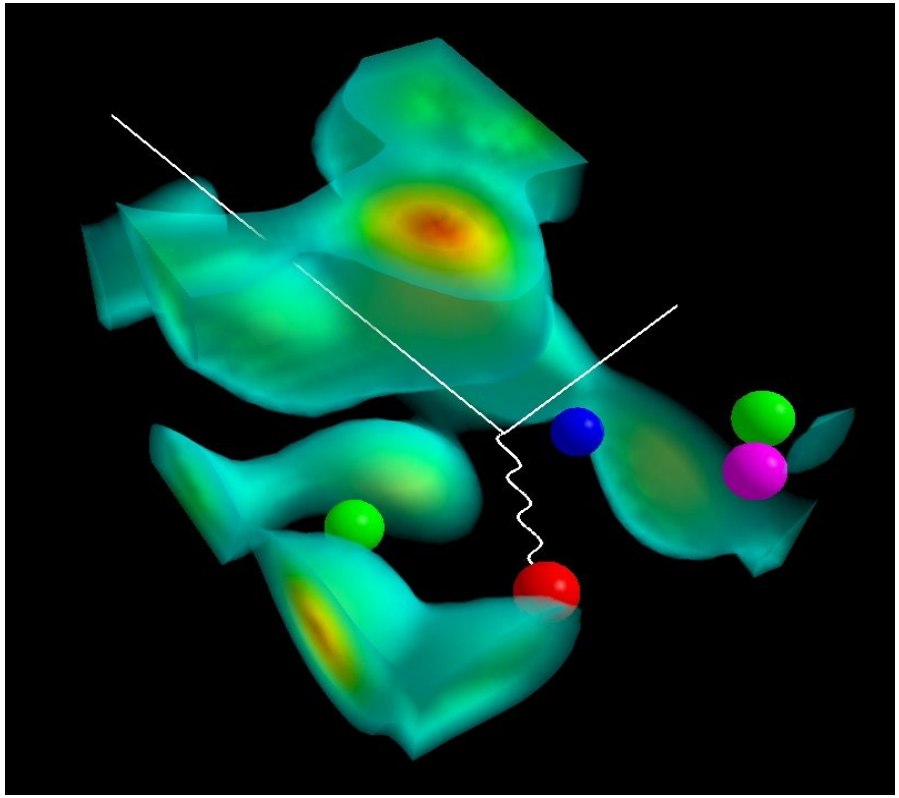
New Insights into the Baryon Spectrum



Anthony W. Thomas

INPC 2025 : Daejeon Korea

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SPECIAL RESEARCH CENTRE FOR THE SUBATOMIC STRUCTURE OF MATTER

Outline

I. Excitations of the nucleon octet are a vital piece of the challenge to understand how hadrons are made in QCD

II. New Insight into the Quark Model

- The Λ(1405) IS a $\overline{K}$ - N bound state
- The Roper IS generated by πN-σN-πΔ re-scattering
- But not all states are dynamically generated e.g. N(1535)!

III. The Quark Model is not so bad!

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A brief and very incomplete history

- Dalitz began to develop quark model before quarks were “real”
- Chodos, Jaffe, Johnson &Thorn, 1974+: MIT bag
- De Rujula, Georgi & Glashow, 1975: color hyperfine gluon exchange
- Isgur and Karl, 1977+: developed “shell model” description of baryons using harmonic oscillator with color hyperfine
  - but no role for chiral symmetry
- Capstick et al. : relativistic extension of Isgur-Karl

History cont.

- **Dyson-Schwinger equations: Roberts *et al.***
- **Coupled channels work: ANL, Jülich, Kyoto, Tokyo IT, Valencia…**
- **χPT; holographic "QCD"**
- **Hundreds of others (apologies) with an excellent recent review by Dȍring, Haidenbauer, Mai and Sato, *arXiv: 2505.02745 (123 pp.)***
- **Tremendous amount of data from CLAS**
- **Recent reviews of experimental data:**
  - **Thiel *et al*., *PPNP 125 (2022) 103949***
  - **Ireland *et al., PPNP 111 (2019) 103752***

Spectroscopy

- How do excited states emerge from QCD ?
- What are the fundamental degrees of freedom ?
- Lattice QCD provides extremely valuable information

Resonances are very complicated – but the lattice is less so….

- Everything is stable on the lattice
  – an eigenstate of the QCD Hamiltonian

- Whereas real resonances decay like crazy.....

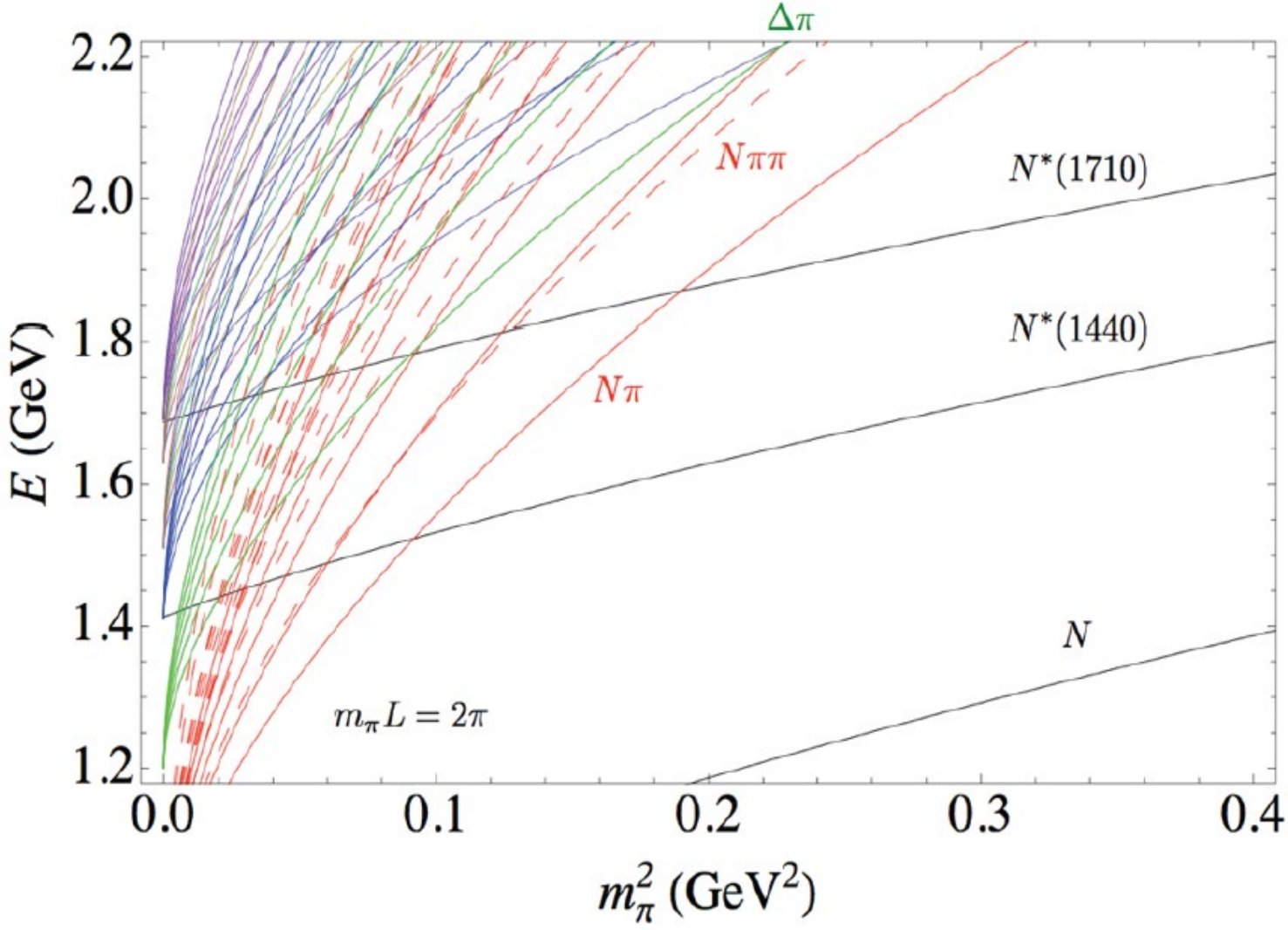
- Lüscher has a method to derive phase shifts at discrete energies when there is one open channel

- That approach has been generalized to coupled channels by Hansen and Sharpe (Phys. Rev. D86 (2012) 016007) and Lellouch and Lüscher (Comm. Math. Phys., 219 (2011) 31), as well as moving frames, e.g. Li et al., (Phys. Rev. D 103 (2021) 094518) BUT it becomes very complicated

- Excited states involve <u>many</u> coupled channels

In General: Multiple open channels

Illustration of multiple non-interacting energy levels on a lattice



and then there are: σN, ωN, ρN etc....

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Lüscher procedure does not solve all our problems

- **This procedure gives no more information than we get from a phase shift analysis of experimental data**

- **It provides no new insight into the nature of the excited states**

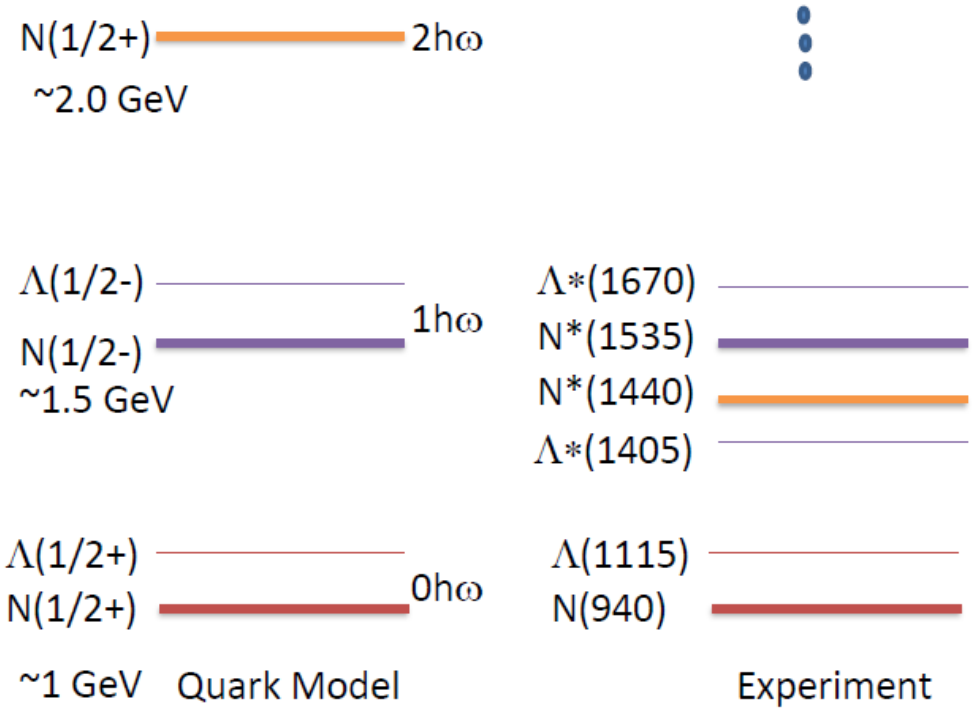
Longstanding Challenges

Quark Model looks WRONG

First +ve parity state occurs before first –ve parity state

Excited state with a heavier strange quark occurs below even the first nucleon excited state

A number of expected excited states are not observed



The Λ(1405)

- We have unambiguous evidence that it is a $\overline{K}$ - N bound state! 50 years after speculation by Dalitz *et al*.
- To be fair Dalitz had no quark model then so there was not much else it could be at that time.
- Rather than the Lüscher method we apply Hamiltonian Effective Field Theory
  - shown to be equivalent for phase shifts*
  - BUT also provides information on eigenstates
- Fit experimental data and carry out a Hamiltonian analysis of lattice data
- Examine the strange magnetic form factor of Λ(1405)

* Wu et al., Phys. Rev. C 90 (2014) 5, 055206

First calculation after QCD incorporating chiral symmetry

PHYSICAL REVIEW D VOLUME 31, NUMBER 5 1 MARCH 1985

S-wave meson-nucleon scattering in an SU(3) cloudy bag model

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R. C. Barrett
Physics Department, University of Surrey, Guildford GU2 5XH, United Kingdom
(Received 8 June 1984)

The cloudy bag model (CBM) is extended to incorporate chiral SU(3)×SU(3) symmetry, in order to describe *S*-wave KN and $\bar{K}N$ scattering. In spite of the large mass of the kaon, the model yields reasonable results once the physical masses of the mesons are used. We use that version of the CBM in which the mesons couple to the quarks with an axial-vector coupling throughout the bag volume. This version also has a meson-quark contact interaction with the same spin-flavor structure as the exchange of the octet of vector mesons. The present model strongly supports the contention that the $\Lambda^*(1405)$ is a $\bar{K}N$ bound state.

But now, through the lattice, we can use QCD itself

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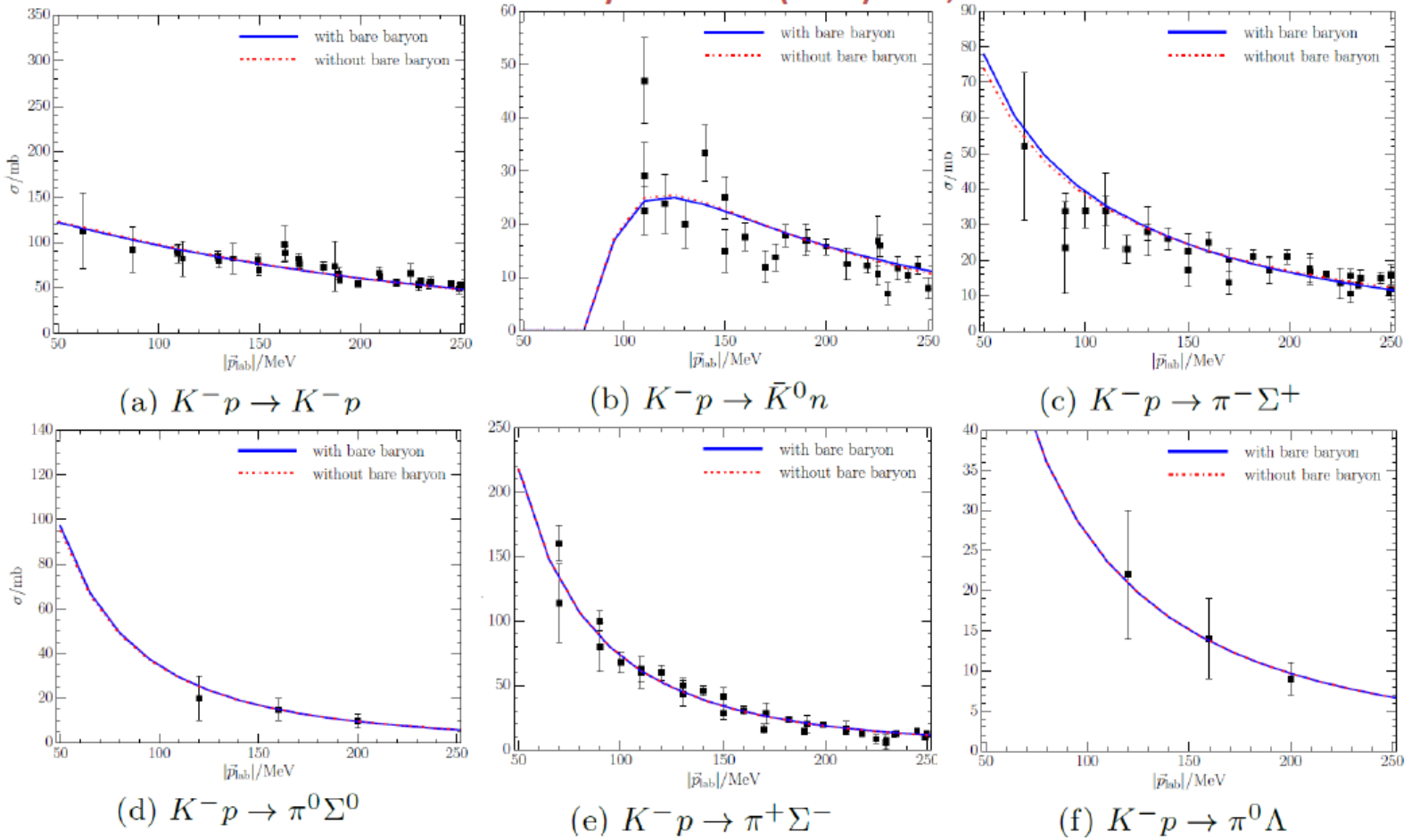
Hamiltonian fit to existing data

Zhan-wei Liu etc. Phys.Rev. D95 (2017) no.1, 014506



Include πΣ, $\bar{K}$N, ηΛ and KΞ channels

Similar work by Valencia, Bonn, JLab and other groups

Find the same two-pole structure as other analyses

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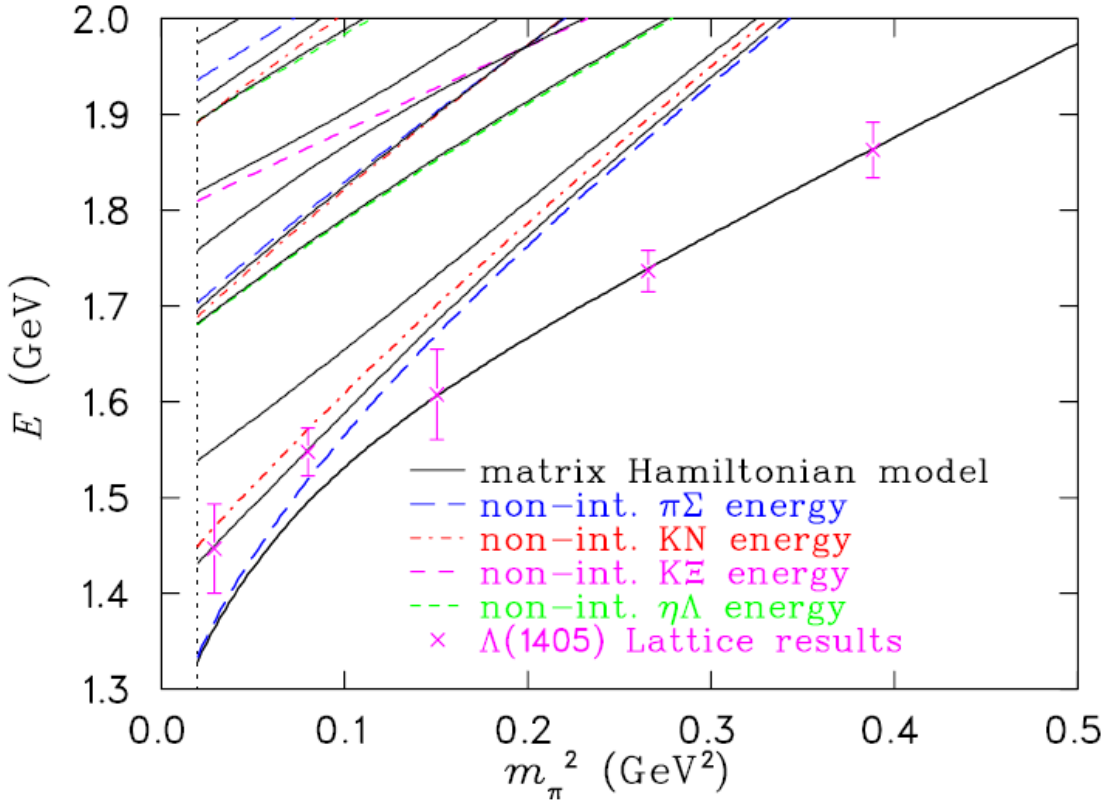
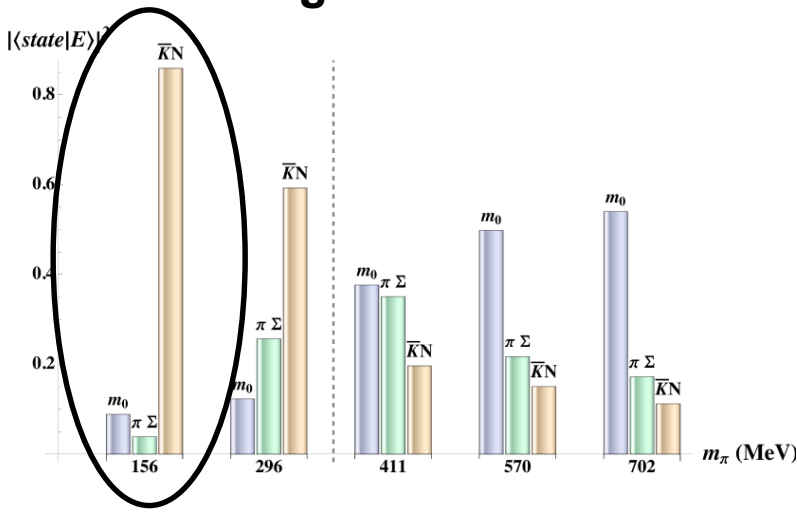
Low lying negative parity state : Λ(1405)

Clear evidence that it is a Kbar-N bound state

Dashed lines: free energy levels; solid: eigenstates of Hamiltonian



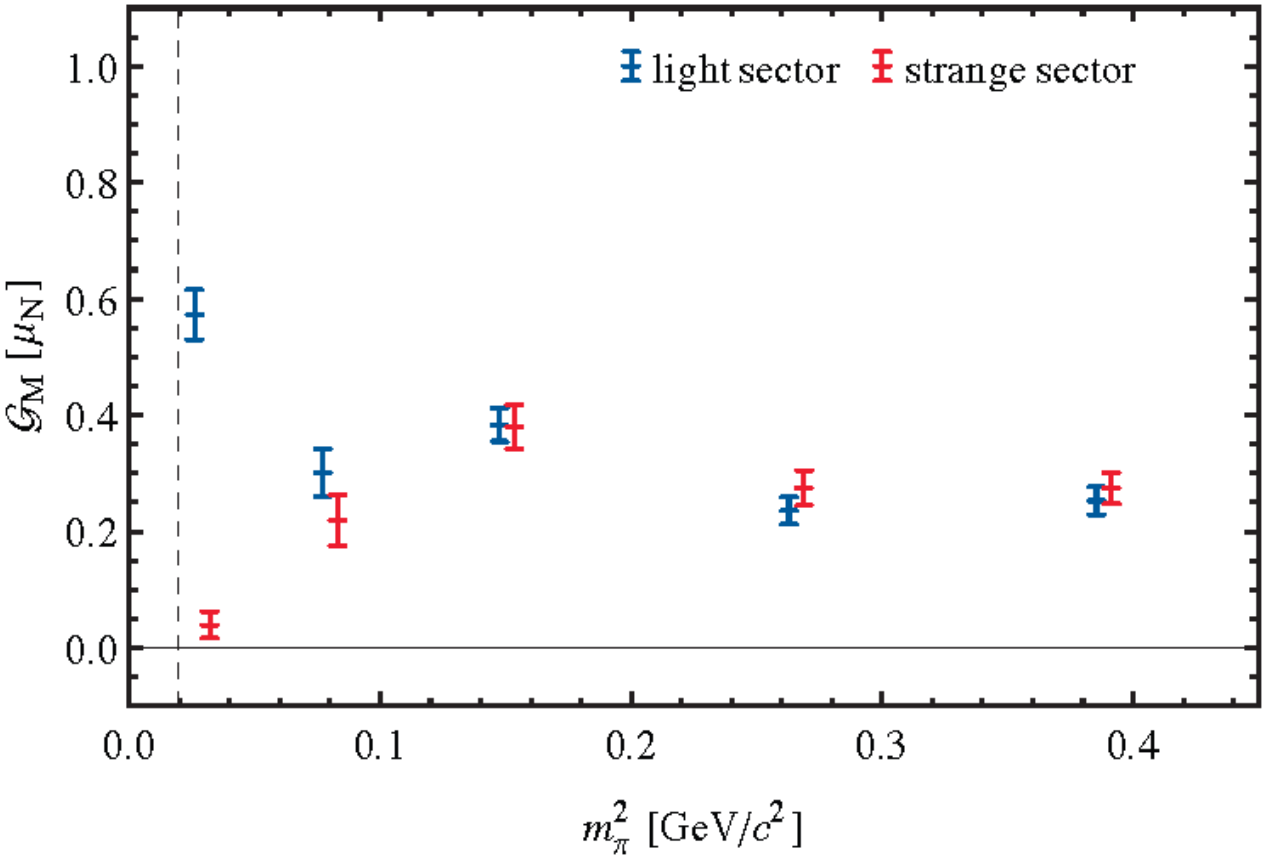
Hamiltonian approach allows one to examine the eigenstates:



Hall, Leinweber, Menadue, Young, AWT
– Phys. Rev. Lett. 114 (2015) 13

Lattice Magnetic Form Factor Calculations

- **Calculation of the individual quark contributions to the magnetic form factor confirms that it is a $\overline{K}$ - N bound state**



Only an L=0 $\overline{K}$ - N state gives vanishing strange moment

Hall *et al.*, Phys. Rev. *D* 95 (2017) 5, 054510

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Note that Lattice QCD allows

us to study hadron structure IN QCD as a

function of quark mass – a powerful tool*

*AWT, Nucl. Phys. B Proc. Suppl. 119 (2003) 50

Roper Resonance (Discovered 1964)

Investigating the nature of $N(1535)$ and $\Lambda(1405)$ in a quenched chiral quark model

Yue Tan[a], Zi-Xuan Ma[b], Xiaoyun Chen[c], Xiaohuang Hu[d], Youchang Yang[e], Qi Huang[b], Jialun Ping[b]*
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[c]*College of Science, Jinling Institute of Technology, Nanjing 211169, People's Republic of China*
[d]*Department of Physics, Changzhou Vocational Institute of Engineering, Changzhou 213164, People's Republic of China and*
[e]*Department of Physics, Guizhou University of Engineering Science, Bijie 551700, People's Republic of China.*
(Dated: March 7, 2025)

In this work, we systematically study $N(1440)$, $N(1535)$, and $\Lambda(1405)$ in both the quenched three-quark and five-quark frameworks using the Gaussian Expansion Method (GEM) within the chiral quark model. Our calculations show that $N(1535)$ can be reproduced as a three-quark state ($N(1P)$), while $N(1440)$ and $\Lambda(1405)$ cannot be accommodated as the three-quark candidates, ($N(2S)$ and $\Lambda(1P)$), respectively. In the five-quark framework, we find that the ΛK state for $N(1535)$ can not form a bound state, while in the $N\bar{K}$ channel there will $\Lambda(1405)$ form a shallow bound state. Based

2503.04586

For the Roper this mirrors conclusion of Jülich group two decades ago

PHYSICAL REVIEW C, VOLUME 62, 025207

What is the structure of the Roper resonance?

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[1]*Institut für Kernphysik, Forschungszentrum Jülich GmbH, D-52425 Jülich, Germany*
[2]*Department of Physics and INT, University of Washington, Seattle, Washington 98195*
(Received 30 November 1999; published 24 July 2000)

We investigate the structure of the nucleon resonance $N^*(1440)$ (Roper) within a coupled-channel meson exchange model for pion-nucleon scattering. The coupling to $\pi\pi N$ states is realized effectively by the coupling to the σN, $\pi\Delta$, and ρN channels. The interaction within and between these channels is derived from an effective Lagrangian based on a chirally symmetric Lagrangian, which is supplemented by well known terms for the coupling of the Δ isobar, the ω meson, and the ''σ,'' which is the name given here to the strong correlation of two pions in the scalar-isoscalar channel. In this model the Roper resonance can be described by meson-baryon dynamics alone; no genuine $N^*(1440)$ (three quark) resonance is needed in order to fit πN phase shifts and inelasticities.

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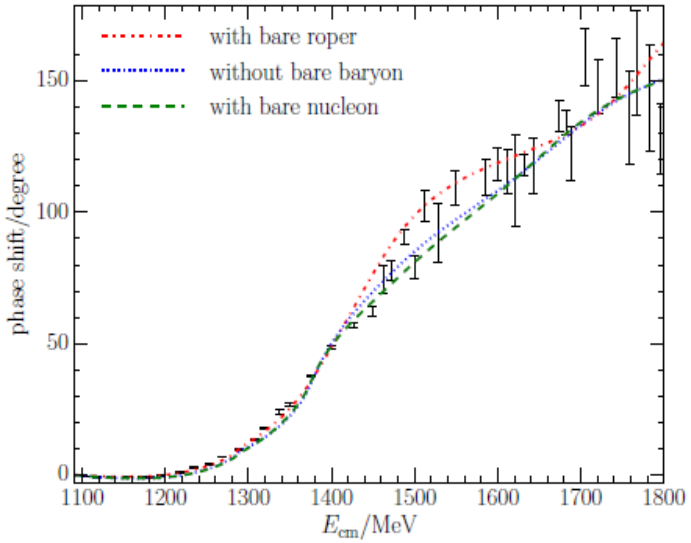
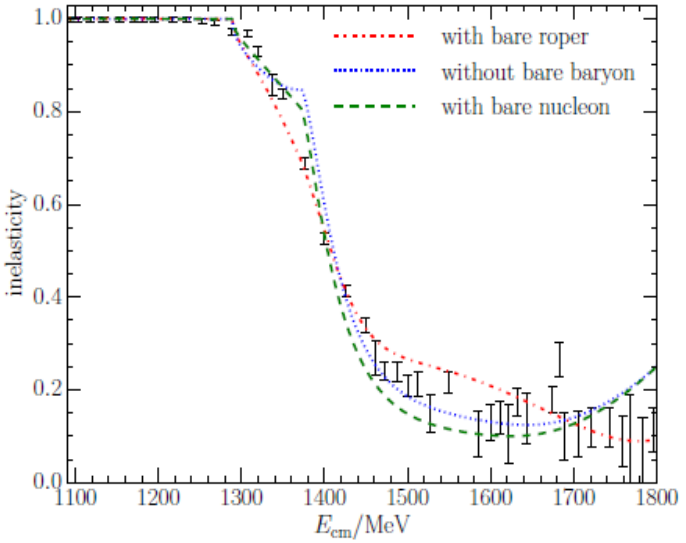
Roper Resonance and Lattice QCD

This has long been a challenge for the quark model, as it is the 1st positive parity excited state and lies below the N(1535), the 1st negative parity state

Bare Roper Case: $m_0 = 2.03$ GeV

- Consider πN, $\pi\Delta$ and σN channels, dressing a bare state.
- Fit to phase shift and inelasticity





- Fit yields a pole at 1380 $- i\,87$ MeV.
- Compare PDG estimate $1365 \pm 15 - i\,95 \pm 15$ MeV.

Wu et al., Phys. Rev. D 97 (2018) 9, 094509
- use Hamiltonian Effective Field Theory (HEFT)

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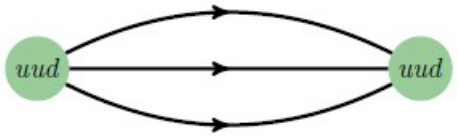
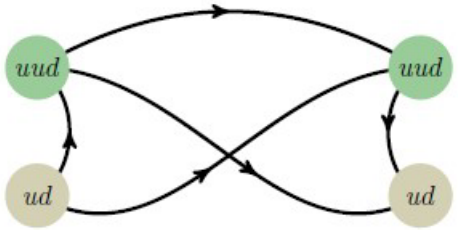
3-quark correlator mainly yields a localized state

This is useful in understanding the nature of a resonance



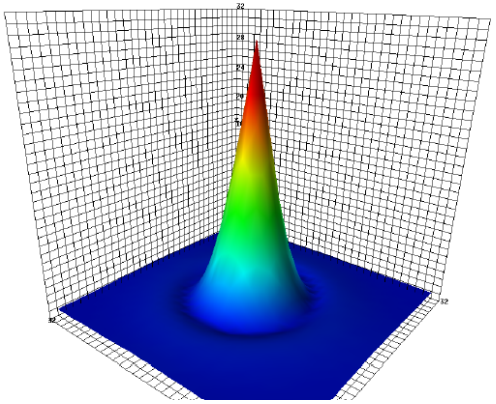
Need explicit meson-baryon correlator to efficiently excite composite states

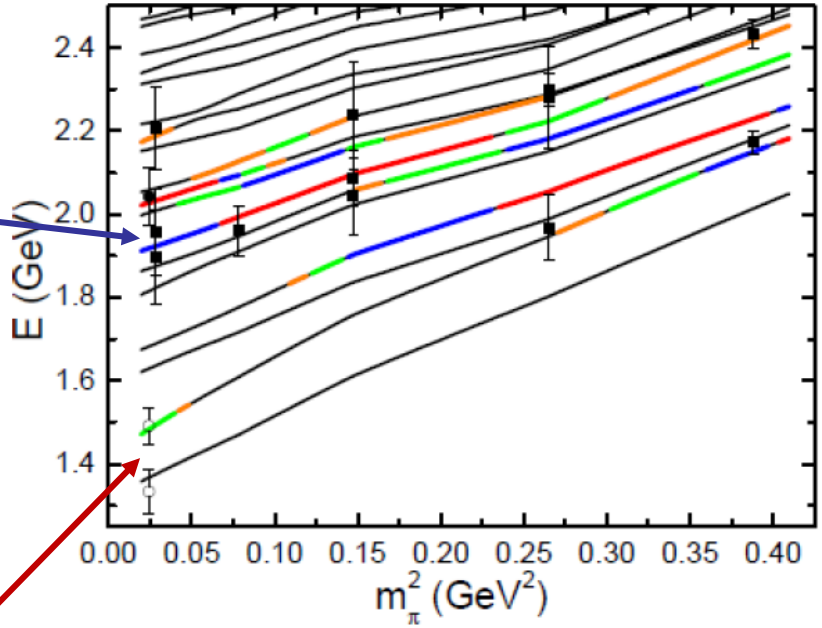


Leskovec, Lang, Padmanath and Prelovsek, 1806.02363 only saw a state in Roper region with explicit Nσ included

Comparison of HEFT Results with Lattice Energy Levels

- Blue indicates high "bare state" (i.e. 3-quark) content: this matches the lowest excited state found with a 3-quark interpolating field and looks Like a 2s state





- Lattice calculations of Padmanath, Lang et al., Phys. Rev. D 95, 014510 (2017), using baryon-meson interpolating fields, especially Nσ, crucial to finding state In Roper mass region

- Matched by Hamiltonian levels but with little or no 3-quark content

The first scenario with a bare state for P11 around the pole at 2.0 GeV can fit both Lattice data and experimental data well, it indicates that N*(1440) seems a molecule state, and first radial excitation of nucleon should be around 2.0 GeV.

Wu et al., Phys. Rev. D 97 (2018) 9, 094509

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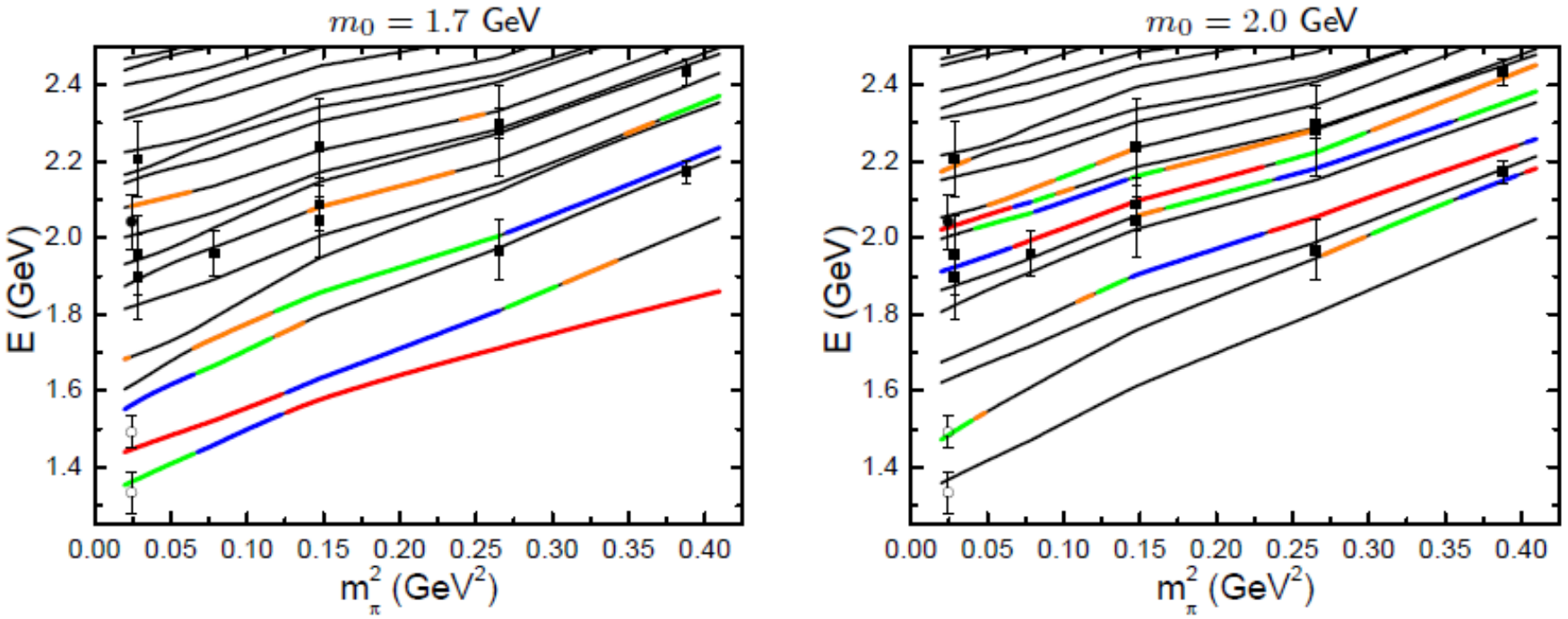
To emphasise the point

Two different descriptions of the Roper resonance



(left) Meson dressings of a quark-model like core.
(right) Resonance generated by strong rescattering in meson-baryon channels.

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Slide from Derek Leinweber

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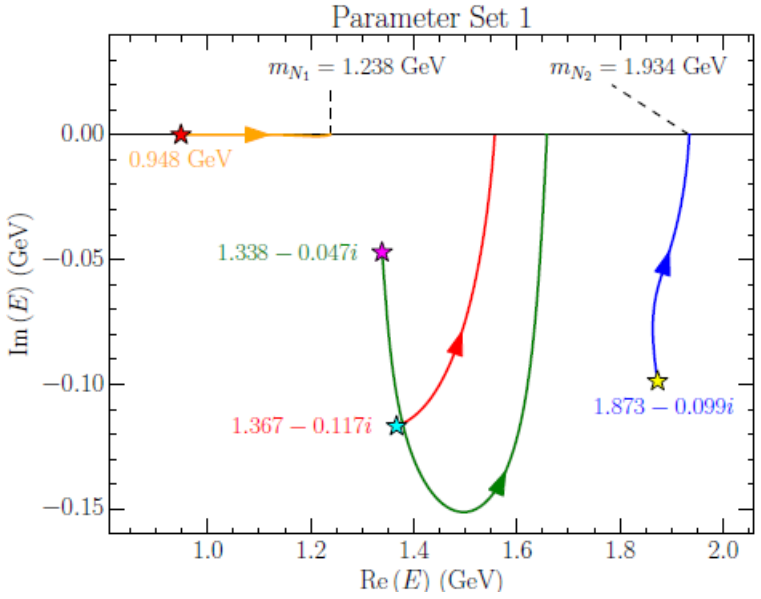
Clear conclusion is that the Roper is dynamically generated by coupling to the Nσ and Δπ channels

Movement of Poles is Model Dependent

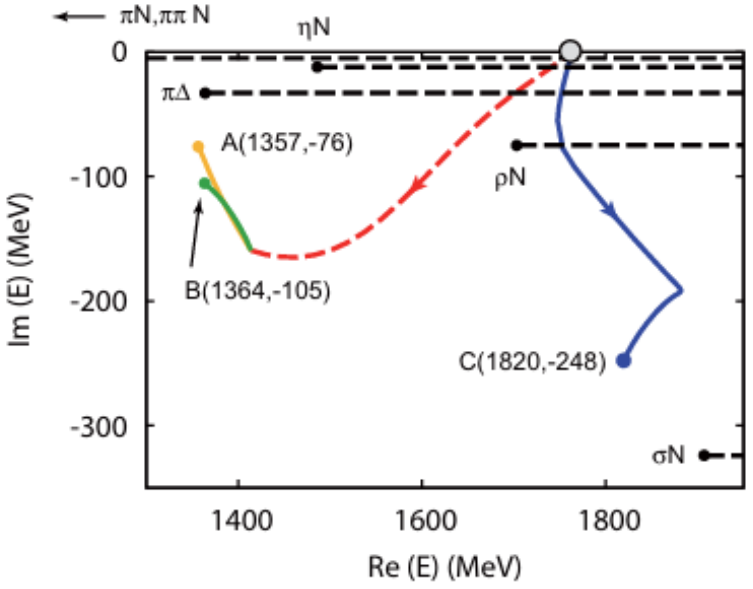
Set	Location	Our value (GeV)	PDG value (GeV)
1	$\{u,p,p,u\}$	$1.367-0.117i$	$1.365(15)-0.095(15)i$
	$\{u,u,p,u\}$	$1.338-0.047i$	
	$\{u,u,u,u\}$	$1.873-0.099i$	$1.860(40)-0.115(25)i$
2	$\{u,p,p,u\}$	$1.363-0.120i$	$1.365(15)-0.095(15)i$
	$\{u,u,p,u\}$	$1.344-0.056i$	
	$\{u,u,u,u\}$	$1.867-0.107i$	$1.860(40)-0.115(25)i$

{πN, πΔ, ηN, σN}

As interactions are turned off the movement of poles is unpredictable



Owa et al., 2503.09940



Suzuki et al., PRL 111 (2010) 042302

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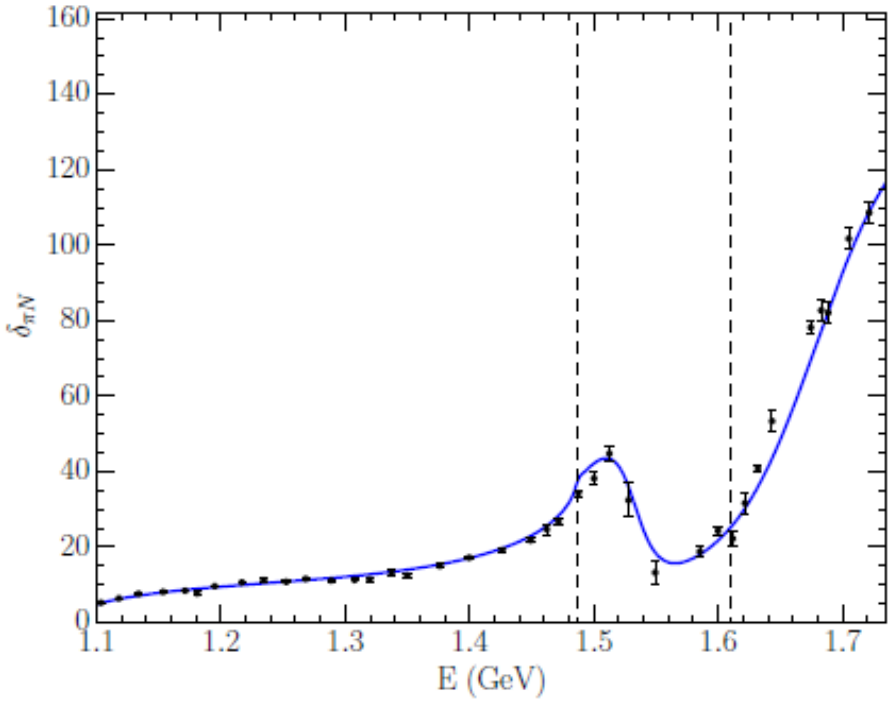
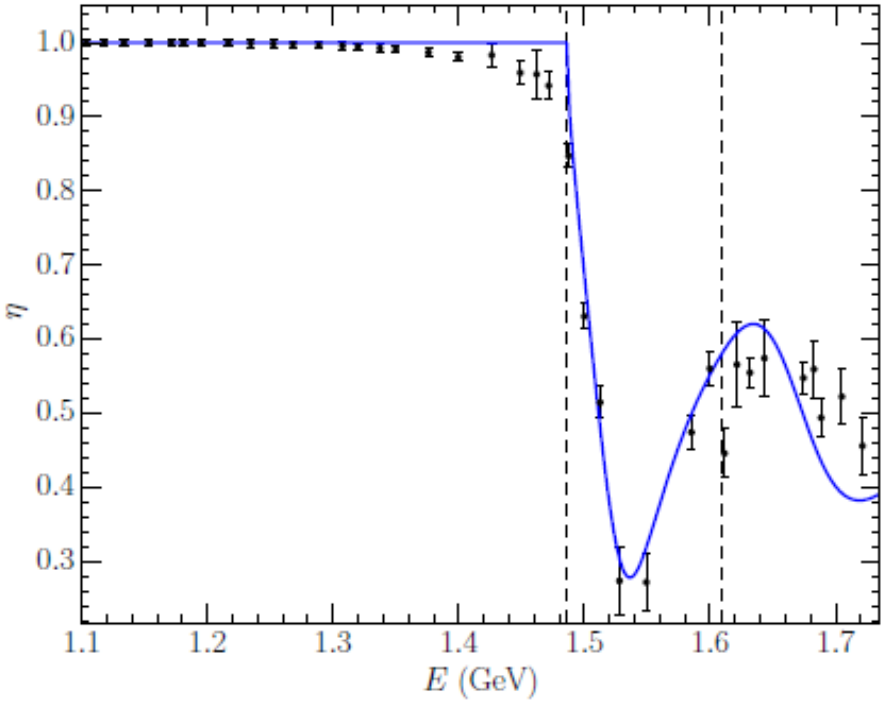
Are all states like this?

NO: The N(1535) *is* a 3-quark state

First construct a Hamiltonian to accurately describe experimental data

Pole at 1531 ± 29 - i 88 ± 2 MeV



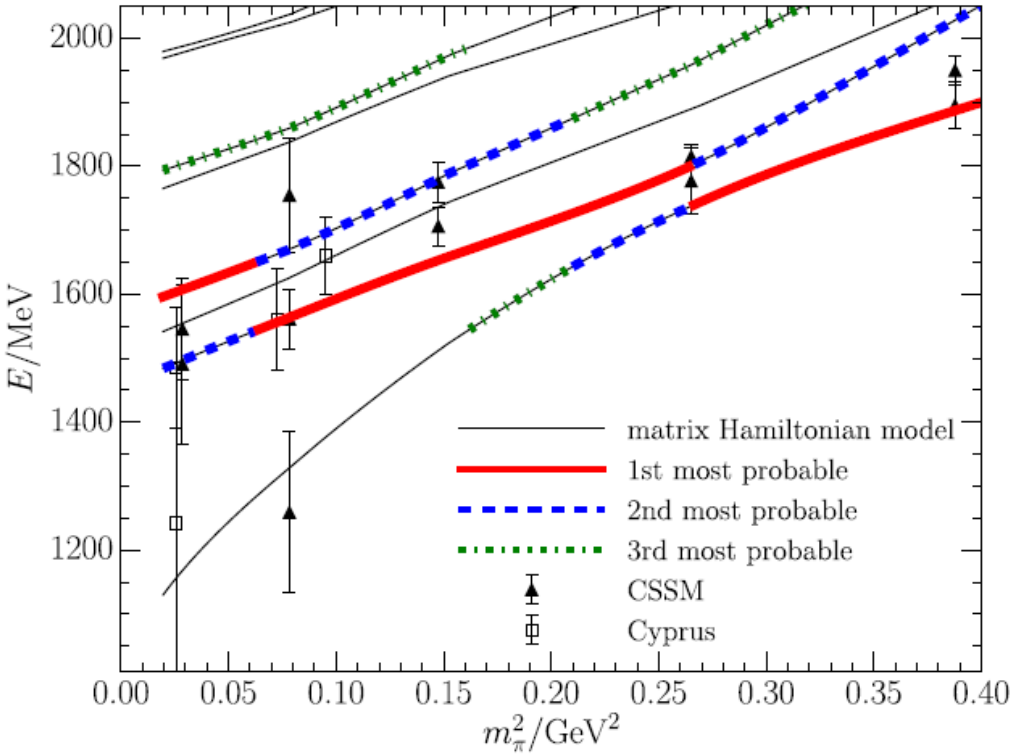


- WI08 single-energy data from SAID.
- Vertical lines indicate the opening of the ηN and $K\Lambda$ thresholds.

Liu *et al.*, PRL 116 (2016) 082004

Analysis of eigenstates

Lattice 3-quark states (L=2.9 fm) match bare-state dominated eigenstates in HEFT calculation



Same conclusion regarding Roper and N(1535) in Tan et al., 2503.04586

Liu *et al.*, PRL 116 (2016) 082004

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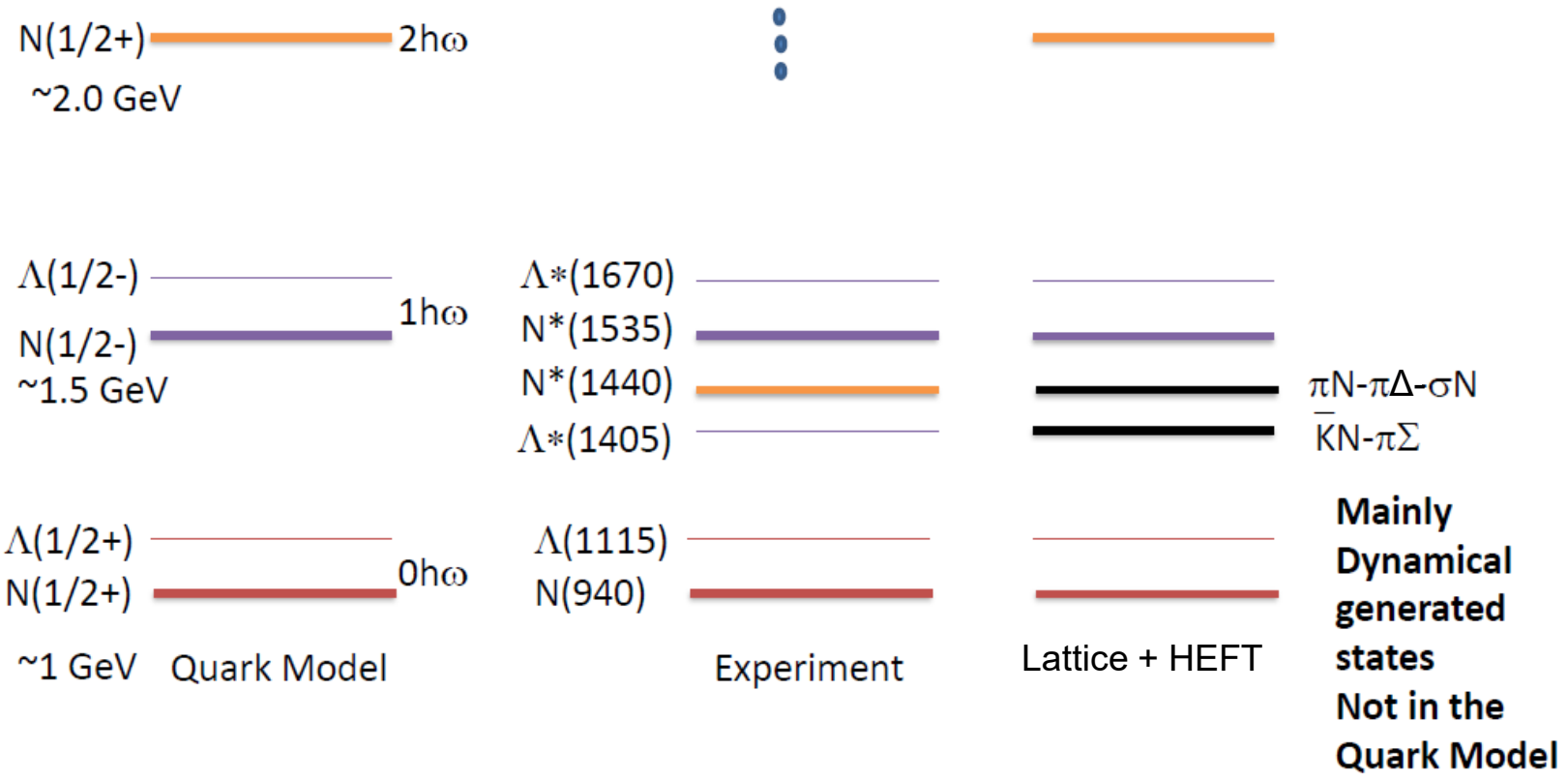
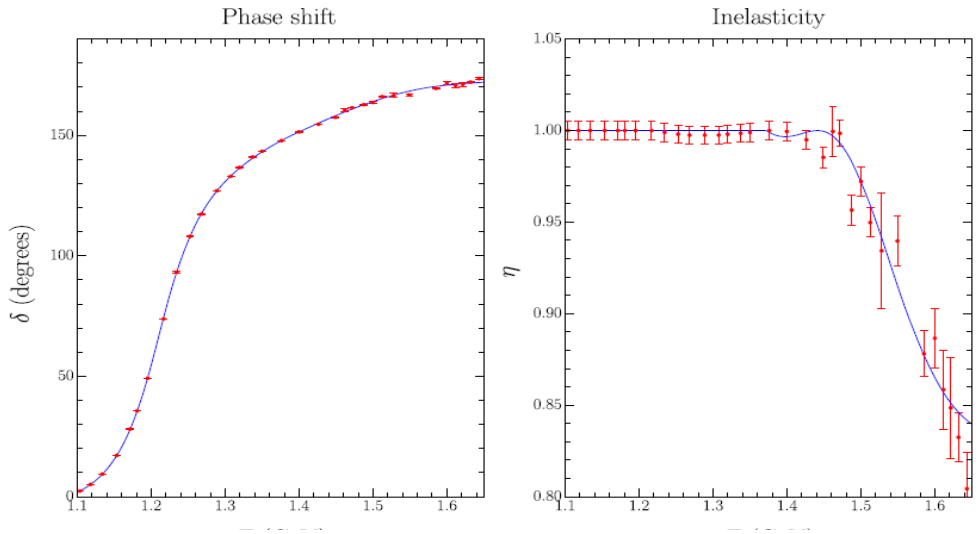
Once the nature of key states becomes clear

the quark model makes sense

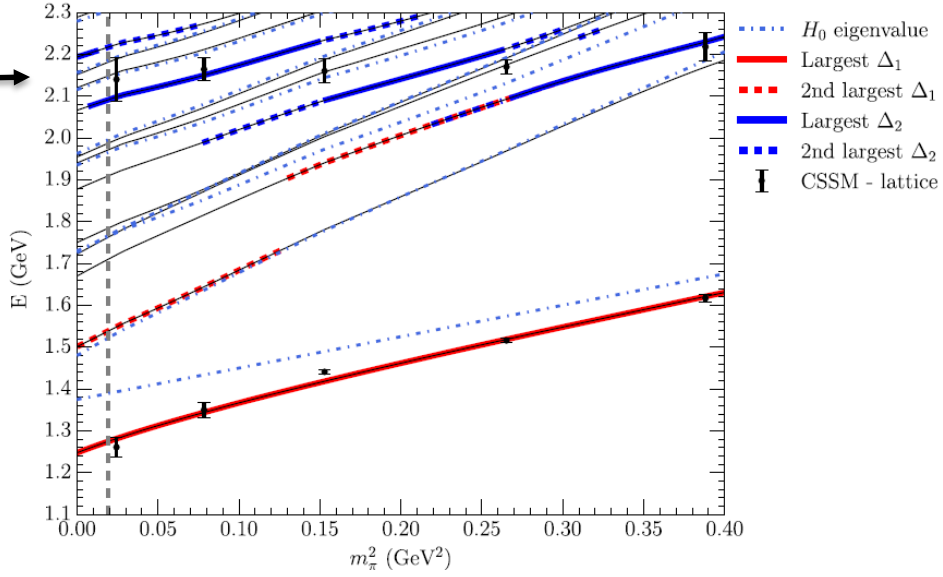


Wu, Leinweber et al., Physical Review D97, 094509 (2018)

Like the Roper the Δ(1600) is also dynamically generated



State dominated by 3-quark (2s) excitation near 2.1 GeV



Summary

- New techniques (HEFT/coupled channels) applied to lattice QCD provide hitherto unimagined insights into hadron structure

- Neither the Λ(1405) nor the Roper are predominantly three-quark states

- The quark model has new life with ordering of major shells as expected

- These insights may well resolve "missing state" problem - as 3-quark states occur at higher energy

Acknowledgements: Derek Leinweber, Zhan-Wei Liu, Jon Hall, Curtis Abell, Jiajun Wu, Waseem Kamleh, Liam Hockley, Shiryo Owa

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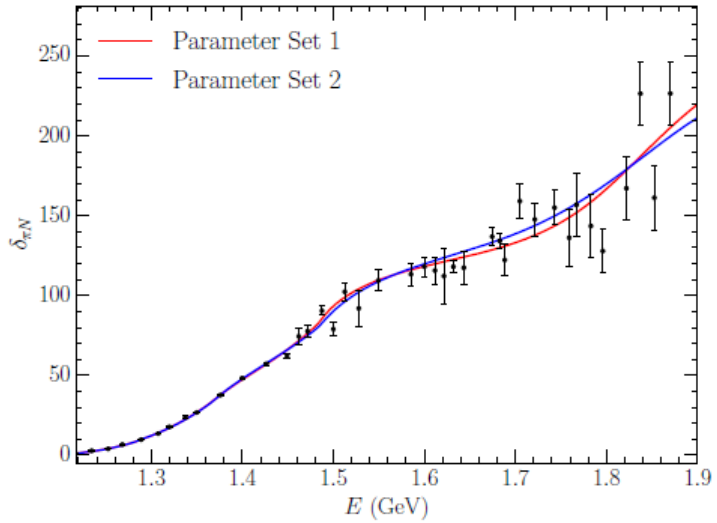
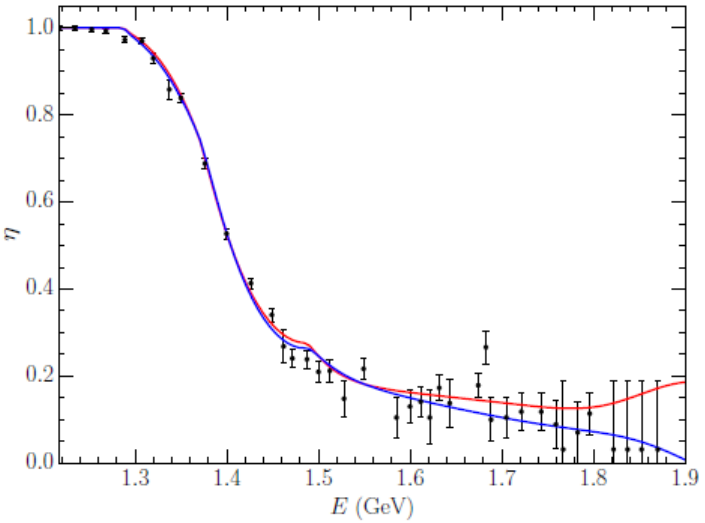
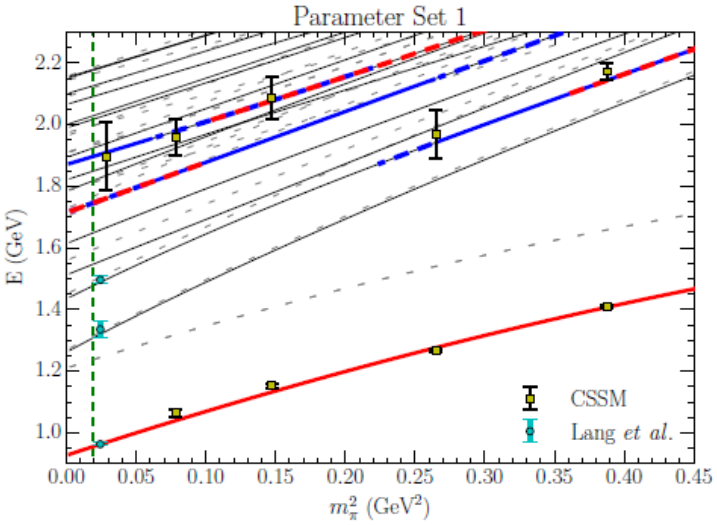
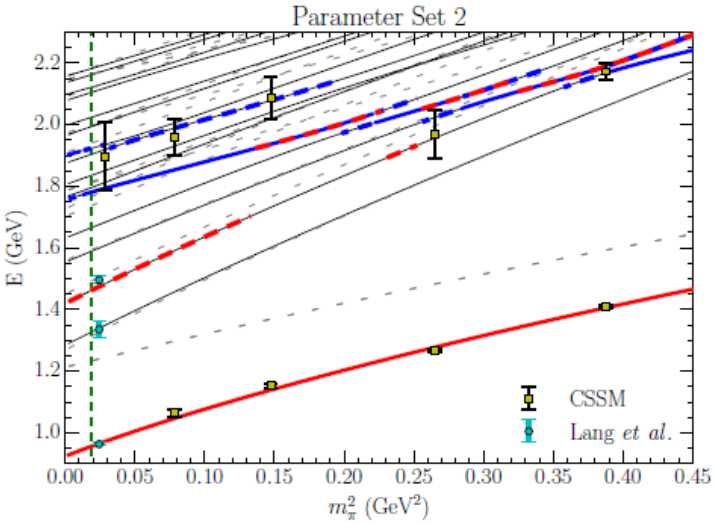
More Detail on Roper









Owa et al., 2503.09945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