De Sitter & String Theory

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Some of my interests...

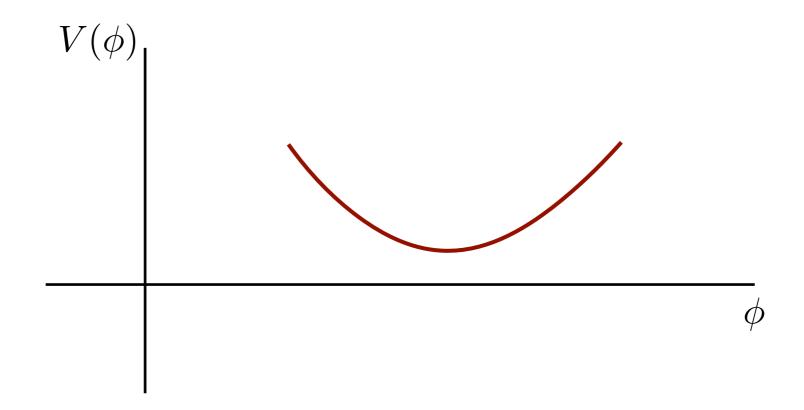
- Axions (in string theory)
 - Lectures June '20
- Euclidean wormholes
 - Journal club Nov '20
- Model building in string theory
 - ▶ IBS-IFT workshop Oct '20
 - Today

Outline

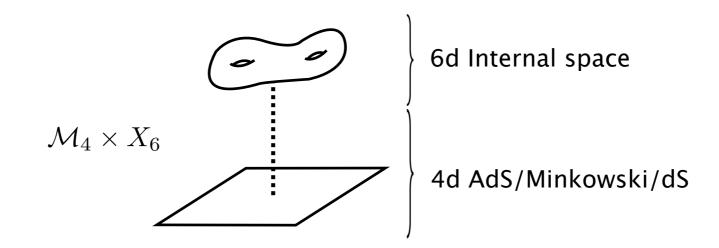
• Why is dS so hard to find?

The KKLT proposal

Why is dS so hard to find?



• Consider a compactification of string theory: $10d \rightarrow 4d$



No free parameters: coupling constants in 4d are vevs of scalar fields (moduli), e.g.

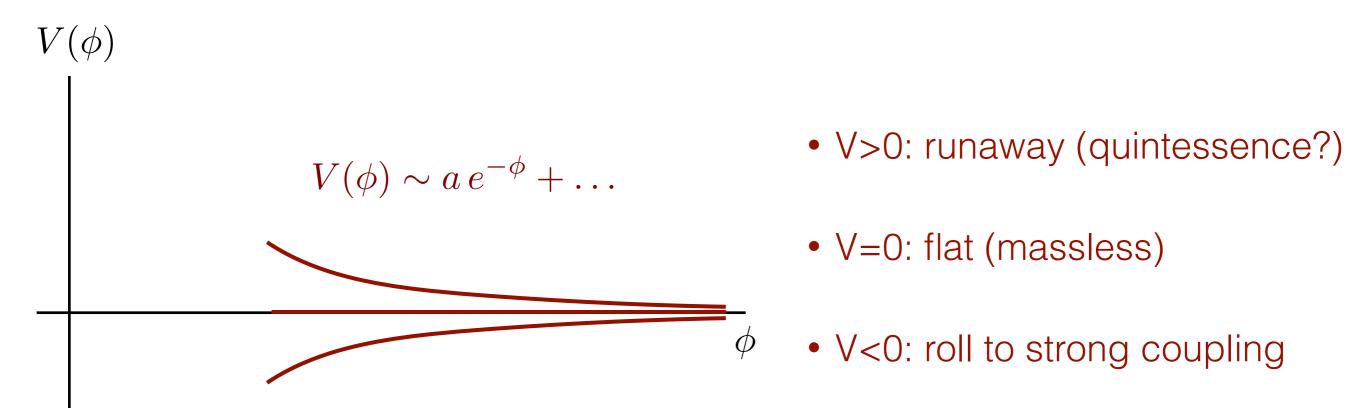
$$g_s = \langle e^{-\phi} \rangle$$
, $\operatorname{Vol}_{X_6} = \langle \phi \rangle$

Most computations rely on perturbation theory $\langle \phi \rangle \gg 1$

$$V(\phi) = a_0 + \frac{a_1}{\phi} + \frac{a_2}{\phi^2} + \dots$$
 or $V(\phi) = b_0 + b_1 e^{-\phi} + b_2 e^{-2\phi} + \dots$

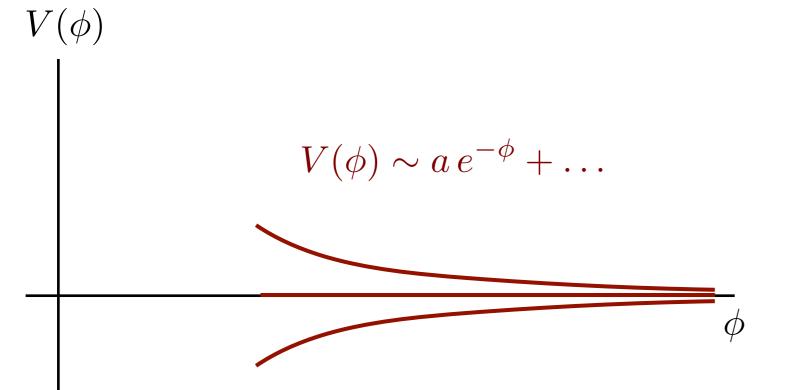
- It is notoriously difficult to obtain string dS vacua.
- The difficulty can be traced back to the Dine-Seiberg problem:

At weak coupling $(\phi \to \infty)$, vacuum energy vanishes



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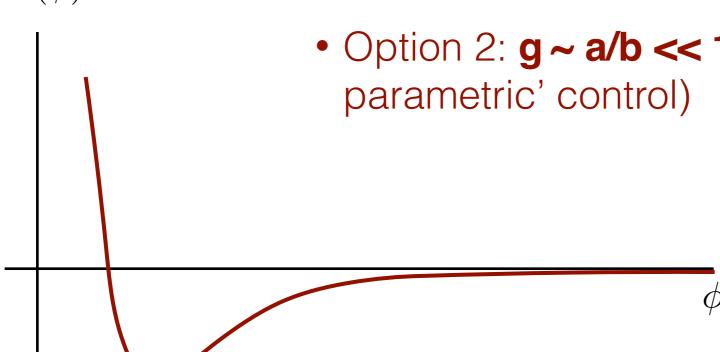
No stable vacuum exists at parametrically weak coupling!

 To find a minimum, one needs higher order corrections in the potential, but then perturbativity is endangered

$$V(\phi) = -a e^{-\phi} + b e^{-2\phi} + \dots \qquad \Longrightarrow \qquad g = e^{-\langle \phi \rangle} = \frac{a}{2b}$$



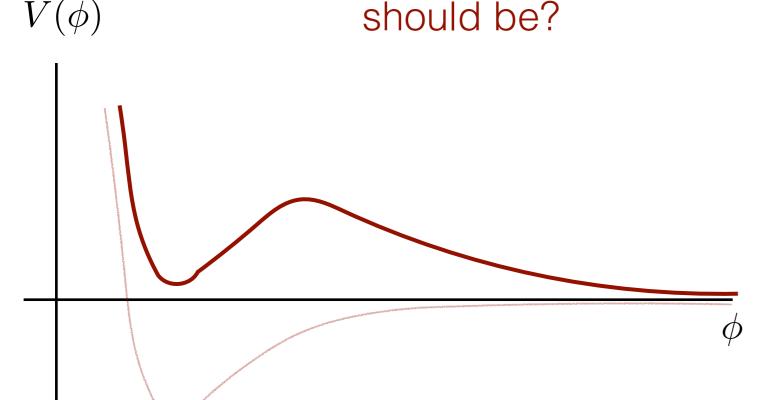
Option 2: g ~ a/b << 1 AdS at small coupling ('non-parametric' control)



 To find a minimum, one needs higher order corrections in the potential, but then perturbativity is endangered

$$V(\phi) = a e^{-\phi} - b e^{-2\phi} + c e^{-3\phi} + \dots$$

 With one more term, one can obtain potentials with dS minima at g << 1. Not ideal, but who said it should be?



Moduli stabilization and constructions of vacua (AdS, Mink or dS) exploit this mechanism

Recent suggestion: the potential must satisfy asymptotically

$$|\nabla V(\phi)| \ge \alpha V$$

$$\alpha \sim \mathcal{O}(1), \, \phi \to \infty$$

Obied, Ooguri, Spodyneiko, Vafa '18

This behaviour arises naturally in string theory, and is required asymptotically by swampland conjectures.

Ooguri, Palti, Shiu, Vafa '18

• **De Sitter swampland conjecture:** this must hold (with minor qualifications) throughout moduli space, forbidding dS vacua.

dS vacua (KKLT, LVS)

VS.

dS swampland conjecture

Necessarily complicated Vacuum energy

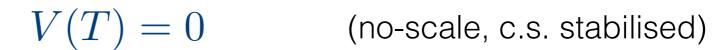
Simple but speculative Quintessence

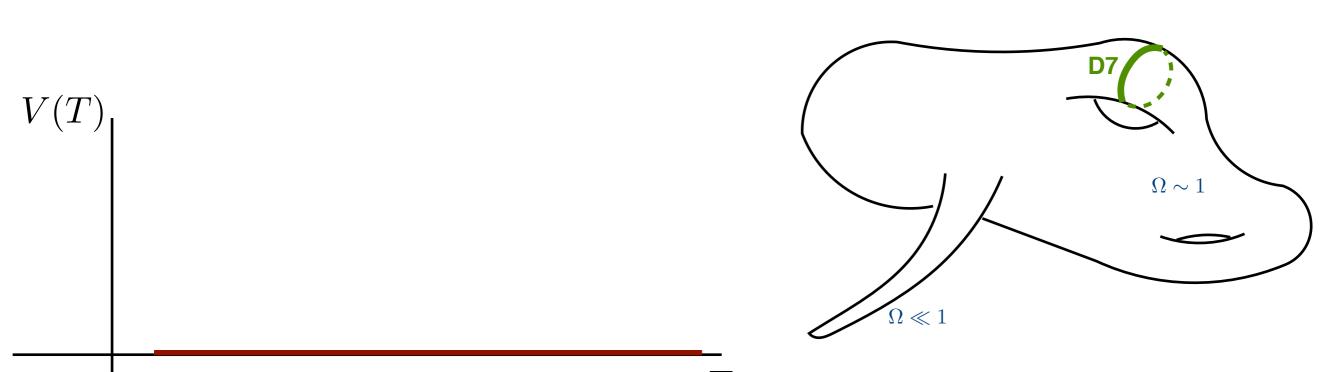
The KKLT proposal

Kachru, Kallosh, Linde, Trivedi '03

Type IIB compactification with one Kahler (volume) modulus T.

Step 0: warped compactification to Minkowski with O7-planes, D3/D7-branes and fluxes





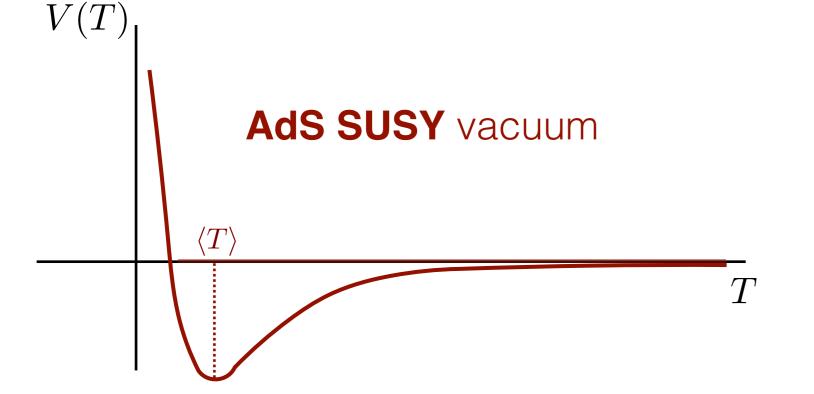
$$T ds_{10}^2 = \Omega^2(y) \eta_{\mu\nu} dx^{\mu} dx^{\nu} + g_{mn} dy^m dy^n$$

Giddings, Kachru, Polchinski '01

Type IIB compactification with one Kahler (volume) modulus T.

Step 1: take into account non-perturbative effects, e.g. gaugino condensation on D7-branes

$$\langle \lambda \lambda \rangle \sim e^{-T} \implies V(T) \sim \frac{1}{T} e^{-2T} - \frac{2}{T^2} W_0 e^{-T}$$



W₀: flux superpotential

$$W_0 \sim e^{-\langle T \rangle}$$

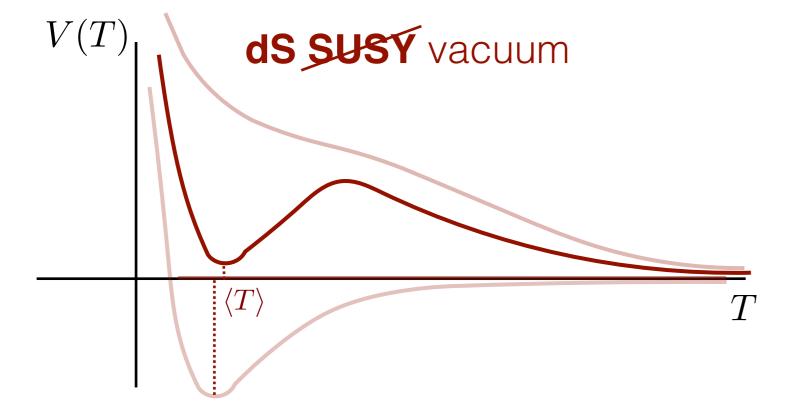
Computational control:

$$\langle T \rangle > 1 \implies W_0 \ll 1$$

Type IIB compactification with one Kahler (volume) modulus T.

Step 2: introduce anti-D3-brane at the tip of the throat

$$V(T) \sim \frac{1}{T}e^{-2T} - \frac{2}{T^2}W_0e^{-T} + \frac{\mu_3}{T^2}$$



μ₃ <<1: warped D3-tensionIf μ₃ not small enough⇒ runaway

 The KKLT proposal combines string theory ingredients (fluxes, non-perturbative effects & anti-branes) in a clever way

Individually, each ingredient is relatively well understood

It is their combination in a single setup that is poorly controlled

 Despite thorough scrutiny since proposed, it has resisted strong criticism rather well.

On the other hand, no explicit construction obtained so far

 If (strong) no-dS conjecture holds they should be pathological Renewed interest and attacks on KKLT and LVS.