

# Superfluid optomechanics for dark matter direct detection

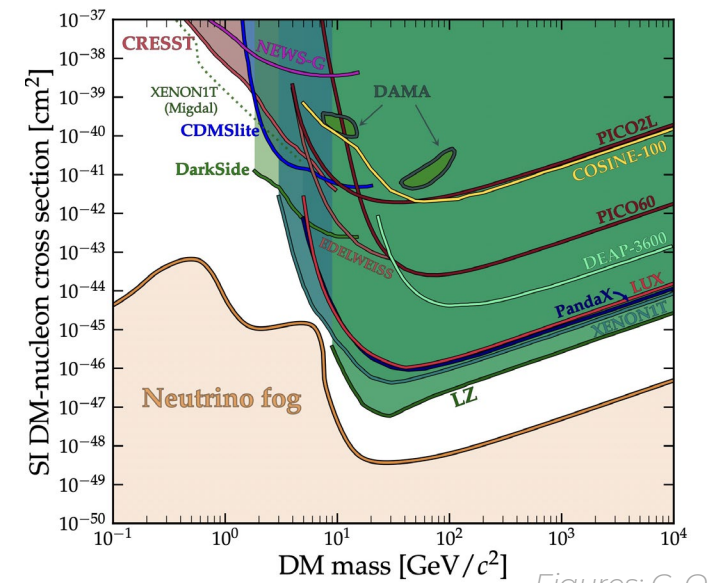
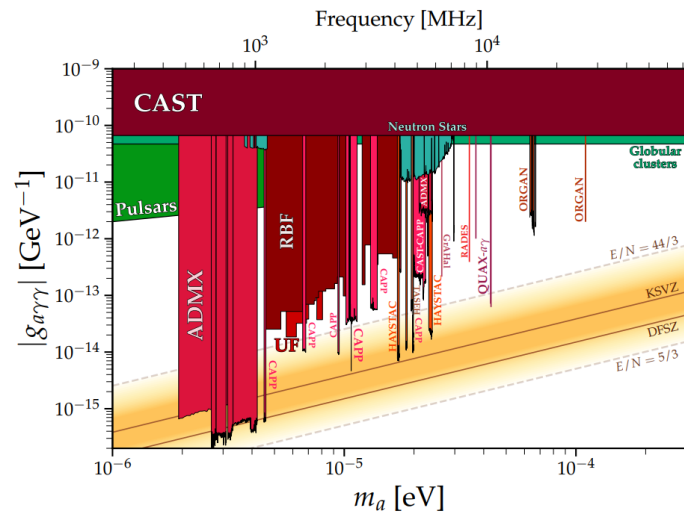
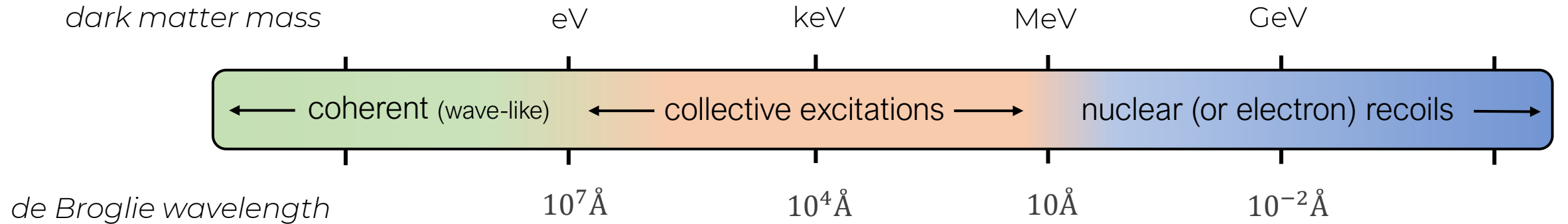
Peter Cox

*The University of Melbourne*

*with C. Baker, W. Bowen, M. Dolan, M. Goryachev, G. Harris*

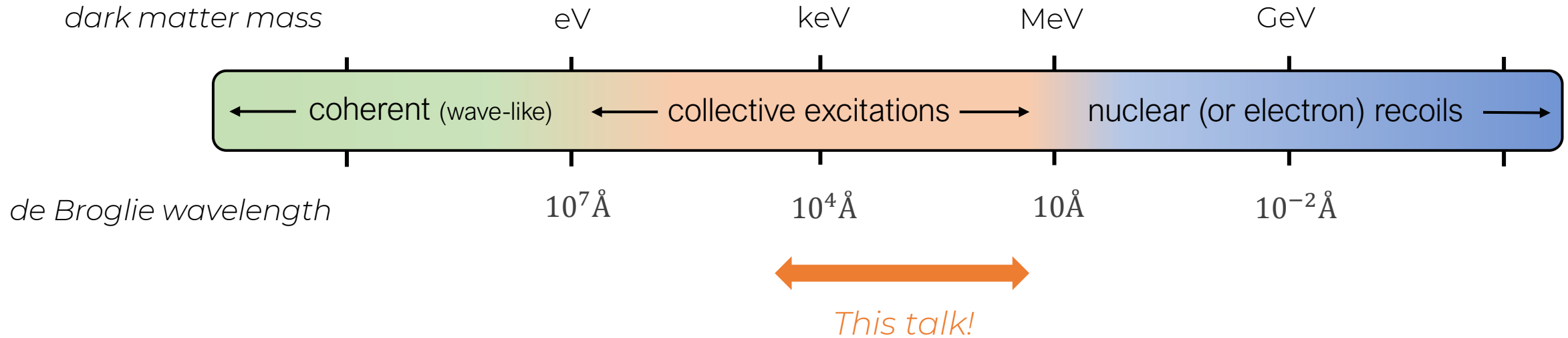


# Direct detection regimes



Figures: C. O'Hare

# Direct detection regimes



Direct detection in the keV – MeV regime:

- Relevant degrees of freedom are collective modes, not atoms
- Need new detectors with sub-eV thresholds

# Superfluid $^4\text{He}$

- Long-lived/stable collective modes (phonons/rotons)

How to detect these low-energy phonons?

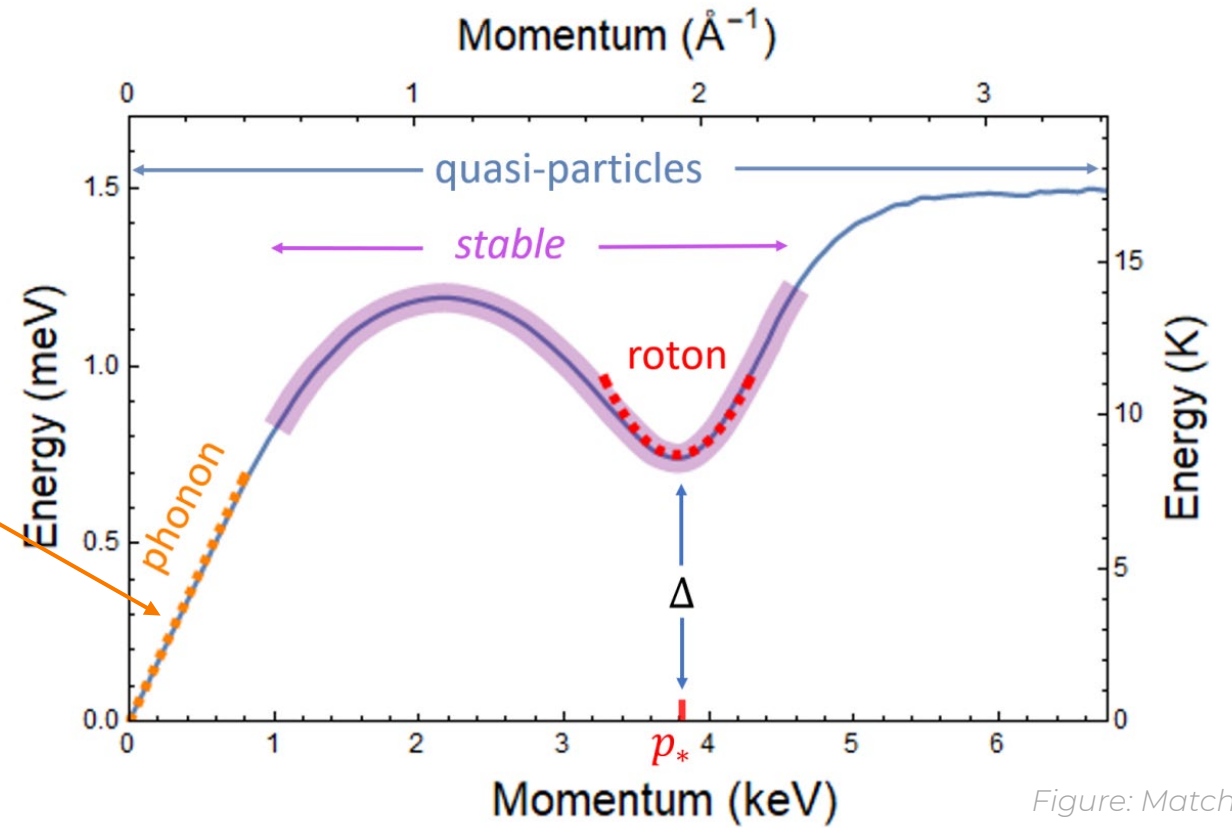


Figure: Matchev+ '21

# Superfluid $^4\text{He}$

- Long-lived/stable collective modes (phonons/rotons)

How to detect these low-energy phonons?

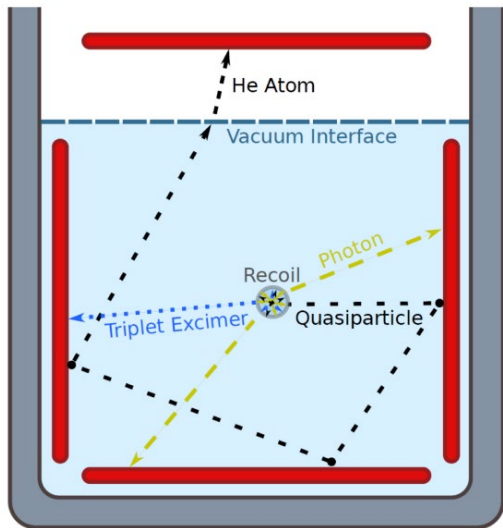


Figure: Hertel+ '18

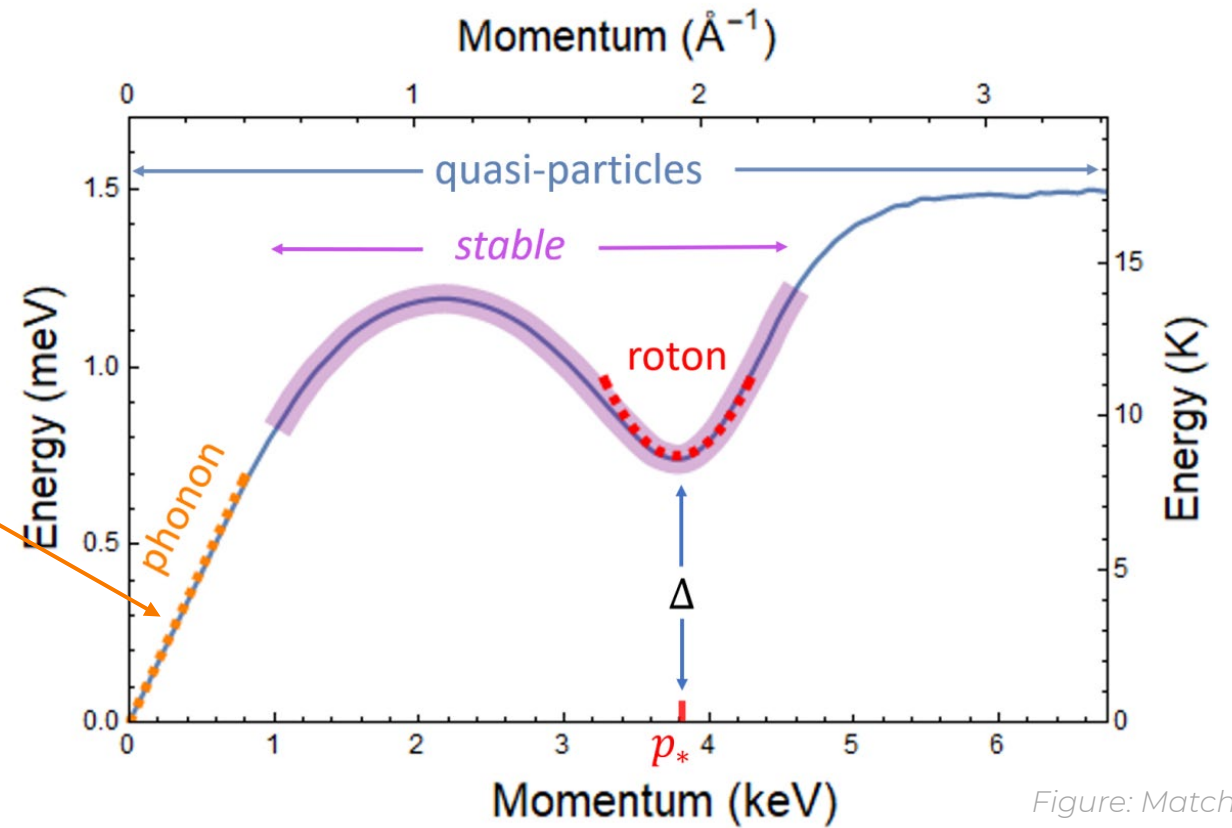


Figure: Matchev+ '21

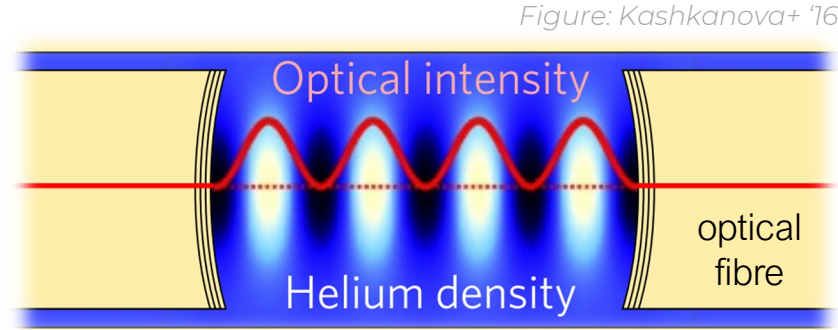
- Existing proposals use quantum evaporation signal in bulk superfluid (*HeRALD*, *DELIGHT*)

# Single phonon detection

Basic idea: Superfluid optomechanical systems are *single phonon detectors*

Coupling between acoustic (density) modes and optical modes:

➔ convert  $\sim \mu\text{eV}$  phonons into detectable  $\sim \text{eV}$  photons



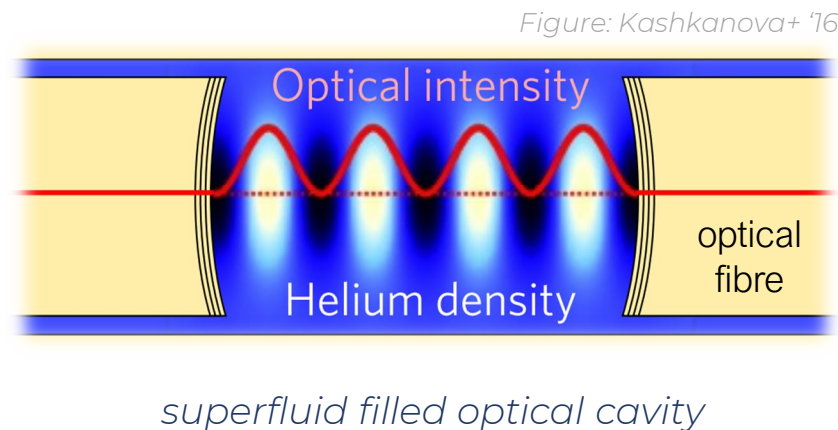
*superfluid filled optical cavity*

# Single phonon detection

Basic idea: Superfluid optomechanical systems are *single phonon detectors*

Coupling between acoustic (density) modes and optical modes:

➡ convert  $\sim \mu\text{eV}$  phonons into detectable  $\sim \text{eV}$  photons



Optomechanical systems have demonstrated  $\mu\text{eV}$  phonon counting (e.g. Patil et. al. '22)

# Superfluid optomechanics

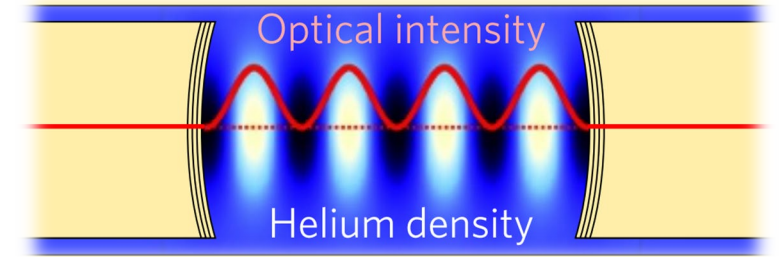
At quantum level described by the Hamiltonian:

$$H_{\text{OM}} = -g_0(a_{\gamma_1} a_{\gamma_2}^\dagger b_m + a_{\gamma_1}^\dagger a_{\gamma_2} b_m^\dagger)$$

photons      phonon

optomechanical coupling

Figure: Kashkanova+ '16





# Superfluid optomechanics

At quantum level described by the Hamiltonian:

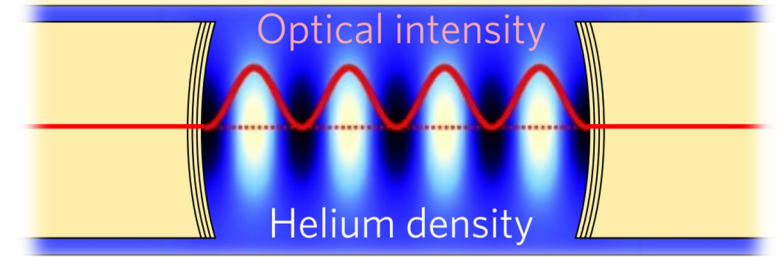
$$H_{\text{OM}} = -g_0(a_{\gamma_1} a_{\gamma_2}^\dagger b_m + a_{\gamma_1}^\dagger a_{\gamma_2} b_m^\dagger)$$
$$\rightarrow -g_0 \sqrt{N_1} (a_{\gamma_2}^\dagger b_m + a_{\gamma_2} b_m^\dagger)$$

Diagram illustrating the Hamiltonian components:

- photons (orange arrows pointing to  $a_{\gamma_1}$  and  $a_{\gamma_2}$ )
- phonon (blue arrow pointing to  $b_m$ )

pump laser enhances small  $g_0$

Figure: Kashkanova+ '16



Acoustic mode energy/wavelength related to optical modes:

$$\Omega_m = \pm (\omega_{\gamma_1} - \omega_{\gamma_2})$$

$$\lambda_m \approx \lambda_\gamma / 2$$

# Phonon lasing

Superfluid optomechanical systems as dark matter detectors:

- ✓ exceptional low-energy sensitivity ( $\sim \mu\text{eV}$ )
  - ✗ narrow-band detector (single phonon energy)
- ➡ Very low dark matter scattering rate due to restricted phase space

# Phonon lasing

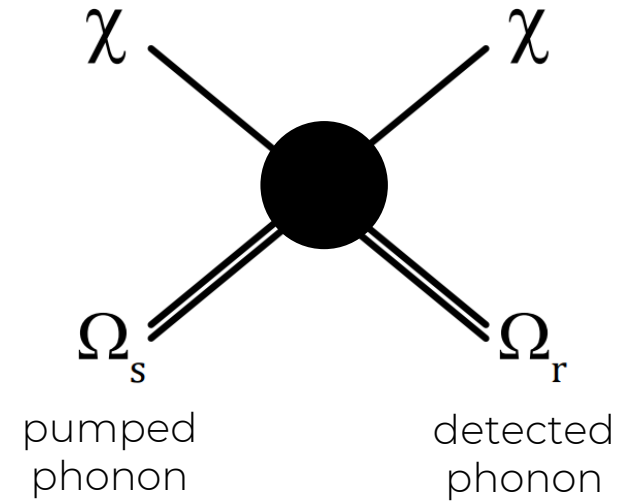
Superfluid optomechanical systems as dark matter detectors:

- ✓ exceptional low-energy sensitivity ( $\sim \mu\text{eV}$ )
- ✗ narrow-band detector (single phonon energy)

➔ Very low dark matter scattering rate due to restricted phase space

*Solution:* Phonon lasing

➔ Stimulated scattering rate proportional to phonon occupation number

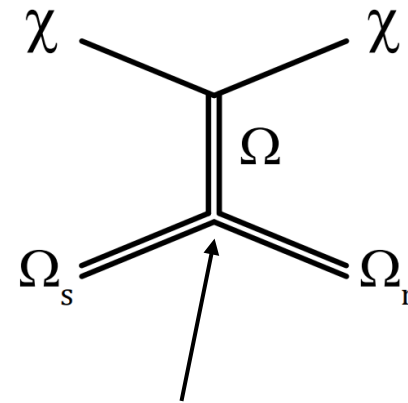
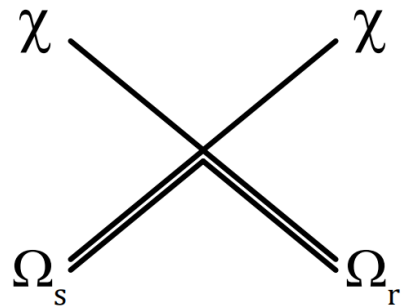


# Dark matter – phonon scattering

Low-energy phonons in superfluid described by effective field theory (EFT)

- Phonons are Nambu-Goldstone bosons of U(1) symmetry breaking

Two contributions to stimulated scattering process:



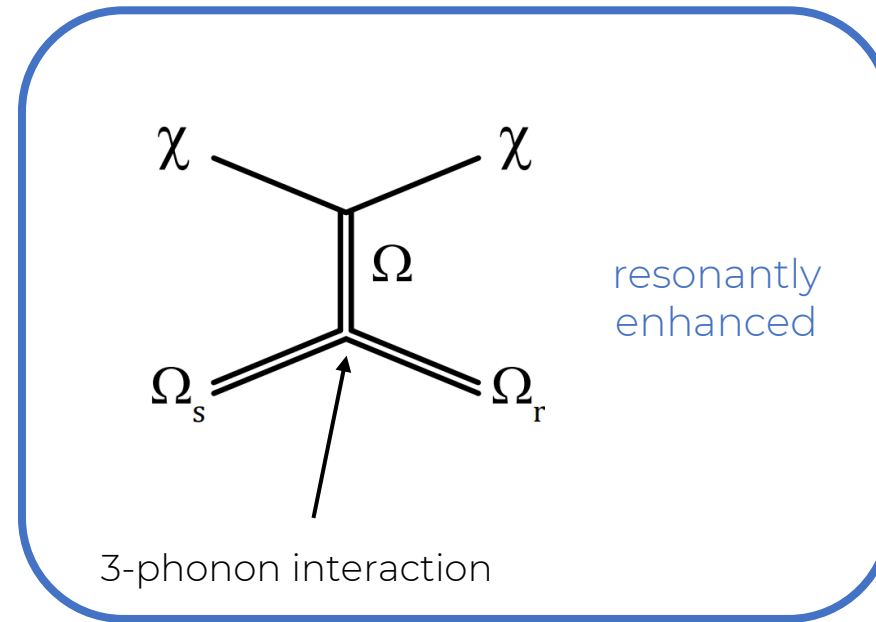
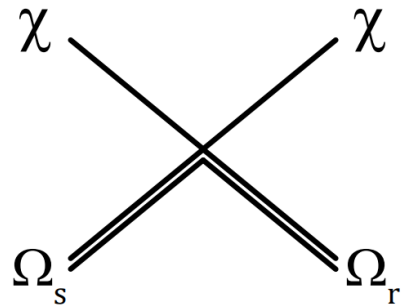
3-phonon interaction

# Dark matter – phonon scattering

Low-energy phonons in superfluid described by effective field theory (EFT)

- Phonons are Nambu-Goldstone bosons of U(1) symmetry breaking

Two contributions to stimulated scattering process:



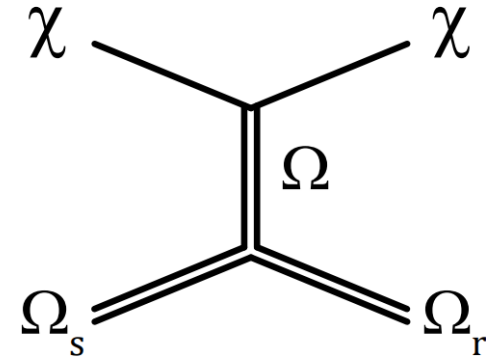
# Scattering rate

Initial state phonon number density

$$R \propto \frac{\rho_\chi \sigma_\chi n}{m_\chi^3} n_s Q^2 \frac{\Omega_r \Omega_s}{m_{\text{He}}^2 c_s^4} \frac{(1 + \gamma_G)^2}{1}$$

Acoustic quality factor:  $Q = \Omega_m / \Gamma_m$

3-phonon interaction  
(Grüneisen parameter)



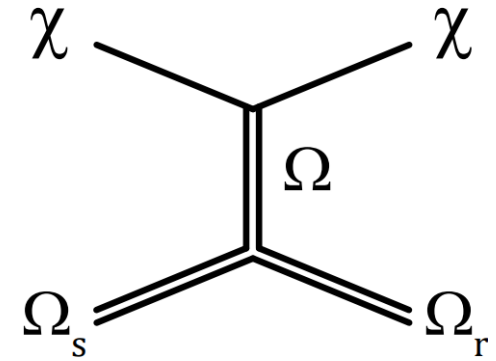
# Scattering rate

Initial state phonon number density

$$R \propto \frac{\rho_\chi \sigma_\chi n}{m_\chi^3} n_s Q^2 \frac{\Omega_r \Omega_s}{m_{\text{He}}^2 c_s^4} \frac{(1 + \gamma_G)^2}{1}$$

Acoustic quality factor:  $Q = \Omega_m / \Gamma_m$

3-phonon interaction  
(Grüneisen parameter)

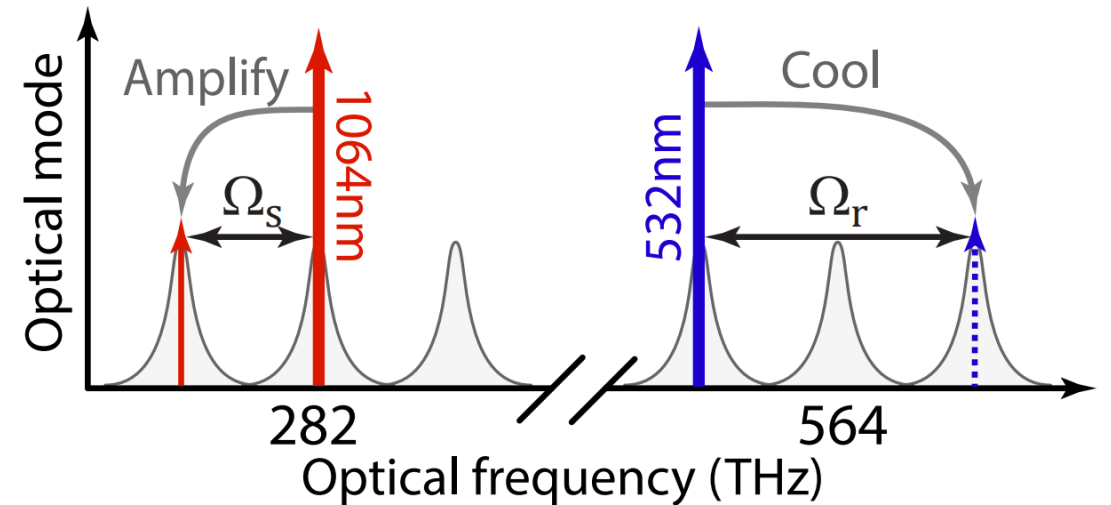
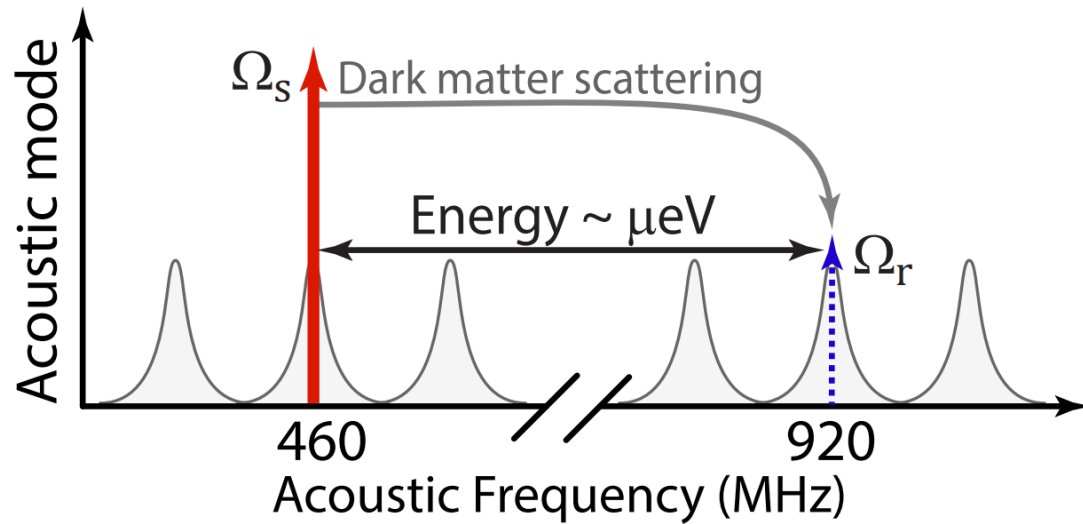
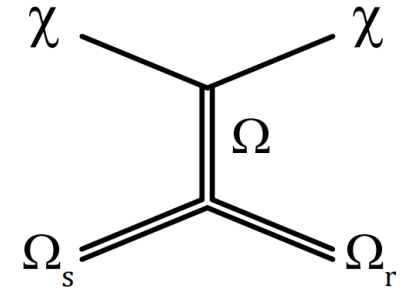


Scattering is between specific initial *and final* phonon states:

- I. Scattering is at fixed momentum transfer:  $q = (\Omega_r - \Omega_m) / c_s$
- II. Event rate is *independent* of detector volume!

# Optomechanical detection

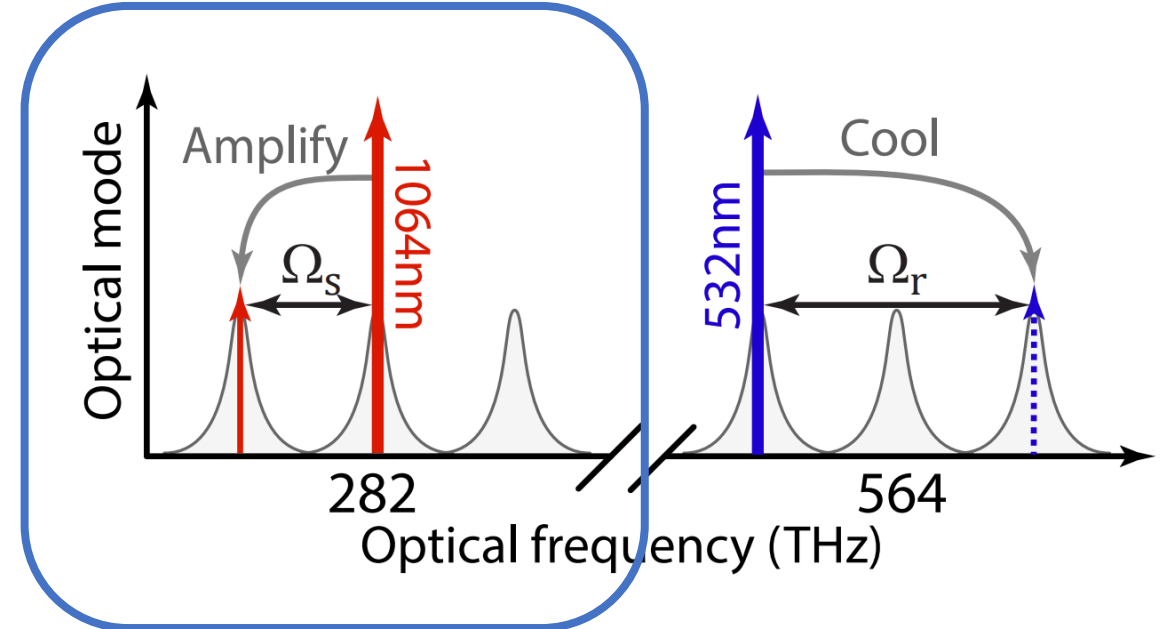
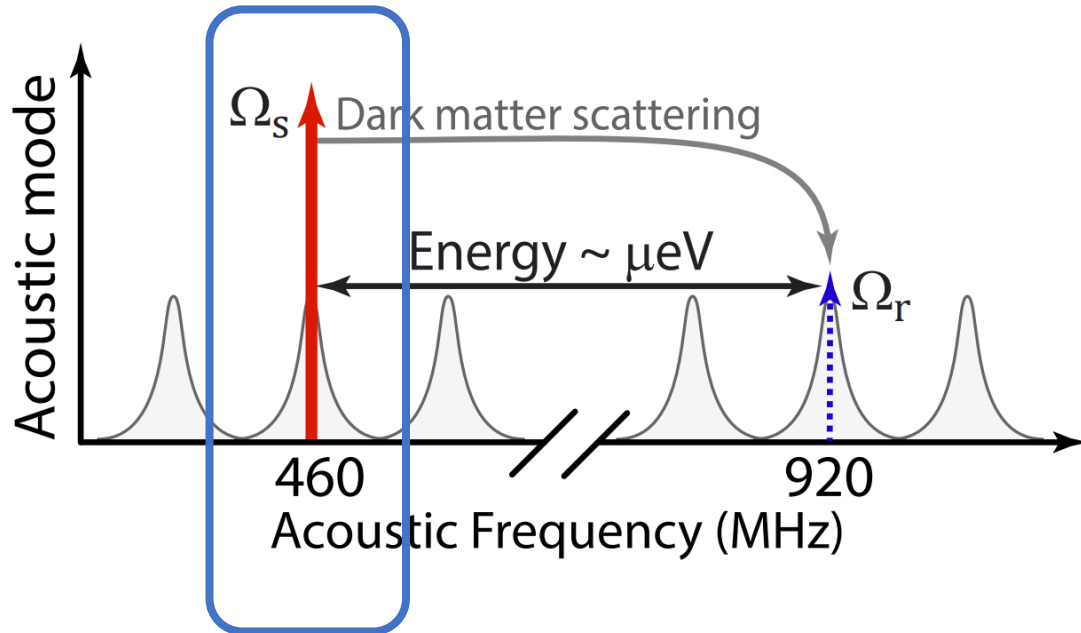
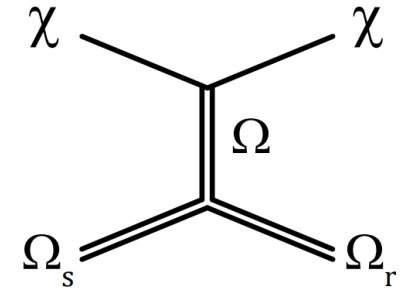
Dark matter detector requires optomechanical control of *two* acoustic modes





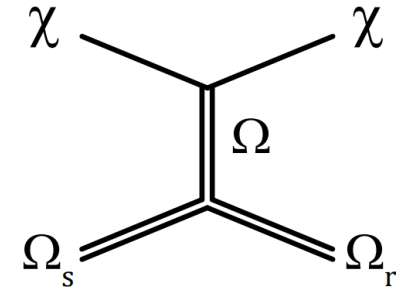
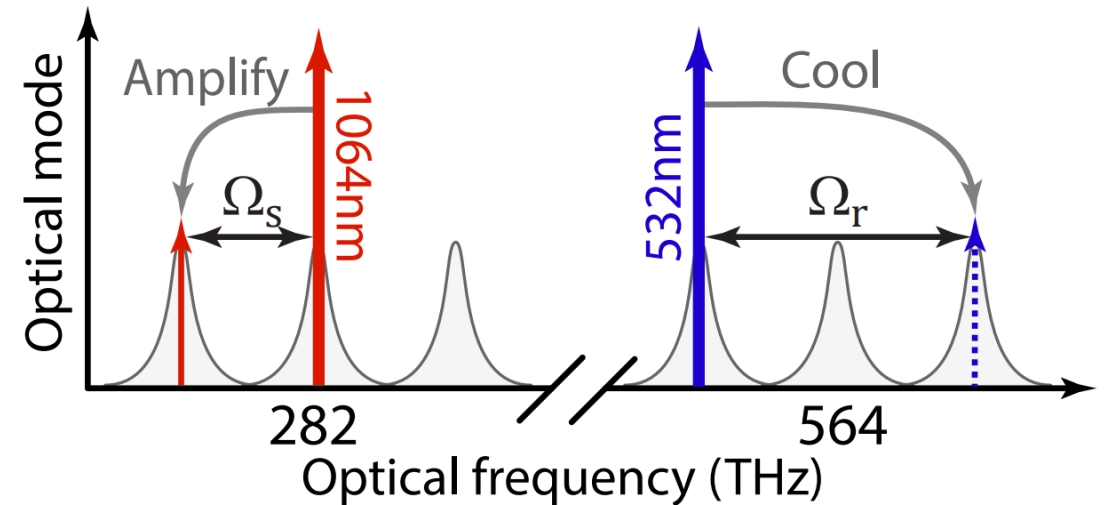
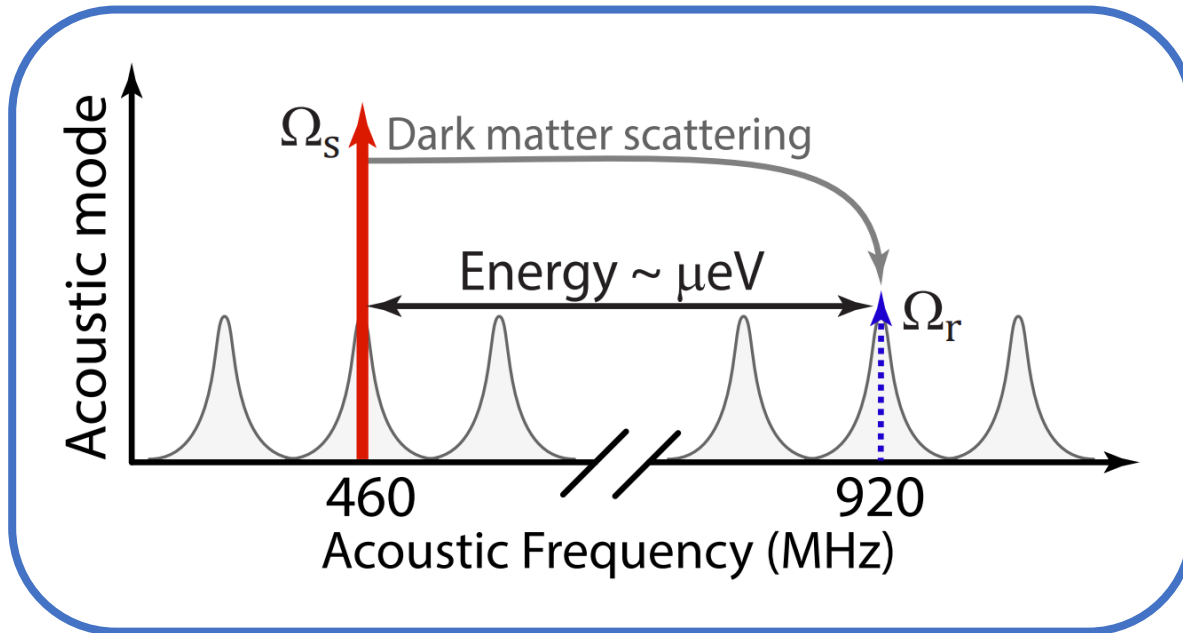
# Optomechanical detection

- 1 Lower-energy phonon mode  $\Omega_s$  populated via optomechanical interaction



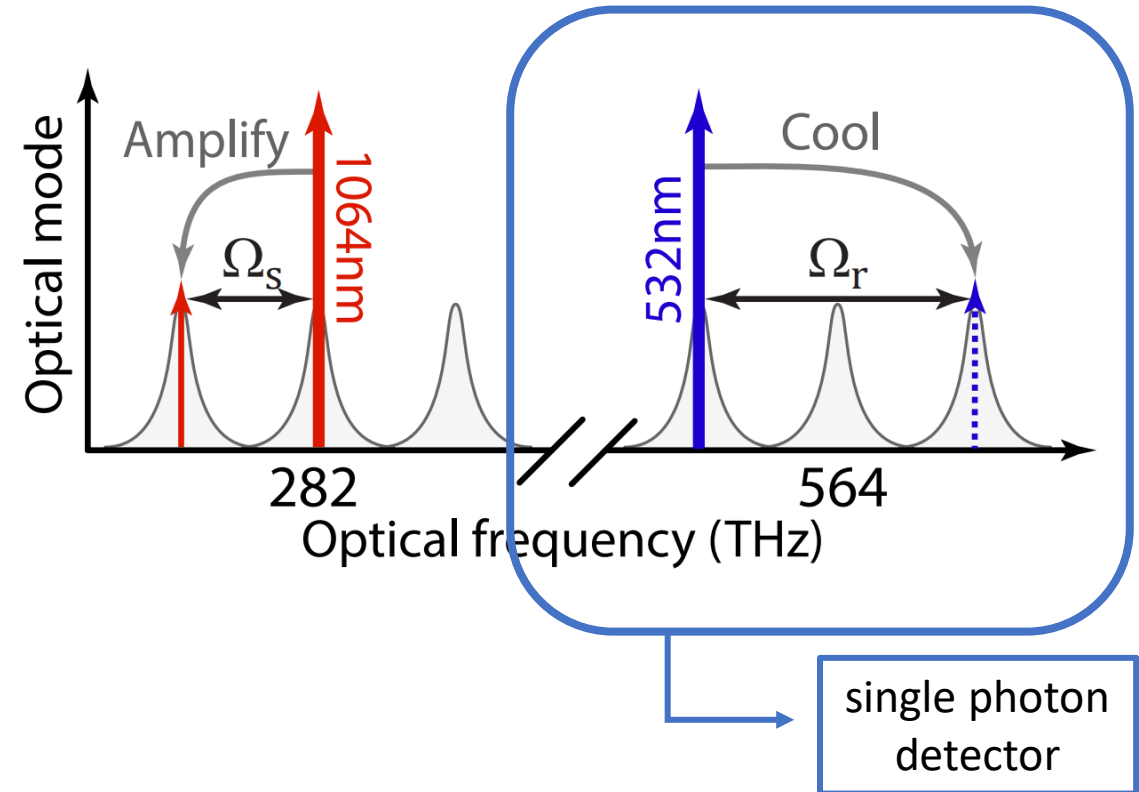
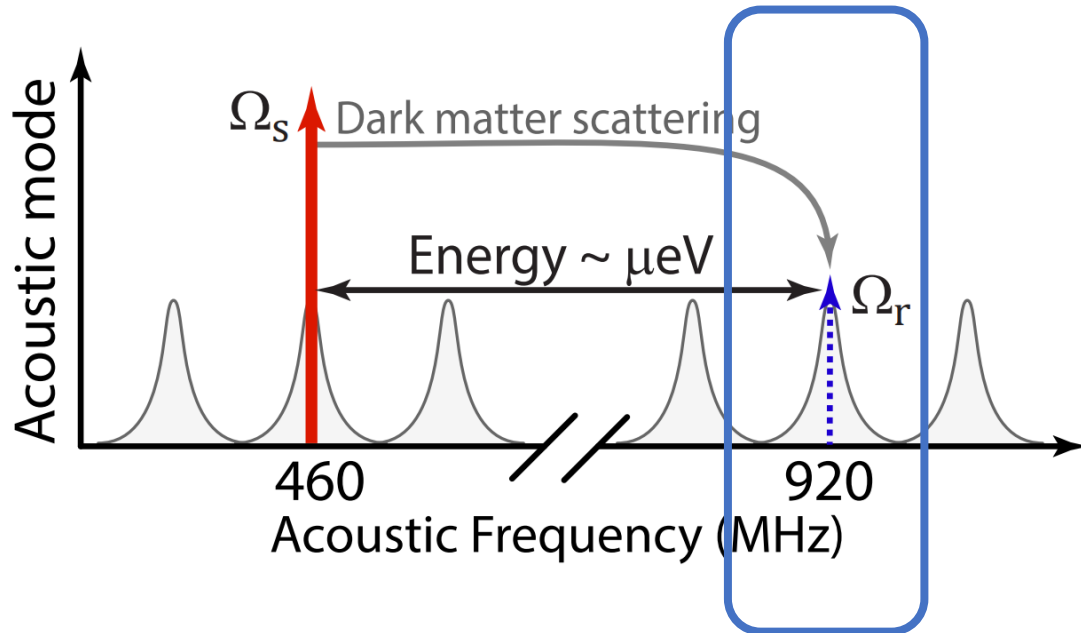
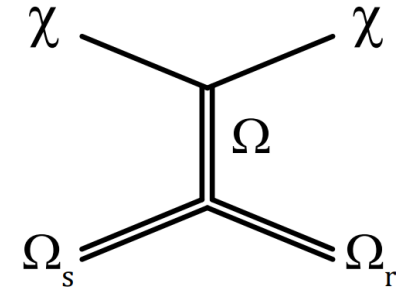
# Optomechanical detection

- 2 Stimulated dark matter scattering excites higher energy phonon mode  $\Omega_s \rightarrow \Omega_r$

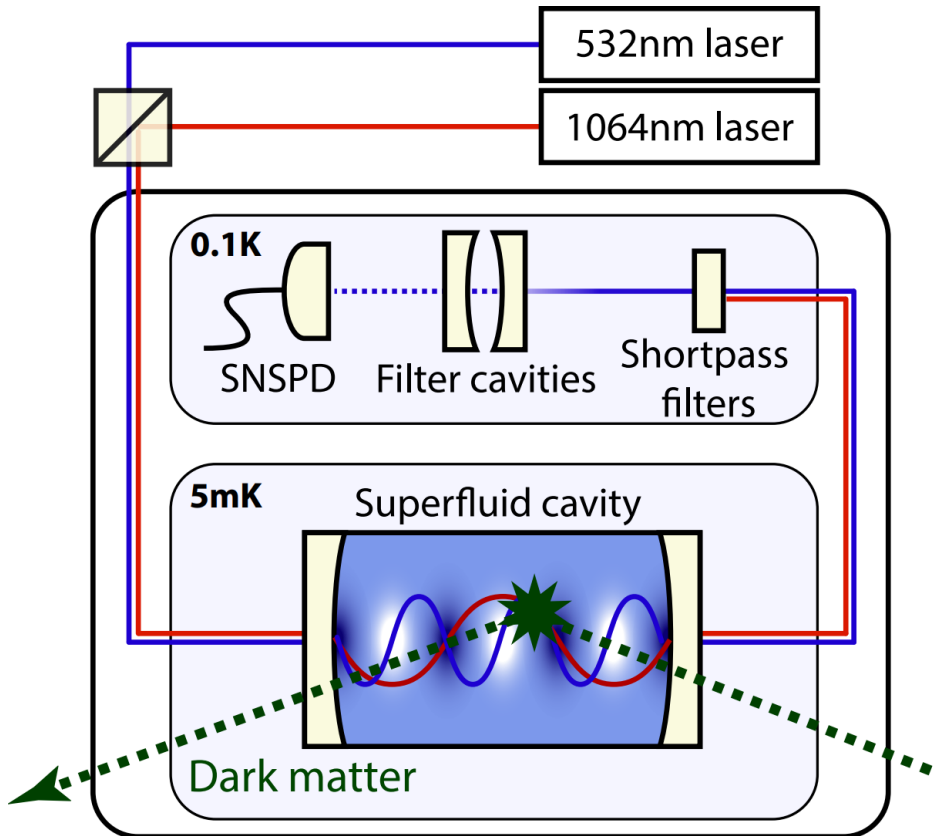


# Optomechanical detection

- 3 Optomechanical conversion of  $\Omega_r$  phonon to photon that is detected with SNSPD



# Schematic experiment

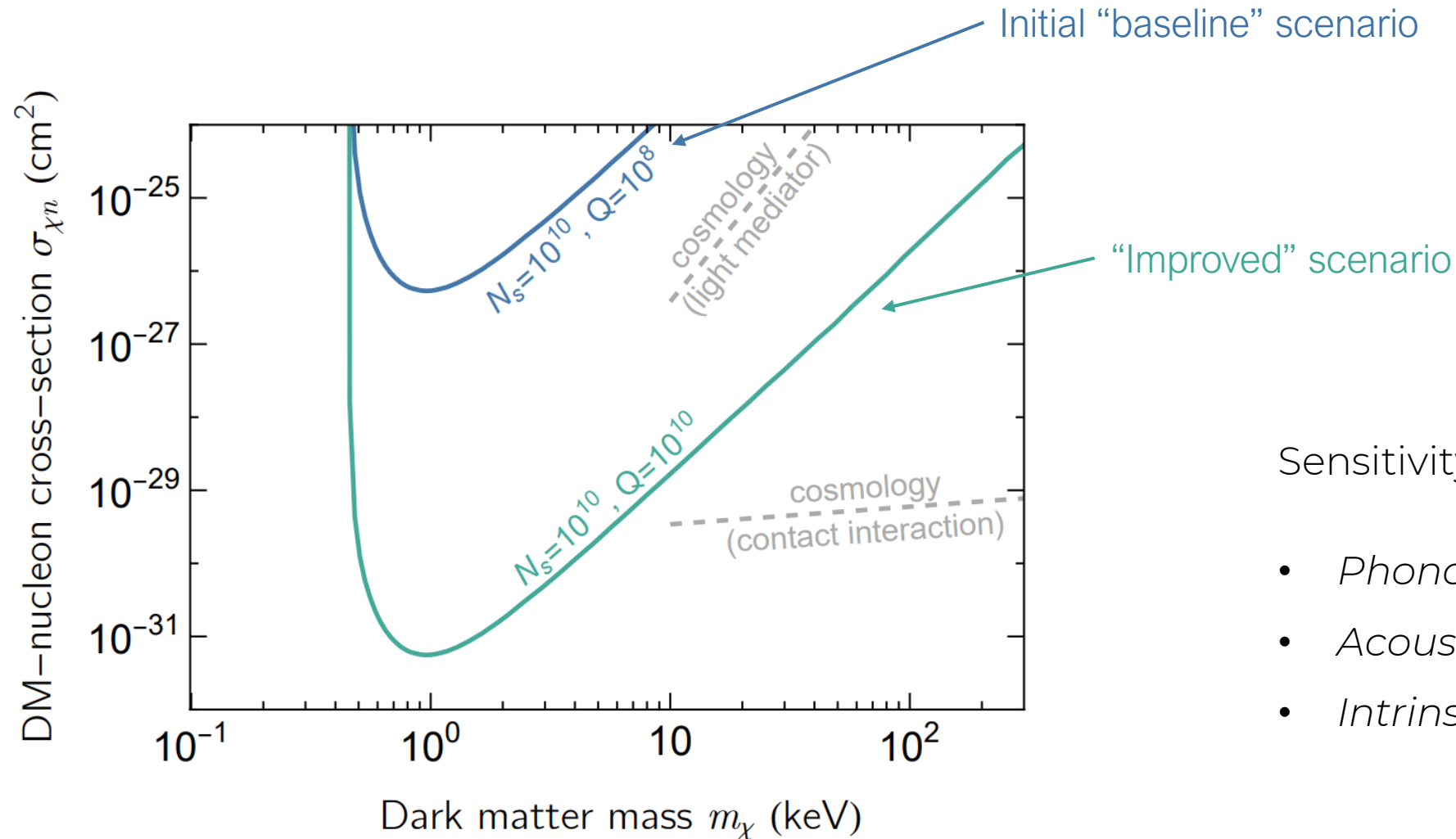


Important detector backgrounds:

- *Thermal phonons*  
( $10^{-5}$  Hz at  $T = 4\text{mK}$  and  $Q = 10^{10}$ )
- *Dark counts in SNSPD*  
( $\sim 6 \times 10^{-6}$  Hz)
- *Photons from pump lasers*  
(especially 532nm, suppressed with filter cavities)

Expected background rate  $\sim 1$  event/day

# Projected Sensitivity

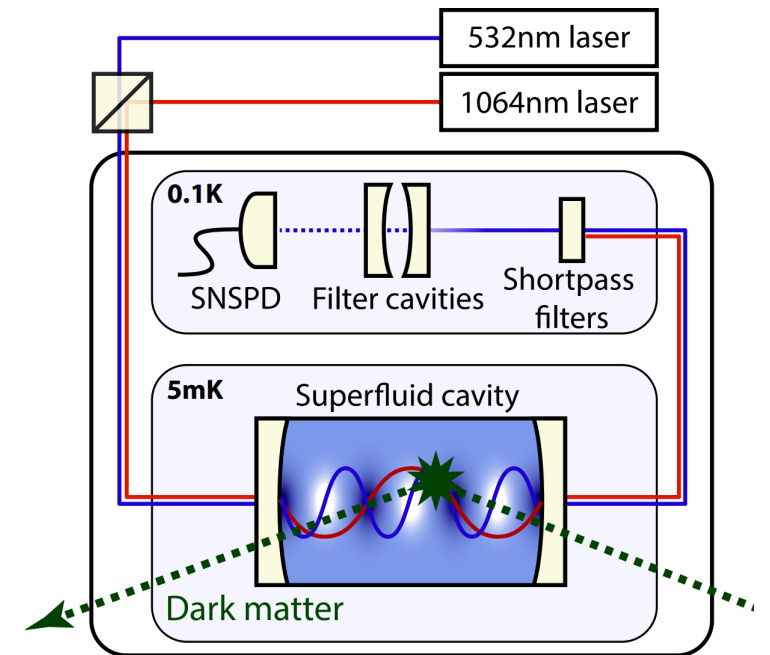


Sensitivity primarily determined by:

- Phonon occupation  $N_s$
- Acoustic  $Q$ -factor ( $\Omega_m/\Gamma_m$ )
- Intrinsic background rate

# Outlook

- Superfluid optomechanical systems are *single phonon detectors*
- Conversion of  $\sim\mu\text{eV}$  phonons to  $\sim\text{eV}$  photons
- Phonon lasing enhances dark matter event rate
- Sensitive to dark matter in the keV mass range with cross-sections as low as  $\mathcal{O}(10^{-32}) \text{ cm}^2$



# Backup

# Optical asymmetry

Optical mode spacing (FSR) can be engineered to select amplification/cooling of acoustic modes:

