

# QUANTUM POSITION VERIFICATION AND TIME-CONSTRAINED STATE DISCRIMINATION

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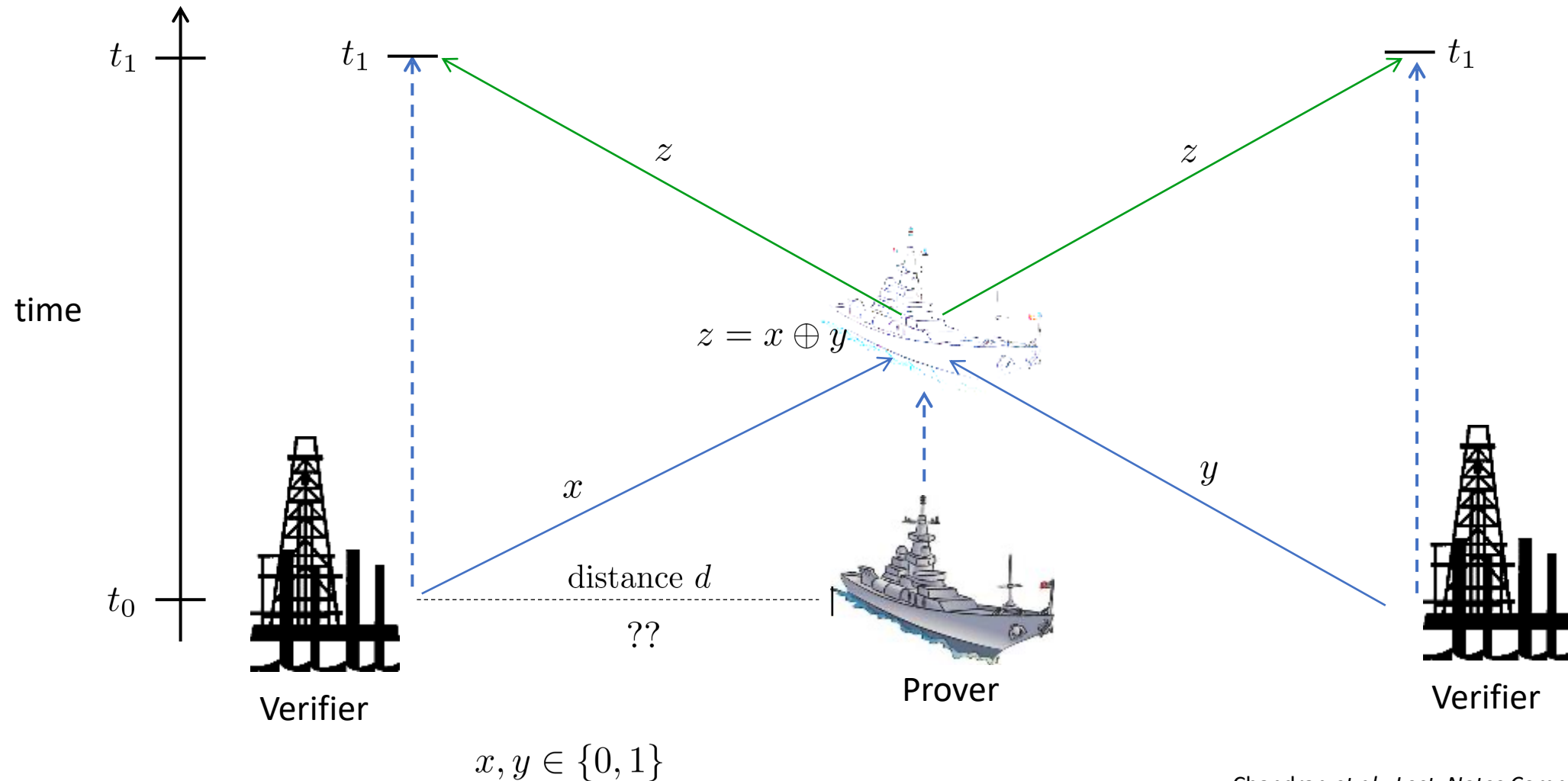
Philip Lunel

# Outline of talk

1. Introduce the cryptographic task of quantum position verification (QPV)
2. Make a connection to AdS/CFT correspondence (one slide)
3. Describe QPV as a quantum state discrimination problem under restricted communication
4. Compare QPV product state discrimination protocols using **classical** versus **quantum** communication

# Classical Position Verification

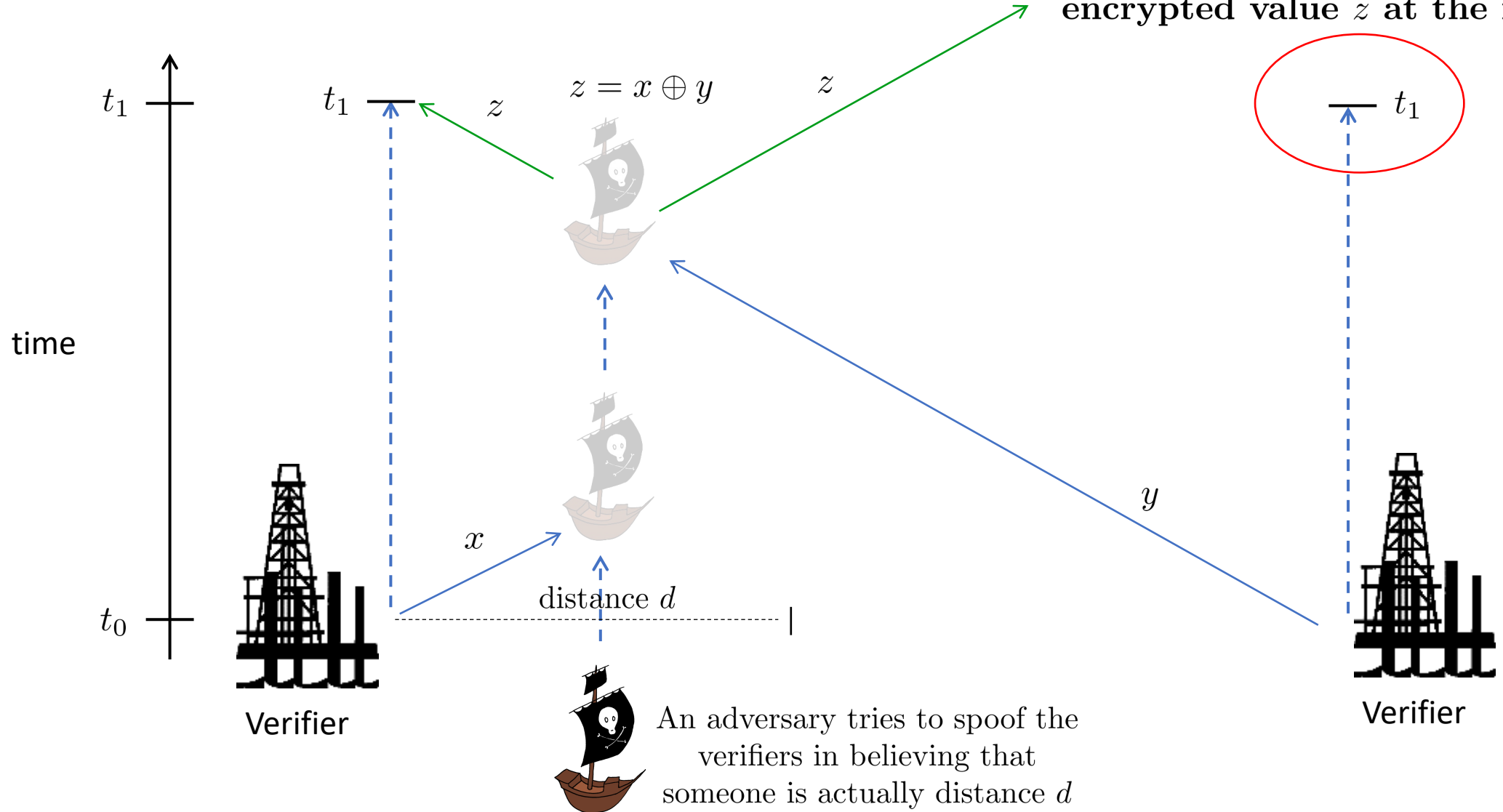
- **Goal:** Verify that the prover is truly a distance  $d$  from the leftmost verifier.



# Classical Position Verification

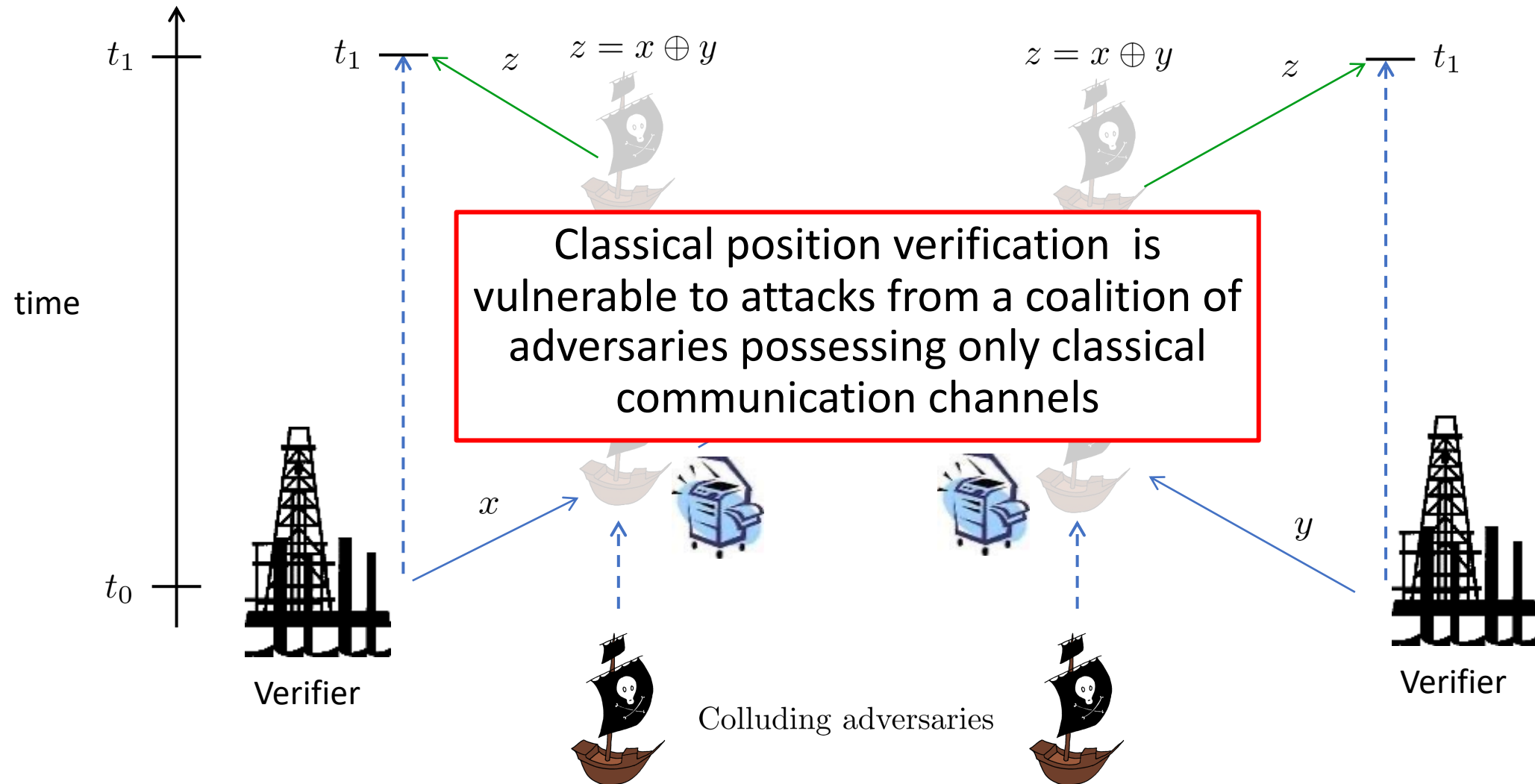
- Why does this work?

The second verifier won't receive the encrypted value  $z$  at the right time!



# Classical Position Verification

- Not so fast...



# Quantum Position Verification (QPV)

$$x, y \in \{0, 1\}$$

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$$

$$H|0\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) =: |+\rangle$$

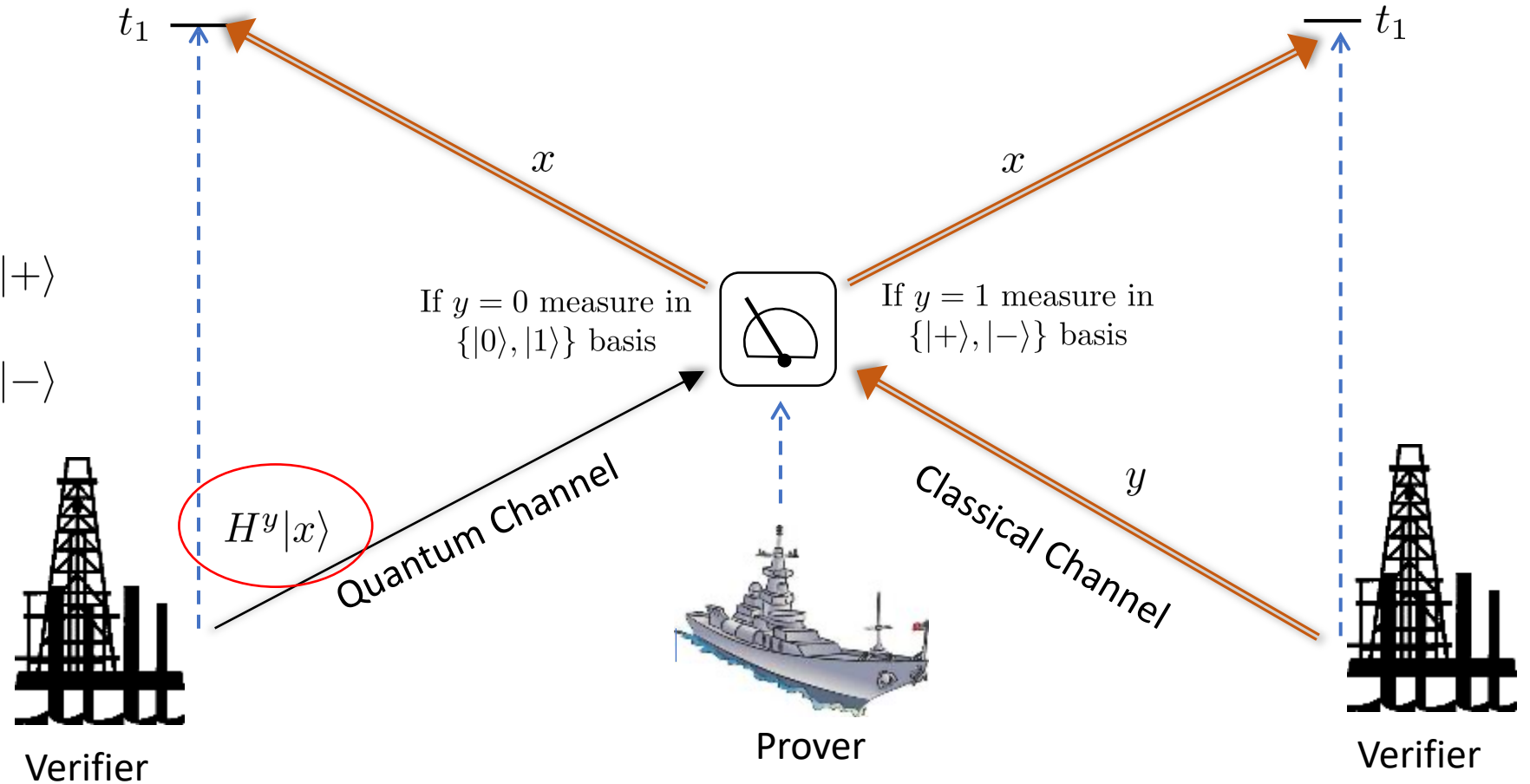
$$H|1\rangle = \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle) =: |-\rangle$$

$$H^0|0\rangle = |0\rangle$$

$$H^0|1\rangle = |1\rangle$$

So the verifier sends one of the BB84 states:

$$\{|0\rangle, |1\rangle, |+\rangle, |-\rangle\}$$



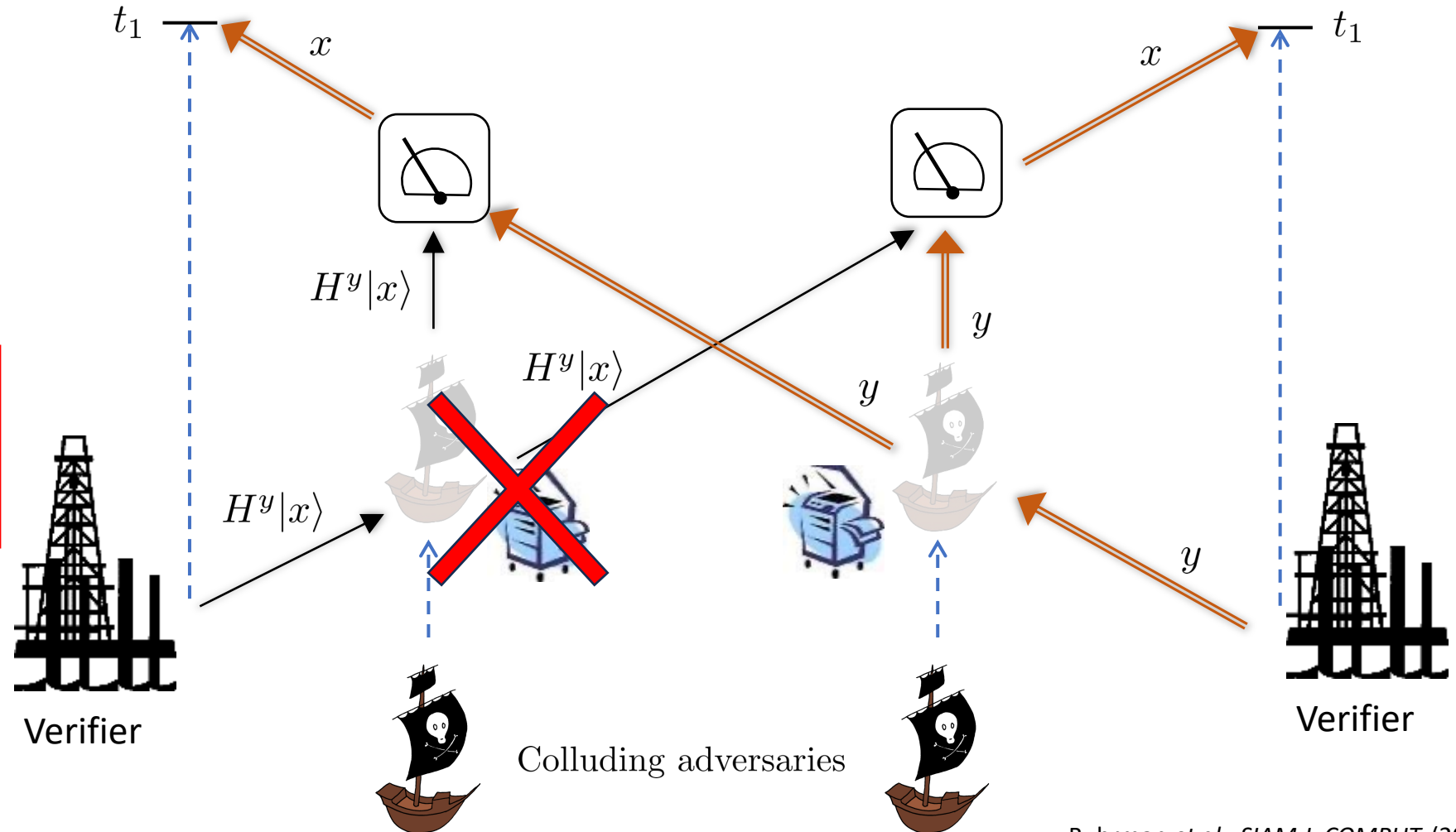
# Quantum Position Verification (QPV)

- Is it secure? What about the classical attack?

Not possible due to  
quantum no-cloning  
principle!

$$H^y|x\rangle \rightarrow H^y|x\rangle \otimes H^y|x\rangle$$

Not physically  
realizable

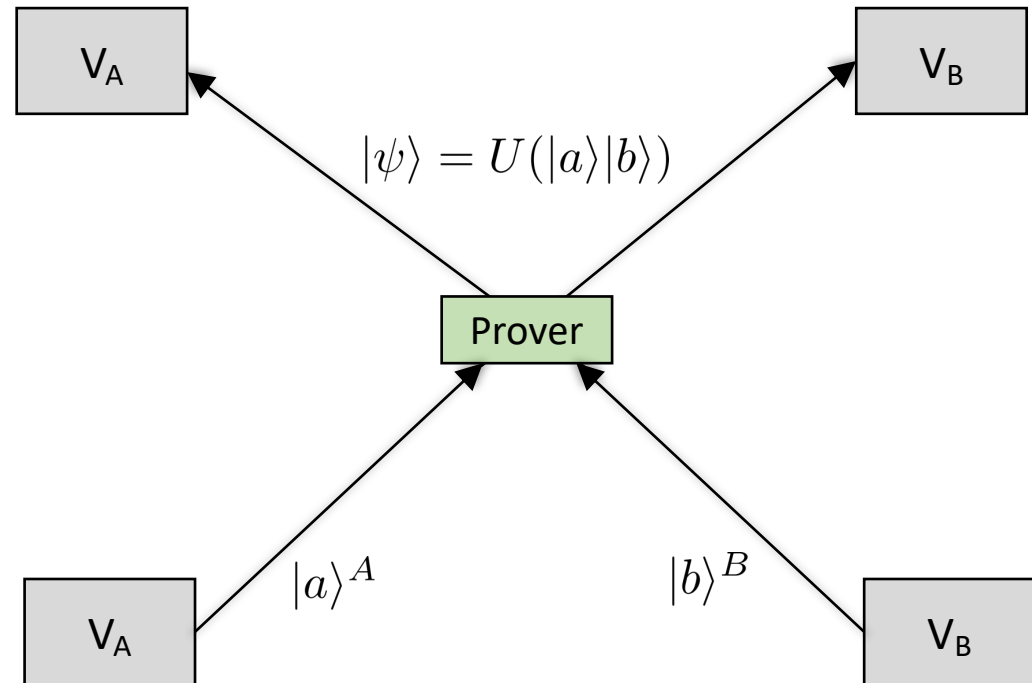


# Quantum Position Verification (QPV)

- Is it secure?
- Cloning attacks fail, but what about others?
- The prover is asked to perform some quantum computation  $U$  on the inputs  $|a\rangle|b\rangle$  and return  $|\psi\rangle = U(|a\rangle|b\rangle)$  within the correct time.
- **Intuition:** The unitary  $U$  should not be local:

$$U \neq U^A \otimes U^B.$$

A general QPV protocol



Honest QPV computation



# Quantum Position Verification (QPV)

- Is it secure?
- The dishonest adversaries attempt to implement the computation  $U$  in spatially separated laboratories.
- Additionally they can have some pre-shared entanglement.

This is sometimes referred to as  
“instantaneous nonlocal computation”.

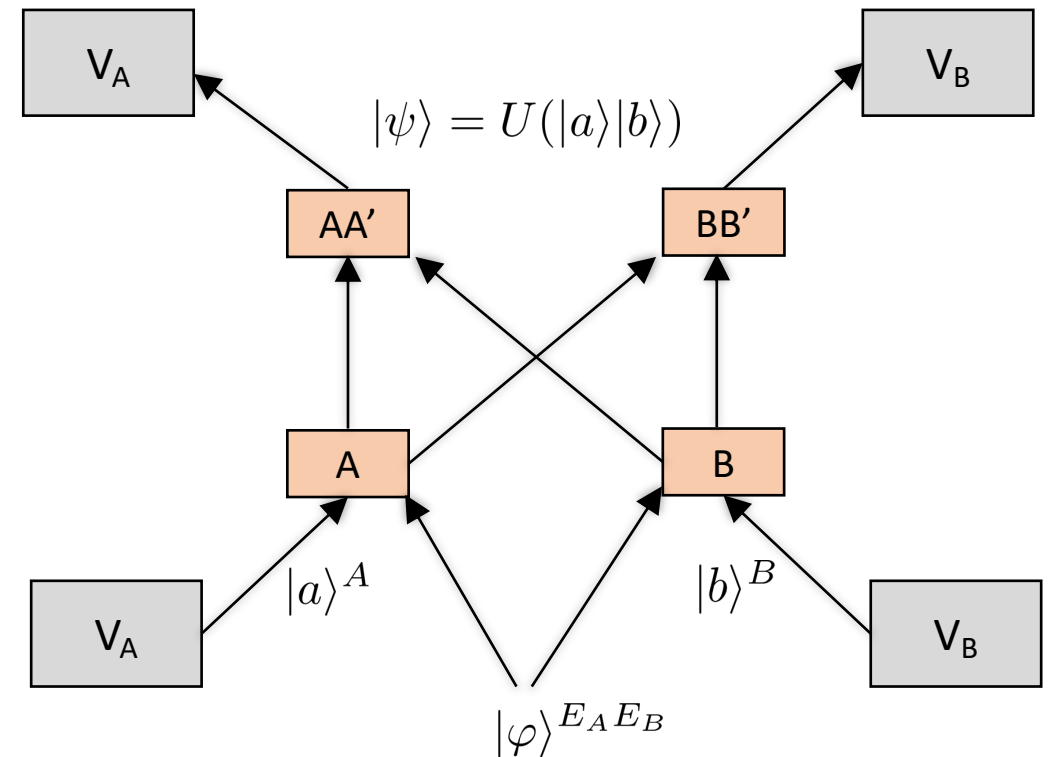
## Bad news:

Given enough entanglement, the adversaries can implement *any* quantum computation  $U$ .

## Good news:

There are certain computations that cannot be implemented unless the adversaries have a linear amount of entanglement.

## A general QPV protocol



## Dishonest QPV computation

# Quantum Position Verification (QPV)

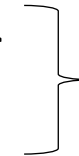
- **Fundamental Question:** How much entanglement is needed to break quantum position verification?

- The best known attack is based on “Port-Based Teleportation”:

Beigi and König, *New J. Phys.* (2011)

For  $n$  qubits sent from the verifier, QPV becomes insecure if adversaries use  $O(\exp(n))$  ebits.

- It is unknown whether attacks are possible that use a sub-exponential number of ebits for the hackers.

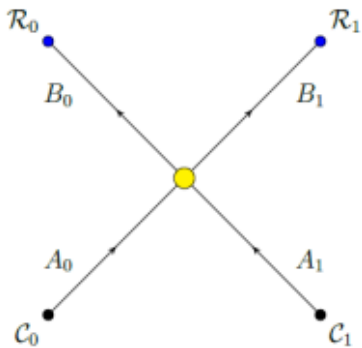


Open problem!!

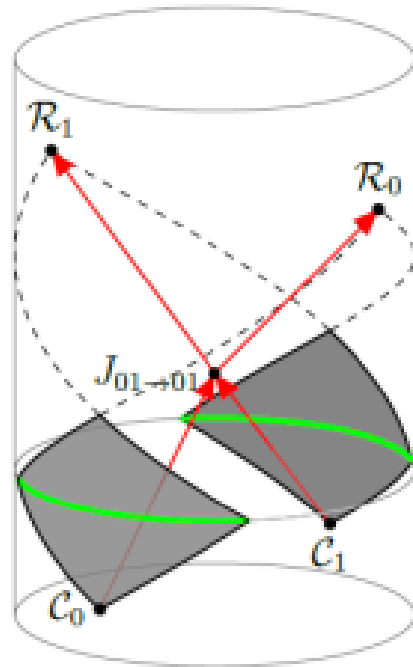
- Assuming a high entanglement cost for hacking, secure quantum position verification might be a good candidate for implementation on first-generation quantum networks.
- However, the time delay in measurement by the prover could be problematic.

# Connection to AdS/CFT correspondence

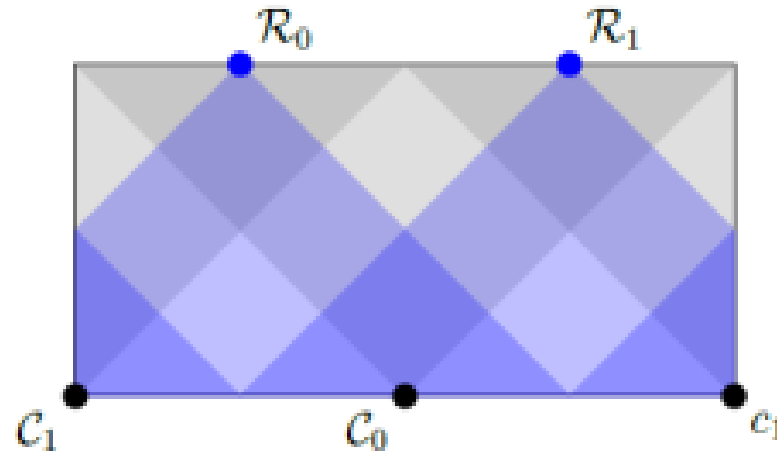
- The AdS/CFT correspondence proposes that quantum gravity in  $(d + 1)$ -dim anti-de Sitter space (**the bulk**) can be described by a  $d$ -dim non-gravitational conformal field theory (**the boundary**).



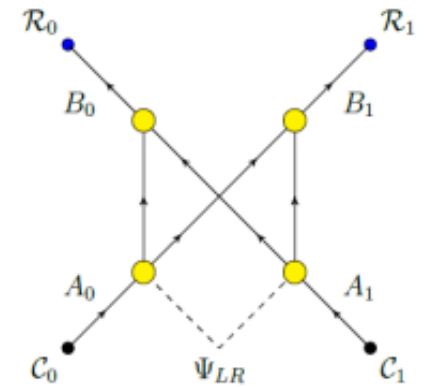
Honest QPV  
computation



The computation  
in the bulk



The computation on the boundary



Dishonest QPV  
computation

**Puzzle:** There is no causally consistent spacetime point for the computation to occur.

**Proposed solution:** Bulk computation = Honest QPV  
Boundary computation = Dishonest QPV

Please read Alex May's paper:  
*Quantum* **6**, 864 (2022).

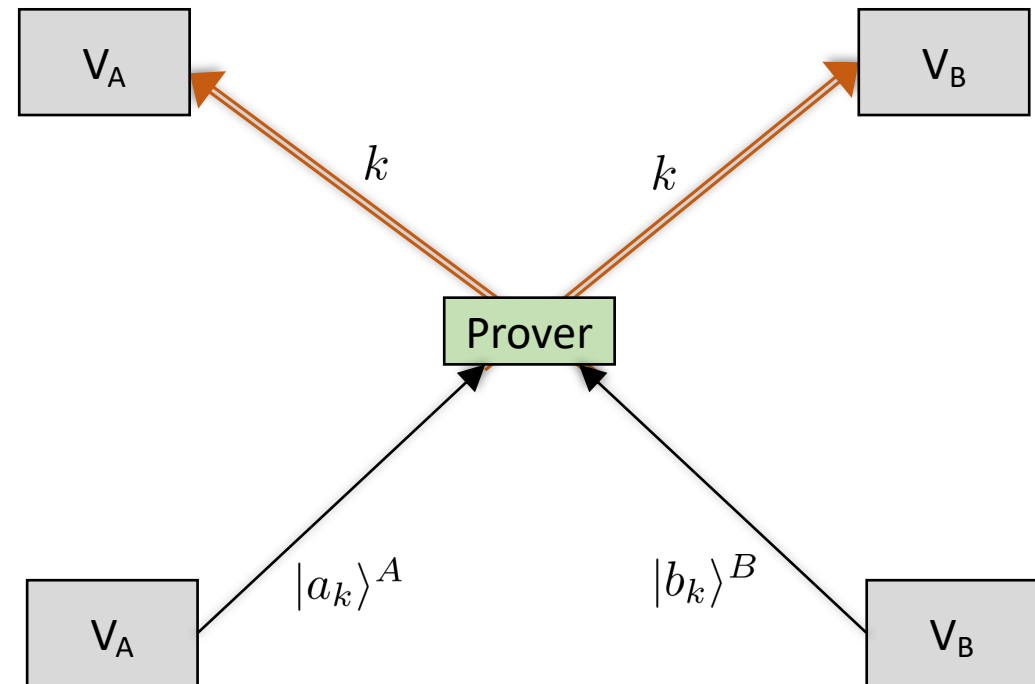
# QPV and state discrimination

- Let us consider a family of QPV protocols based on state discrimination.

A family of  
orthogonal  
states

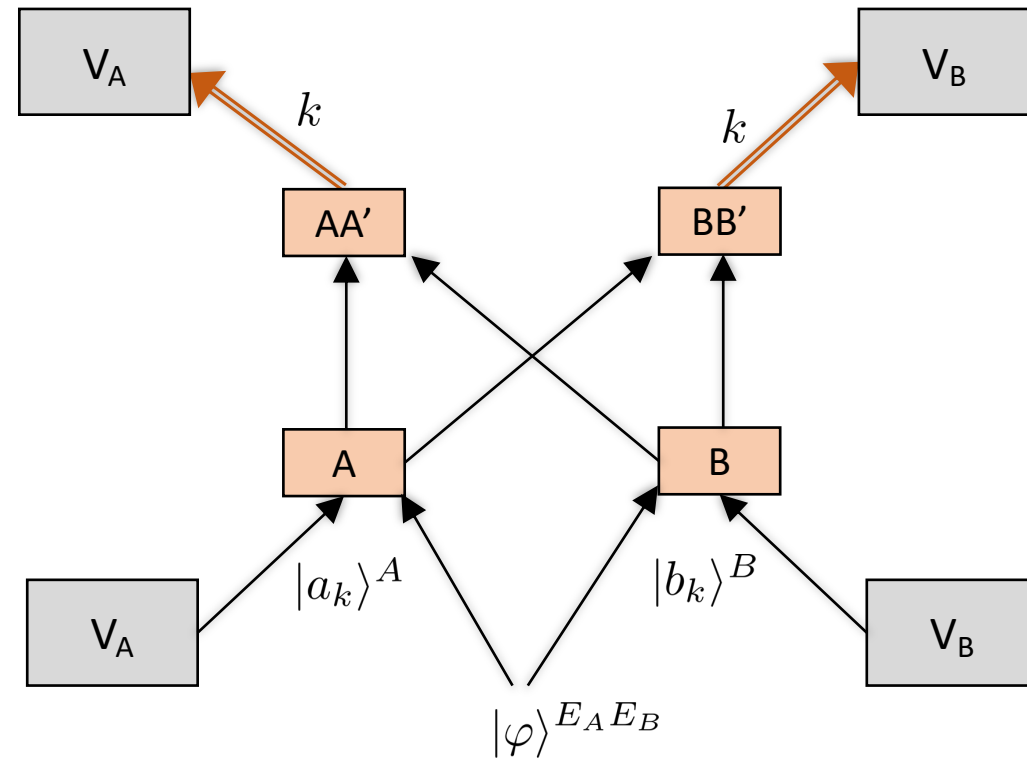
$$\left\{ \begin{array}{l} |\psi_1\rangle^{AB} = |0\rangle^A \otimes |0\rangle^B \\ |\psi_2\rangle^{AB} = |1\rangle^A \otimes |0\rangle^B \\ |\psi_3\rangle^{AB} = |+\rangle^A \otimes |1\rangle^B \\ |\psi_4\rangle^{AB} = |-\rangle^A \otimes |1\rangle^B \end{array} \right.$$

- The prover needs to identify which bipartite state  $|\psi_k\rangle^{AB}$  was sent by the verifiers.



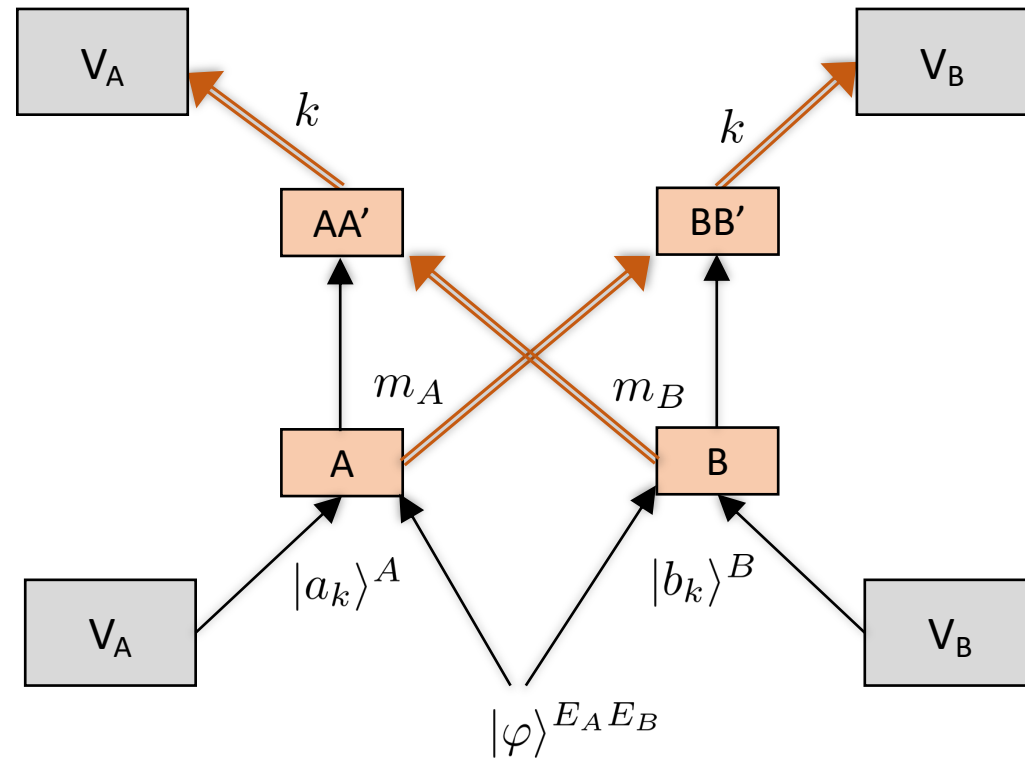
# QPV and state discrimination

- In order to be secure, the orthogonality of the encoded states  $|\psi_k\rangle$  must be sufficiently nonlocal.
- They should not be distinguishable by  
    local operations and  
    simultaneous communication.
- Different adversarial models to consider:
  - Local operations and simultaneous quantum communication (**LOSQC**)
  - Entanglement-assisted local operations and simultaneous quantum communication (**eLOSQC**)



# QPV and state discrimination

- In order to be secure, the orthogonality of the encoded states  $|\psi_k\rangle$  must be sufficiently nonlocal.
- They should not be distinguishable by **local operations** and **simultaneous communication**.
- Different adversarial models to consider:
  - Local operations and simultaneous classical communication (**LOSCC**)
  - Entanglement-assisted local operations and simultaneous classical communication (**eLOSCC**)

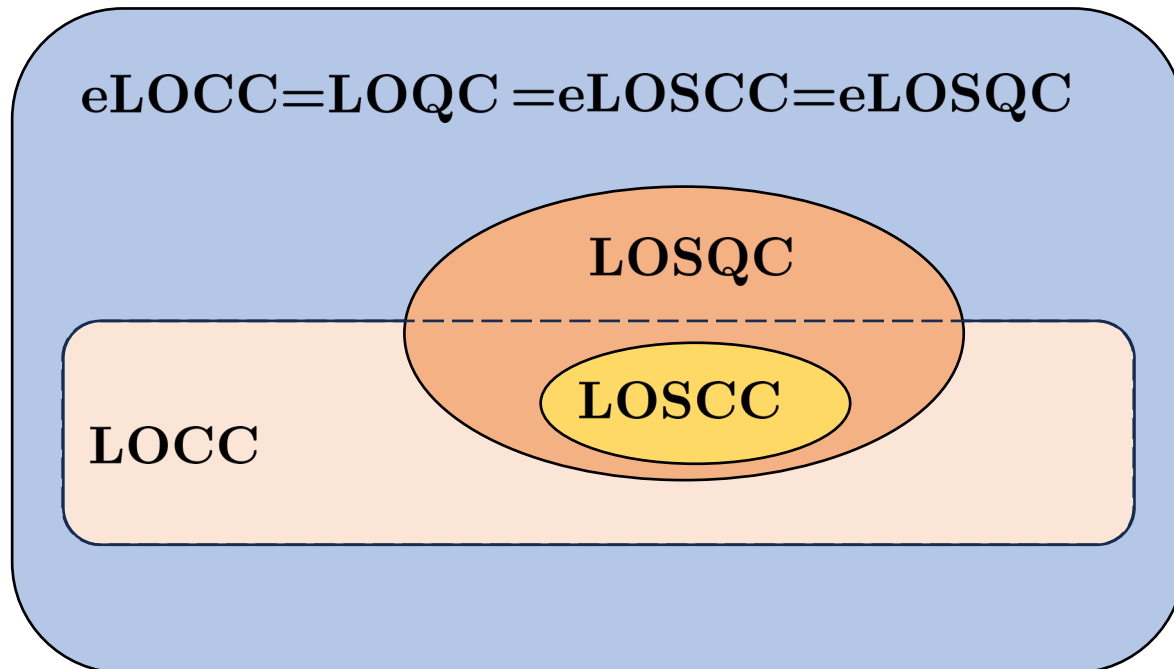


# Different operational classes

- These should be compared to standard:

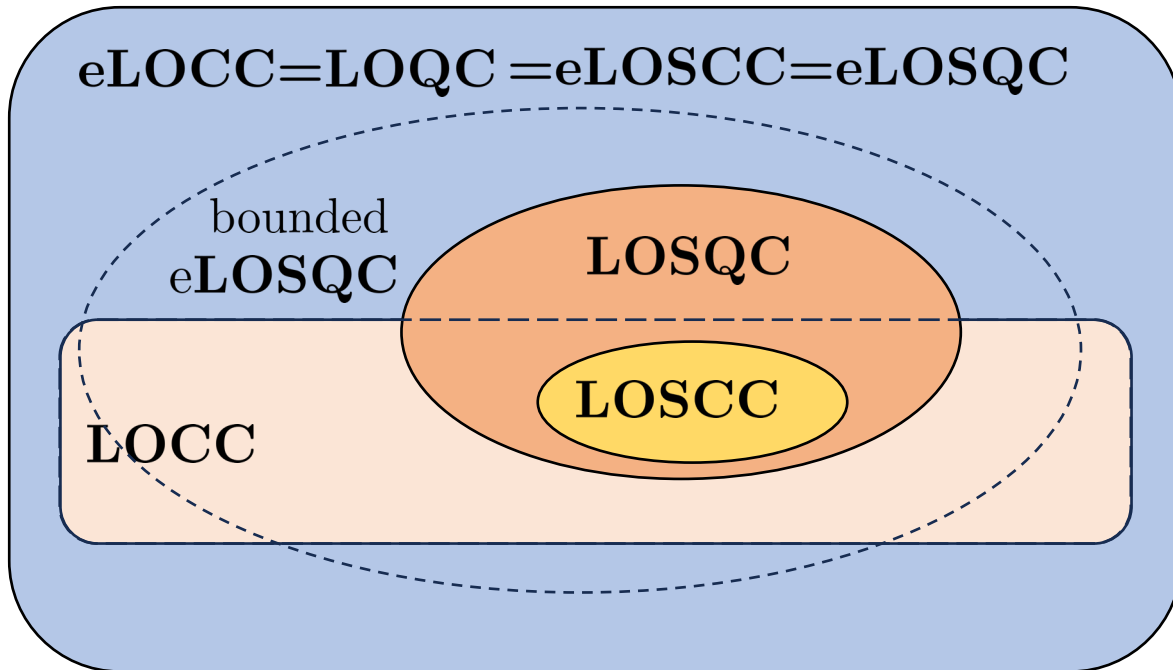
- Local operations and classical communication (**LOCC**)
- Entanglement-assisted local operations and classical communication (**eLOCC**)
- Local operations and quantum communication (**LOQC**)

Unrestricted  
classical communication

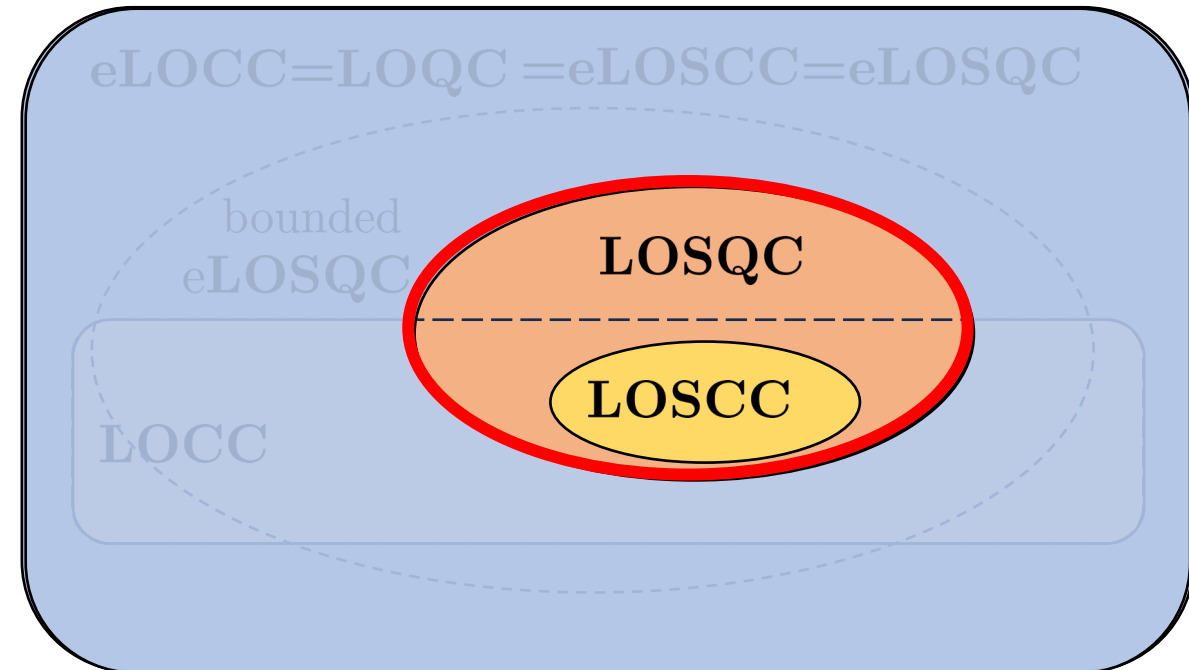


- Only a few papers (motivated by QPV) have explored the relationship between these operational classes.

# Different operational classes



- The intermediate regime of **bounded entanglement** is where most QPV analysis sits.
- Every family of orthogonal  $\{|\psi_k\rangle\}_k$  that is difficult to discriminate using a class of operations constitutes a good QPV scheme under attacks from that class.



- The **no pre-shared entanglement** model is the simplest to analyze, but even in this scenario relatively little is known.

- Simplify the problem even further:  
How well can a family of orthogonal **product states**

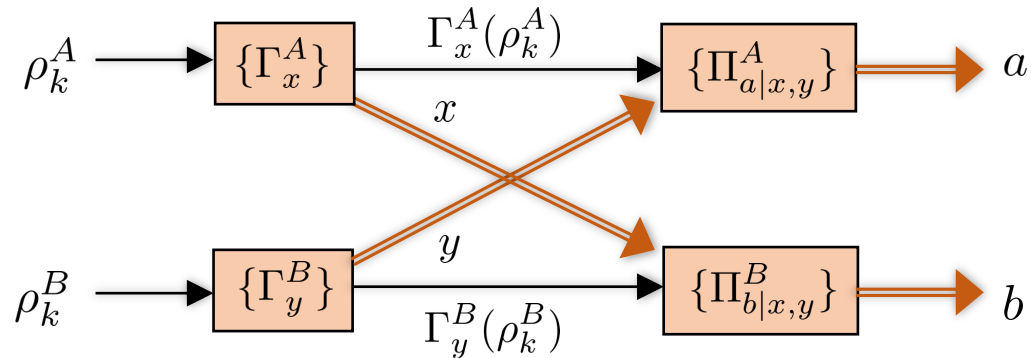
$$\{|\psi_k\rangle = |a_k\rangle^A |b_k\rangle^A\}_k$$

be distinguished by LOSCC and LOSQC?

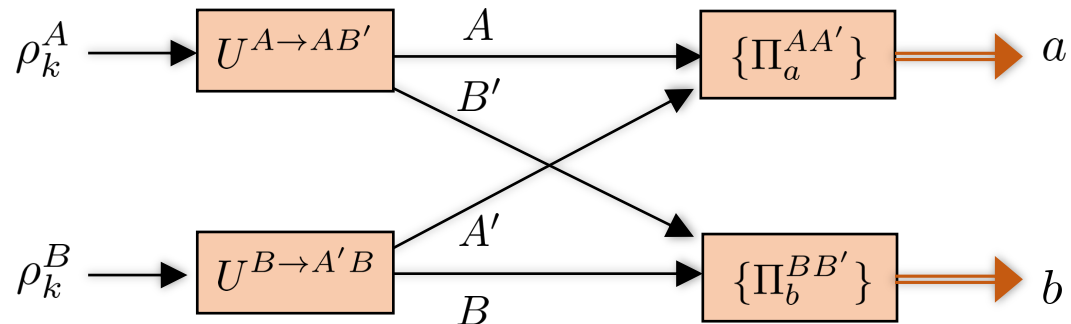


# The structure of LOSCC and LOSQC protocols

- The structure of LOSCC protocols:



- The structure of LOSQC protocols:



- Does the quantum communication help?

Yes!

Consider the symmetric and anti-symmetric projectors:

$$\rho_+^{AB'} = \frac{1}{3}(|\Phi^+\rangle\langle\Phi^+| + |\Psi^+\rangle\langle\Psi^+| + |\Phi^-\rangle\langle\Phi^-|)$$

$$\rho_-^{AB'} = |\Psi^-\rangle\langle\Psi^-|$$

Take two copies:  $\rho_+^{AB'} \otimes \rho_+^{A'B}$

$$\rho_-^{AB'} \otimes \rho_-^{A'B}$$

- Perfectly distinguishable by LOSQC but not LOSCC.

Allerstorfer, Buhrman, Speelman, Lunel, arXiv:2208.04341.

- But these involve distinguishing entangled states.

What about for product states?

# Distinguishing orthogonal product states

- This problem has a rich history in quantum information theory.
  - Any  $2 \otimes 2$  family of orthogonal product states can be perfectly distinguished by LOCC.

$$\left\{ \begin{array}{ll} |\psi_1\rangle = |0\rangle \otimes |\theta\rangle & |\psi_3\rangle = |1\rangle \otimes |\phi\rangle \\ |\psi_2\rangle = |0\rangle \otimes |\theta^\perp\rangle & |\psi_4\rangle = |1\rangle \otimes |\phi^\perp\rangle \end{array} \right.$$

Walgate and Hardy, **PRL** 89, 147901 (2002).

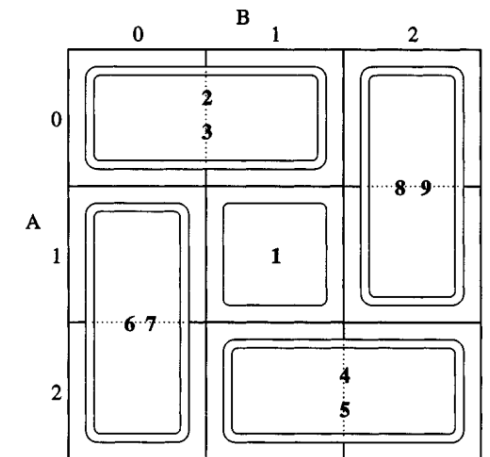
- Any  $2 \otimes n$  family of orthogonal product states can be perfectly distinguished by LOCC.

Bennett, DiVincenzo, Mor, Shor, Smolin, Terhal, **PRL** 82, 5385 (1999).

- There exists orthogonal product state that cannot be distinguished by LOCC

**“Nonlocality without entanglement”**

$$\left\{ \begin{array}{lll} |\psi_1\rangle = |1\rangle \otimes |1\rangle & |\psi_4\rangle = |2\rangle \otimes |1+2\rangle & |\psi_7\rangle = |1-2\rangle \otimes |0\rangle \\ |\psi_2\rangle = |0\rangle \otimes |0+1\rangle & |\psi_5\rangle = |2\rangle \otimes |1-2\rangle & |\psi_8\rangle = |0+1\rangle \otimes |2\rangle \\ |\psi_3\rangle = |0\rangle \otimes |0-1\rangle & |\psi_6\rangle = |1+2\rangle \otimes |0\rangle & |\psi_9\rangle = |0-1\rangle \otimes |2\rangle \end{array} \right.$$



Bennett, DiVincenzo, Fuchs, Mor, Rains, Shor, Smolin, Wootters, **PRA** 59, 1070 (1999).

# Distinguishing orthogonal product states

**Proposition** [I.George, R. Allerstorfer, P. Lunel, E.C.]:

- For perfect discrimination of  $2 \otimes 2$  orthogonal product states, LOSQC=LOSCC and the states must have the form:

$$\left\{ \begin{array}{ll} |\psi_1\rangle = |0\rangle \otimes |0\rangle & |\psi_3\rangle = |1\rangle \otimes |0\rangle \\ |\psi_2\rangle = |0\rangle \otimes |1\rangle & |\psi_4\rangle = |1\rangle \otimes |1\rangle \end{array} \right.$$

- A  $2 \otimes n$  family of orthogonal product states can be perfectly distinguished by LOSC iff it has the form:

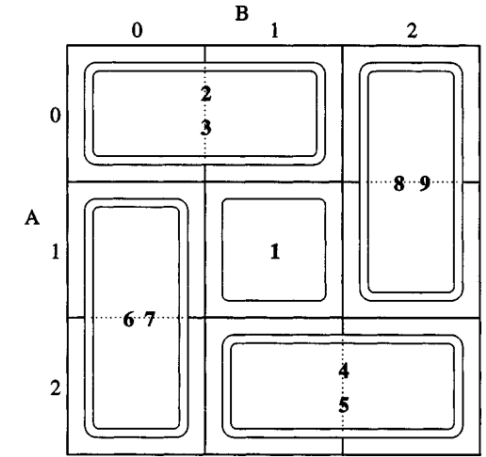
$$\left\{ \begin{array}{l} |0\rangle^A \otimes |j\rangle^B \\ |1\rangle^A \otimes (x_j |j\rangle + y_j |j+1\rangle)^B \\ |g_i\rangle^A \otimes |i\rangle^B \quad \text{for } i > 2m+1 \end{array} \right\} \quad \text{for } j \in \{0, 2, 4, \dots, 2m\}$$

- For arbitrary dimensions, the necessary and sufficient conditions are unknown! (**Open problem**)

# Distinguishing orthogonal product states

- But what about the sausage states?

$$\left\{ \begin{array}{lll} |\psi_1\rangle = |1\rangle \otimes |1\rangle & |\psi_4\rangle = |2\rangle \otimes |1+2\rangle & |\psi_7\rangle = |1-2\rangle \otimes |0\rangle \\ |\psi_2\rangle = |0\rangle \otimes |0+1\rangle & |\psi_5\rangle = |2\rangle \otimes |1-2\rangle & |\psi_8\rangle = |0+1\rangle \otimes |2\rangle \\ |\psi_3\rangle = |0\rangle \otimes |0-1\rangle & |\psi_6\rangle = |1+2\rangle \otimes |0\rangle & |\psi_9\rangle = |0-1\rangle \otimes |2\rangle \end{array} \right.$$

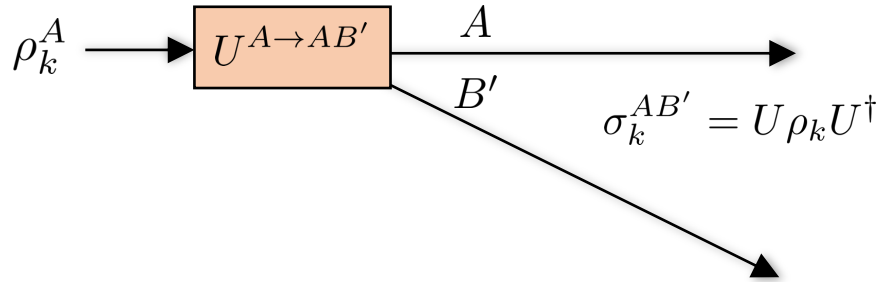


- These states cannot be distinguished by LOSCC.
- They also cannot be distinguished by LOSQC (see theorem below).
- What about two copies of the states:  $\{|\psi_k\rangle^{\otimes 2} = |a_k\rangle^{\otimes 2} \otimes |b_k\rangle^{\otimes 2}\}$ ?  $\Rightarrow$ 
  - Distinguishable by LOSQC
  - Distinguishable by LOSCC

## Conjecture:

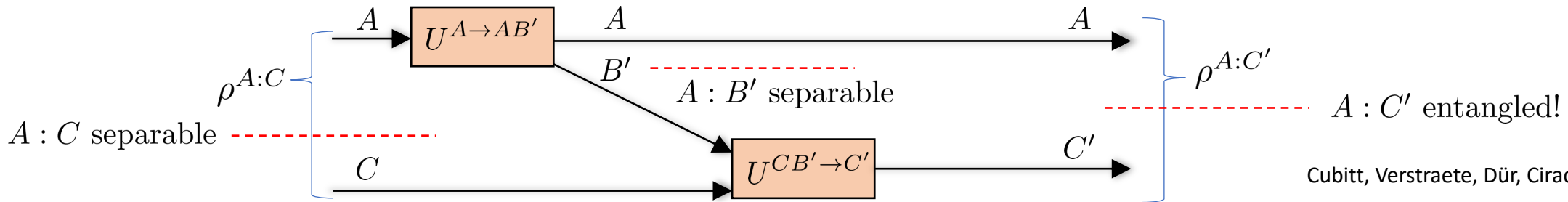
Two copies of any set of orthogonal product states is sufficient for LOSCC discrimination (or at least the ensemble must have a large number of states).

# LOSQC is more powerful than LOSCC



- Distinguish between two types of quantum communication:
  - **Separable communication**, i.e.  $\sigma_k^{AB'}$  is separable for all  $k$ .
  - **Entangled communication**, i.e.  $\sigma_k^{AB'}$  is entangled for some  $k$ .

- Separable communication can be used to perform non-classical tasks, like entanglement distribution.



Cubitt, Verstraete, Dür, Cirac, **PRL** 91, 037902 (2003).

**Theorem** [I.George, R. Allerstorfer, P. Lunel, E.C.]:

The four states can be perfectly distinguished by LOSQC only if entangled communication is used:

$$\left\{ \begin{array}{ll} |\psi_1\rangle = |0\rangle \otimes |0+1\rangle & |\psi_3\rangle = |1\rangle \otimes |0+2\rangle \\ |\psi_2\rangle = |0\rangle \otimes |0-1\rangle & |\psi_4\rangle = |1\rangle \otimes |0-2\rangle \end{array} \right.$$

# LOSQC state discrimination with error

- Perfect state discrimination is interesting from a fundamental perspective, but not for practical QPV.
- **QPV question:**

Given an ensemble  $\{|\psi_k\rangle\}_k$ , what is the smallest error probability in state discrimination using LOSQC?

**Theorem** [I.George, R. Allerstorfer, P. Lunel, E.C.]:

Let  $\{|\psi_k\rangle^{AB} = |a_k\rangle^A |b_k\rangle^B\}_k$  be an ensemble of product states that contains four states of the form

$$\begin{aligned} |\psi_0\rangle^{AB} &= |a_0\rangle^A |b_0\rangle^B, \\ |\psi_1\rangle^{AB} &= |a_1\rangle^A |b_1\rangle^B, \\ |\psi_2\rangle^{AB} &= |a_2\rangle^A (\cos \theta |b_0\rangle + e^{i\phi} \sin \theta |b_1\rangle)^B \\ |\psi_3\rangle^{AB} &= |a_3\rangle^A (\cos \theta |b_0\rangle - e^{i\phi} \sin \theta |b_1\rangle)^B, \end{aligned}$$

with  $\langle a_0 | a_1 \rangle \neq 0$ . Suppose Alice and Bob can identify each state with at least probability  $1 - \epsilon$  using some LOBQC protocol. Then

$$2\epsilon + \frac{4\sqrt{\epsilon(1-\epsilon)}}{|\langle a_0 | a_1 \rangle|^2} + \sqrt{1 - |\langle a_2 | a_3 \rangle|^2} > 1.$$

# LOSQC state discrimination with error

**Example:** Generalized BB84 states:

$$|\psi_0\rangle^{AB} = |0\rangle^A \otimes |0\rangle^B,$$

$$|\psi_1\rangle^{AB} = |0\rangle^A \otimes |1\rangle^B,$$

$$|\psi_2\rangle^{AB} = |1\rangle^A \otimes (\cos \theta |0\rangle + e^{i\phi} \sin \theta |1\rangle)^B$$

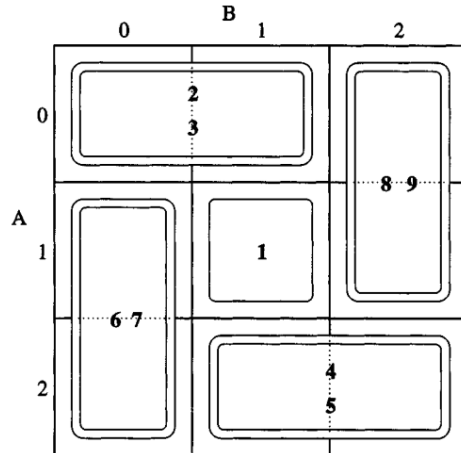
$$|\psi_3\rangle^{AB} = |1\rangle^A \otimes (\cos \theta |0\rangle - e^{i\phi} \sin \theta |1\rangle)^B$$

The LOSQC error probability  $P_{err}$  is lower bounded as:  $P_{err} > \frac{1}{4} \left( \frac{1}{2} - \frac{1}{\sqrt{5}} \right) \approx 1.3\%.$

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- But what about the sausage states?

**Example:**



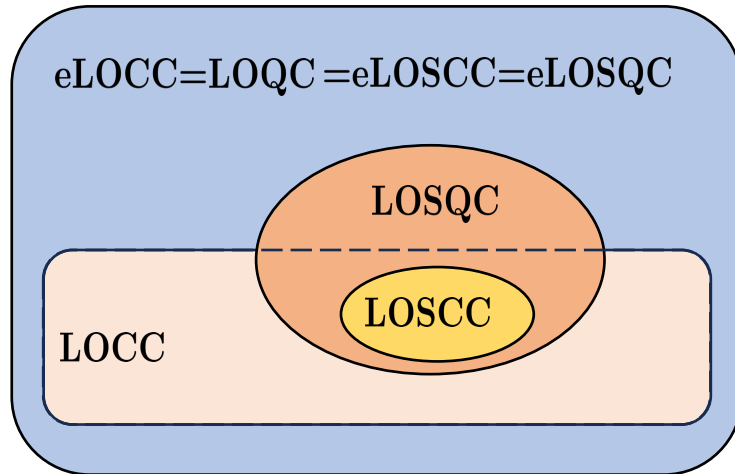
The LOSQC error probability  $P_{err}$  is lower bounded as:

$$P_{err} > \frac{1}{9} \left( \frac{1}{2} - \frac{2}{\sqrt{17}} \right) \approx .16\%.$$

# Open problems and future directions

- What are the necessary and sufficient conditions for product state discrimination under LOCC and LOSQC?
- **Copy complexity:** How many copies of an ensemble state do Alice and Bob need before they can perfectly discriminate by LOCC?

$$\{|\psi_k\rangle^{\otimes n} = |a_k\rangle^{\otimes n} \otimes |b_k\rangle^{\otimes n}\}$$



- What families of states are distinguishable by LOSQC but not LOCC?
- **Most important question for QPV:**  
What are the entanglement costs for state discrimination under eLOSCC and eLOSQC?

- **Example: BB84 states:**

$$\left\{ \begin{array}{ll} |\psi_1\rangle^{AB} = |0\rangle^A \otimes |0\rangle^B & |\psi_3\rangle^{AB} = |+\rangle^A \otimes |1\rangle^B \\ |\psi_2\rangle^{AB} = |1\rangle^A \otimes |0\rangle^B & |\psi_4\rangle^{AB} = |-\rangle^A \otimes |1\rangle^B \end{array} \right.$$

One ebit suffices for perfect discrimination

Lo and Lau **PRA** 83, 012322 (2011).



**Thank You!**

