## Gravitational Wave Reflectivity from an Ultralight Bosonic Cloud

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- Ultralight bosons/bosonic cloud formation
- Probing ultralight boson mass
- Echo signatures of bosonic cloud due to reflectivity
- Concluding thoughts

- Ultralight boson: dark matter candidate, scalar
- mass  $\approx O(10^{-22}) \mathrm{eV}$
- Superradiant condition for Kerr BH 0  $< \omega_B < m\Omega_H$
- Negative energy and angular momentum flux at horizon
- Energy and angular momentum of BH goes to bosonic cloud
- $\bullet~\mbox{For massive boson}$   $\rightarrow~\mbox{bound state}$
- "Gravitational atom"



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- Stochastic background<sup>1</sup>
- "Holes" in spin-mass plane<sup>2</sup>
- Our focus: echoes from bosonic cloud

<sup>1</sup>Leo Tsukada et al. "First search for a stochastic gravitational-wave background from ultralight bosons". In: *Physical Review D* 99.10 (2019), p. 103015. <sup>2</sup>Richard Brito, Vitor Cardoso, and Paolo Pani. "Superradiance". In: *Lect. Notes Phys* 906.1 (2015), pp. 1501–06570.

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#### Echoes in GW searches



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• Can bosonic cloud produce a similar effect?

<sup>3</sup>Jahed Abedi, Hannah Dykaar, and Niayesh Afshordi. "Echoes from the Abyss: Tentative evidence for Planck-scale structure at black hole horizons". In: *Physical Review D* 96.8 (2017), p. 082004.

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### Calculating the Bosonic Cloud Reflectivity

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Bosonic cloud wavefunction for l = m = 1 mode

$$\varphi(t,r,\theta,\phi) = \sqrt{\frac{3M_s}{4\pi i M}} (M\mu)^3 \mu r e^{-M\mu^2 r/2} \sin\theta \cos(\phi - \mu t), \qquad (1)$$

- M: BH mass
- $M_s$ : Total mass of bosonic cloud extracted in superradiance from BH
- $\mu = \frac{m_s}{\hbar}$ : boson frequency for boson mass  $m_s$ Stress tensor of wavefunction given by

$$T_{\mu\nu} = \nabla_{\mu}\varphi\nabla_{\nu}\varphi - \frac{1}{2}g_{\mu\nu}\mathcal{L}$$
<sup>(2)</sup>

• How does stress tensor modify gravitational waveforms?

### **Teukolsky Equation**

- Black hole perturbations method: Teukolsky equation
- Curvature perturbation
- Master equation for GW waveform
- Separable

$$\begin{split} \left[\frac{(r^2+a^2)^2}{\Delta} - a^2\sin^2\theta\right] \frac{\partial^2\psi}{\partial t^2} + \frac{4Mar}{\Delta}\frac{\partial^2\psi}{\partial t\partial\varphi} + \left[\frac{a^2}{\Delta} - \frac{1}{\sin^2\theta}\right]\frac{\partial^2\psi}{\partial\varphi^2} \\ &- \Delta^{-s}\frac{\partial}{\partial r}\left(\Delta^{s+1}\frac{\partial\psi}{\partial r}\right) - \frac{1}{\sin\theta}\frac{\partial}{\partial\theta}\left(\sin\theta\frac{\partial\psi}{\partial\theta}\right) - 2s\left[\frac{a(r-M)}{\Delta} + \frac{i\cos\theta}{\sin^2\theta}\right]\frac{\partial\psi}{\partial\varphi} \\ &- 2s\left[\frac{M(r^2-a^2)}{\Delta} - r - ia\cos\theta\right]\frac{\partial\psi}{\partial t} + (s^2\cot^2\theta - s)\psi = 4\pi\Sigma T \,. \end{split}$$

• Project stress-tensor with null vectors, set a = 0

$$T_{\bar{m}\bar{m}} = T_{\mu\nu}\bar{m}^{\mu}\bar{m}^{\nu} \approx \frac{1}{2\omega^2}\psi_4[\partial_{\alpha}\varphi\partial^{\alpha}\varphi + \mu^2\varphi^2]$$
(3)

$$\bar{m}^{\mu} = \frac{1}{r\sqrt{2}}(0, 0, 1, -i\sin^{-1}\theta)$$
(4)

• 
$$\Sigma = r^2 + a^2 \cos^2 \theta$$

Eigenfunctions

$$\psi_{4} = \rho^{4} \sum_{l,m} \int \frac{\mathrm{d}\tilde{\omega}}{\sqrt{2\pi}} Y_{lm}(\theta,\phi) e^{-i\tilde{\omega}t} R_{lm\tilde{\omega}}(r)$$
(5)

- *l* = *m* = 2 mode
- Extract Teukolsky source term

$$\tilde{T}(t,r,\Omega) \sim -\frac{1}{4}\rho^{8}\bar{\rho}\Delta^{2}\hat{J}_{+}[\rho^{-4}\hat{J}_{+}(\rho^{-2}\bar{\rho}T_{\bar{m}\bar{m}})]$$
(6)

• 
$$\hat{J}_{+} = \partial_r - \Delta^{-1}(r^2\partial_t)$$
  
•  $R(r,\omega) \propto e^{i\omega r}$   
• Ansatz:  $\frac{\partial R}{\partial r} \sim i\omega R(r)$ 

- Radial and angular equation decouple
- Radial: modified wave equation with extra potential  $\frac{\partial^2 R}{\partial r^2} + (\omega^2 - V(r) - Q(r))R(r,\omega) = 0$
- V(r) usual Schwarzschild potential in Teukolsky equation
- Q(r) extra contribution from bosonic cloud
- How does this potential affect waveforms?
- GW Echoes/scattering

(7)

### Conceptual Idea



Asymptotic solutions

$$\Psi(t,x) = \left\{ \begin{array}{ll} e^{-i\omega x} + \mathcal{R}(\omega)e^{i\omega x}, & \text{for } x \to \infty \\ e^{-i\omega x}, & \text{for } x \to -\infty \end{array} \right\}$$
(8)

• Full solution sum of free wave solution plus potential dependant term

$$\tilde{\Psi}(\omega, x) = \tilde{\Psi}_0(\omega, x) + \sum_{n=1}^{\infty} \tilde{\Psi}_n(\omega, x)$$
(9)

• Powers in  $\frac{V}{\omega^2}$ 

• To first order in boson cloud potential

$$\mathcal{R}(\omega) \sim \int \mathrm{d}x' \frac{e^{-i\omega x'}}{2i\omega} \frac{r'^2}{\Delta} Q(r')$$
 (10)

- Integrate from light ring r = 3M to  $\frac{1}{M^2 \mu^3}$
- Peak of Q at  $r_{peak} \sim rac{1}{M\mu^2}$
- Numerical calculation: vary  $\mu, \omega$

### Reflectivity

• Fixed QNM frequency



•  $\mu^8$  dominates

# Reflectivity( $\omega$ )

• Re[ $\mathcal{R}(\omega)$ ] • Fixed  $M\mu = 0.1$ 



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- Invert into time-domain for waveform
- How does potential change echo model?
- Draw constraints from LIGO events
- Check degeneracy with other models
- Other signatures?

# The End

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Image: A mathematical states of the state