

Anapole Moment of Neutralino Dark Matter

Merlin Reichard

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based on A. Ibarra, MR, R. Nagai [JHEP 01 (2023) 086]
(+ A. Ibarra, MR, G. Tomar *work in progress*)

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15.06.2023



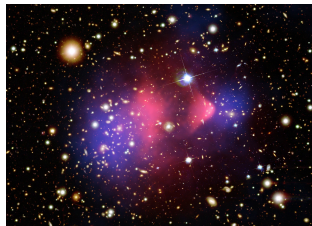
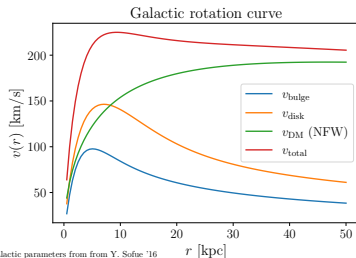
Is dark matter really »dark«?

BlueWillow: "a dark matter particle interacting with the
electromagnetic field"

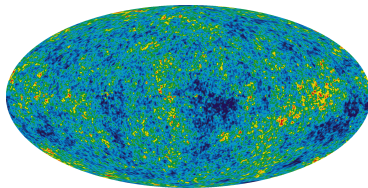


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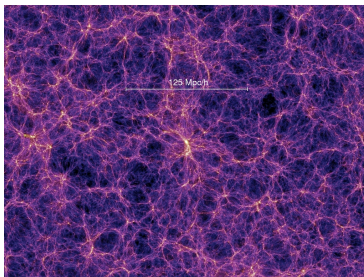
Motivation: DM Evidence



Credit: NASA



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Credit: Springel et al. '05 (Millennium Simulation)

Motivation: DM Direct Searches

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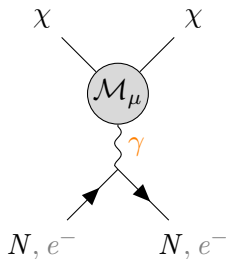
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- Information from particle physics (DM mass, cross section) and astrophysics (local abundance, velocity distribution)
- For example: DM interacting electromagnetically

Anapole Moment: The Setup

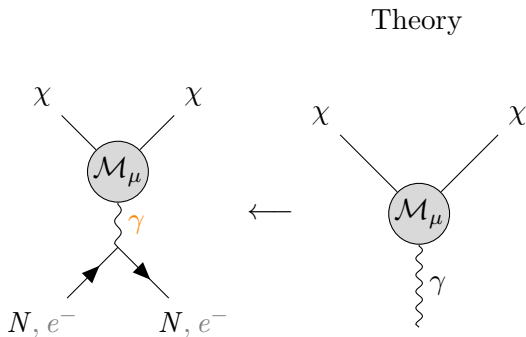
Effective **photon** interaction



[See also hep-ph/0003010, 1401.6457, 1503.01500, 1608.00642, 1808.04112, 1806.07896, 2003.04039, ...]

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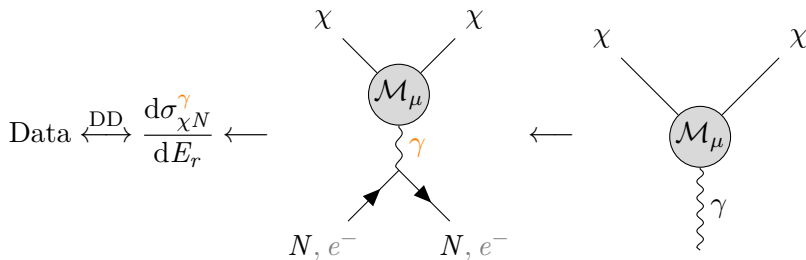
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Experiment

Theory



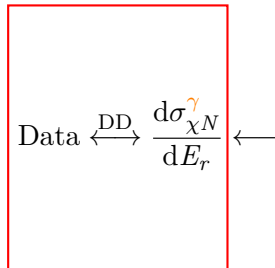
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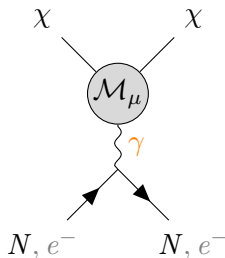
Anapole Moment: The Setup

[see also talk by Xuyang Ning about Panda-X results (DM-1 on Monday)]

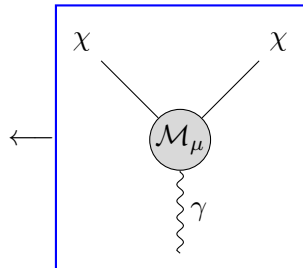
Effective **photon** interaction

Experiment


$$\text{Data} \xleftrightarrow{\text{DD}} \frac{d\sigma_{\chi N}^{\gamma}}{dE_r}$$



Theory



We want to compare **DD limits** with **model predictions**

[See also hep-ph/0003010, 1401.6457, 1503.01500, 1608.00642, 1808.04112, 1806.07896, 2003.04039, ...]

Electromagnetic Vertex Decomposition

- $\mathcal{M}_\mu(q) =$
 $(\gamma_\mu - q_\mu \not{q}/q^2) [f_Q(q^2) + f_A(q^2) q^2 \gamma_5] - i\sigma_{\mu\nu} q^\nu [f_M(q^2) + if_E(q^2) \gamma_5]$

[Broggini, Giunti, Studenikin '12]

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$$f_Q(0) = q \quad \text{charge}$$

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- If χ Majorana: only $\mathcal{A}$ non-vanishing! $\mathcal{L} = \frac{\mathcal{A}}{2} \bar{\chi} \gamma_\mu \gamma_5 \chi \partial_\nu F^{\mu\nu}$

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- $\mathcal{A}$ violates separately P and C

[Broggini, Giunti, Studenikin '12]

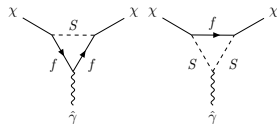
Anapole Moment: Calculation

- χ - Fermion - Scalar interaction:

$$\mathcal{L}_{\text{FFS}} = \bar{\chi} [c_L P_L + c_R P_R] S^* f + \text{h.c.}$$

$\rightarrow \mathcal{A}_S$ scalar contribution

[Kopp et al. '14; Garny et al. '15; Sandick et al. '16]



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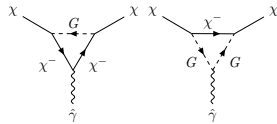
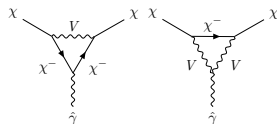
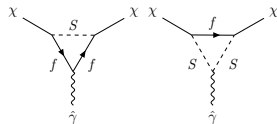
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$$\mathcal{L}_{\text{FFV}} = \bar{\chi} \gamma^\mu [v_L P_L + v_R P_R] \chi^- V_\mu^+ + \text{h.c.}$$

+ Goldstones

→ $\mathcal{A}_V$ vector contribution



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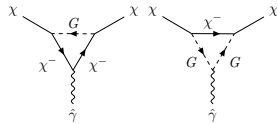
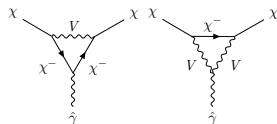
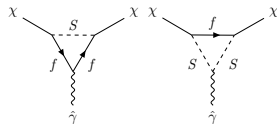
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- Diagrams with vectors → unphysical result!



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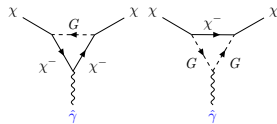
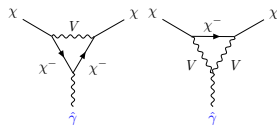
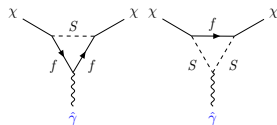
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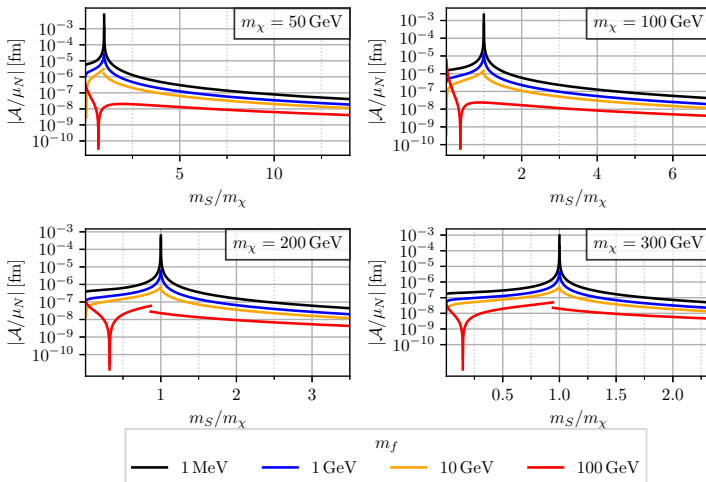
→ $\mathcal{A}_V$ vector contribution

- Diagrams with vectors → unphysical result!
- Pinch Technique/Background Field Method:
 $\gamma \rightarrow \hat{\gamma}, \xi_{\text{BFM}} = 1$



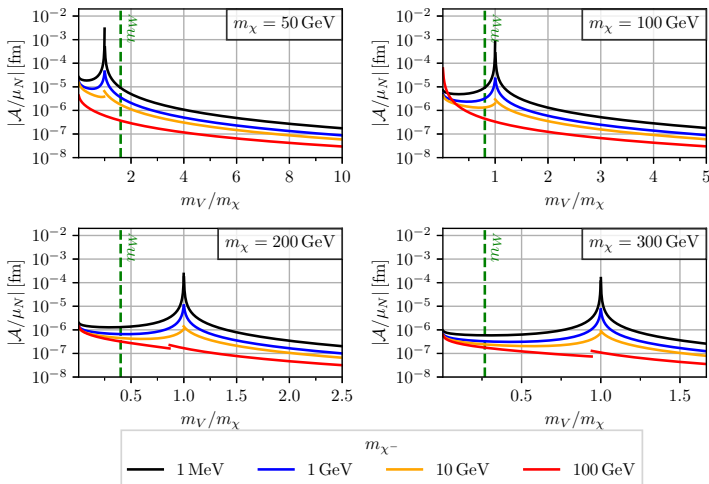
[Works by Cornwall, Papavassiliou, Bernabeu, Rosado, Vidal, Binosi, Cabral-Rosetti, ...
 See review Binosi & Papavassiliou '09]

Anapole Moment: Model-Independent Results (Scalar)



for $c_L = 1$, $c_R = 0$

Anapole Moment: Model-Independent Results (Vector)



for $v_L = 1, v_R = 0$

Model-dependent Goldstone loops are not included here

$$\mathcal{L}_{\text{FFS}} = \bar{\chi} [c_L P_L + c_R P_R] S^* f + \text{h.c.}$$

How is the anapole moment of $\tilde{\chi}_1^0$ generated?

- Reminder: $\mathcal{A}$ breaks parity ($c_L \neq c_R$, $v_L \neq v_R$)
- Interactions w/ sfermions
- Interactions with $\tilde{\chi}^\pm$ for non-trivial mixing

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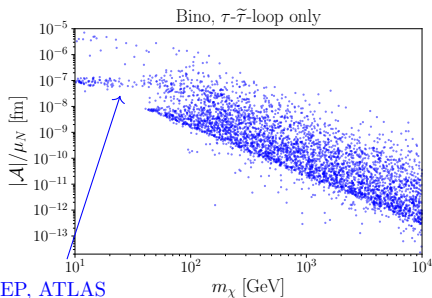
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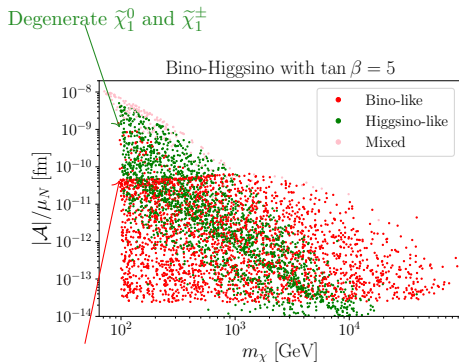
Two possibilities to enhance $\mathcal{A}$:

- 1) Light sfermions: $\mathcal{A}$ generated by scalar contribution $\mathcal{A}_S$
- 2) Mixing in the EW-ino sector: $\mathcal{A}$ generated by vector contribution $\mathcal{A}_V$

MSSM: Two Simplified Scenarios



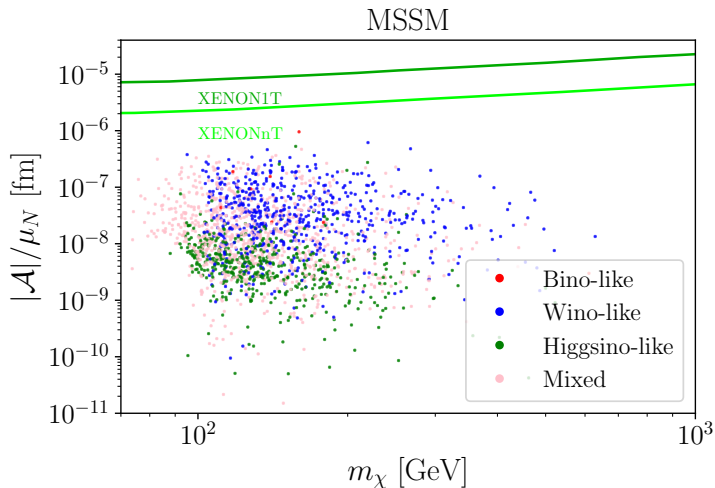
(a) Only $\mathcal{A}_S$



Heavy Chargino

(b) Only $\mathcal{A}_V$

MSSM: Full Scan and DD limits



Constraints from LEP, $\Gamma(Z \rightarrow \text{inv.})$ [micrOMEGAS], Higgs sector [HiggsBounds, HiggsSignals], LHC [SModelS]
and flavour physics [SuperIso, GM2Calc], but **no DM constraints**

Electromagnetic Vertex Decomposition: Dirac

Reminder:

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- Formfactors: charge, anapole, magnetic dipole, electric dipole
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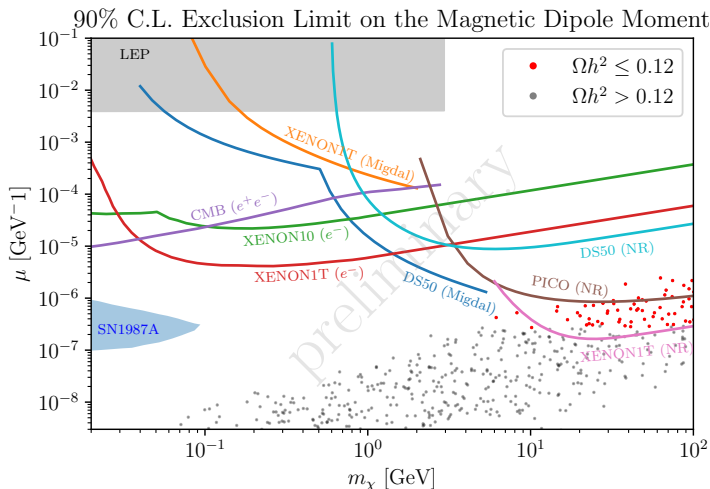
$$f_M(0) = \mu \quad \text{magnetic dipole moment}$$

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- If χ Dirac: dipole operators are allowed $\mathcal{L} = \frac{\mu}{2} \bar{\chi} \sigma_{\mu\nu} \chi F^{\mu\nu}$

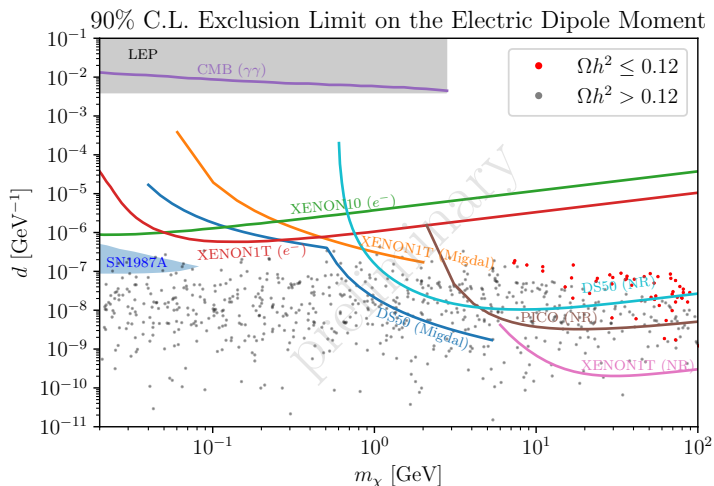
[Broggini, Giunti, Studenikin '12]

The Dirac Case: Magnetic Dipole (work in progress)



[A. Ibarra, **MR**, G. Tomar, *work in progress*]

The Dirac Case: Electric Dipole (work in progress)



[A. Ibarra, **MR**, G. Tomar, *work in progress*]

- We calculated two model-independent contributions: *scalar* $\mathcal{A}_S$ and *vector* $\mathcal{A}_V$
- Naive calculation of $\mathcal{A}_V$ leads to non-physical result $\rightarrow$ BFM
- Simplified Models: $\mathcal{A}_S \sim \mathcal{O}(10^{-6} \mu_N \text{ fm})$ and $\mathcal{A}_V \sim \mathcal{O}(10^{-8} \mu_N \text{ fm})$
- Anapole is not experimentally accessible (yet)
- Different situation for the dipole moments: already experimentally testable

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Thank you for your attention



Backup: Analytical Anapole Expressions

- Scalar:

$$\mathcal{A}_S = -\frac{e}{96\pi^2 m_\chi^2} Q_f \left[|c_L|^2 - |c_R|^2 \right] \mathcal{F}_S\left(\frac{m_f}{m_\chi}, \frac{m_S}{m_\chi}\right)$$

- Vector:

$$\mathcal{A}_V = \frac{e}{96\pi^2 m_\chi^2} \left\{ 2 \left[|v_L|^2 - |v_R|^2 \right] \mathcal{F}_V\left(\frac{m_{\chi^-}}{m_\chi}, \frac{m_V}{m_\chi}\right) + \left[|c_L^G|^2 - |c_R^G|^2 \right] \mathcal{F}_S\left(\frac{m_{\chi^-}}{m_\chi}, \frac{m_V}{m_\chi}\right) \right\}$$

- Loop functions:

$$\mathcal{F}_X(\mu, \eta) = \frac{3}{2} \log\left(\frac{\mu^2}{\eta^2}\right) + (3\eta^2 - 3\mu^2 + n_X)f(\mu, \eta),$$

with $n_V = -7$, $n_S = 1$

Backup: Loop Functions

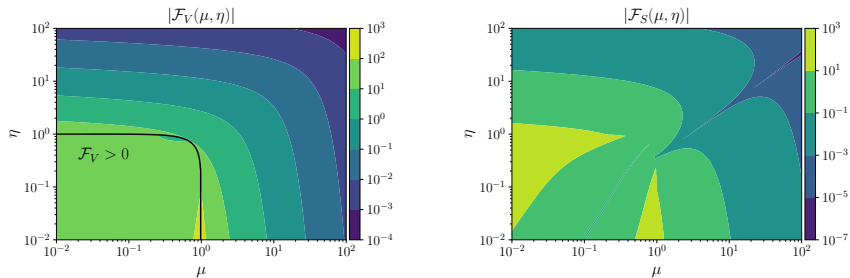
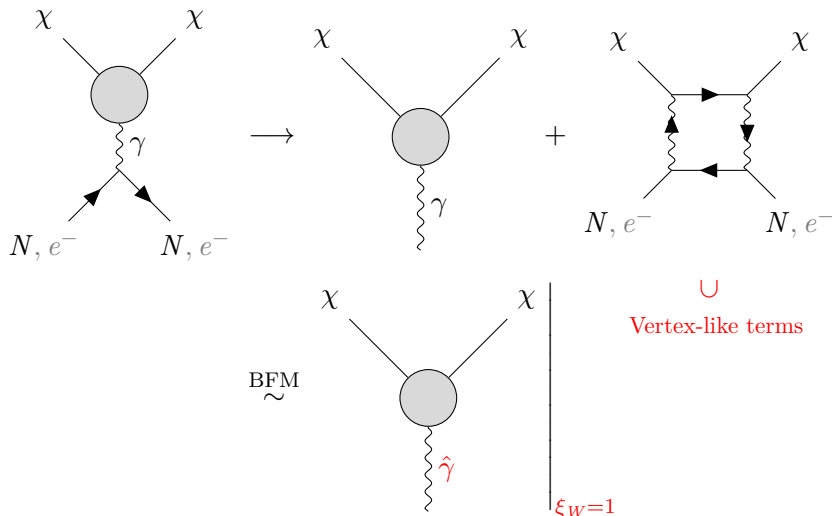


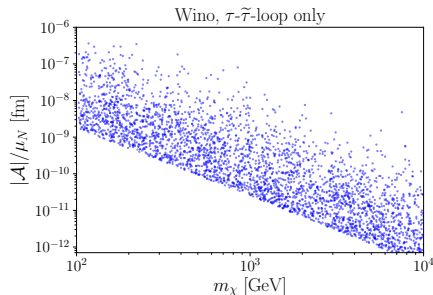
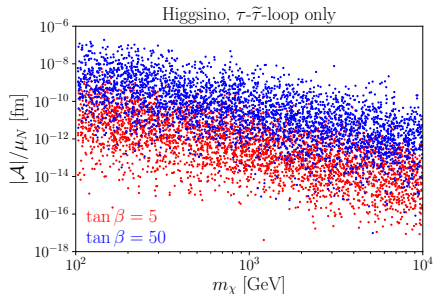
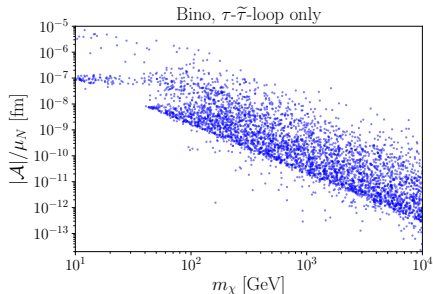
Figure: Contour plot of absolute values of $\mathcal{F}_V(\mu, \eta)$ (left) and $\mathcal{F}_S(\mu, \eta)$ (right).

Backup: BFM and Pinch Technique

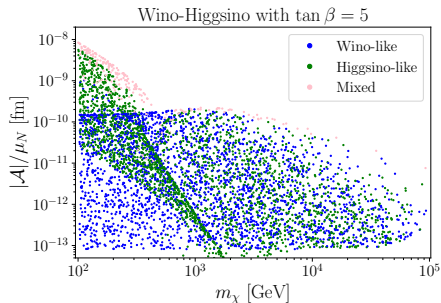
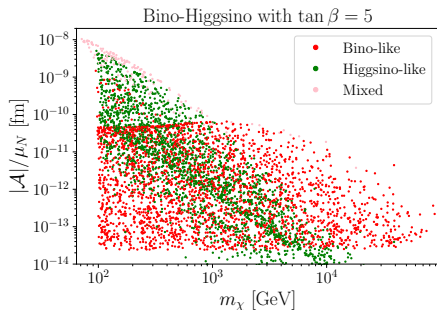


[Hasimoto, Kodaira, Yasui, Sasaki '94; Papavassiliou '94; Denner, Dittmaier, Weiglein '94; Review by Binosi & Papavassiliou '09]

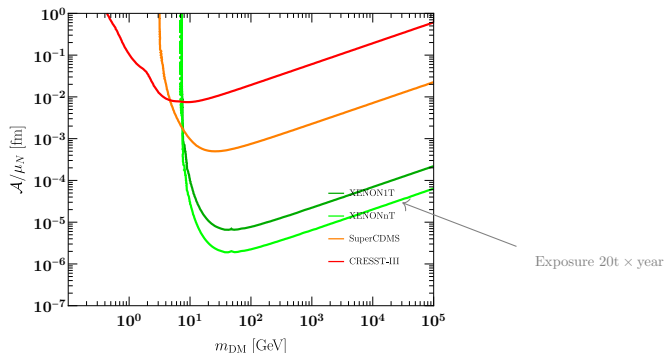
Backup: Pure EW-inos & light sfermions



Backup: Mixed EW-inos & heavy sfermions

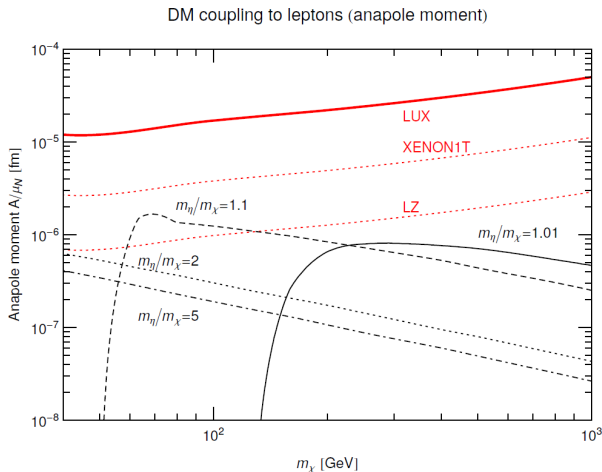


Backup: Scattering Rate for the Anapole Interaction



$$\frac{d\sigma}{dE_R} = \alpha_{EM} \mathcal{A}^2 \left[Z^2 \left(2m_T - \left(1 + \frac{m_T}{m_\chi} \right)^2 \frac{E_R}{v^2} \right) F_Z^2(q^2) + \frac{1}{3} \frac{m_T}{m_\chi^2} \left(\frac{\bar{\mu}_T}{\mu_N} \right)^2 \frac{E_R}{v^2} F_D^2(q^2) \right]$$

Relic density in simplified model

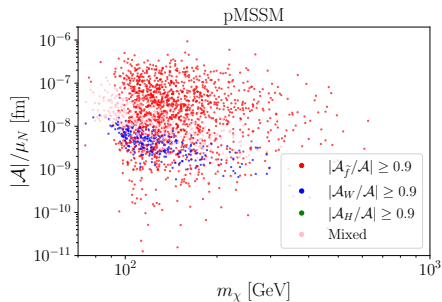
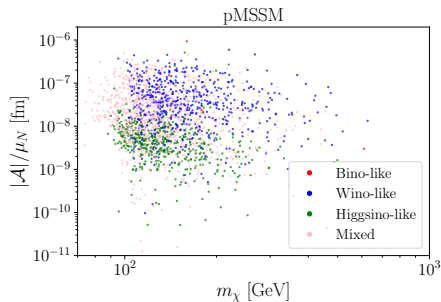


Garny et al. '15

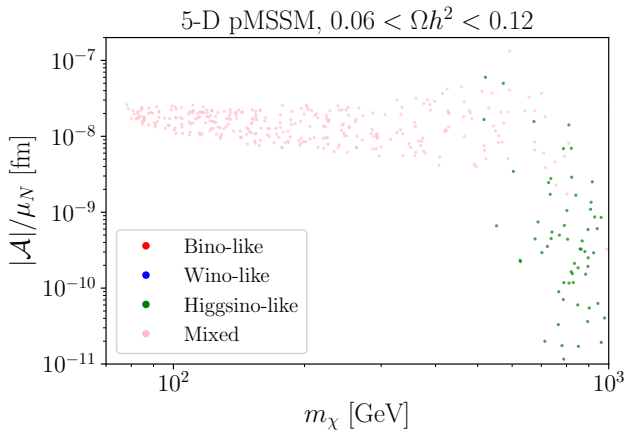
Backup: pMSSM Parameter Space

Parameter	Full	5-D
M_1	[100, 2000] GeV	[100, 2000] GeV
M_2	[100, 2000] GeV	[100, 2000] GeV
M_3	[2000, 5000] GeV	3000 GeV
$A_{t,b,\tau}$	[-4000, 4000] GeV	4000, 0, 0 GeV
m_A	$[10^3, 10^5]$ GeV	5000 GeV
$\tan \beta$	[3, 50]	50
μ	[100, 2000] GeV	[100, 2000] GeV
$m_{\tilde{\ell}_{L,R}}$	[100, 2000] GeV	[100, 1000], 3000 ($\tilde{\tau}$) GeV
$m_{\tilde{q}_{L1,2}}$	[400, 2000] GeV	5×10^4 GeV
$m_{\tilde{u}_{R1,2}}, m_{\tilde{d}_{R1,2}}$	[400, 2000] GeV	5×10^4 GeV
$m_{\tilde{q}_{L3}}$	[300, 2000] GeV	10^4 GeV
$m_{\tilde{u}_{R3}}, m_{\tilde{d}_{R3}}$ GeV	[300, 2000] GeV	10^4 GeV

Backup: MSSM Full Scan



Backup: 5-D pMSSM scan including $\Omega_\chi h^2$



Backup: MSSM Couplings I

To Charginos

$$v_L^j = -g N_{12} U_{j1}^* - g \frac{1}{\sqrt{2}} N_{13} U_{j2}^*, \quad (1a)$$

$$v_R^j = -g N_{12}^* V_{j1} + g \frac{1}{\sqrt{2}} N_{14}^* V_{j2}, \quad (1b)$$

$$c_L^{G,j} = g \cos \beta \left[N_{13}^* U_{j1}^* - \frac{1}{\sqrt{2}} U_{j2}^* (N_{12}^* + \tan \theta_W N_{11}^*) \right], \quad (1c)$$

$$c_R^{G,j} = -g \sin \beta \left[N_{14} V_{j1} + \frac{1}{\sqrt{2}} V_{j2} (N_{12} + \tan \theta_W N_{11}) \right], \quad (1d)$$

$$c_L^{H,j} = -g \sin \beta \left[N_{13}^* U_{j1}^* - \frac{1}{\sqrt{2}} U_{j2}^* (N_{12}^* + \tan \theta_W N_{11}^*) \right], \quad (2a)$$

$$c_R^{H,j} = -g \cos \beta \left[N_{14} V_{j1} + \frac{1}{\sqrt{2}} V_{j2} (N_{12} + \tan \theta_W N_{11}) \right]. \quad (2b)$$

Backup: MSSM Couplings II

To Sfermion/fermions

$$c_L^{i,1} = G^{f_{iL}} \cos \theta_{\tilde{f}_a} + H^{f_{iR}} \sin \theta_{\tilde{f}_a}, \quad (3a)$$

$$c_R^{i,1} = G^{f_{iR}} \sin \theta_{\tilde{f}_a} + H^{f_{iL}} \cos \theta_{\tilde{f}_a}, \quad (3b)$$

$$c_L^{i,2} = -G^{f_{iL}} \sin \theta_{\tilde{f}_a} + H^{f_{iR}} \cos \theta_{\tilde{f}_a}, \quad (3c)$$

$$c_R^{i,2} = G^{f_{iR}} \cos \theta_{\tilde{f}_a} - H^{f_{iL}} \sin \theta_{\tilde{f}_a}, \quad (3d)$$

with

$$G^{f_{iL}} = -\sqrt{2}g \left[T_{3L}^{f_i} N_{12}^* + \tan \theta_W (Q_{f_i} - T_{3L}^{f_i}) N_{11}^* \right], \quad (4a)$$

$$G^{f_{iR}} = \sqrt{2}g \tan \theta_W Q_{f_i} N_{11}, \quad (4b)$$

$$H^{f_{iL}} = -\frac{g}{\sqrt{2}m_W} m_{f_i} \times \begin{cases} N_{14}/\sin \beta, & f_i = u\text{-type} \\ N_{13}/\cos \beta, & f_i = d\text{-type}, \ell \end{cases} \quad (4c)$$

$$H^{f_{iR}} = H^{f_{iL}*}. \quad (4d)$$