Recent Developments in **Supersymmetric Dualities**

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Rabbit or Duck?

- Duality—the same phenomenon, multiple descriptions
- Essential roles in various fields, such as high energy and condensed matter physics
- Many different kinds known—today: IR duality of QFT

Welche Thiere gleichen ein= ander am meisten?



Kaninchen und Ente.

"Rabbit and Duck," Fliegende Blätter (23 Oct 1892)

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- Introduction
 - Rabbit or duck?
 - Renormalization
 - Web of dualities Example I, II, III
- What are the most fundamental dualities?
- Example \bullet
- Conclusion



Renormalization

Renormalization in QFT

- **Couplings in QFT** varying along the scale
- Different dynamics at different scales
- Rich low-energy dynamics from an almost free theory
- "QCD" in the Standard Model—an asymptotically-free theory (i.e., free in the high energy) exhibiting confinement in the low energy
- Understanding a QFT -> understanding physics at different scales



Renormalization in QFT

- A conformal field theory at each end of the RG-flow
- Invariant under the scale transformation (e.g., a free theory)
- Interesting emergent phenomena in strongly coupled CFTs
- E.g., symmetry enhancement, duality, gravity, ...
- Still difficult to solve; nevertheless, more universal control
- Especially with supersymmetry





SCFTs and Dualities

- SCFTs QFTs preserving a **superconformal** symmetry
- Interacting SCFTs: deformation of free theories (Lagrangian constructions), string theory
 constructions (strongly coupled), or
- SCFTs from SCFTs
 - Adding interactions
 - Gauging symmetries
 - Compactification
- Fixed points of RG-flows
- An Infra-Red (IR) duality—the same IR SCFT from multiple UV SCFTs



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Web of dualities

Example I: 3D Gauge-Scalar Duality

The 3d Maxwell theory

$$Z_{Maxwell} = \int \mathscr{D}A_{\mu} \exp\left[\int d^3x \,\frac{1}{4e^2} \,F_{\mu\nu} F^{\mu\nu}\right]$$

• The field-strength satisfies the Bianchi identity:

• Integrating over $F_{\mu\nu}$, $Z_{Maxwell}$ can be written as follows:

$$Z_{Maxwell} = \int \mathscr{D}F_{\mu\nu} \mathscr{D}\gamma \exp\left[\int d^3x \left(\frac{1}{4e^2}F_{\mu\nu}F^{\mu\nu} - \frac{i}{4\pi}\gamma \epsilon^{\mu\nu\rho}\partial_{\mu}F_{\nu\rho}\right)\right]$$
$$= \int \mathscr{D}\gamma \exp\left[\int d^3x \frac{e^2}{8\pi^2}\partial_{\mu}\gamma \partial^{\mu}\gamma\right]$$

 $\epsilon^{\mu\nu\rho}\,\partial_{\mu}F_{\nu\rho}=0$

• A simple duality example between the 3d Maxwell theory and a scalar theory, both of which are free.

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$$P \partial_{\mu} F_{\nu\rho} = 0$$

 \blacksquare A conserved current $J^{\mu}_{top} = \frac{1}{4\pi} \epsilon^{\mu\nu\rho} F_{\mu\nu}$

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$$= \int \mathscr{D}\gamma \exp\left[\int d^3x \frac{e^2}{8\pi^2} \partial_{\mu}\gamma \partial^{\mu}\gamma\right] \qquad \qquad J^{\mu}_{top} = \frac{ie^2}{(2\pi)^2} \partial^{\mu}\gamma \qquad \clubsuit \qquad \gamma + \alpha$$



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Example II: 3D Vortex-Particle Duality

Generalization to interacting theories: the 3d bosonization

$$i\overline{\chi}\gamma_{\mu}D_{a}^{\mu}\chi - \frac{1}{2\pi}Bda - \frac{1}{4\pi}BdB \qquad \longleftrightarrow \qquad D_{B}\phi^{2} - \phi^{4}$$

• A number of new dualities derived; e.g., the vortex-particle duality

$$D_{b}\phi^{2} - \phi^{4} + \frac{1}{2\pi}bdC \qquad \longleftrightarrow \qquad i\overline{\chi}\gamma_{\mu}D_{a}^{\mu}\chi - \frac{1}{2\pi}bda - \frac{1}{4\pi}bdb + \frac{1}{2\pi}bdC \qquad \longleftrightarrow \qquad D_{C}\hat{\phi}^{2} - \hat{\phi}^{4}$$

- Adding $\frac{1}{2\pi}BdC$ + gauging $B \rightarrow b$ + the (time-reversed) bosonization
- bosonization is more fundamental than the vortex-particle duality.

• The bosonization -> the vortex-particle duality, but not the other way around; namely, the

- Generalization to supersymmetric theories
- The simplest case:

 - A free hypermultiplet



Multiple flavors \bullet

$$1$$
 N_f





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 \ni a pair of chiral multiplets with charge 1 & -1 each







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- Generalization to supersymmetric theories
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Multiple flavors \bullet

$$(1)$$
 N_f

Flavor $U(1) \times SU(N_f)$

 \ni a pair of chiral multiplets with charge 1 & -1 each













Web of IR Dualities

Hwang-Pasquetti-Sacchi duality Comi-Hwang-Marino-Pasquetti-Sacchi duality

Intriligator-Seiberg duality **Dorey-Tong duality** Gaiotto-Witten duality of $T_{\rho}^{\sigma}[G]$

. . .

3d bosonization Particle-vortex duality Karch-Robinson-Tong duality

. . .

2d bosonization

Fateev-Zamolodchikov^2 duality

. . .

Hori-Vafa duality Hori-Kapustin duality

. . .





Web of IR Dualities



Seiberg-like dualities



Web of IR Dualities

Mirror-like dualities

What are the most fundamental dualities?

Most Fundamental Dualities?

- Why do those dualities work?
- No single (interacting) duality has been proven. (Not surprising since we don't know how to handle strongly coupled QFTs in general.)
- Are they really different?—Own microscopic mechanism each or any universal mechanism?
- If the latter is the case, what are the most fundamental dualities?

Duality of Theories vs Duality of Fields

- A QFT duality relates a "theory" to another "theory".
- Many of them are motivated by string theory, but no systematic construction in the QFT language.
- **Proposal:** dualize a "field" rather than a theory!
- The duality of a "theory" can be obtained by gauging symmetries (and adding interactions if necessary) as in the previous examples.

- - Consisting of (bi-)fundamental matter fields
 - Interactions fixed by the $\mathcal{N} = 4$ supersymmetry

Based on \bullet

• Today let's focus on mirror symmetry of 3d $\mathcal{N} = 4$ linear quiver gauge theories:

- R. Comi, CH, F. Marino, S. Pasquetti, M. Sacchi, "The SL(2, Z) dualization algorithm at work," JHEP 06 (2023) 119, [arXiv:2212.10571]. - CH, S. Pasquetti, M. Sacchi, "Rethinking mirror symmetry as a local duality on fields," Phys.Rev.D 106 (2022) 10, 105014, [arXiv:2110.11362]. - L. E. Bottini, CH, S. Pasquetti, M. Sacchi, "4d S-duality wall and SL(2, \mathbb{Z}) relations," JHEP 03 (2022) 035, [arXiv:2110.08001].

Basic Ingredients: S-Wall

• The S-wall theory: T[U(N)]



- The S-duality domain-wall theory of the 4d $\mathcal{N} = 4$ SYM (Giaotto-Witten 08)
- $U(1)^{N-1}$ topological symmetry + background U(1) coupled via mixed CS + $U(N)_{Y}$ flavor symmetry
- Enhanced $U(N)_X \times U(N)_Y$ symmetry in the IR



$$N$$
 \overrightarrow{Y}

Basic Ingredients: I-Wall

• The identity-wall theory: two S-walls glued by gauging common U(N)

$$\vec{X}$$
 N $- N$ \vec{Y} = $\left(\begin{array}{c} N \\ N \end{array} \right)$

The partition function proportional to the delta function

$$\sim \sum_{\sigma \in S_N} \prod_{j=1}^N \delta\left(X_j - Y_{\sigma(j)}\right)$$

 \bullet

 $U(N)_X \times U(N)_Y \rightarrow U(N)_D$



 $\langle \mathfrak{M} \rangle \neq 0$, where \mathfrak{M} is a (monopole) operator in the $U(N)_X \times U(N)_Y$ bifund. rep., breaking

• Dualization of a $U(N) \times U(N)$ bifundamental hypermultiplet



- A bifundamental hypermultiplet -> a fundamental (twisted) hypermultiplet
- Topological $U(1) \rightarrow \text{flavor } U(1)$
- A QFT version of the S-transformation in IIB string theory, exchanging NS5 and D5



Building Blocks of 3D Mirror Symmetry

• Dualization of a **fundamental** hypermultiplet (+ an identity wall)



- Obtained by using the I-wall property
- A fundamental hypermultiplet -> a bifundamental (twisted) hypermultiplet
- Flavor U(1) -> topological U(1)



Generalization: Mass-Deformed S- & I-Walls



 \bullet



• Mass-deformed S-wall: S-wall + mass terms breaking $U(N)_Y \rightarrow U(M) \times U(1)$

$$Y_{M+j} = V + \frac{N - M + 1 - 2j}{2} (iQ - 2m_A), \qquad j = 1, ..., N - 2$$





(Generalized) Building Blocks of 3D Mirror

Dualization of a $U(N) \times U(M)$ bifundamental hypermultiplet ullet



hypermultiplet dualized by the mass-deformed S-wall



• A bifundamental hypermultiplet between different ranks = a fundamental (twisted)

Swap Fundamental and Bifundamental

- the Hanany-Witten brane move in IIB string theory.
- Nothing but Higgs mechanism \bullet



The mass-deformed S- & I-walls satisfy an interesting property resembling





Example

• Mirror symmetry of the 3d $\mathcal{N} = 4 U(2)$ theory with 4 flavors



- $U(1) \times SU(4)$ vs $SU(2) \times U(1)^3$
- The global symmetry enhanced to $SU(2) \times SU(4)$ in the IR
- Monopoles and mesons exchanged under the duality



- Dualization Algorithm
 - Chop the quiver into the basic blocks.
 - Dualize each block and glue them back.
 - Carry out the (QFT version of) Hanany-Witten move until I-walls disappear.











The 1st HW move







The 2nd HW move



$$Z_1 \quad --- \quad W_1$$
$$Z_2 \quad --- \quad W_2$$

The permutations belong to the Weyl of U(2)





Collapsing I-walls



$$Z_1 \quad --- \quad W_1$$
$$Z_2 \quad --- \quad W_2$$

The permutations belong to the Weyl of U(2)





One more thing...

Where do these building blocks come from?

Where do these building blocks come from?

- The 3d mirror building blocks are derived from the Aharony duality!
- The Aharony duality:



- A new bridge between mirror-like dualities and Seiberg-like dualities
- The same relation in 4d





Seiberg-like dualities



Web of IR Dualities

Mirror-like dualities



Seiberg-like dualities



Web of IR Dualities

Mirror-like dualities



- Generalization \bullet
 - 3d $SL(2,\mathbb{Z})$ dualities
 - 4d mirror and $PSL(2,\mathbb{Z})$ dualities
 - Some Seiberg-like dualities
- These building blocks can be derived from the Aharony duality.
- Different looking dualities can be derived from common fundamental dualities. \bullet -> Microscopic mechanism of these dualities would be universal.
- dimensional SCFTs? Holography?

• We have found a new properties of T[U(N)], which can be used to find building blocks of 3d mirror symmetry.

Deeper understanding of the duality building blocks from other perspectives? E.g., connection to higher





Congratulation CTPU!