

Nucleon axial charge and form factor from lattice QCD

Sungwoo Park

Lawrence Livermore National Laboratory, CA, USA

In collaboration with Rajan Gupta (LANL) and Junsik Yoo (LANL)
INPC 2025, May 30th

LLNL-PRES-2006451

This work was performed under the auspices of the U.S. Department of Energy by Lawrence Livermore National Laboratory under Contract DE-AC52-07NA27344.

• Quasi-elastic scattering (QE)

 $v(\bar{v})$

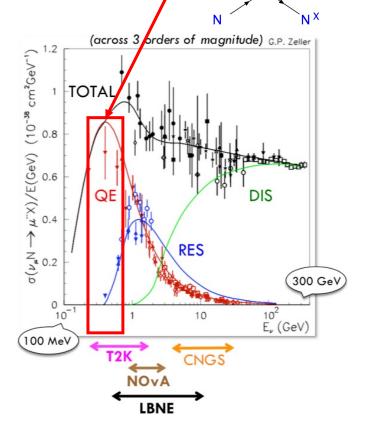
⁻(|+)

Neutrino Oscillation Experiment

- Upcoming flagship neutrino oscillation experiment such as DUNE (US) and HyperK (Japan),
 - Quasi-elastic (QE) neutrino-nucleon scattering is the dominant interaction process

$$\frac{d\sigma}{dQ^2} \propto f(Q^2, [VFF], [AFF], \cdots)$$

- Weak interaction (V-A): Low statistics
 - Vector form factor (VFF)
 - → high-statistics electron scattering experiments
 - Axial vector form factor (AFF)
 - → Lattice QCD calculation is simple at QE processes and can help constraining experimental cross-section
 - \rightarrow Must provide a complete parametrization function (e.g. z-expansion) of $G_A(Q^2)$ including a covariance matrix of parameters.



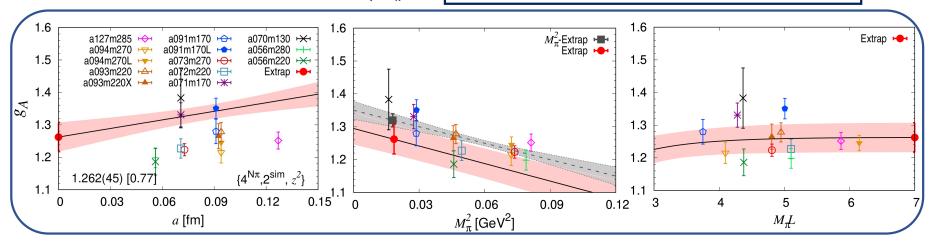
Major Uncertainties in Lattice QCD calculations

- Finite lattice spacing a (UV cut-off effect)
- Chiral fit to get value at physical pion mass
- Finite Volume

$$g(a, M_{\pi}, M_{\pi}L) = c_1 + c_2 a + c_3 M_{\pi}^2 + c_4 \frac{M_{\pi}^2 e^{-M_{\pi}L}}{\sqrt{M_{\pi}L}}$$

- Statistical errors
- Excited state contaminations
- Renormalization

Chiral-Continuum Finite volume extrapolation of nucleon axial charge g_A [NME preliminary]



This work: two new MILC $M_\pi^{ m Phys}$ ensembles

- $N_f = 2 + 1 + 1$ dynamical fermion flavors (isospin symmetric u,d quark masses)
- Gauge ensemble generated by MILC collaboration using HISQ (Highly Improved Staggered Quark) seq quark action

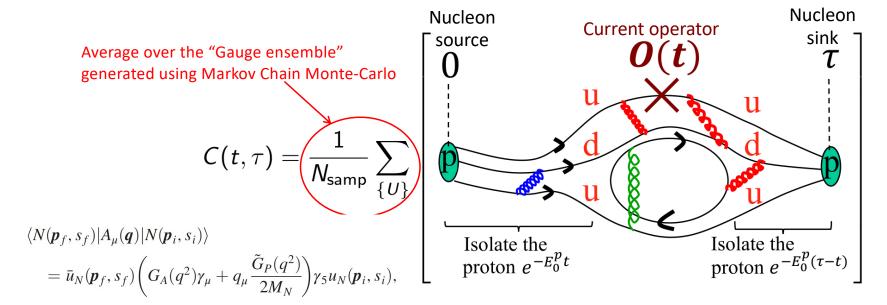
ID	β	a	M_{π}	M_N	L	$T \mid$	$M_{\pi}L $	Lattices	$N_{ m HP}$	$N_{ m LP}$	au
		(fm)	(MeV)	(MeV)							
a09m135	6.3	0.087(1)	134(1)	947.1(7.6)	64	96	3.80	5,497	16,491	527,712	$\{10,12,14,16,18\}$
a06m135	6.7	0.057(0)	136(0)	932.2(8.1)	96	192	3.78	3,990	15,960	430,920	{16,18,20,22,24}

~4x improved statistics

- Improved gauge link smearing, mass parameter tuning
- Larger source-sink time separation
- The largest momentum transfer $\vec{q}=2\pi\vec{n}/L$ is doubled
 - spacelike 4-momentum transfer $Q^2 = -q^2$ max has increased from 0.45 GeV² to 0.82 GeV²

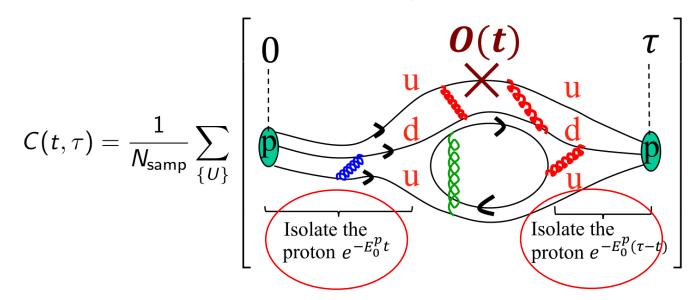
Calculation of Nucleon Matrix Element

• Nucleon matrix element $(\langle p'|A|p\rangle$, to be decomposed to Form Factors) are extracted from the <u>3-point correlation function</u> $C(t,\tau) \equiv \langle N^p(\tau)O(t)\overline{N}^p(0)\rangle$:

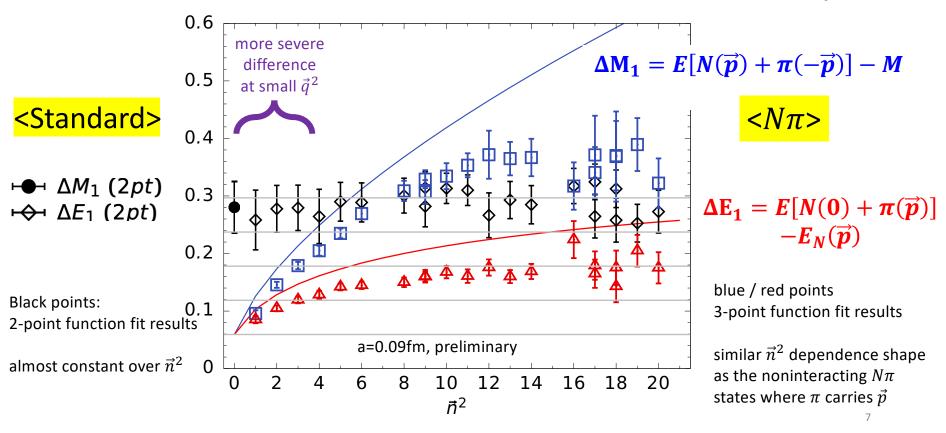


Excited State (ES) effect is more severe at M_{π}^{Phys}

- Nucleon signal/noise decays $\propto e^{-(E-1.5M_{\pi})\tau}$ with Euclidean time τ .
- Nucleon operator creates ground state nucleons (N) plus all excited states with the same quantum number, including $N\pi$, $N\pi\pi$, $N\rho$, $N^*(1440)$, $N^*(1710)$,



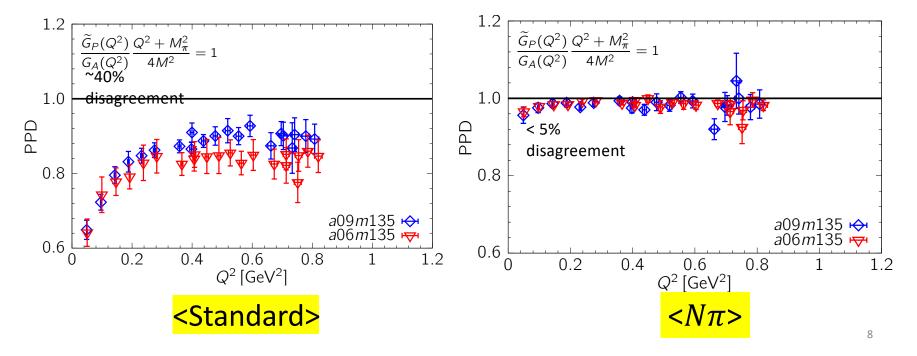
1st excited state masses from two different analysis



Checking Pion Pole Dominance (PPD) hypothesis

• Relates Induced pseudoscalar and the axial form factors

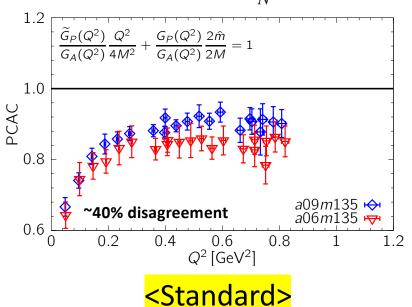
$$(\widetilde{G_P}(Q^2) = \frac{4M^2}{Q^2 + M_\pi^2} G_A(Q^2))$$

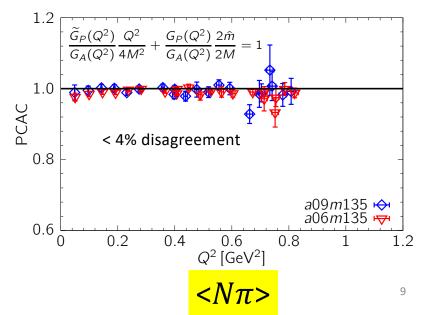


PCAC (Partially Conserved Axial Current) relation

- $\partial_{\mu}A_{\mu}(x) = 2\widehat{m}P(x)$ where $\widehat{m} = Z_{m}m_{ud}Z_{P}Z_{A}^{-1}$
- Applied to nucleon ground state, it relates the 3 nucleon form factors

$$2M_NG_A(Q^2)-rac{Q^2}{2M_N} ilde{G}_P(Q^2)=2\hat{m}G_P(Q^2)$$
 Generalized Goldberger-Treiman Relation





2024 Comparison of lattice results of G_A

-0.45170

-0.24394

-0.45170

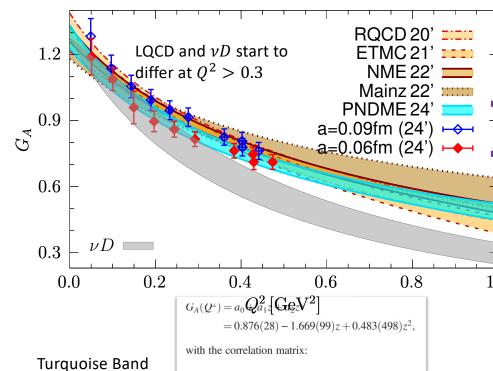
-0.02966

-0.02966

-0.24394

1.0

RQCD 20' JHEP 05 126 ETMC 21' PRD 103, 034509 NME 22' PRD 105, 054505 Mainz 22' PRD 106, 074503 PNDME 24' PRD 109, 014503



PNDME 24'

Two physical pion mass data in PNDME 24'

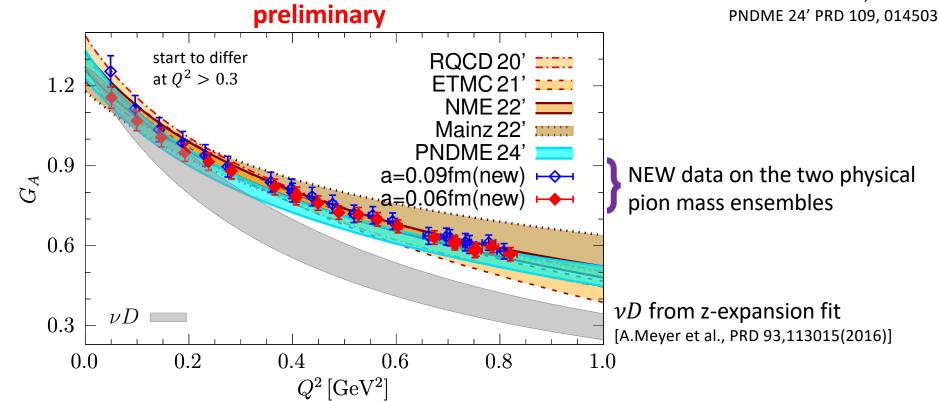
νD from z-expansion fit

[A.Meyer et al., PRD 93,113015(2016)]

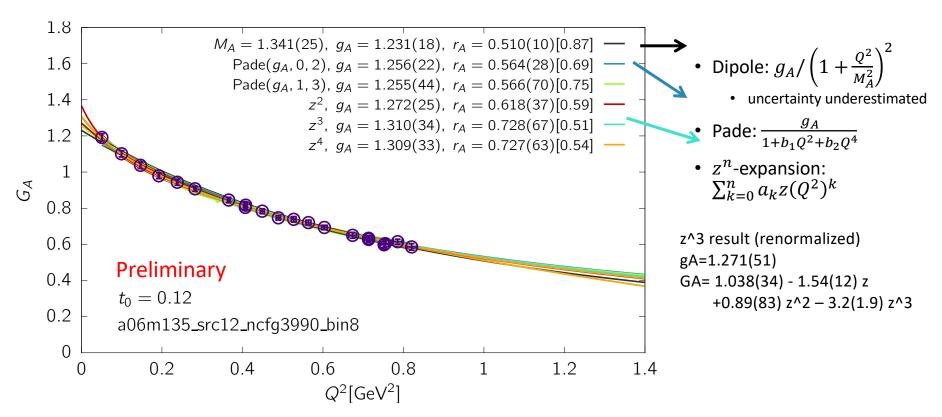
- The deuterium bubble chamber experiment data (70-80s) have smal statistics (10^3),
 - Unknown correction added. Cannot access original data.
 - Unresolved nuclear (deuterium) correction
 - Dipole fit underestimated G_A uncertainty by x10

Our new data for G_A

RQCD 20' JHEP 05 126 ETMC 21' PRD 103, 034509 NME 22' PRD 105, 054505 Mainz 22' PRD 106, 074503 PNDME 24' PRD 109, 014503



G_A^{u-d} : Examined Dipole, Pade and z-expansion fits



Summary

- Calculating the isovector axial form factor $G_A^{u-d}(Q^2)$ as part of a comprehensive analysis of nucleon structure.
- New data for G_A^{u-d} on 2 physical pion mass ensembles over $0.04 < Q^2 \lesssim 1 \text{ GeV}^2$ with significantly improved statistics and systematics control. This enables a more reliable chiral-continuum extrapolation.
- Result: a z-expansion parametrization with the correlation matrix of the fit parameters.
- $G_A^{u-d}(Q^2)$ is important input in the calculation of the quasi-elastic neutrino-nucleon scattering cross-section. (Needed for DUNE to meet precision goal)

PNDME and NME members

- Tanmoy Bhattacharya (LANL)
- Vincenzo Cirigliano (INT)
- Rajan Gupta (LANL)
- Emanuele Mereghetti (LANL)
- Boram Yoon (NVIDIA)
- Junsik Yoo (LANL)
- Yong-Chull Jang (BNL)
- Sungwoo Park (LLNL)
- Santanu Mondal (MSU)
- Huey-Wen Lin (MSU)
- Balint Joo (ORNL)
- Frank Winter (Jlab)

References

PNDME (clover-on-HISQ formulation)

• Charges: Gupta et al, PRD.98 (2018) 034503

• AFF: Gupta et al, PRD 96 (2017) 114503

Jang et al, PRL 124 (2020) 072002 Jang et al, PRD 109 (2024) 014503

• VFF: Jang et al, PRD 100 (2020) 014507

• $\sigma_{\pi N}$ Gupta et al, PRL 127 (2021) 242002

• d_n from Θ -term Bhattacharya et al, PRD 103 (2021) 114507

d_n from qEDM Gupta et al, PRD 98 (2018) 091501
 Moments of PDFs Mondal et al, PRD 102 (2020) 054512

Proton spin: Lin et al, PRD 98 (2018) 094512

• Flavor diag. charges: Park et al, arXiv:2503.07100

NME (clover-on-clover formulation)

• Charges, VFF, AFF: Park et al, PRD 105 (2022) 054505

• Moments of PDFs Mondal et al, JHEP 04 (2021) 044

We thank MILC Collaboration for providing the 2+1+1 flavor HISQ lattices.

The calculations used the CHROMA software suite.

We thank DOE for computer time allocations at NERSC (ERDCAP) and OLCF (INCITE HEP133, ALCC HEP145).

We thank the USQCD collaboration and institutional Computing at LANL for computer time.

BACKUP

Lepton-nucleon scattering

$$\frac{d\sigma}{dQ^{2}} \begin{pmatrix} \nu_{l} + n \to l^{-} + p \\ \bar{\nu}_{l} + p \to l^{+} + n \end{pmatrix}$$

$$= \frac{M^{2}G_{F}^{2}\cos^{2}\theta_{c}}{8\pi E_{\nu}^{2}} \left\{ A(Q^{2}) \pm B(Q^{2}) \frac{(s-u)}{M^{2}} + C(Q^{2}) \frac{(s-u)^{2}}{M^{4}} \right\},$$

$$A(Q^{2}) = \frac{(m^{2} + Q^{2})}{M^{2}} \left[(1+\tau)F_{A}^{2} - (1-\tau)F_{1}^{2} + \tau(1-\tau)F_{2}^{2} + 4\tau F_{1}F_{2} - \frac{m^{2}}{4M^{2}} \left((F_{1} + F_{2})^{2} + (F_{A} + 2F_{P})^{2} - 4\left(1 + \frac{Q^{2}}{4M^{2}}\right)F_{P}^{2} \right) \right],$$

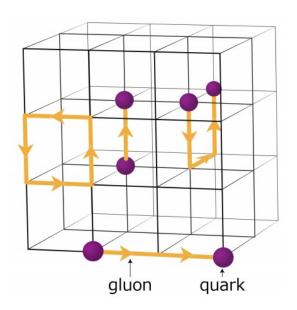
$$B(Q^{2}) = \frac{Q^{2}}{M^{2}}F_{A}(F_{1} + F_{2}),$$

$$C(Q^{2}) = \frac{1}{4}(F_{A}^{2} + F_{1}^{2} + \tau F_{2}^{2}).$$

 F_A = axial form factor \tilde{F}_P = induced pseudoscalar $G_E = F_1 - \tau F_2$ Electric $G_M = F_1 + F_2$ Magnetic $\tau = Q^2/4M^2$ $M=M_n=M_p\approx 939$ MeV $m=M_\pi$

Lattice QCD

[Formulated by K. Wilson (1974). Numerical computation field opened by M. Creutz (1979)]



Lattice QCD is QCD defined on a 4-dimensional Euclidean space-time lattice

- Finite lattice spacing: (a)
- Quark fields (q, \bar{q}) , Gauge fields (gluons): (U_{μ})
- Perturbative & Numerical (nonperturbative) calculations

The simulation allows **ab initio** calculations of nonperturbative QCD interactions of quarks and gluons using the **Feynman path integral** formulation of QFT.

Major systematic errors coming from:

- Finite lattice spacing a (UV cut-off effect)
- Chiral fit to get value at physical pion mass
- Finite Volume

- Statistical errors
- Excited state contaminations
- Renormalization