

Cosmological Constant UNIVERSE

Higgs Mass Squared

 $\frac{2}{mh}|H|^2$

WEAK FORCE, STRUCTURE OF NUCLEI, COMPLEX CHEMISTRY, ...

SYMMETRY

$$m_h^2 \sim y_t^2 M_{\rm Pl}^2$$

SYMMETRY

$$m_h^2 \sim y_t^2 M_{\rm Pl}^2$$

Selection Rules of Spacetime Dilations

(assuming masses at the Planck scale)

SYMMETRY

$$m_h^2 \sim y_t^2 M_{\rm Pl}^2$$

Selection Rules of the Higher-Spin Symmetry of Free Scalars

See R. Rattazzi @ GGI: https://www.ggi.infn.it/talkfiles/slides/slides5297.pdf

Cosmological Constant Theory ~ 10¹²⁰ Experiment

Higgs Mass Squared

 $\frac{2}{mh}|H|^2$

Cosmological Constant

Higgs Mass Squared

This is not a well-posed problem in the Standard Model EFT

Theory ~ 10¹²⁰ Experiment

Cosmological Constant

Higgs Mass Squared

The two quantities are not calculable

Theory ~ 10¹²⁰ Experiment

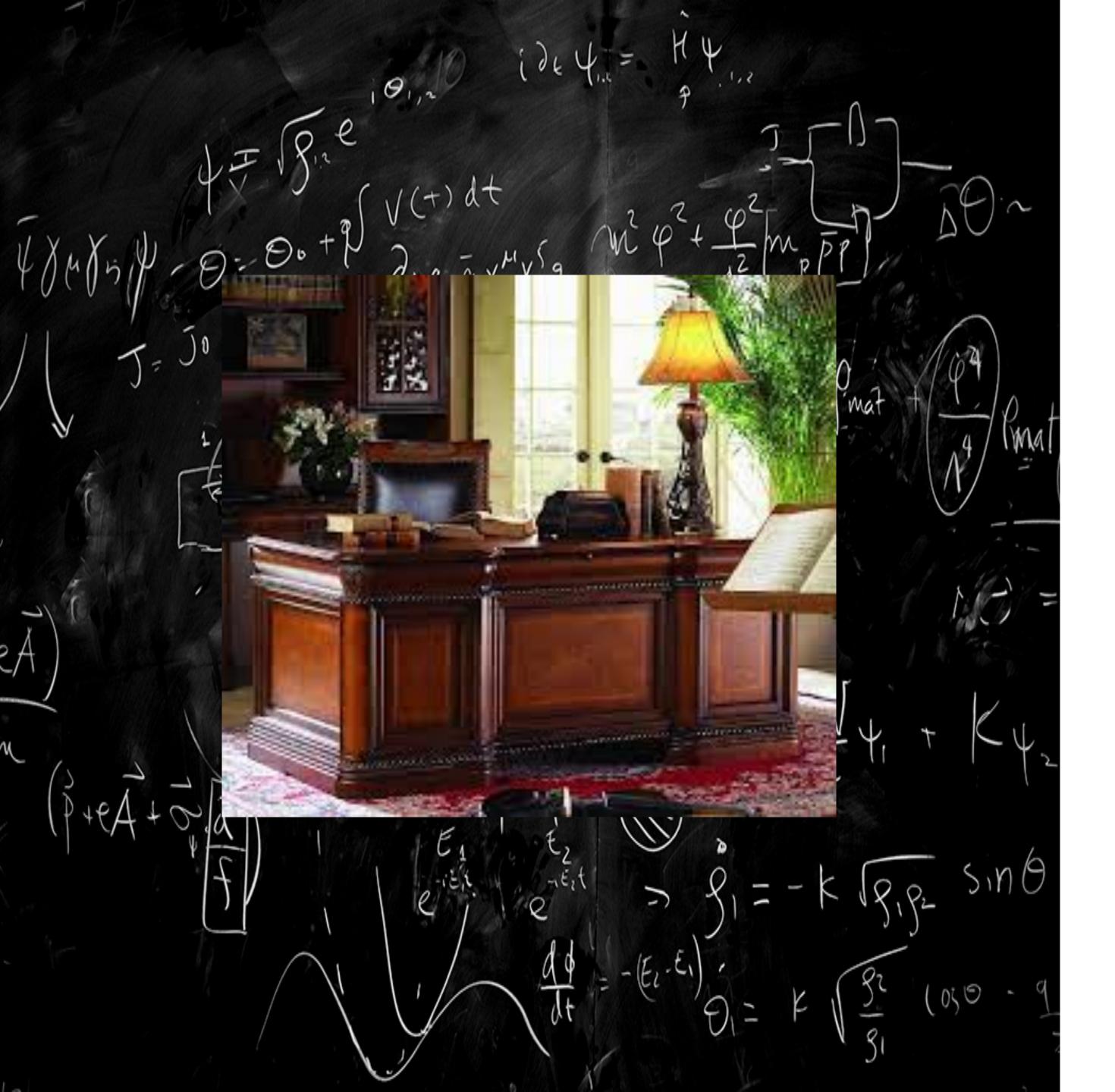
Cosmological Constant

Higgs Mass Squared

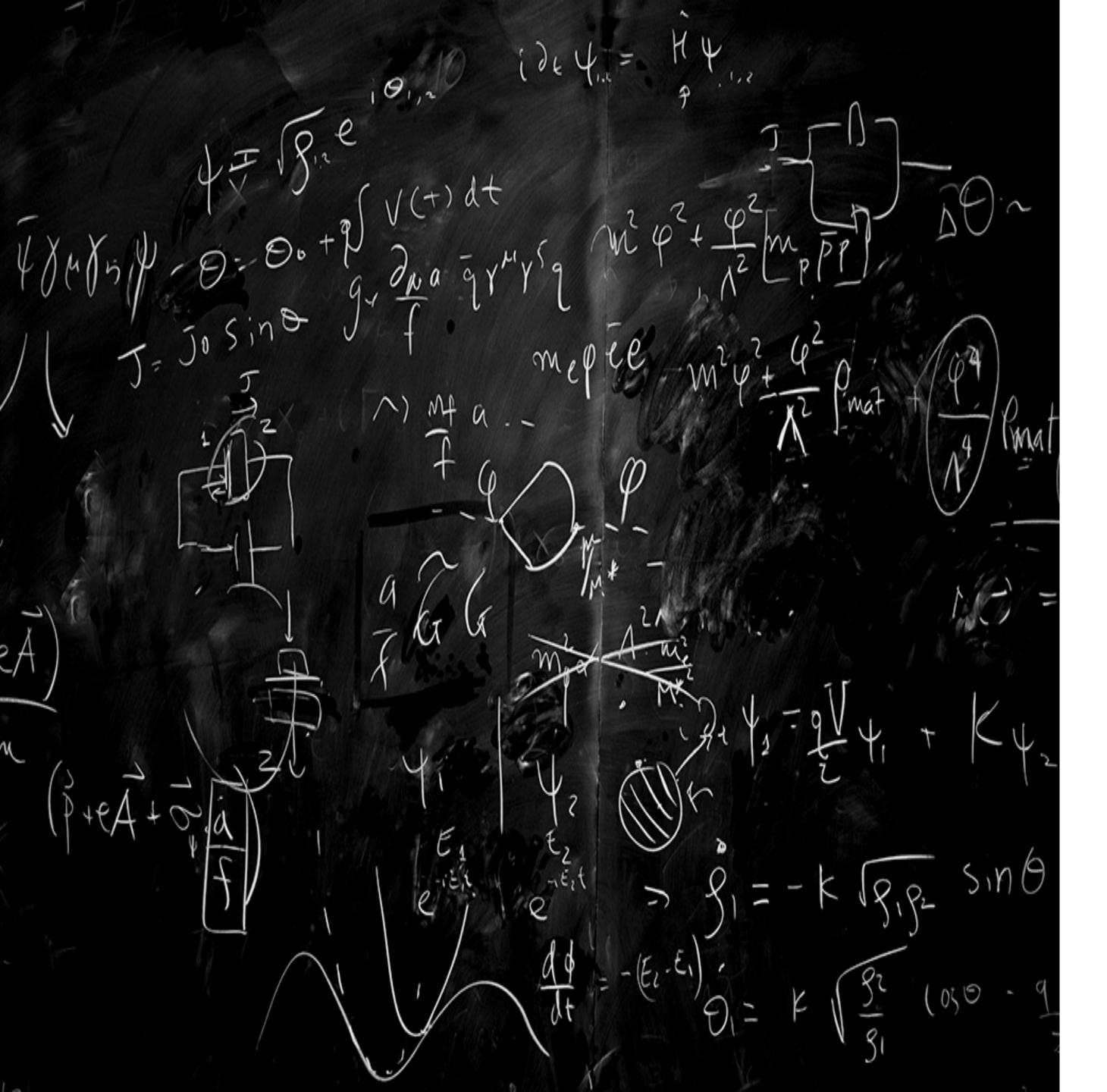
Precious clue about fundamental physics in the UV

Theory ~ 10¹²⁰ Experiment

WHAT KIND OF PHYSICIST ARE YOU?



The Theorist



The Theorist

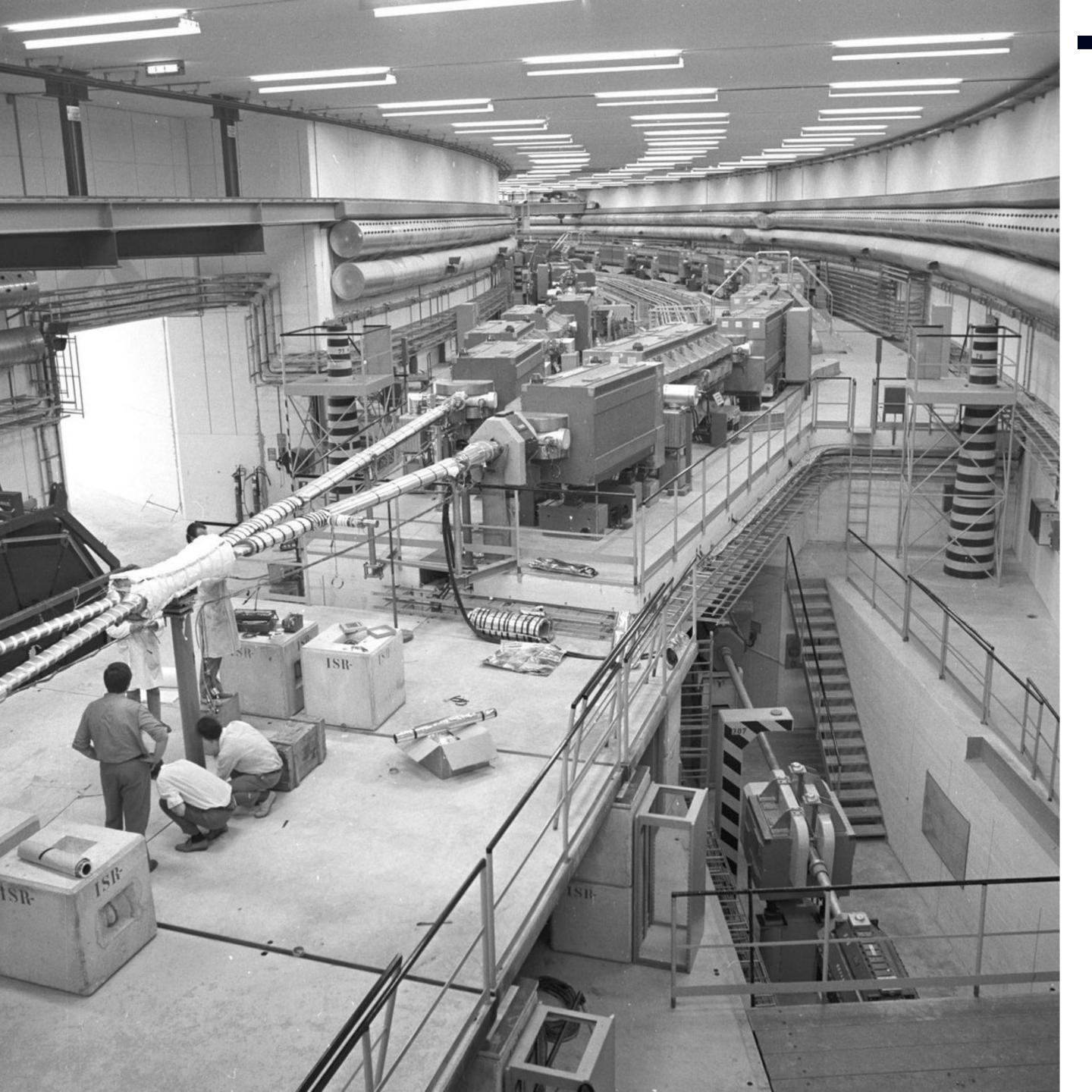
Planck
String

SUSY

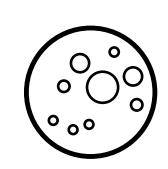
SIV



The Experimentalist

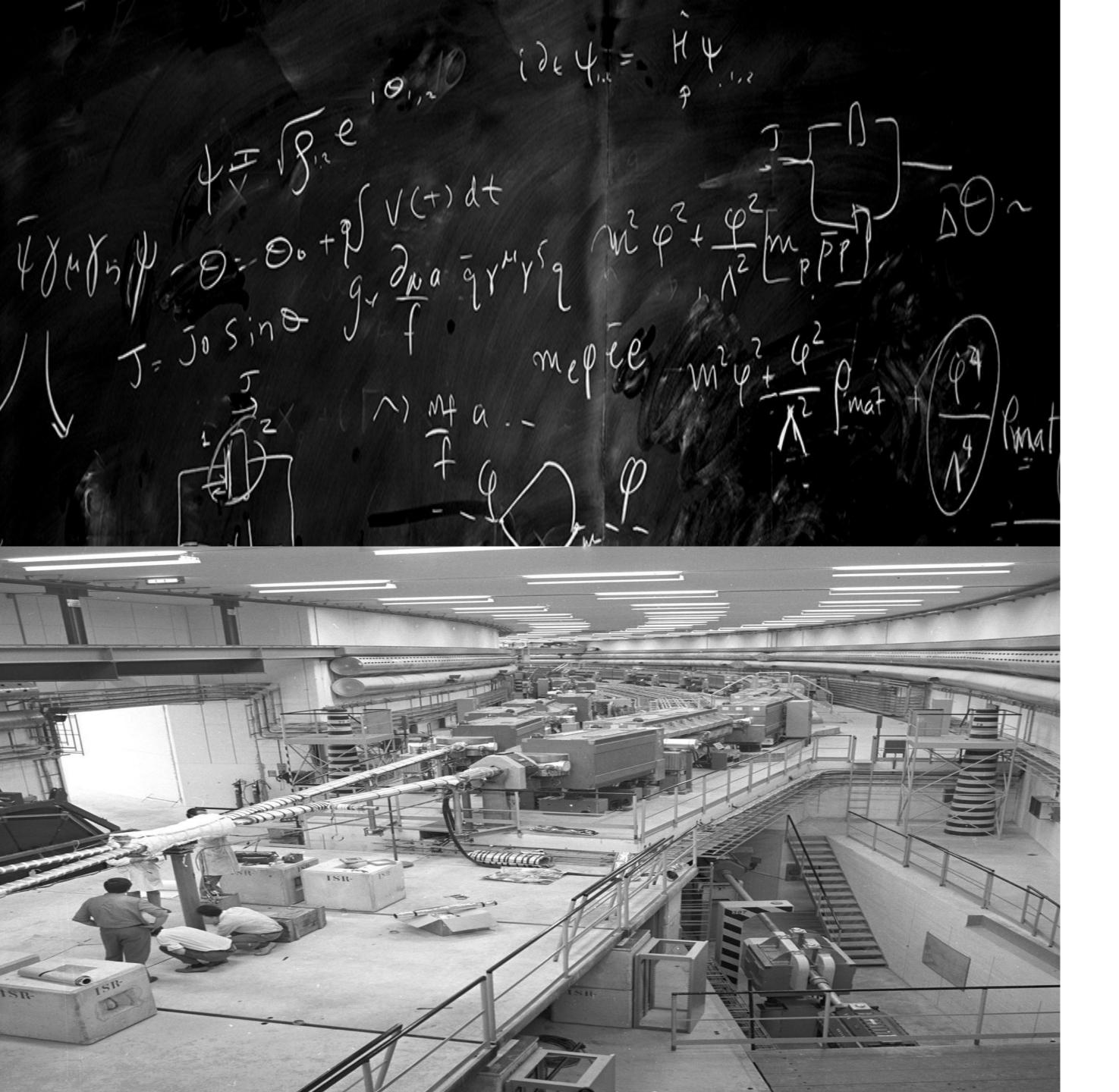


The Experimentalist

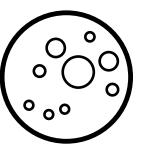


Mysterious QG Blob

SM



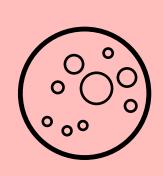
The Theorist who wants to be an Experimentalist



Mysterious QG Blob

SUSY or confining sector

SM



Mysterious QG RIOh

SM



 Cosmological Constant and Higgs mass are inputs (can not be calculated, only measured)

2. UV/IR Mixing

IR constraints from UV consistency (swampland, ...)

SM

Planck
String
SUSY

Mysterious
QG
Blob
SUSY
SUSY

SM SM

Mysterious Planck QG Blob String SUSY SUSY

Well-posed question in the low energy theory

Caveat: it could sill be answered in the UV (landscape+swampland, ...)

EXAMPLE: HIGGS VEV IN THE MSSM

$$v^{2} = \frac{2}{g^{2} + g'^{2}} \left(\frac{|m_{H_{d}}^{2} - m_{H_{u}}^{2}|}{\sqrt{1 - \sin(2\beta)^{2}}} - m_{H_{u}}^{2} - m_{H_{d}}^{2} - 2|\mu|^{2} + \frac{y_{t}^{2}}{16\pi^{2}} m_{\widetilde{t}}^{2} f(\frac{m_{\widetilde{t}_{1}}^{2}}{m_{\widetilde{t}}^{2}}, \frac{m_{\widetilde{t}_{2}}^{2}}{m_{\widetilde{t}}^{2}}) \right) + \dots$$

FINE-TUNING 101

A physical observable can be computed as the sum of multiple unrelated contributions

$$\mathcal{O} = O_1 + O_2 + \dots$$

At least two of them are much larger than its observed value

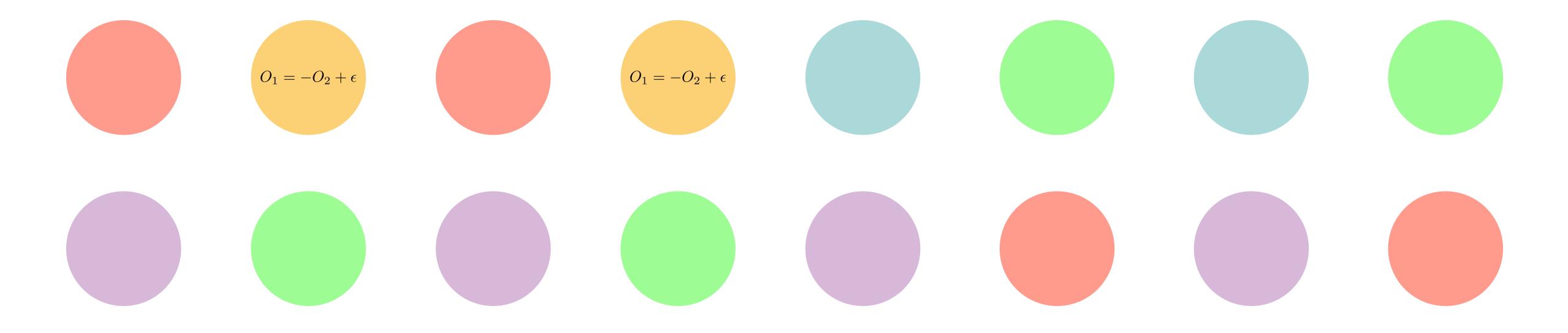
$$\mathcal{O}_{\mathrm{obs}} \ll |O_{1,2}|$$

FINE-TUNING 101

Is there a symmetry without a landscape?

$$O_1 = -O_2 + \epsilon$$

Is there a landscape?



FINE-TUNING 101

Is there a symmetry without a landscape?

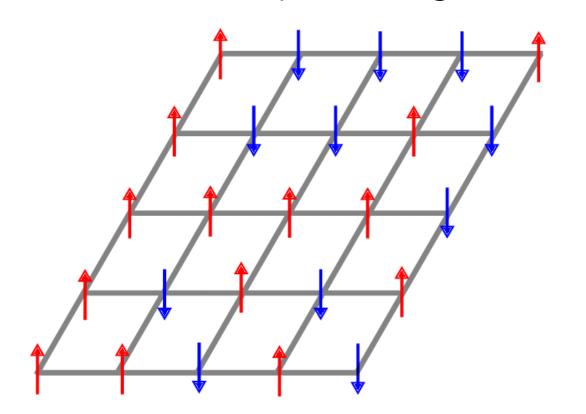
$$O_1 = -O_2 + \epsilon$$

Is there a landscape?

Example:

Prepare Ising Model

Scan Temperature



$$T - T_c \simeq 10^{-30}$$

The scalar is much lighter than the lattice spacing

STATUS OF THE FIELD

Is there a symmetry without a landscape?

$$O_1 = -O_2 + \epsilon$$



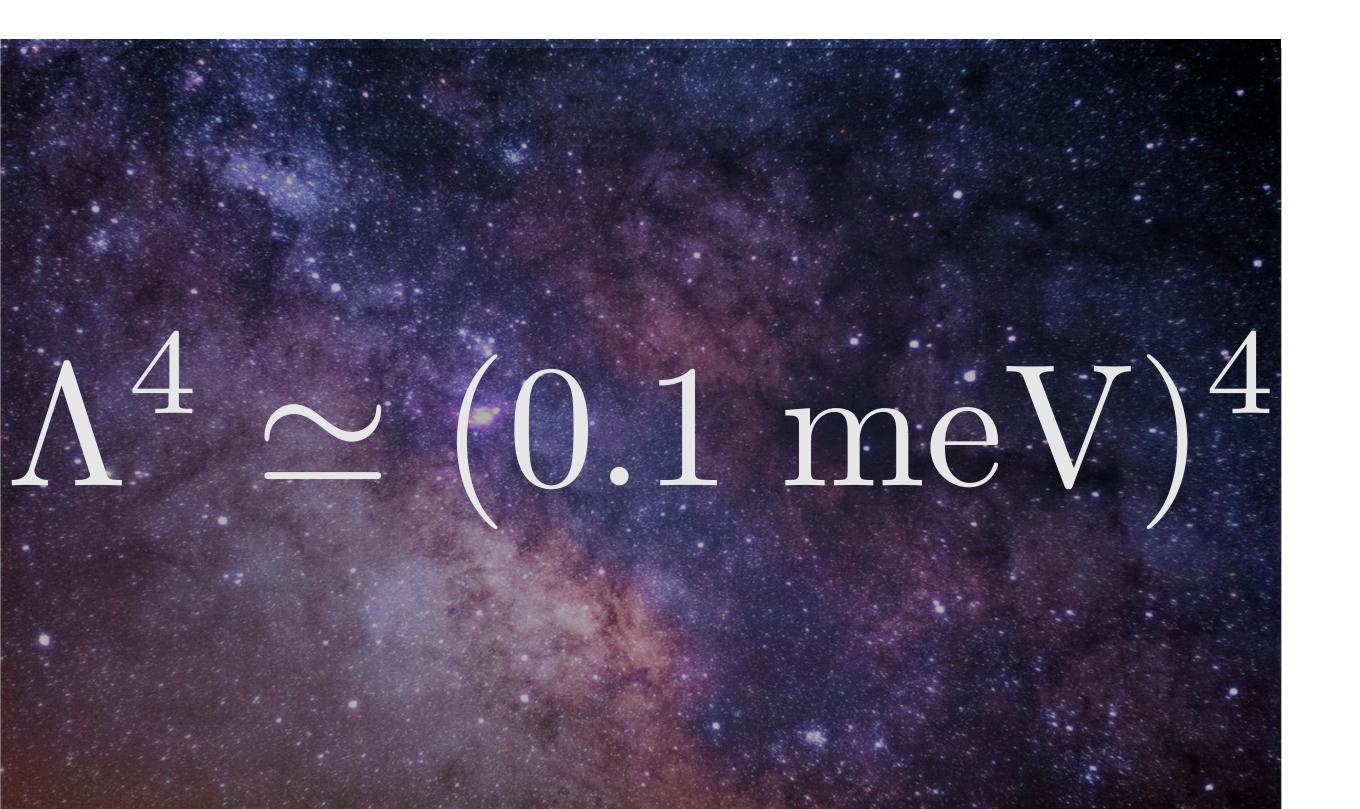
Traditional Approach: Factorize the problems



STATUS OF THE FIELD

Is there a symmetry without a landscape?

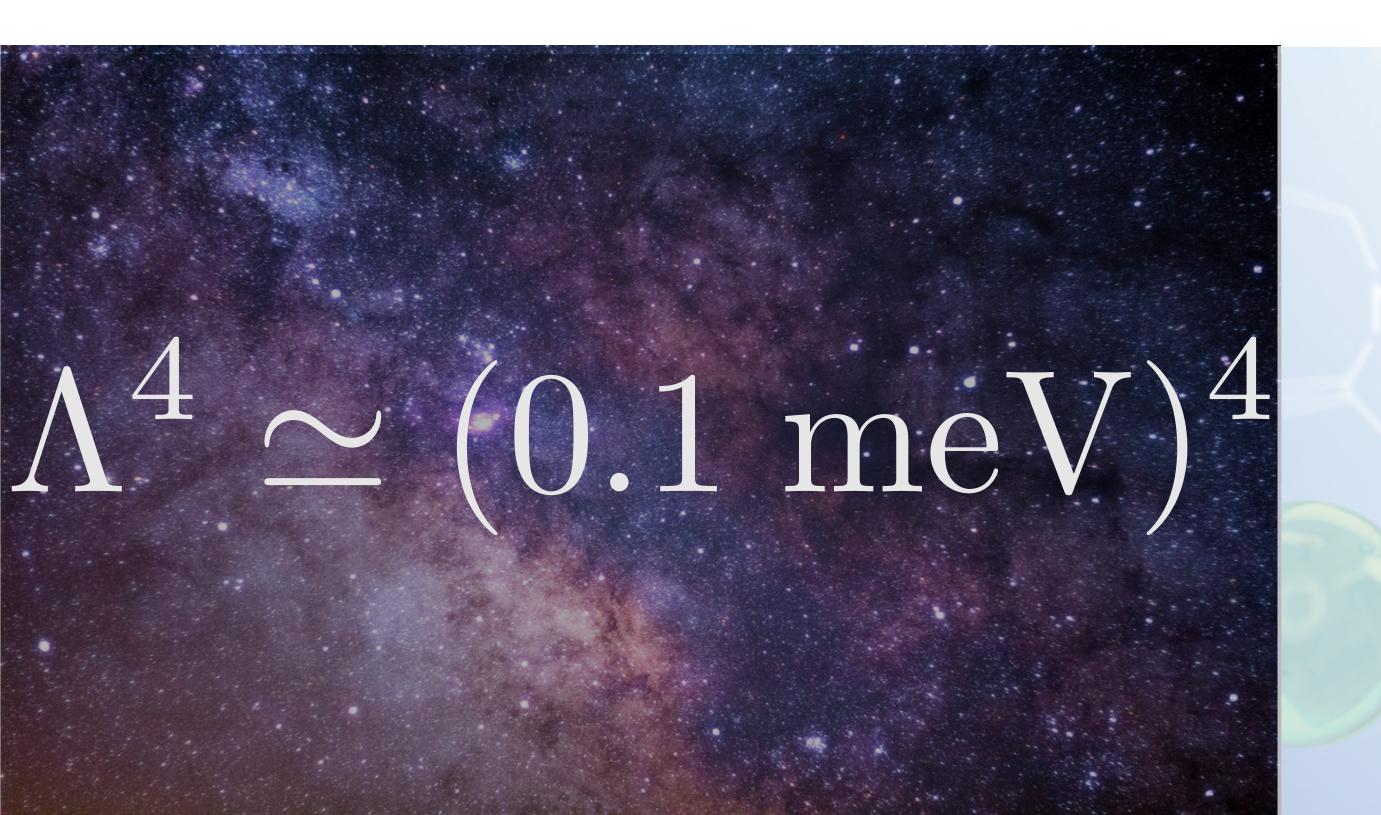
$$O_1 = -O_2 + \epsilon$$

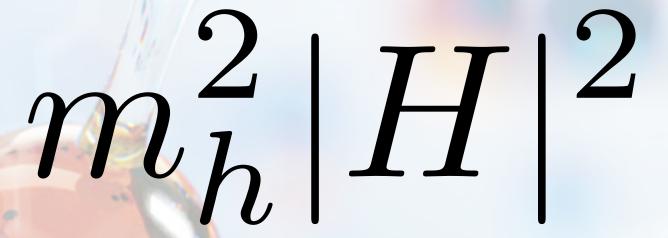


STATUS OF THE FIELD

Is there a symmetry without a landscape?

$$O_1 = -O_2 + \epsilon$$





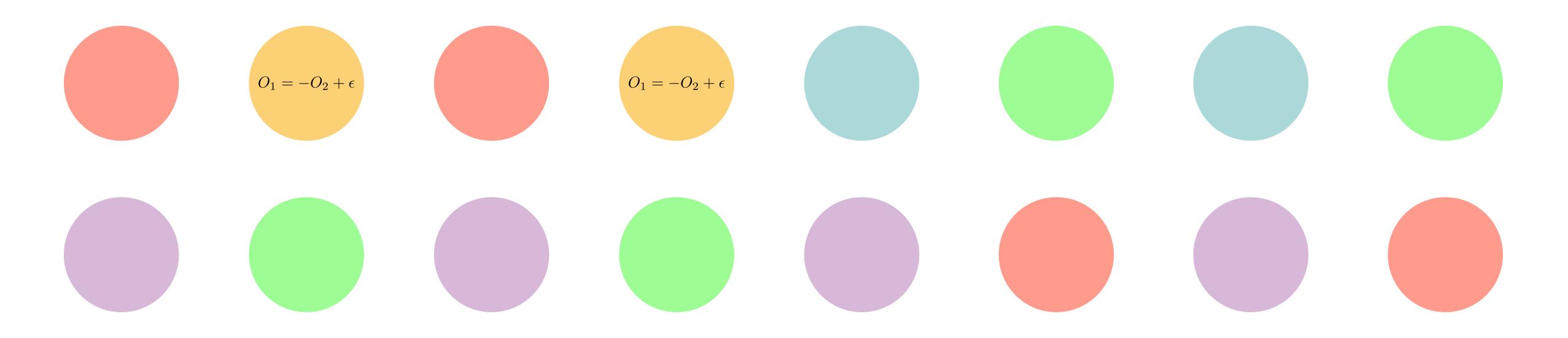
We have been looking for these simple and elegant solutions for more than 40 years

We have been looking for these simple and elegant solutions for more than 40 years

It is a good time to consider seriously more creative alternatives

A LANDSCAPE

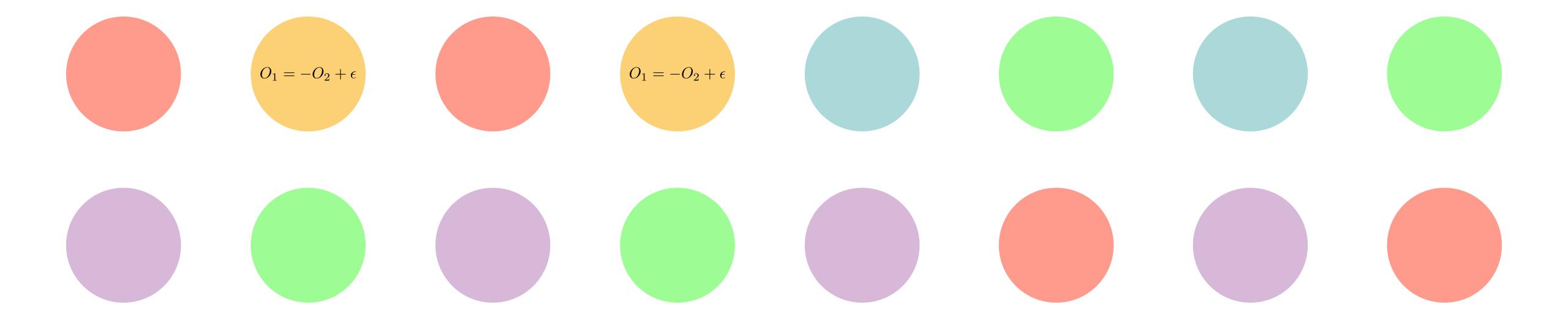
Is there a landscape?



A LANDSCAPE

- 1. One day it can be tested experimentally
- 2. Currently our most concrete explanation for the CC
- 3. It probably exists independently of the two problems

Is there a landscape?



EXAMPLE: HIGGS VEV IN THE MSSM

$$v^{2} = \frac{2}{g^{2} + g'^{2}} \left(\frac{|m_{H_{d}}^{2} - m_{H_{u}}^{2}|}{\sqrt{1 - \sin(2\beta)^{2}}} - m_{H_{u}}^{2} - m_{H_{d}}^{2} - 2|\mu|^{2} + \frac{y_{t}^{2}}{16\pi^{2}} m_{\widetilde{t}}^{2} f(\frac{m_{\widetilde{t}_{1}}^{2}}{m_{\widetilde{t}}^{2}}, \frac{m_{\widetilde{t}_{2}}^{2}}{m_{\widetilde{t}}^{2}}) \right) + \dots$$

$$m_{H_u}^2 \gg m_h^2$$

Is there a problem? You first need to know the distribution of these parameters in the landscape

Change of perspective:



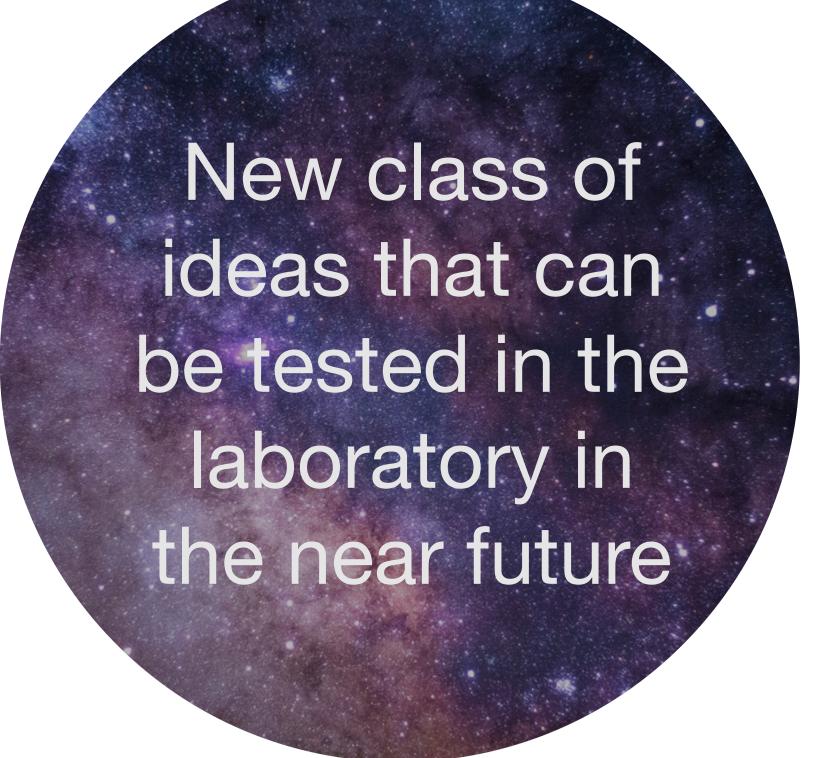


Can we find the origin of the weak scale early in the history of the Universe?

Historically:



Recently:



Historically:



Recently:



EXAMPLE: ANTHROPICS

Anthropic Arguments: a symmetry and a landscape

$$\Lambda_{
m CC} \leftrightarrow
ho_{
m MR} \left(rac{\delta
ho}{
ho}
ight)^3 \ll M_{
m Pl}^4$$
 [Weinberg '87]

$$v \leftrightarrow \Lambda_{\mathrm{QCD}} \ll M_{\mathrm{Pl}}$$
 [Agrawal, Barr, Donoghue, Seckel '97]

EXAMPLE: ANTHROPICS

Anthropic Arguments: a symmetry and a landscape

$$\Lambda_{\rm CC} \leftrightarrow \rho_{
m MR} \left(rac{\delta
ho}{
ho}
ight)^3 \ll M_{
m Pl}^4$$

We do not know yet, but easy to achieve

$$v \leftrightarrow \Lambda_{\rm QCD} \ll M_{\rm Pl}$$

Approximate scale invariance

Historically:



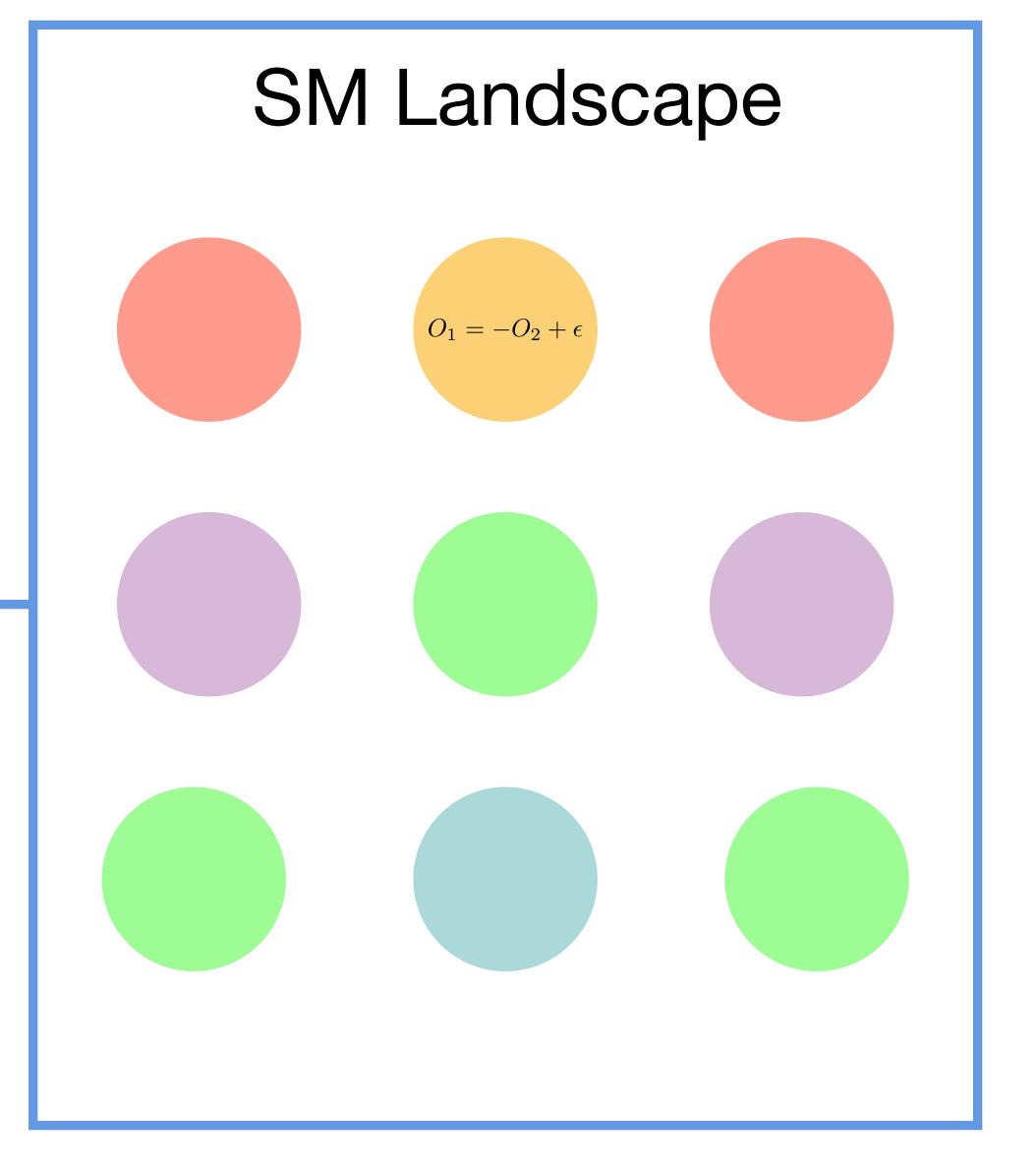
Recently:

New class of ideas that can be tested in the laboratory in the near future

Early History of the Universe

Symmetric Sector

$$\Lambda_S \ll M_{\rm Pl}$$

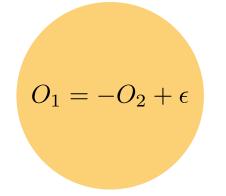


Late times

Symmetric Sector

$$\Lambda_S \ll M_{\rm Pl}$$

SM Landscape



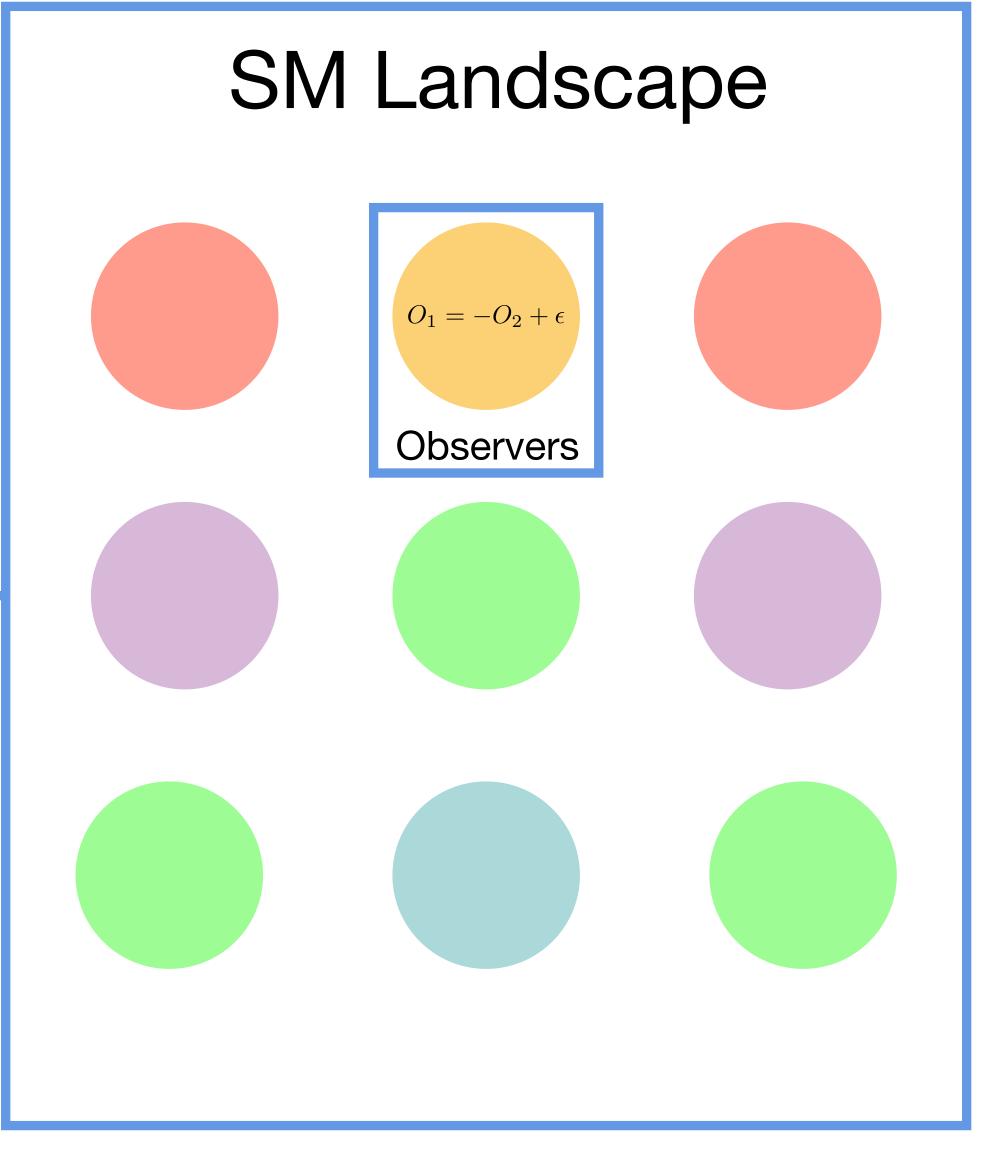
An event triggered by the symmetric sector selects the observed

$$m_h^2$$

"Anthropic" Selection

Symmetric Sector

 $\Lambda_S \ll M_{\rm Pl}$

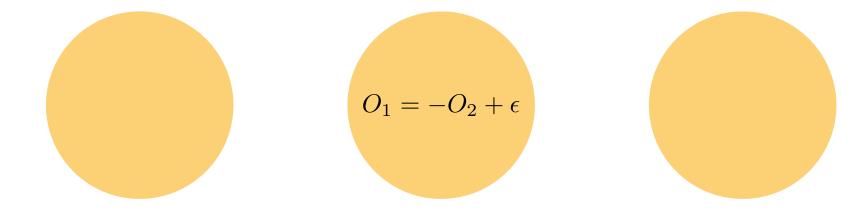


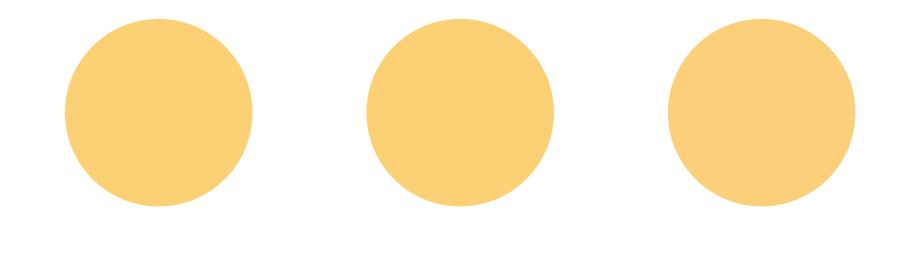
"Statistical" Selection

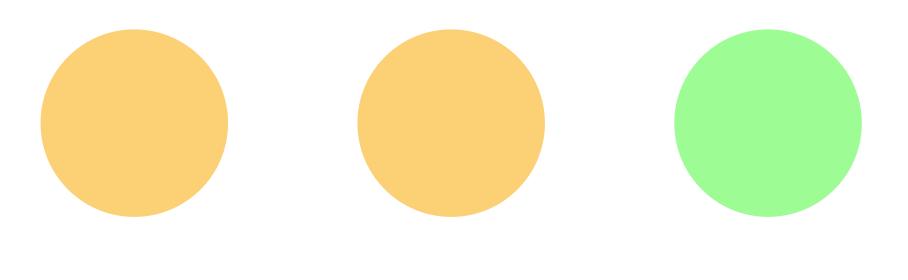
Symmetric Sector

$$\Lambda_S \ll M_{\rm Pl}$$









"Dynamical" Selection

SM Landscape

$$O_1 = -O_2 + \epsilon$$

Symmetric Sector

$$\Lambda_S \ll M_{\rm Pl}$$

From now on: focus on the Higgs vev

From now on: focus on the Higgs vev

Anthropic Selection

Statistical Selection

Dynamical Selection

[Agrawal, Barr, Donoghue, Seckel '97],[Arvanitaki, Dimopoulos, Gorbenko, Huang, Van Tilburg '16],[Arkani-Hamed, RTD, Kim, '20],[Giudice, Kehagias, Riotto, '20],

[Dvali, Vilenkin '03], [Dvali '04], [Geller, Hochberg, Kuflik, '18], [Giudice, McCullough, You, '21],

. . .

[Graham, Rajendran, Kaplan, '15], [Arkani-Hamed, Cohen, RTD, Kim, Pinner, '16], [Csaki, RTD, Geller, Ismail, '20], [Strumia, Teresi, '20], [RTD, Teresi, '21],

. . .

EXAMPLE: ANTHROPIC

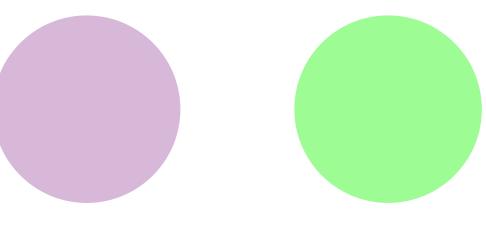
[Agrawal, Barr, Donoghue, Seckel '97]

For complex chemistry we need a Higgs vev not too far from the QCD scale

Symmetric Sector

$$\Lambda_{
m QCD} \ll M_{
m Pl}$$

SM Landscape



QCD

 Y_qQHq^c

"Friendly"

String Landscape?

[Dvali, Vilenkin '03], [Dvali '04]

$$F_4 = dA_3$$

[Dvali, Vilenkin '03], [Dvali '04]

$$F_4 = dA_3$$

$$S \supset \int d^4x \sqrt{-g} \left(\frac{F_4^2}{48} + M_{\rm Pl}^2 (-1 + \frac{F_4^2}{M_{\rm Pl}^2} + \ldots) |\phi|^2 + \ldots \right) + q(\phi) \int d^3\xi A_{\mu\nu\rho} \frac{\partial x^\mu}{\partial \xi^a} \frac{\partial x^\nu}{\partial \xi^b} \frac{\partial x^\rho}{\partial \xi^c} \epsilon^{abc}$$

[Dvali, Vilenkin '03], [Dvali '04]

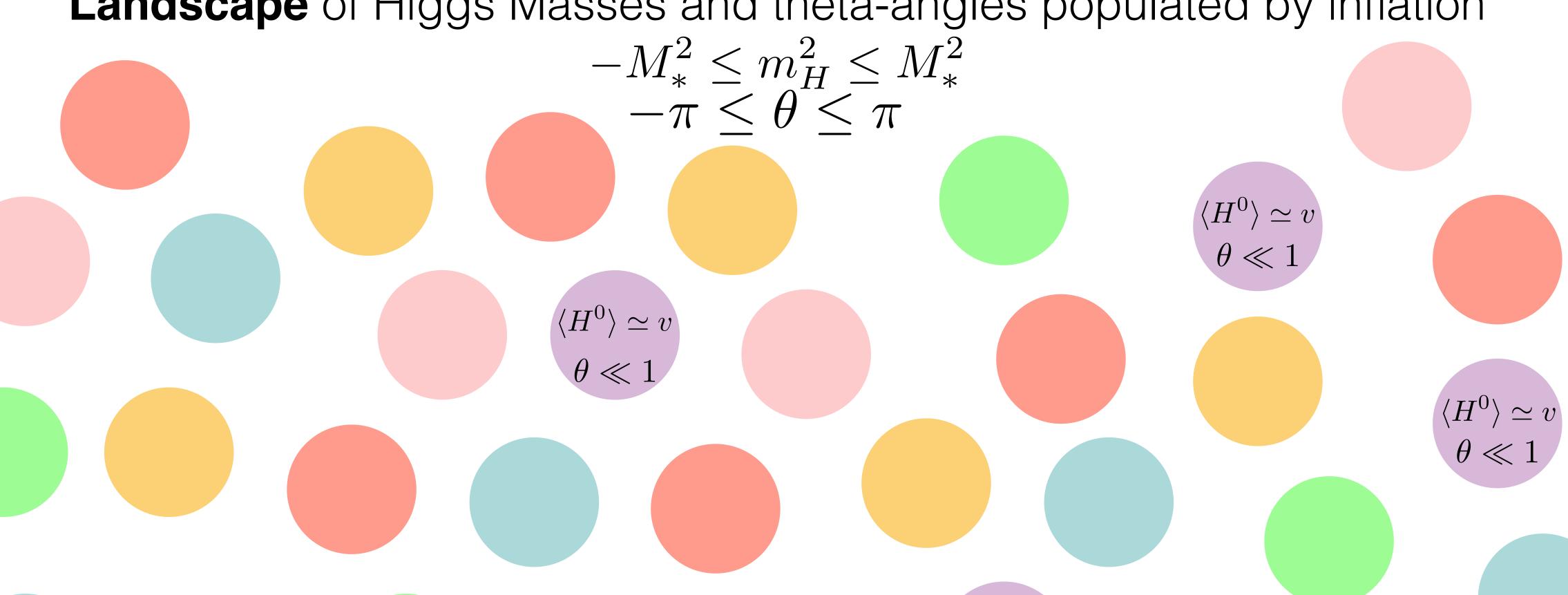
$$q(\phi) = \frac{\phi^N}{M_{\rm Pl}^{N-2}}$$

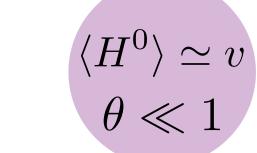
$$\Delta \langle \phi \rangle^2 / \langle \phi \rangle^2 \sim \langle \phi \rangle^{N-2}$$

At every step the brane charge is smaller -> most vacua are at small vev

[Bloch, Csaki, Geller, Volansky '19], [Csaki, Geller, RTD, Ismail, '20], [RTD, Teresi '21]

Landscape of Higgs Masses and theta-angles populated by inflation







 $\theta \ll 1$

[Bloch, Csaki, Geller, Volansky '19], [Csaki, Geller, RTD, Ismail, '20], [RTD, Teresi '21]

After reheating and a time

$$t_c \sim 1/H(\Lambda_{\rm QCD}) \sim 10^{-5} s$$

All patches where the Higgs vev

$$\langle H^0 \rangle \simeq v$$
 $\theta \ll 1$

$$\langle H^0 \rangle \equiv h$$

Is outside of a certain range

$$h_{\min} \lesssim h \leq h_{\text{crit}}$$

And theta is large

$$\theta \le \theta_{\rm max}$$

$$\langle H^0 \rangle \simeq$$
 $\theta \ll 1$

crunch



$$\langle H^0 \rangle \simeq v$$
$$\theta \ll 1$$

$$\langle H^0 \rangle \simeq v$$
 $\theta \ll 1$

[Bloch, Csaki, Geller, Volansky '19], [Csaki, Geller, RTD, Ismail, '20], [RTD, Teresi '21] Only universes with the observed value of the weak scale can live longer than EW time. **Today the multiverse looks like**:

$$\langle H^0 \rangle \simeq v$$
 $\theta \ll 1$

$$\langle H^0 \rangle \simeq v$$
 $\theta \ll 1$

$$\langle H^0 \rangle \simeq v$$

$$\theta \ll 1$$

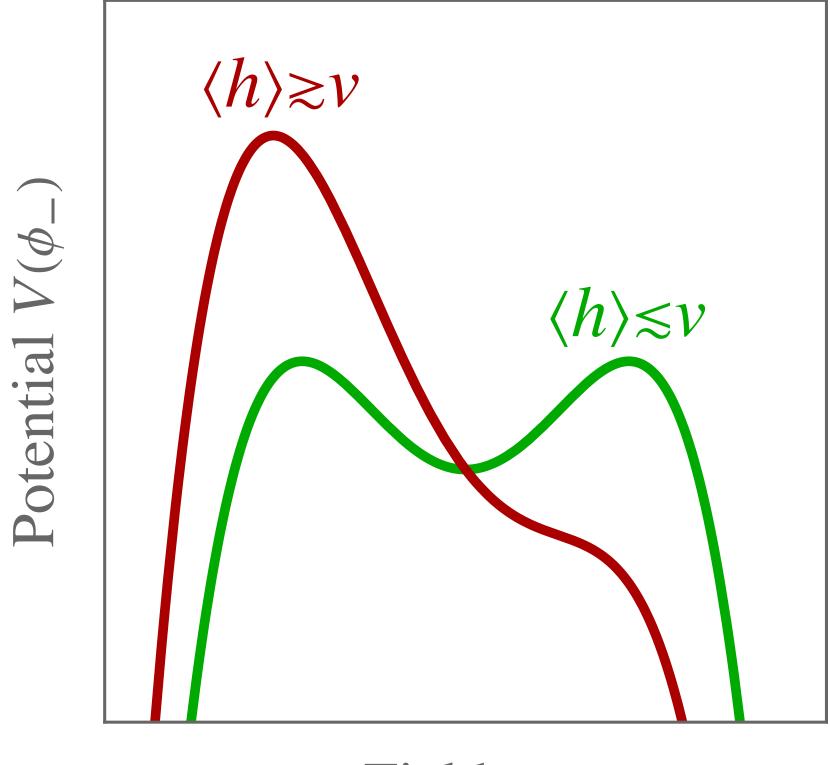
$$\langle H^0 \rangle \simeq i$$
 $\theta \ll 1$

[RTD, Teresi '21]

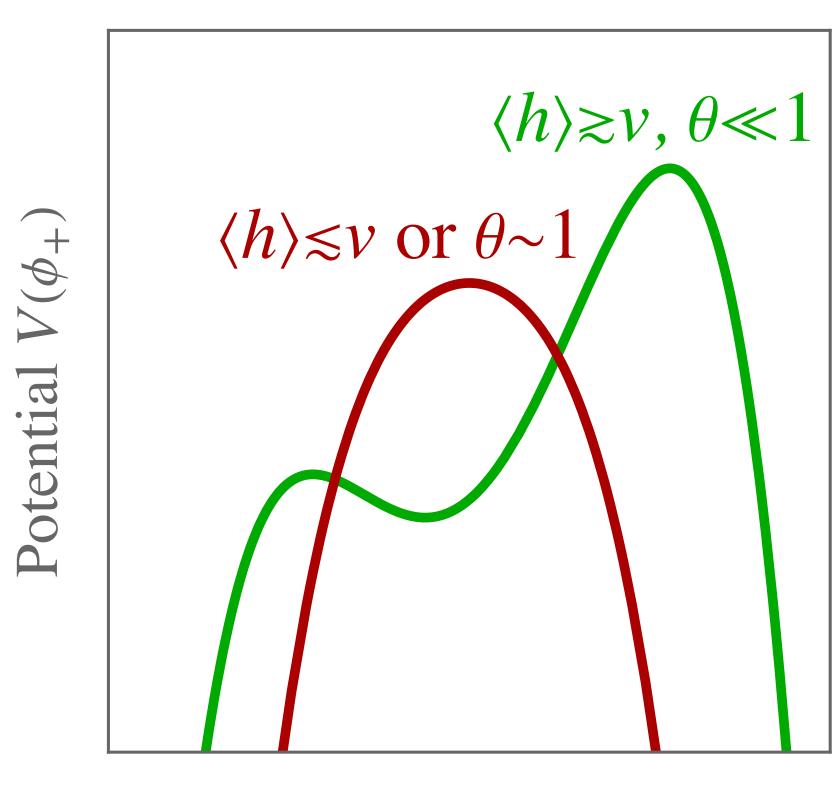
$$V = \mp m_{\pm}^2 \phi_{\pm}^2 - \lambda_{\pm} \phi_{\pm}^4 + \frac{\alpha_s}{8\pi} \left(\theta + \frac{\phi_+}{F_+} + \frac{\phi_-}{F_-} \right) G\widetilde{G}$$

[RTD, Teresi '21]

$$V = \mp m_{\pm}^2 \phi_{\pm}^2 - \lambda_{\pm} \phi_{\pm}^4 + \frac{\alpha_s}{8\pi} \left(\theta + \frac{\phi_+}{F_+} + \frac{\phi_-}{F_-} \right) G \widetilde{G}$$



Field ϕ_{-}



Field ϕ_+

[RTD, Teresi '21]

$$V = \mp m_{\pm}^2 \phi_{\pm}^2 - \lambda_{\pm} \phi_{\pm}^4 + \frac{\alpha_s}{8\pi} \left(\theta + \frac{\phi_+}{F_+} + \frac{\phi_-}{F_-} \right) G \widetilde{G}$$

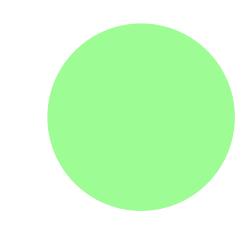
Symmetric Sector

$$m_{\pm} \ll M_{
m Pl}$$

ϕ_{\pm}

$$\phi_{\pm}G\widetilde{G}$$

SM Landscape





Anthropic Selection

Does not require new physics with couplings to the SM stronger than gravitational

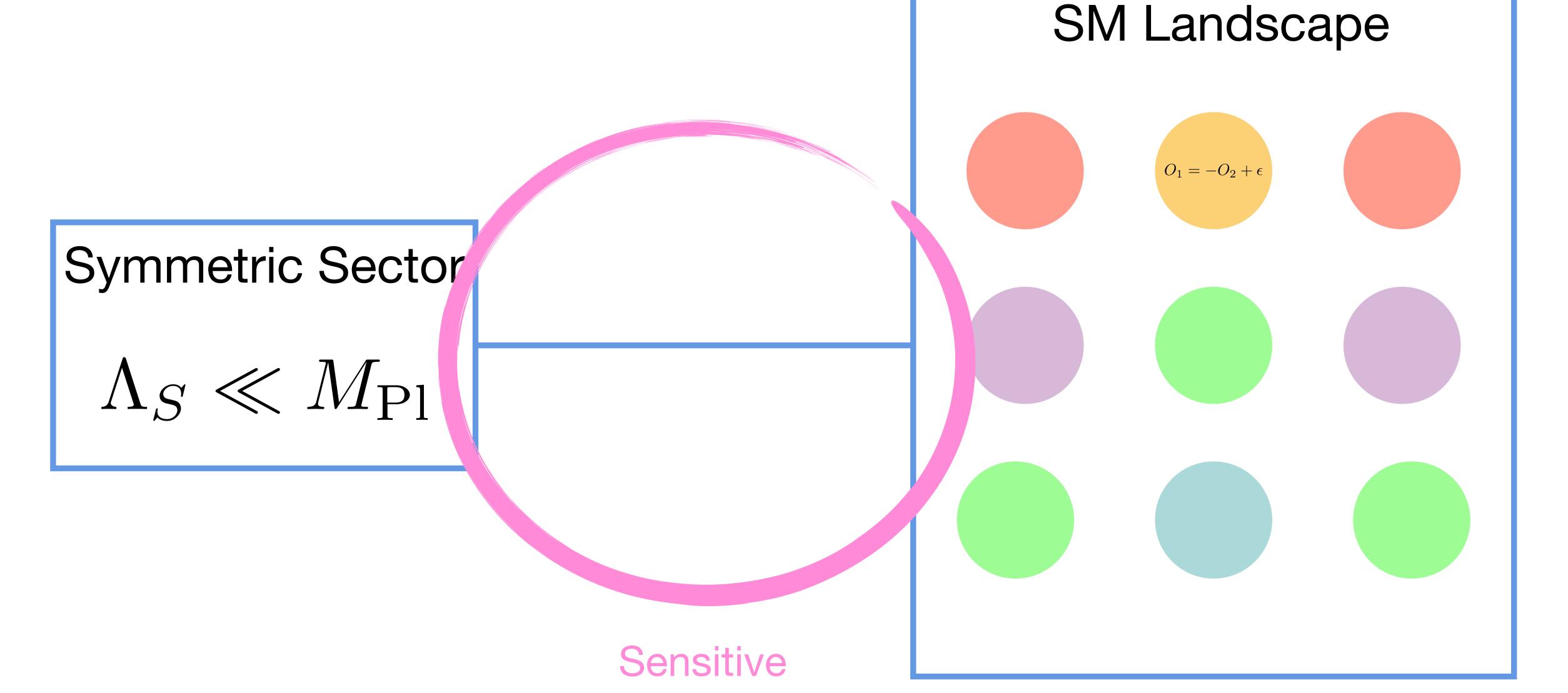
Statistical Selection

Does not require new physics with couplings to the SM stronger than gravitational

Dynamical Selection

Typically visible!

"Dynamical" Selection



to the Higgs vev

HIERARCHY 102

[Arkani-Hamed, RTD, Kim '20]

Does anything change in the SM as we vary $\langle h \rangle$?

1. Obviously the spectrum

2. If we look at local operators we discover the hierarchy problem:

$$\langle h^{\dagger} h \rangle \sim \Lambda_H^2$$

HIERARCHY 102

Does anything change in the SM as we vary $\langle h \rangle$?

$$\frac{\alpha_s}{8\pi} (\xi \phi + \theta) \operatorname{Tr} \left[G \tilde{G} \right]$$

$$m_{\pi}^2 f_{\pi}^2 \sqrt{1 - \frac{4m_u m_d}{(m_u + m_d)^2} \sin^2 (\xi \phi + \theta)}$$

$$(y_u + y_d) v f_{\pi}^3 \left(\theta \xi \phi + \xi^2 \phi^2 + \ldots \right)$$

HIERARCHY 102

Does anything change in the SM as we vary $\langle h \rangle$?

$$\xi \phi \text{Tr} \left[\widetilde{G} \widetilde{G} \right]$$



Important Pheno Message:

Axion-Like phenomenology can be related to the hierarchy problem

TRIGGER PHENOMENOLOGY

 \widetilde{GG}

ALPs

TRIGGER PHENOMENOLOGY

$$\widetilde{GG}$$

$$F\widetilde{F} + yLHE^c$$

Vector-like Leptons

TRIGGER PHENOMENOLOGY

 \widetilde{GG}

ALPs

$$F\widetilde{F} + yLHE^c$$

Vector-like Leptons

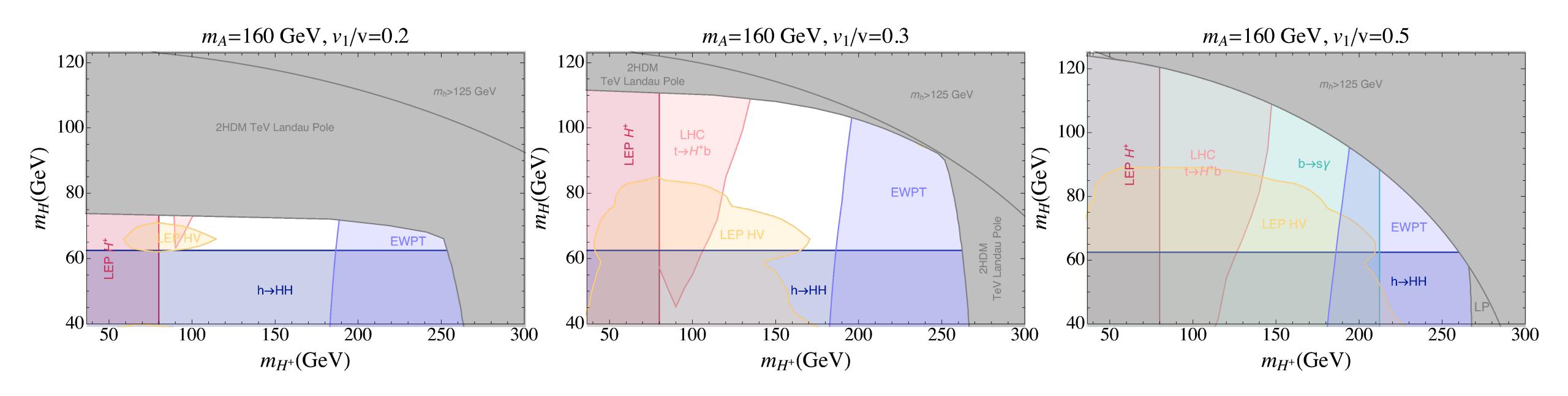
 H_1H_2

Type-0 2HDM
See [Arkani-Hamed, RTD, Kim, '20]

EXAMPLE: TYPE-0 2HDM

[Arkani-Hamed, RTD, Kim '20]

N.B. It is extremely hard to find a viable BSM trigger



Sharp target which can't be decoupled!

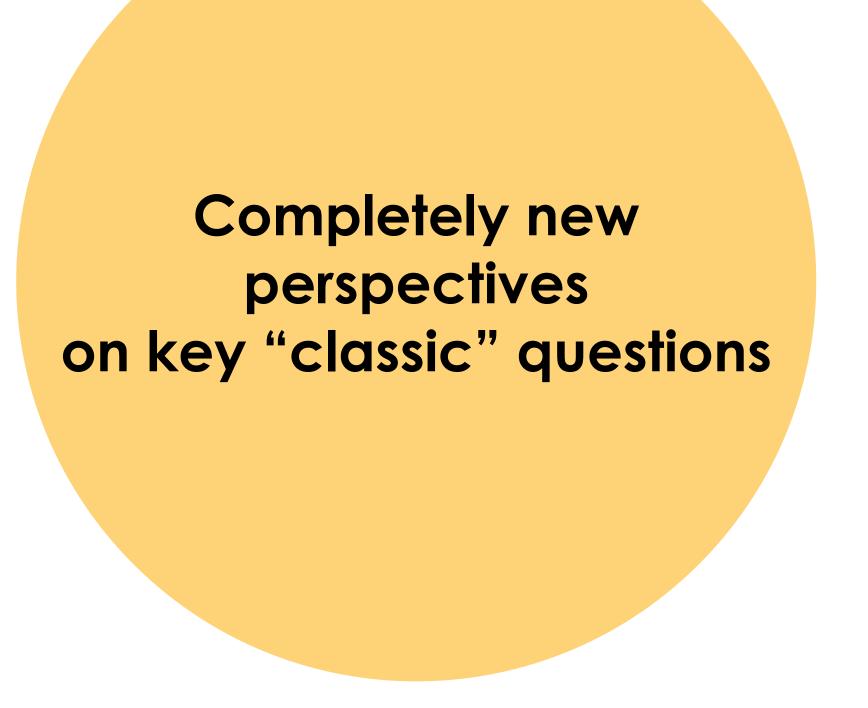
$$m_{\mathrm{NP}} \lesssim m_h$$

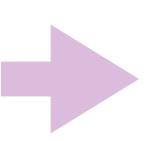
Change of perspective:





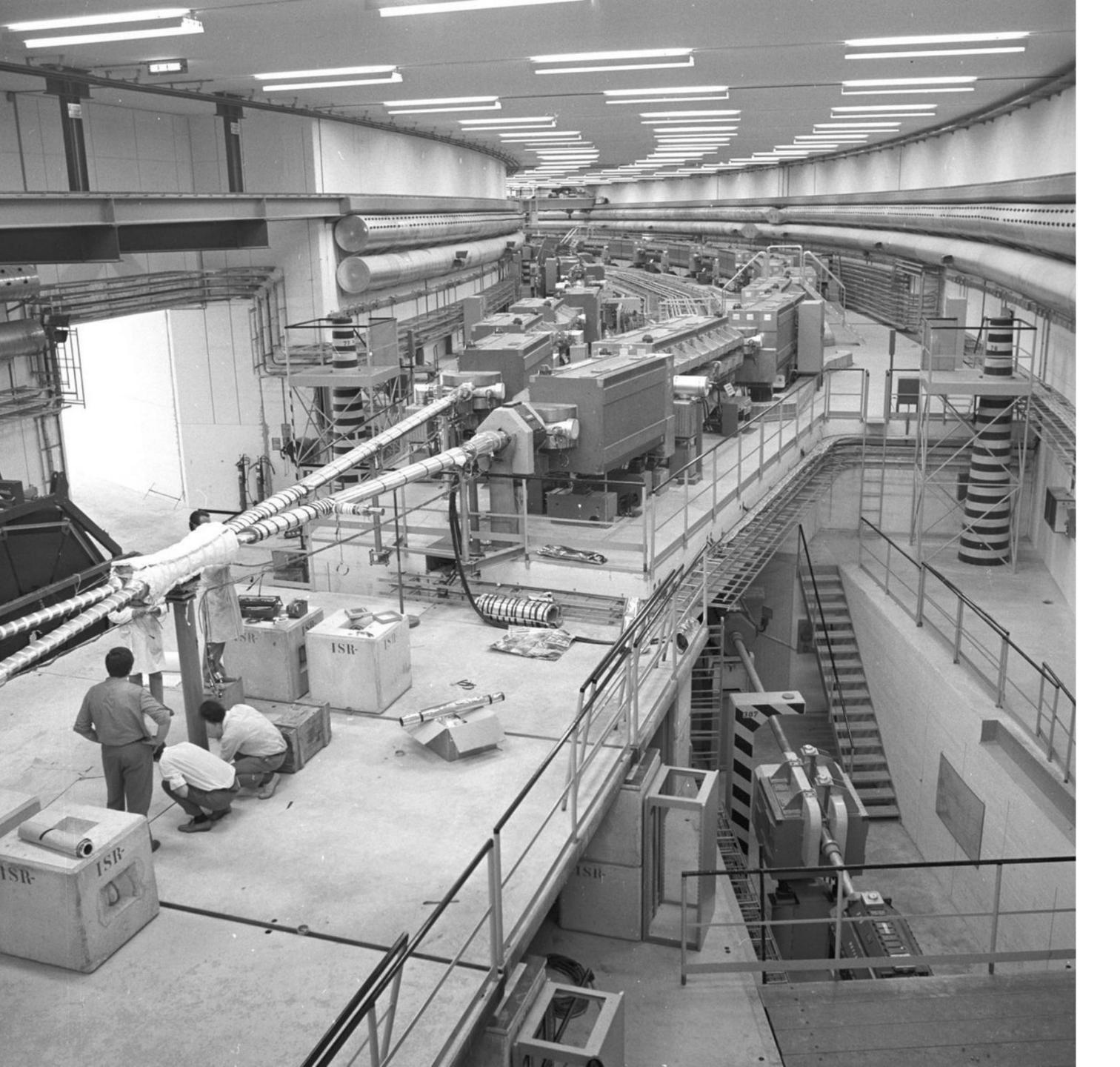
Can we find the origin of the weak scale early in the history of the Universe?





New Experimental Targets

BACKUP



CERN 1971

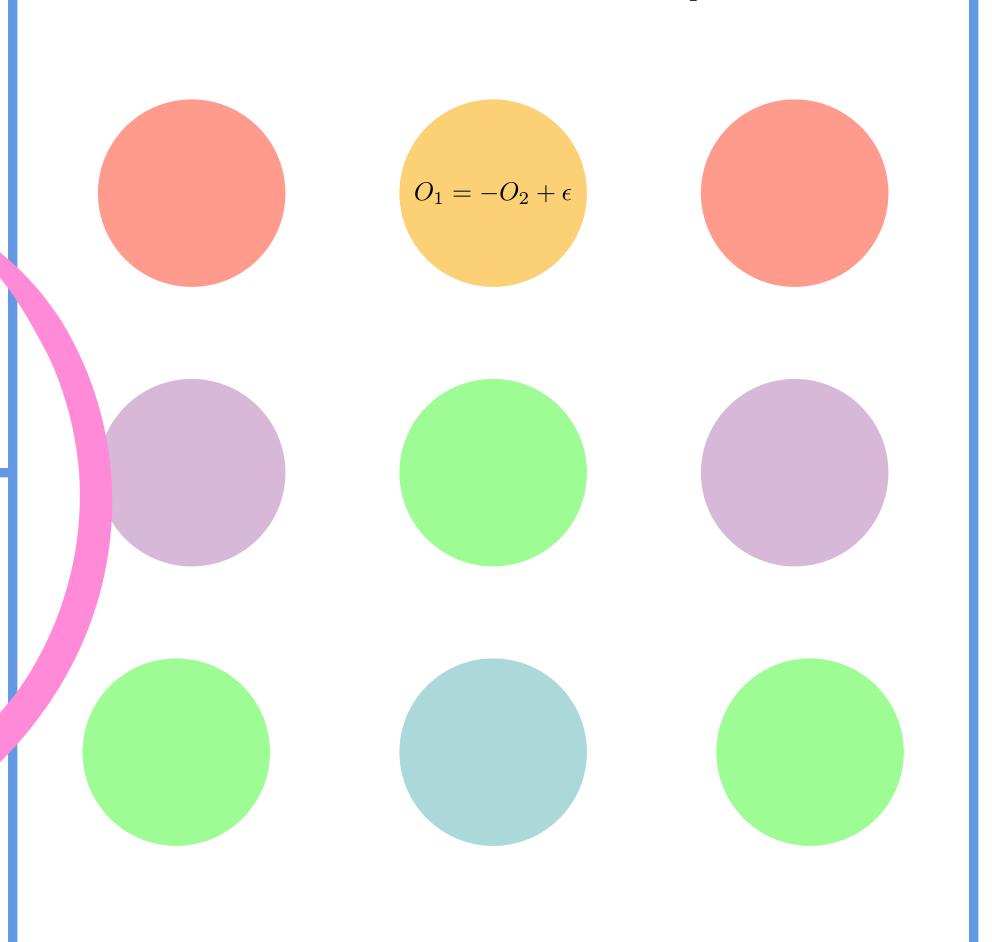
A change in theoretical perspective can win (or lose) you two Nobel Prizes



SM Landscape



$$\Lambda_S \ll M_{
m Pl}$$



[Dvali, Vilenkin '03], [Dvali '04]

$$F_4 = dA_3$$

$$S \supset \int d^4x \sqrt{-g} \left(\frac{F_4^2}{48} + M_{\rm Pl}^2 (-1 + \frac{F_4^2}{M_{\rm Pl}^2} + \ldots) |\phi|^2 + \ldots \right) + q(\phi) \int d^3\xi A_{\mu\nu\rho} \frac{\partial x^\mu}{\partial \xi^a} \frac{\partial x^\nu}{\partial \xi^b} \frac{\partial x^\rho}{\partial \xi^c} \epsilon^{abc}$$

Large initial "Electric Field" (Brown-Teitelboim)

$$F_4^2 \sim M_{\rm Pl}^4$$

[Dvali, Vilenkin '03], [Dvali '04]

$$F_4 = dA_3$$

$$S \supset \int d^4x \sqrt{-g} \left(\frac{F_4^2}{48} + M_{\rm Pl}^2 (-1 + \frac{F_4^2}{M_{\rm Pl}^2} + \ldots) |\phi|^2 + \ldots \right) + q(\phi) \int d^3\xi A_{\mu\nu\rho} \frac{\partial x^\mu}{\partial \xi^a} \frac{\partial x^\nu}{\partial \xi^b} \frac{\partial x^\rho}{\partial \xi^c} \epsilon^{abc}$$

Very slow process: requires eternal inflation!

Branes can be spontaneously nucleated (tunnelling)

[Dvali, Vilenkin '03], [Dvali '04]

$$F_4 = dA_3$$

$$S \supset \int d^4x \sqrt{-g} \left(\frac{F_4^2}{48} + M_{\rm Pl}^2 (-1 + \frac{F_4^2}{M_{\rm Pl}^2} + \ldots) |\phi|^2 + \ldots \right) + q(\phi) \int d^3\xi A_{\mu\nu\rho} \frac{\partial x^\mu}{\partial \xi^a} \frac{\partial x^\nu}{\partial \xi^b} \frac{\partial x^\rho}{\partial \xi^c} \epsilon^{abc}$$

$$\Delta F_4 = q(\phi)$$

Branes can be spontaneously nucleated (tunnelling)

[Dvali, Vilenkin '03], [Dvali '04]

$$F_4 = dA_3$$

$$S \supset \int d^4x \sqrt{-g} \left(\frac{F_4^2}{48} + M_{\rm Pl}^2 (-1 + \frac{F_4^2}{M_{\rm Pl}^2} + \ldots) |\phi|^2 + \ldots \right) + q(\phi) \int d^3\xi A_{\mu\nu\rho} \frac{\partial x^{\mu}}{\partial \xi^a} \frac{\partial x^{\nu}}{\partial \xi^b} \frac{\partial x^{\rho}}{\partial \xi^c} \epsilon^{abc}$$

The scalar mass is scanned

$$\Delta F_4 = q(\phi)$$

Cosmological Constant

Higgs Mass Squared

1. The two quantities are not calculable

2. Scale of gravity?

Theory ~ 10¹²⁰ Experiment

Theory~10³⁴ Experiment

Cosmological Constant

Higgs Mass Squared

1. The two quantities are not calculable

2. Scale of gravity?

3. Planck Scale = QFT Mass Scale?

Theory ~ 10¹²⁰ Experiment

Theory~10³⁴ Experiment

[Arkani-Hamed, RTD, Kim '20]

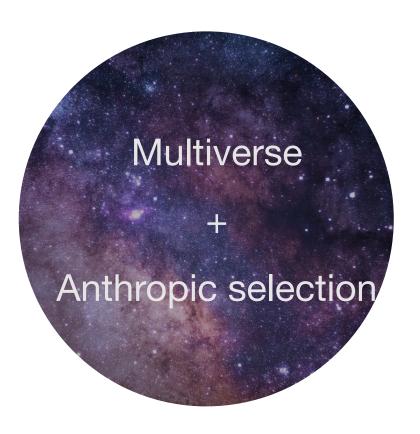
Does anything change in the SM as we vary $\langle h \rangle$?

Does anything change in the SM as we vary $\langle h \rangle$?

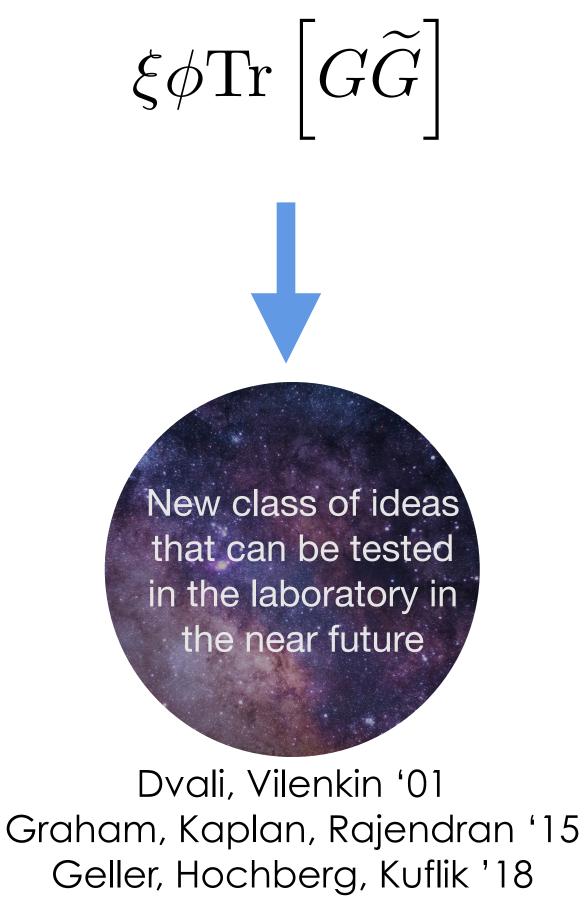
1. Obviously the spectrum

Does anything change in the SM as we vary $\langle h \rangle$?

1. Obviously the spectrum



Does anything change in the SM as we vary $\langle h \rangle$?



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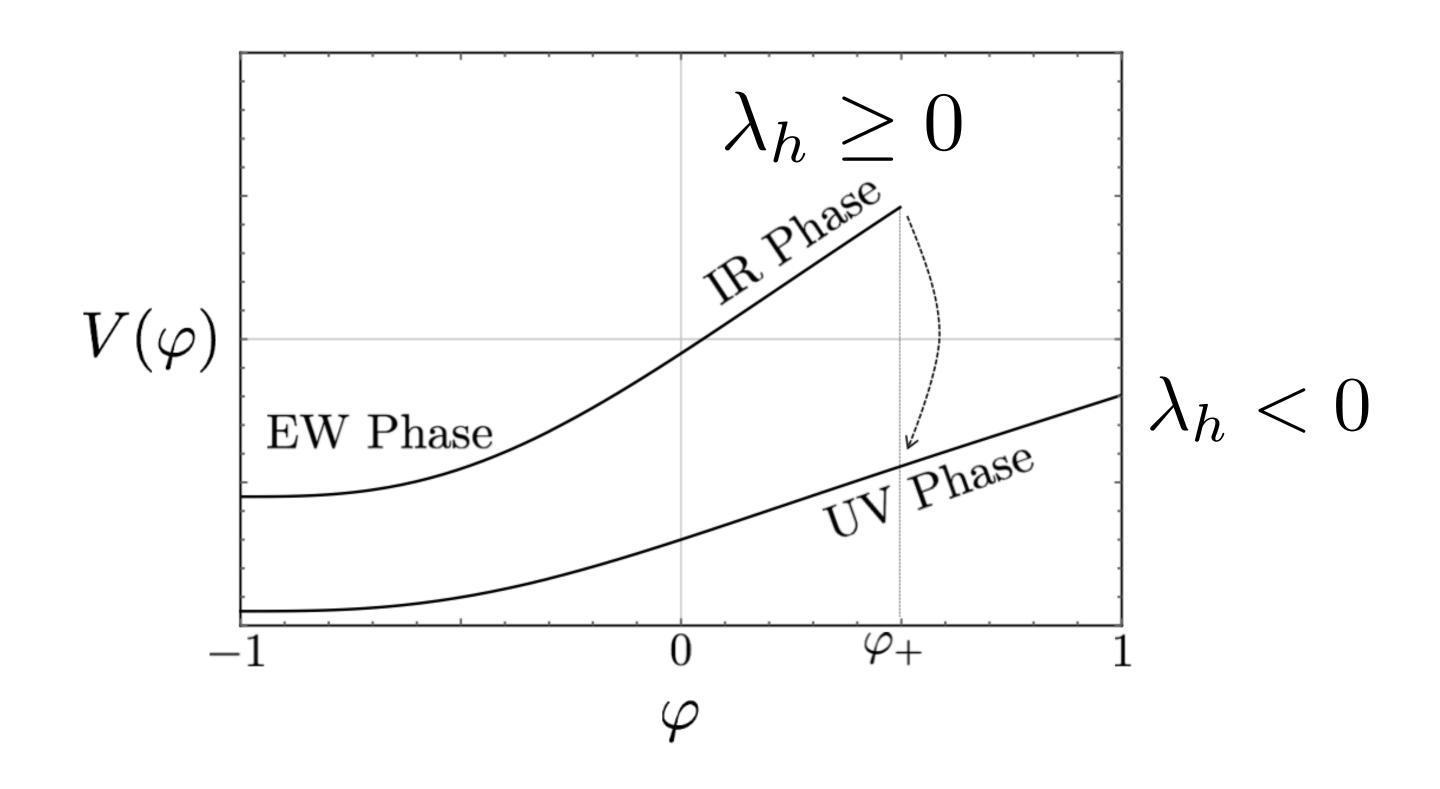
[Giudice, McCullough, You '21]

Scalar dominated by quantum dynamics during inflation

Solve Fokker-Planck Equation

In most gauges you will find that the volume is dominated by quantum dynamics for critical points of some potentials (measure problem)

[Giudice, McCullough, You '21]



Can select Higgs vev corresponding to zero quartic

[Graham, Kaplan, Rajendran '15],

$$V(\phi) = g\phi + ... + (M^2 + g\phi + ...)|H|^2 + \frac{\phi}{f}G\tilde{G}$$

[Graham, Kaplan, Rajendran '15],

$$V(\phi) = g\phi + \ldots + (M^2 + g\phi + \ldots)|H|^2 + \frac{\phi}{f}G\widetilde{G}$$

$$g\phi |H|^2$$
 Landscape of Higgs vevs

[Graham, Kaplan, Rajendran '15],

$$V(\phi) = g\phi + \ldots + (M^2 + g\phi + \ldots)|H|^2 + \frac{\phi}{f}G\widetilde{G}$$

$$\frac{\phi}{f}G\widetilde{G}$$
 Selection

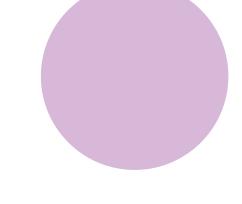
[Graham, Kaplan, Rajendran '15],

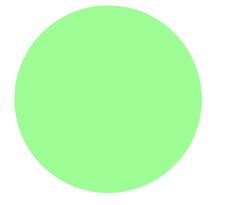
$$V(\phi) = g\phi + \dots + (M^2 + g\phi + \dots)|H|^2 + \frac{\phi}{f}G\tilde{G}$$

Symmetric Sector

$$g \ll M_{\rm Pl}^3$$

SM Landscape







$$\phi G\widetilde{G}$$

$$g\phi |H|^2$$

[Dvali, Vilenkin '03], [Dvali '04]

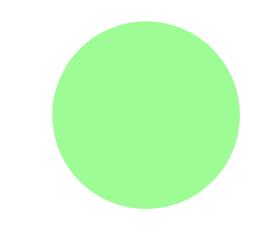
$$F_4 = dA_3$$

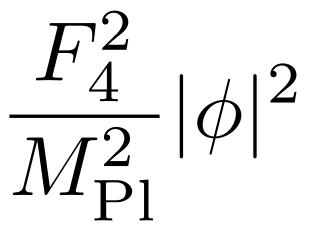
Symmetric Sector $q(\phi) \lesssim M_{\rm Pl}^2$

$$q(\phi) \lesssim M_{\rm Pl}^2$$



$$\frac{\phi^N}{M_{\rm Pl}^{N-2}} \int_{2+1} A_3$$





Planck String	Mysterious QG Blob	Mysterious QG Blob
SUSY	- SUSY	
SM	- SM	SM