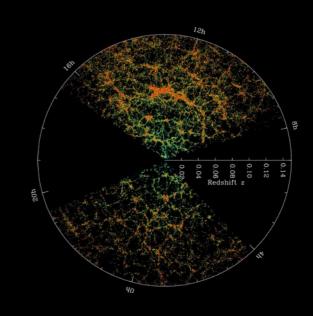
Signals of New Physics in the Large Scale Structure Data

Yuhsin Tsai

U. Maryland / U. Notre Dame

Particle Physics in Computing Frontier 12/12/2019



There's no machine learning in this talk

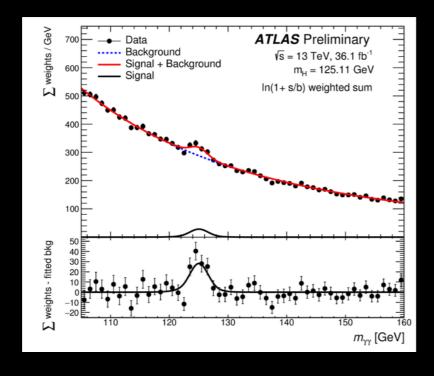
But I will explain:

What is the Large Scale Structure (LSS) of the universe?

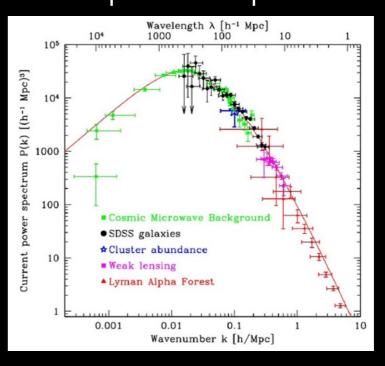
How can BSM physics show up?

What can BSM signals look like?

collider data



matter power spectrum



The plan:

Large Scale Structure (LSS) of the universe

LSS signal from "homogeneously-distributed" matter

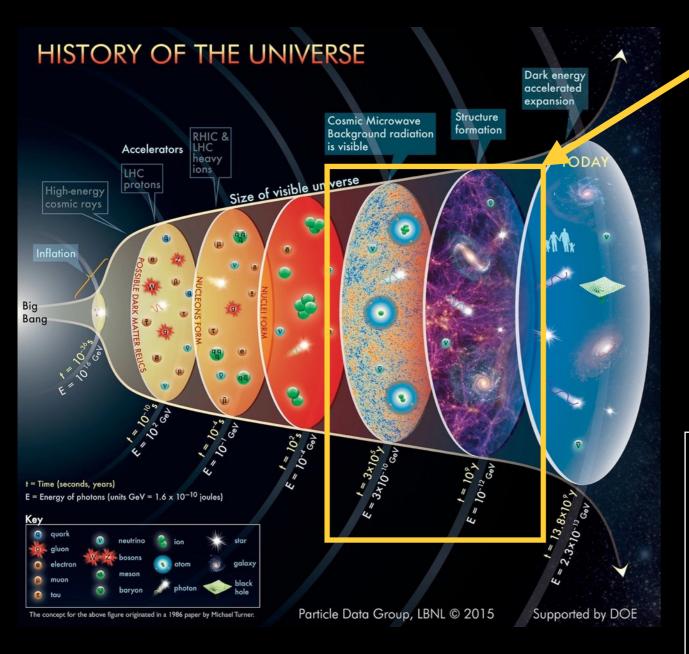
- example 1: Neutral Naturalness models
- example 2: signals of neutrino decay

Discussion:

 defects in CMB, LSS and the DM thermal distribution anisotropy in gravitational wave background...

Large Scale Structure of the universe

Time scale of the structure formation

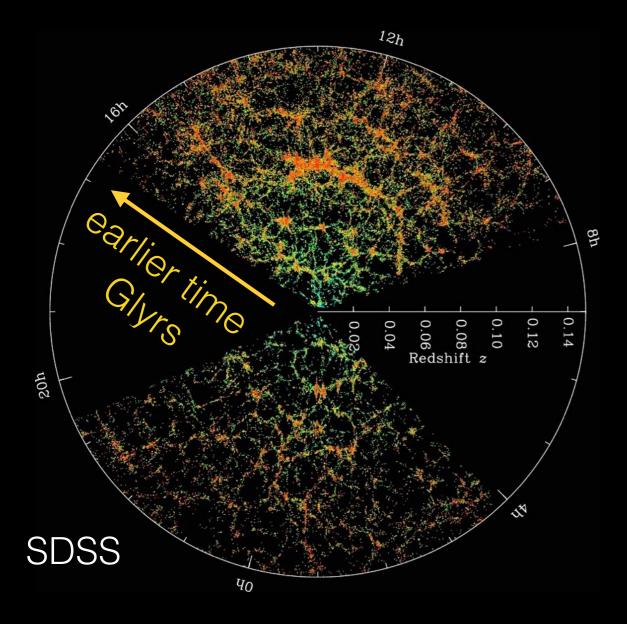


Cold DM clumps since matter-radiation equilibrium (T ~ eV, z ~ 3000) and the structure formation continues until CC dominates (T ~ 10^-4 eV, z ~ 1)

New physics can modify LSS by

- change initial condition
- change expansion history
- change DM clumping

Several ways to see LSS, e.g., galaxy survey



it's a 3D map

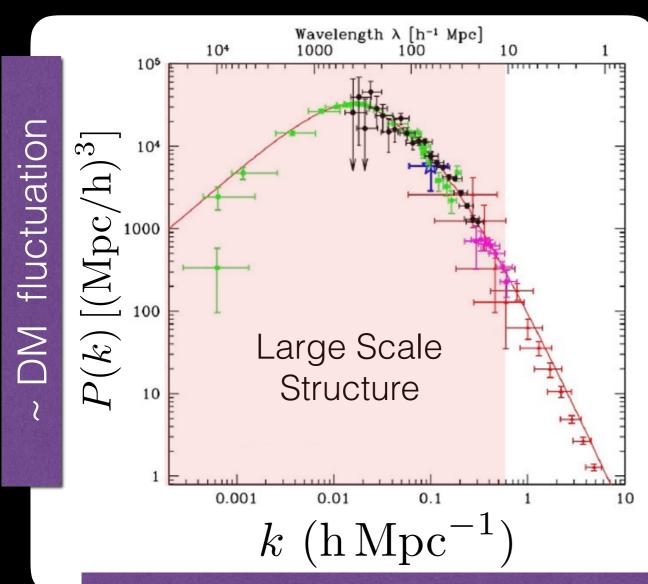
galaxy distribution is determined by cold dark matter distribution

$$\frac{\delta N_{\rm galaxy}}{N_{\rm galaxy}} \propto \frac{\delta \rho_{\rm cdm}}{\rho_{\rm cdm}} \gg 10^{-5}$$

density fluctuation is larger than CMB fluctuation due to structure formation

Matter power spectrum

$$P(k) \sim k^{-3} \delta_m(k)^2 \qquad k \sim L^{-1}$$

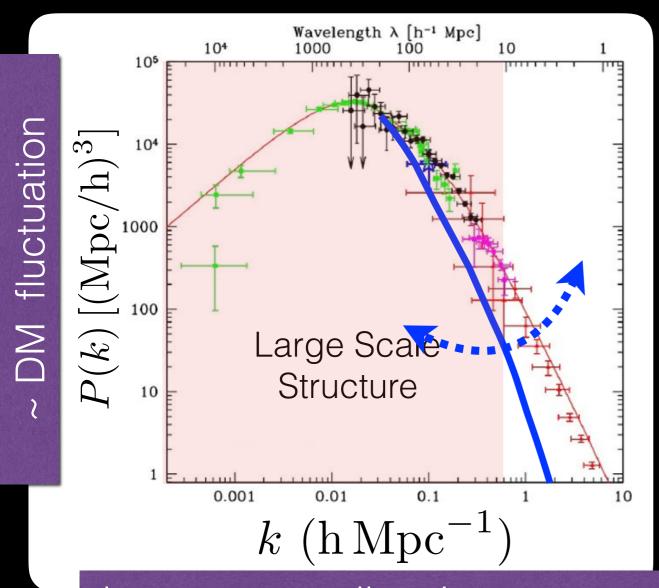


larger -> smaller size structure

Matter power spectrum

$$P(k) \sim k^{-3} \delta_m(k)^2$$

$$k \sim L^{-1}$$

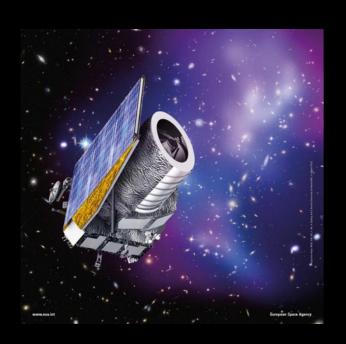


New physics can change the structure *everywhere* in the universe

current measurements have about 10% level precision

larger -> smaller size structure

Higher precision measurements will come soon



Euclid (2022)



LSST (2023)



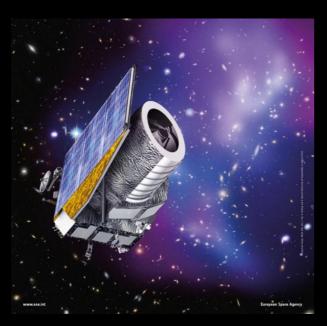
DESI (2019)



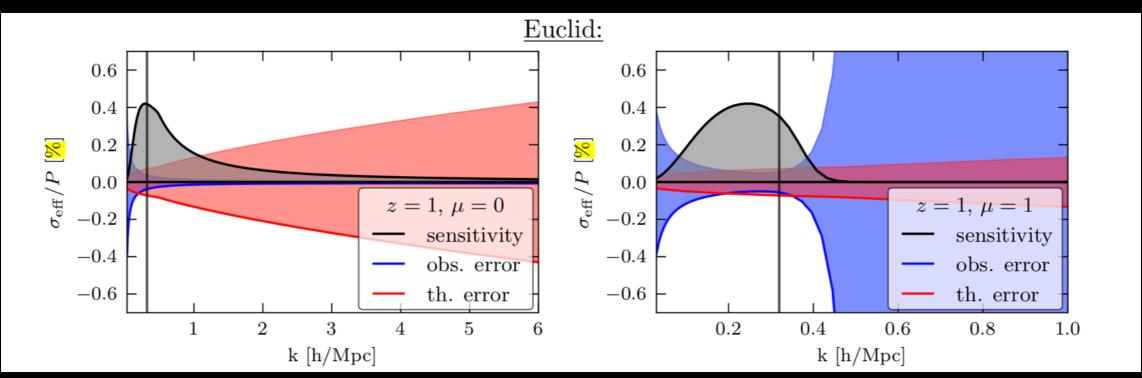
will be sensitive to ~ 0.1 - 1% change in matter power spectrum (mainly by increasing statistics)

CMB-S4 (202?)

E.g., Euclid experiment

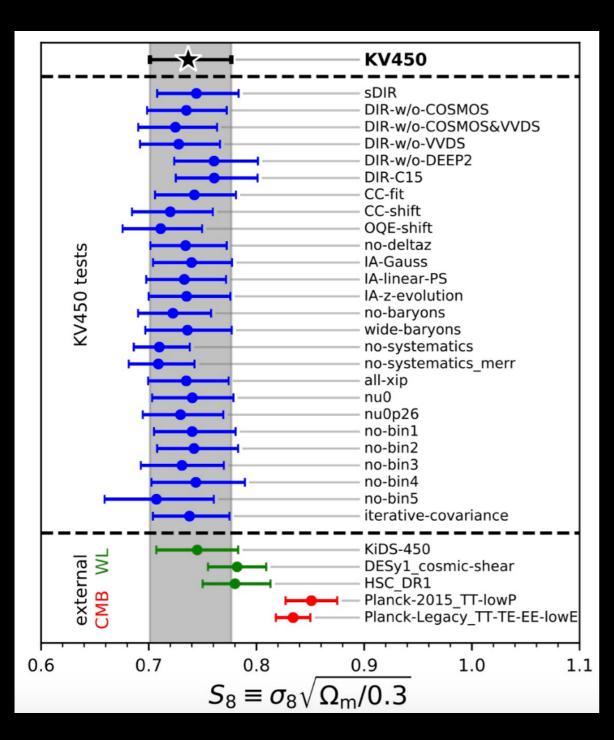


Euclid (plan to launch at 2022) $0.45 \sim < z \sim < 2.05$



Sprenger, Archidiacono, Clesse, Lesgourgues (2018)

sigma 8 (S8) anomaly?



all the late time density contrast measurements get smaller results than

Planck fit by assuming LCDM (2~3 sigma level)

LSS signals from homogeneously distributed matter

L $ar{
ho}_{
m cdm}$ $ar{
ho}_{
m cdm}$ $ar{
ho}_{
m cdm}$

~ primordial fluctuation

$$t_{eq} + \Delta t$$

 I_{ρ} I_{ρ

matter falls into deeper gravitational well

Newtonian physics

$$\frac{\mathrm{d}^2 x}{\mathrm{d} t^2} \sim \frac{G \, \delta m_{\mathrm{cdm}}}{L^2}$$

$$\frac{\Delta x}{L} \sim \frac{\delta \rho_{\rm cdm}}{\bar{\rho}_{\rm cdm}} = \delta_{\rm cdm}$$

$$\frac{\mathrm{d}^2 \delta_{\mathrm{cdm}}}{\mathrm{d} t^2} \sim G \, \bar{\rho}_{\mathrm{cdm}} \, \delta_{\mathrm{cdm}}$$

$t \rightarrow a$

We don't quite know the physical time of structure formation, but we know it mainly begins at matter-radiation equilibrium

$$\rho_m(a_{eq}) = \rho_r(a_{eq})$$

and the physical time depends on energy density

$$dt = \frac{da}{aH(a,\bar{\rho}_{tot})} \approx H_0^{-1} \left(\frac{\rho_c}{\bar{\rho}_{tot}}\right)^{\frac{1}{2}} \sqrt{a} da$$

Larger total energy, shorter physical time for structure formation

$$\frac{\mathrm{d}^2 \delta_{\mathrm{cdm}}}{\mathrm{d} t^2} \sim G \, \bar{\rho}_{\mathrm{cdm}} \, \delta_{\mathrm{cdm}}$$
 $\mathrm{d} t \to \mathrm{d} a$

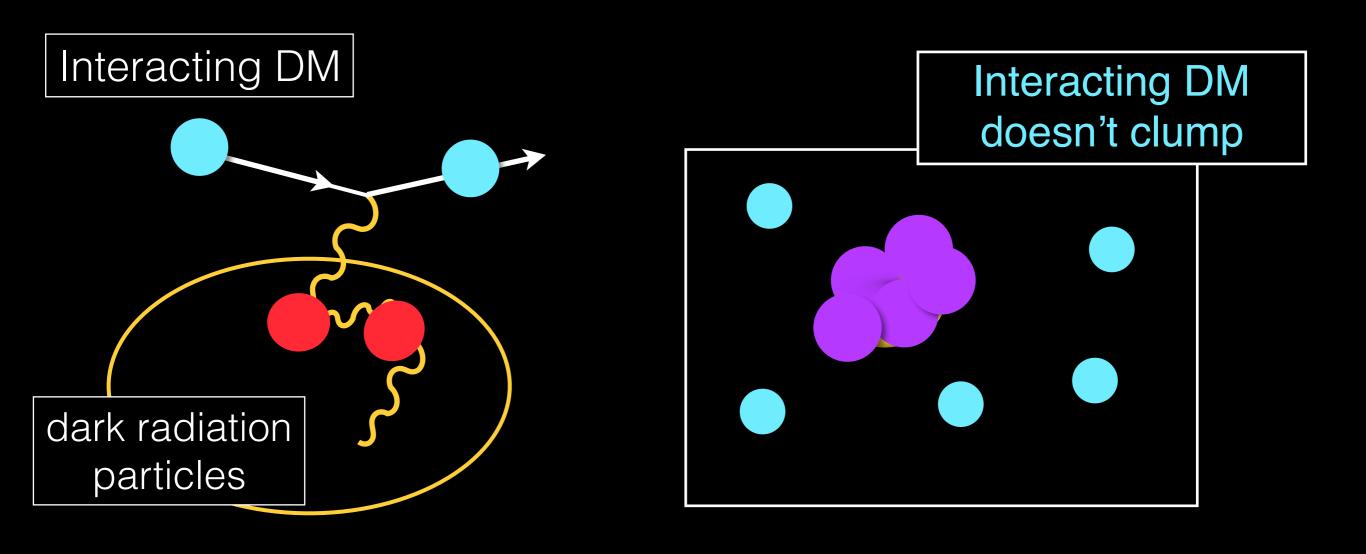
$$dt \to da$$

$$\left[\frac{\mathrm{d}^2}{\mathrm{d}a^2} - \frac{3}{2a}\frac{\mathrm{d}}{\mathrm{d}a} - \frac{1}{2a^2}\right]\delta_{\mathrm{cdm}}(a) \sim \frac{1}{a^2}\left(\frac{\bar{\rho}_{\mathrm{cdm}}}{\bar{\rho}_{\mathrm{tot}}}\right)\delta_{\mathrm{cdm}}$$

$$\delta_{\text{cdm}}(a_f) = \delta_{\text{cdm}}(a_i) \left(\frac{a_f}{a_i}\right)^{1 - \frac{3}{5}} \frac{\left(1 - \frac{\bar{\rho}_{\text{cdm}}}{\bar{\rho}_{\text{tot}}}\right)}{a_i}$$

If having matter that doesn't clump => slow down the growth!

Matter-radiation interaction can suppress LSS

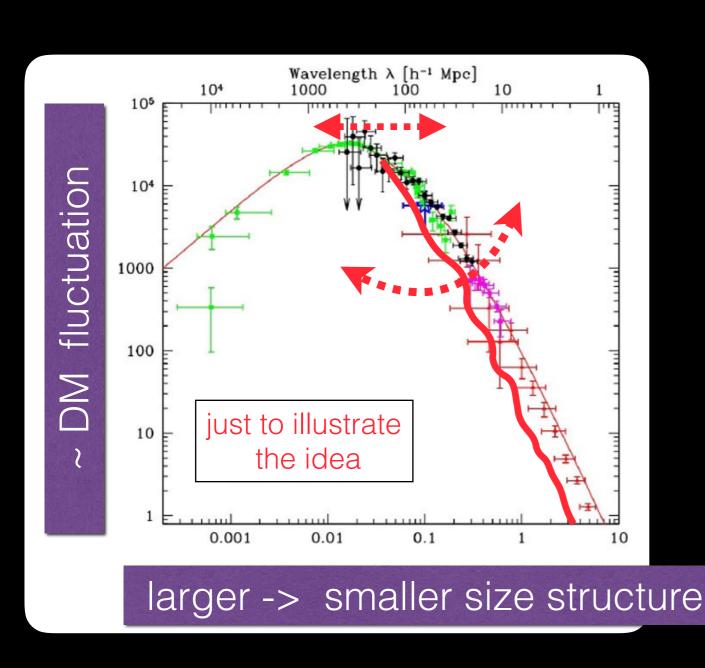








Matter-radiation interaction can suppress LSS



a scale-dependent suppression

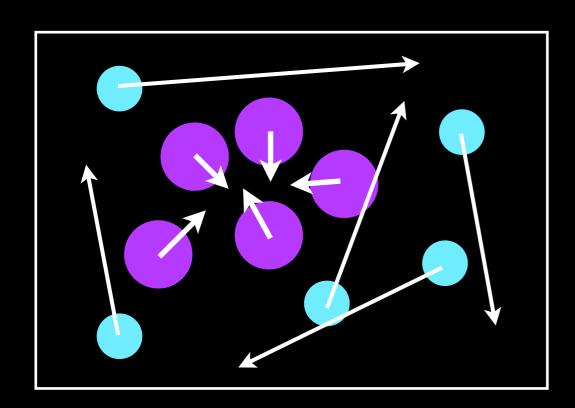
Dark Acoustic Oscillations (Cyr-Racine el al (2013))

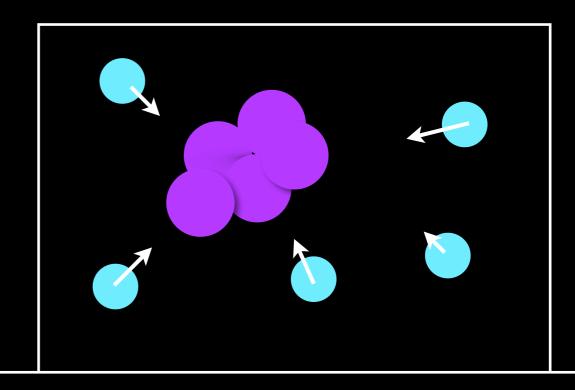
see e.g., Chacko, Cui, Hong, Okui, **YT** (2016), Prilepina, **YT** (2017), Chacko, Curtin, Geller, **YT** (2018), Dessert, Kilic, Trendafilova (2018)

Warm DM (or massive neutrinos) suppress LSS

neutrinos look massless when
dark matter begins to clump

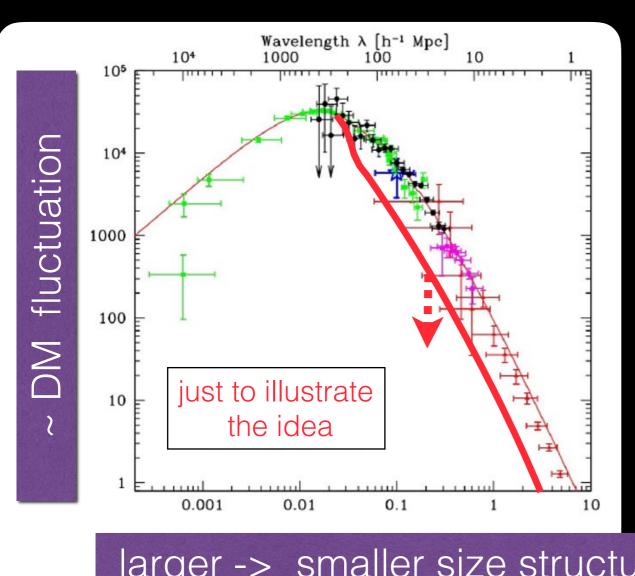
neutrinos finally slow down dark matter already forms structure





they don't help structure formation

Warm DM (or massive neutrinos) suppress LSS



reduce structure before neutrinos become non-relativistic

larger -> smaller size structure

see e.g., Chacko, Dev, Du, Poulin, YT (2019) and the reference therein

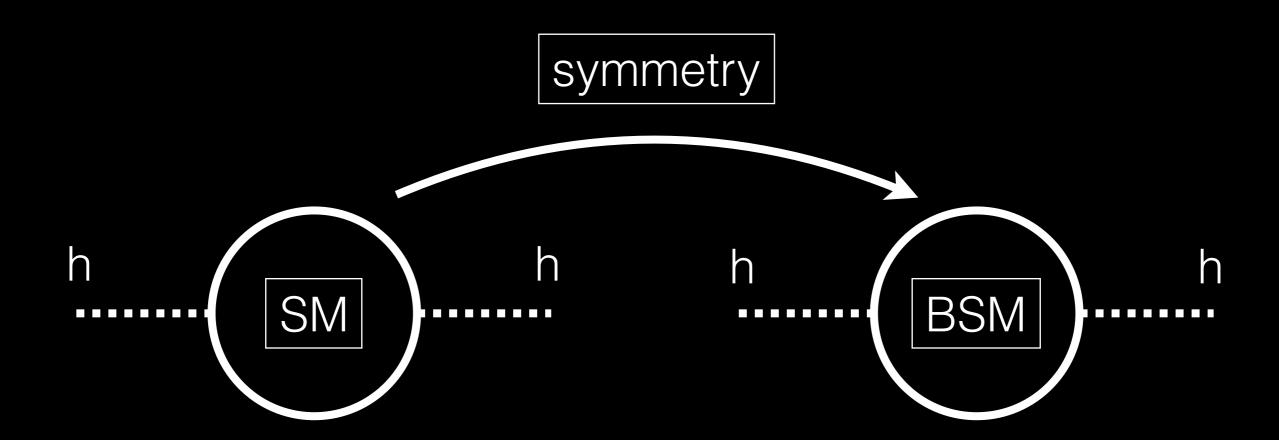
Eg.1: LSS signals of Hidden Naturalness model

— an example of Mirror Twin Higgs model

(other examples: *N*naturalness model, ...)

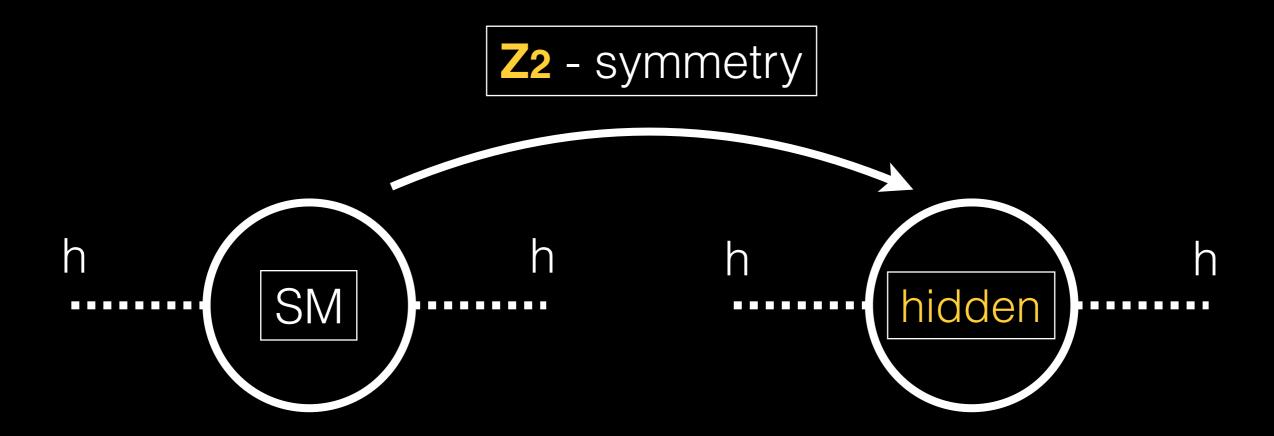
Common feature: (almost) no collider signatures

=> have to study them in cosmology



The Mirror Twin Higgs model (MTH)

Chacko, Goh, Harnik (2005)



contain SM-like particles in the "twin"-sector

BUT with lower-T and slightly larger VEV (3<f/v<5)

A long time ago, when T ~ eV

Mirror Twin Sector

GARDIANS OF THE ELECTROWEAK FORCE

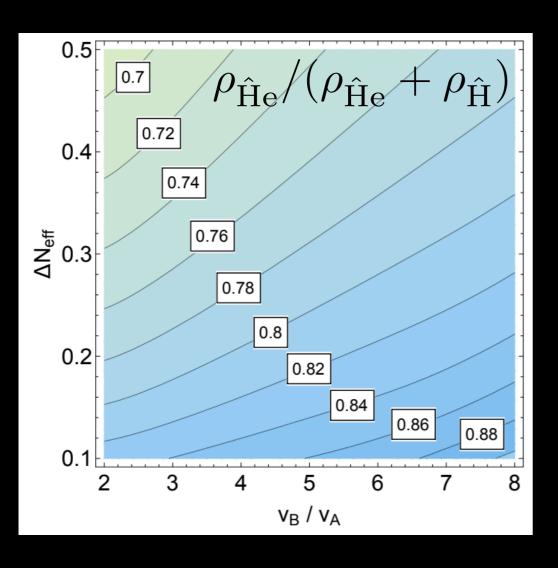
A long time ago, in a hidden universe that is so close to us

There are twin particles maintaining the stability of the Universe

mirror baryons (p/n/e)
dark radiation (photon/neutrinos)

We know the twin helium/hydrogen composition

Chacko, Curtin, Geller, YT (2018)

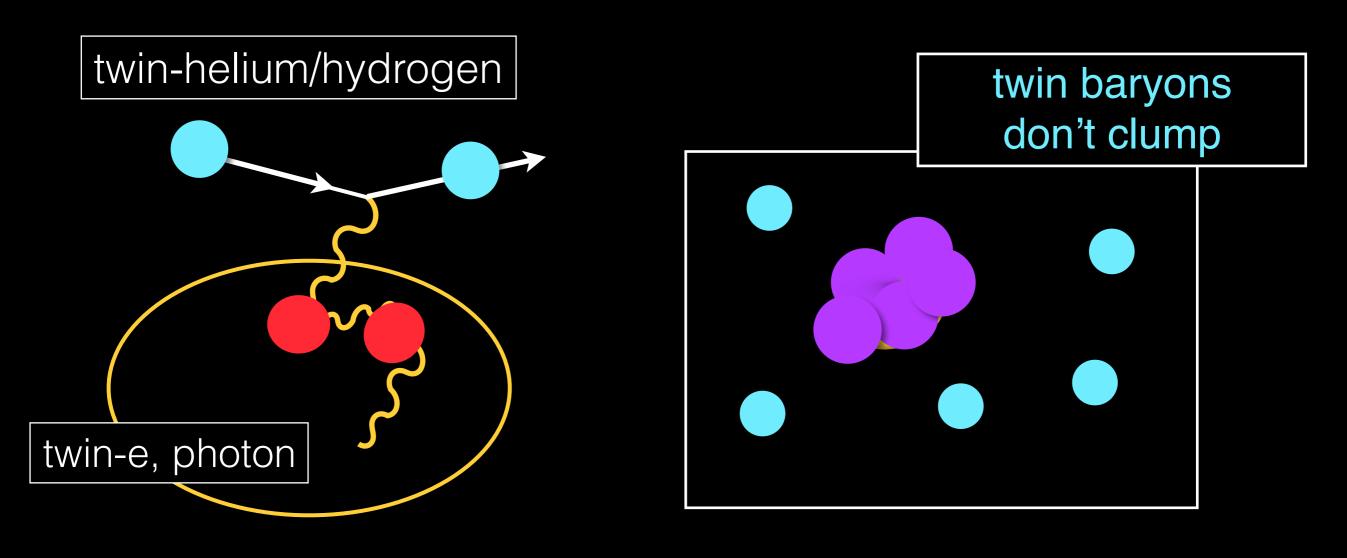


Mirror: ~ 75% mass is in mirror He

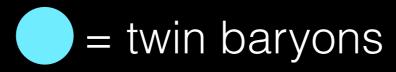
SM: ~ 75% mass is in Hydrogen

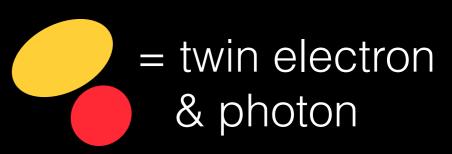
twin helium dominates the twin baryon acoustic oscillations

Twin Acoustic Oscillations



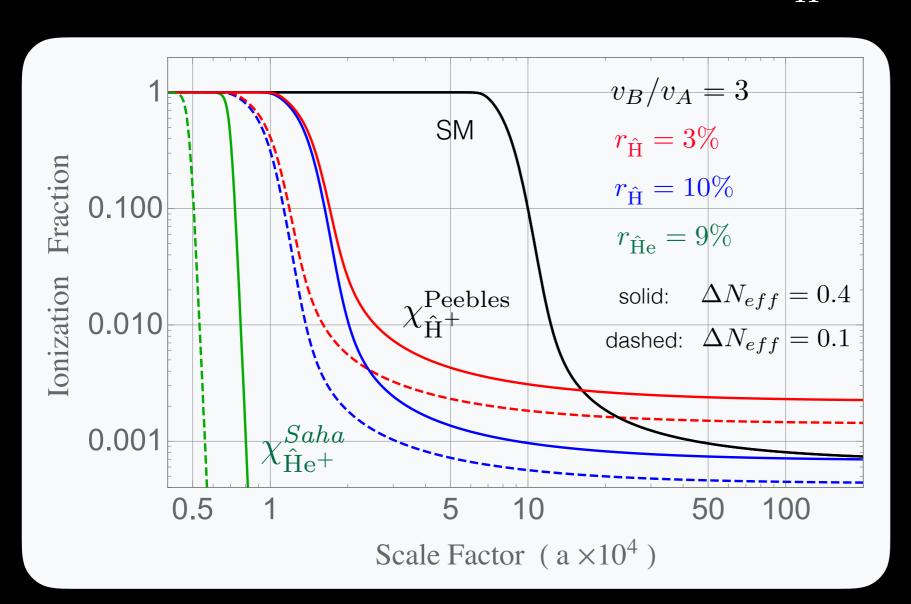






Twin-recombination process

$$H^{+} + e^{-} \to H^{0} + \gamma + (\gamma)$$
 $\frac{n_{H^{+}} n_{e^{-}}}{n_{H^{0}}} \sim \left(\frac{m_{e} T}{2\pi}\right)^{3/2} e^{-\frac{13.6 \text{ eV}}{T}}$



taking more precise energy transitions

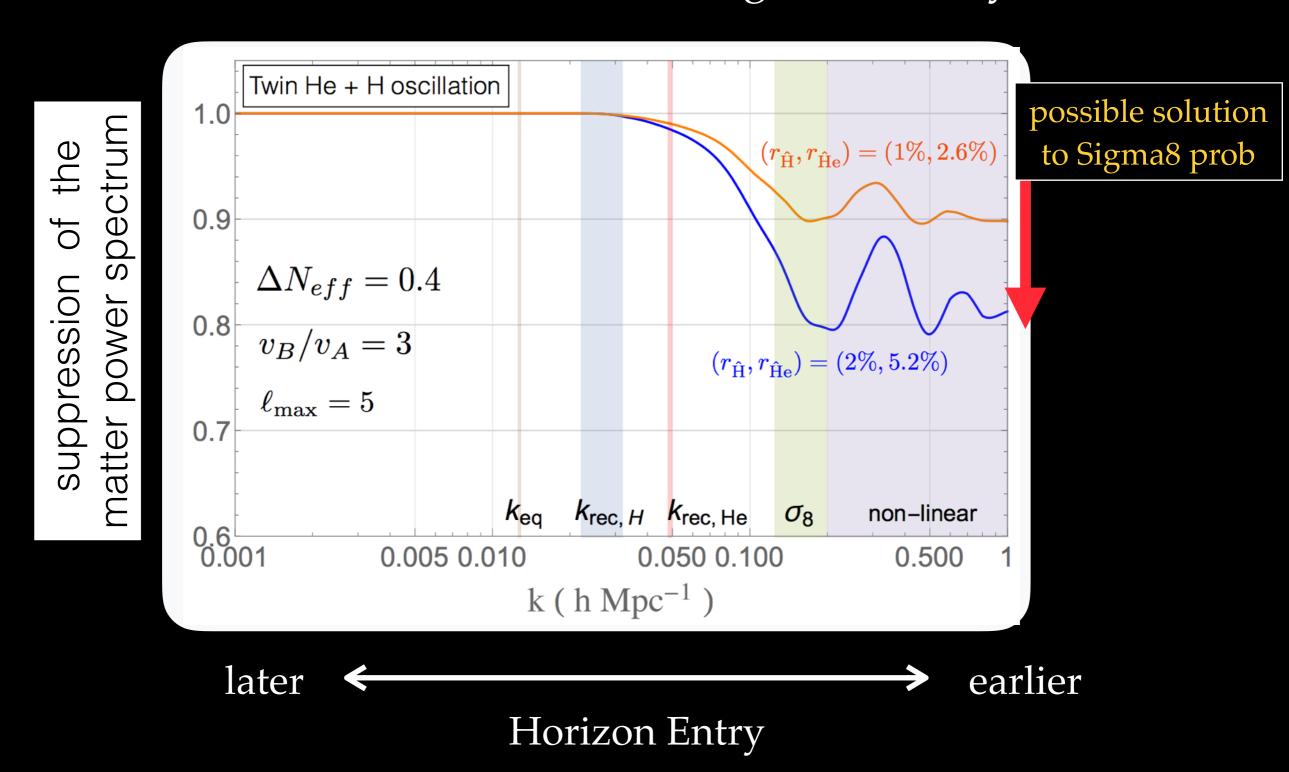
into account (Peebles)

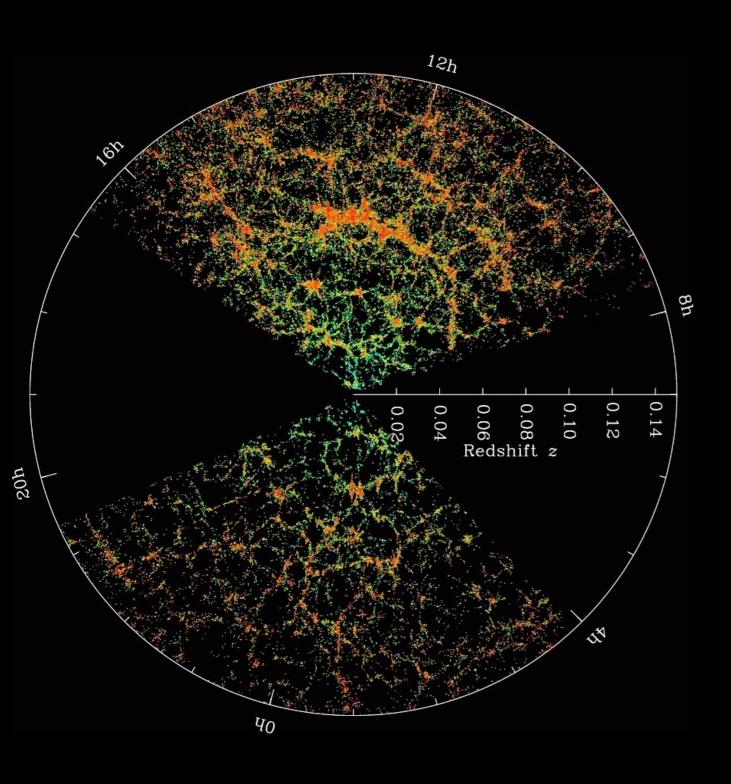
Saha's eq

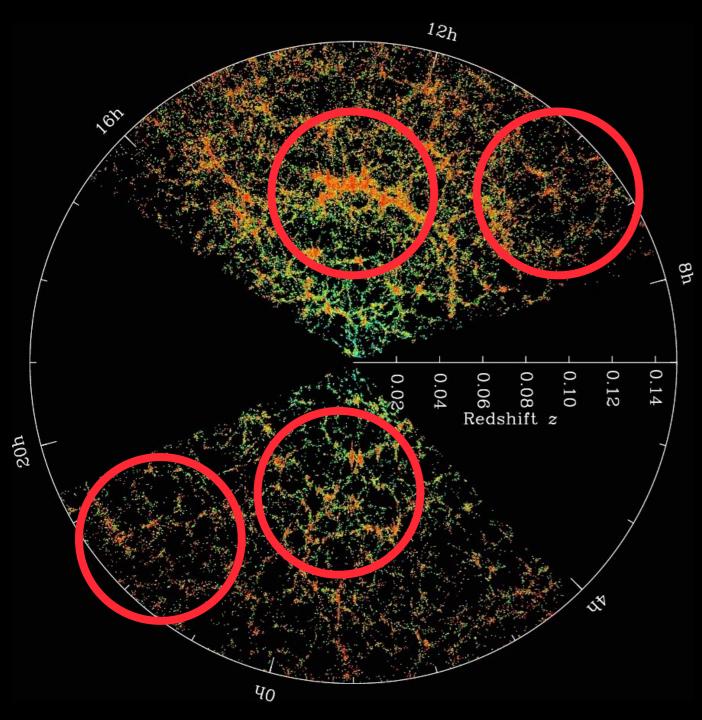
Chacko, Curtin, Geller, YT (2018)

Suppression of the Large Scale Structure

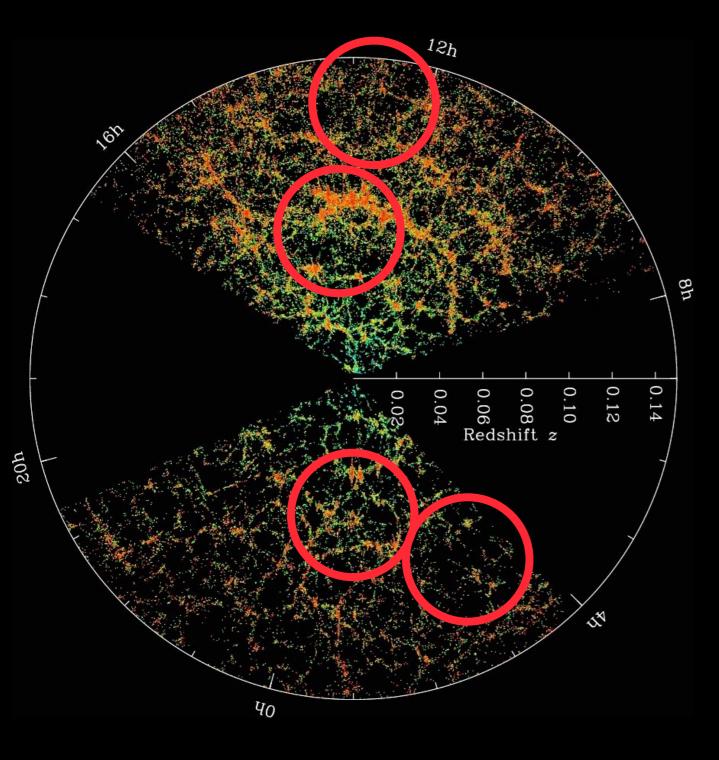
(with some fraction of DM being mirror baryons)



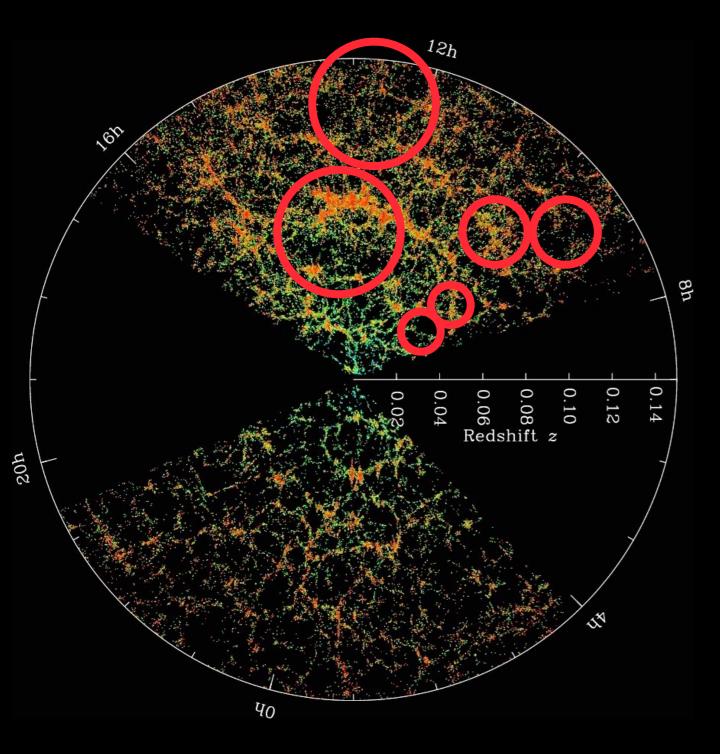




bigger structure (before recomb) density contrast as expected



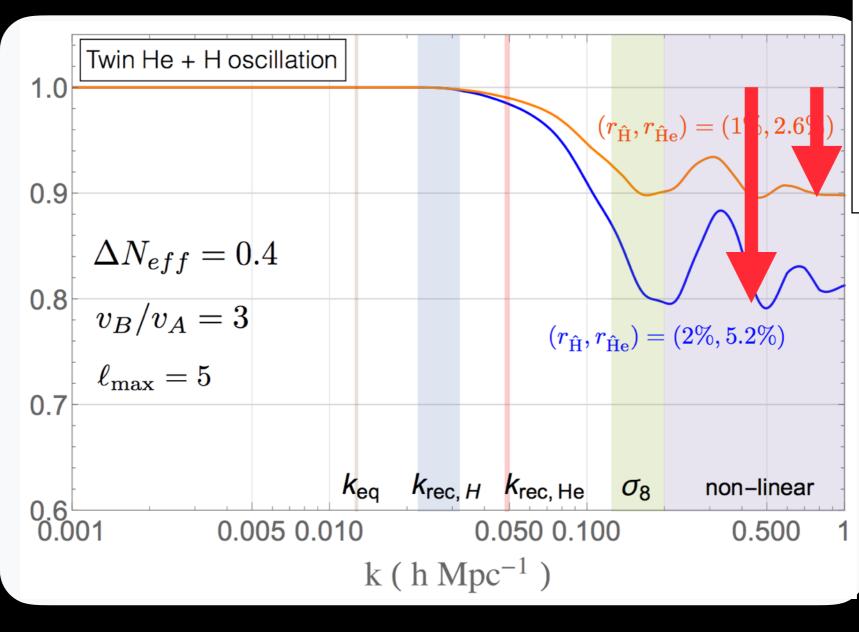
smaller structure (after recomb)
density contrast is smaller
than expected



smaller structure (after recomb)
density constrast oscillates
when looking at smaller and
smaller size of the structure

I. measure the mirror matter density

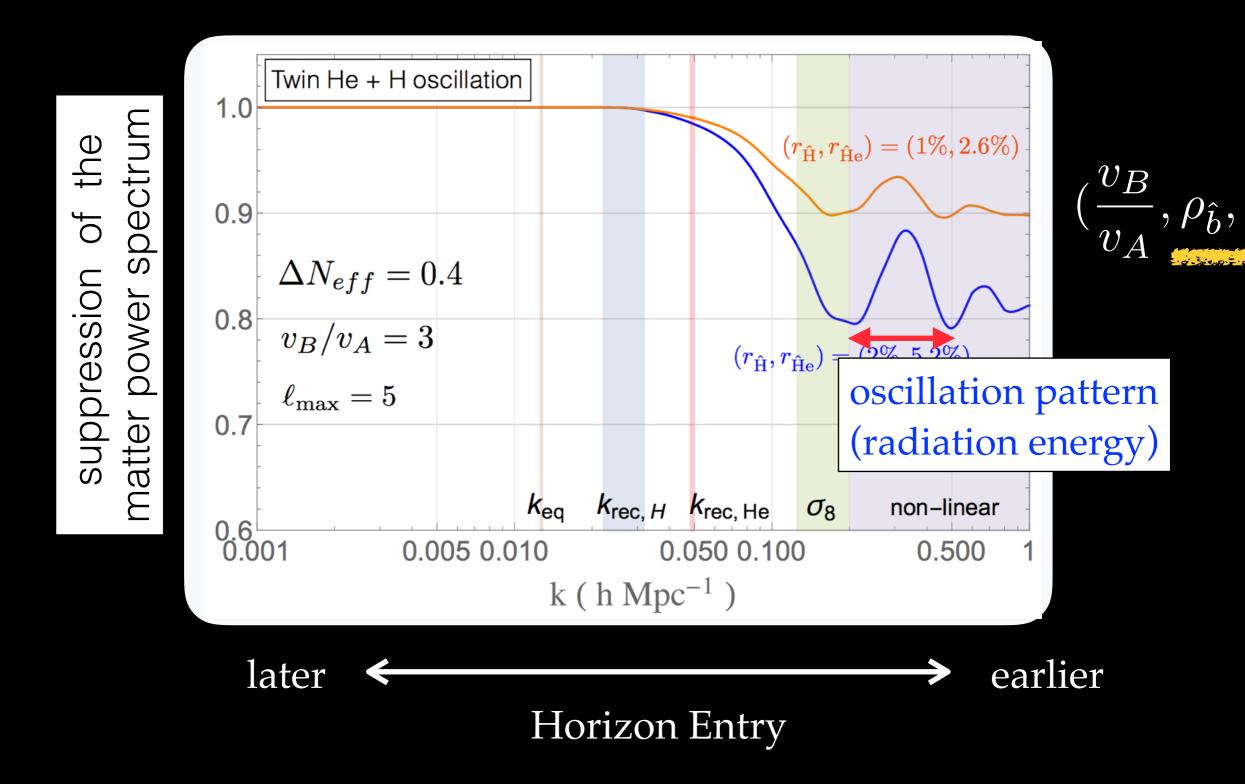
suppression of the matter power spectrum



suppression
due to mirror
oscillations
(mass fraction)

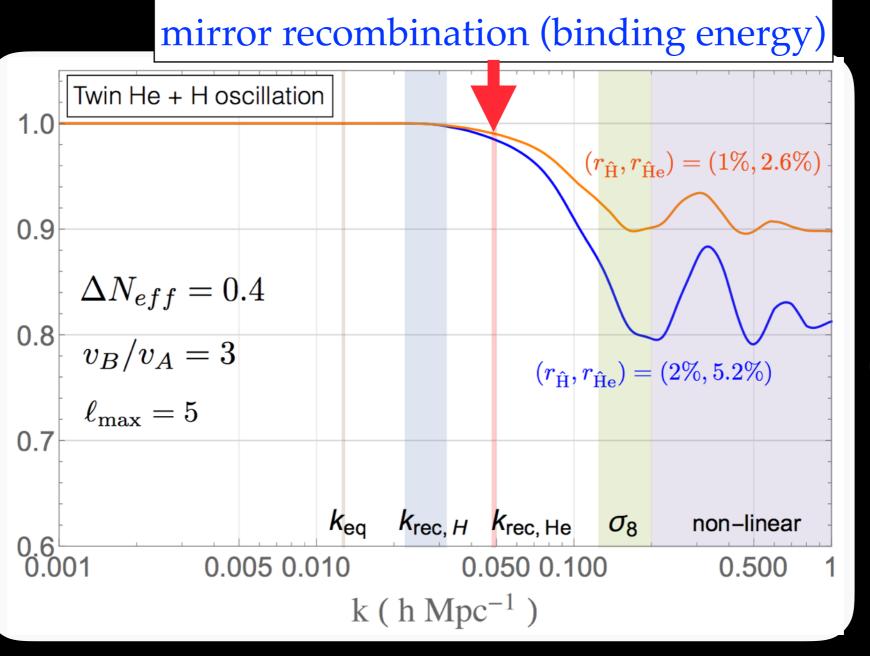
$$(\frac{v_B}{v_A}, \rho_{\hat{b}}, \hat{T})$$

II. measure mirror photon energy



III. measure binding energy of mirror atoms





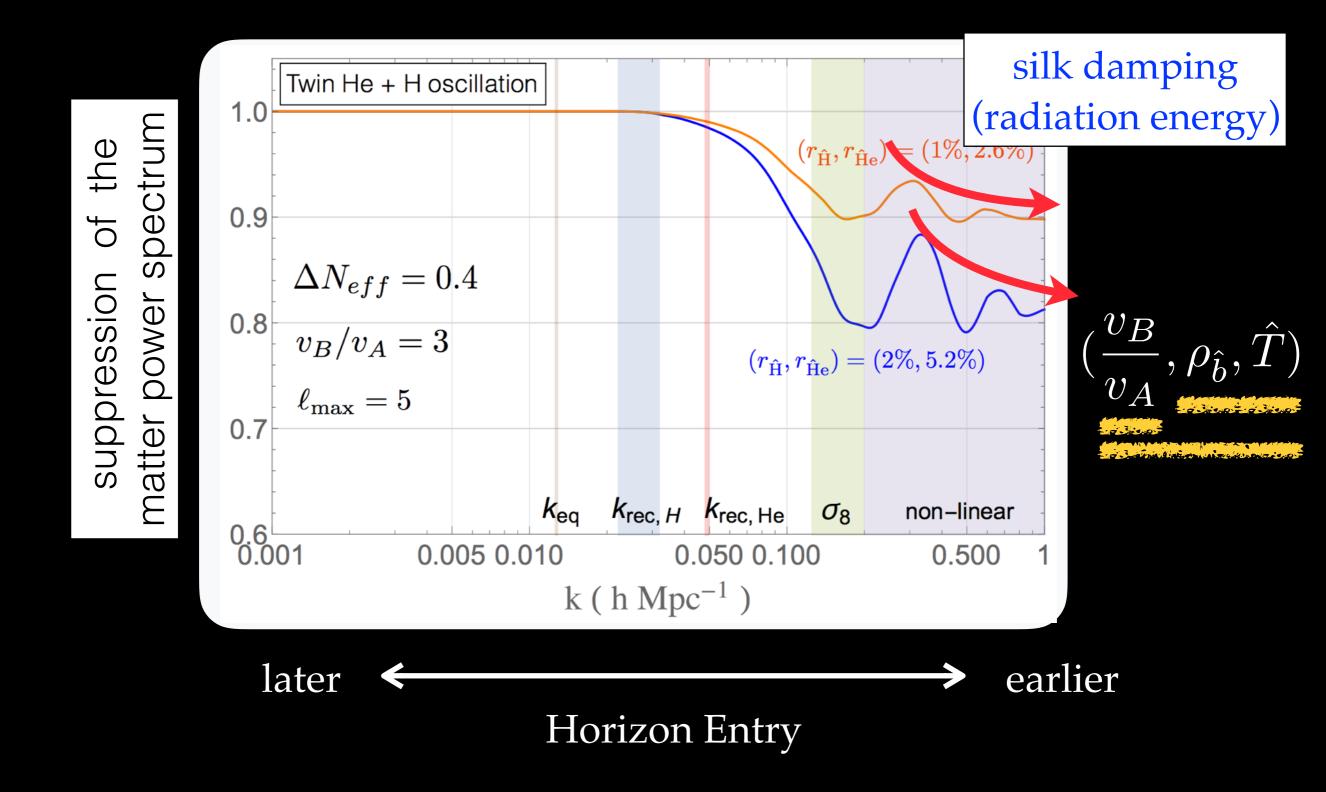
 $(rac{v_B}{v_A},
ho_{\hat{b}},\hat{T})$

later

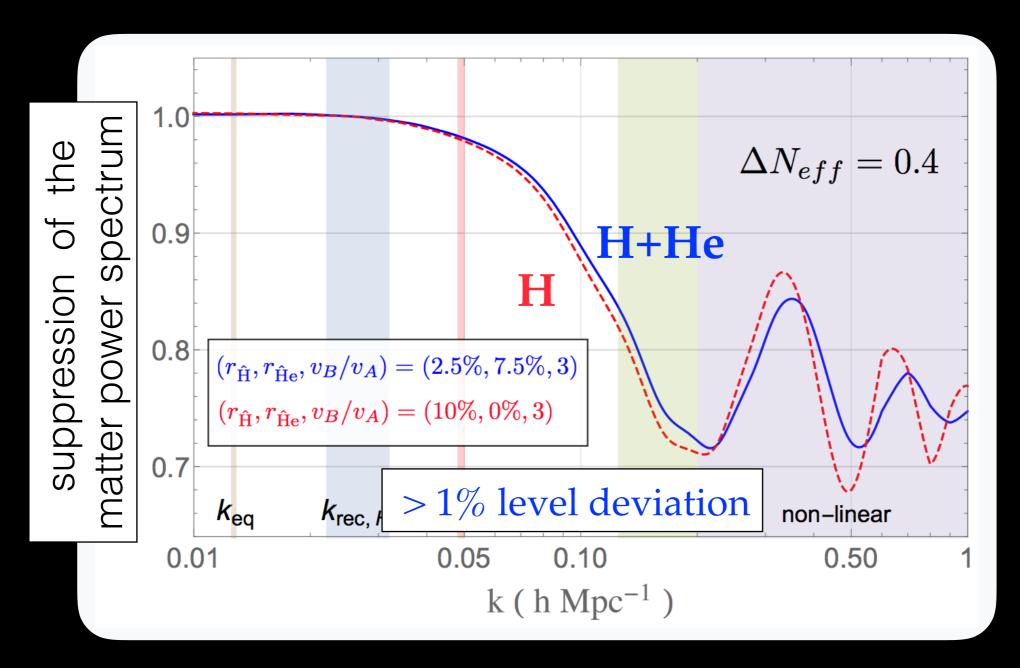
Horizon Entry

earlier

IV. dark photon - dark electron coupling



IV. identify the composition of mirror atoms



Here we choose the hydrogen abundance to match the average suppression at large k-modes

Short summary:

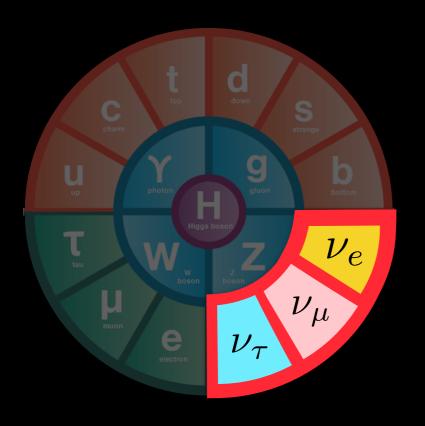
Even only couple to SM gravitationally,
MTH particles generate LSS signals in the form of
dark acoustic oscillations that stops after twin-recombination

Current bound requires ~< 10% of DM being MTH-baryons

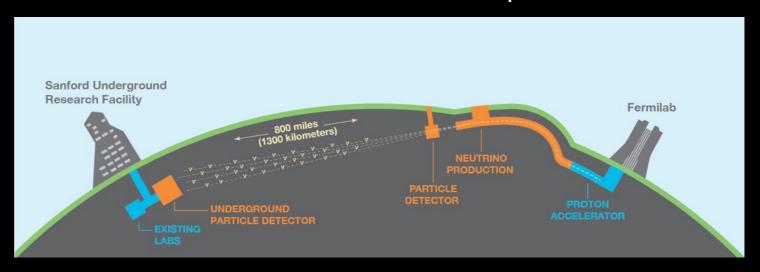
Can identify various details of the model

Example 2: neutrino mass / lifetime

We have very successful neutrino experiments to measure the **difference** between neutrino masses



Neutrino oscillation experiment



but what's the absolute value of their masses?

The best bound on neutrino mass comes from cosmology

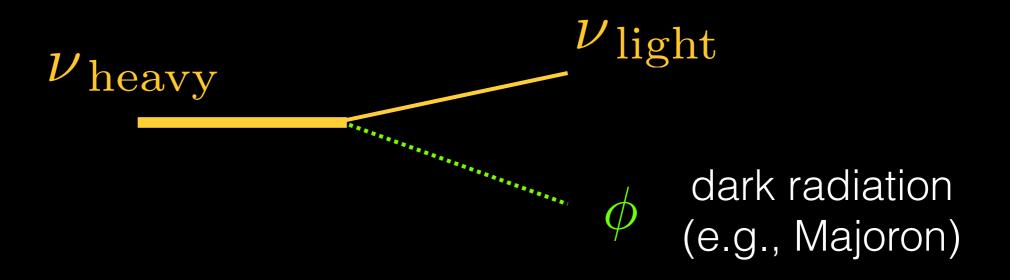
From Planck 2018 data

$$\sum m_{\nu} < 0.24 \, {\rm eV}$$

(~0.12 eV if including BAO)

This assumes neutrinos are stable particles

Neutrinos may not be as stable as predicted in the SM e.g., models explain the tiny neutrino mass



Can we measure the lifetime of neutrinos?

How stable are SM neutrinos?

Existing bounds on neutrino lifetime are very weak (for decay into invisible particles)

```
long-based line experiments \tau > 10^{-14}\,\mathrm{sec}
long-based line experiments \tau > 10^{-14}\,\mathrm{s
```

e.g., Archidiacono and Hannestad (2014) Escudero and Fairbairn (2019)

How stable are SM neutrinos?

Existing bounds on neutrino lifetime are very weak (for decay into invisible particles)

```
long-based line experiments \tau > 10^{-14}\,\mathrm{sec}

long-based line experiments \tau > 10^{-14}\,\mathrm{sec}

supernovae \tau > 8\,\mathrm{hrs} e.g., Frieman et al (1988)

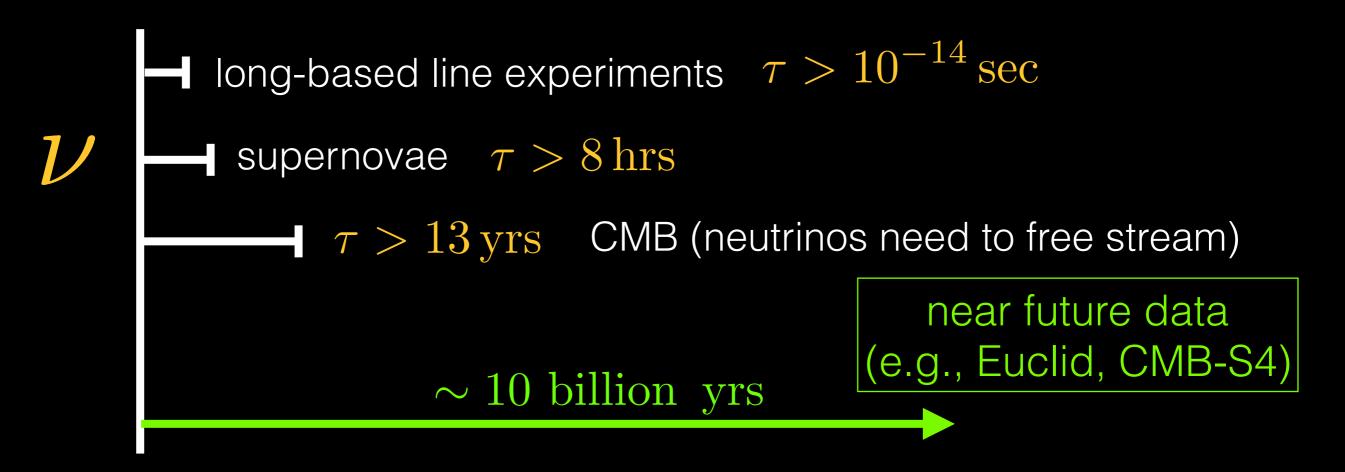
\tau > 13\,\mathrm{yrs} CMB (neutrinos need to free stream)
```

neutrino mass / lifetime are very hard to measure

- can we improve the bounds?
- can we probe neutrino decay?

How stable are SM neutrinos?

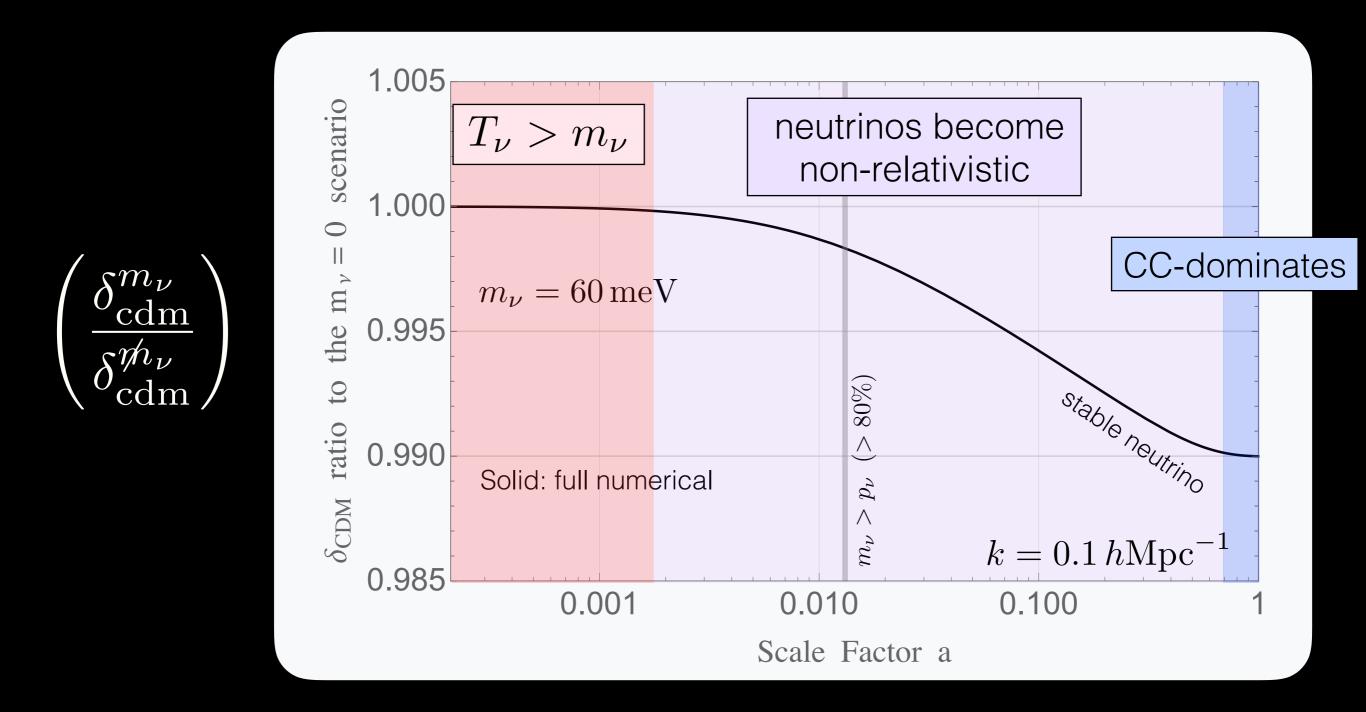
Existing bounds on neutrino lifetime are very weak (for decay into invisible particles)



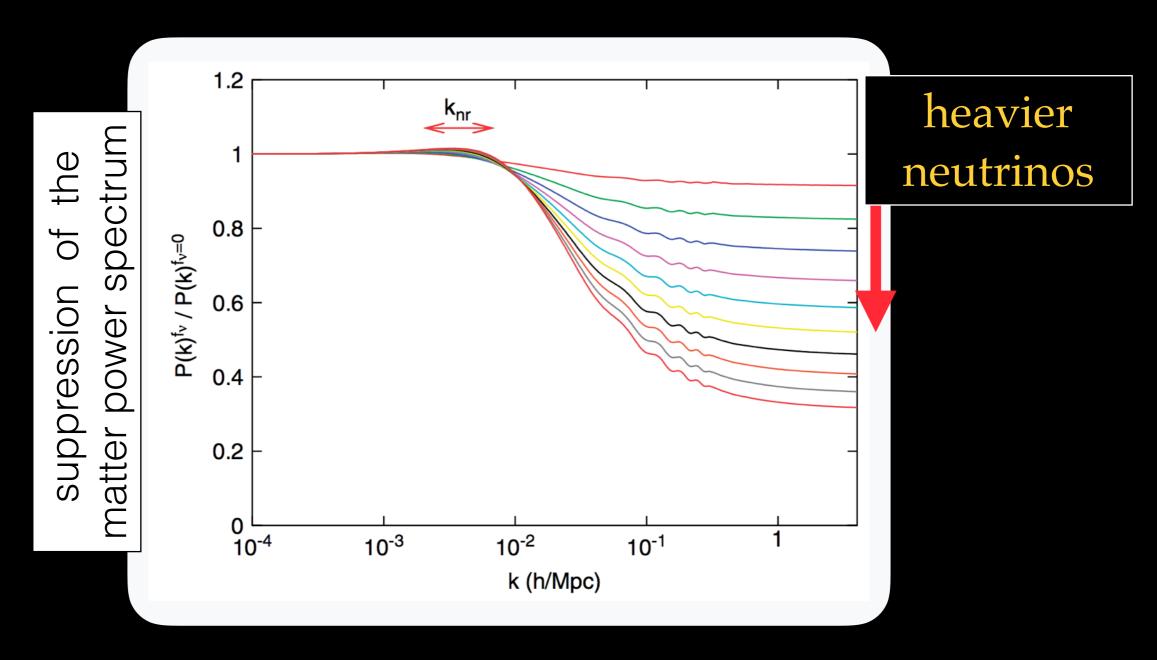
Chacko, Dev, Du, Poulin, YT (2019 one more soon), also Serpico (2007)

stable neutrinos with mass

Ratio of perturbation $\delta_{\mathrm{cdm}} = \frac{\delta \rho_{\mathrm{cdm}}}{\rho_{\mathrm{cdm}}}$ in redshift



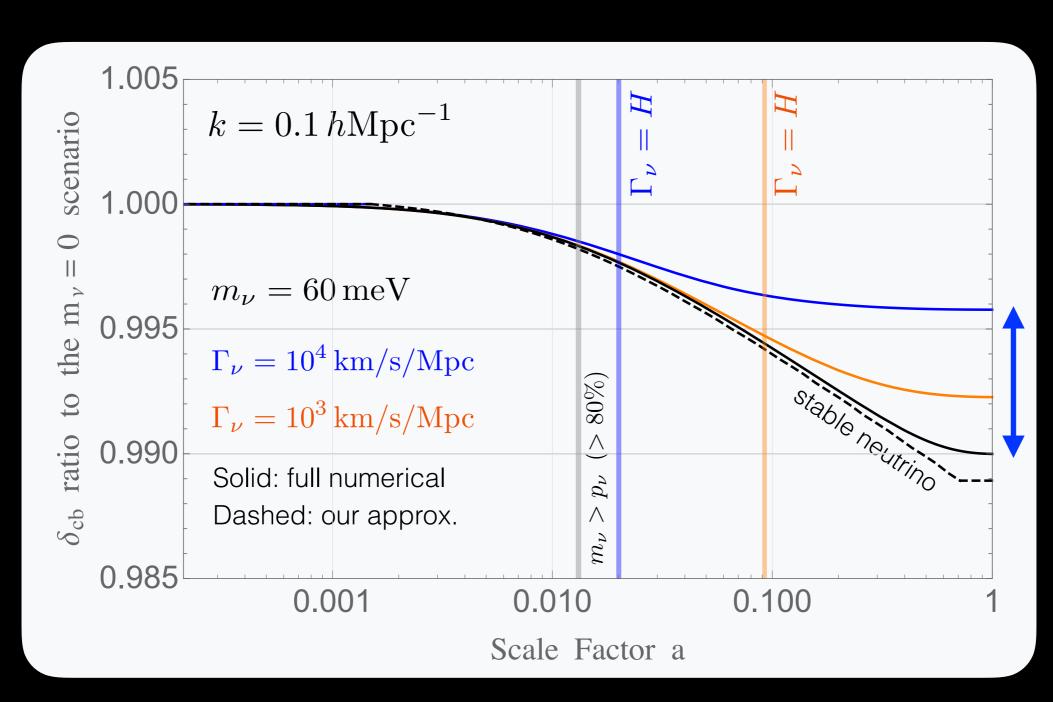
Suppression of the matter power spectrum



Lesgourgues and Pastor (2006)

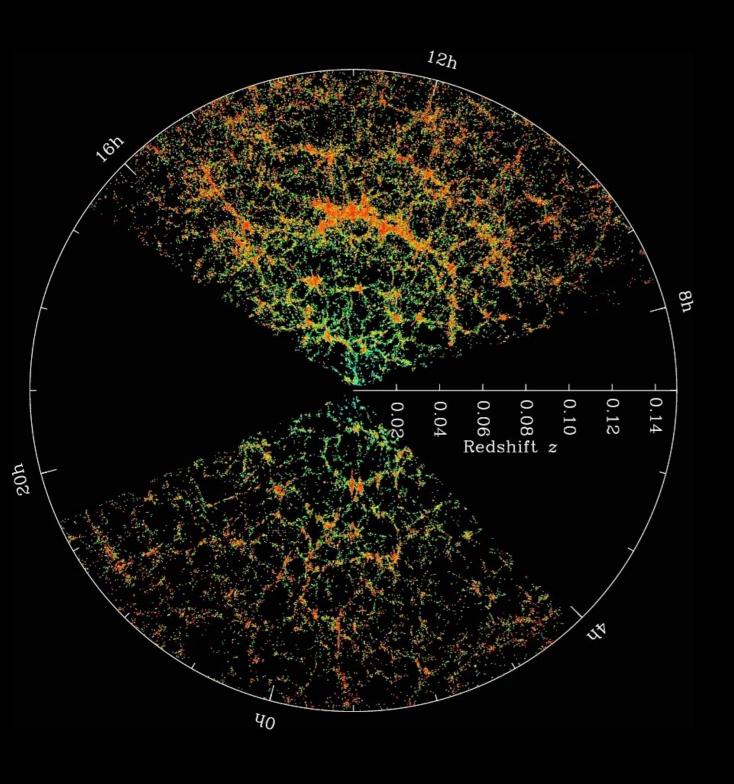
What if neutrinos decay?

``Larger" density contrast than expected

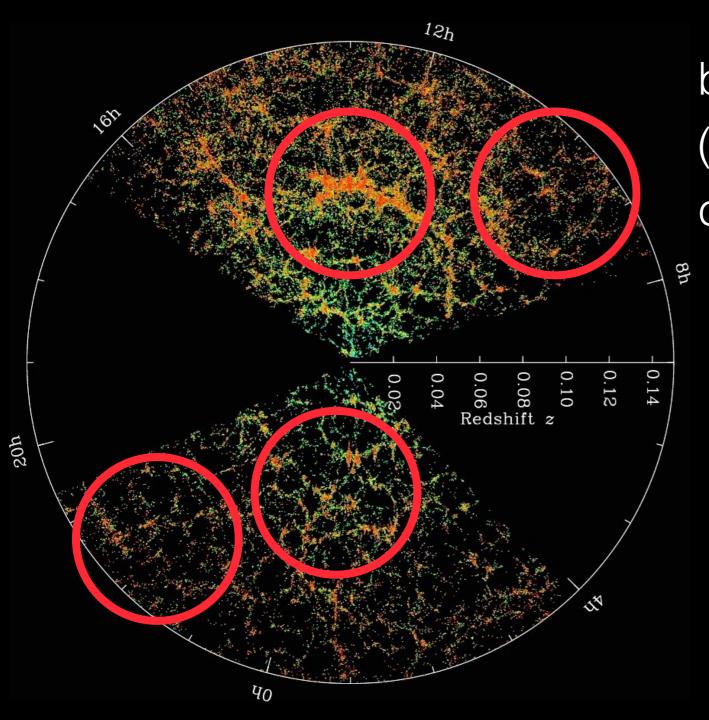


 $\rm km/s/Mpc \approx (10^3 \, Gyrs)^{-1}$

A cartoon picture of signal from stable massive neutrino

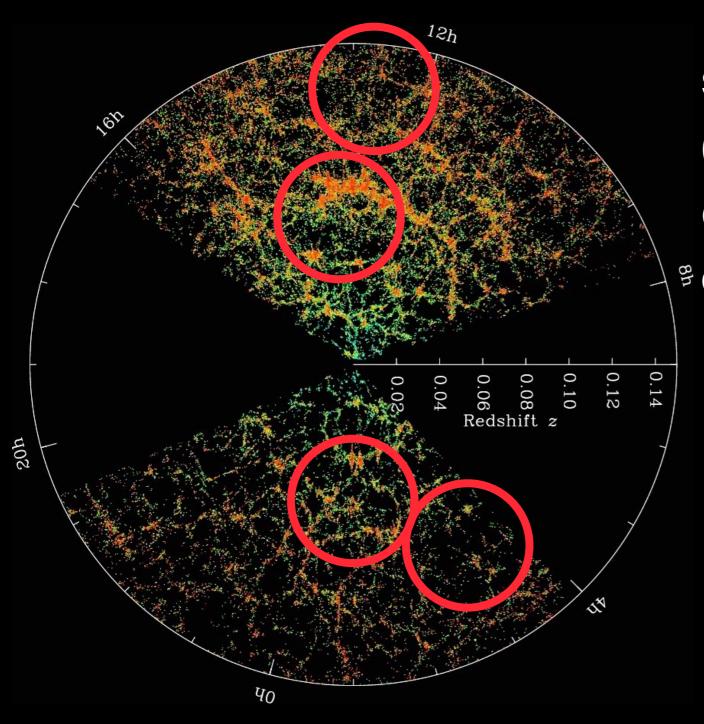


A cartoon picture of signal from stable massive neutrino



bigger structure
(after Nu is non-relativistic)
density contrast as expected

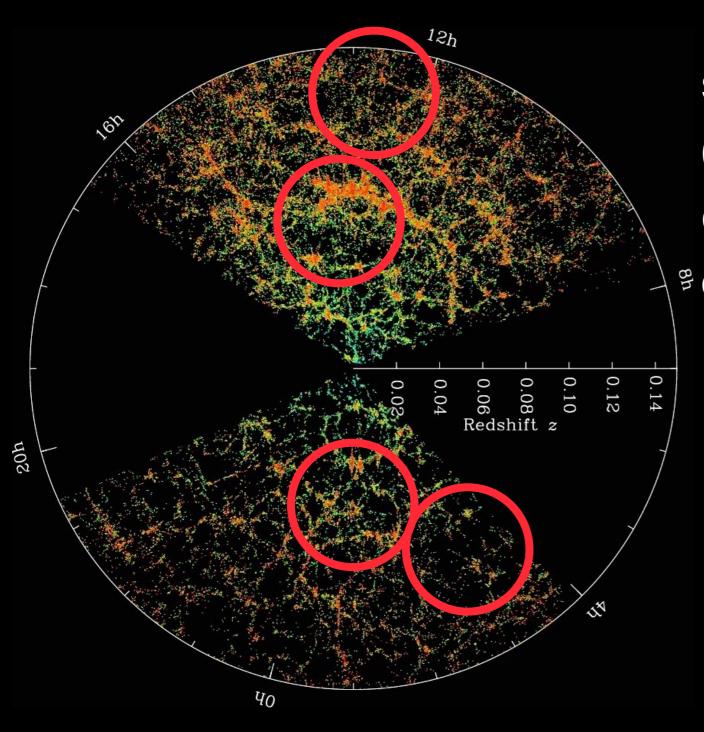
A cartoon picture of signal from stable massive neutrino



smaller structure
(since Nu was still relativistic)
density contrast is smaller than

gexpected

But if neutrinos decay!

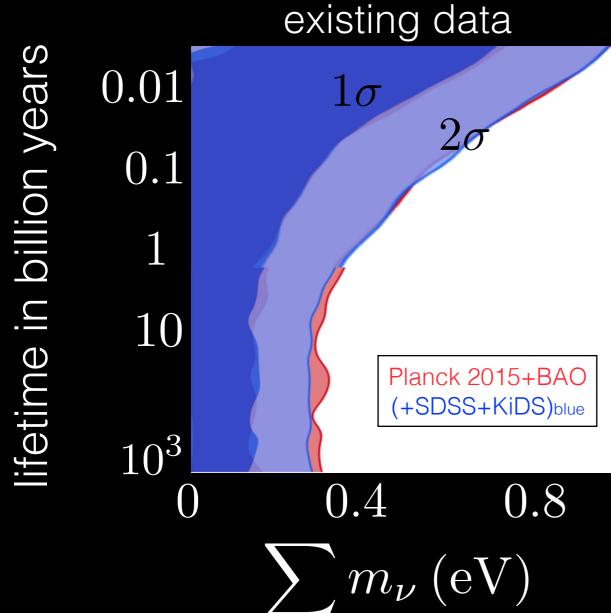


smaller structure
(since Nu was still relativistic)
density contrast is larger than
separated for massive neutrino

neutrino decay gives larger density perturbation compare to the stable Nu case

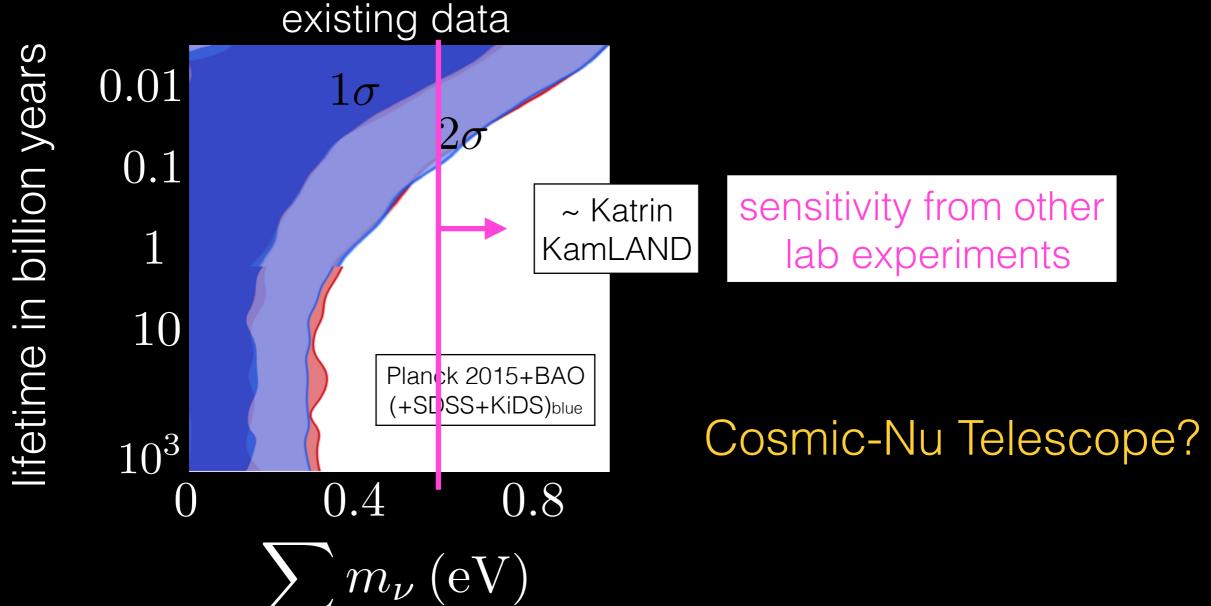
Current Large Scale Structure & Planck 2015

Chacko, Dev, Du, Poulin, **YT** (1909.05275)



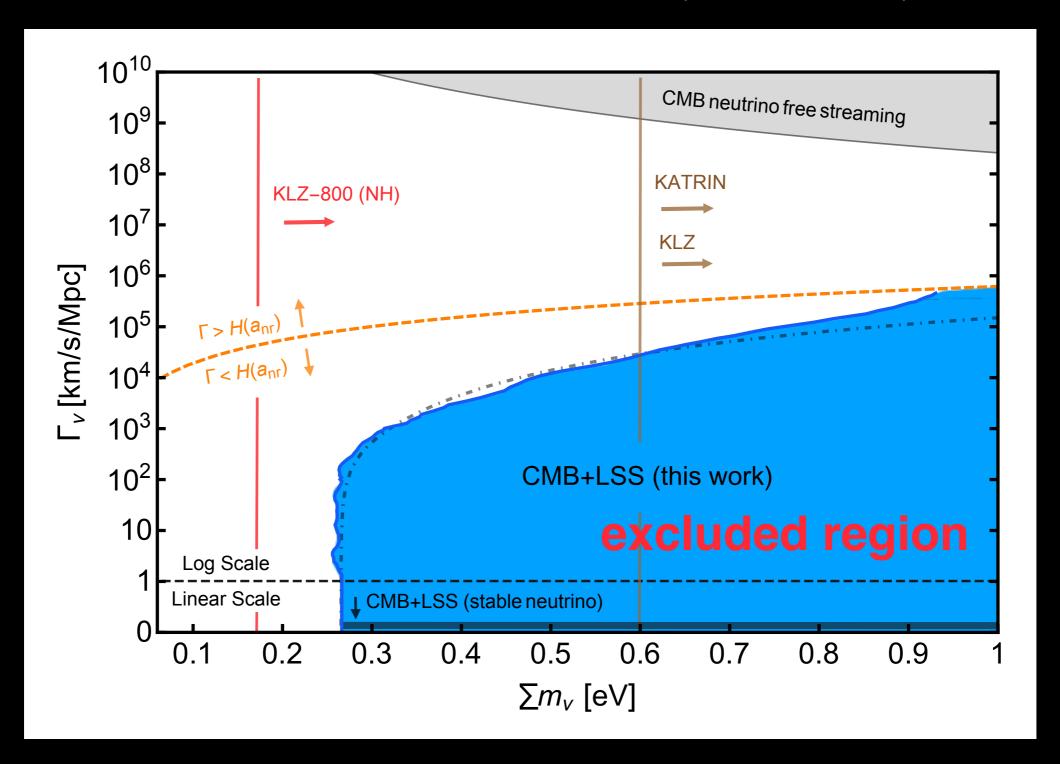
Current Large Scale Structure & Planck 2015

Chacko, Dev, Du, Poulin, **YT** (1909.05275)



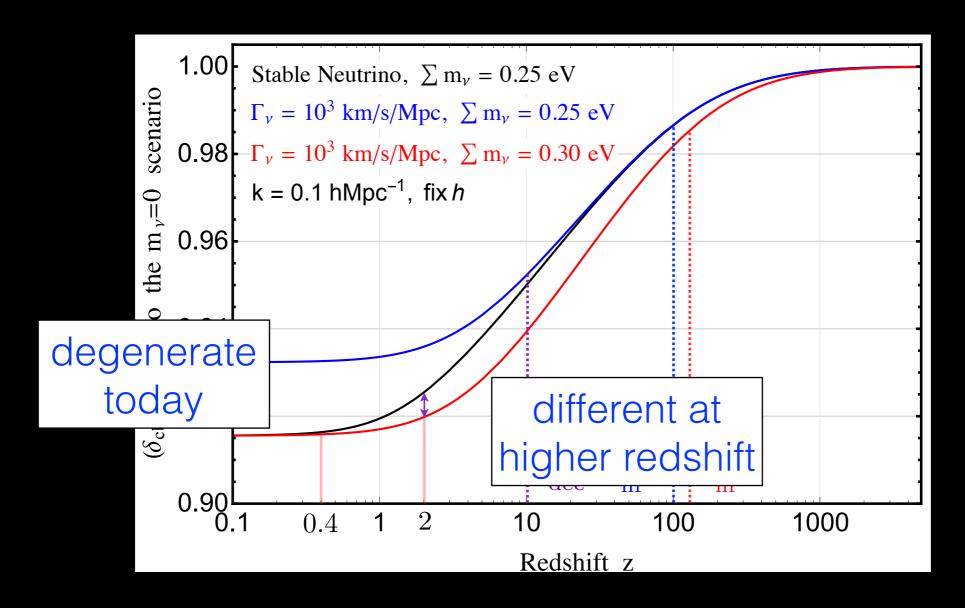
Current Large Scale Structure & Planck 2015

Chacko, Dev, Du, Poulin, **YT** (1909.05275)



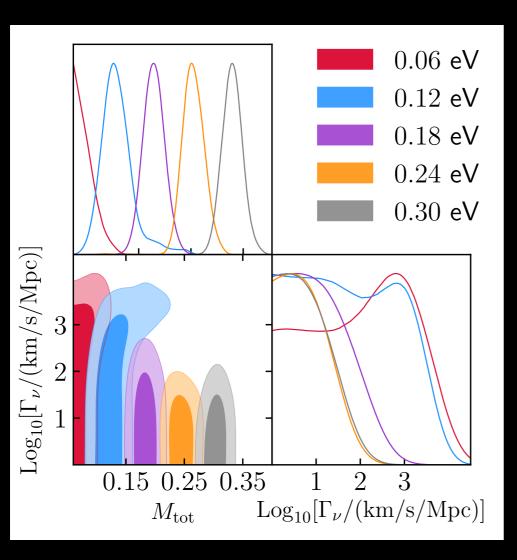
A degeneracy between mass & lifetime

Heavier neutrino decays can fake lighter neutrino signal

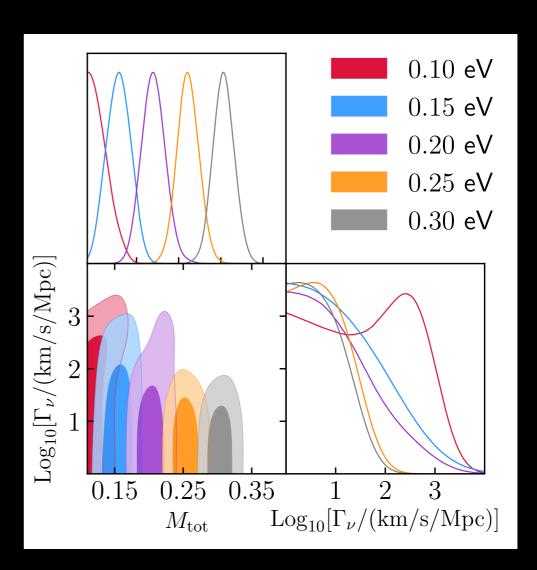


This is where the "time" (redshift) dependent measurement helps

e.g., near future Euclid data can break the degeneracy and set robust lifetime bound



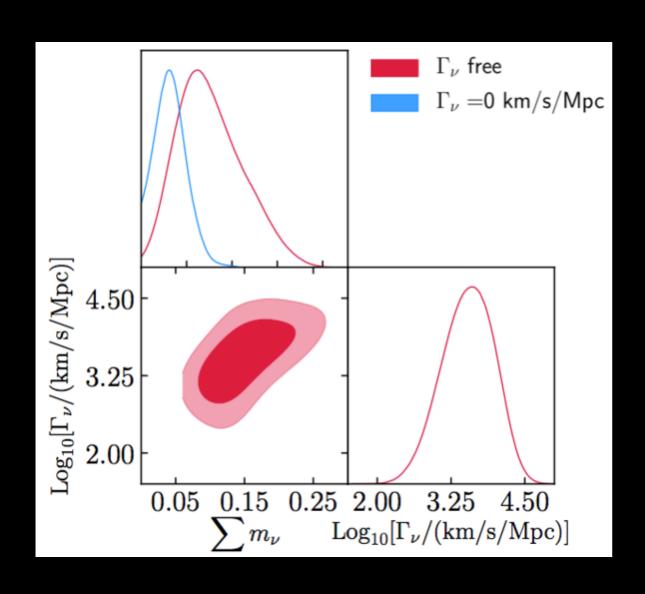
normal hierarchy

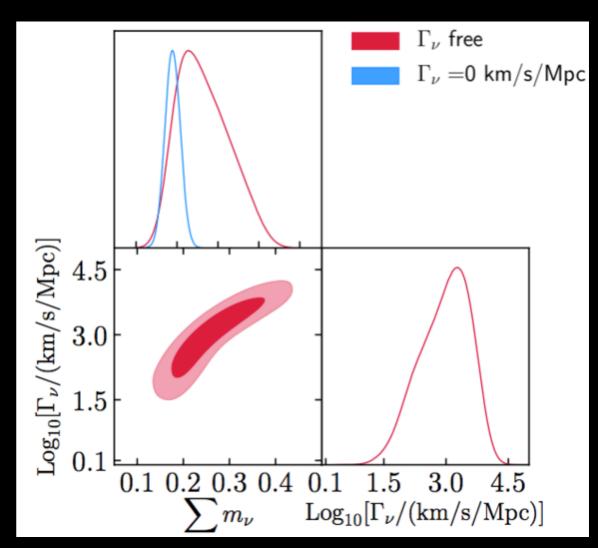


inverted hierarchy

Chacko, Dev, Du, Poulin, YT (preliminary result, to appear)

can even "measure" the decay rate if Nu decay





Chacko, Dev, Du, Poulin, YT (preliminary result, to appear)

Short summary:

massive neutrinos suppress the structure

LSS measurement can constrain/measure nu-mass

neutrino decay "increases" the structure

will be able to measure/constrain neutrino lifetime to the age of the universe time scale

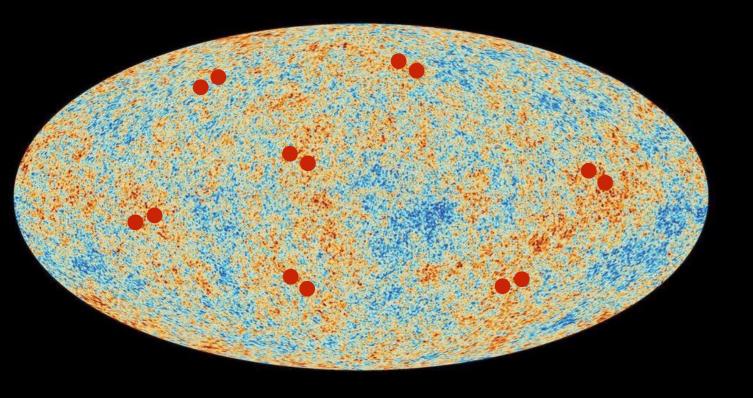
Conclusion

New physics that interacts very weakly, or even only gravitationally, to visible particles may be studied using precision cosmological data

As the much better quality data is coming to us, we need to get prepared and know what to look at

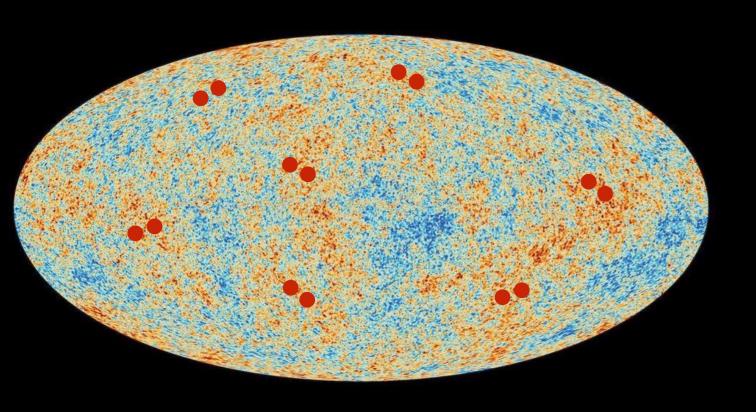
Discussion

e.g., pair-produced hot/cold spots



see also Maldacena (2015)

e.g., pair-produced hot/cold spots

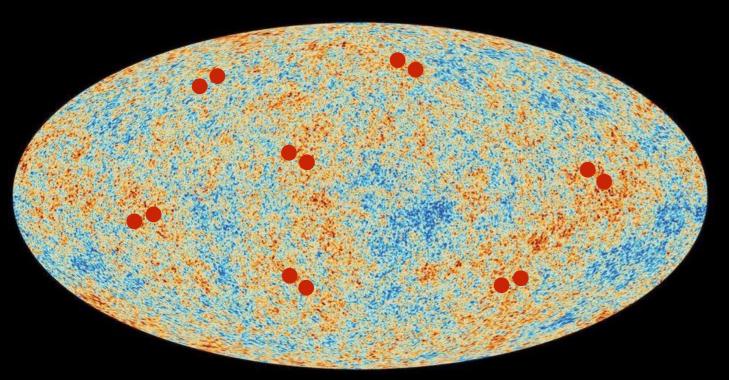


cannot see this in power spectrum (CI^TT) maybe machine learning?

see also Maldacena (2015)

e.g., pair-produced hot/cold spots

How to make this?



- heavy particle production during inflation
- mass sources larger perturbation

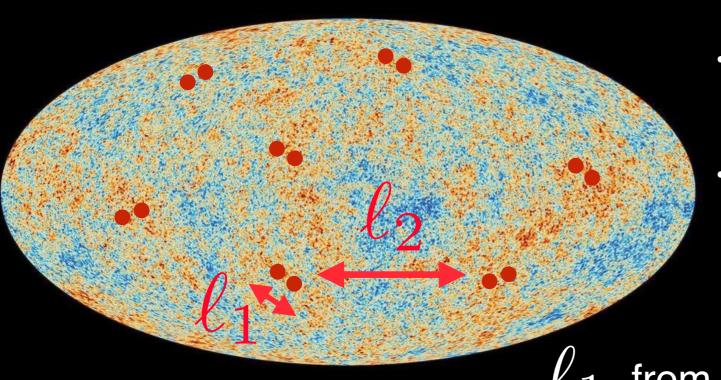
$$\mathcal{L} \subset \int \frac{d\eta}{H} m(\eta) \partial_{\eta} \xi$$

need time-dependent mass $m \gg H \gg m \sim H \gg m \sim Mpl \gg H$

$$\langle \xi \rangle \sim \left(\frac{m(\eta = x)}{\sqrt{\epsilon} M_{pl}} \right) \left(\frac{H}{\sqrt{\epsilon} M_{pl}} \right)$$

e.g., pair-produced hot/cold spots

How to make this?



- heavy particle production during inflation
- mass sources larger perturbation

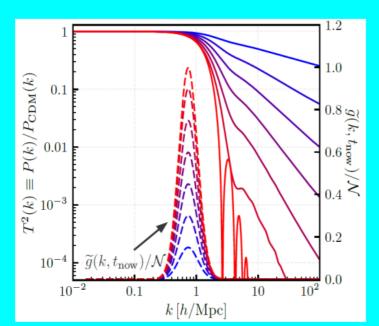
 ℓ_1 from horizon size at particle production

 ℓ_2 determined by the production probability

Transfer function => DM phase space distribution

Pushing this further,

we can even conjecture a specific function η !





$$\left| \frac{d \log T^2}{d \log k} \right| \approx [F(k)]^2 + \frac{3}{2}F(k)$$

approximate relation holds to very high precision!

Our conjecture then takes the non-trivial form

$$\frac{\widetilde{g}(k)}{\mathcal{N}} \approx \frac{1}{2} \left(\frac{9}{16} + \left| \frac{d \log T^2}{d \log k} \right| \right)^{-1/2} \left| \frac{d^2 \log T^2}{(d \log k)^2} \right|$$

This would allow us to "resurrect" g(k) from the transfer function $T^2(k)$!

slide credit: Keith Dienes

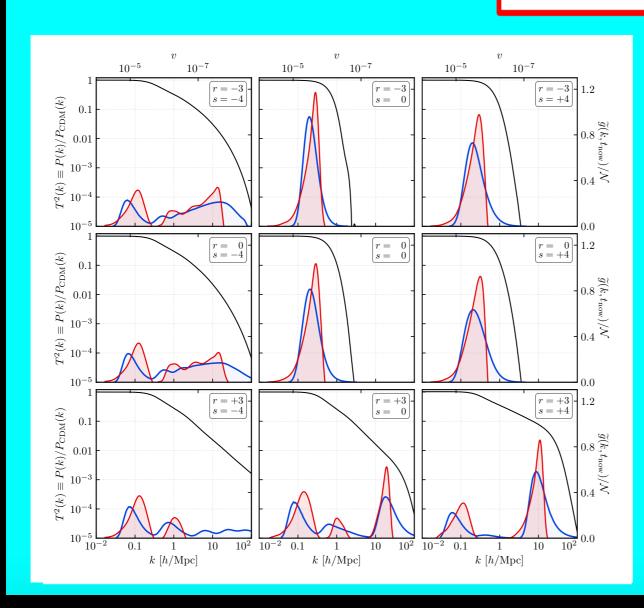
Dienes, Huang, Kost, Su, Thomas

Transfer function => DM phase space distribution

Finally, to what extent can we "resurrect" the dark-matter phase-space distribution from the transfer function?

Recall our conjecture....

$$\frac{\widetilde{g}(k)}{\mathcal{N}} \approx \frac{1}{2} \left(\frac{9}{16} + \left| \frac{d \log T^2}{d \log k} \right| \right)^{-1/2} \left| \frac{d^2 \log T^2}{(d \log k)^2} \right|$$



Blue outline = original k-space DM distribution

<u>Pink shaded</u> = reconstruction directly from transfer function

slide credit: Keith Dienes

Dienes, Huang, Kost, Su, Thomas