IBS-ICTP Workshop on Axion-Like Particles 29 October 2021

Axion-like particles from primordial black holes shining through the Universe

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Based on arXiv:2107.03420 with D. Montanino, A. Mirizzi and F. Capozzi

- 1 ALPs, primordial black holes and Hawking spectra
- 2 ALP-photon conversions
- 3 Astrophysical and cosmological signatures
- Massive ALP decay
- Conclusions

Axion-like particles

- Axions first proposed as a solution to the strong-CP problem via the Peccei-Quinn mechanism¹
- Axion-like (neutral pseudoscalar) particles postulated in several BSM theories
- Various thermal and non-thermal production mechanisms for a cosmic axion background²

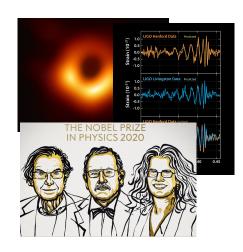


¹Peccei and Quinn 1977

²Dror et al. 2021 [2101.09287]

Primordial black holes

- Density fluctuations in the Early Universe collapsed below their Schwarzschild radius³
- "Hot topic" in modern cosmology
- Invoked to explain SMBHs, LIGO/VIRGO data, dark matter...



³Zel'dovich and Novikov 1967; Hawking 1971; Carr and Hawking 1974

PBH formation and evolution

PBH mass at formation

$$M \sim m_P^2 t_f \sim 10^{15} \left(\frac{t_f}{10^{-23} \,\mathrm{s}} \right) \mathrm{g}$$

- PBHs behave as matter: $ho_{\rm BH}/
 ho_r \sim a$
- Early matter domination is possible
- Broad mass range: between inflation and BBN⁴ $10\,\mathrm{g} \lesssim M \lesssim 10^9\,\mathrm{g}$

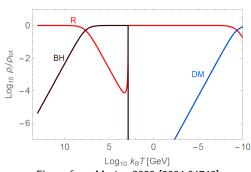


Figure from Masina 2020 [2004.04740]

⁴Papanikolaou et al. 2020 [2010.11573]

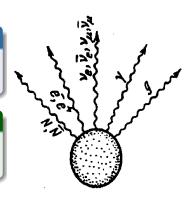
Hawking radiation

Hawking emission rate (Q = J = 0)

$$\frac{\mathrm{d}N_s}{\mathrm{d}\omega\,\mathrm{d}t} = \frac{1}{2\pi}\frac{\Gamma_s(\omega,M)}{e^{\omega/T_{\mathrm{BH}}}-(-1)^{2s}} \ [\mathrm{Hawking} \ 1975]$$

Black-hole temperature

$$T_{\rm BH} = \frac{m_{\rm P}^2}{8\pi M} \simeq 10^7 \left(\frac{10^6 \,\rm g}{M}\right) \,\rm GeV$$



- ullet "Democratic" emission for $m < T_{
 m BH}$: all SM states plus one ALP
- Non-thermal ALP production mechanism

Graybody factors

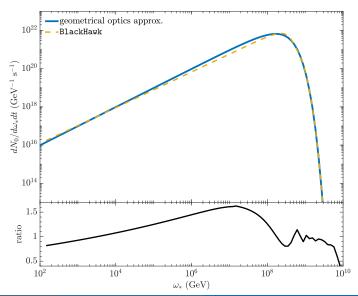
- Graybody factors $\Gamma_s(\omega,M)$: complicated functions, depending on the spin s of radiated particles
- Usually computed numerically (e.g. via BlackHawk⁵)
- ullet However, for high $T_{
 m BH}$ (i.e. high ω) "geometrical optics" holds

Geometrical optics approximation

$$\Gamma_s(\omega, M) \simeq \frac{27M^2\omega^2}{m_P^4}$$

⁵Arbey and Auffinger 2021 [1905.04268]

Geometrical optics approximation



Mass loss rate and lifetime

Mass loss rate

$$\frac{\mathrm{d}M}{\mathrm{d}t} = -\sum_{s} \int \mathrm{d}\omega \, \omega \frac{\mathrm{d}N_{s}}{\mathrm{d}\omega \mathrm{d}t} \equiv -\frac{m_{P}^{4} f_{\mathrm{ev}}}{M^{2}}$$

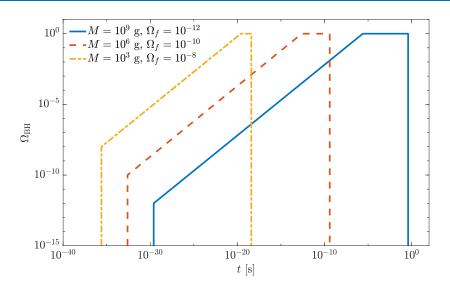
Evolution of PBH mass

$$M(t) = M \left(1 - \frac{t}{\tau_{\rm BH}} \right)^{1/3}$$

PBH lifetime

$$\tau_{\rm BH} = \frac{1}{3f_{\rm ev}} \frac{M^3}{m_P^4} \simeq 4.16 \times 10^{-1} \left(\frac{M}{10^9 \,\mathrm{g}}\right)^3 \,\mathrm{s}$$

Evolution of PBH density



Hawking ALP spectra

 ALP spectra obtained by integrating Hawking emission rate over PBH lifetime in an expanding matter-dominated Universe

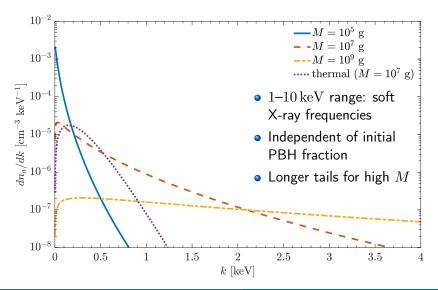
Spectral shape factor

$$\mathcal{I}(x) = x^2 \int_0^1 \frac{\theta^{-2} (1 - \theta)^{2/3}}{\exp\left[x\theta^{-2/3} (1 - \theta)^{1/3}\right] - 1} d\theta$$

- Mass spectrum of PBHs simply assumed monochromatic⁶
- Spectra calculated at PBH evaporation, then redshifted to the present epoch

⁶Other PBH mass spectra reviewed in Carr et al. 2020 [2002.12778]

Hawking ALP spectra



ALP extra radiation

- Ultrarelativistic ALPs contribute to extra or dark radiation
- Contribution measured as a deviation from the standard effective number of neutrinos $N_{
 m eff}=3.046$

ALP contribution to extra radiation

$$\Delta N_{\text{eff}} = \frac{8}{7} \left(\frac{11}{4}\right)^{4/3} \frac{\rho_a(T_0)}{\rho_{\text{CMB}}(T_0)} \simeq 0.042 \left(\frac{100}{g_S(T_*)}\right)^{1/3}$$

• Present experimental value⁷: $N_{\rm eff} = 2.99 \pm 0.17$

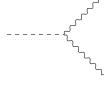
⁷Aghanim et al. 2020 [1807.06209]

ALP-photon coupling

Signatures of this model from the ALP-photon coupling $g_{a\gamma}$:

• Ultralight ALPs ($m_a \lesssim 10^{-9} \, \mathrm{eV}$): $a \to \gamma$ conversions (cosmological magnetic fields required, $B_0 \sim 1 \, \mathrm{nG}$)

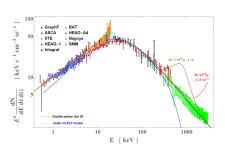
• "Heavy" ALPs ($m_a \gtrsim 10 \, \mathrm{eV}$): $a \to \gamma \gamma$ decays

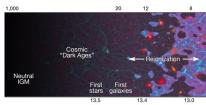


Astrophysical and cosmological signatures

These two processes inject high-energy photons in the Universe contributing to:

- The present-day cosmic X-ray background (CXB), directly measured by experiments
- ullet The history of reionization, measured by the Thomson optical depth au





Bounds on ALP-photon coupling

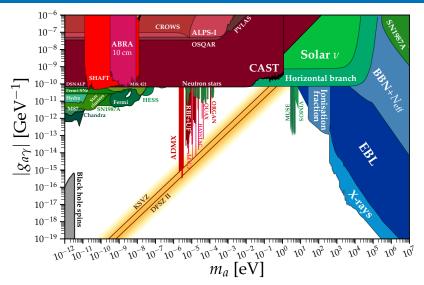


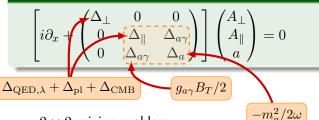
Figure from https://github.com/cajohare/AxionLimits

ALP-photon mixing in flat spacetime

ALP-photon Lagrangian

$$\mathcal{L} = \underbrace{\frac{1}{2} \big(\partial_{\mu} a \partial^{\mu} a + m_{a}^{2} a^{2} \big)}_{\text{ALP kinetic term}} \underbrace{-\frac{1}{4} F_{\mu\nu} F^{\mu\nu}}_{\text{e.m. term}} \underbrace{-\frac{g_{a\gamma}}{4} F_{\mu\nu} \tilde{F}^{\mu\nu} a}_{a\gamma \text{ interaction term}} + \begin{pmatrix} \text{higher orders} \\ \text{orders} \end{pmatrix}$$

E-L equations in a constant magnetic field⁸



Axion-like particles from PBHs

• 2×2 mixing problem

Mixing parameters

$$\Delta_{a\gamma} \simeq 1.52 \times 10^{-3} \left(\frac{g_{a\gamma}}{10^{-12} \,\text{GeV}^{-1}} \right) \left(\frac{B_T}{10^{-9} \,\text{G}} \right) \,\text{Mpc}^{-1}$$

$$\Delta_a \simeq -7.8 \times 10^5 \left(\frac{m_a}{10^{-10} \,\text{eV}} \right)^2 \left(\frac{\omega}{1 \,\text{keV}} \right)^{-1} \,\text{Mpc}^{-1}$$

$$\Delta_{\text{pl}} \simeq -1.1 \times 10^{-2} \left(\frac{\omega}{1 \,\text{keV}} \right)^{-1} \left(\frac{n_e}{10^{-7} \,\text{cm}^{-3}} \right) \,\text{Mpc}^{-1}$$

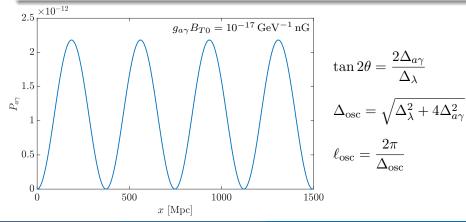
$$\Delta_{\text{QED}} \simeq 4.1 \times 10^{-18} \left(\frac{\omega}{1 \,\text{keV}} \right) \left(\frac{B_T}{10^{-9} \,\text{G}} \right)^2 \,\text{Mpc}^{-1}$$

$$\Delta_{\text{CMB}} \simeq 2.62 \times 10^4 \left(\frac{T}{1 \,\text{eV}} \right)^4 \left(\frac{\omega}{1 \,\text{keV}} \right) \,\text{Mpc}^{-1}$$

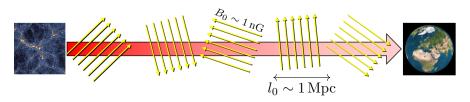
Mixing in a constant field

Oscillation probability

$$P(x) = \sin^2(2\theta)\sin^2\left(\frac{\Delta_{\rm osc}x}{2}\right)$$



ALP-photon conversions in primordial magnetic fields



- Expansion of the Universe
- Full 3-dimensional problem
- Photon absorption by the intergalactic medium

electron density

Photon absorption rate⁹

H/He density

$$\Gamma(\omega) = \sigma_{\rm H}^{\rm PE}(\omega) n_{\rm H} + \sigma_{\rm He}^{\rm PE}(\omega) n_{\rm He} + \sigma_{\rm KN}(\omega) n_e$$

Photoelectric effect on H/He

Compton scattering

⁹Evoli et al. 2016 [1602.08433]

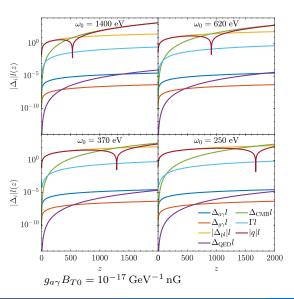
Redshift of parameters

- A sequence of redshift parameters is built as $z_{n+1} = z_n + \Delta z_n$
- Δz_n : how much time relativistic ALPs take to cross a domain of size $l_n = l_0/(1+z_n)$

Time-redshift relation

$$\frac{\mathrm{d}z}{\mathrm{d}t} = -(1+z)H_0\sqrt{\Omega_{\Lambda} + \Omega_m(1+z)^3}$$

Redshift of parameters



$$\Delta_a^{(n)} = \Delta_a^{(0)} (1 + z_n)^{-1}$$

$$\Delta_{a\gamma}^{(n)} = \Delta_{a\gamma}^{(0)} (1 + z_n)^2$$

$$\Delta_{\text{QED}}^{(n)} = \Delta_{\text{QED}}^{(0)} (1 + z_n)^5$$

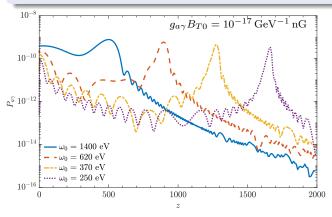
$$\Delta_{\text{pl}}^{(n)} = \Delta_{\text{pl}}^{(0)} (1 + z_n)^2$$

$$\Delta_{\text{CMB}}^{(n)} = \Delta_{\text{CMB}}^{(0)} (1 + z_n)^5$$

Conversion probability

Recursive probability formula [Evoli et al. 2016]

$$P_{a\gamma}^{(n+1)} = \left[P_{a\gamma}^{(n)} + (\Delta_{a\gamma}^{(n)} l_n)^2 \operatorname{sinc}^2 \left(\frac{q_n l_n}{2} \right) \right] e^{-\Gamma_n l_n}$$



Accounts for:

- Expansion
- Conversions
- Absorption

The cosmic X-ray background

- Since the 1960s, several space experiments have detected a diffuse electromagnetic background radiation
- Referred to as the cosmic X-ray background (CXB) in the 0.1–100 keV energy band
- Most sources in the low-energy part resolved (e.g. active galactic nuclei)
- Less complete picture at high energies



The Chandra X-Ray Observatory

The cosmic X-ray background

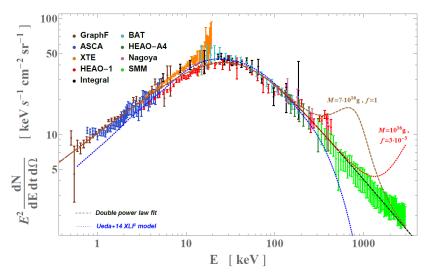


Figure from Ballesteros et al. 2020 [1906.10113]

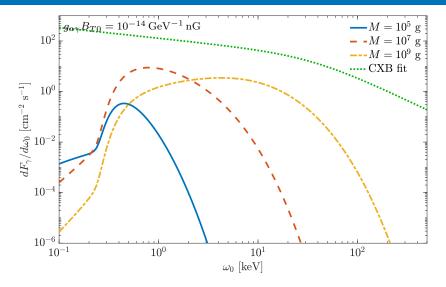
X-ray fluxes from ALP conversions

Present X-ray flux from conversions

$$\frac{\mathrm{d}F_{\gamma}}{\mathrm{d}\omega_0}(\omega_0, t_0) = \frac{\mathrm{d}F_a}{\mathrm{d}\omega_0}(\omega_0, t_0)P_{a\gamma}(\omega_0, t_0)$$

- Conversion probability evolved until present-day using the recursive formula
- Mixing parameter $g_{a\gamma}B_{T0}$ constrained by imposing that the obtained flux does not exceed the measured one

X-ray fluxes from ALP conversions



Soft X-ray excess from galaxy clusters

- Several galaxy clusters exhibit excess luminosity in the soft X-ray spectrum
- Possibly explained by a "cosmic axion background" 10 converting into photons inside the cluster magnetic fields ($\sim \mu G$)
- Similar results can be reproduced with PBHs

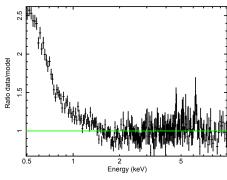
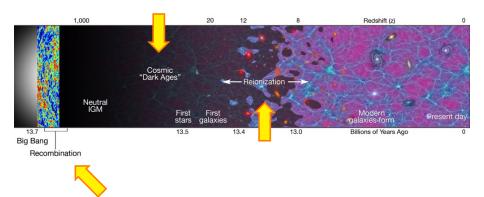


Figure from Petrucci et al. 2007 [0706.0134]

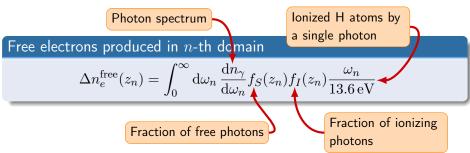
¹⁰Conlon and Marsh 2013 [1305.3603], Angus et al. 2014 [1312.3947]

Ionization history of the Universe



Reionization from ALP conversions

 ALPs convert into high-energy photons which may ionize neutral atoms after recombination



The ionization fraction

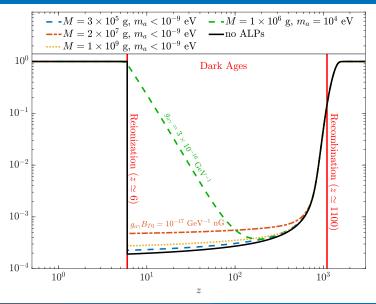
Ionization fraction

$$X_e(z_n) = \frac{n_e^{\text{free}}(z_n)}{n_H(z_n)} = X_e^0(z_n) + \sum_{i=1}^n \frac{\Delta n_e^{\text{free}}(z_i)}{n_H(z_i)}$$

- The standard evolution of the ionization fraction of the Universe $X_e^0(z)$ can be computed numerically 11
- It is modified by the injection of extra free electrons

¹¹https://www.cfa.harvard.edu/~sasselov/rec/

The ionization fraction



Optical depth

 The increased ionization contributes to the optical depth of the Universe

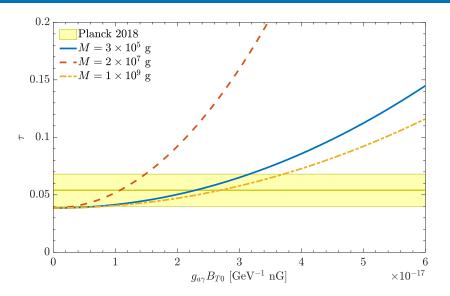
Thomson optical depth

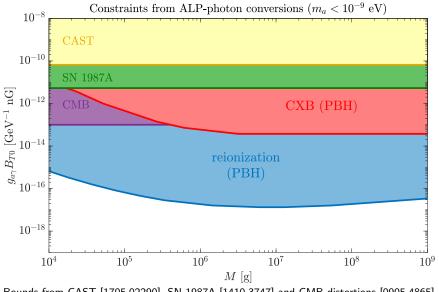
$$\tau = \int_0^\infty \mathrm{d}z \left| \frac{\mathrm{d}t}{\mathrm{d}z} \right| \sigma_T n_e^{\text{free}}(z)$$

• Impose that the obtained value of au does not exceed the measured one 12 to constrain $g_{a\gamma}B_{T0}$

¹²Aghanim et al. 2020 [1807.06209]

Optical depth





Bounds from CAST [1705.02290], SN 1987A [1410.3747] and CMB distortions [0905.4865] reported assuming $B_0=1\,\mathrm{nG}$

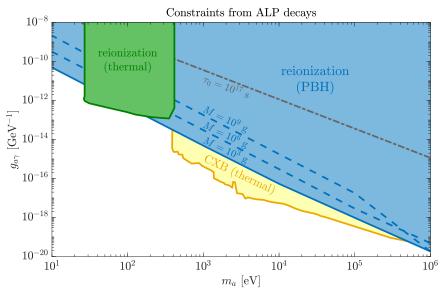
Decay of massive ALPs

$a ightarrow \gamma \gamma$ decay rate

$$\Gamma_{a\gamma} = \frac{m_a^3 g_{a\gamma}^2}{64\pi} \simeq 7.55 \times 10^{-40} \left(\frac{m_a}{1 \text{ eV}}\right)^3 \left(\frac{g_{a\gamma}}{10^{-17} \text{ GeV}^{-1}}\right)^2 \text{ s}^{-1}$$

- Decay is relevant for $\Gamma_{a\gamma}^{-1}$ < age of the Universe (including Lorentz boosts)
- Focus on "heavy" ALPs
- Photon spectrum obtained by solving a Boltzmann equation including absorption and production terms





Bounds for thermal production from Cadamuro and Redondo 2012 [1110.2895]

Conclusions

- PBH domination is a possible occurrence in the Early Universe
- Several observable signatures if PBHs emit axion-like particles
- Stringent constraints on ALP-photon mixing in this scenario
- Further developments: include gravitons (e.g. from spinning PBHs) and graviton-photon conversions in the picture

Selected references

More recent papers about PBHs and Hawking radiation:

- Carr et al. 2020 [2002.12778]
- Hooper et al. 2019 [1905.01301], Hooper et al. 2020 [2004.00618]
- Masina 2020 [2004.04740], Auffinger et al. 2021 [2012.09867], Masina 2021 [2103.13825]
- Arbey et al. 2021 [2104.04051]

About axion-like particles from PBHs:

Bernal et al. 2021 [2107.13575], [2110.04312]

