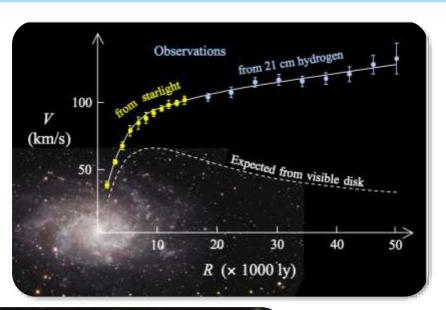
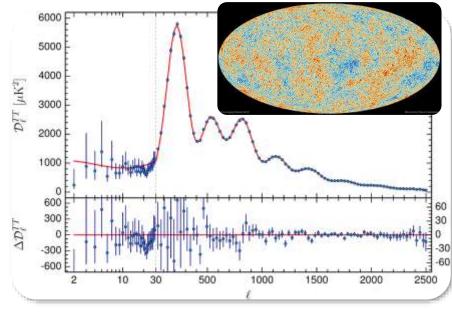
Searching for Super-Light Dark Matter with a Graphene Josephson Junction Detector

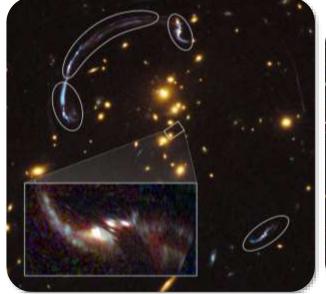
with D. Kim, K.C. Fong & G.-H. Lee [arXiv: 2002.07821]



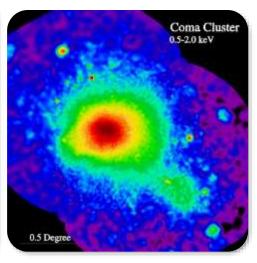
Observational Evidence for DM













Classic Solution*: WIMP

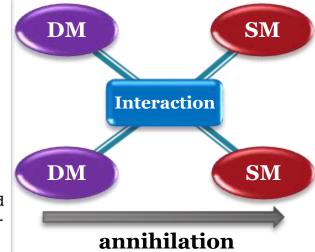
Cosmological Lower Bound on Heavy-Neutrino Masses

Beniamin W. Lee^(a) Fermi National Accelerator Laboratory, (b) Batavia, Illinois 60510

and

Steven Weinberg(c) Stanford University, Physics Department, Stanford, California 94305 (Received 13 May 1977)

The present cosmic mass density of possible stable neutral heavy leptons is calculated in a standard cosmological model. In order for this density not to exceed the upper limit of 2×10^{-29} g/cm³, the lepton mass would have to be greater than a lower bound of the order of 2 GeV.



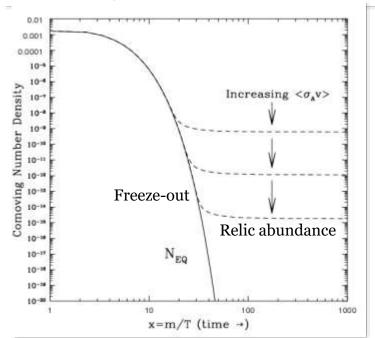
> Correct thermal relic abundance:

 $\Omega h^2 \sim \frac{0.1 \ pb}{\langle \sigma v \rangle}$ with $\langle \sigma v \rangle \sim \frac{\alpha_X^2 m_\chi^2}{M^4}$ (M: dark scale/mediator)

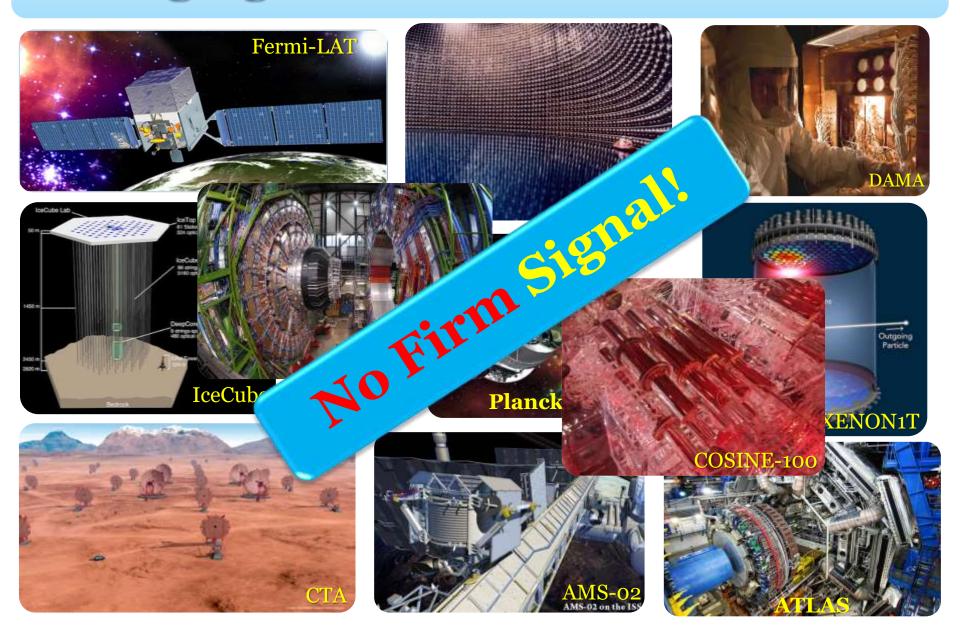
- ➤ Weak coupling → naturally weak scale mass:
 - ~1 GeV 10 TeV mass range favored
 - → weak scale (new) physics

Of course, **axion** is another classic solution.

(Luca Visinelli's talk)



Diverging Efforts for WIMP Searches

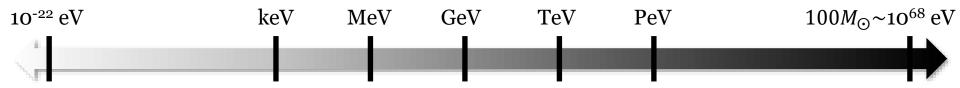


Only WIMP?

No!



DM Landscape: A Very Wide Mass Range



DM Landscape: A Very Wide Mass Range



Ultralight
Scalar field
(QCD) Axion
Hidden photon
Fuzzy
non-thermal
→ Luca

Superlight Light

WDM particle

Sterile v DM

axino can be

may be thermal

thermal

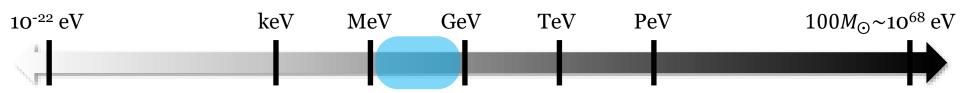
wellmotivated
extensively
studied
Thermal

Javier

Superheavy
/ Composite
WIMPzilla
Q-ball
Fermi-ball
Dark-quark
nugget
non-thermal
JeongPyong

MACHO /
Primordial
BHs
non-particle
Daniele,
Takahiro

Light DM Sector



Light particle

DM

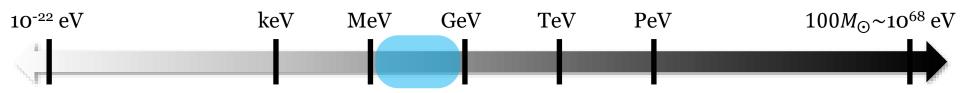
- For heavy mediator, $\langle \sigma v \rangle \sim \frac{\alpha_X^2 m_\chi^2}{M^4}$
- For weak scale physics,

sub-GeV DM overproduction

- → New mediator $< M_W$ for freeze-out or freeze-in, ...
- Various light DM/mediator scenarios:
 - ✓ MeV DM for GC 511 keV line observation [astro-ph/0309686]
 - ✓ Secluded MeV DM [arXiv:0711.3528, 0711.4866]
 - ✓ Sommerfeld enhancement for e^+ excess [arXiv:0810.0713]
 - ✓ $(g-2)_{e,\mu}$: ~2 3 σ discrepancy [arXiv:1806.10252]
 - ✓ New ν interactions for the MiniBooNE excess [arXiv:1807.09877]
 - ✓ Solutions of Yukawa coupling hierarchy prob. [arXiv:1905.02692]
 - **√** ...

- New DM relic determination mechanisms:
 - ✓ Assisted Freeze-Out [arXiv:1112.4491]
 - ✓ Cannibal DM [arXiv:1602.04219, 1607.03108]
 - ✓ Co-Decaying [arXiv:1105.1652, 1607.03110]
 - ✓ Semi-Annihilation [arXiv:0811.0172, 1003.5912]
 - ✓ SIMP [arXiv:1402.5143]
 - ✓ ...

Light DM Sector



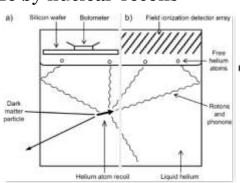
- ❖ $E_k \sim mv^2 < O(\text{keV})$ with $v \sim 10^{-3}$: Light $< E_r^{th}$ of typical DM direct detectors DM for nuclear recoils
- New ideas for $low E_r^{th} w/e$ -recoil are required!
 - ✓ Ionization by e-recoils (semiconductor)

 [arXiv:1108.5383, 1509.01598]
 - ✓ Ejection of e's (graphene, C-nanotube)

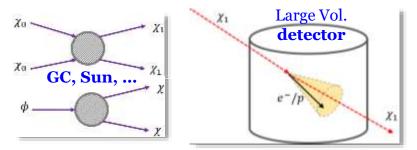
 [arXiv:1606.08849, 1706.02487, 1808.01892]
 - ✓ Evaporation of He by nuclear-recoils

[arXiv:1706.00117] a)

√ ...

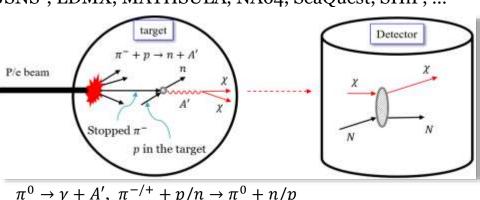


Cosmogenic boostedDM searches: COSINE-100, DUNE/ProtoDUNE, IceCube, SK/HK/KNO, ...

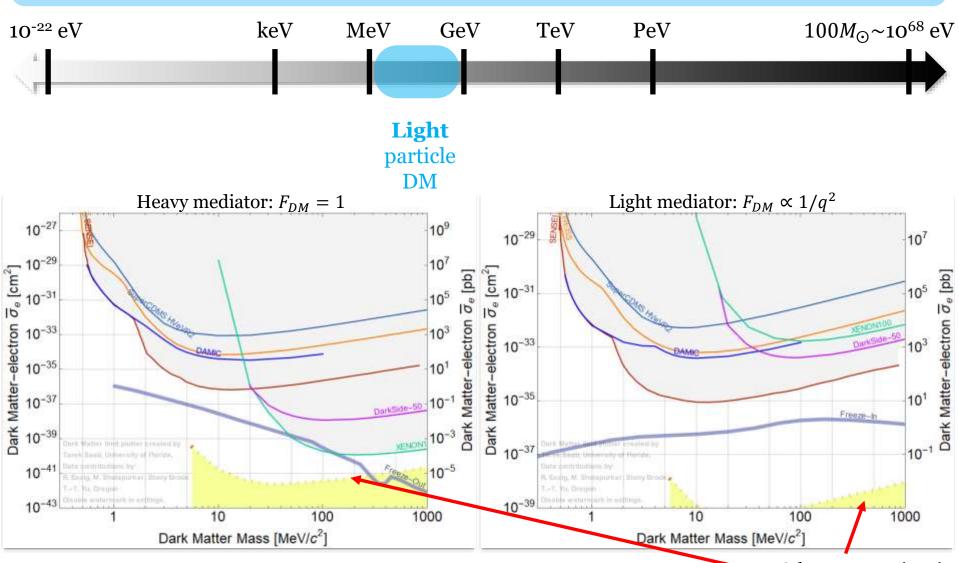


❖ Beam-produced light DM/mediator searches:

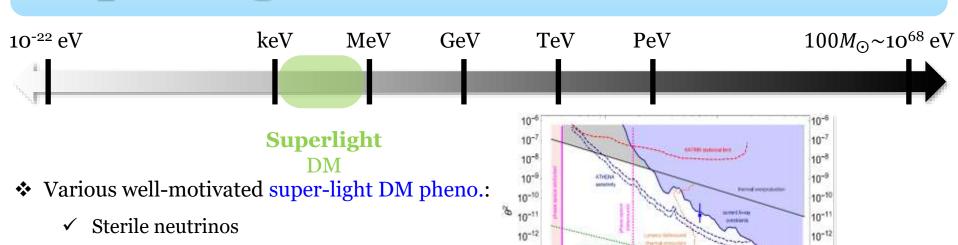
Babar, BDX, Belle-II, CCM, COHERENT, DUNE, FASER, JSNS², LDMX, MATHSULA, NA64, SeaQuest, SHiP, ...



Light DM: Direct Search Current Status



Super-Light DM: Main Focus



10-13

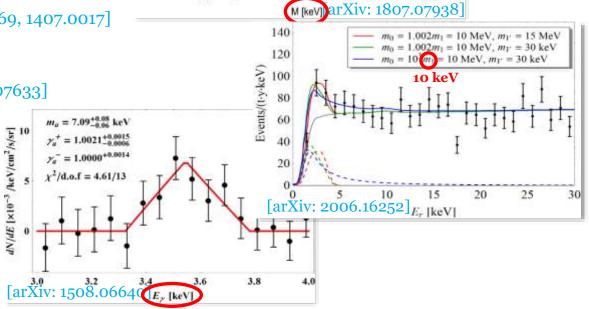
10⁻¹⁴

0.5

- ✓ Mirror ν DM [hep-ph/9505385]
- ✓ Axino/gravitino [arXiv: 0902.0769, 1407.0017]
- ✓ Axion-like particles[arXiv:0912.0015, 1407.0017, 1510.07633]

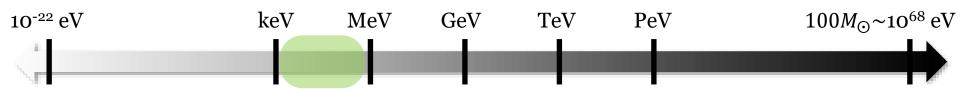
[hep-ph/9303287, astro-ph/9810076]

- ✓ Super-light dark gauge bosons
 [arXiv:1105.2812, 1201.5902]
- ✓ Decaying DM for 3.5 keV line [arXiv:1403.1536, 1508.06640]
- ✓ keV DM for XENON1T, ...



10-13

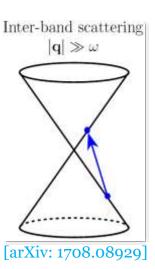
Super-Light DM: Main Focus

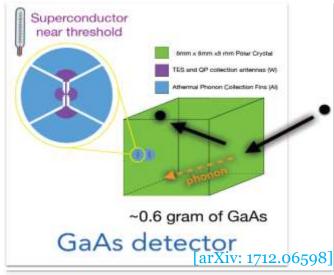


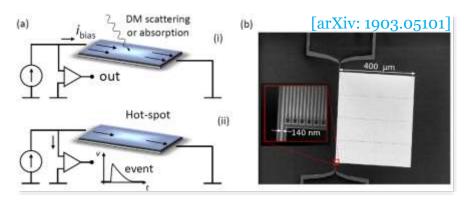
- **♦** $E_k \sim mv^2 < O(eV)$ Superlight DM
- \rightarrow Very low E_r^{th} required!
- New ideas for <u>very low E_r^{th} w/ e-recoil</u>:
 - ✓ Superconductor target w/ TES or MKID

 [arXiv:1504.07237, 1512.04533]
 - ✓ Superfluid He w/ TES or MKID
 [arXiv:1604.08206, 1611.06228]
 - ✓ 3D Dirac materials [arXiv:1708.08929]
 - ✓ Polar materials w/ TES or MKID

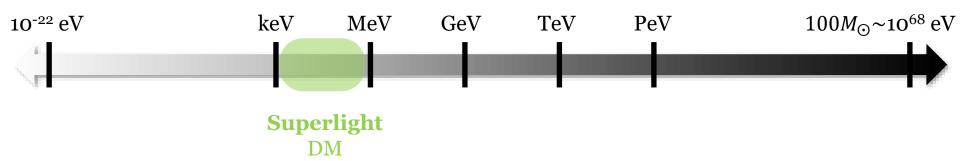
 [arXiv:1712.06598, 1807.10291]
 - ✓ Superconducting-nanowire single-photon detector [arXiv:1903.05101]
 - **√** ...







Super-Light DM: Technologies



- ❖ Transition edge sensor (TES): X-ray ~ near-IR, E_{th} ~ sub-eV [Superconducting Devices in Quantum Optics (2016)]
- ❖ Microwave kinetic inductance device (MKID): X-ray ~ far-IR, E_{th} ~0(10 meV)

 [Annual Review of Condensed Matter Physics (2012)]
- ❖ Superconducting-nanowire single-photon detector (SNSPD): UV ~ mid-IR, E_{th} ~0(100 meV) [Techno. (2018)]

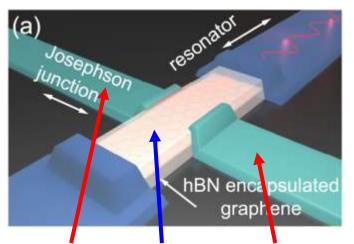
Well-developed in the laboratory in their respective E-bands.

But for the sensitivity to $E_{th} \leq meV$, further R&D is needed!

We proposed a new super-light DM direct detection strategy adopting the graphene-based Josephson junction* (GJJ) microwave single photon detector.

* A "state-of-the-art" technology: much lower $E_{th} \sim O(0.1 \text{ meV})$

Graphene Josephson Junction Device



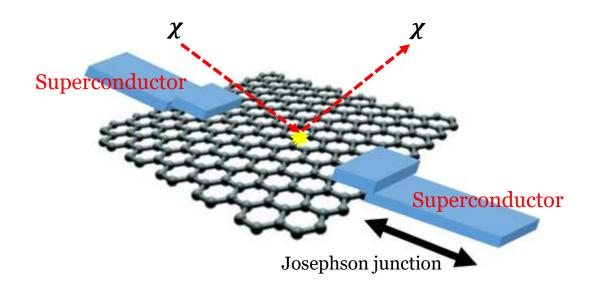
The device consists of a sheet of mono-layer graphene two sides of which are joined to superconductor, forming a superconductor-normal metal-superconductor Josephson junction.

Superconductor-Graphene-Superconductor (SGS)

- ❖ A GJJ single-photon detector was proposed, covering from near-IR to microwave.

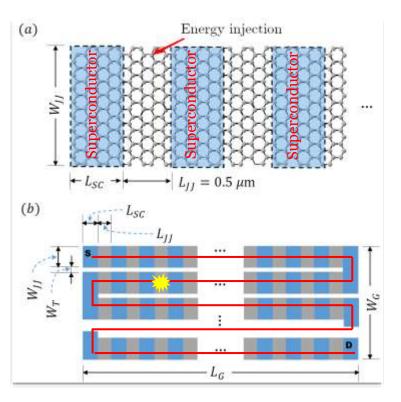
 [Phys. Rev. Applied (2017)]
- ❖ K.C. Fong, G.-H. Lee & their collaborators have demonstrated experimentally that the GJJ microwave bolometer can have sensitivity to E~0.1 meV energy deposit. [Nature (2020)]
- ❖ Currently, a GJJ single-photon detector is under testing in the laboratory.

Detection Principle



- I. DM scatters off (π -bond) free electrons, transferring some fraction of its incoming E_k .
- II. The recoiling e heats up & thermalizes with nearby e's rapidly via e-e interactions.
- III. The JJ is triggered: the temperature rise switches the zero-voltage of JJ to resistive state.
 - $E_k \sim mv^2 \sim 1 \text{ meV for } m_{DM} = 1 \text{ keV}$
 - → The GJJ device can posses the sensitivity to the signal induced even by sub-keV DM.

Conceptual Design Proposal



- I. Single graphene strip (a): the assembly of a graphene strip & a number of superconducting material strips → an array of SC-graphene-SC-graphene-SC-w (SGSGS···).
- II. Each sequence of SGS represents a single GJJ device.
- III. Full detector unit (b): all GJJs are connected in series so that even a single switched GJJ allows the series resistance measured between S & D to switch from 0 to a finite value.
- * E_{th} is determined by the strip width W_{JJ} : $W_{JJ} = 3 \, \mu \text{m} \, (30 \, \mu \text{m}) \rightarrow E_{th} \approx 0.1 \, \text{meV} \, (1 \, \text{meV})$.
- ❖ A large-scale detector can be made of a stack of such detector units.

To calculate experimental sensitivities, we should consider the scattering between DM traveling in 3D & free electrons living in 3D but confined in 2D graphene layer.

Calculating Signal Rates

- Goal: The event rate of DM scattering off free electrons in a 2-dimensional graphene sheet.
- ❖ Key point: An electron is still **confined** in the 2D graphene even after the collision.
 - \rightarrow No significant momentum change along the surface-normal (z-axis) direction.
- ❖ We will calculate the number of events/unit detector mass/unit run time:

$$n_{\text{eve}} = \frac{N_{\text{eve}}^{\text{total}}}{M_T t_{\text{run}}}$$

 $(N_{\text{eve}}^{\text{total}}: \text{total number of events}, M_T: \text{total detector mass}, t_{\text{run}}: \text{total time exposure})$

Calculation Procedure I

$$\star n_{\text{eve}} = \frac{N_{\text{eve}}^{\text{total}}}{M_T t_{\text{run}}} = \frac{1}{M_T t_{\text{run}}} \int_{E_r > E_{\text{th}}} dE_r \frac{dN_{\text{eve}}}{dE_r}$$

$$= \frac{1}{M_T t_{\text{run}}} \int_{E_r > E_{\text{th}}} \int dE_r \, dv_\chi \, f_{\text{MB}}(v_\chi) \frac{d}{dE_r} N_e \sigma_{e\chi} v_{\text{rel}} \frac{\rho_\chi}{m_\chi} t_{\text{run}}$$

$$= \int_{E_r > E_{\text{th}}} dE_r \, dv_\chi f_{\text{MB}}(v_\chi) \frac{dn_e^{3D} \sigma_{e\chi} v_{\text{rel}}}{dE_r} \frac{1}{\rho_T^{3D}} \frac{\rho_\chi}{m_\chi}$$

$$= \int_{E_r > E_{\text{th}}} dE_r \, dv_\chi f_{\text{MB}}(v_\chi) \frac{dn_e^{3D} \sigma_{e\chi} v_{\text{rel}}}{dE_r} \frac{1}{\rho_T^{3D}} \frac{\rho_\chi}{m_\chi}$$

$$= \int_{E_r > E_{\text{th}}} dE_r \, dv_\chi f_{\text{MB}}(v_\chi) \frac{dn_e^{3D} \sigma_{e\chi} v_{\text{rel}}}{dE_r} \frac{1}{\rho_T^{3D}} \frac{\rho_\chi}{m_\chi}$$

$$= \int_{E_r > E_{\text{th}}} dE_r \, dv_\chi f_{\text{MB}}(v_\chi) \frac{dn_e^{3D} \sigma_{e\chi} v_{\text{rel}}}{dE_r} \frac{1}{\rho_T^{3D}} \frac{\rho_\chi}{m_\chi}$$

$$= \int_{E_r > E_{\text{th}}} dE_r \, dv_\chi f_{\text{MB}}(v_\chi) \frac{dn_e^{3D} \sigma_{e\chi} v_{\text{rel}}}{dE_r} \frac{1}{\rho_T^{3D}} \frac{\rho_\chi}{m_\chi}$$

$$= \int_{E_r > E_{\text{th}}} dE_r \, dv_\chi f_{\text{MB}}(v_\chi) \frac{dn_e^{3D} \sigma_{e\chi} v_{\text{rel}}}{dE_r} \frac{1}{\rho_T^{3D}} \frac{\rho_\chi}{m_\chi}$$

$$= \int_{E_r > E_{\text{th}}} dE_r \, dv_\chi f_{\text{MB}}(v_\chi) \frac{dn_e^{3D} \sigma_{e\chi} v_{\text{rel}}}{dE_r} \frac{1}{\rho_T^{3D}} \frac{\rho_\chi}{m_\chi}$$

$$= \int_{E_r > E_{\text{th}}} dE_r \, dv_\chi f_{\text{MB}}(v_\chi) \frac{dn_e^{3D} \sigma_{e\chi} v_{\text{rel}}}{dE_r} \frac{1}{\rho_T^{3D}} \frac{\rho_\chi}{m_\chi}$$

$$= \int_{E_r > E_{\text{th}}} dE_r \, dv_\chi f_{\text{MB}}(v_\chi) \frac{dn_e^{3D} \sigma_{e\chi} v_{\text{rel}}}{dE_r} \frac{1}{\rho_T^{3D}} \frac{\rho_\chi}{m_\chi}$$

$$= \int_{E_r > E_{\text{th}}} dE_r \, dv_\chi f_{\text{MB}}(v_\chi) \frac{dn_e^{3D} \sigma_{e\chi} v_{\text{rel}}}{dE_r} \frac{1}{\rho_T^{3D}} \frac{\rho_\chi}{m_\chi}$$

2D nature of graphene
$$\frac{N_e}{M_T} = \frac{N_e/V}{M_T/V} = \frac{n_e^{3D}}{\rho_T^{3D}}$$
$$= \frac{N_e/(A\Delta l)}{M_T/(A\Delta l)} = \frac{n_e^{2D}}{\rho_T^{2D}}$$

graphene

 $= \int_{E_r > E_{\text{th}}} dE_r dv_{\chi \parallel} f_{\text{MB}}(v_{\chi \parallel}) \frac{dn_e^{2D} \sigma_{e\chi} v_{\text{rel} \parallel}}{dE_r} \frac{1}{\sigma_e^{2D}} \frac{\rho_{\chi}}{m}$

$$=2\int \frac{d^{3}p_{e,i}}{(2\pi)^{3}}(2\pi)\underline{\delta(p_{e,i}^{z}-p_{e,f}^{z})}f_{e,i}(E_{e,i})$$

✓
$$f_{e,i}(E_{e,i})=1/\{1+\exp(\frac{E_{e,i}-\mu}{T})\}, (\mu\sim E_F)$$

→ Fermi-Dirac distribution function

Consistent with the assumption of no significant

momentum change along the surface-normal direction

Calculation Procedure II

- * Graphene-surface-parallel DM velocity profile: $f_{\text{MB}}(v_{\chi\parallel}) = \frac{2(e^{-v_{\chi\parallel}^2/v_0^2} e^{-v_{\text{esc}}^2/v_0^2})}{\sqrt{\pi}v_0 \text{erf}(v_{\text{esc}}/v_0) 2v_{\text{esc}}e^{-v_{\text{esc}}^2/v_0^2}}$
 - → We take a plane-projection of a modified Maxwell-Boltzmann distribution.
- Event rate on a (sufficiently thin) 2D material: $\langle n_e^{\rm 2D} \sigma_{e\chi} v_{\rm rel} \rangle = \int \frac{d^3 p_{\chi,f}}{(2\pi)^3} \frac{|\mathcal{M}|^2}{16\pi m_e^2 m_\chi^2} S_{\rm 2D}(E_r, q)$
- * Structure function for the 2D system: $S_{2D}(E_r, q)$

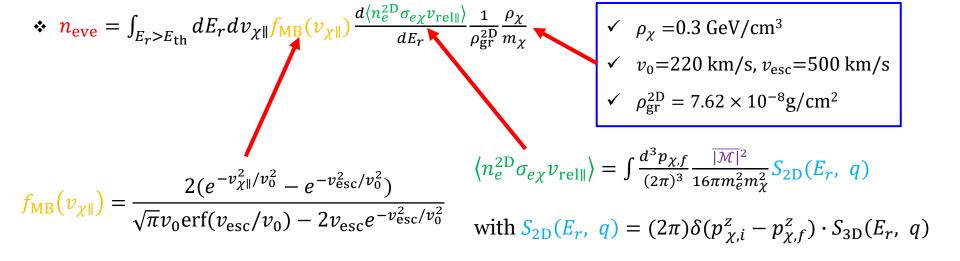
$$= 2 \int \frac{d^{3}p_{e,i}}{(2\pi)^{3}} \int \frac{d^{3}p_{e,f}}{(2\pi)^{3}} (2\pi) \delta(p_{e,i}^{z} - p_{e,f}^{z}) (2\pi)^{4} \delta^{(4)}(p_{\chi,i} + p_{e\,i} - p_{\chi,f} - p_{e,f}) f_{e,i}(E_{e,i}) \{1 - f_{e,f}(E_{e,f})\}$$

$$= (2\pi) \delta(p_{\chi,i}^{z} - p_{\chi,f}^{z}) \cdot \frac{1}{2\pi^{2}} \int d^{3}p_{e,i} \delta(E_{r} + E_{\chi,i} - E_{\chi,f}) f_{e,i}(E_{e,i}) \{1 - f_{e,f}(E_{e,f})\}$$

$$= (2\pi) \delta(p_{\chi,i}^{z} - p_{\chi,f}^{z}) \cdot S_{3D}(E_{r}, q)$$

→ The Pauli blocking effects(=phase space suppression) are encoded in the structure function. The analytic expression for $S_{3D}(E_r, q)$ is available in the non-relativistic limit [astro-ph/9710115, 1512.04533].

Calculation Procedure III



* We assume that DM interacts with electrons via an exchange of mediator ϕ as done in many of the preceding studies:

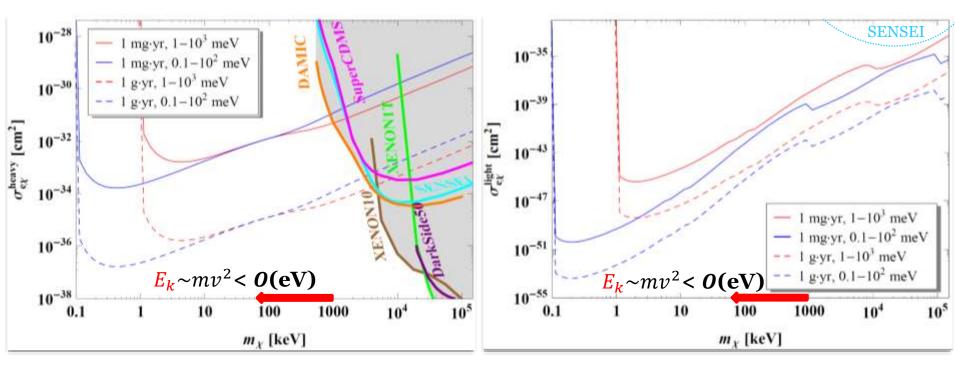
$$\sigma_{e\chi} \approx \frac{g_e^2 g_\chi^2}{\pi} \frac{\mu_{e\chi}^2}{(m_\phi^2 + q^2)^2} \implies \sigma_{e\chi}^{\text{heavy}} \approx \frac{g_e^2 g_\chi^2}{\pi} \frac{\mu_{e\chi}^2}{m_\phi^4} \text{ for } (m_\phi^2 \gg q^2) \ \& \ \sigma_{e\chi}^{\text{light}} \approx \frac{g_e^2 g_\chi^2}{\pi} \frac{\mu_{e\chi}^2}{q^4} \text{ for } (m_\phi^2 \ll q^2)$$

- * The matrix element $\overline{|\mathcal{M}|^2}$ is related to the scattering cross section as $\sigma_{e\chi} = \frac{\overline{|\mathcal{M}|^2}}{16\pi \, m_e^2 m_\chi^2} \mu_{e\chi}^2$.
- From the linear dispersion of graphene: $E_F = v_F \sqrt{\pi n_c}$ with $v_F \sim 10^8 \text{cm/s} \& n_c \sim 10^{12}/\text{cm}^2$.

Expected Sensitivities

Heavy mediator: $F_{DM} = 1$

Light mediator: $F_{DM} \propto 1/q^2$ with $q_{ref} = \alpha_e m_e$



- ✓ We required N_{eve} =3.6 under the negligible background assumption.
- ✓ The proposed GJJ DM detector can improve the minimum detectable DM mass ($m_{\rm DM}$ ~0.1 keV) by more than 3 orders of magnitude over the ongoing/existing experiments.
- ✓ Even capable of probing sub-keV DM with great expected reaches.

Future Plans & Conclusions

- > We have proposed a class of new DM detectors, adopting the GJJ device which has been implemented & demonstrated experimentally.
- > For the scattering between DM moving in 3D space & e's confined in 2D graphene, we (for the first time) built an effective model and computed the event rate.
 - → Signal rate depends on the DM flow direction!
- ➤ The proposed detector is capable of sensing sub-keV (warm) DM scattering off electrons due to its outstanding $E_{\rm th} \sim 0.1$ meV. → Improving the minimum detectable DM mass ($m_{\rm DM} \sim 0.1$ keV) by more than 3 orders of magnitude.

We are planning to do the First Experiment with the Existing GJJ Device samples.

