



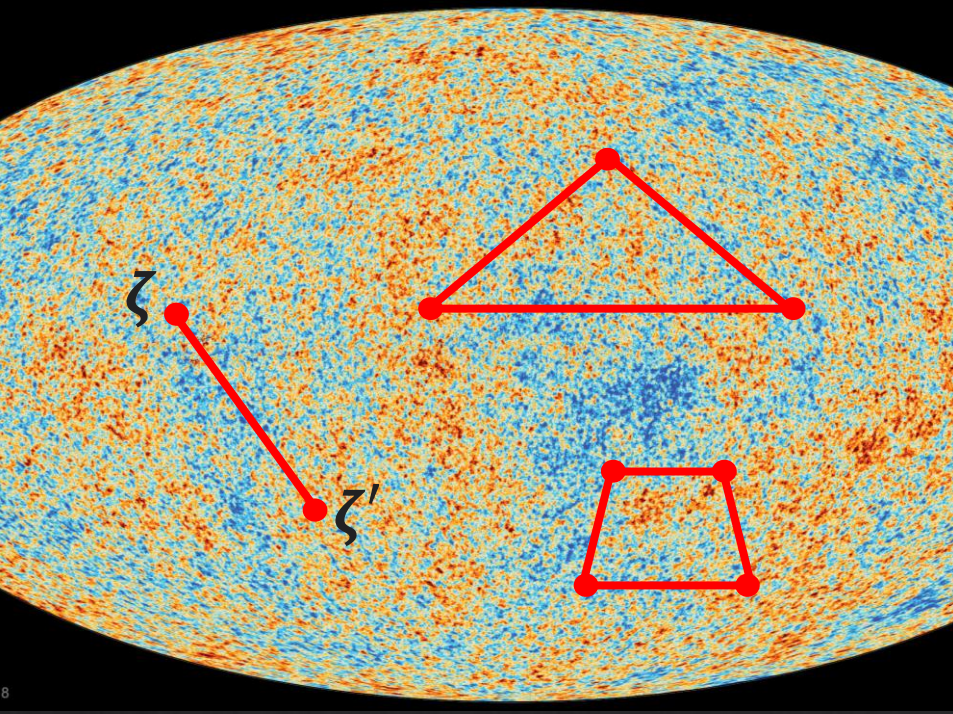
Disentangling mass spectra of multiple fields in cosmological collider

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Based on JHEP 04 (2021) 127
with Masahide Yamaguchi

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Inflationary observables :
Correlation functions of ζ , γ_{ij}

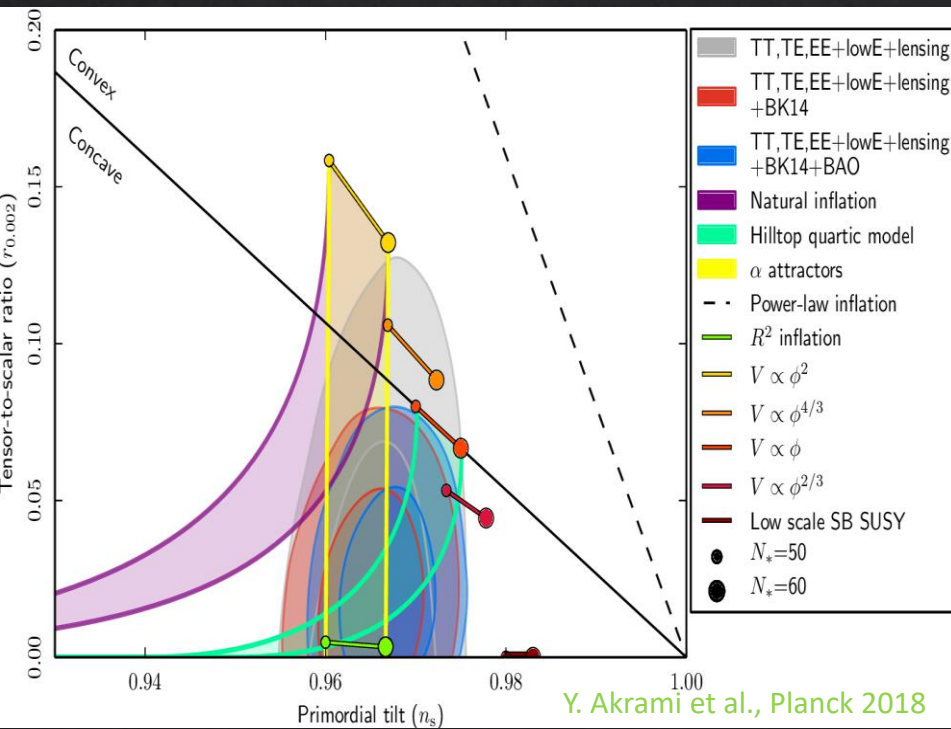
Strong restrictions on inflation models
(high energy physics)

* 2 pt. correlation & tilt

➔ e.g., Inflaton potential

* 3 or higher pt. correlations

➔ e.g.,
multi field inflation scenario,
higher derivative couplings,
...

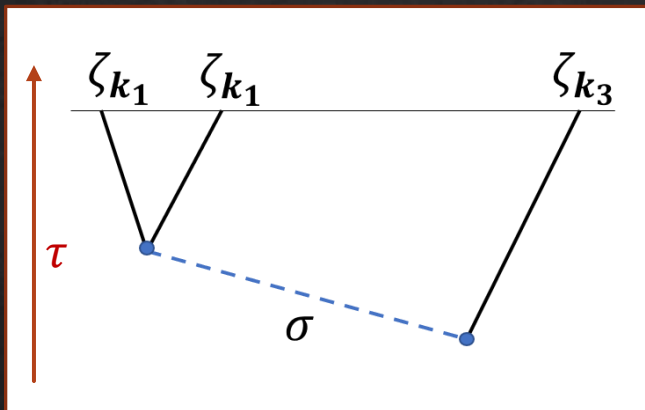


Y. Akrami et al., Planck 2018

Cosmological collider

- * Higher pt. functions contain robust information (mass & spin) of heavy ($m_\sigma \sim H$) particles σ

e.g., a scalar σ



oscillation

$$S (\propto \langle \zeta^3 \rangle) \sim C(\mu) \left(\frac{k_l}{k_s} \right)^{2i\mu} + \text{c.c.}$$

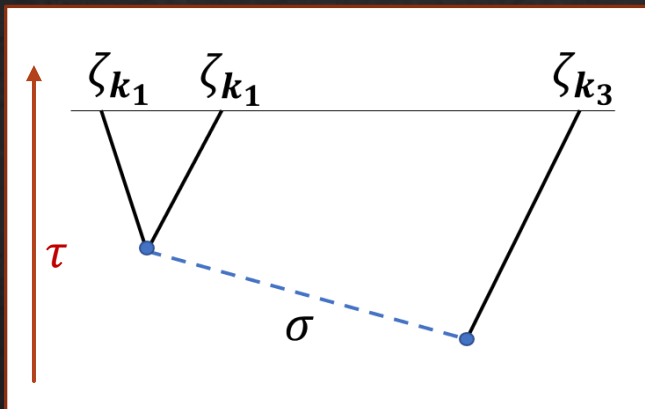
$$\mu \equiv \sqrt{\left(\frac{m_\sigma}{H} \right)^2 - \frac{9}{4}}$$

squeezed limit ($k_l \equiv k_3 \ll k_{1,2} \equiv k_s$)

Cosmological collider

- * Higher pt. functions contain robust information (mass & spin) of heavy ($m_\sigma \sim H$) particles σ

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$$S (\propto \langle \zeta^3 \rangle) \sim C(\mu) \left(\frac{k_l}{k_s} \right)^{2i\mu} + \text{c.c.}$$

$$\mu \equiv \sqrt{\left(\frac{m_\sigma}{H} \right)^2 - \frac{9}{4}}$$

- * Explore very high scale :
 $m_\sigma \sim H \sim 10^{13} \text{ (GeV)} \gg$ energy scale of terrestrial experiment

Many works

- * Quasi single field inflation (Chen, Wang, '10)
- * Signature of SUSY (Baumann, Green, '12)
- * General int. by EFT approach (Noumi, Yamaguchi, Yokoyama, '13)
- * Spinning particle (Arkani-Hamed, Maldacena, '15)
- * Standard model particles, Neutrino (Chen, Wang, Xianyu, '16)
- ...

Question

✓ **single** particle with a definite mass and spin



Q. What about **multiple** particles case?

Question

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Q. What about **multiple** particles case?

Natural setup in realistic situation

* SM particles acquire masses $\sim H$ by loop effects Chen, Wang, Xianyu, '16

* Supergravity

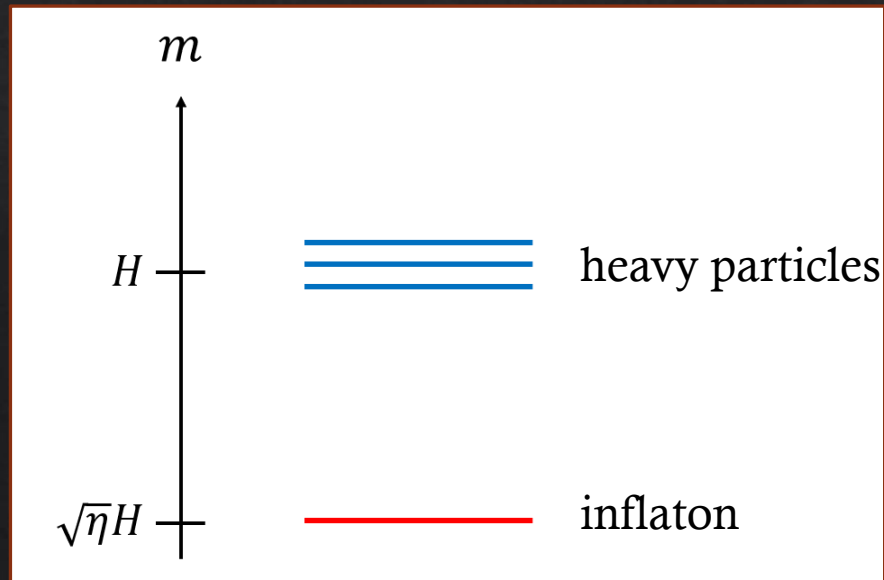
Hubble induced masses for scalar particles at tree level Stewart '95

Question

✓ **single** particle with a definite mass and spin



Q. What about **multiple** particles case?



Question

✓ **single** particle with a definite mass and spin



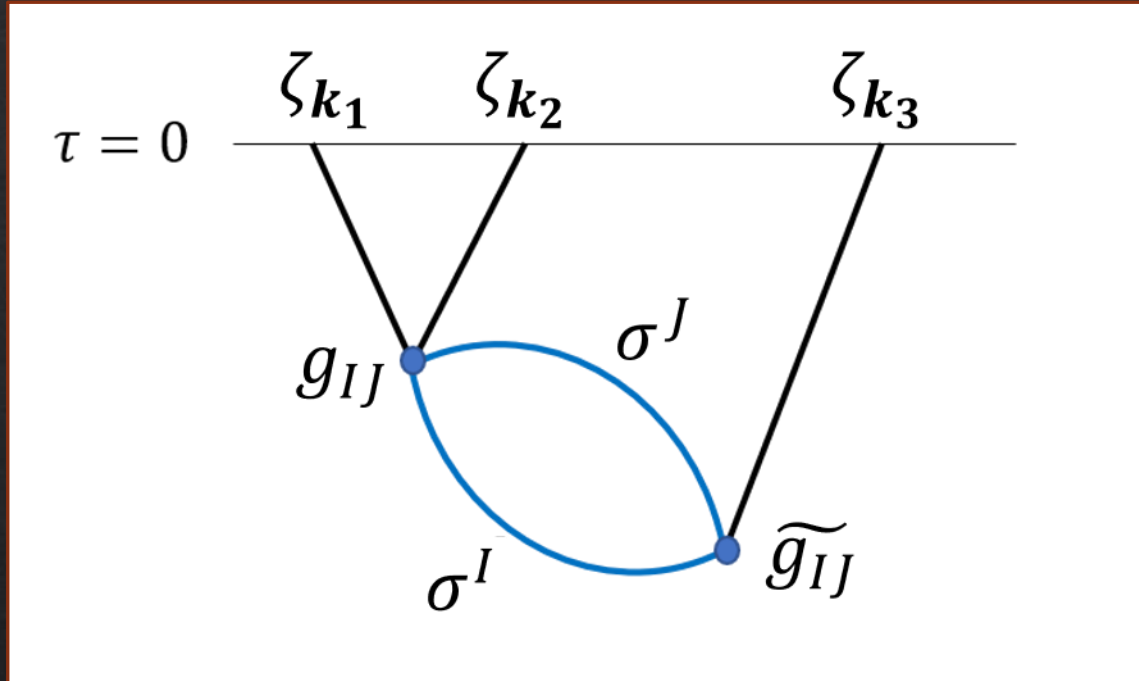
Q. What about **multiple** particles case?

More specifically,

Q. Can we distinguish single and multi cases?
(specific signal in multi case?)

Q. How to disentangle degenerate mass spectra?

Setup



massive scalars : σ^I ($I = 1, \dots, n$)
 $g_{IJ}, \widetilde{g}_{IJ}$: non-diagonal in general
 $g_{IJ} = \widetilde{g}_{IJ}$ for simplicity

Result (in squeezed limit)

$$S = \sum_{I,J} S_{IJ}$$

$I = J$

$$S_{II} \propto g_{II}^2 C(\mu_I, \mu_I) \left(\frac{k_l}{k_s} \right)^{2i\mu_I} + \text{const.} + \text{c.c.}$$

← same as single
but **superposed**

$I \neq J$

$$S_{IJ} \propto g_{IJ}^2 C(\mu_I, \mu_J) \left(\frac{k_l}{k_s} \right)^{i(\mu_I + \mu_J)} + g_{IJ}^2 C(\mu_I, -\mu_J) \left(\frac{k_l}{k_s} \right)^{i(\mu_I - \mu_J)} + \text{c.c.}$$

← **mixing**
(new effect)

Result (in squeezed limit)

$$S = \sum_{I,J} S_{IJ}$$

$I = J$

$$S_{II} \propto g_{II}^2 C(\mu_I, \mu_I) \left(\frac{k_l}{k_s} \right)^{2i\mu_I} + \text{const.} + \text{c.c.}$$

High frequency

$$\mu^I \equiv \sqrt{\left(\frac{m^I}{H} \right)^2 - \frac{9}{4}}$$

$I \neq J$

$$S_{IJ} \propto g_{IJ}^2 C(\mu_I, \mu_J) \left(\frac{k_l}{k_s} \right)^{i(\mu_I + \mu_J)} + g_{IJ}^2 C(\mu_I, -\mu_J) \left(\frac{k_l}{k_s} \right)^{i(\mu_I - \mu_J)} + \text{c.c.}$$

High frequency

low frequency (modulation)

* easily distinguished

* specific to multi particles

Example : $n=2$ (two fields)

$$S = S_{11} + S_{22} + S_{12}$$

- Case (i) : without mixing ($g_{12} = 0$)
- Case (ii) : with mixing ($g_{12} \neq 0$)

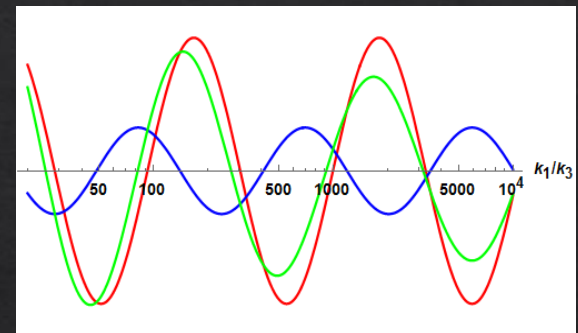
Note :

- * Boltzmann suppression factor in $C(\mu^I, \mu^J)$
- * When $\mu^1 \ll \mu^2 \rightarrow |S_{22}| \ll |S_{12}| \ll |S_{11}|$

Case (i) : Without mixing ($g_{12} = 0$ & $g_{11} = g_{22}$)

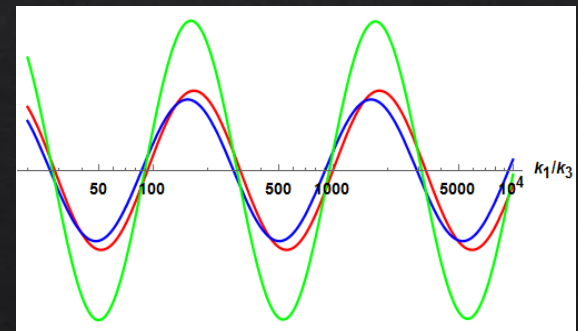
$S = S_{11} + S_{22}$: just superposition of two independent modes with μ^l

- ① $m_1 \ll m_2$
→ light mode dominates
→ $S_{\text{total}} \sim S_{11}$



$$(m_1, m_2)/H = (2, 2.5)$$

- ② $m_1 \sim m_2$
→ $S_{11} \sim S_{22}$
→ $S_{\text{total}} \sim 2 \times S_{11}$

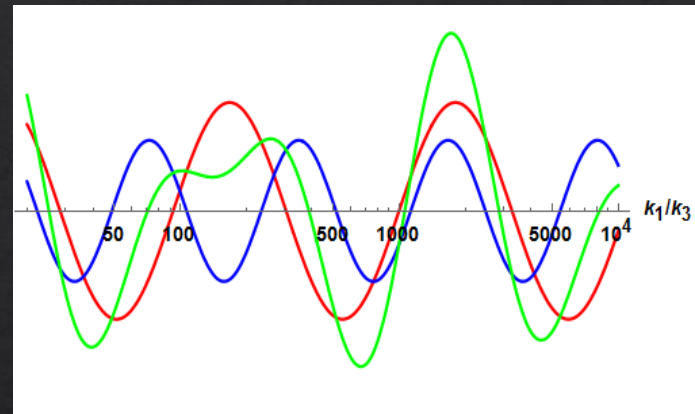


$$(m_1, m_2)/H = (2, 2.1)$$

Nothing interesting happens though a slight difference appears

Case (i) : Without mixing ($g_{12} = 0$ & $g_{11} \ll g_{22}$)

③ $m_1 \ll m_2$, but $g_{11} \ll g_{22}$
→ g_{22} compensates suppression



$(m_1, m_2)/H = (2, 2.5)$ and $g_{11}/g_{22} = 0.1$

- * Nontrivial superposition
= proof of multi particles with hierarchical couplings
- * However, such hierarchy seems to be unnatural.

Case (ii) : With mixing ($g_{11} = g_{22} = g_{12}$)

$$S = S_{11} + S_{22} + S_{12}$$

① when $\mu^1 \ll \mu^2 \rightarrow |S_{22}| \ll |S_{12}| \ll |S_{11}|$
($S_{\text{total}} \sim S_{11}$)

② $\mu^1 \sim \mu^2$ degenerate case

Case (ii) : With mixing ($g_{11} = g_{22} = g_{12}$)

$$S_{12} \propto g_{12}^2 C(\mu_1, \mu_2) \left(\frac{k_l}{k_s} \right)^{i(\mu_1 + \mu_2)} + g_{12}^2 C(\mu_1, -\mu_2) \left(\frac{k_l}{k_s} \right)^{i(\mu_1 - \mu_2)} + \text{c.c.}$$

High frequency

low frequency
(modulation)

Case (ii) : With mixing ($g_{11} = g_{22} = g_{12}$)

$$S_{12} \propto g_{12}^2 \underbrace{C(\mu_1, \mu_2)}_{\text{High frequency}} \left(\frac{k_l}{k_s} \right)^{\boxed{i(\mu_1 + \mu_2)}} + g_{12}^2 \underbrace{C(\mu_1, -\mu_2)}_{\text{low frequency (modulation)}} \left(\frac{k_l}{k_s} \right)^{\boxed{i(\mu_1 - \mu_2)}} + \text{c.c.}$$

In degenerate limit ($\mu_1 \sim \mu_2 = \mu$)

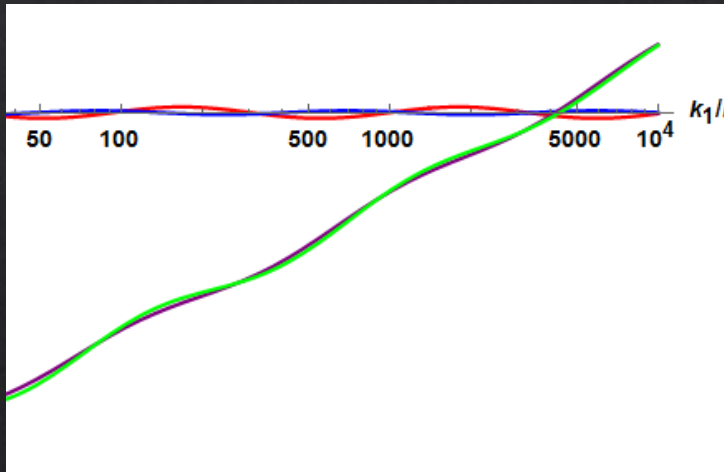
$$\left| \underbrace{C(\mu_1, \mu_2)}_{\text{High frequency}} / \underbrace{C(\mu_1, -\mu_2)}_{\text{low frequency (modulation)}} \right| \sim 2 \times 10^{-2} \times \mu^{-5/2} \ll 1$$



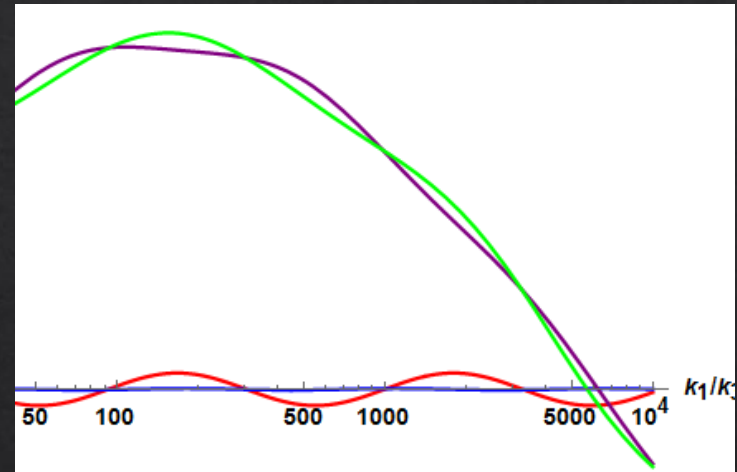
The total signal

~ characterized by the low frequency mode (large wavelength)

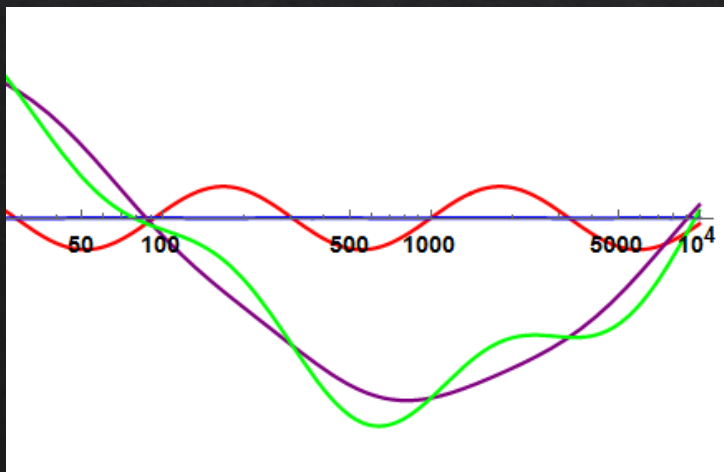
Case (ii) : With mixing ($g_{11} = g_{22} = g_{12}$)



$(m_1, m_2)/H = (2, 2.1)$



$(m_1, m_2)/H = (2, 2.3)$



$(m_1, m_2)/H = (2, 2.5)$

- * $S_{\text{total}} \sim S_{12} !!$
- * large wavelength
- * disentangle the spectra

Summary

Multi particles effects on cosmological collider

- ✓ Superposition

- * Nontrivial waveforms appear when there exists a hierarchy in m^I and g_{IJ} (rare case)

- ✓ Mixing effect

- * Exists in general
- * S is characterized by wavelength with $1/|\mu^1 - \mu^2|$, when $\mu^1 \sim \mu^2$
- * Easily specified \rightarrow existence of degenerate particles
- * Depends on details of interaction (e.g., does not exist in derivative interactions)

Thank you very much

Appendix

Interaction

$$\mathcal{L}_{\text{int}}(\phi, \sigma) = f(\phi, X) g_{IJ} \sigma^I \sigma^J$$

$X \equiv (\partial\phi)^2$, neglect it just for simplicity
 g_{IJ} : real and symmetric const.



$$\begin{aligned}\mathcal{L}_{\text{int}}^3 &= a^3 c_3 \delta\phi g_{IJ} \sigma^I \sigma^J \\ \mathcal{L}_{\text{int}}^4 &= a^3 c_4 (\delta\phi)^2 g_{IJ} \sigma^I \sigma^J\end{aligned}$$

$$\begin{aligned}c_3 &\equiv f_\phi \\ c_4 &\equiv f_{\phi\phi}/2\end{aligned}$$

Assume
they are slowly varying \sim const.