UNITARITY VIOLATION FROM NONSTANDARD HIGGS COUPLINGS



Spencer Chang (U. Oregon) w/ Markus Luty 1902.05556+ongoing also see Falkowski & Rattazzi 1902.05936 and earlier work by Belyaev et.al. 1212.3860 IBS CTPU Workshop 12/12/19

PINNING DOWN HIGGS PROPERTIES



Post-discovery, major goal of LHC and future colliders is measuring Higgs properties so we can test EWSB mechanism, mass generation, but also to look for new physics beyond the Standard Model

\sqrt{s} = 13 TeV, 24.5 - 79.8 fb⁻¹ **ATLAS** $m_H = 125.09 \,\text{GeV}, |y_{\perp}| < 2.5$ 68% CL: 95% CL: κ_{z} κ_{W} κ_t κ_b K_{τ} κ_g κ_{γ} B_{inv} B_{undet} B_{BSM} 0.5 -0.51.5 CERN-EP-20 | 9-097 Parameter value Standard Model values

HIGGS COUPLINGS MEASUREMENTS

Fits to

σ x Branching Ratios,
for Higgs couplings
have 10-25%
errors and currently
agree with SM value

HIGGS COUPLINGS IN FUTURE

| kappa-0 | HL-LHC | LHeC | HE- | -LHC | 1 | ILC | | l | CLIC | | CEPC | FCC | C-ee | FCC-ee/eh/hh |
|------------------------|--------|-------|-----|------|------|------|------|------|-------|------|------|-------------|-------------|--------------|
| марра о | | 21100 | S2 | S2' | 250 | | 1000 | 380 | 15000 | 3000 | | 240 | | |
| κ _W [%] | 1.7 | 0.75 | 1.4 | 0.98 | 1.8 | 0.29 | 0.24 | 0.86 | 0.16 | 0.11 | 1.3 | 1.3 | 0.43 | 0.14 |
| κ_{Z} [%] | 1.5 | 1.2 | 1.3 | 0.9 | 0.29 | 0.23 | 0.22 | 0.5 | 0.26 | 0.23 | 0.14 | 0.20 | 0.17 | 0.12 |
| κ_g [%] | 2.3 | 3.6 | 1.9 | 1.2 | 2.3 | 0.97 | 0.66 | 2.5 | 1.3 | 0.9 | 1.5 | 1.7 | 1.0 | 0.49 |
| κ _γ [%] | 1.9 | 7.6 | 1.6 | 1.2 | 6.7 | 3.4 | 1.9 | 98∗ | 5.0 | 2.2 | 3.7 | 4.7 | 3.9 | 0.29 |
| $\kappa_{Z\gamma}$ [%] | 10. | 1-1 | 5.7 | 3.8 | 99∗ | 86× | 85∗ | 120∗ | 15 | 6.9 | 8.2 | 81 ★ | 75 ★ | 0.69 |
| κ_c [%] | _ | 4.1 | - | - | 2.5 | 1.3 | 0.9 | 4.3 | 1.8 | 1.4 | 2.2 | 1.8 | 1.3 | 0.95 |
| κ_t [%] | 3.3 | 1-1 | 2.8 | 1.7 | - | 6.9 | 1.6 | _ | _ | 2.7 | _ | _ | _ | 1.0 |
| κ_b [%] | 3.6 | 2.1 | 3.2 | 2.3 | 1.8 | 0.58 | 0.48 | 1.9 | 0.46 | 0.37 | 1.2 | 1.3 | 0.67 | 0.43 |
| κ_{μ} [%] | 4.6 | 1-1 | 2.5 | 1.7 | 15 | 9.4 | 6.2 | 320∗ | 13 | 5.8 | 8.9 | 10 | 8.9 | 0.41 |
| κ_{τ} [%] | 1.9 | 3.3 | 1.5 | 1.1 | 1.9 | 0.70 | 0.57 | 3.0 | 1.3 | 0.88 | 1.3 | 1.4 | 0.73 | 0.44 |

Taken from Higgs@FutureColliders report (1905.03764)

HIGHER ORDER HIGGS COUPLINGS

HL-LHC and future colliders will be sensitive to higher order Higgs couplings, including self-interactions and quadratic couplings to W, top

$$(1+\delta_3)\frac{m_h^2}{2v}h^3 + (1+\delta_4)\frac{m_h^2}{8v^2}h^4$$

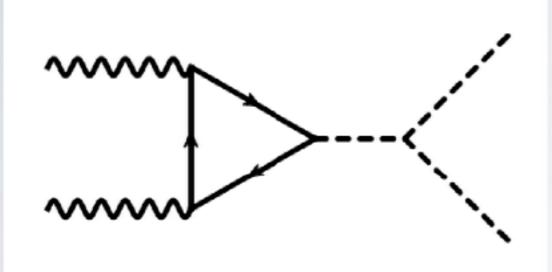
• • •

$$\left(\frac{1}{2}m_Z^2Z^2 + m_W^2W^2\right)\left(1 + \delta_{hhVV}\right)\frac{h^2}{v^2}$$

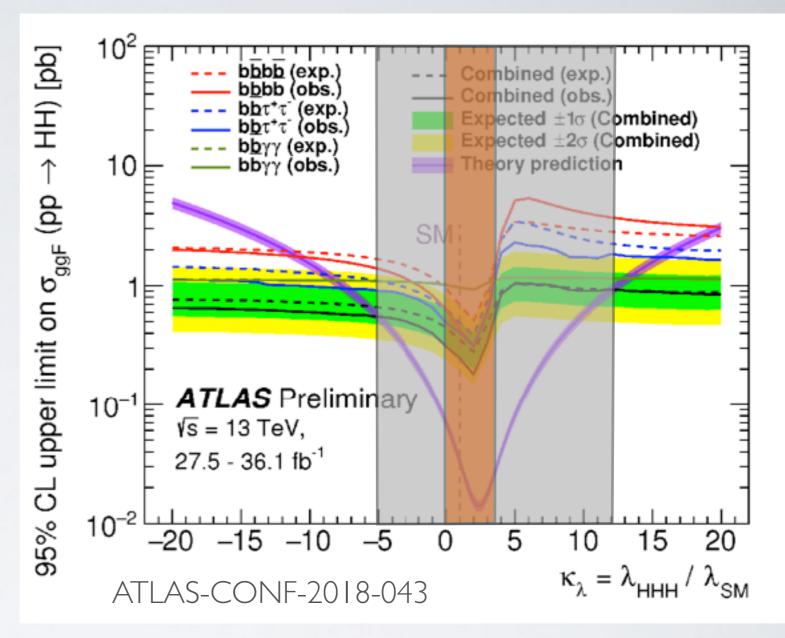
. . .

$$-\frac{m_t}{2v^2}c_2h^2\bar{T}T$$

TRILINEAR SEARCH



Trilinear probed by search for Double Higgs production



Currently only sensitive to O(10) variations, but projections estimate trilinear sensitivity to \sim [-0.2,3.6] at LHC w/ 3 ab⁻¹ and 20-30% at future colliders

TRIPLE HIGGS PROCESS

Papaefstathiou and Sakurai See also Chien et.al.

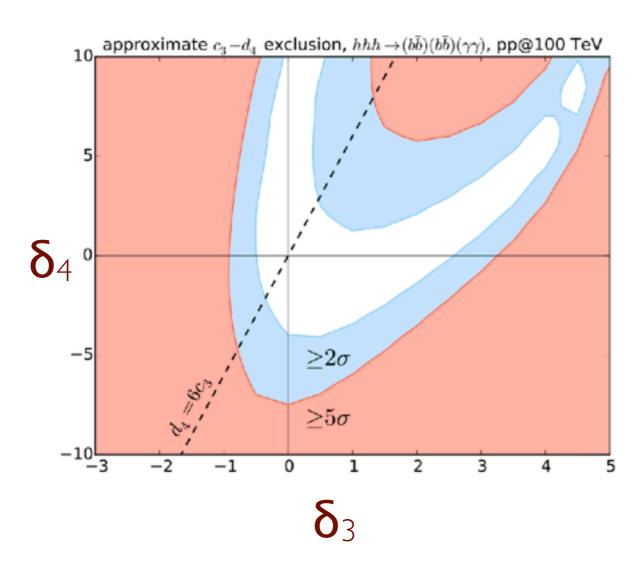
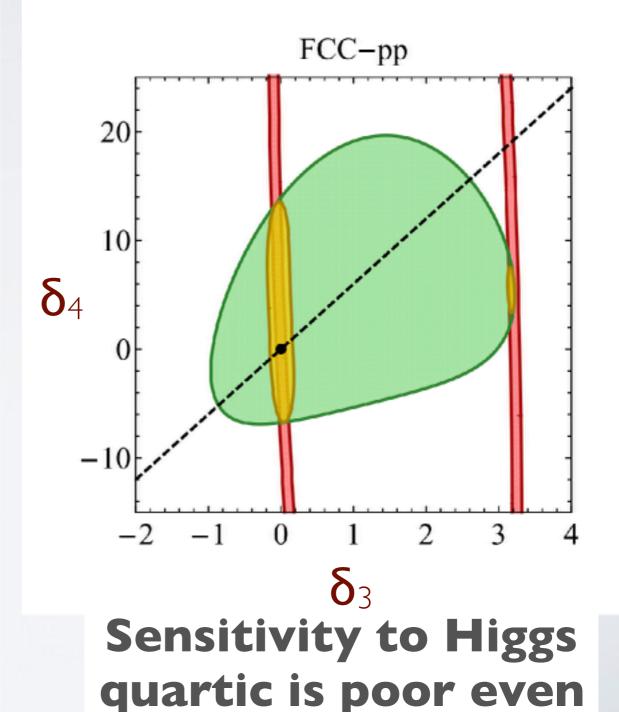


FIG. 6: The approximate expected 2σ (blue) and 5σ (red) exclusion regions on the c_3-d_4 plane after 30 ab⁻¹ of integrated luminosity, derived assuming a constant signal efficiency, calculated along the $d_4 = 6c_3$ line in $c_3 \in [-3.0, 4.0]$.

hh and hhh at one loop e.g. Bizon et.al.



in optimistic cases

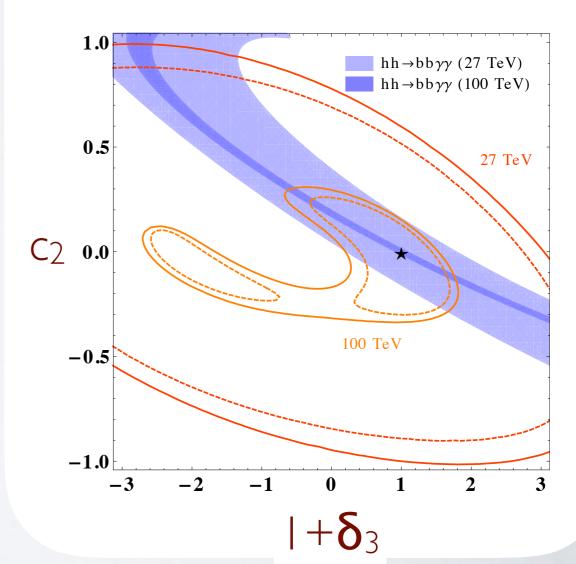
W/Z AND TOP COUPLINGS TO HH

VVhh measured in VBF DiHiggs to 4b's Bishara et.al. (1611.03860)

| | 68% probability interval on | | | | |
|--------------------|-----------------------------|-----------------------------|--|--|--|
| | $1 \times \sigma_{\rm bkg}$ | $3 \times \sigma_{\rm bkg}$ | | | |
| LHC ₁₄ | [-0.37, 0.45] | [-0.43, 0.48] | | | |
| HL-LHC | [-0.15, 0.19] | [-0.18, 0.20] | | | |
| FCC ₁₀₀ | [0, 0.01] | [-0.01, 0.01] | | | |

Sensitivity to O(.I-I) for quadratic Higgs couplings

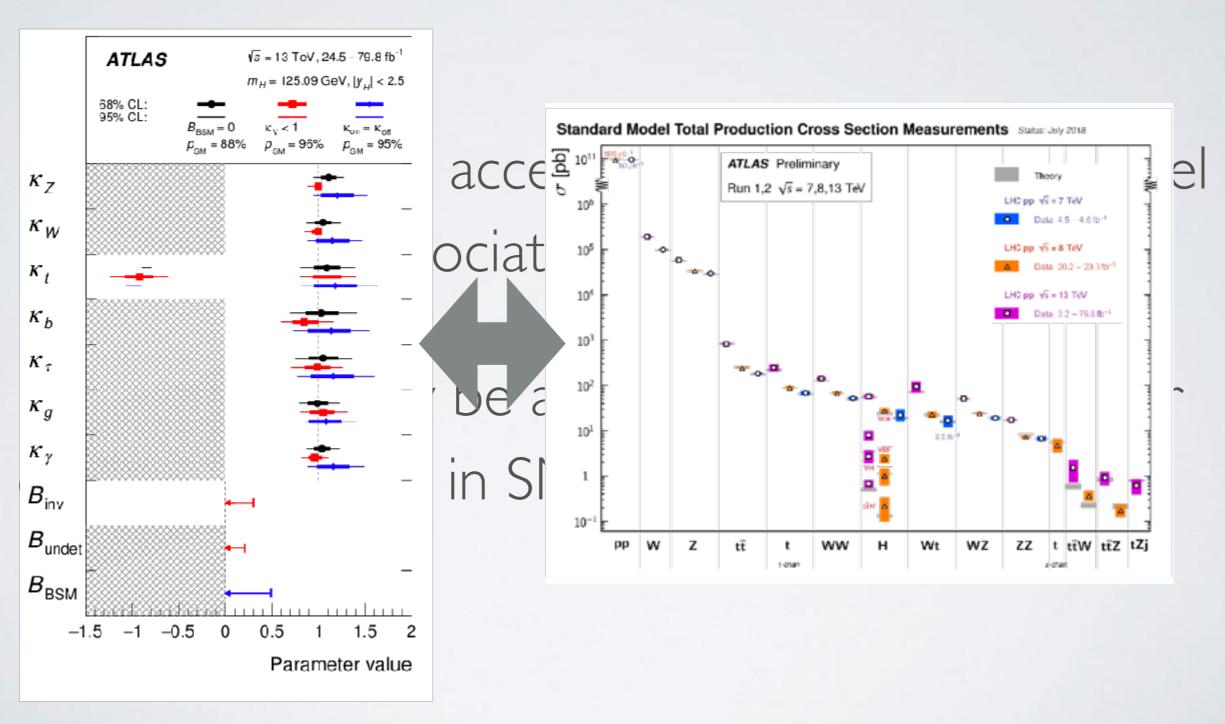
tthh coupling probed by tthh production
Li et.al. (1905.03772)



SUMMARY OF HIGGS COUPLING PROSPECTS

- Currently many Higgs linear couplings constrained to 10-25% level, improved to few % in future
- Higher order couplings: VVhh more promising by almost order of magnitude than h³, tthh. O(I) at HL-LHC, O(I0%) at future colliders

What do we do if we find a significant deviation from the SM prediction?



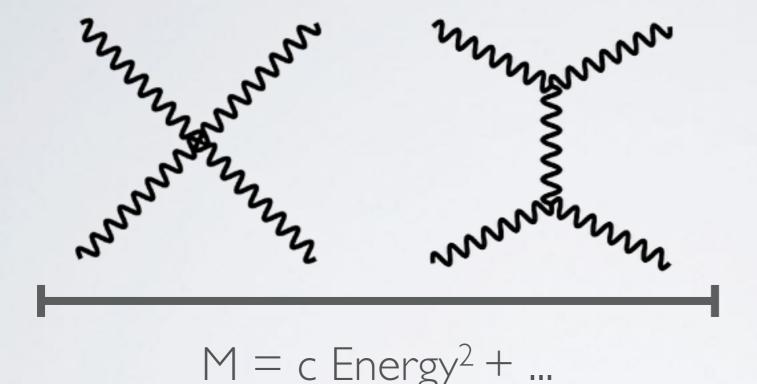
NEW PHYSICS SCALE BOUND FROM UNITARITY VIOLATION (W/ LUTY)

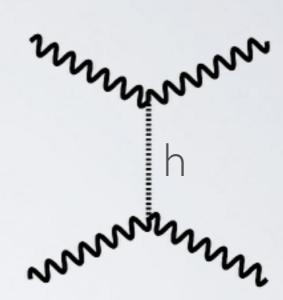


The Standard Model is a precise deck of cards, modifications (due to higher dimensional operators) lead to problems at high energies.

Unitarity violating processes give interesting processes to measure and put bound on new physics scale

CLASSIC EXAMPLE SCATTERING $Z_L Z_L \Leftrightarrow W^+_L W^-_L$





$$M = -c Energy^2 + ...$$

Higgs exchange cancels high energy growth if its couplings are SM-like, matrix element is Unitary if m_H ≤ ITeV (Lee, Quigg, Thacker)

GENERAL HIGGS COUPLINGS

Higgs Effective Field Theory (HEFT) parameterizes most general Higgs couplings phenomenologically

$$V = \frac{1}{2}m_h^2 h^2 + \lambda_{hhh} h^3 + \lambda_{hhhh} h^4 + \lambda_{hhhh} h^5 + \cdots$$

$$V \to \frac{1}{2}m_h^2 X^2 + \lambda_{hhh} X^3 + \lambda_{hhhh} X^4 + \lambda_{hhhh} X^5 + \cdots$$

 $SU(2) \times U(1)$ invariant form uses an nonanalytic field

$$X \equiv \sqrt{2|H|^2} - v = \sqrt{(v+h)^2 + \vec{G}^2 - v}$$
$$= h + \frac{1}{2v}\vec{G}^2 - \frac{1}{2v^2}h\vec{G}^2 + \cdots$$

VAND TOP COUPLINGS

Use a nonanalytic Higgs doublet
$$P = \frac{H}{\sqrt{|H|^2}} = \begin{pmatrix} 0 \\ 1 \end{pmatrix} + \begin{pmatrix} \sqrt{2}G^+/v \\ iG^0/v \end{pmatrix} + \cdots$$

$$(m_W^2 W^2 + \frac{1}{2} m_Z^2 Z^2) \left(1 + 2(1 + \delta_{hVV}) \frac{h}{v} + (1 + \delta_{hhVV}) \frac{h^2}{v^2} + c_3 \frac{h^3}{v^3} \right)$$

$$\rightarrow |DP|^2 \left(\delta_{hVV} vX + \frac{1}{2} \delta_{hhVV} X^2 + \frac{c_3}{2v} X^3 \right)$$

$$-m_{t}\overline{T}T\left[1+(1+\delta\kappa_{t})\frac{h}{v}+\frac{1}{2}c_{2}\left(\frac{h}{v}\right)^{2}+\frac{1}{6}c_{3}\left(\frac{h}{v}\right)^{3}\right]$$

$$\rightarrow -m_{t}\overline{T}_{R}P\epsilon\left(\frac{T_{L}}{B_{L}}\right)\left[\delta\kappa_{t}\frac{X}{v}+\frac{1}{2}c_{2}\left(\frac{X}{v}\right)^{2}+\frac{1}{6}c_{3}\left(\frac{X}{v}\right)^{3}+\cdots\right]+h.c.$$

OUR GENERAL UNITARITY VIOLATION APPROACH

 $|P, \alpha \rangle$ Define states of total momentum P w/ other properties α (e.g. # Higgses)

Properly normalized
$$\langle P', \alpha' | P, \alpha \rangle = (2\pi)^4 \delta(P-P') \delta_{\alpha\alpha'}$$

Leads to bounds $|T_{\alpha\alpha'}| \leq 1$

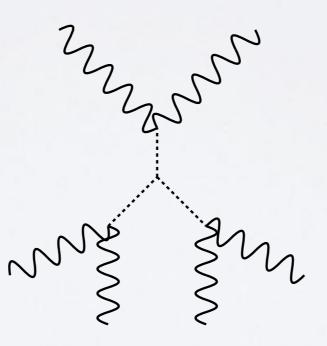
$$\langle P', \alpha' | T | P, \alpha \rangle = (2\pi)^4 \delta(P - P') T_{\alpha \alpha'}$$

Allows us to go beyond 2 to 2 processes and set better bounds

TRILINEAR UNITARITY VIOLATION

Modifying trilinear from SM value automatically leads to Unitarity violation at high energies

win sm



Example: $Z_L Z_L Z_L \Leftrightarrow Z_L Z_L Z_L$

Cancellation to get

M ~ I/Energy²

requires SM

trilinear value!

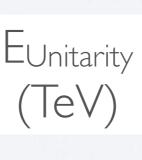
HIGGS TRILINEAR MODIFICATION

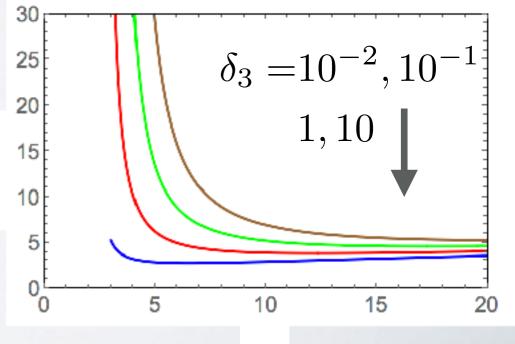
$$\frac{m_h^2}{2v}\delta_3 X^3 = \frac{m_h^2}{2v}\delta_3 \left(\sqrt{(v+h)^2 + \vec{G}^2} - v\right)^3$$

$$\supset \sum_{m} \delta_3(-1)^m \frac{3m_h^2}{4v^m} \vec{G}^2 h^m$$

Goldstone Equivalence
Theorem says
Goldstone scattering
gives high energy
longitudinal W,Z
scattering

Unitarity violating scale for $Z_L h^{m/2} \iff Z_L h^{m/2}$ is ~5 TeV for m ~ 10-15 (nondecoupling effect)





MODEL DEPENDENCE OF TERMS

$$X^{3} \sim h^{3} + \vec{G}^{2}(h^{2} + h^{3} + \cdots) + \vec{G}^{4}(h + h^{2} + \cdots) + \vec{G}^{6}(1 + h + \cdots)$$

$$+ \vec{G}^{8}(1 + h + \cdots) + \vec{G}^{10}(1 + h + \cdots) + \cdots ,$$

$$X^{4} \sim h^{4} + \vec{G}^{2}(h^{3} + h^{4} + \cdots) + \vec{G}^{4}(h^{2} + h^{3} + \cdots) + \vec{G}^{6}(h + h^{2} + \cdots)$$

$$+ \vec{G}^{8}(1 + h + \cdots) + \vec{G}^{10}(1 + h + \cdots) + \cdots ,$$

$$X^{5} \sim h^{5} + \vec{G}^{2}(h^{4} + h^{5} + \cdots) + \vec{G}^{4}(h^{3} + h^{4} + \cdots) + \vec{G}^{6}(h^{2} + h + \cdots)$$

$$+ \vec{G}^{8}(h + h^{2} + \cdots) + \vec{G}^{10}(1 + h + \cdots) + \cdots ,$$

(Schematic without coefficients, but we know cancellations can occur due to SMEFT description)

Terms circled can only come from trilinear!

STANDARD MODEL EFT (SMEFT)

Nonanalytic nature of HEFT around v = 0 reflects a nonlocal EFT for Higgs doublet in ultraviolet

SMEFT instead looks at the most general EW gauge invariant analytic EFT for H

$$Y \equiv |H|^2 - \frac{v^2}{2} = vh + \frac{1}{2} \left(h^2 + \vec{G}^2 \right)$$

$$V(Y) = \lambda_{SM} Y^2 + c_3 Y^3 + c_4 Y^4 + \cdots$$

SMEFTTRILINEAR

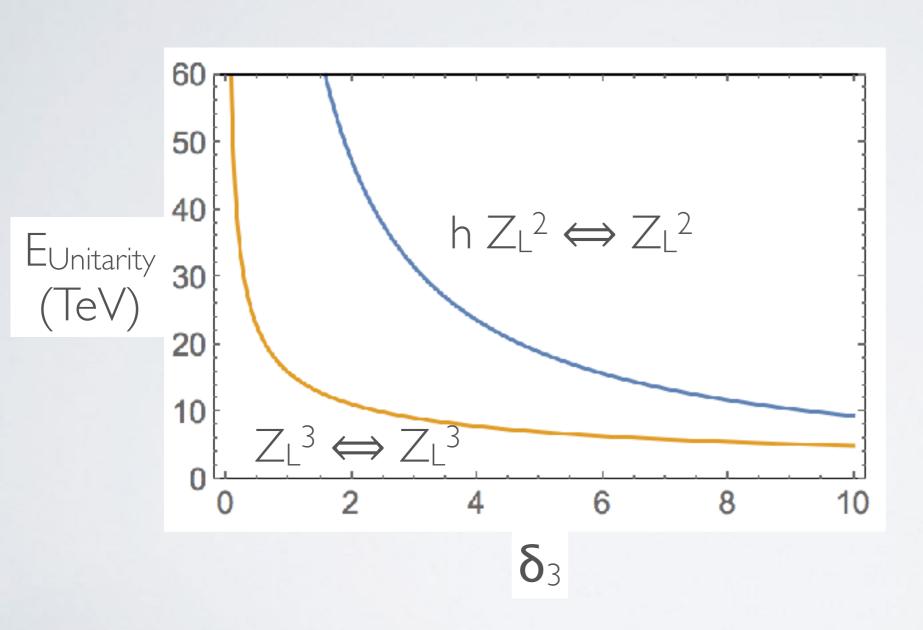
$$V = V_{SM} + \frac{m_h^2}{2v^4} \delta_3 Y^3 + \cdots$$

$$= V_{SM} + \frac{m_h^2}{2v} \delta_3 h^3 + \frac{3m_h^2}{4v^2} \delta_3 h^4 + \frac{3m_h^2}{8v^3} \delta_3 h^5 + \frac{m_h^2}{16v^4} \delta_3 h^6 + \cdots$$

This operator only leads to contact interactions up to 6-body, so it proves that high multiplicities can be cancelled by choosing correlated h⁴,h⁵,h⁶ couplings

These potential cancellations suggest one should focus on processes that only depend on trilinear (i.e. G⁶, hG⁵)

MODEL INDEPENDENT VIOLATION



These couplings only depend on trilinear modifications and give much higher bounds than ~5 TeV, but also decouple as coupling vanishes, giving a conservative bound

Coupled channel analysis leads to improved bound $13\,\text{TeV}/\sqrt{\delta_3}$

W/Z COUPLINGS (IN PROGRESS)

$$Y|D_{\mu}H|^2$$

tells us that only up to 4-body processes are dependent solely on δ_{hVV}

e.g. WW \rightarrow ZZ, gives E< 4.7 TeV/ $\sqrt{(\delta_{hVV}/0.05)}$

However, if only hhVV is modified, then need a higher dimensional operator, which says now 6-body processes are interesting

$$|Y^2|D_\mu H|^2$$

e.g. WWZ \rightarrow WWZ, gives E< 2.6 TeV/ δ ^{1/4}hhVV

TOP COUPLINGS (IN PROGRESS)

$$Y\bar{T}_RH\epsilon Q_L$$

tells us that only up to 5-body processes are dependent solely on δ_{htt}

e.g.
$$t_R \bar{b}_R \to W^+ W^+ W^-$$
 gives E< 33 TeV/ $\sqrt{(\delta_{\rm htt}/0.05)}$

However, if only hhtt is modified, then 7-body processes are interesting

$$Y^2 \bar{T}_R H \epsilon Q_L$$

e.g.
$$t_R \bar{b}_R W^- \to W^4$$
 gives E< 4 TeV/c2 1/4

UNITARITY BOUND SUMMARY

| Higgs | hVV | htt | | | |
|------------------------|--|--|--|--|--|
| Trilinear | hhVV | hhtt | | | |
| 13 TeV/√δ ₃ | $0.7~{\rm TeV/V}(\delta_{\rm hVV}/0.05)$ $2.6~{\rm TeV/\delta^{1/4}_{hhVV}}$ | 33 TeV/√(δ _{htt} /0.05) 4 TeV/c ₂ ^{1/4} | | | |

These suggest, nontrivial V or top couplings is kinematically accessible at 14-27 TeV hadron collider, while trilinear requires 100 TeV. (For HEFT, bounds would be few TeV.)

COLLIDER PROBES

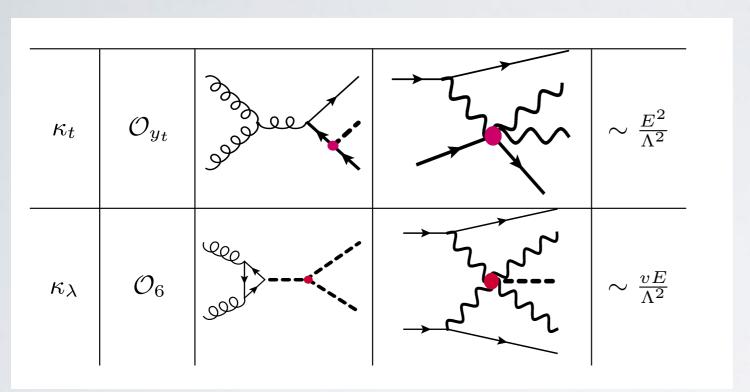
SM behavior

Probing high energy processes can test energy growth, complementary sensitivity to Higgs couplings

New resonances possible, but not guaranteed.
E.g. Higgs not discovered in VB scattering

ENERGY

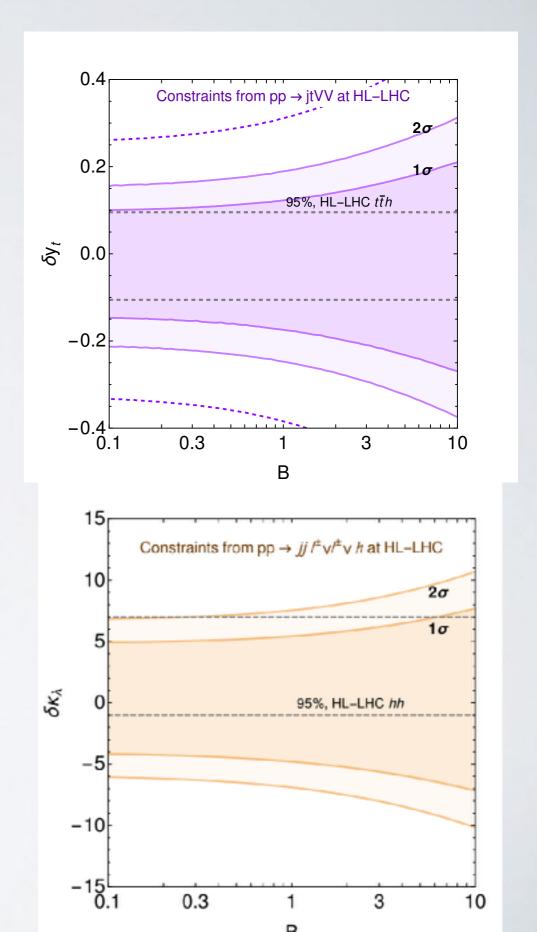
COLLIDER TESTS OF UNITARITY VIOLATION



Searching for Unitarity violating processes (solid) has similar sensitivities to coupling measurement (dashed) for tth, hhh

Extension to tthh and VVhh?

Henning et.al. 1812.09299



COLLIDER PROBES QUESTIONS

Have we exhausted interesting processes? Unitarity violation only suggests sensitivity for very high energies, so is there another way to identify processes?

Also, if new physics is at accessible scale, need to systematically explore UV completions

(e.g. Historical lesson, W/Z and Higgs weren't discovered by looking at unitarity violating processes, appeared at much lower energy than then needed to)

CONCLUSIONS

- Precision Higgs couplings could discover a deviation from SM, suggesting new physics at some energy scale
- Unitarity violation places a bound on this energy scale and identifies interesting correlated processes to measure
- Higgs trilinear modifications lead to Unitarity violation at high energies (~ 5 13 TeV for $\delta_3 \sim 1$ depending on assumptions)
- VVhh, tthh is of interest at HL-LHC and future colliders,
 Unitarity bound is ~3-5 TeV for O(I) coupling

Thanks for your attention!

EXTRA SLIDES

DISCUSSION QUESTIONS

- Possible to systematically model build UV completions of these coupling modifications? If so, what are the correlated signals?
- Other technique besides Equivalence theorem that identifies interesting processes that are modified by nonstandard couplings?

UNITARITY CONSTRAINTS ON NON-DERIVATIVE COUPLINGS

$$\frac{\lambda}{n_1! \cdots n_r!} \phi_1^{n_1} \cdots \phi_r^{n_r}$$

Consider s-wave scattering

$$\phi_1^{k_1} \cdots \phi_r^{k_r} \leftrightarrow \phi_1^{n_1 - k_1} \cdots \phi_r^{n_r - k_r}$$

If we properly normalize these states, then unitarity bounds are easy (see paper for more details)

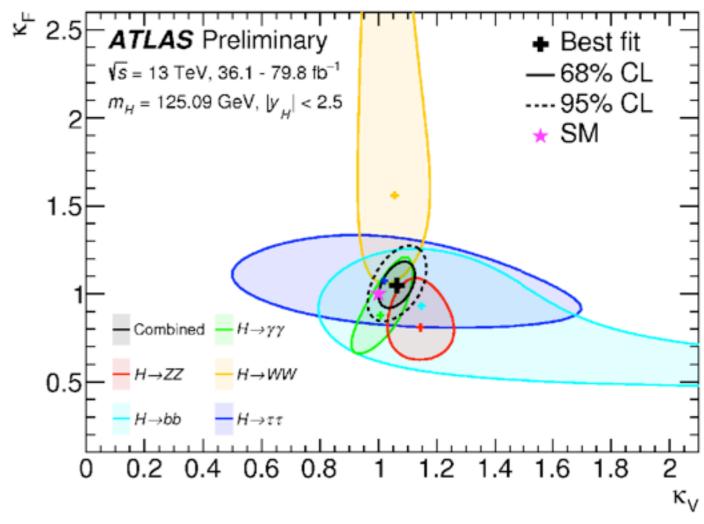
Unitarity constraints from this amplitude requires

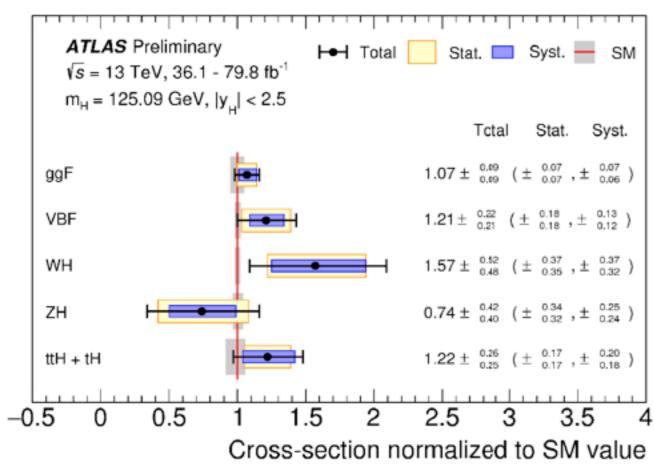
$$E \le 4\pi \left[\frac{64\pi^2}{\lambda^2} \left(k_1! \cdots k_r! \left(k - 1 \right)! \left(k - 2 \right)! \right) \left((n_1 - k_1)! \cdots (n_r - k_r)! \left(n - k - 1 \right)! \left(n - k - 2 \right)! \right) \right]^{\frac{1}{2n - 8}}$$
 where $n \equiv n_1 + \cdots + n_r, k \equiv k_1 + \cdots + k_r$

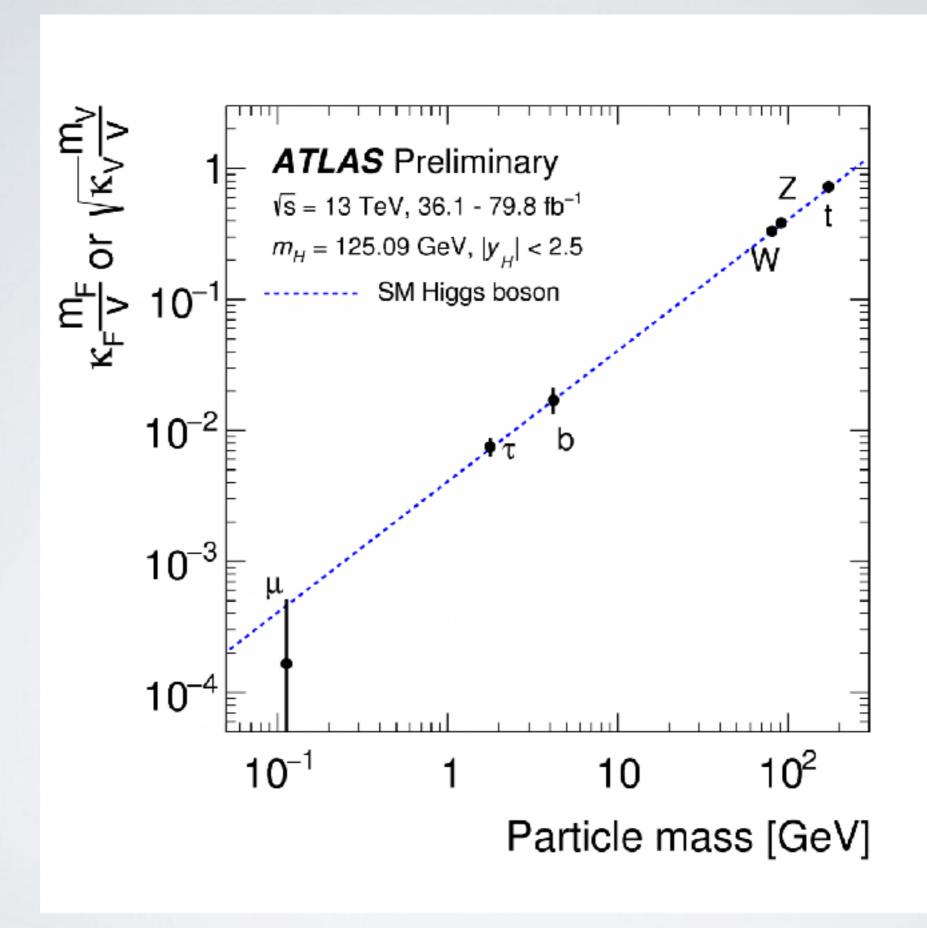
SCALAR STATES

$$|P, k_1, \dots, k_r\rangle = C_{k_1 \dots k_r} \int d^4x \, e^{-iP \cdot x} \prod_{i=1}^r \left[\phi_i^{(-)}(x) \right]^{k_i} |0\rangle$$

$$\frac{1}{|C_{k_1\cdots k_r}|^2} = \frac{1}{(k-1)!(k-2)!} \frac{1}{8\pi} \left(\prod_{i=1}^r k_i!\right) \left(\frac{E}{4\pi}\right)^{2k-4}$$







INCLUDING QUARTIC

$$V \supset \frac{m_h^2}{8v^2} (1 + \delta_4) h^4 + \frac{m_h^2}{4v^3} (\delta_4 - 3\delta_3) h^3 \vec{G}^2 + \frac{3m_h^2}{16v^4} (\delta_4 - 5\delta_3) h^2 \vec{G}^4 + \frac{m_h^2}{16v^5} (\delta_4 - 6\delta_3) h \vec{G}^6 + \frac{m_h^2}{128v^6} (\delta_4 - 6\delta_3) \vec{G}^8.$$

| Process | Unitarity Violating Scale |
|---------------------------------------|--|
| $h^2 Z_L \leftrightarrow h Z_L$ | $66.7 \text{ TeV}/ \delta_3 - \frac{1}{3}\delta_4 $ |
| $hZ_L^2 \leftrightarrow Z_L^2$ | $94.2 \text{ TeV}/ \delta_3 $ |
| $hW_LZ_L \leftrightarrow W_LZ_L$ | $141 \text{ TeV}/ \delta_3 $ |
| $hZ_L^2 \leftrightarrow hZ_L^2$ | $9.1 \text{ TeV}/\sqrt{ \delta_3 - \frac{1}{5}\delta_4 }$ |
| $hW_L Z_L \leftrightarrow hW_L Z_L$ | $11.1 \text{ TeV}/\sqrt{ \delta_3 - \frac{1}{5}\delta_4 }$ |
| $Z_L^3 \leftrightarrow Z_L^3$ | $15.7 \text{ TeV}/\sqrt{ \delta_3 }$ |
| $Z_L^2 W_L \leftrightarrow Z_L^2 W_L$ | $20.4 \text{ TeV}/\sqrt{ \delta_3 }$ |
| $hZ_L^3 \leftrightarrow Z_L^3$ | $6.8 \text{ TeV}/ \delta_3 - \frac{1}{6}\delta_4 ^{\frac{1}{3}}$ |
| $hZ_L^2W_L \leftrightarrow Z_L^2W_L$ | $8.0 \text{ TeV}/ \delta_3 - \frac{1}{6}\delta_4 ^{\frac{1}{3}}$ |
| $Z_L^4 \leftrightarrow Z_L^4$ | $6.1 \text{ TeV}/ \delta_3 - \frac{1}{6}\delta_4 ^{\frac{1}{4}}$ |