

Cosmic Axion Background: the QCD axion as a hot relic

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talk @ PASCOS 2021.

¹In collaboration with R.Z. Ferreira, F. Rompineve, F. D'Eramo, F. Arias-Aragon, J.L. Bernal, L.Merlo.

QCD Axion through N_{eff}

Axions as Hot
Relics

The QCD
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Axions via
Gluons

Axion via
Quarks

Axion via
Leptons

Axions via
Pions

The **QCD Axion** (a) is a very light particle that

- Solves the “**Strong CP problem**” via coupling to gluons

$$\mathcal{L}_a = \frac{\alpha_s}{8\pi} \frac{a}{f} G_{\mu\nu} \tilde{G}^{\mu\nu}$$

- boundary term sensitive to QCD Instantons,
 - 1 Induces a potential $V(a) \propto \cos(a/f)$;
 - 2 \implies Drives \cancel{CP} to zero
 - 3 \implies Axion mass $m_a \approx 0.57 \left(\frac{10^7 \text{ GeV}}{f} \right) \text{ eV}$

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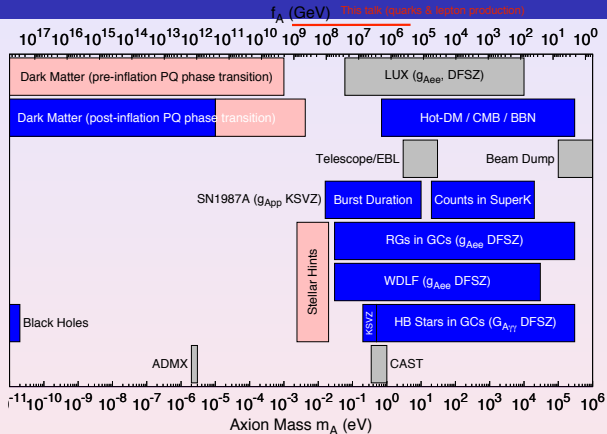
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- Bounds on $f \Leftrightarrow$ bounds on m_a

Axion: constraints



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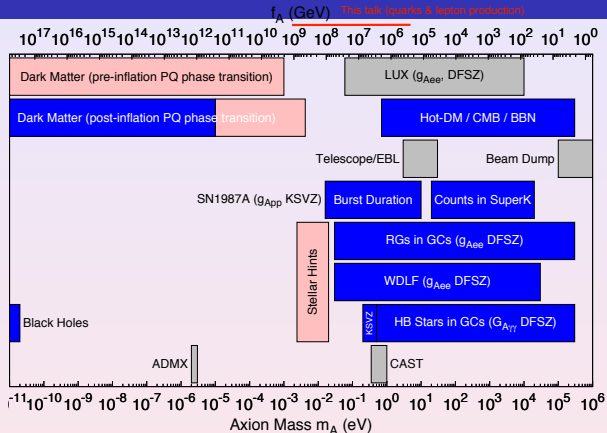
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- **Caveat:** Constraints based on individual couplings with e, γ , nucleons... Expected $\mathcal{O}(1/f)$, but model dependent.

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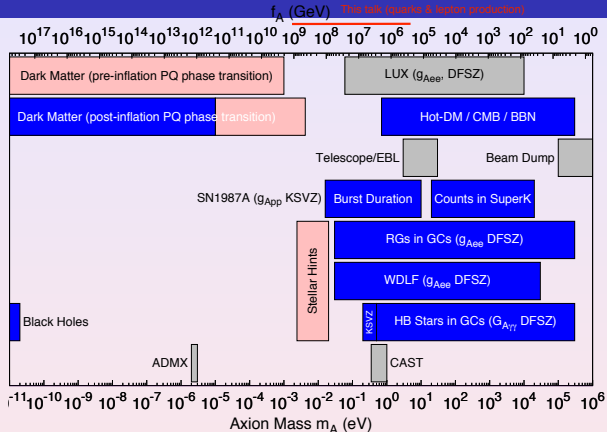
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- **Caveat:** Constraints based on individual couplings with e, γ , nucleons... Expected $\mathcal{O}(1/f)$, but model dependent.
- Small $m_a \ll \mathcal{O}(\text{eV}) \implies$ acts as **Dark Radiation**, visible in CMB (*Cosmic Axion Background*)

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- **Axion** effective lagrangian:
 - ① May couple with continuous shift symmetry with **all SM**
 - ② **Only** breaking: Instanton-induced (tiny) mass

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- Due to $\frac{\alpha_s}{8\pi} \frac{a}{f} G_{\mu\nu} \tilde{G}^{\mu\nu}$ QCD Axions can be produced by gluon scatterings in the Early Universe

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- Can be produced at high T and decouples at $T \lesssim T_{DEC}$
→ **hot relic (dark radiation)**
(M.Turner, 1987; Masso, F. Rota, and G. Zsembinski, 2003, Salvio, Strumia, Xue, 2014)

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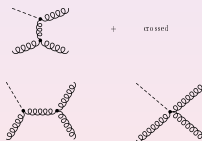


Figure: (Massò et al. Phys.Rev. D66 (2002).).

$$\Gamma_s \equiv \langle \sigma v \rangle \cdot n_g^{EQ} = \left(\frac{\alpha_s}{2\pi f} \right)^2 g_s^2 \cdot T^3$$

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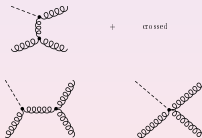


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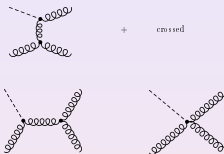


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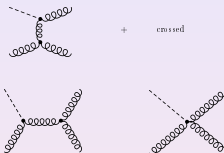


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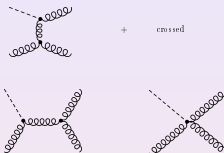


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- Example:

① $f = 10^8 \text{ GeV} \implies T_{DEC} \approx \text{TeV}$

② $f = 10^9 \text{ GeV} \implies T_{DEC} \approx 100 \text{ TeV}$

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- \implies **Observable by CMB** (and BBN)
(mostly affects expansion rate, Matter-Radiation equality...)

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- $N_{\text{eff}} = 3.046 + \Delta N_{\text{eff}}$

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- $$\Delta N_{\text{eff}} \approx \frac{13.6}{g_{*,\text{DEC}}^{4/3}}.$$

ΔN_{eff} diluted by $g_{*,DEC}$

- Abundance ΔN_{eff} diluted if total number of relativistic species in the plasma $g_{*,DEC}$ is large

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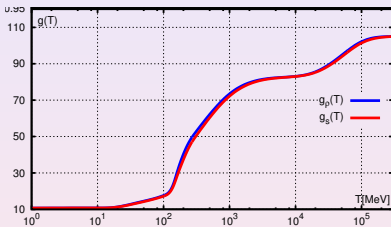
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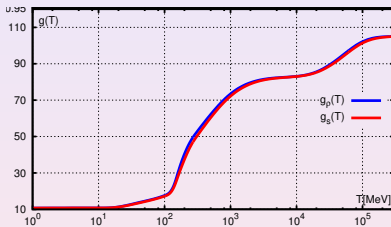
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- Abundance ΔN_{eff} diluted if total number of relativistic species in the plasma $g_{*,\text{DEC}}$ is large



- $$\Delta N_{\text{eff}} \approx \frac{13.6}{g_{*,\text{DEC}}^{4/3}}$$
- If $T_{\text{DEC}} \gg 100 \text{ GeV}$, $\Rightarrow g_{*,\text{DEC}} \geq 106.75$
- $\Rightarrow \Delta N_{\text{eff}} \lesssim 0.027$ (only upper bound!)
(marginally detectable, 1σ , by CMB-Stage 4 experiments)

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- If $f \lesssim 10^9\text{-}10^{10}$ GeV dominant channels can be via **quarks & leptons**² with **$T_{\text{DEC}} \leq \text{Electroweak scale}$**

²A.N. & R.Z.Ferreira, PRL 2018; D'Eramo, Ferreira, A.N., Bernal JCAP 2018, F. Arias-Aragón et al. JCAP 2021.

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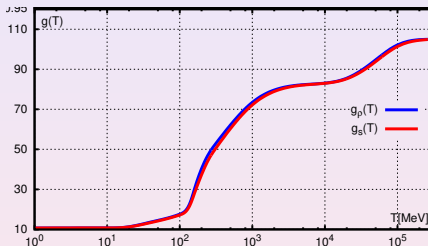
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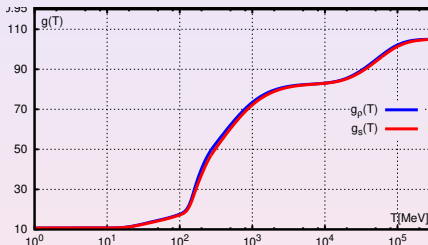
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ADVANTAGES:

- 1 g_*^{SM} is smaller \Rightarrow **larger N_{eff}**
- 2 Here we are confident on $g_*^{\text{SM}} \Rightarrow$ **Precise predictions**
- 3 Lower $f \Rightarrow$ more accessible by **direct** searches
(**CAST, IAXO**)

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- If a is directly coupled to SM **heavy quarks** (c, b, t):

$$\mathcal{L}_{a-q} = \partial_\mu a \sum_i \frac{c_i}{2f} \bar{q}_i \gamma^\mu \gamma^5 q_i,$$

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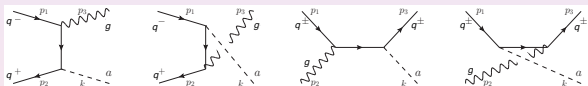
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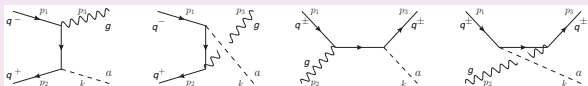
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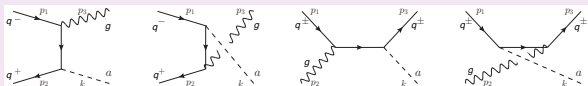
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- Indeed:
 - This coupling can be **rotated away** $q \rightarrow e^{i \frac{c_i a}{2f} \gamma^5} q$
 - But it **reappears in the mass term** $m_q \bar{q} e^{i \frac{c_i a}{f} \gamma^5} q$

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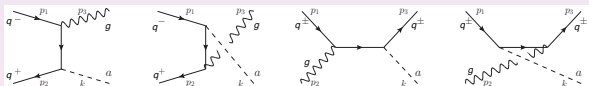
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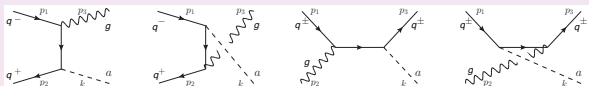
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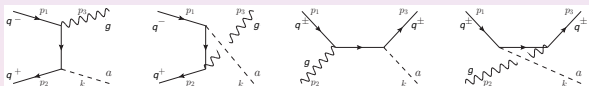
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- Axions produced **dominantly** via quarks

$$1 \text{ GeV} \lesssim T \lesssim 100 \text{ GeV}$$

- Range** $10^9 \text{ GeV} \gtrsim f/c_i \gtrsim 10^7 \text{ GeV}$ ³
(partly in tension with SN bounds, if all $c_i = 1$)

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- Ratio peaks at $T \approx m_q$

- Axions produced **dominantly** via quarks

$$1 \text{ GeV} \lesssim T \lesssim 100 \text{ GeV}$$

- Range** $10^9 \text{ GeV} \gtrsim f/c_i \gtrsim 10^7 \text{ GeV}$ ³
(partly in tension with SN bounds, if all $c_i = 1$)

- Interesting for **direct detection** (e.g. **IAXO**),
 $m_a \approx 10^{-1} \sim 10^{-3} \text{ eV}$,

³R.Ferreira & A.N., PRL 2018. See also Turner PRL 1987, Brust et al. JHEP 2013, Baumann et al. PRL 2016.

QCD Axion through N_{eff}

Axions as Hot
Relics

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The QCD
Axion

Axions via
Gluons

Axion via
Quarks

Axion via
Leptons

Axions via
Pions

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QCD Axion through N_{eff}

Axions as Hot
Relics

The QCD
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Axions via
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Pions

- $g_{*,DEC}$ is smaller at $1 \text{ GeV} \lesssim T \lesssim 100 \text{ GeV}$
- **Prediction:** larger $N_{\text{eff}} \lesssim 0.045$ (*Not just upper bound!*)

QCD Axion through N_{eff}

Axions as Hot
Relics

The QCD
Axion

Axions via
Gluons

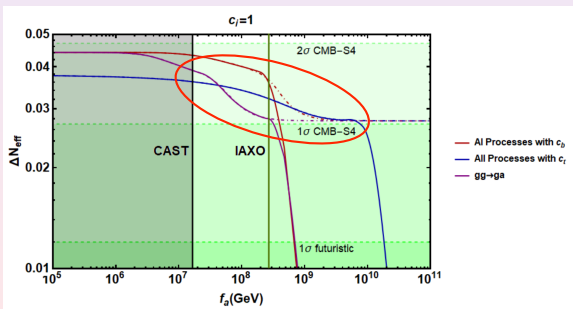
Axion via
Quarks

Axion via
Leptons

Axions via
Pions

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- Prediction:** larger $N_{\text{eff}} \lesssim 0.045$ (*Not just upper bound!*)
- Solving Boltzmann equations for n_a :**

(R.Ferreira & A.N., PRL 2018; F.Arias-Aragon et al. JCAP, 2021)



$$10^9 \text{ GeV} \gtrsim f/c_i \gtrsim 10^7 \text{ GeV}, \quad 5 \times 10^{-3} \text{ eV} \lesssim m_a \lesssim 0.5 \text{ eV}$$

($c_i = 1$, for QCD Axion)

QCD Axion through N_{eff}

Axions as Hot
Relics

The QCD
Axion

Axions via
Gluons

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Axion via
Leptons

Axions via
Pions

- Potentially larger for c -quark: $N_{\text{eff}} \lesssim 0.05 - 0.06$
(but uncertain)

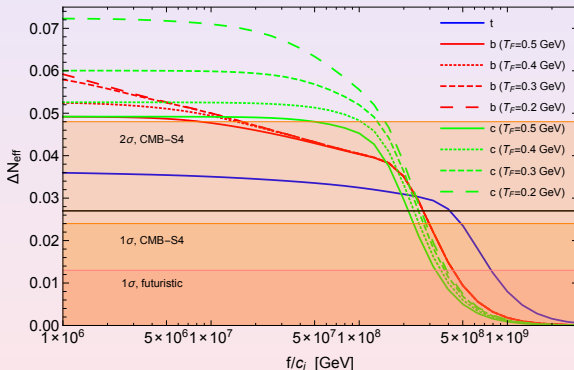


Figure: R.Ferreira & A.N., PRL 2018.

Hot Axions via Leptons

Axions as Hot
Relics

The QCD
Axion

Axions via
Gluons

Axion via
Quarks

Axion via
Leptons

Axions via
Pions

- The same can be done with **leptons** (μ and τ)⁴
- a -electron uninteresting (strongly constrained)

⁴F.D'Eramo, A.N., R.Z.Ferreira, J.L.Bernal, JCAP 2018.

Hot Axions via Leptons

Axions as Hot
Relics

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Axions via
Gluons

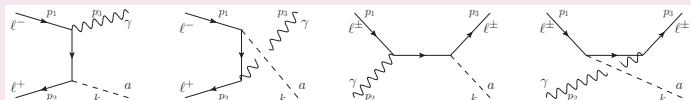
Axion via
Quarks

Axion via
Leptons

Axions via
Pions

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$$\mathcal{L}_{a-\ell} = \partial_\mu a \sum_i \frac{c_i}{2f} \bar{\ell}_i \gamma^\mu \gamma^5 \ell_i,$$



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Hot Axions via Leptons

Axions as Hot Relics

The QCD Axion

Axions via Gluons

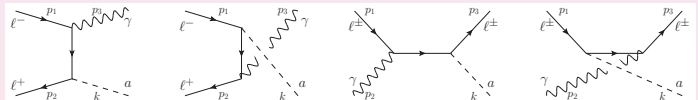
Axion via Quarks

Axion via Leptons

Axions via Pions

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- Slightly smaller f/c_ℓ
- Ratio peaks at $T \approx m_\ell \implies$ **Larger N_{eff}**

⁴F.D'Eramo, A.N.,R.Z.Ferreira, J.L.Bernal, JCAP 2018.

Hot Axions via Lepton Scatterings

Axions as Hot
Relics

The QCD
Axion

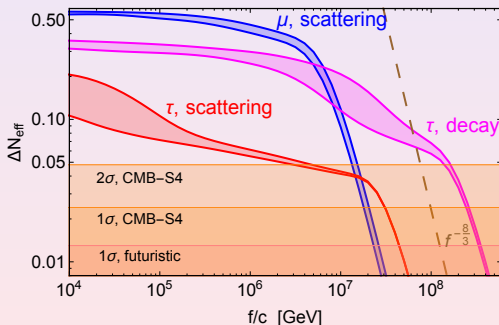
Axions via
Gluons

Axion via
Quarks

Axion via
Leptons

Axions via
Pions

- Smaller $f/c_i \lesssim \text{few} \cdot 10^7 \text{ GeV}$
- Ratio peaks at $T \approx m_\ell \Rightarrow \text{Larger } N_{\text{eff}}$



- Caveat: μ scattering constrained by SN cooling at $f/c_\mu \gtrsim 10^8 \text{ GeV}$ (Bolling et al. PRL 2020, Croon et al. JHEP 2021)

Hot Axions via Lepton Decays

Axions as Hot
Relics

- $a - \ell$ interaction can be **flavor non-diagonal**

$$\mathcal{L}_{a-\ell} = \partial_\mu a \sum_{\ell \neq \ell'} \bar{\ell}' \gamma^\mu (\mathcal{V}_{\ell'\ell} + \mathcal{A}_{\ell'\ell} \gamma^5) \ell + \text{h.c.} ,$$

- Decays $\tau \rightarrow \mu + a, \tau \rightarrow e + a$

The QCD
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Axion via
Quarks

Axion via
Leptons

Axions via
Pions

Hot Axions via Lepton Decays

Axions as Hot
Relics

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Axion via
Quarks

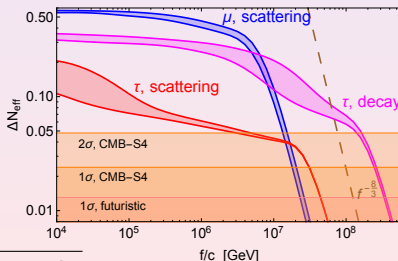
Axion via
Leptons

Axions via
Pions

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$$(c_{\ell\ell'} \equiv \sqrt{\mathcal{V}_{\ell'\ell}^2 + \mathcal{A}_{\ell'\ell}^2})$$

Hot Axions via Lepton Decays

Axions as Hot
Relics

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Axion via
Quarks

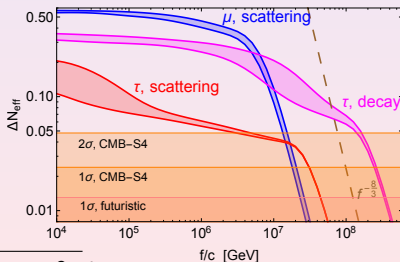
Axion via
Leptons

Axions via
Pions

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$$(c_{\ell\ell'} \equiv \sqrt{\mathcal{V}_{\ell'\ell}^2 + \mathcal{A}_{\ell'\ell}^2})$$

- More efficient** than scatterings (**larger f/c**)

Hot Axions via quark Decays

Axions as Hot
Relics

The QCD
Axion

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Gluons

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Axion via
Leptons

Axions via
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- a -quarks interaction can be also **flavor non-diagonal**

Hot Axions via quark Decays

Axions as Hot
Relics

The QCD
Axion

Axions via
Gluons

Axion via
Quarks

Axion via
Leptons

Axions via
Pions

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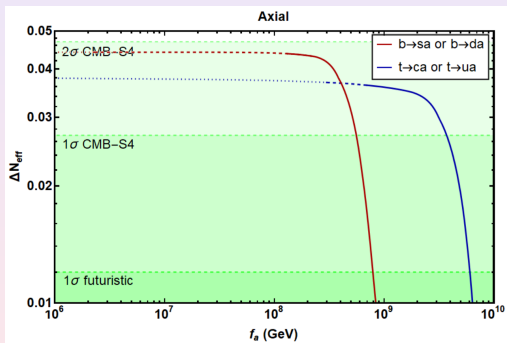


Figure: F.Arias-Aragon et al. JCAP 2021.

- **More efficient** than scatterings (**larger $f/c \lesssim 10^{10}$ GeV**)

Axion-Pion coupling

Axions as Hot
Relics

- **DFSZ model**: couples to u-type and d-type quarks,
- **KSVZ model**: no coupling to SM fermions

The QCD
Axion

Axions via
Gluons

Axion via
Quarks

Axion via
Leptons

Axions via
Pions

$$\mathbf{DFSZ} : \quad c_u^0 = \frac{1}{3} \cos^2(\beta), \quad c_d^0 = \frac{1}{3} \sin^2(\beta),$$

$$\mathbf{KSVZ} : \quad c_u^0 = c_d^0 = 0,$$

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The QCD
Axion

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Axions via
Gluons

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Axion via
Quarks

- Coupling to pions:

Axion via
Leptons

$$\mathcal{L}_{a\pi} = \frac{c_{a\pi}}{f_\pi} \frac{\partial_\mu a}{f} \left[2\partial^\mu \pi^0 \pi^+ \pi^- - \pi_0 (\partial^\mu \pi^+ \pi^- - \pi^+ \partial^\mu \pi^-) \right],$$

Axions via
Pions

where

$$c_{a\pi} = -\frac{1}{3} c_u^0 - c_d^0 - \frac{1-z}{1+z}. \quad z \equiv \frac{m_u}{m_d} \simeq 0.47_{-0.07}^{+0.06},$$

Axion-Pion coupling

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Relics

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$$\text{KSVZ : } c_{a\pi} \simeq 0.12_{-0.018}^{+0.023},$$

$$\text{DFSZ : } c_{a\pi} \simeq 0.12_{-0.018}^{+0.023} - \frac{1}{9} \cos(2\beta).$$

CMB Bounds on DFSZ

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Relics

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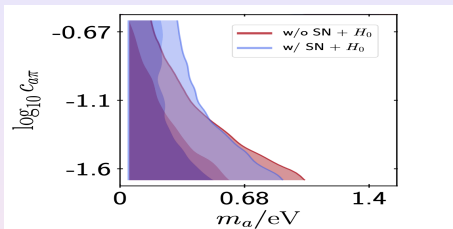


Figure: Constraints due to pion production Planck 18 + BAO (+ Pantheon + SH0ES H_0)

CMB Bounds on DFSZ

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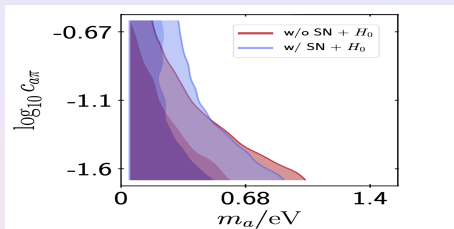


Figure: Constraints due to pion production Planck 18 + BAO (+ Pantheon + SH0ES H_0)

For DFSZ-II: muon production is also relevant for $c_{a\pi} \lesssim \mathcal{O}(0.1)$:

DFSZ-I	Planck 18+BAO (+SN+ H_0)
$c_{a\pi} = 0.225$	$m_a \leq 0.20$ (0.29) eV
$c_{a\pi} = 0.0225$	$m_a \leq 0.84$ (0.82) eV
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CMB Bounds on DFSZ

Axions as Hot Relics

The QCD Axion

Axions via Gluons

Axion via Quarks

Axion via Leptons

Axions via Pions

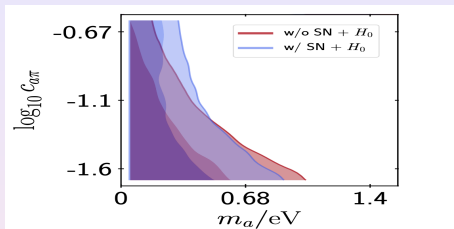


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Caveat! Pion cross-section calculation should break down at

$T \gtrsim 60$ MeV (Di Luzio et al. 2021, arXiv 2101.10330.)

Conclusions

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- 1 If $f \lesssim \mathcal{O}(10^9)$ GeV, coupling with **quarks and leptons** (with $c_i = \mathcal{O}(1)$) dominates over $\frac{\alpha_s}{8\pi} \frac{a}{f} G\tilde{G}$
- 2 Efficiency peaks at $T \approx m_f$

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(*maybe higher for c -quark?)

Conclusions

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Conclusions

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- 8 **Future CMB experiments** will tell in a few years, plus **direct detection** (*e.g. IAXO*)