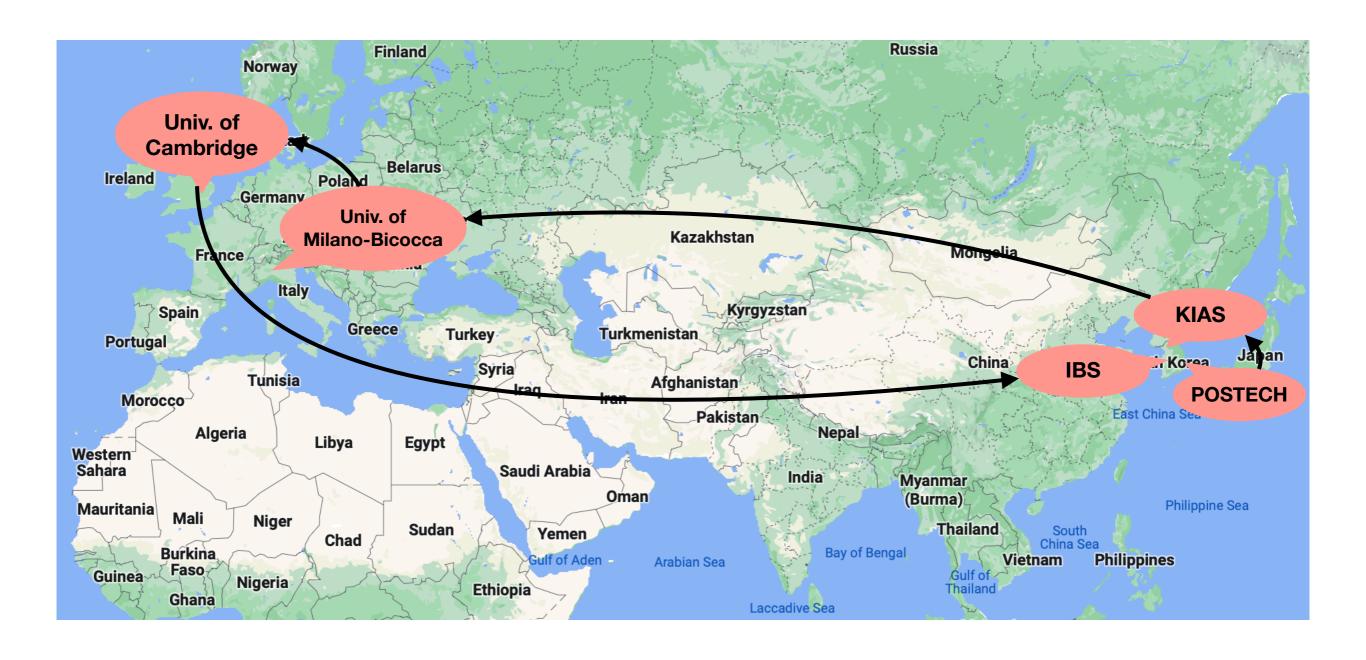
From Strongly Coupled SQFT to Quantum Gravity

Chiung Hwang

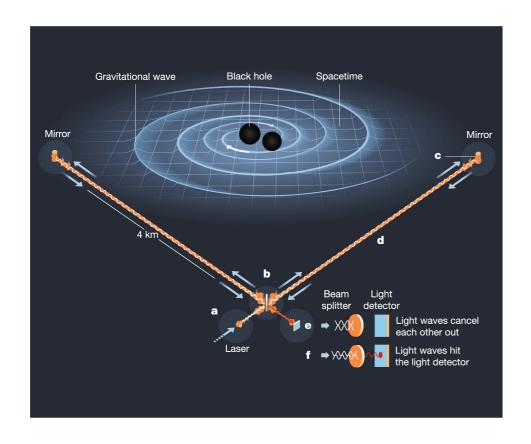
Who Am I?

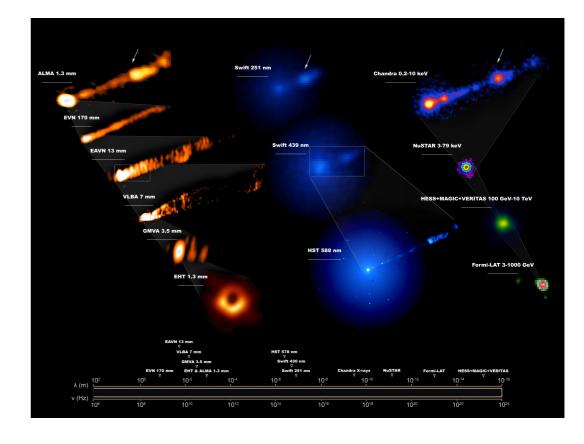
- Chiung Hwang
- Joined CTPU in October as a Young Scientist Fellow
- My research interest is holography and duality of quantum field theory.

Academic Journey



Physics of Black Holes

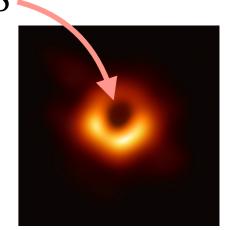




- Still, a black hole is a mysterious object.
- A (classical) black hole vs. the 2nd law of thermodynamics

Entropy
$$S$$

$$\Delta S \ge 0$$



- The observable entropy seems to decrease.
- A violation of the 2nd law or non-zero entropy of a black hole?
- Signals non-zero temperature of a black hole—thermal radiation?

 Semi-classical analysis—the quantum fluctuation of fields around a classical black hole background

- Black holes do radiate! [Hawking 74]
- Thermal properties of a black hole

$$T_H = \frac{1}{8\pi M k_B} \qquad S_{BH} = \frac{1}{4} k_B A$$

- The *microstates* of a black hole? $e^{S/k_B} \sim e^{10^{44}}$ for 10 M_{\odot}
- On the other hand, everything is squeezed into a single point in a Schwarzschild black hole.
- Quantum nature of spacetime itself

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Quantum gravity and black hole thermodynamics go side by side.

Black Hole Thermodynamics and Holographic Duality

Black Hole in the Anti-de Sitter Space

The Hawking temperature

$$T_H = \frac{\hbar c^3}{8\pi GM k_B} \sim \frac{1}{M}$$

- A black hole in the flat spacetime -> negative specific heat
- The Anti-de Sitter (AdS) spacetime:

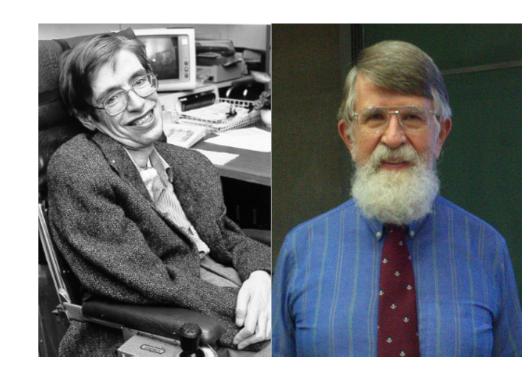
the maximally symmetric solution to Einstein's equation with negative cosmological constant

Positive specific heat, thermodynamically stable

Black Hole in the Anti-de Sitter Space

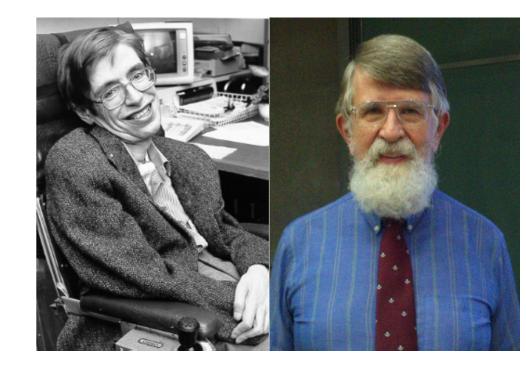
- Hawking-Page phase transition (1983)
 - $T < T_0$: thermal AdS in equilibrium
 - $T_0 < T < T_1$: thermal AdS preferred
 - $T_1 < T < T_2$: black hole preferred
 - $T_2 < T$: no equilibrium without a black hole





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Nontrivial phase structure with the 1st order transition

AdS is a nice playground to study black hole thermodynamics.

Holographic Duality

- A groundbreaking property of the AdS gravity [Maldacena 97]
- Also called: the AdS/CFT correspondence, the gauge/gravity duality, the Maldacena duality, ...
- The most cited paper in INSPIRE (18,000 times)

The Large N Limit of Superconformal field theories and supergravity

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Abstract

We show that the large N limit of certain conformal field theories in various dimensions include in their Hilbert space a sector describing supergravity on the product of Anti-deSitter spacetimes, spheres and other compact manifolds. This is shown by taking



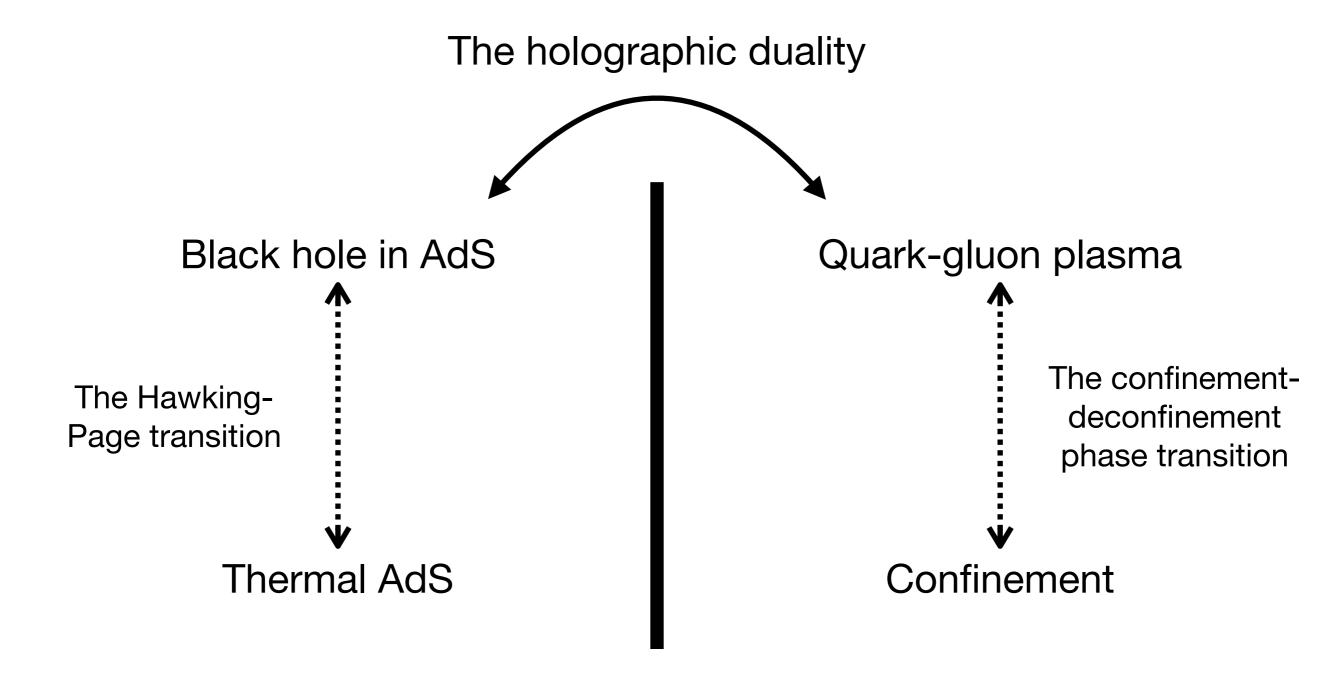
Holographic Duality

Duality between d-dimensional QFT and d+1-dimensional gravity

3d Gauge Theory	4d AdS Gravity
More (Less) Gluons	Small (Large) Quantum Correction
Strong (Weak) Coupling	Small (Large) Higher Cuvature Correction

- Quantum gravity described by QFT
- Strongly coupled QFT described by gravity

Black Hole from Non-Perturbative Dynamics of QFT



Gravity side:

 $\mathcal{N}=8$ gauged supergravity

- An AdS black hole with (M,Q,J) satisfying $M=4Q+J=\tilde{Q}+J$
- Black hole entropy [Choi-CH-Kim-Nahmgoong 18]

$$S = \frac{\pi}{2} \sqrt{\tilde{Q}^2 - \frac{8L^2}{G} \frac{J}{\tilde{Q}}}$$

QFT side:

• 2+1d U(N) gauge theory with matter fields

(Dual of) the 2+1d maximally supersymmetric U(N) Yang-Mills theory

The grand canonical partition function

$$Z = \operatorname{Tr}_{E = \tilde{Q} + J} e^{-\Delta \tilde{Q} - \omega J} \quad \blacksquare$$

Hard to compute for strongly coupled theories

Witten (1982)

$$I = \operatorname{Tr}_{E = \tilde{Q} + J} (-1)^{F} e^{-\Delta \tilde{Q} - \omega J}$$



Topological, exactly computable even if the theory is strongly coupled

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Topological, exactly computable even if the theory is strongly coupled

Explore phases of the 2+1d gauge theory dual to 4-dimensional AdS gravity!

Confinement

- 1. The confinement phase [Bhattacharya-Minwalla 08, S.Kim 10]
- No isolated quark, only colorless bound states
- An order parameter for the confinement: the Polyakov loop

$$P \sim \operatorname{Tr} e^{\int A} = \sum_{k=1}^{N} e^{u_k}$$
 # of gluons: $N \times N$ # of quarks: N

The free energy of an isolated quark

$$F_{quark} \sim -\log \langle P \rangle$$

• In the confinement phase with $\left\langle u_k \right\rangle = 2\pi i k/N$

- 2. The deconfinement phase [Choi-CH-Kim 19]
- A saddle point with $\langle P \rangle \neq 0$?
- A long-standing problem over 10 years

- Non-perturbative monopoles play a crucial role.
- A new saddle point with the monopole condensation

$$\langle P \rangle \sim e^{\sqrt{N}}$$

2. The deconfinement phase [Choi-CH-Kim 19]

The grand canonical partition function and the entropy

$$\log I \approx i \frac{4\sqrt{2}N^{\frac{3}{2}}}{3} \frac{\Delta^2}{\omega} \qquad S = \frac{\pi}{2} \sqrt{\tilde{Q}^2 - \frac{64}{9} \frac{N^3 J}{\tilde{Q}}}$$

Holographic dictionary:
$$\frac{L^2}{G} = \frac{2\sqrt{2}}{3}N^{\frac{3}{2}}$$

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Quantum mechanical derivation of the black hole entropy

The crucial effect of the monopole condensation

• # of gluons: N^2



The naive dof in the deconfinement:

 N^2

• # of quarks: N

The holographic dictionary predicts

$$\log I \sim \frac{1}{G} \sim N^{\frac{3}{2}}$$

• N^2 vs $N^{\frac{3}{2}}$?

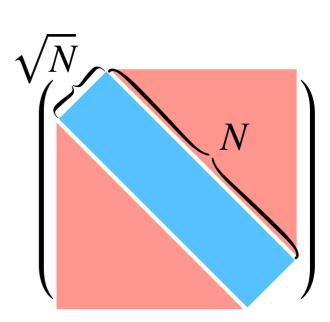
• Effective mass of $N \times N$ gluons due to monopole condensation (Higgs mechanism)

$$M_{ab} \sim N^{-\frac{1}{2}} |a-b|$$

•
$$|a-b| \lesssim N^{\frac{1}{2}}$$
: $M_{ab} \lesssim 1$

•
$$|a-b| \gg N^{\frac{1}{2}}$$
: $M_{ab} \gg 1$





ullet The free energy of N quarks

$$N \times F_{quark} \sim -N \log \langle P \rangle \sim -N^{\frac{3}{2}}$$

- A number of results reporting $N^{\frac{3}{2}}$ for other non-perturbative observables in the literature [Gang-**CH**-Kim-Park 11, Cheon-Gang-**CH**-Nagaoka-Park 11, ...]
- The first physical understanding of the origin of $N^{\frac{3}{2}}$
- The condensation of the monopole operator in the deconfinement phase
 - -> partial deconfinement rather than full deconfinement
- Nontrivial prediction for strong dynamics of the 2+1d U(N) gauge theory
- The 1st order phase transition between the partial deconfinement and the confinement expected
- Exactly (a supersymmetric version of) the Hawking-Page transition between black hole and thermal AdS phases in dual gravity

What Is Next?

- Holographic duality relates thermodynamics of gravity to strong dynamics of QFT.
- Black hole entropy by counting quantum states of the dual QFT

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What are those states holographically dual to a black hole?

Classification of Black Hole States

- Much harder than counting
- Need some simplification
- More (super)symmetric states?
- Already classified and no deconfinement

4D MSYM & AdS5 Black Holes

- Instead, let's consider a 3+1d maximally supersymmetric Yang-Mills theory.
- Dual to 5-dimensional AdS
- The states preserving more supersymmetry are not completely classified.

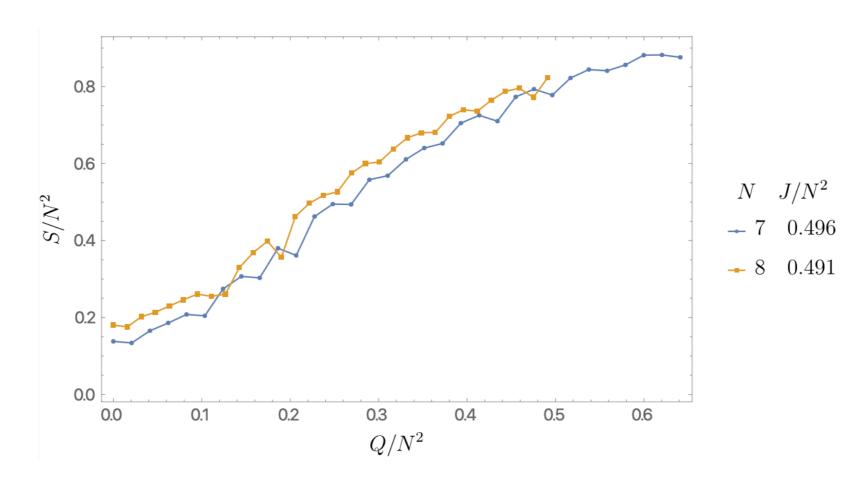
A New Black Hole?

- The known supersymmetric black hole solution in AdS5 preserves a minimal amount of supersymmetry.
- A black hole preserving more supersymmetry?
- More constraints on dual QFT states -> could be useful for a classification

- No known solution on the gravity side
- Deconfinement on the QFT side?

Black Hole with More Supersymmetry

- Work in progress with O. Dias, P. Mitra, J. Santos
- Explore the deconfinement phase ($\sim N^2$) in more supersymmetric sectors



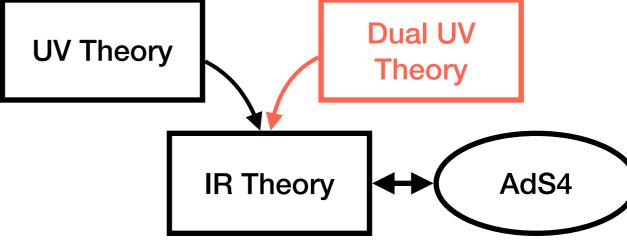
Black Hole with More Supersymmetry

- Predicts a new black hole—a guideline where to investigate to find a new black hole solution
- What is entropy?
- More constrained QFT states—a classification of black hole states?

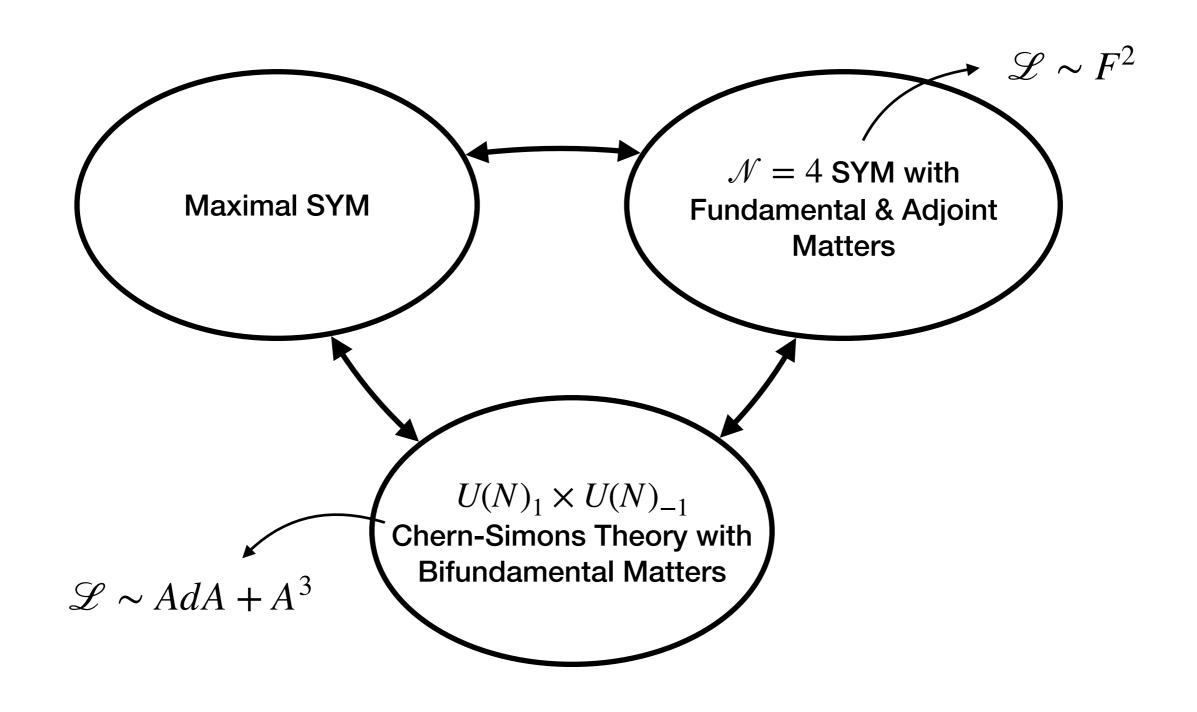
Duality—3D MSYM Theory Revisited

3D MSYM Revisited

- Recall the 2+1d maximally supersymmetric Yang-Mills theory
- Flows to a strongly interacting (conformal) field theory, which is dual to AdS4 gravity
- UV symmetry ≠ IR symmetry
- Not appropriate to study the dual black hole entropy
- Instead, a dual description flowing to the same IR CFT can be used!



3D MSYM Revisited



Mirror Symmetry

- A special case of Mirror Symmetry [Intriligator-Seiberg 96]
- A simple example: 2+1d supersymmetric QED
 - $\mathcal{N}=4$ supersymmetric U(1) gauge theory with N_f fundamental matters
 - $\mathcal{N}=4$ supersymmetric $U(1)^{N_f-1}$ gauge theory with N_f-2 bifundamental and two fundamental matters
- More non-Abelian Yang-Mills examples

Mirror Symmetry

- Those dualities are motivated by String Theory
- Many nontrivial field-theoretic tests
- But no field-theoretic derivation
- Understanding at a more elementary level?

Theory vs. Fields

- A (Lagrangian) theory is defined in terms of fields & interactions
- Interactions are fixed by $\mathcal{N}=4$ supersymmetry

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Can we dualize each elementary field rather than the whole theory?

Building Blocks of Mirror Symmetry

- Find new dualities between a theory of a single free field and an interacting gauge theory
- Depends on the representation of the field
 - A fundamental field
 - A bifundamental field
 - An adjoint field

CH-Pasquetti-Sacchi 21, Bottini-CH-Pasquetti-Sacchi 21

Work in progress

 A large class of mirror dual pairs can be derived from the fundamental duality & the bifundamental duality only

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Building blocks of mirror symmetry!

What Is Next?

- Application to holographic theories?
- E.g., the dualities of the MSYM
 - Dualization of an adjoint field?
 - Chern-Simons instead Yang-Mills?
- Better understand the duality map of the states
- Hints of black hole states?

Conclusion

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- The holographic duality relates gravity and QFT.
- Quantum aspects of black holes from non-perturbative dynamics of holographic QFTs
- Many useful tools: supersymmetry, duality, ...

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A journey to a better/new understanding of strongly coupled QFT and quantum gravity!

Thank you