

# Probing axions with the measurements of photon's birefringence

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ダークマターの正体は何か？

広大なディスカバリースペースの網羅的研究

文部科学省  
科学研究費助成事業  
学術変革領域研究  
(2020-2024)

What is dark matter? - Comprehensive study of the huge discovery space in dark matter

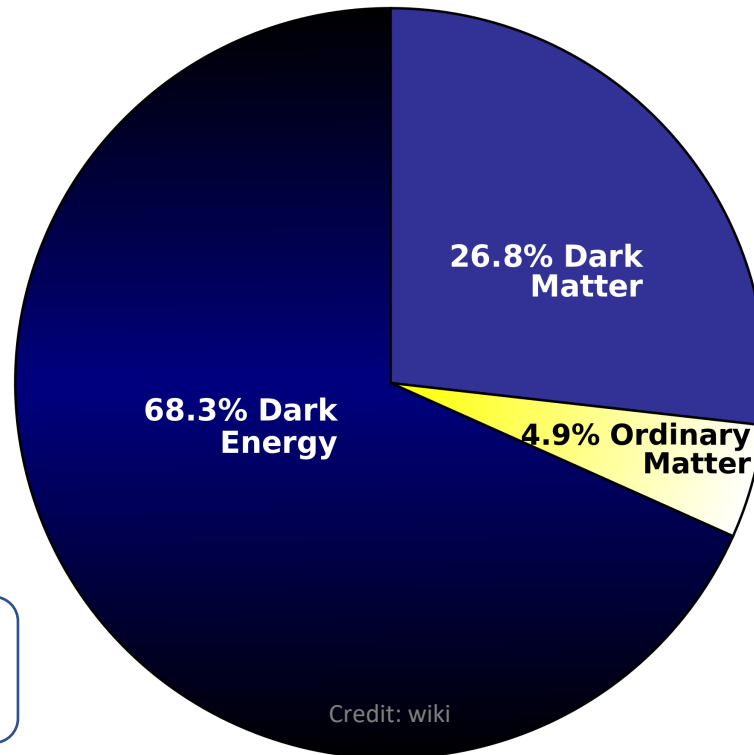


# Dark sectors in our universe



Credit: higgstan.com

Cosmological Constant?  
Quintessence?



Credit: higgstan.com

WIMP?  
Axion?  
MACHO?  
PBHs?

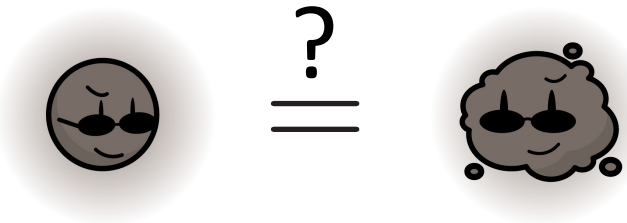
(standard model)

We know only 5 % in our universe!

# Question

Are these sectors independent components?

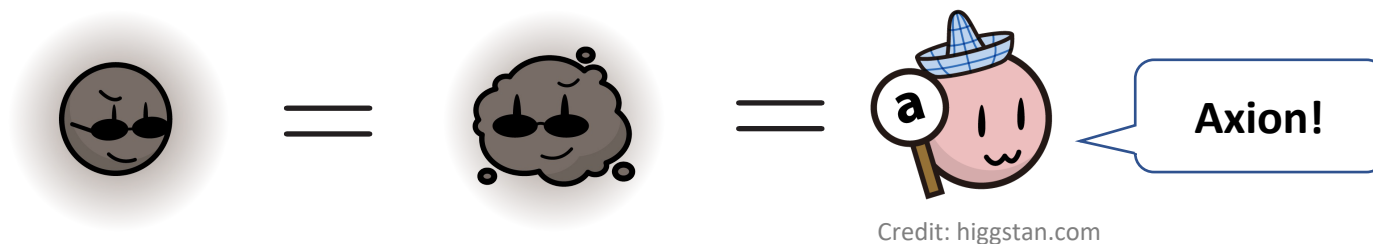
What if they are related by a common origin?



**Unified models for dark matter and dark energy have been developed**

- **Generalized Chaplygin gas:** Bento, Bertolami & Sen (2002); Makler, Oliveira, Waga (2003); ...
- **k-essence:** Scherrer (2004); Giannakis & Hu (2005); ...
- **Fast transition models:** Bruni, Lazkoz & Fernandez (2013); Leanizbarrutia, Fernandez & Tereno (2017); ...
- ...

# Motivation



Axion could be a candidate for both dark sectors!

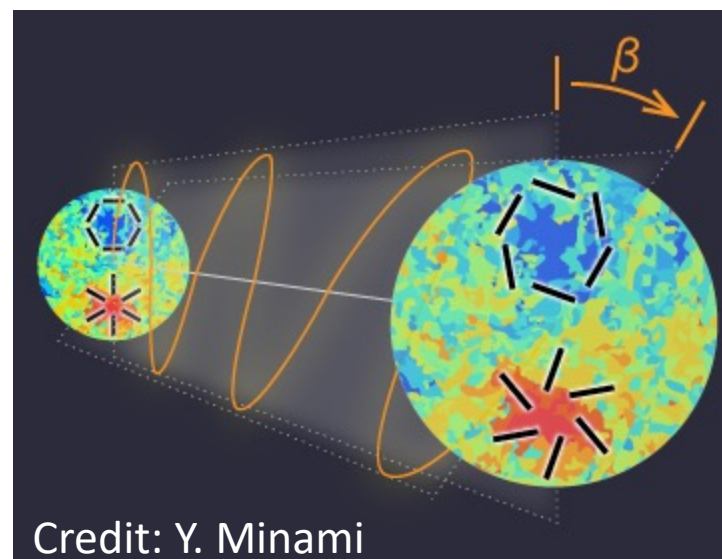
## Overview of talk

### Background

Introduce what is axion? how to search for it?, mainly focusing on the search for **axion-photon conversions**

### Our study

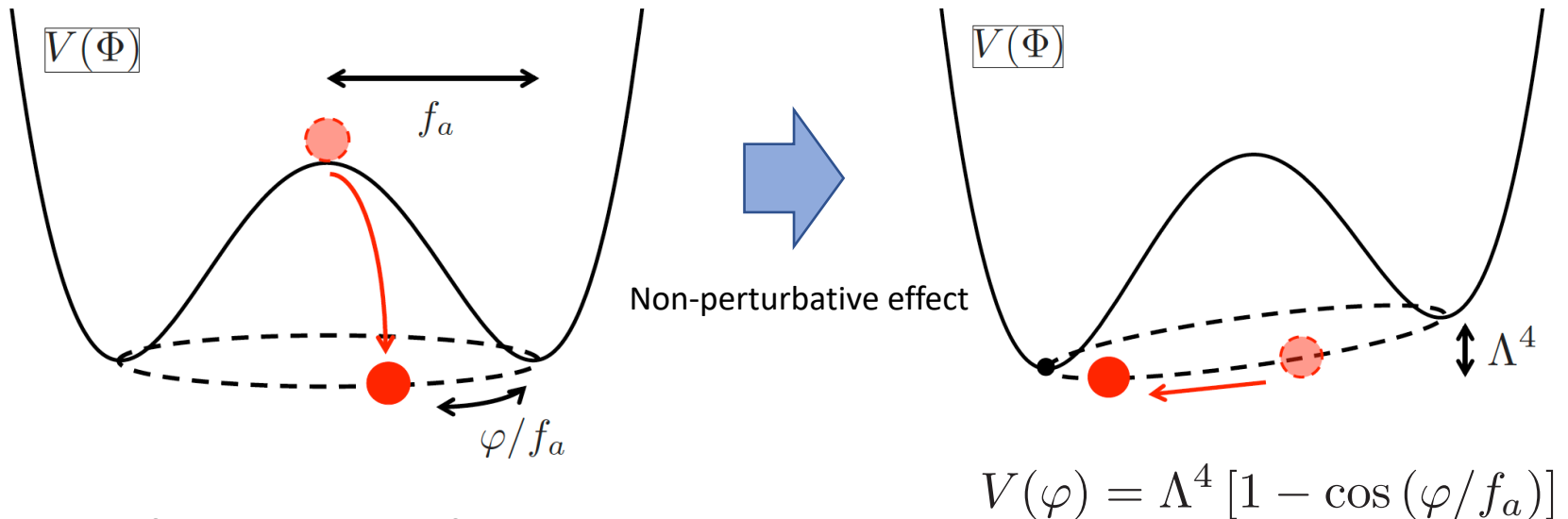
Introduce a new search method using the **photon's birefringence effect** and corresponding observational progress



# What is axion?



A (pseudo) Nambu-Goldstone boson of global  $U(1)$  symmetry



**Two characteristic scales:**

$f_a$  : decay constant

$m_a = \Lambda^2/f_a$  : mass

# When people say axion...

QCD axion *Peccei & Quinn (1977); Weinberg, Wilczek (1978); ...*

- Suggested to solve strong CP problem
- Mass & decay constant are related to each other in QCD energy scale

$$m_a = \frac{\Lambda_{\text{QCD}}^2}{f_a} \simeq 6 \mu\text{eV} \left( \frac{10^{12} \text{GeV}}{f_a} \right) \quad (\text{typical}) \text{ QCD axion window}$$
$$10^9 \text{ GeV} \lesssim f_a \lesssim 10^{12} \text{ GeV}$$

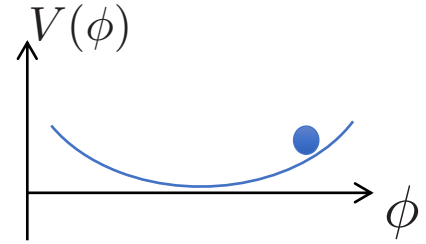
Axion-like particle (ALP) *Svrcek & Witten (2006); Arvanitaki+ (2010); ...*

- Predicted by theories beyond the standard model (e.g. string theory)
- Mass & decay constant are treated as independent parameters

↑ Main focus of this talk

# Cosmological axion evolution

$$\ddot{\phi} + 3H\dot{\phi} + m^2\phi = 0 \quad (\text{background evolution})$$



■ In early universe ( $m < H$ ),  $\phi \simeq \text{const.}$  (frozen due to the Hubble friction)

■ In late universe ( $m > H$ ),  $\phi \simeq a^{-3/2}\phi_0 \cos(mt)$  (start to oscillate)

■ After oscillation begins, axion behaves as a pressureless matter fluid

$$\rho = \frac{1}{2}\dot{\phi}^2 + \frac{1}{2}m^2\phi^2 \simeq \frac{\rho_0}{a^3}$$

$$P = \frac{1}{2}\dot{\phi}^2 - \frac{1}{2}m^2\phi^2 \simeq \frac{P_0}{a^3} \sin(2mt) \sim 0$$

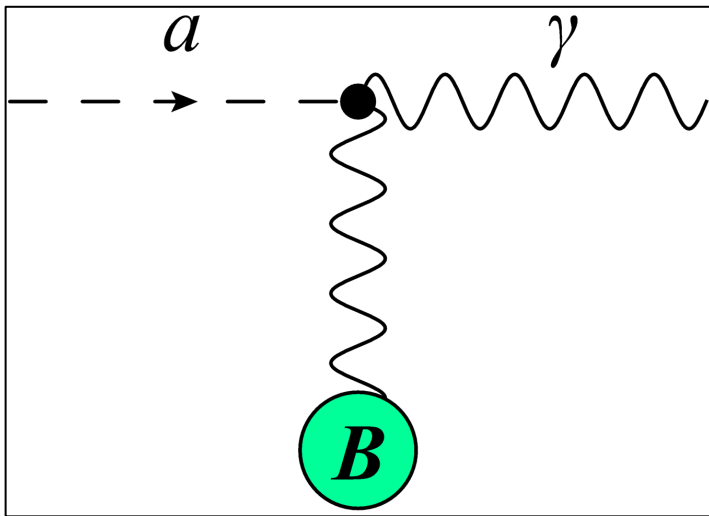
(pressure) < (gravity)  
→ could behave as dark matter

■ If  $m \lesssim H_0 \sim 10^{-33} \text{ eV} \rightarrow$  could behave as dark energy (quintessence)

# Conventional Axion Search

- Axion generically couples to photon via the topological term

$$\frac{g_{a\gamma}}{4} a F_{\mu\nu} \tilde{F}^{\mu\nu} = g_{a\gamma} a \mathbf{E} \cdot \mathbf{B}$$



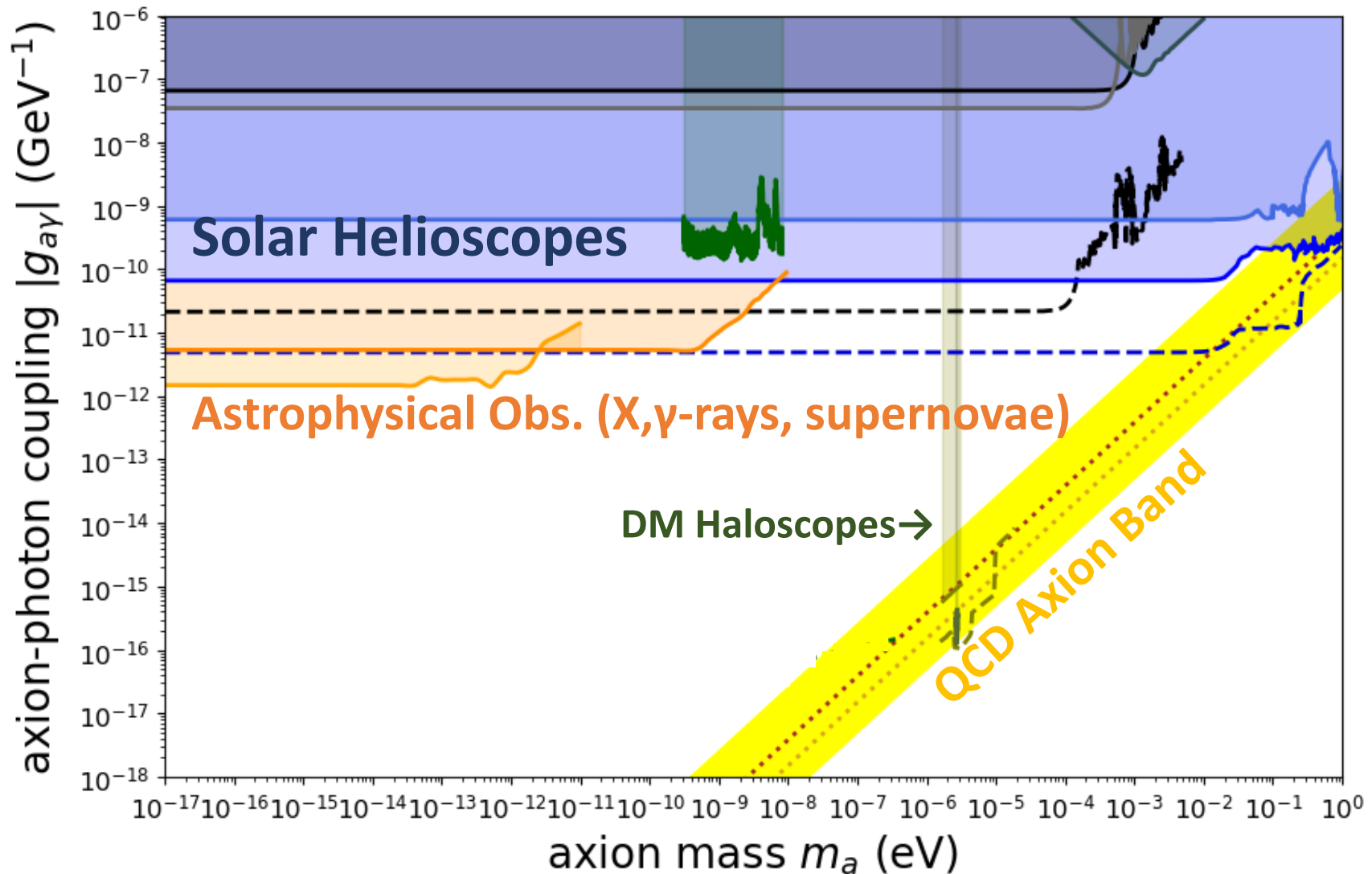
$a$  : axion

$\gamma$  : photon

$B$  : magnetic field

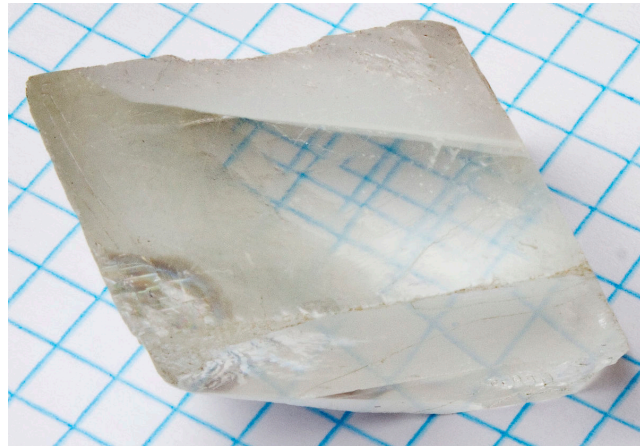
- Axion is converted into photon under the background magnetic field (“axion-photon conversion” or “Primakoff effect”)

# Overview of Target Spaces



# New search methods for axions

- Photon's birefringence effect-



Axion as a birefringent material

# Photon's birefringence by axion

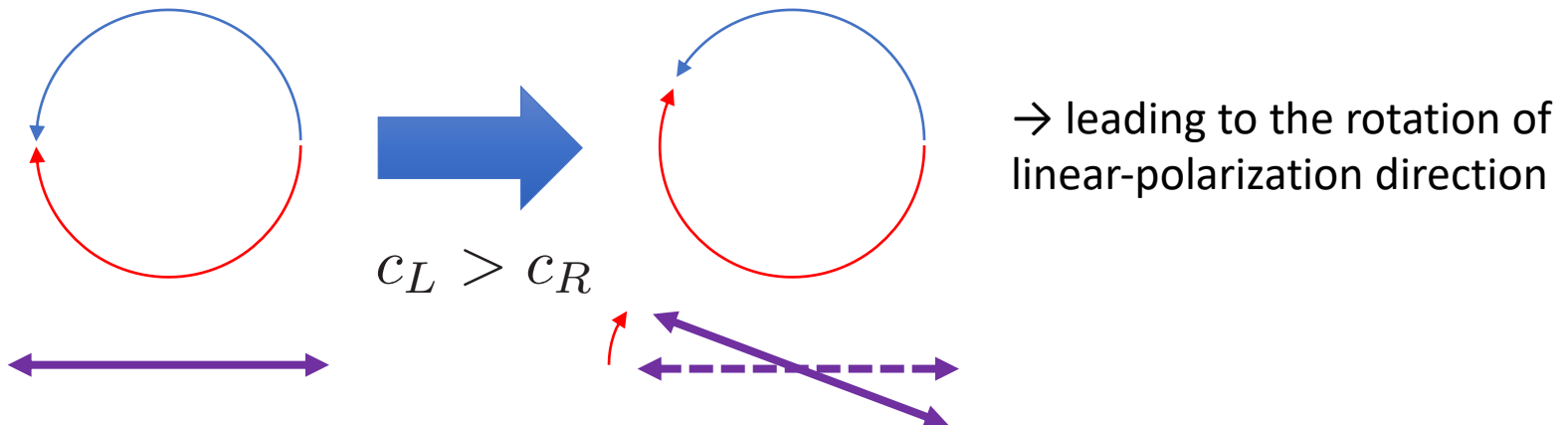
*Carroll, Field & Jackiw (1990); Harari & Sikivie (1992); Carroll (1998); ...*

**Axion behaves as a birefringent material in our universe**

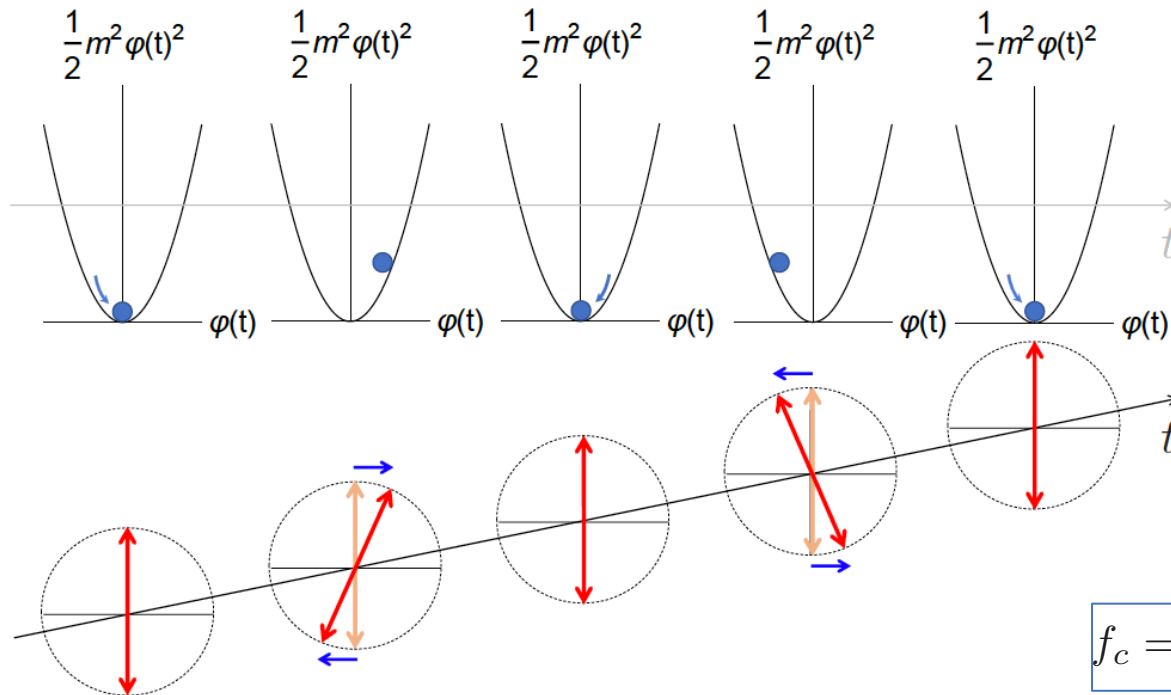
- Axion differentiates the phase velocities of circular-polarized photon

$$\mathcal{L} \supset \frac{1}{4} g_{a\gamma} \varphi F_{\mu\nu} \tilde{F}^{\mu\nu}, \quad F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$$

$$\text{Dispersion relation: } \ddot{A}_k^{\text{L/R}} + \omega_{\text{L/R}}^2 A_k^{\text{L/R}} = 0, \quad c_{\text{L/R}} \equiv \frac{\omega_{\text{L/R}}}{k} = \sqrt{1 \pm \frac{g_{a\gamma} \dot{\varphi}}{k}}$$



# Birefringence by axion DM

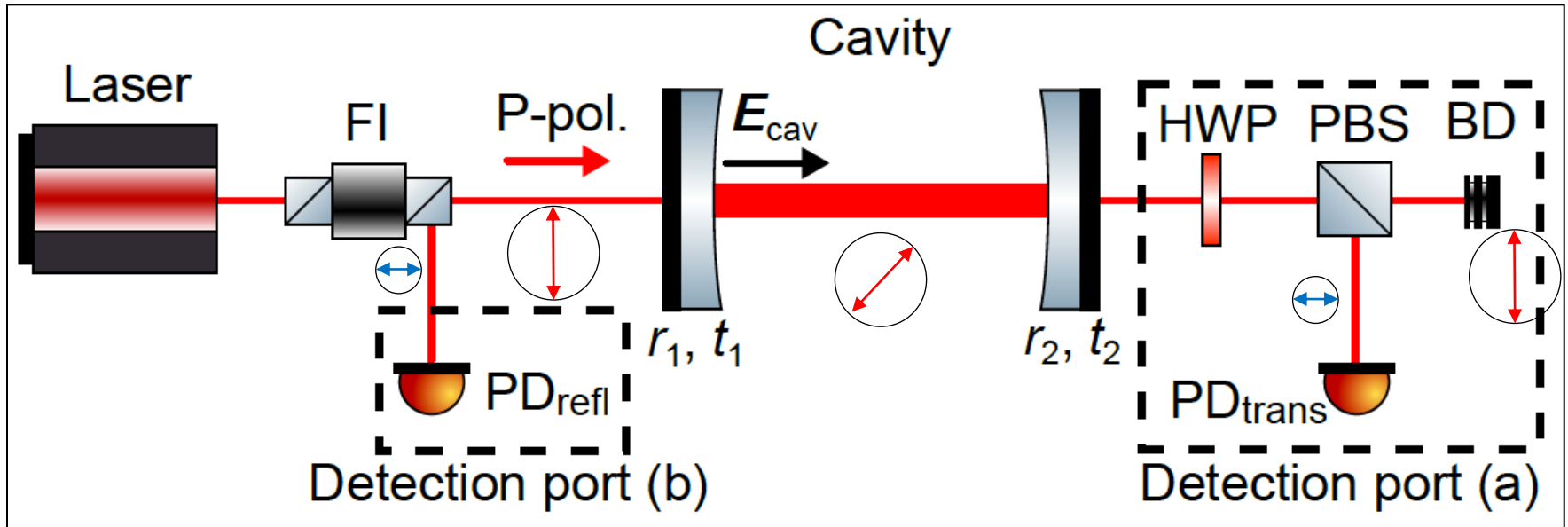


$$f_c = \frac{m_a}{2\pi} \simeq 2.4 \text{Hz} \left( \frac{m_a}{10^{-14} \text{eV}} \right)$$

- Axion DM induces the polarization rotation **oscillating in time** with a frequency of axion mass:
- Possible to observe by several experimental/astrophysical approaches!

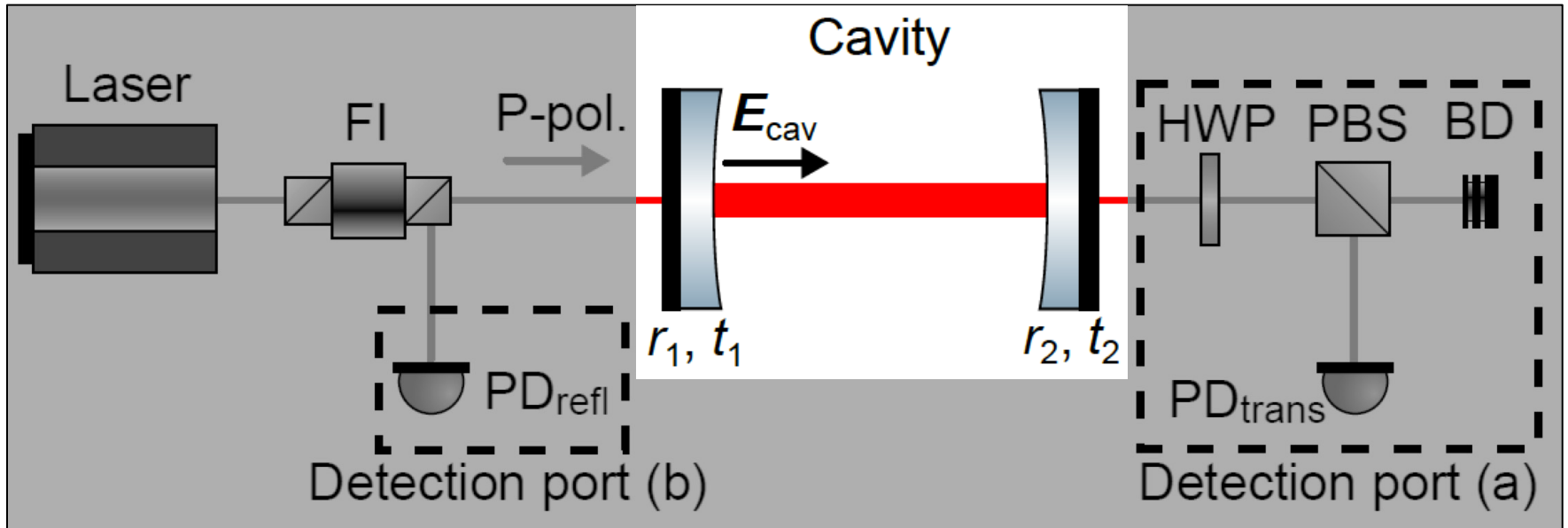
# Birefringence in Resonant Cavity

*Nagano, Fujita, Michimura, IO (2019);...*



- As a carrier wave, we input the linearly-polarized monochromatic laser light.
- The linear cavity consists of front and end mirrors. The cavity is kept to resonate with a phase measurement.
- The signal is detected in detection port (a) or (b) as polarization modulation.

# Inside Cavity



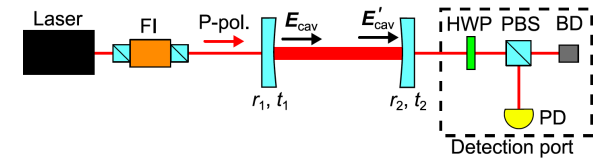
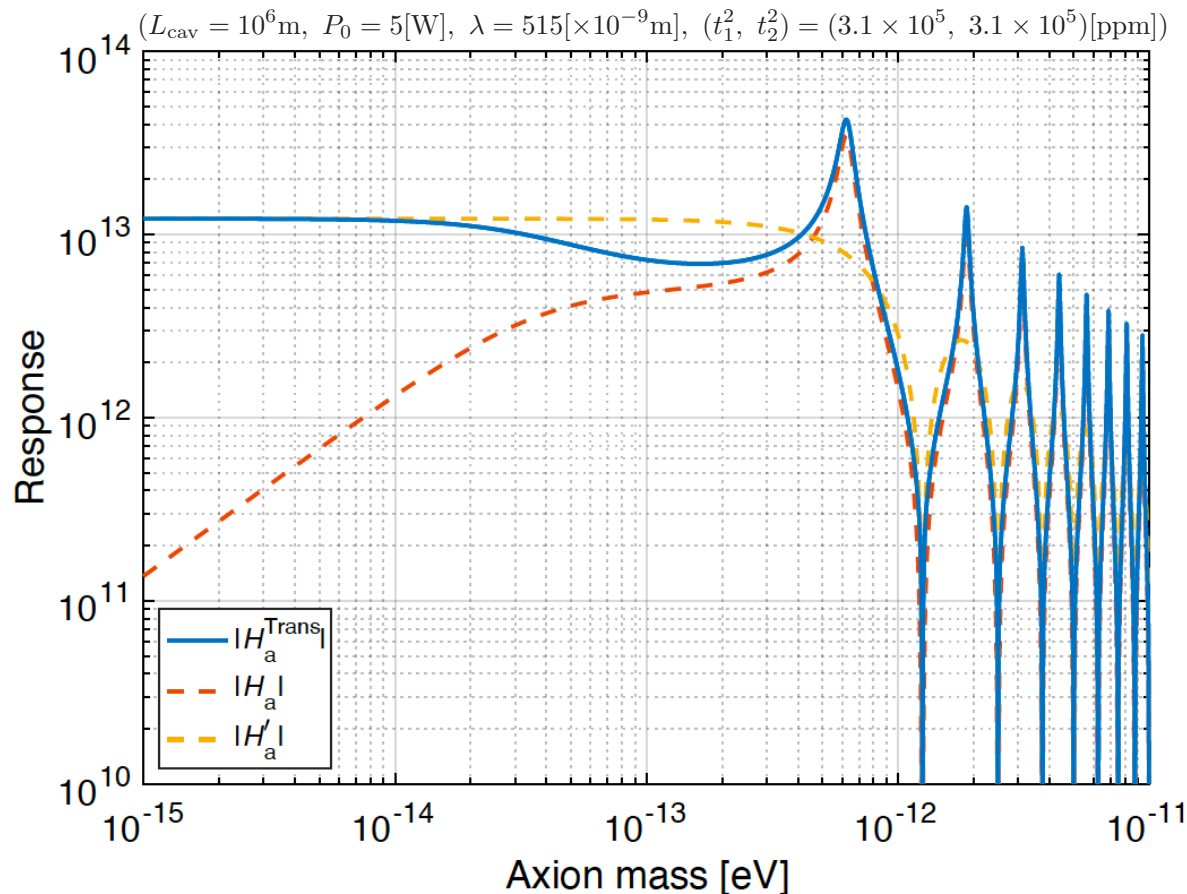
- When the resonant condition is met, the beam is accumulated inside cavity:

$$\begin{aligned}
 \mathbf{E}_{\text{cav}}(t) &= \frac{t_1 E_0 e^{ikt}}{1 - r_1 r_2} \begin{pmatrix} \mathbf{e}^L & \mathbf{e}^R \end{pmatrix} \begin{pmatrix} 1 + i\delta\phi(t) & 0 \\ 0 & 1 - i\delta\phi(t) \end{pmatrix} \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} \\
 &= \frac{t_1}{1 - r_1 r_2} [\mathbf{E}^p(t) - \delta\phi \mathbf{E}^s(t)] \quad (2kL = 2\pi\mathbb{N}) \\
 &\gg 1 \quad (\because r_i \simeq 1)
 \end{aligned}$$

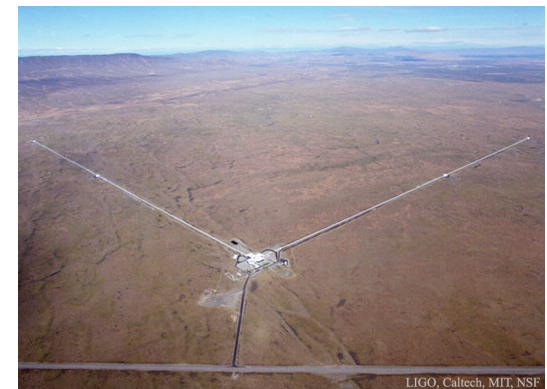
$$r_i^2 + t_i^2 = 1$$

# Frequency response in signal

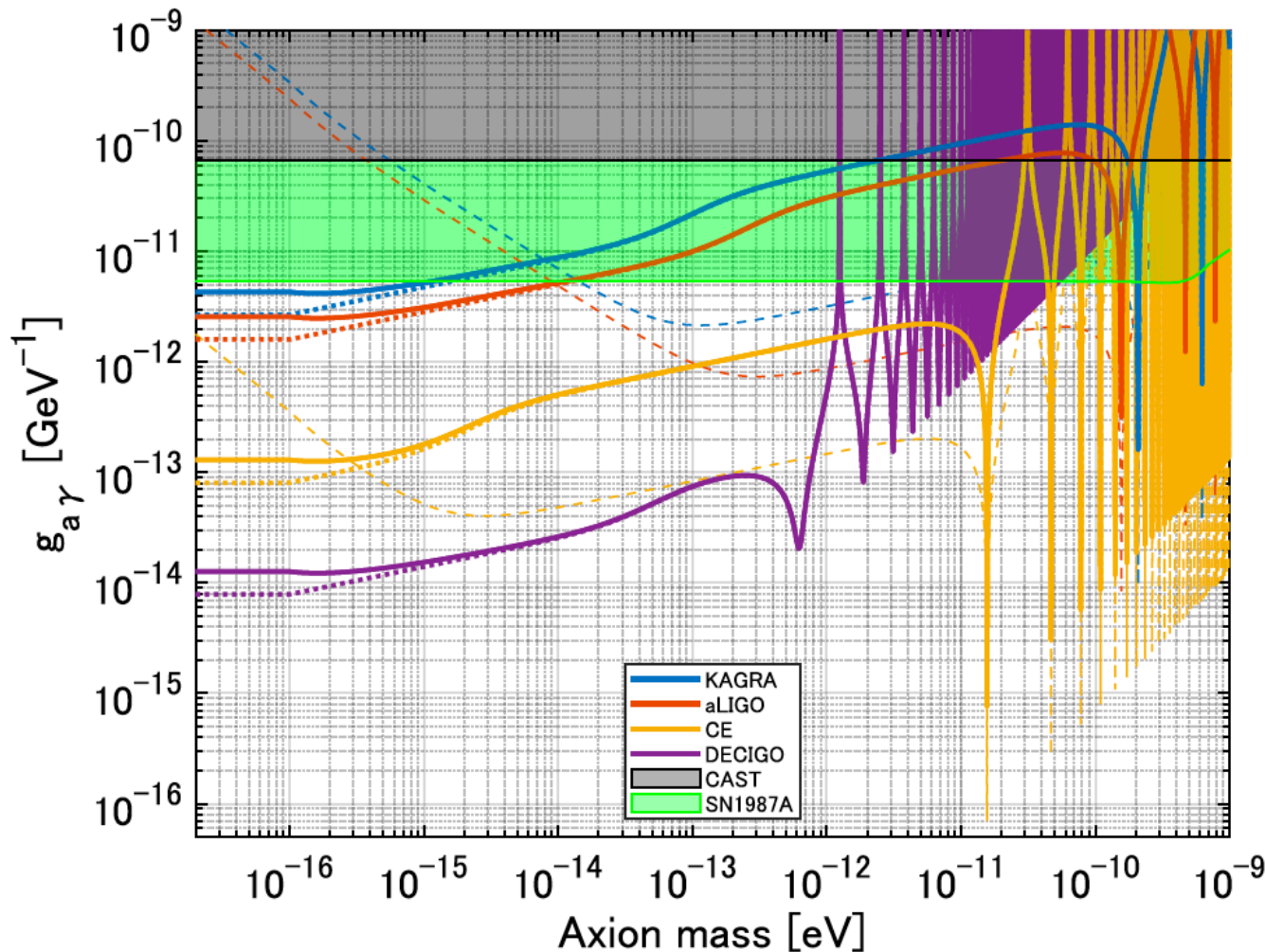
$$\delta\phi(t) = \int_{-\infty}^{\infty} \delta c(m) H_a(m) e^{imt} \frac{dm}{2\pi} \quad H_a(m) : \text{response function}$$



Longer cavity length  
&  
Higher mirror quality  
is better  
→ **GW detector!**



# Axion DM search with GW detector



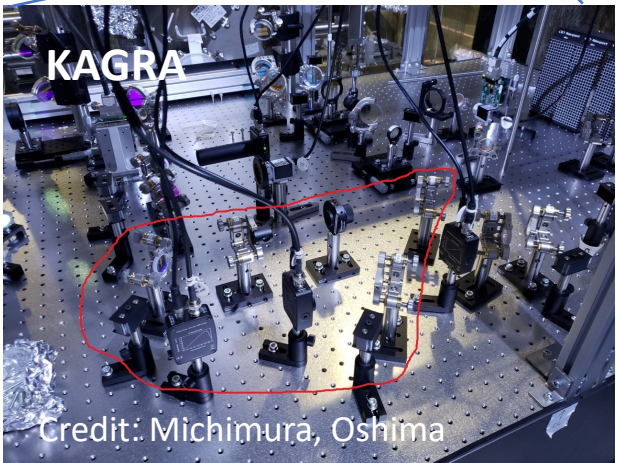
■ Assuming

- Axion is 100% in DM
- 1-year observation
- shot noise is primary

■ sensitive to

$$10^{-17} \text{eV} \lesssim m_a \lesssim 10^{-11} \text{eV}$$

(~mHz)                      (~kHz)



# Optical cavity experiment

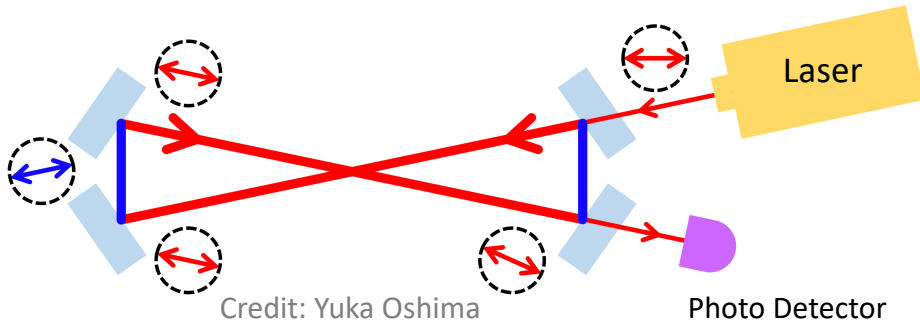
**DANCE:** Dark matter Axion search with ring Cavity Experiment



*Ando, Fujimoto, Fujita, Kume, Michimura,  
Morisaki, Nagano, Nishizawa, IO, Oshima,  
Wang*

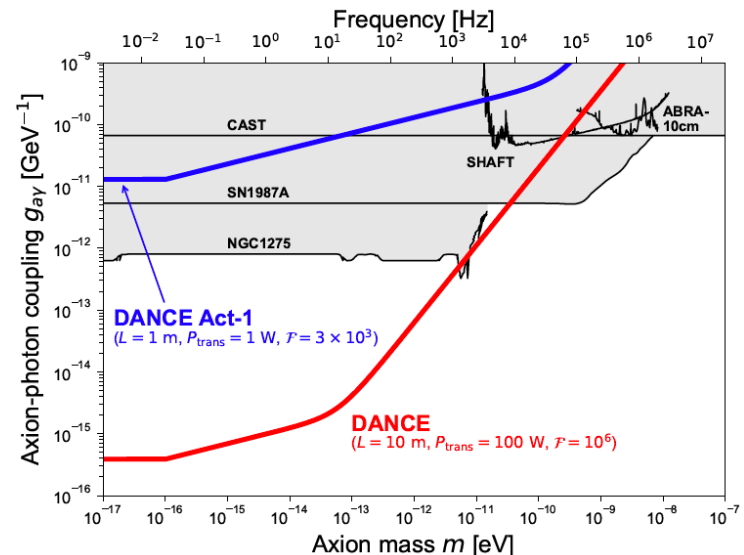
@Ando Lab. (Hongo Campus)

*IO, Fujita, Michimura (2018);...*



- A bowtie-like mirror configuration can sustain the rotational orientation of photon polarization

Goal of sensitivity curves



# Optical cavity experiment

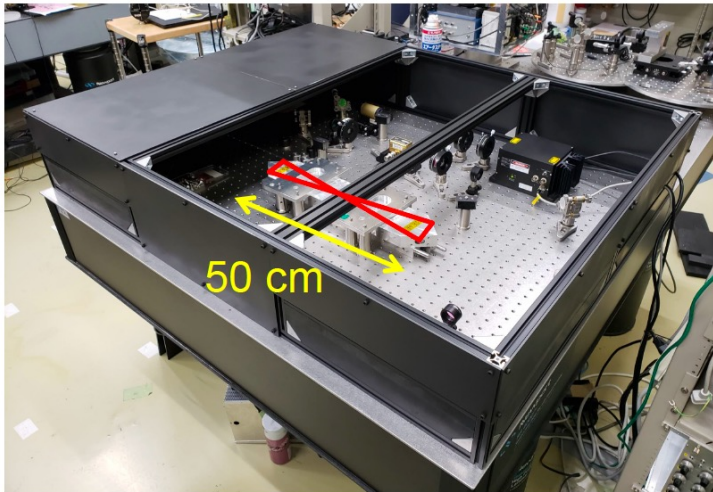
**DANCE:** Dark matter Axion search with riNg Cavity Experiment



*Ando, Fujimoto, Fujita, Kume, Michimura,  
Morisaki, Nagano, Nishizawa, IO, Oshima,  
Wang*

@Ando Lab. (Hongo Campus)

## DANCE Act-1 (prototype experiment)

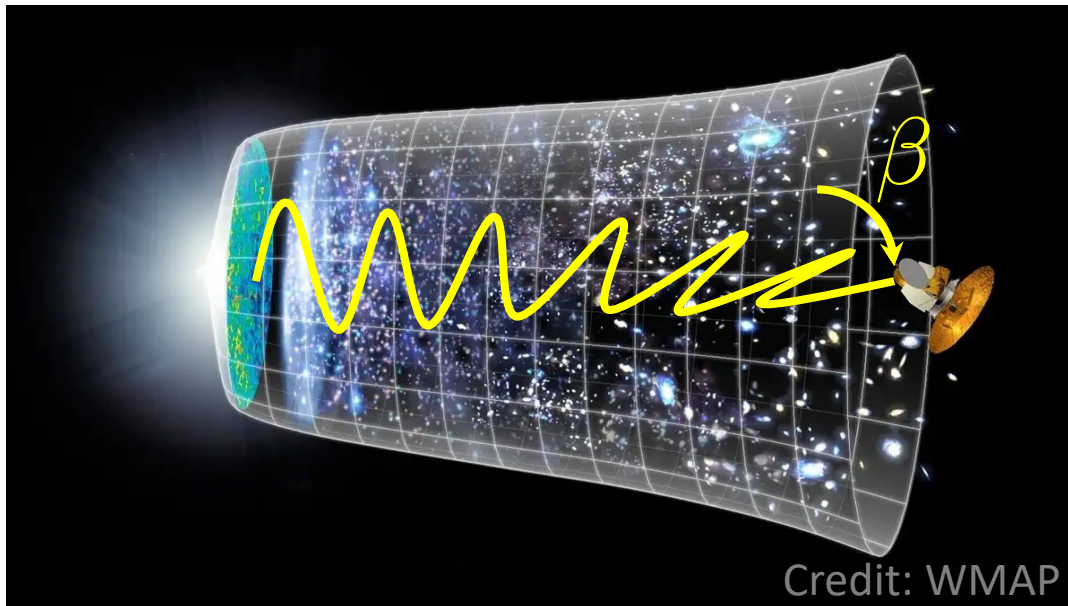
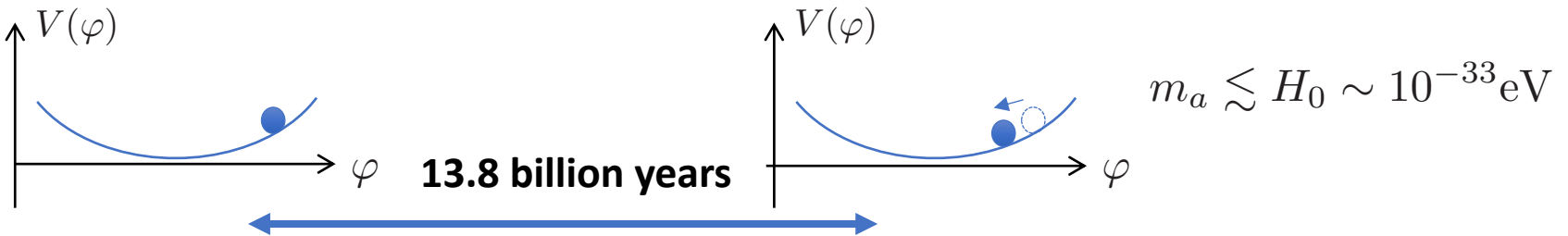


- ✓ First observation in May 2021
- ✓ Constructing further data analysis environment

# CMB Birefringence by axion DE

(Fukugita & Yanagida (1994); Friemann+ (1995); J.E.Kim+ (1999); ...)

- Axion with mass smaller than current Hubble scale behaves as a dark energy



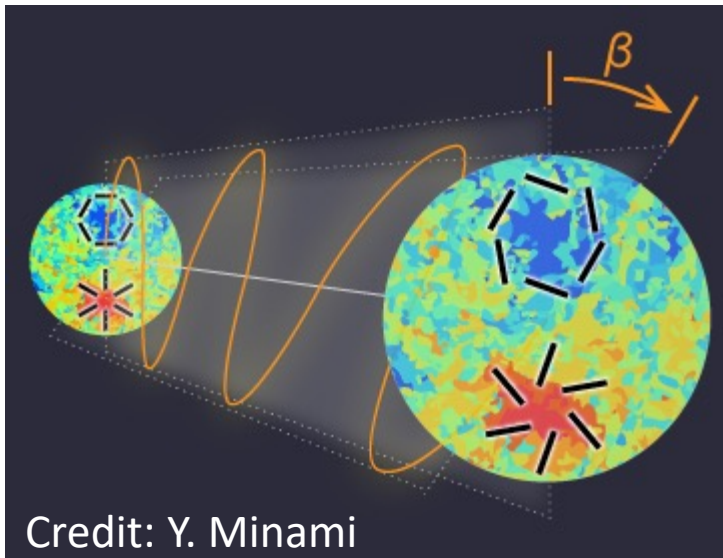
- If axion is responsible for dark energy, it makes the polarization plane of CMB rotate from the last-scattering-surface

Rotation angle:

$$\beta = \frac{g_{\phi\gamma}}{2} \Delta\phi \equiv \frac{g_{\phi\gamma}}{2} (\phi_0 - \langle\phi_{\text{LSS}}\rangle)$$

# Generation EB correlation function

*Lue, Wang & Kamionkowski (1999); Feng+ (2005,2006); Liu, Lee & Ng (2006); ...*



Parity-violating interaction

$$\text{e.g. } \mathcal{L}_{\text{int}} = \frac{1}{4} g_{a\gamma} \varphi F_{\mu\nu} \tilde{F}^{\mu\nu}$$

produces the parity-odd EB correlation

$$C_{\ell}^{EB,o} = \frac{1}{2} \sin(4\beta) \left( C_{\ell}^{EE,\text{CMB}} - C_{\ell}^{BB,\text{CMB}} \right) + \cos(4\beta) C_{\ell}^{EB,\text{CMB}}$$

↑ measured value

↑ usually assume 0

**History of measurements (WMAP, Planck, ACT,...)**

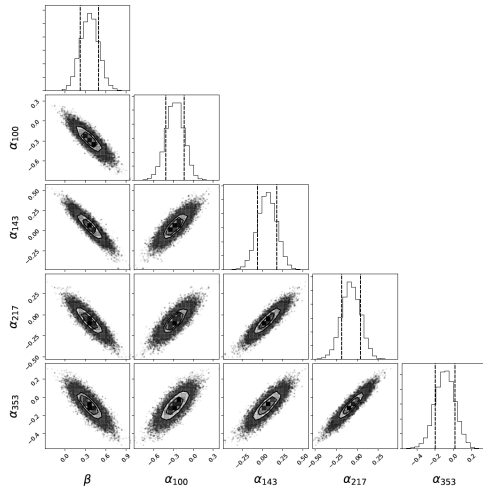
Non-zero  $\langle EB \rangle$  has been detected.

But, not reliable estimates due to the miscalibration of instrumental angle “ $\alpha$ ” .

# Foreground-based calibration

**Minami & Komatsu (2020);**

calibrate  $\alpha$  by using the polarized emission from the galactic foregrounds and measures the intrinsic birefringence angle  $\beta$  by Planck 2018 PR3 data



Angles	Results (deg)
$\beta$	$0.35 \pm 0.14$
$\alpha_{100}$	$-0.28 \pm 0.13$
$\alpha_{143}$	$0.07 \pm 0.12$
$\alpha_{217}$	$-0.07 \pm 0.11$
$\alpha_{353}$	$-0.09 \pm 0.11$

$$\beta = 0.35 \pm 0.14 \text{ deg } (2.4\sigma)$$

**Upgrade! (with Planck PR4 data):**

$$\beta = 0.36 \pm 0.11 \text{ deg } (3.3\sigma)$$

**Upgrade! (with Planck + WMAP data):**

$$\beta = 0.34 \pm 0.09 \text{ deg } (3.6\sigma)$$

# CMB Birefringence by axion DE

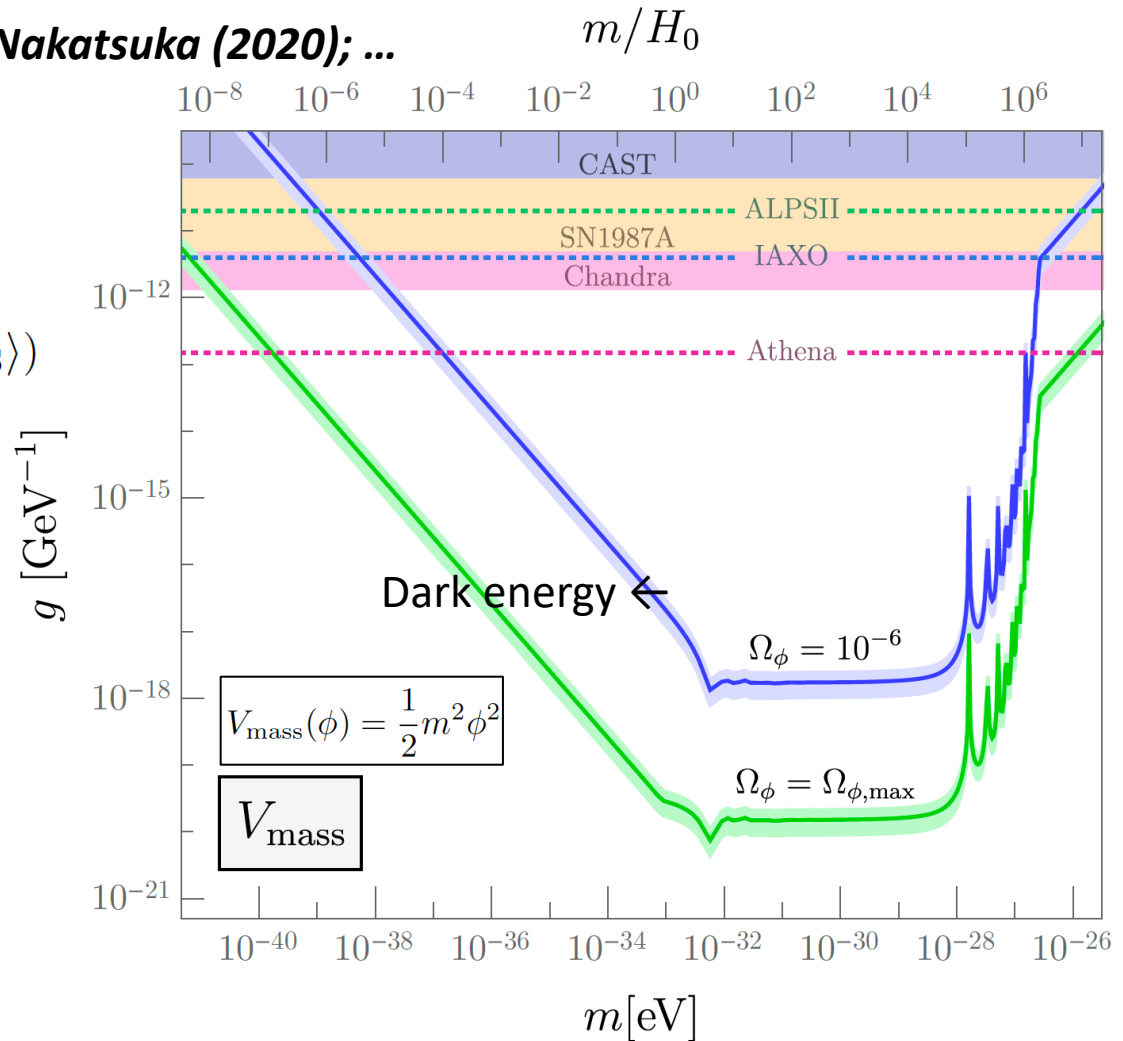
*Fujita, Minami, Murai & Nakatsuka (2020); ...*

➤ From this relationship

$$\beta = \frac{g_{\phi\gamma}}{2} \Delta\phi \equiv \frac{g_{\phi\gamma}}{2} (\phi_0 - \langle\phi_{\text{LSS}}\rangle)$$

we can constrain the  
parameter space of axion-  
photon coupling w.r.t.  
axion mass

➤ Support the presence of  
axion as dark energy

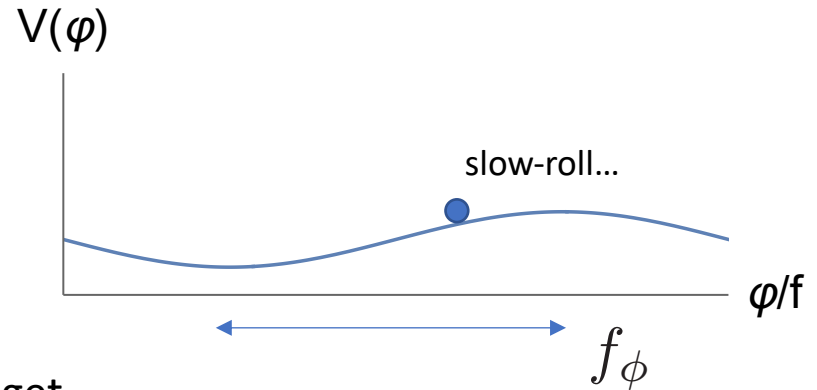


# Some issues of single-field model

*Friemann+ (1995); ...*

- Consider a **nearly flat** axion cosine potential

$$V(\phi) = m_\phi^2 f_\phi^2 \left[ 1 - \cos \left( \frac{\phi}{f_\phi} \right) \right]$$



- To satisfy the constraint on EoS parameter, we get

$$f_\phi \simeq 14 M_{\text{Pl}} \left( \frac{\Omega_\phi}{0.69} \right)^{1/2} \left( \frac{m_\phi/H_0}{0.1} \right)^{-1} > M_{\text{Pl}}$$

requires a **super-Planckian** decay constant or a fine-tuning of initial axion displacement

- To get the measured  $\beta$ , a **large anomaly coefficient** is required

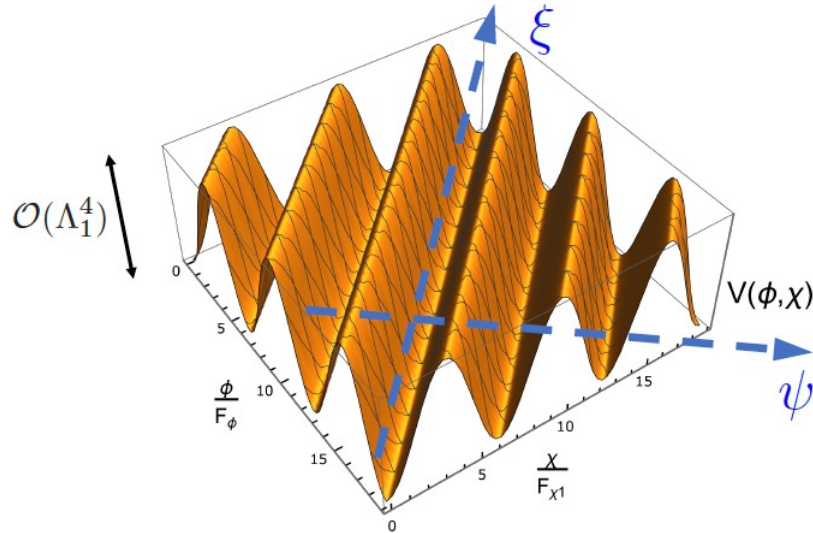
$$g_{\phi\gamma} = \frac{\alpha}{2\pi} \frac{c_{\phi\gamma}}{f_\phi}$$

$$|c_{\phi\gamma}| \simeq 7.5 \times 10^3 \left( \frac{\beta}{0.35\text{deg}} \right) \left( \frac{m_\phi/H_0}{0.1} \right)^{-2} \gg 1$$

# Two-fields axion model

Kim (1999)(2000), ...

Kim, Nilles & Peloso (2005), ...



## Potential

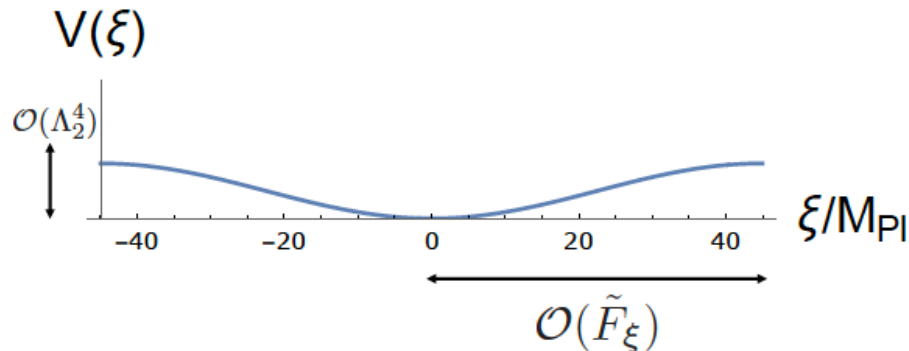
$$V(\phi, \chi) = \Lambda_1^4 \left[ 1 - \cos \left( \frac{\phi}{F_{\phi 1}} + \frac{\chi}{F_{\chi 1}} \right) \right] + \Lambda_2^4 \left[ 1 - \cos \left( \frac{\phi}{F_{\phi 2}} + \frac{\chi}{F_{\chi 2}} \right) \right]$$

$$(\Lambda_1^4 \gg \Lambda_2^4, F_i < M_{\text{Pl}})$$

$$\xi = \frac{F_\phi}{\sqrt{F_\phi^2 + F_{\chi 1}^2}} \phi - \frac{F_{\chi 1}}{\sqrt{F_\phi^2 + F_{\chi 1}^2}} \chi, \quad \psi = -\frac{F_{\chi 1}}{\sqrt{F_\phi^2 + F_{\chi 1}^2}} \phi - \frac{F_\phi}{\sqrt{F_\phi^2 + F_{\chi 1}^2}} \chi$$

- Nearly-flat direction can be realized by an alignment of the **multiple axion potentials**:

$$F_{\phi 1} = F_{\phi 2} \equiv F_\phi, \quad F_{\chi 2} = F_{\chi 1}(1 + \epsilon) \quad \epsilon \ll 1 : \text{misalignment parameter}$$



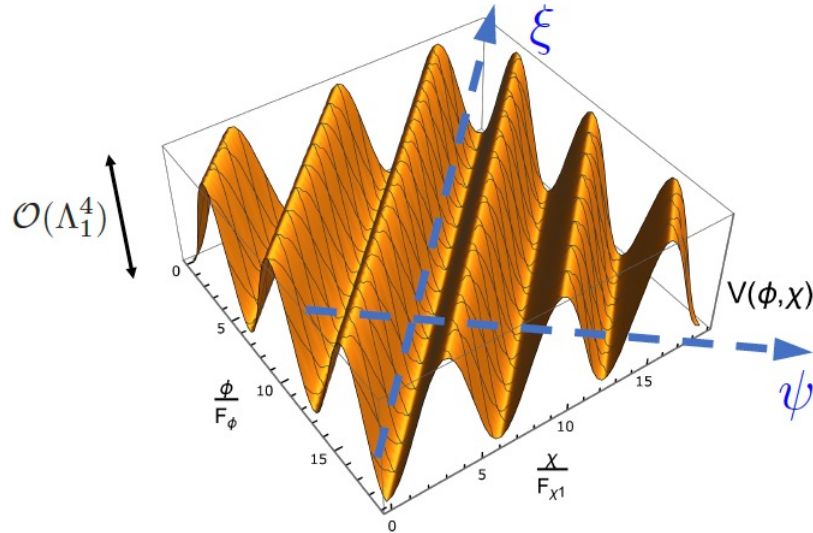
## Effective field range of axion

$$\tilde{F}_\xi = \frac{\sqrt{F_\phi^2 + F_{\chi 1}^2}}{\epsilon} \gg M_{\text{Pl}}$$

# Two-fields axion model

*Kim (1999)(2000), ...*

*Kim, Nilles & Peloso (2005), ...*



## Potential

$$V(\phi, \chi) = \Lambda_1^4 \left[ 1 - \cos \left( \frac{\phi}{F_{\phi 1}} + \frac{\chi}{F_{\chi 1}} \right) \right] + \Lambda_2^4 \left[ 1 - \cos \left( \frac{\phi}{F_{\phi 2}} + \frac{\chi}{F_{\chi 2}} \right) \right]$$

$$(\Lambda_1^4 \gg \Lambda_2^4, F_i < M_{\text{Pl}})$$

$$\xi = \frac{F_\phi}{\sqrt{F_\phi^2 + F_{\chi 1}^2}} \phi - \frac{F_{\chi 1}}{\sqrt{F_\phi^2 + F_{\chi 1}^2}} \chi, \quad \psi = -\frac{F_{\chi 1}}{\sqrt{F_\phi^2 + F_{\chi 1}^2}} \phi - \frac{F_\phi}{\sqrt{F_\phi^2 + F_{\chi 1}^2}} \chi$$

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Linear combinations of two-fields provide two **(nearly) massless** & **massive** fields

**Dark energy**

**Dark matter**

$\xi$

$\psi$

# Introduce axion-photon couplings

*IO (2021)*

- The interactions of photon to the (original) axion fields are given by

$$\mathcal{L} \supset \frac{\alpha}{8\pi} \left( \frac{\phi}{F_{\phi\gamma}} + \frac{\chi}{F_{\chi\gamma}} \right) F_{\mu\nu} \tilde{F}^{\mu\nu}$$

- In terms of  $(\psi, \xi)$ , the effective coupling constants are obtained:

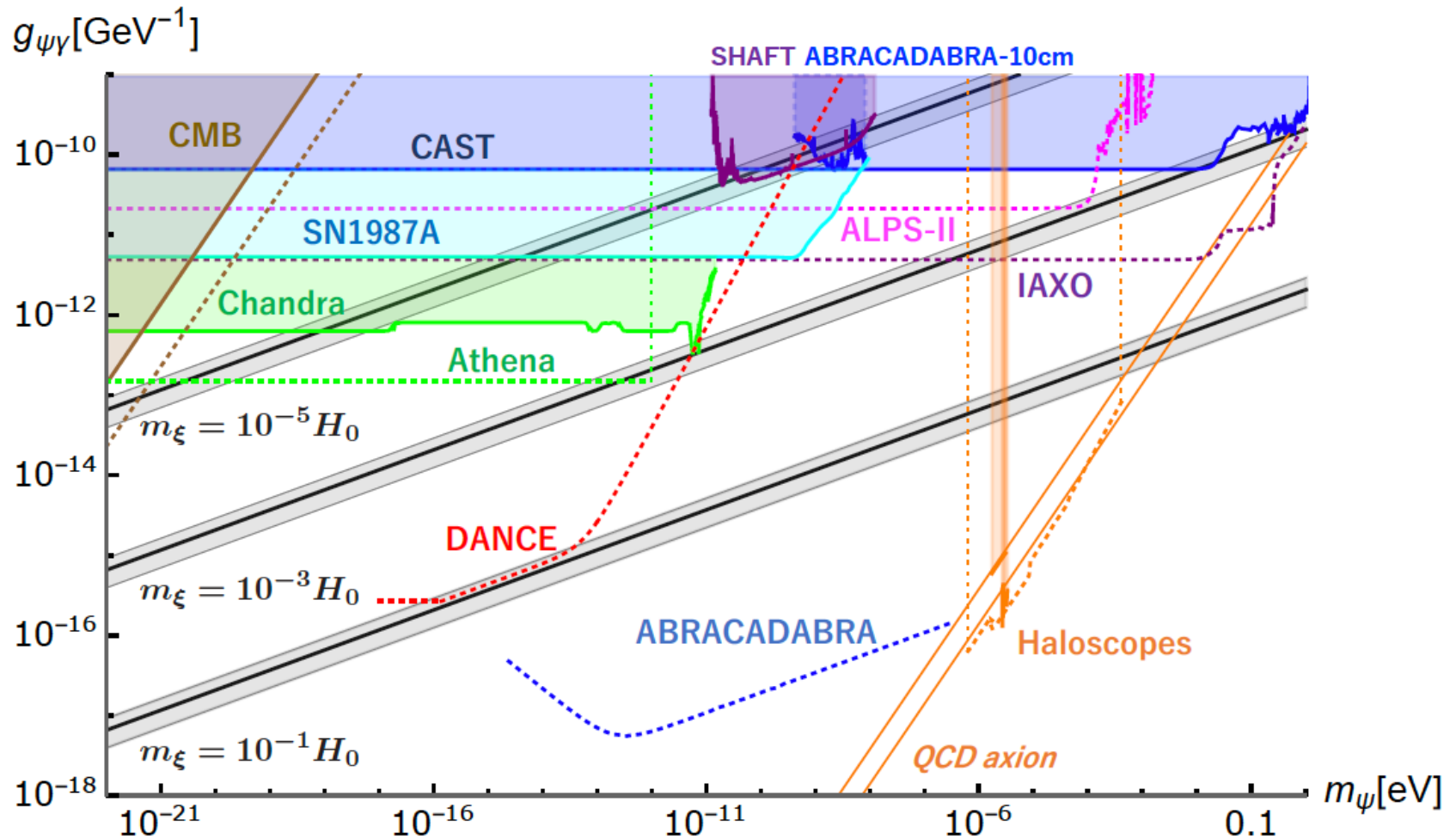
$$g_{\xi\gamma} = \frac{\alpha}{2\pi} \frac{c_{\xi\gamma}}{\tilde{F}_{\xi}} , \quad c_{\xi\gamma} \equiv \frac{1}{\epsilon} \left( \frac{F_{\phi}}{F_{\phi\gamma}} - \frac{F_{\chi 1}}{F_{\chi\gamma}} \right) ,$$
$$g_{\psi\gamma} = \frac{\alpha}{2\pi} \frac{c_{\psi\gamma}}{\tilde{F}_{\psi}} , \quad c_{\psi\gamma} \equiv - \left( \frac{F_{\phi}}{F_{\chi\gamma}} + \frac{F_{\chi 1}}{F_{\phi\gamma}} \right) \frac{F_{\phi} F_{\chi 1}}{F_{\phi}^2 + F_{\chi 1}^2}$$

(release 3)

- $g_{\xi\gamma}$  (dark energy) is fixed by the measured birefringence angle  $\beta = 0.35 \pm 0.14$

→ the parameter space of  $g_{\psi\gamma}$  (dark matter) is also constrained

# Parameter space of axion DM



# Summary & Outlook

- Axions are the promising candidate for the dark sector of our universe. Up to now, a variety of astrophysical/experimental methods has been proposed by using axion-photon conversions.
- Photon's birefringence measurements potentially develop a new frontier of the axion search! We have suggested an experimental scheme to search for axion dark matter with optical cavities used in such as gravitational wave detectors.
- A recent measurement of CMB birefringence gives us a tantalizing hint for the axion physics, especially axion as a dark energy. Based on a multiple axion scenario, this observable can connect the constraints on axion as dark energy and dark matter.
- More exciting discovery to come in next decades!?

**감사드립니다**