Lsp coan. + Higgsino DM



A/H Resonances



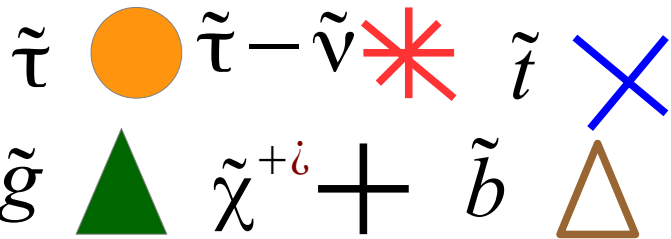
Higgsino DM



LHC excluded



Coannihilations:



CONCLUSIONS

We have identified different patterns of soft SUSY-breaking terms at the GUT scale, with different dark matter predictions and the constraints from LHC searches.

We have calculated the SUSY spectra for the different gauge groups, finding that the models predict different spectra for the same LSP mass, connecting possible future observations with the structure of the underlying unified theory.

The see-saw mechanism to explain neutrino data may imply cLFV. The $\mu \rightarrow e \gamma$ prediction may be complementary (or alternative) signal to SUSY spectroscopy at the LHC.

We have found that SU(5) and FSU(5) GUT's (gaugino universality), PS (gaugino non universality) lead to very different predictions for dark matter LHC and cLFV experiments, and thus are distinguishable in future searches. PS predicts stop-LSP, gluino-LSP coannihilations that are absent in the other groups and can be explored by LHC searches.

The LHC searches for generic missing ET, charginos and stops are quite complementary, and future LHC runs will be able to constrain the models in several different ways.

Models with compressed SUSY spectrum are difficult to test, however these are the natural scenarios with a suitable DM relic density. Furthermore, some of these scenarios also favor the detection of cLFV signals in the current experiments. The most promising one will be the chargino-LSP coannihilation class of models.