



Flow-induced soft hair and the modification to the Hawking temperature

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Outline

- Motivation: Information loss paradox
- Soft hair on Schwarzschild black hole Hawking, Perry, Strominger (2016)
 - The conserved charge of BMS symmetry
- Soft hair on Vaidya black hole Chu, Koyama (2018)
- Two extensions to time-dependent supertranslation
- Hawking temperature of a dynamical black hole
- Conclusion

Information loss paradox

- Classical black hole → Information is hidden forever
- Black hole thermodynamics → Hawking radiation
- No hair theorem → **M, Q, J are enough to describe Hawking radiation**
- Unitarity → Entropy must come from somewhere (Entanglement)
- Pure thermal radiation → **Information loss paradox**
- We need **more DOFs** (secret channel) for BH?
- Evade no-hair theorem → Soft hair on black hole
 - How does the soft hair store black hole information?
- **Soft hair induced by a continuous flow in Vaidya spacetime**

BMS symmetry

- Bondi gauge $g_{rr} = g_{rA} = 0$, $\det(g_{AB}) = \det(r^2\gamma_{AB})$
- Ignoring boost, rotation & uniform translation, the most general **coordinate transformation** preserving the **asymptotically flat** condition is

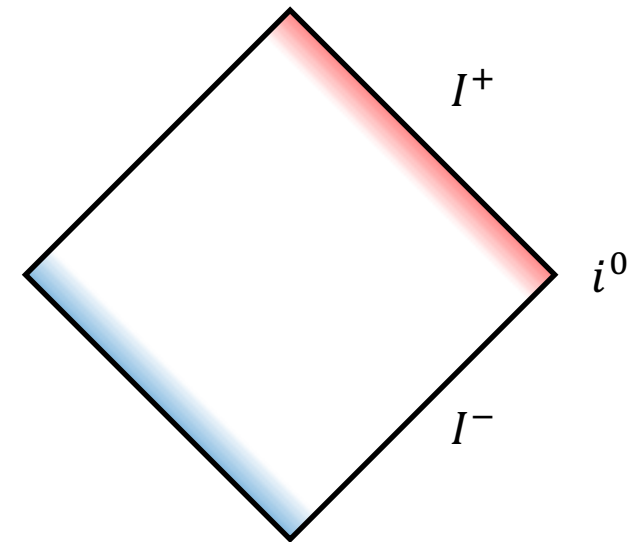
$$\delta v = f(\Theta^A), \delta r = -\frac{1}{2}D^2 f + O(r^{-1}), \delta\Theta^A = \frac{1}{r}D^A f + O(r^{-2})$$

- The generating vector field for BMS symmetry

$$\zeta_f = f\partial_v - \frac{1}{2}D^2 f\partial_r + \frac{1}{r}D^A f\partial_A$$

- The generator of BMS transformation

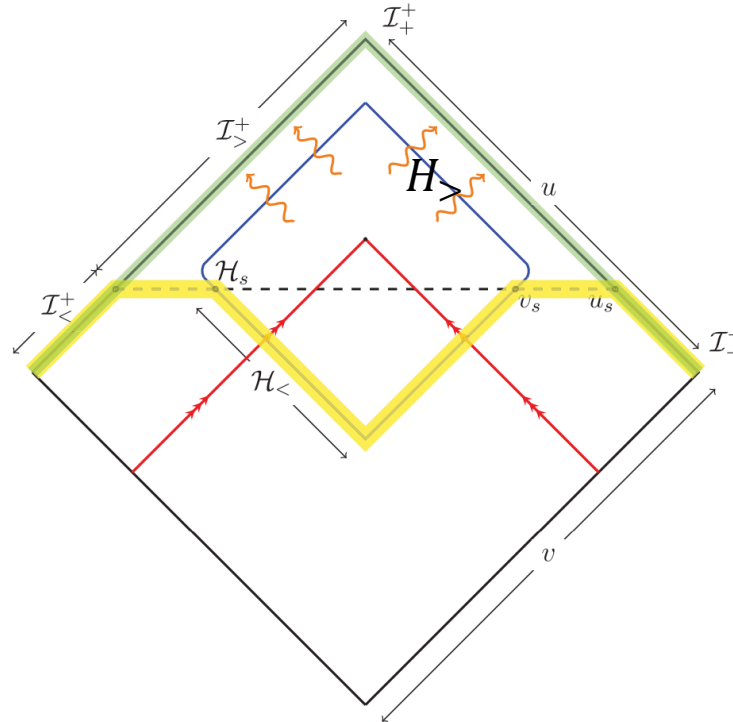
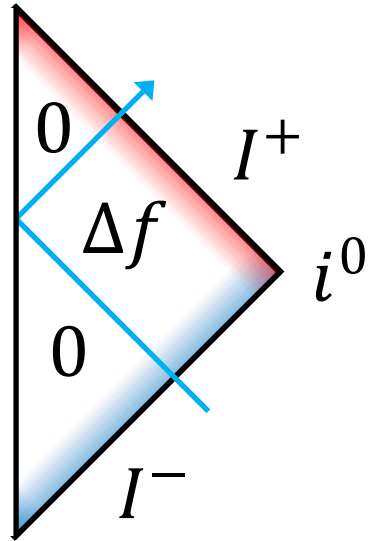
$$Q^- [f] \equiv \int_{I^-} d^2\Theta f(\Theta) \int_{-\infty}^{\infty} dv \frac{1}{32\pi G} N_{AB} N^{AB} + T_{vv}^M$$



Bondi, Metzner, Sachs (1962)

Scattering Process with Black Hole

$$Q_f^{I^+} = Q_f^{I^-}$$



$$Q_f^{I^+} \approx Q_f^{H>} + Q_f^{I^+<}$$

$$Q_f^{H<} \rightarrow Q_f^{H>}$$

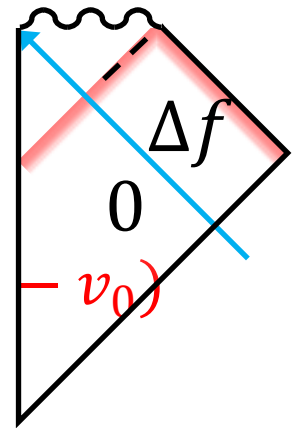
Unknown process...

$$Q_f^{I^-} = Q_f^{H<} + Q_f^{I^-<}$$

Hawking, Perry, Strominger (2016)

Implant Soft hair on black hole

- Supertranslated Schwarzschild geometry
- $ds^2 = -\left(V - \frac{MD^2f}{r^2}\right)dv^2 + 2dvdr - D_A(2Vf + D^2f)dvd\Theta^A + \left(r^2\gamma_{AB} + 2rD_AD_Bf - r\gamma_{AB}D^2f\right)d\Theta^Ad\Theta^B$, $V = 1 - 2M/r$
- **How to implant soft hair?** → Send a shock-wave from I^-
- Metric transforms as $g_{\mu\nu} \rightarrow g_{\mu\nu} + \theta(v - v_0)\mathcal{L}_f g_{\mu\nu}$
- The constraint equation $\partial_v m = \frac{1}{4}D^AD^BN_{AB} + \left(\mu + \frac{1}{4}D^2(D^2 + 2)f\right)\delta(v - v_0)$
- The perturbation of the metric is $h_{\mu\nu} = \theta(v - v_0)\left[\mathcal{L}_f g_{\mu\nu} + \frac{2\mu}{r}\delta_\mu^v\delta_\nu^v\right]$
- **Supertranslation \subset diffeomorphism** → observable only on the null shell
- Extension of time dependency?



Hawking, Perry, Strominger (2016)

Super-translated Vaidya metric

- Consider an accreting black hole
- Vaidya metric: $-Vdv^2 + 2dvdr + r^2\gamma_{AB}d\Theta^A d\Theta^B$, $V = 1 - 2M(v)/r$
- The associate energy-momentum tensor $T_{vv} = \frac{M'(v)}{4\pi r^2}$
- Super-translated Vaidya metric:
- $ds^2 = -\left(V - \frac{2M'f}{r} - \frac{MD^2f}{r^2}\right)dv^2 + 2dvdr - D_A(2Vf + D^2f)dvd\Theta^A + \left(r^2\gamma_{AB} + 2rD_AD_Bf - r\gamma_{AB}D^2f\right)d\Theta^A d\Theta^B$

Time-dependent supertranslation

- A **coordinate transformation** on spacetime: tST transformation
- $g_{\mu\nu} \rightarrow g_{\mu\nu} + \mathcal{L}_{f(v)}g_{\mu\nu}$, $T_{\mu\nu} \rightarrow T + \mathcal{L}_{f(v)}T_{\mu\nu}$
- The tSTed Vaidya spacetime
- $ds^2 = -\left(V - \frac{2M'f}{r} - \frac{MD^2f}{r^2} + (D^2 + 2V)f'\right)dv^2 + 2(1 + f')dvdr - D_A(D^2f + 2Vf + 2rf')dvd\Theta^A + (r^2\gamma_{AB} + 2rD_AD_Bf - r\gamma_{AB}D^2f)d\Theta^Ad\Theta^B$
- $\mathcal{L}_{f(v)}T_{vv} = \frac{M'f'}{2\pi r^2} + \frac{M''f}{4\pi r^2} + \frac{M'D^2f}{4\pi r^3}$, $\mathcal{L}_{f(v)}T_{vA} = \frac{M'D_Af}{4\pi r^2}$
- Recover BMS transformation when $f' = 0$
- What's the effect of tST transformation on physical observable?

Doppler and Unruh effect

- tST transformation can be recast as a dressing process on a scalar fields and the Fourier basis up to $O(f)$ near the horizon

$$\phi(v) = \int_0^\infty \frac{dp}{\sqrt{2\pi p}} (a_p e^{-ipv} + a_p^\dagger e^{ipv}), \phi(u) = \int_0^\infty \frac{dp}{\sqrt{2\pi p}} (b_p e^{-ipu} + b_p^\dagger e^{ipu})$$

- Consider almost radially outgoing modes in the Eikonal limit $q = \omega du + Y_A d\Theta^A + O(f^2)$

- We obtain the dressing : $e^{i\omega f} : \hat{b}_\omega \approx b_\omega$, where $b_\omega = b_q, q_u = \omega, f = f(\partial_\omega)$

- Imposing the adiabatic condition $Mf''(u) \ll f'(u)$, the modified Hawking spectrum is

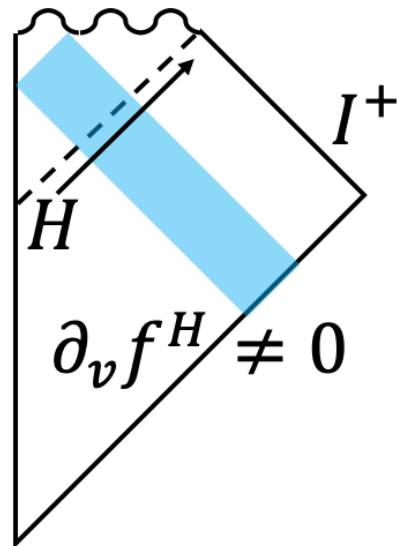
$$\langle g(q) b_q^\dagger b_q \rangle = \langle g(q) \hat{b}_q^\dagger \hat{b}_q \rangle |_{q \rightarrow (1+f')q, \kappa \rightarrow \kappa - f''}$$

- **tST is a coordinate transformation on I preserves all physical observable up to a non-trivial Bogoliubov transformation**

- Notice that the dressing of incoming modes is irrelevant to the Hawking radiation. The black hole information requires another channel to restore.

Flow-induced supertranslation

- Can supertranslation has non-trivial effect on Hawking radiation?
- Consider an anisotropic null flow into the black hole
- The metric is transformed by a flow-induced supertranslation (FST), which is **not** a coordinate transformation.
- $g_{\mu\nu} \rightarrow g_{\mu\nu} + \mathcal{L}_f g_{\mu\nu}|_{f \rightarrow f(v)}$
- $T_{\mu\nu} \rightarrow T_{\mu\nu}[g_{\mu\nu} \rightarrow g_{\mu\nu} + \mathcal{L}_f g_{\mu\nu}|_{f \rightarrow f(v)}]$
- **FST is the only form for sub-leading part of the initial condition at I^- that leads to a stable asymptote** (Christodoulou, Klainerman)

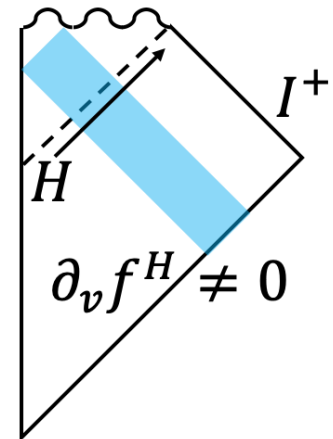


Linear response relation

- BMS transformation: $T_{\mu\nu} \rightarrow T_{\mu\nu} + \mathcal{L}_f T_{\mu\nu}$
- Shock-wave induced BMS transformation (special case of FST) :
- $T_{\mu\nu} \rightarrow T_{\mu\nu} + \mathcal{L}_{\theta(v-v_0)f} T_{\mu\nu} + T_{\mu\nu}^S$
- FST transformation:
- $T_{\mu\nu} \rightarrow T_{\mu\nu} + \underbrace{\mathcal{L}_{f(v)} T_{\mu\nu}}_{T_{\mu\nu}^{tST}} + \underbrace{\int_{v_0=-\infty}^{v_0=v} T_{\mu\nu}^S(v_0)}_{T_{\mu\nu}^{FST}}$
- $T_{\mu\nu}^{FST} - T_{\mu\nu}^{tST} = \int_{v_0=-\infty}^{v_0=v} T_{\mu\nu}^S(v_0), T_{\mu\nu}^S \equiv (T_{\mu\nu}^{FST} - T_{\mu\nu}^{tST})|_{f(v) \rightarrow f'(v_0)\theta(v-v_0)}$
- Dynamically generate linearized conserved charge of BMS symmetry

Exact conserved charge

- The linear response relation between f and T at $O(f)$ is satisfied under a **specific tST transformation** → A linear conserved charge of BMS symmetry.
- How about at $O(f^n)$? → In fact, T has the form $T_1(f') + \partial_t T_2(f')$ → Composed by local quantity f' and a boundary term!
- There exist a conserved charge along the direction $v \rightarrow \nabla^\mu (T_{\mu\nu} n_r^\nu) = 0$ where $n_r^\nu \partial_\nu = \partial_r$ is the null congruence of v .
- n_r is an asymptotic Killing vector field → **the combined transformation of tST and FST generates BMS charge at I^- dynamically!**



Exact supertranslated black hole

- With the combined transformation of tST and FST, we find the appropriate coordinate to study the interaction between Hawking radiation of dynamical black hole and flow-induced soft hair

$$\begin{aligned}
 ds^2 &= g_{\mu\nu} + \mathcal{L}_f g_{\mu\nu} \Big|_{f \rightarrow f(v)} - \mathcal{L}_{f(v)} g_{\mu\nu} \\
 &= - \left((1 - 2f')(1 - 2M/r) - D^2 f' \right) dv^2 + 2(1 - f') dv dr - 4r D_A f' dv d\Theta^A + r^2 \gamma_{AB} d\Theta^A d\Theta^B
 \end{aligned}$$

- In this coordinate, we found the surface gravity is modified as

$$\kappa_{\mathcal{H}} = \frac{1}{4M} \left(1 + \frac{D^2 \psi}{M} - f' + \chi' \right) - f'' + \chi'' + O(f^2)$$

- The lensing potential $\psi(v) = -\frac{1}{4M} \int_0^\infty e^{-\int_v^{v_1} \frac{1}{4M(v_2)} dv_2} D^2 f'(v_1) dv_1 + O(f^2)$
- The anisotropy of the temperature is hard to be eliminated by a coordinate transformation! → Observable soft hair?

Observing soft hair

- To verify the observability of soft hair, we also consider the additional lensing effect on Hawking radiation intensity due to the same anisotropic flow

$$\Delta I_H = \frac{1}{4M} \left(-(D^2 + 1)f(v) + \frac{\int_v^\infty dv_1 D^2 f'(v_1)}{\omega\left(\frac{v_1 - v}{4M}\right)} \right) + O(f^2, M'), \omega^{-1}(x) = x + \ln x$$

and compare the form to $\psi(v) = -\frac{1}{4M} \int_0^\infty e^{-\int_v^{v_1} \frac{1}{4M(v_2)} dv_2} D^2 f'(v_1) dv_1 + O(f^2)$

- The different modification to two observables can isolate the dressing effect and won't be simultaneously compensated by a coordinate transformation!

Conclusions

- The effect of dressing on the Hawking temperature is equivalent to the Doppler effect and Unruh effect generated by a tST transformation.
- We generalized the shock-wave induced supertranslation to an anisotropic null flow (FST) and preserve the asymptotic fall-off conditions.
- The FST transformation together with a suitable choice of tST transformation can dynamically generate linear BMS charge (soft hair) on black hole horizon.
- The flow-induced soft hair can be observed by the Hawking temperature and Hawking radiation intensity that extract the **macroscopic** information of Hawking radiation.

Thank you for your attention