

Preparation for precise mass measurements using MMS at RAON

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RISP, Institute of Basic Science

Outline

- I. Introduction
- II. Proposed experiments
- III. MRTOF Simulator
- IV. Summary

I. Introduction

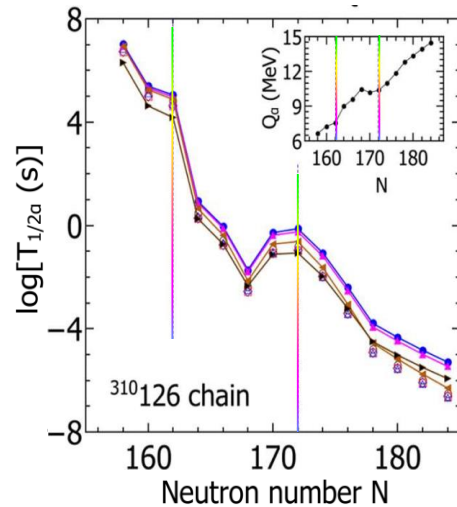
□ Nuclear structure

- Tracking shell closure

2n-separation energy: $S_{2n} = M(Z, N) - M(Z, N+2) - 2m_n$

Shell gap energy: $\delta_{2n} = S_{2n}(Z, N) - S_{2n}(Z, N+2)$

- α -cluster in SHN & α -decay



N.D. Ly et al., Phys.Scr. 96 (2021) 035301.

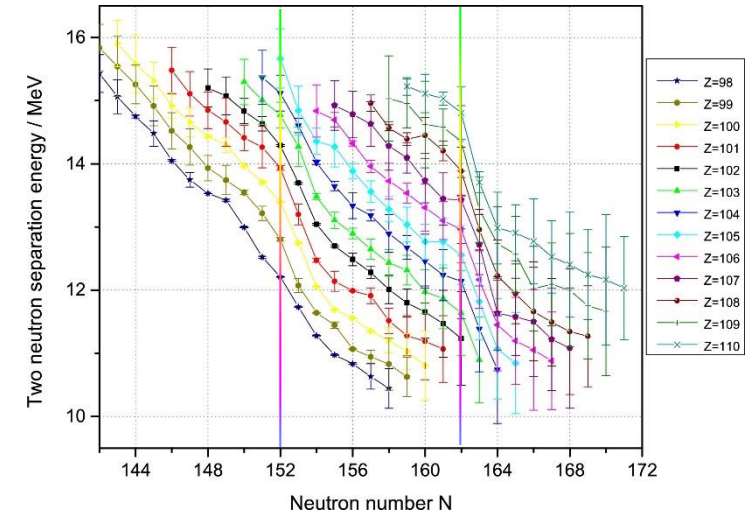
α -preformation: $S_\alpha = \frac{\ln 2}{\nu T_{1/2} P}$

$$P = \exp\left\{-\frac{2}{\hbar} \int_{r_a}^{r_b} \sqrt{2\mu(V_T(r) - Q_\alpha)} dr\right\}$$

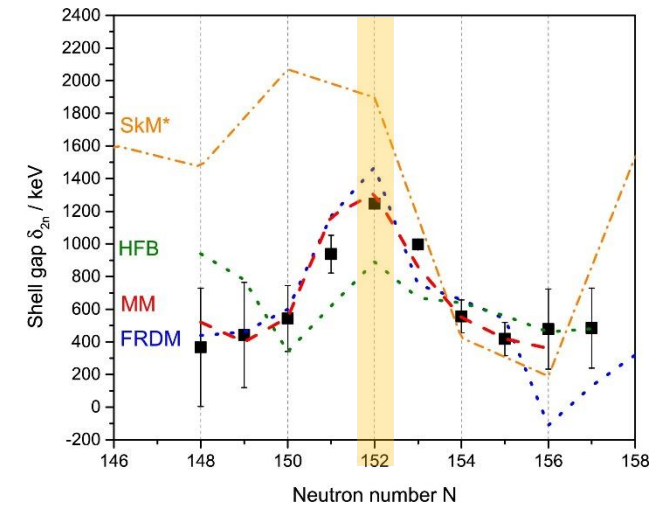
$$\nu = \omega/2\pi = 2E_\nu/h$$

$$E_\nu = Q_\alpha \left[0.056 + 0.039 \exp\left(\frac{4 - A_2}{2.5}\right) \right]$$

$$Q_\alpha = M(Z, N) - M(A-4, Z-2) - M_\alpha$$



M. Block, NPA 944 (2015) 471-491.
Q. Mo et al., PRC90, 024320 (2014)



M. Block, NPA 944 (2015) 471-491.

I. Introduction

□ Astrophysical implications

- Beta-decay rates

$$\ln T_{1/2} = a_6 + \left(\alpha^2 Z^2 - 5 - a_7 \frac{N - Z}{A} \right) \ln(Q - a_8 \delta) + a_9 \alpha^2 Z^2 + \frac{1}{3} \alpha^2 Z^2 \ln(A) - \alpha Z \pi + S(Z, N)$$

$$\begin{aligned} Q_{\beta^-} &= M(Z, N) - M(Z + 1, N - 1) \\ &= E_{\text{bind}}(Z + 1, N - 1) - E_{\text{bind}}(Z, N) + M_n - M_H \end{aligned}$$

[Zhou et al. (2017)]

- Capture rates

$$N_A \langle \sigma v \rangle = 1.54 \times 10^{11} (\mu T_9)^{-3/2} \sum_i (\omega \gamma)_i \exp\left(\frac{-11.605 E_i}{T_9}\right)$$

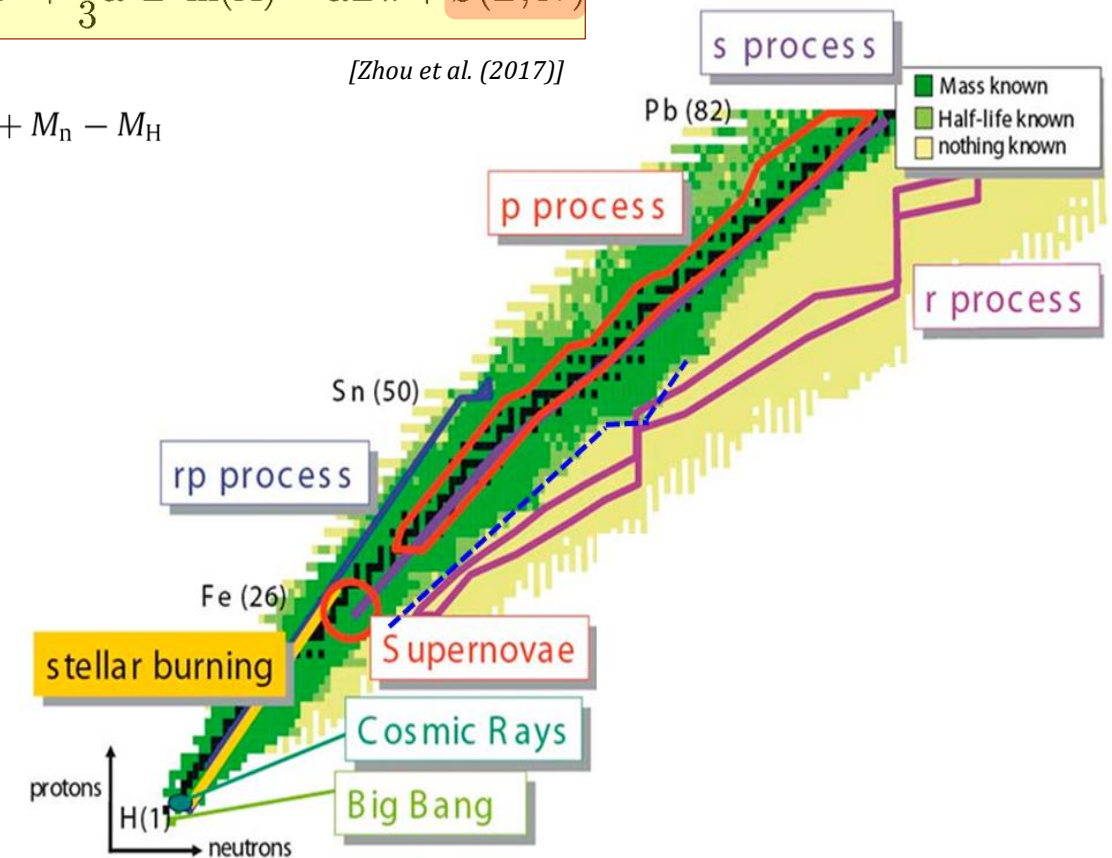
$$E_i = E_{\text{cm}} + Q \longrightarrow Q = (m_I - m_O) \times 931.5 \text{ MeV}$$

- Photodissociation rate in r-process

$$\lambda_\gamma(Z, A) \propto \exp\left[-\frac{S_n(Z, A)}{kT}\right] \langle \sigma v \rangle_{(Z, A-1)}$$

- Saha Equation

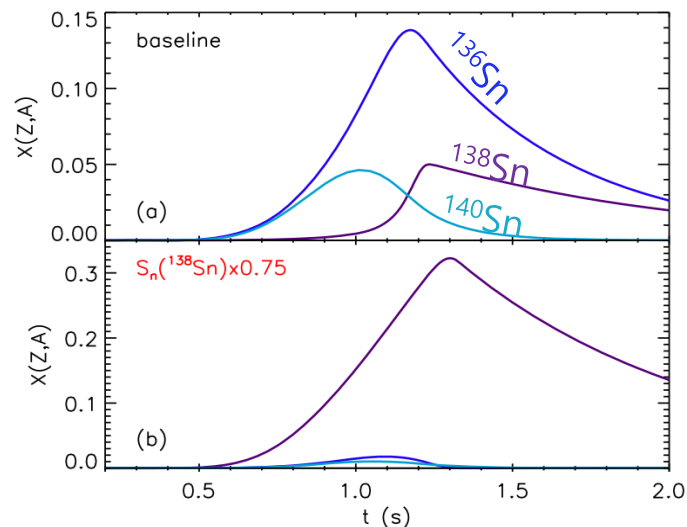
$$\frac{Y(Z, A+1)}{Y(Z, A)} \propto \exp\left[-\frac{S_n(Z, A+1)}{kT}\right]$$



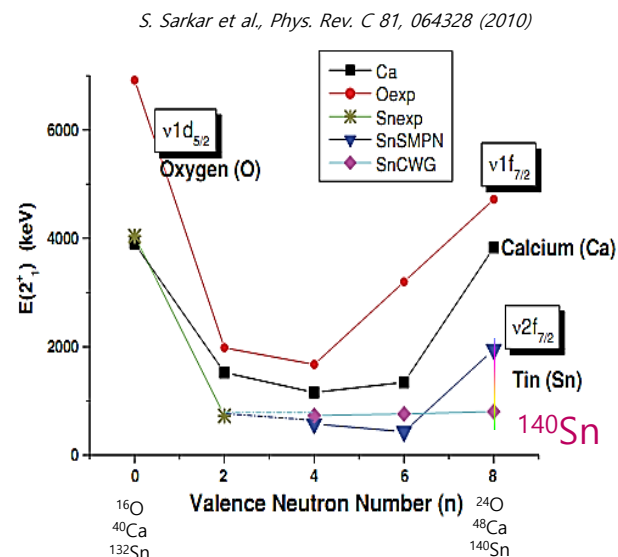
Particles 2020, 3(2), 320-335

II. Proposed mass measurements at RAON

❖ Isotopes of interest (neutron-rich $^{136-140}\text{Sn}$)

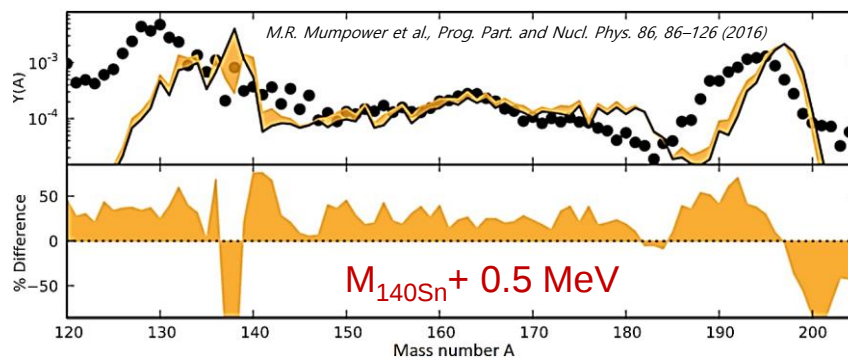


S. Brett et al., Eur. Phys. J. A 48, 184 (2012)



S. Sarkar et al., Phys. Rev. C 81, 064328 (2010)

- $M(^A\text{Sn})$ are essential for nuclear structure & astrophysics
- $M(^{140}\text{Sn})$ is important for:
 - Shell closure at $N = 82+8$?
 - r-process abundance



M.R. Mumpower et al., Prog. Part. and Nucl. Phys. 86, 86-126 (2016)

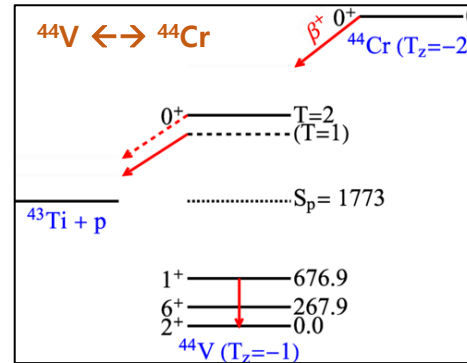
Isotope	m (u)	Δm (μ)	$T_{1/2}$ (ms)
^{136}Sn	135.939990	320	355
^{137}Sn	136.946550	430	249
^{138}Sn	137.951840	430	148
^{139}Sn	138.958730	430	120
^{140}Sn	139.962970	320	

AME2020

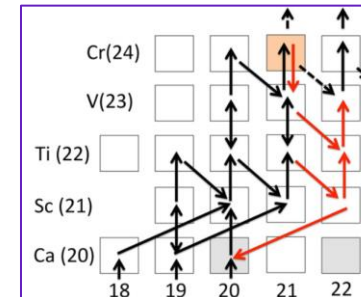
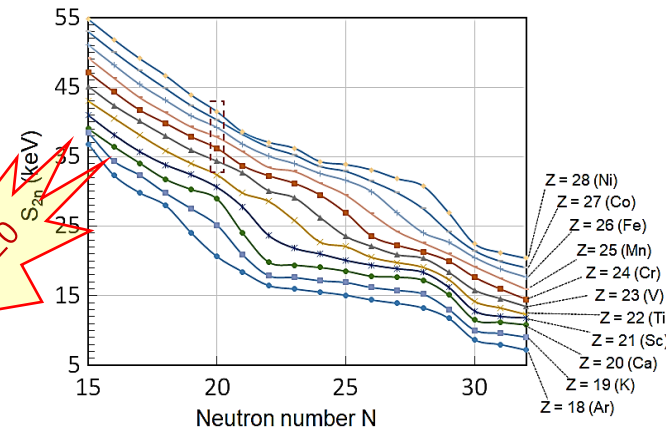
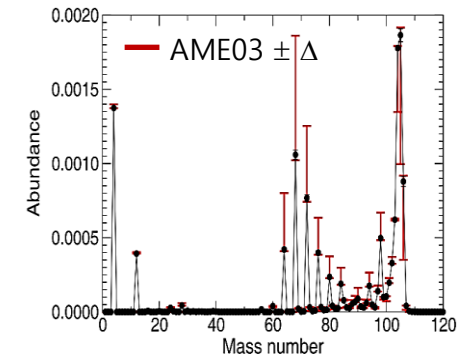
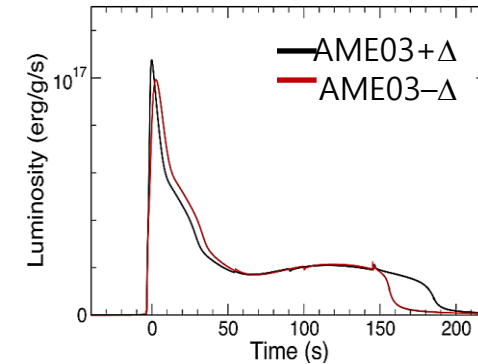
II. Proposed mass measurements at RAON

❖ Isotopes of interest (proton-rich Ti – Fe)

Isotope	Implications	Δm (keV)	$T_{1/2}$ (ms)
^{41}Ti	$^{40}\text{Sc}(p, \gamma)^{41}\text{Ti}$	28	82
^{43}V	$^{42}\text{Ti}(p, \gamma)^{43}\text{V}$ & $N=20$	43	80
^{44}V	$^{44}\text{V}(p, \gamma)^{45}\text{Cr}$ $^{43}\text{Ti}(p, \gamma)^{44}\text{V}$	80	111
^{44}Cr	$N=20$	56	43
^{45}Cr	$^{44}\text{V}(p, \gamma)^{45}\text{Cr}$	37	61
$^{45,46}\text{Mn}$	$N=20$	298, 84	#, 36
^{47}Mn	$^{46}\text{Cr}(p, \gamma)^{47}\text{Mn}$	28	88
$^{46,47}\text{Fe}$	$N=20$	298, 504	13, 22
^{48}Fe	$^{47}\text{Mn}(p, \gamma)^{48}\text{Fe}$	94	50



D. Puentes et al, Phys. Rev. C 101, (2020) 064309



➔ $S_p(^{45}\text{Cr}) = 2.139 \text{ MeV}$ (AME03)

➔ $S_p(^{45}\text{Cr}) = 2.684 \text{ MeV}$
 $\Delta m(^{45}\text{Cr}) = 35 \text{ keV}$

(Yan et al., Astrophys.J.Lett.766, L8)

II. Proposed mass measurements at RAON

❖ Beam productions

Proton beam

- Intensity: 1 μA
- Energy: 70 MeV

Target

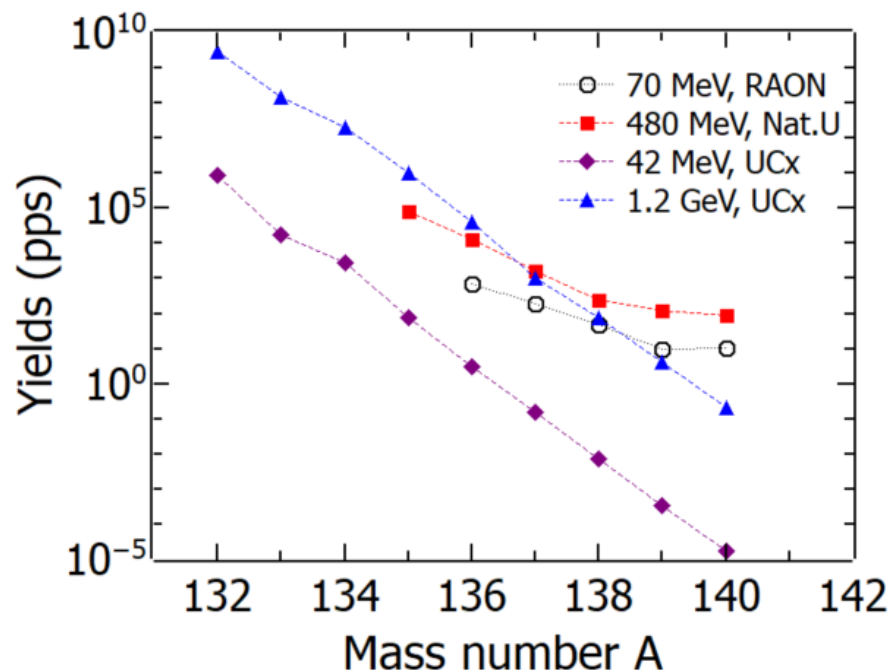
- Materials: UC_x
- Thickness: $18\text{g}/\text{cm}^2$
- Melting point: 2350°C

Driver: SC (Superconductive Cyclotron)

Ion Source: RILIS

(Resonance Ionization Laser Ion Source)
S.J. Park & J.B. Kim, *Hyper. Inter.* 241 (2020) 39

Neutron-rich $^{132-140}\text{Sn}$ isotopes



Estimated production yields of Tin isotopes using proton beams at 70 MeV (opened circles) based on $p+U$ reactions with projectile at 42 MeV (diamonds - HRIBF), 1.2 GeV (triangles - CERN), 480 MeV (squares - TRIUMF).

A	Yields
136	~ 750
137	~ 300
138	~ 80
139	~ 10
140	~ 10

Contaminants

I, Sb & Te

Requirement

UC_x targets !

II. Proposed mass measurements at RAON

❖ Beam productions

Proton beam

- Intensity: 1 μA
- Energy: 70 MeV

Target for $Z = 22 - 24$

- Materials: Ti
- Thickness: 0.05 mol/cm²
- Melting point: ~1660 °C
- Materials: Ni
- Thickness: 0.05 mol/cm²
- Melting point: ~1455 °C
- Max Yields: $10 - 10^4$ (pps/mmol.cm⁻²/μA)

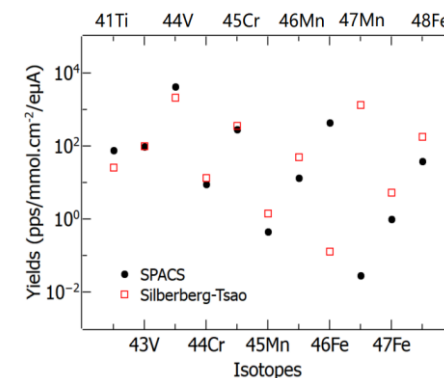
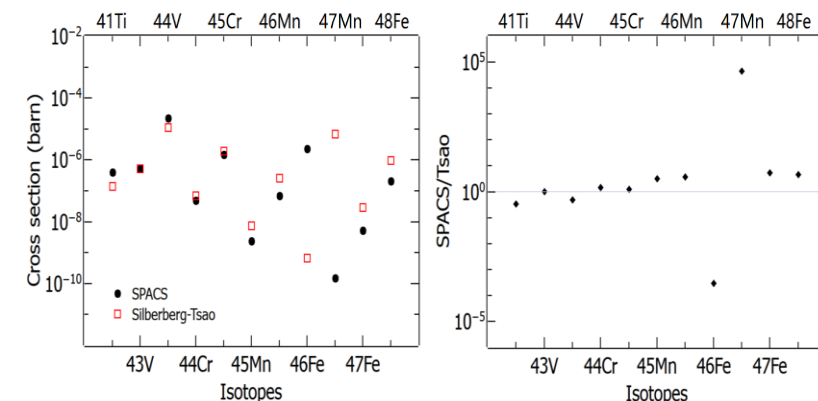
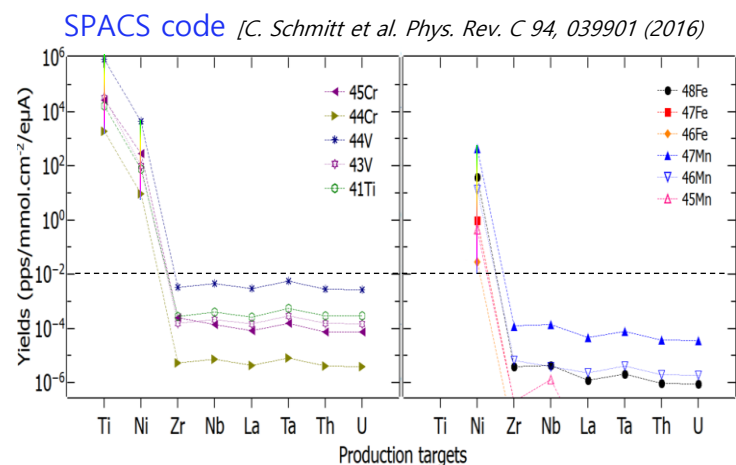
Driver: SC (Superconductive Cyclotron)

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(Resonance Ionization Laser Ion Source)

S.J. Park & J.B. Kim, Hyper. Inter. 241 (2020) 39

Proton-rich isotopes



⁴⁶Fe Min: 0.13 pps (Tsao)
Max: 435 pps (SPACS)

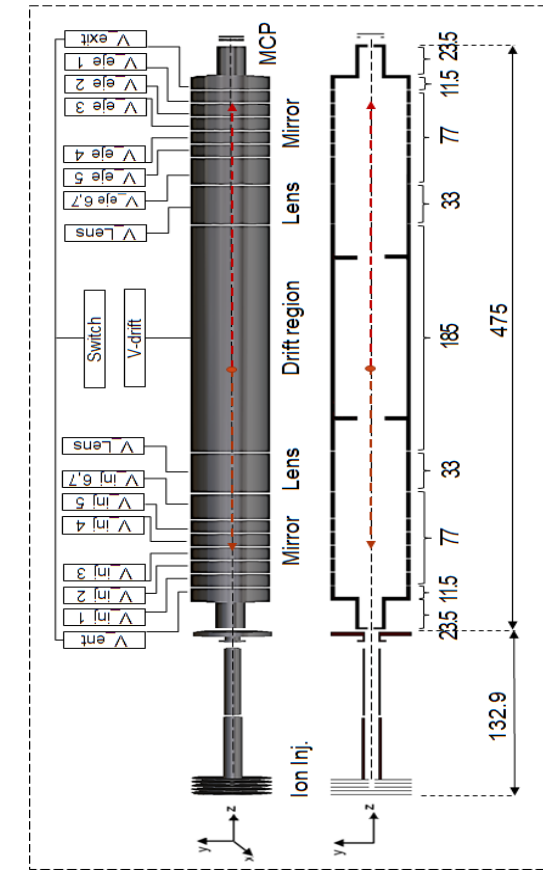
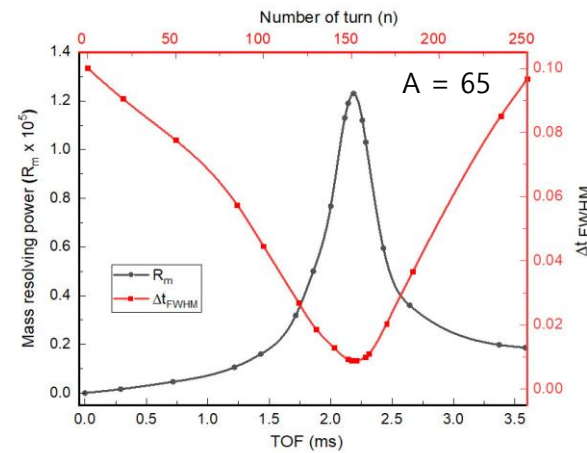
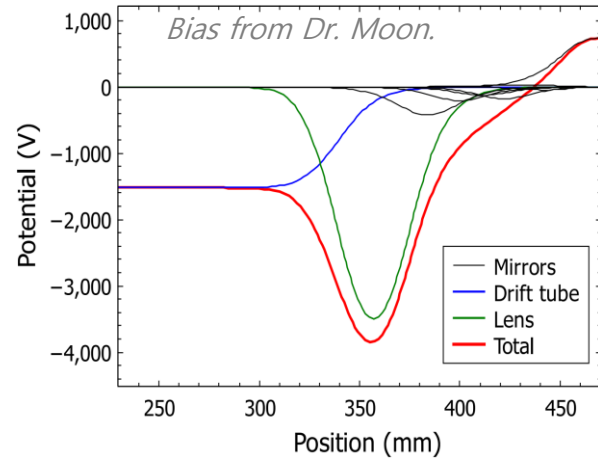
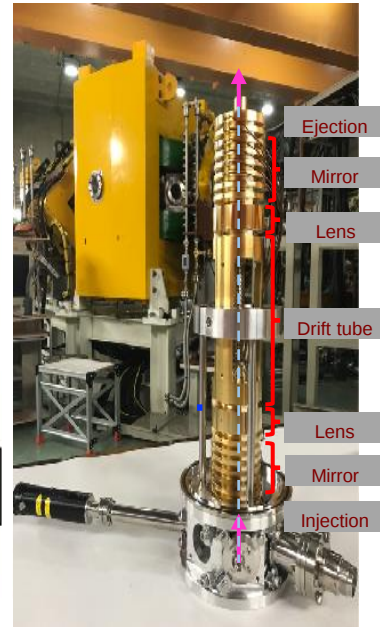
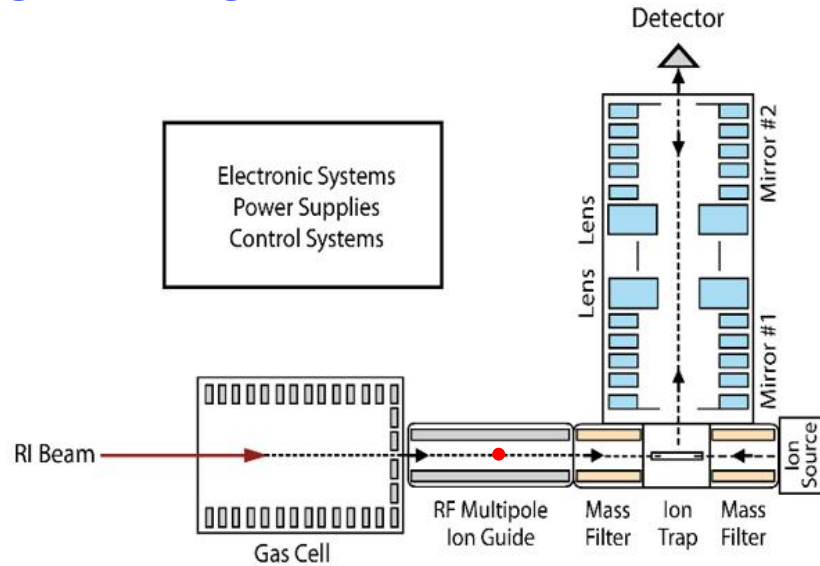
⁴⁷Mn Min: 0.03 pps (SPACS)
Yield_{max}: 1320 pps (Tsao)

Nuclei	⁴⁵ Mn	⁴⁶ Mn	⁴⁷ Mn	⁴⁶ Fe	⁴⁷ Fe	⁴⁸ Fe
mol/cm ²	Ni	Ni	Ni	Ni	Ni	Ni
Yields	4×10 ⁻¹	~10	4×10 ²	3 × 10 ⁻²	~1	~40

Requirement: Ni targets !

II. Proposed mass measurements at RAON

❖ RAON MR-TOF



Dimension from Dr. Moon.

Resolving power:

$$R_m = \sim 1.2 \times 10^5$$

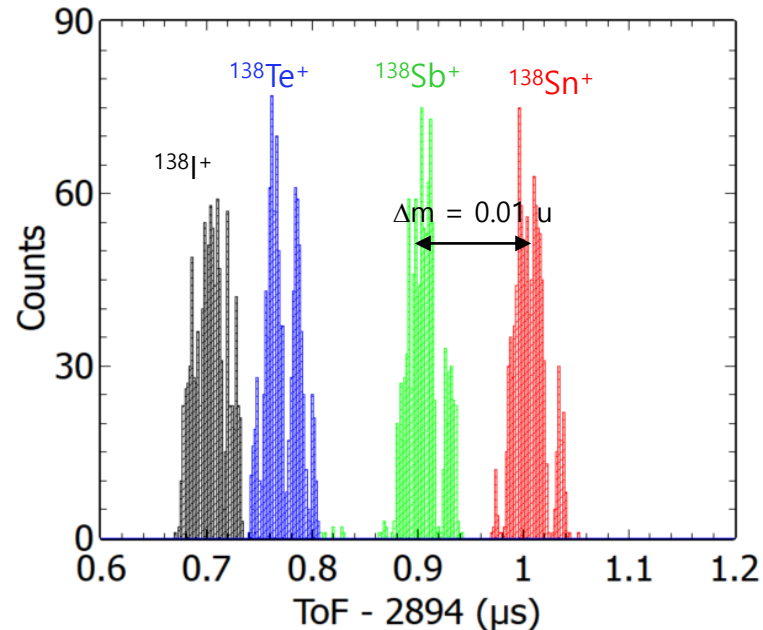
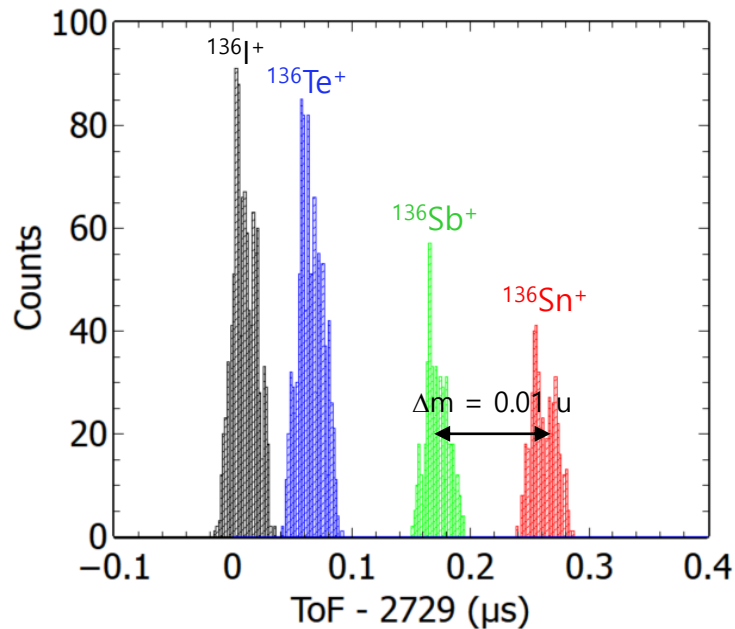
$$\text{FWHM} = \sim 10 \text{ ns}$$

II. Proposed mass measurements at RAON

❖ Separability

Ions	Mass (u)	$T_{1/2}$ (s)
^{136}Sn	135.93970	0.350
^{136}I	135.91460	83.400
^{136}Sb	135.93075	0.923
$^{136}\text{Te}^+$	135.92010	17.630

Ions	Mass (u)	$T_{1/2}$ (s)
^{138}Sn	137.95114	0.150
^{138}I	137.92273	6.230
^{138}Sb	137.94133	0.380
$^{138}\text{Te}^+$	137.92947	1.400



- $R_m = 1.2 \times 10^5$

with

$N = 134$ turns

Sn^+ & Sb^+ ions can be separated.

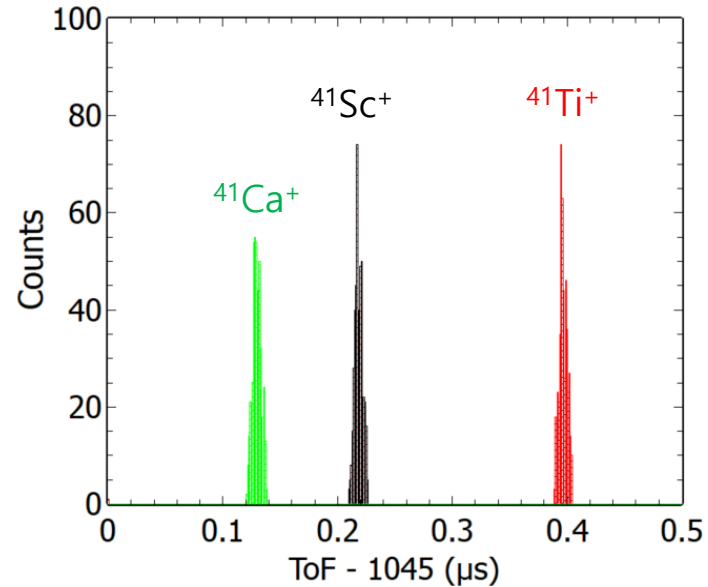
- $\text{ToF} = \sim 3 \text{ ms} \ll 150 \text{ ms}$

→ possible for measuring $^{136,138}\text{Sn}$.

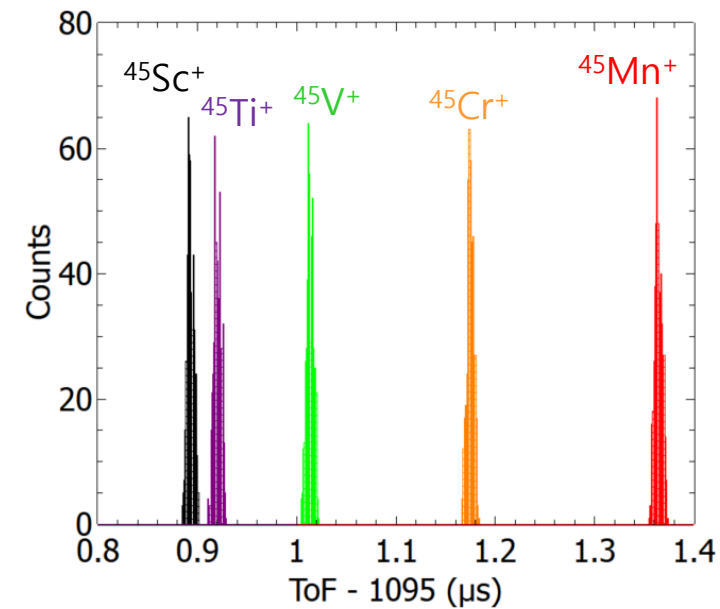
II. Proposed mass measurements at RAON

❖ Separability

Nuclide	Mass (u)	$T_{1/2}$
^{41}Ti	40.98315	82 ms
^{41}Sc	40.96925	596.3 ms
^{41}Ca	40.96228	99.4 ky



Nuclide	Mass (u)	$T_{1/2}$
^{45}Sc	44.95591	Stable
^{45}Ti	44.95812	184.8 m
^{45}V	44.96577	547 ms
^{45}Cr	44.97905	60.9 ms
^{45}Mn	44.99461	18 ms*



II. Proposed mass measurements at RAON

2 Letters of Intent

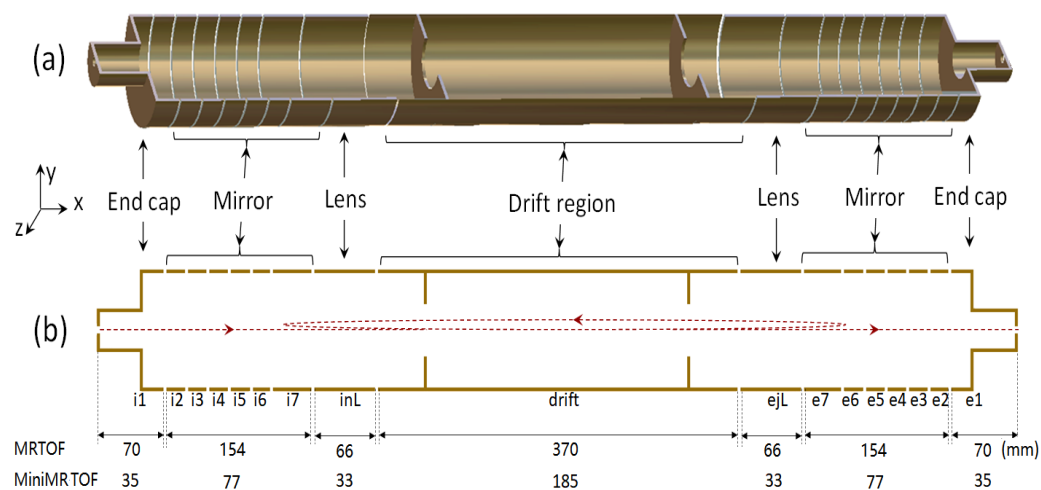
LOI 1: Neutron-rich nuclei

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LOI 2: Proton-rich nuclei

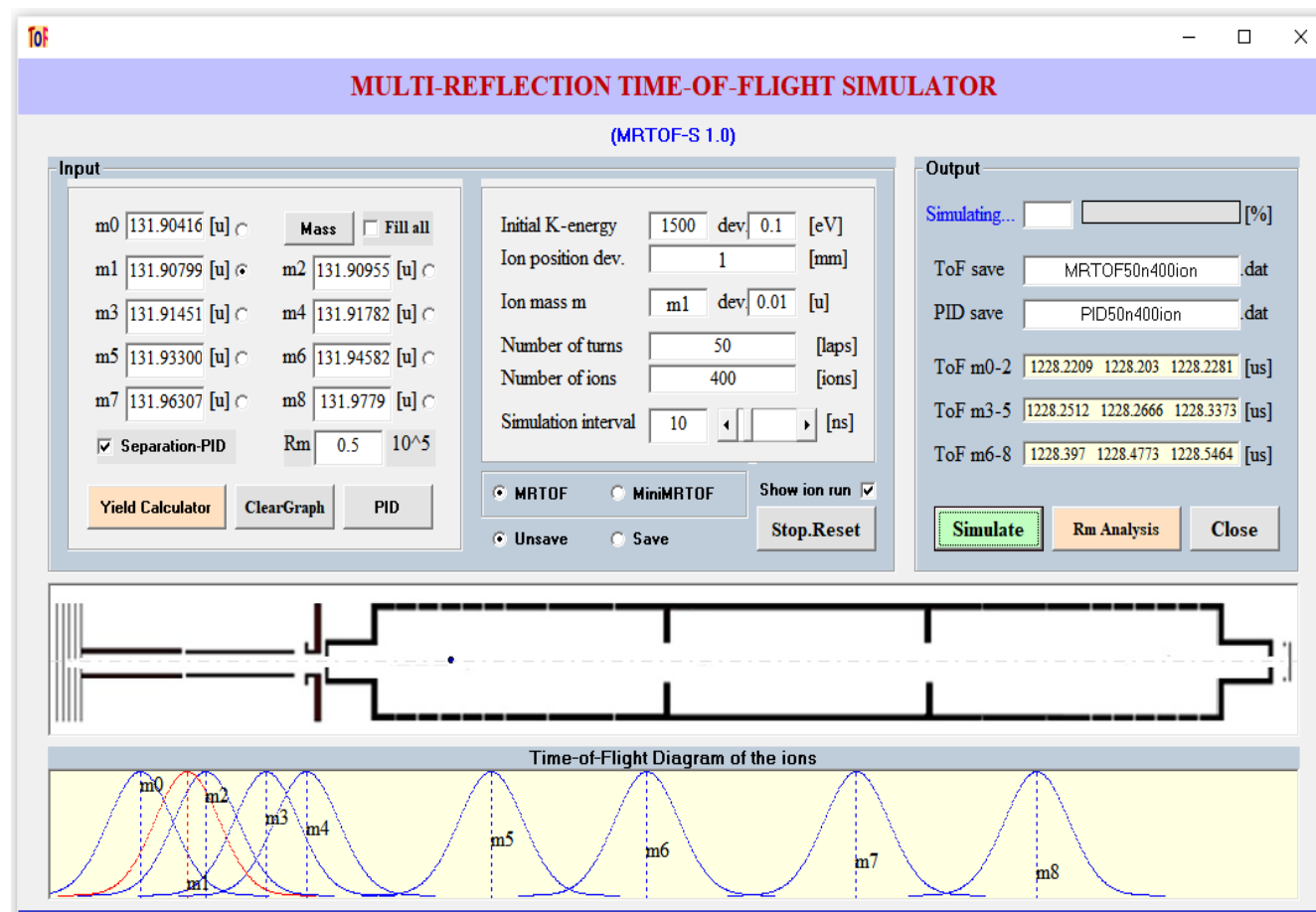
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III. MRTOF Simulator

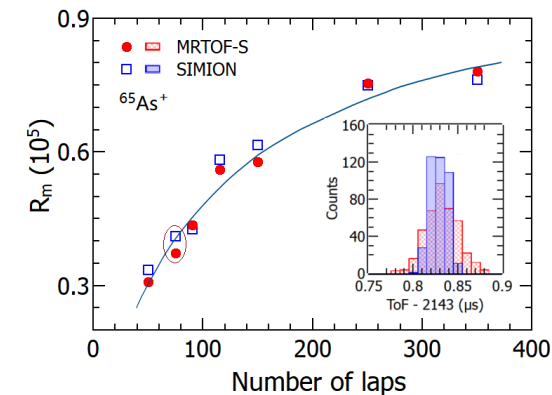
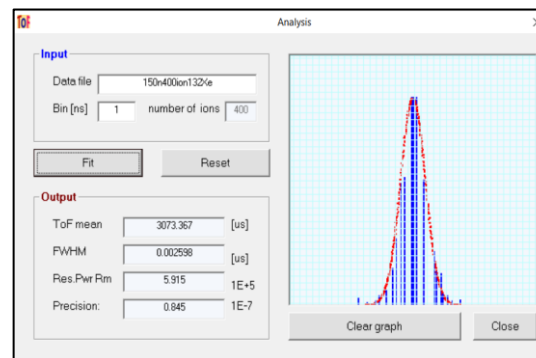
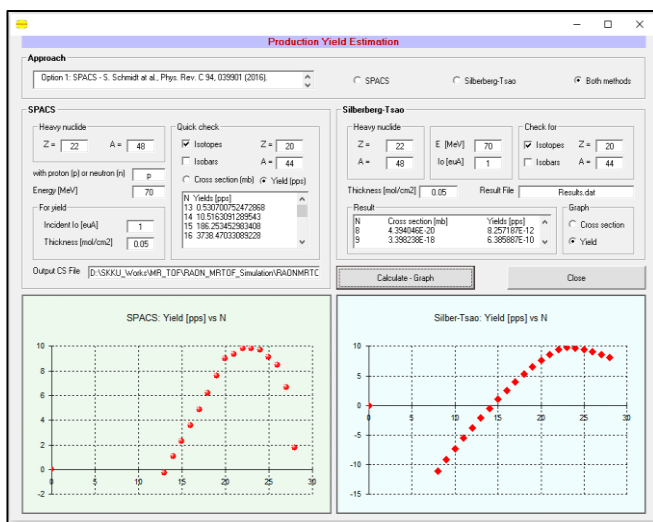
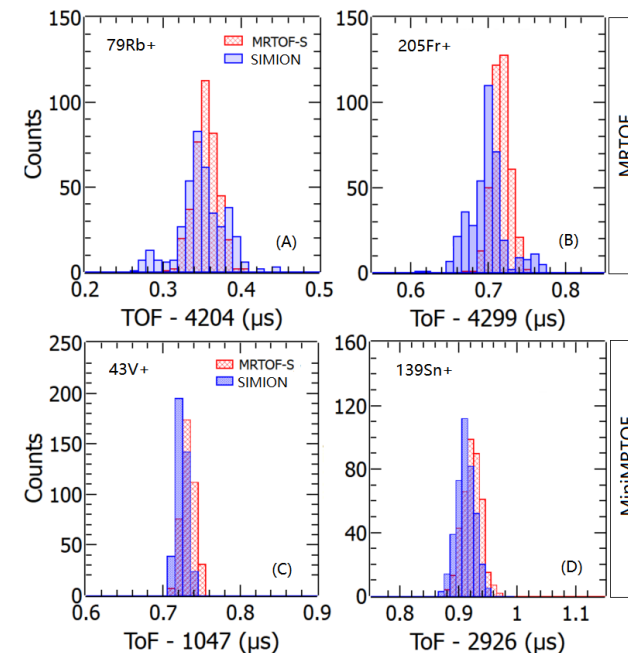
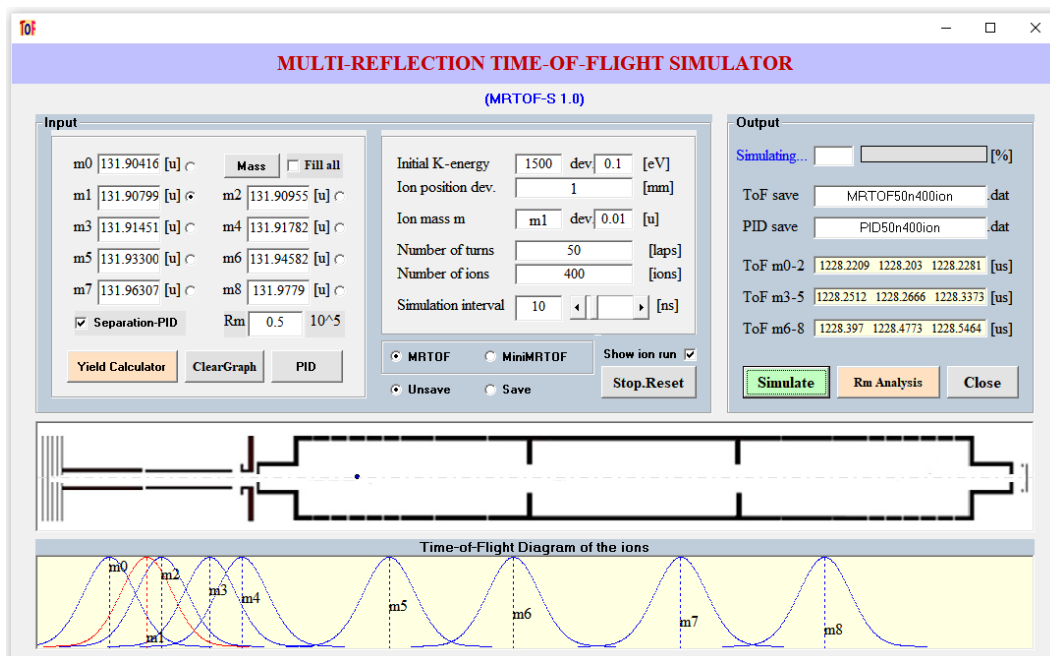


P. Schury *et al.*, Nucl. Instr. Meth. B **335**, 39 (2014)

J. Y. Moon *et al.*, Private communication.



III. MRTOF Simulator



IV. Summary

- Large uncertainty in nuclear mass of the ^{41}Ti , $^{43,44}\text{V}$, $^{44,45}\text{Cr}$, $^{45-47}\text{Mn}$, $^{46-48}\text{Fe}$ (proton-rich) and $^{136-140}\text{Sn}$ (neutron-rich) isotopes
- Precise masses of ^{41}Ti , $^{43,44}\text{V}$, $^{44,45}\text{Cr}$, $^{45-47}\text{Mn}$, $^{46-48}\text{Fe}$ (proton-rich) and $^{136-140}\text{Sn}$ (neutron-rich) isotopes are important for astrophysical calculations and shell closures.
- Ni target for 10^4 pps of Ti-Cr and UC_x target for 10 – 750 pps of $^{136-140}\text{Sn}$.
- RI beam production is necessary for the measurements.
- Mass resolving power of MR-TOF at RAON is about $10^5 \rightarrow$ uncertainty is better than 10 keV.
- Precise mass measurements at RAON are feasible.
- MRTOF Simulator is useful for pre-estimating parameters before experiments.

Thank you for your attention !