

MODULI STABILISATION AND THE HOLOGRAPHIC SWAMPLAND

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(based on JC, Quevedo 1811.06276,
JC, Revello 2006.01021)

MODULI STABILISATION

- String theory
 ↓ **EFT**
- 10-dimensional supergravity with α' corrections
 ↓ **EFT**
- 4-dimensional supergravity of moduli and matter
 ↓ **EFT**
- Integrate out heavy modes to get potential for lightest moduli
 ↓ **EFT**
- Find vacuum as minimum of effective potential

LARGE VOLUME SCENARIO

Balasubramanian, Berglund, JC, Quevedo

- Perturbative corrections to K and non-perturbative corrections to W

$$W = \int G_3 \wedge \Omega + \sum_i A_i e^{-2\pi a_i T_i}$$

$$K = -2 \ln(\mathcal{V} + \xi') + \ln\left(\int \Omega \wedge \bar{\Omega}\right) - \ln(S + \bar{S})$$

- Resulting scalar potential has minimum at *exponentially large* values of the volume

$$V = \frac{A\sqrt{\tau_s} e^{-2a_s \tau_s}}{\mathcal{V}} - \frac{B\tau_s e^{-a_s \tau_s}}{\mathcal{V}^2} + \frac{C}{\mathcal{V}^3}$$

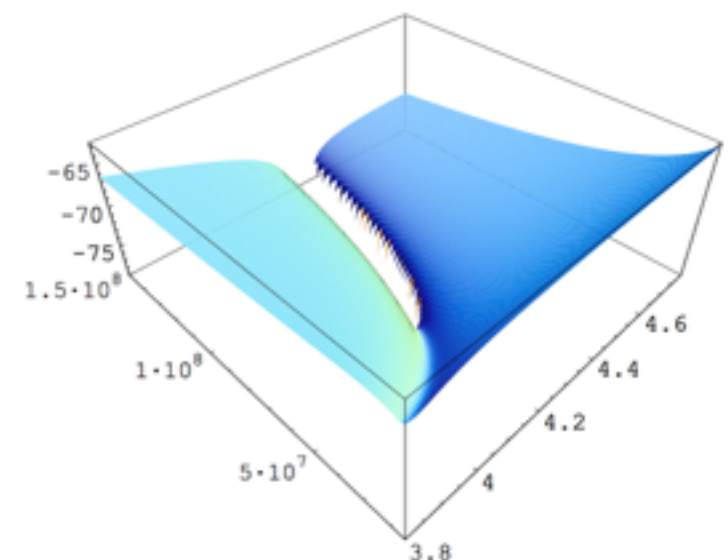


Figure 1: $\ln(V)$ for $P^4_{[1,1,1,6,9]}$ in the large volume limit, as a function of the divisors τ_4 and τ_3 . The void channel corresponds to the region where V becomes negative and $\ln(V)$ undefined. As $V \rightarrow 0$ at infinite volume, this immediately

WHY LVS?

- In LVS, volume is exponentially large - can easily be

$$\mathcal{V} \sim 10^{50} (2\pi\sqrt{\alpha'})^6$$

- This *generates interesting hierarchies* and ensures *superb parametric decoupling* of heavy modes (KK modes, heavy moduli)
- Decoupling also has a clear geometric origin - large volume
 $\langle \mathcal{V} \rangle \sim e^{\xi/g_s}$
- $\mathcal{V} \rightarrow \infty$ limit of LVS also leads to a unique effective theory

LVS HOLOGRAPHY

- LVS effective theory for volume modulus Φ and axion a

$$V_{potential} = V_0 e^{-\lambda \Phi / M_P} \left(- \left(\frac{\Phi}{M_P} \right)^{3/2} + A \right) \quad \left(\lambda = \sqrt{27/2} \right)$$
$$\mathcal{L}_{kinetic} = \frac{1}{2} \partial_\mu \Phi \partial_\mu \Phi + \frac{3}{4} e^{-\sqrt{\frac{8}{3}} \Phi} \partial_\mu a \partial^\mu a$$

- Other terms are subleading in infinite volume limit by $\mathcal{O}\left(\frac{1}{\ln \mathcal{V}}\right)$

LVS HOLOGRAPHY

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- Solve for minimum and expand about it to determine masses and couplings

HOLOGRAPHY

- CFT dimensions of dual operators:

$$\Delta(\Delta - 3) = m_{\Phi}^2 R_{AdS}^2$$

- In infinite volume limit can classify modes as

heavy	$m_{\Phi}^2 \gg R_{AdS}^{-2}, \Delta \rightarrow \infty \quad \text{as } \mathcal{V} \rightarrow \infty$
light	$m_{\Phi}^2 \ll R_{AdS}^{-2}, \Delta \rightarrow 3 \quad \text{as } \mathcal{V} \rightarrow \infty$
interesting	$m_{\Phi}^2 \sim R_{AdS}^{-2}, \Delta \rightarrow \mathcal{O}(1-10) \quad \text{as } \mathcal{V} \rightarrow \infty$

LVS MASS SPECTRUM

- In LVS we have
- **Heavy:** KK modes, complex structure moduli, all Kahler moduli except overall volume
- **Light:** Graviton, overall volume axion
- **Interesting:** overall volume modulus

LVS HOLOGRAPHY

Mode	Spin	Parity	Conformal dimension
$T_{\mu\nu}$	2	+	3
a	0	-	3
Φ	0	+	$8.038 = \frac{3}{2}(1 + \sqrt{19})$

Table 1. The low-lying single-trace operator dimensions for CFT duals of the Large Volume Scenario in the limit $\mathcal{V} \rightarrow \infty$.

In minimal LVS, AdS effective theory has small number of fields which correspond to specific predictions for dual conformal dimensions

No Landscape! (not true of KKLT)

LVS HOLOGRAPHY

- LVS is attractive as it offers a well-motivated Generalised Free Field Theory
- Large volume limit $\mathcal{V} \rightarrow \infty$ gives a unique theory
- Two scalars with fixed and radiatively stable anomalous dimensions
- All AdS interactions are also fixed and radiatively stable

LVS AND THE SWAMPLAND

- LVS effective Lagrangian is

$$V_{potential} = V_0 e^{-\lambda \Phi / M_P} \left(- \left(\frac{\Phi}{M_P} \right)^{3/2} + A \right) \quad \left(\lambda = \sqrt{27/2} \right)$$
$$\mathcal{L}_{kinetic} = \frac{1}{2} \partial_\mu \Phi \partial_\mu \Phi + \frac{3}{4} e^{-\sqrt{\frac{8}{3}} \Phi / M_P} \partial_\mu a \partial^\mu a$$

- This has the expected behaviour that

$$f_a / M_P \rightarrow 0 \quad \text{as} \quad \mathcal{V} \rightarrow \infty$$

LVS AND THE SWAMPLAND

- Now consider this small modification:

$$V_{potential} = V_0 e^{-\lambda \Phi / M_P} \left(- \left(\frac{\Phi}{M_P} \right)^{3/2} + A \right) \quad \left(\lambda = \sqrt{27/2} \right)$$
$$\mathcal{L}_{kinetic} = \frac{1}{2} \partial_\mu \Phi \partial_\mu \Phi + \frac{3}{4} e^{+\sqrt{\frac{8}{3}} \Phi / M_P} \partial_\mu a \partial^\mu a$$

- This coupling is equivalent to axion decay constants that *diverge* in the decompactification limit - must be in the swampland!

$$f_a / M_P \rightarrow \infty \quad \text{as} \quad \mathcal{V} \rightarrow \infty$$

HOLOGRAPHIC SWAMPLAND

- n-point self interactions of volume modulus

$$\mathcal{L}_{n-pt} = (-1)^{n-1} \lambda^n (n-1) \left(-3 \frac{M_P^2}{R_{AdS}^2} \right) \frac{1}{n!} \left(\frac{\delta \Phi}{M_P} \right)^n \left(1 + \mathcal{O} \left(\frac{1}{\ln \mathcal{V}} \right) \right) \quad \left(\lambda = \sqrt{27/2} \right)$$

- Mixed interactions of volume modulus and axion

$$\mathcal{L}_{\Phi^n aa} = \left(-\sqrt{\frac{8}{3}} \right)^n \frac{1}{2n!} \left(\frac{\delta \Phi}{M_P} \right)^n \partial_\mu a \partial^\mu a$$

- The higher-point interaction define 3- and higher point-correlators within a dual CFT

HOLOGRAPHIC SWAMPLAND

- n-point self interactions of volume modulus

$$\mathcal{L}_{n-pt} = (-1)^{n-1} \lambda^n (n-1) \left(-3 \frac{M_P^2}{R_{AdS}^2} \right) \frac{1}{n!} \left(\frac{\delta \Phi}{M_P} \right)^n \left(1 + \mathcal{O} \left(\frac{1}{\ln \mathcal{V}} \right) \right)$$

- Now **modify** interactions of volume modulus and axion $\left(\lambda = \sqrt{27/2} \right)$

$$\mathcal{L}_{\Phi^n aa} = \left(+\sqrt{\frac{8}{3}} \right)^n \frac{1}{2n!} \left(\frac{\delta \Phi}{M_P} \right)^n \partial_\mu a \partial^\mu a$$

- This defines a **perturbation** to the Generalised Free Field CFT with axion decay constants that *diverge* in the decompactification limit - must be in the swampland!

$$f_a / M_P \rightarrow \infty \quad \text{as} \quad \mathcal{V} \rightarrow \infty$$

HOLOGRAPHIC SWAMPLAND

- The problem:

1. Generalised Free Field + (some corrections)

- **consistent theory**

2. Generalised Free Field + (other corrections)

- **swampland!**

Where does the difference lie? Can one correlate any properties of the CFT with this change from the consistent theory to the swampland theory?

HOLOGRAPHIC SWAMPLAND

- 3-pt interactions in AdS theory relate to 3-pt structure functions in CFT
- Signs of 3-pt functions are not determinate - can change by field redefinitions
- Focus on 3-pt interactions only

HOLOGRAPHIC SWAMPLAND

$$\mathcal{L} \supset \frac{g}{M_P R_{AdS}^2} \varphi^3 - \frac{\mu}{M_P} \varphi \partial_\mu a \partial^\mu a$$

- gives rise to CFT structure constants

$$f_{\varphi\varphi}^\varphi = \frac{3g\pi^{d/2}\Gamma(\frac{3\Delta_\varphi-d}{2})}{\Gamma(\Delta_\varphi)^3}.$$

$$f_{aa}^\varphi = \frac{\mu\pi^{d/2}\Gamma(\frac{2\Delta_a+\Delta_\varphi-d}{2})}{\Gamma(\Delta_\varphi)\Gamma(\Delta_a)^2}(\Delta_\varphi - 6),$$

HOLOGRAPHIC SWAMPLAND

- Anomalous dimensions are well-defined; can be related to Mellin amplitude for $2 \rightarrow 2$ scattering

$$\gamma(0, \ell) = - \int_{-i\infty}^{+i\infty} \frac{ds}{2\pi i} M(s, 0) {}_3F_2\left(-\ell, \Delta_1 + \Delta_3 + \ell - 1, \frac{s}{2}; \Delta_1, \Delta_3; 1\right) \\ \times \Gamma\left(\Delta_1 - \frac{s}{2}\right) \Gamma\left(\Delta_3 - \frac{s}{2}\right) \Gamma\left(\frac{s}{2}\right)^2,$$

where $M(s, t)$ is the Mellin amplitude corresponding to the correlator

$$\langle \mathcal{O}_1(x_1) \mathcal{O}_1(x_2) \mathcal{O}_3(x_3) \mathcal{O}_3(x_4) \rangle.$$

HOLOGRAPHIC SWAMPLAND

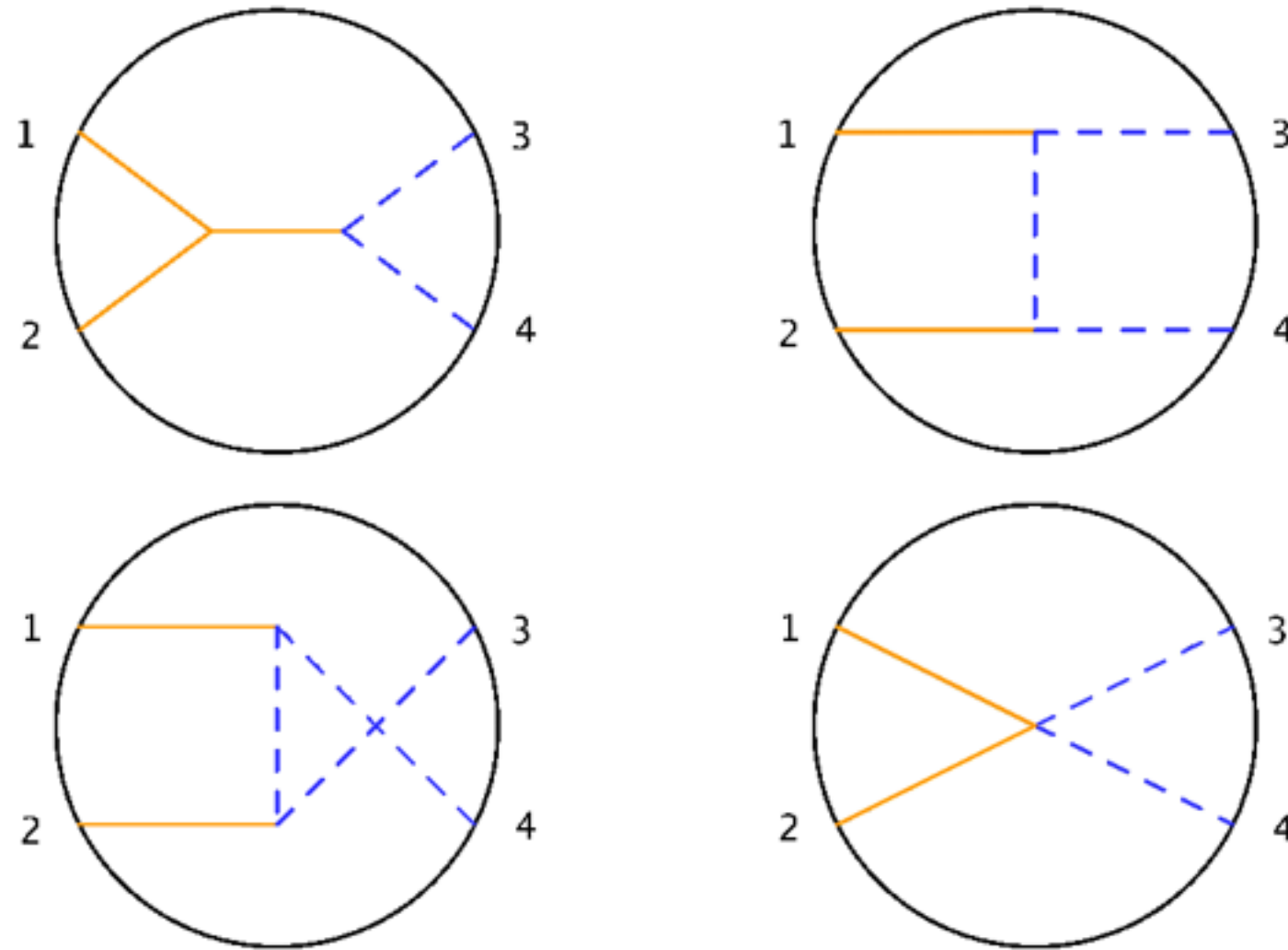


Figure 1: Witten diagrams contributing to the anomalous dimensions of double trace operators of the form $[\mathcal{O}_{\varphi_1} \mathcal{O}_{\varphi_2}]_{(n,\ell)}$. The volume modulus corresponds to the continuous orange lines, and the axion to the blue dashed ones. The last diagram, without internal propagators, only contributes to anomalous dimension for small ℓ .

HOLOGRAPHIC SWAMPLAND

$$\gamma^{\varphi a}(0, \ell) = -2f_{\varphi\varphi\varphi}f_{\varphi aa}\frac{\Gamma(\Delta_a)\Gamma(\Delta_\varphi)^2}{\Gamma\left(\frac{2\Delta_a-\Delta_\varphi}{2}\right)\Gamma\left(\frac{\Delta_\varphi}{2}\right)^3}\frac{1}{\ell^{\Delta_\varphi}} + \mathcal{O}\left(\frac{1}{\ell}\right)$$

Anomalous dimensions are equivalent to binding energies of 2-particle states in AdS

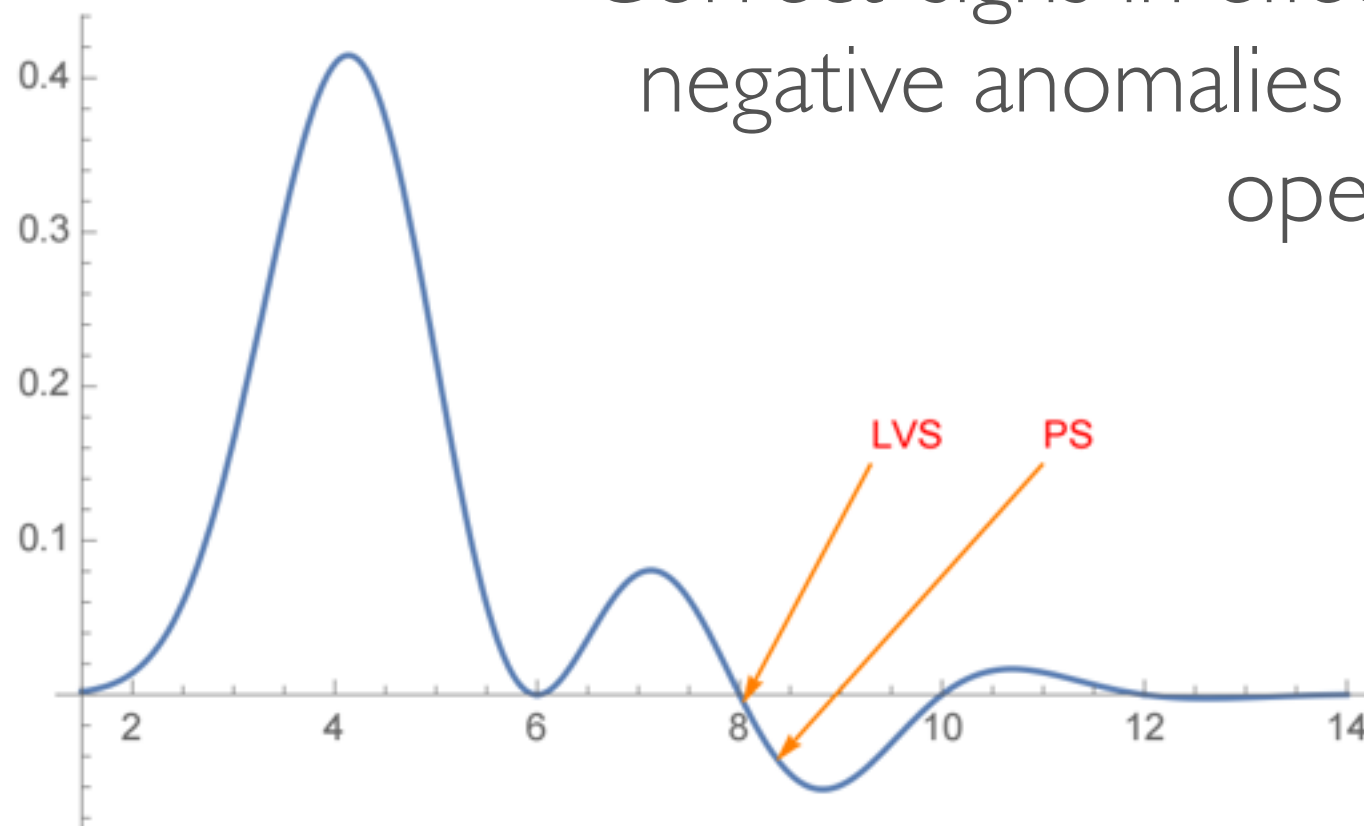
In Mellin amplitude, ‘exchange’ t-channel diagrams provide dominant contribution at large ℓ

HOLOGRAPHIC SWAMPLAND

LVS 'just' gives a negative anomalous dimension for the mixed volume-axion state

$$\gamma^{\varphi a}(0, \ell) \sim -g\mu \frac{(\Delta_\varphi - 6)}{\Gamma(\frac{6-\Delta_\varphi}{2})}.$$

'Correct' signs in effective AdS equivalent to negative anomalies dimension for mixed operator



HOLOGRAPHIC SWAMPLAND

- In LVS context, right signs of 3-pt AdS couplings are equivalent to negative anomalous dimensions for the mixed double-trace operator.
- A similar result holds for perturbative or KKLT stabilisation (qualitatively different as involves a massive axion)

CONNECTION TO REFINED DISTANCE CONJECTURE

$$\frac{1}{2}e^{-\sqrt{\frac{1}{6}}(\varphi-\varphi_0)/M_P}\partial_\mu\psi\partial^\mu\psi - m_\psi^2e^{-\frac{5}{\sqrt{6}}(\varphi-\varphi_0)/M_P}\frac{\psi^2}{2}.$$

KK modes have to couple to light volume modulus in a way that their mass decreases with increasing volume

This fixes the sign of the 3-pt function, again in a way that results in a negative anomalous dimension for the mixed double trace operator

CONCLUSIONS

- For many examples, negative CFT anomalous dimensions appear to correspond to the correct signs in the AdS Lagrangian
- However:
 1. Negativity of anomalous dimensions does not seem to hold for fibred LVS with extra light fibre moduli
 2. Axions couple with different signs to the different fibre moduli, resulting in a mixture of signs
- Are earlier results just a feature of the volume modulus? In progress.....