

Impact of the meson-exchange currents on
the magnetic dipole moments in
odd near doubly magic nuclei
analyzed within nuclear DFT framework

Herlik Wibowo

Collaborators: R. Han, B. C. Backes, G. Danneaux, J. Dobaczewski,
W. Jiang, and M. Kortelainen

The 29th International Nuclear Physics Conference (INPC 2025)
May 26, 2025



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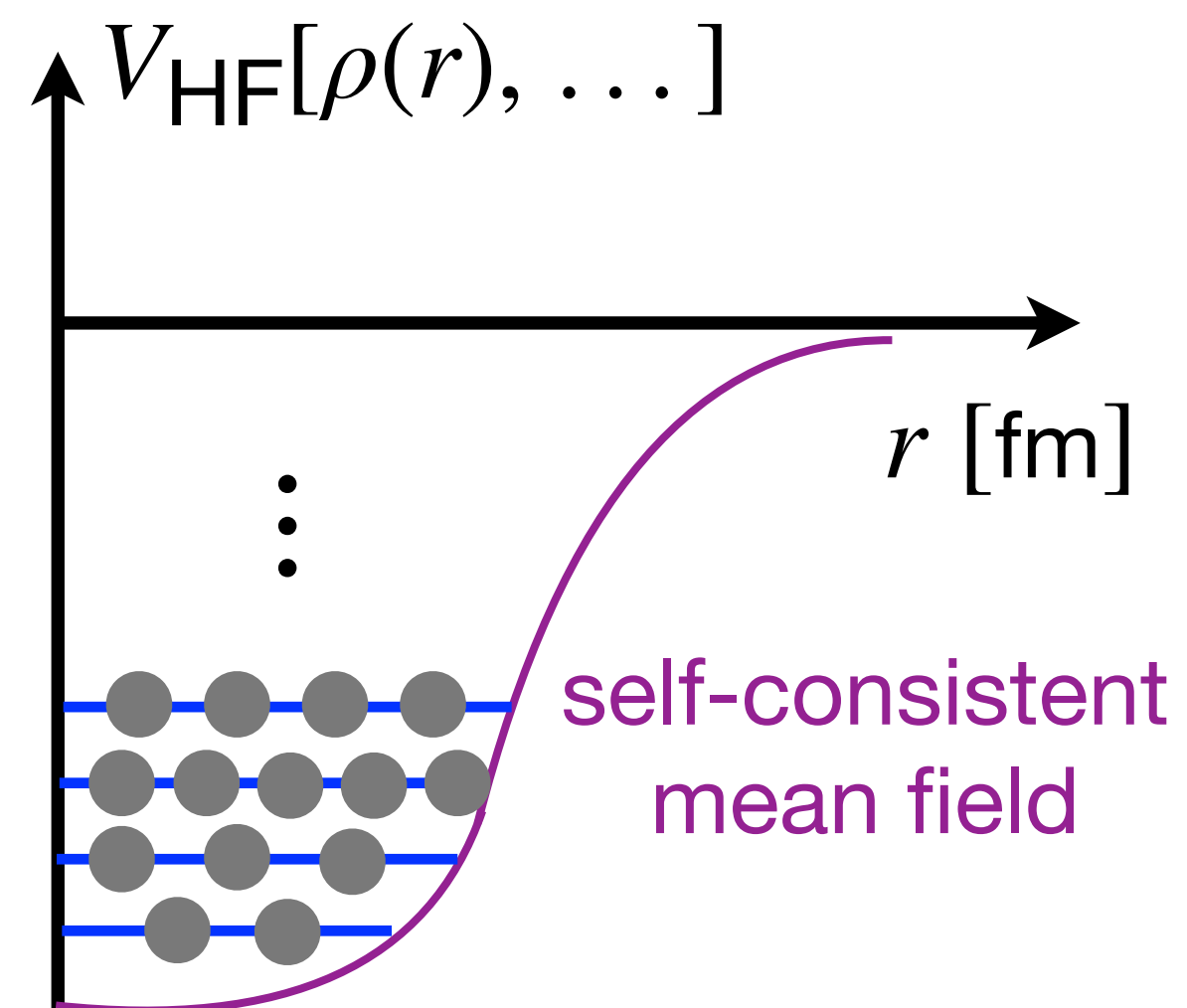
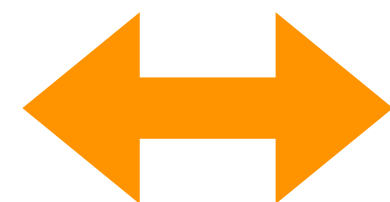
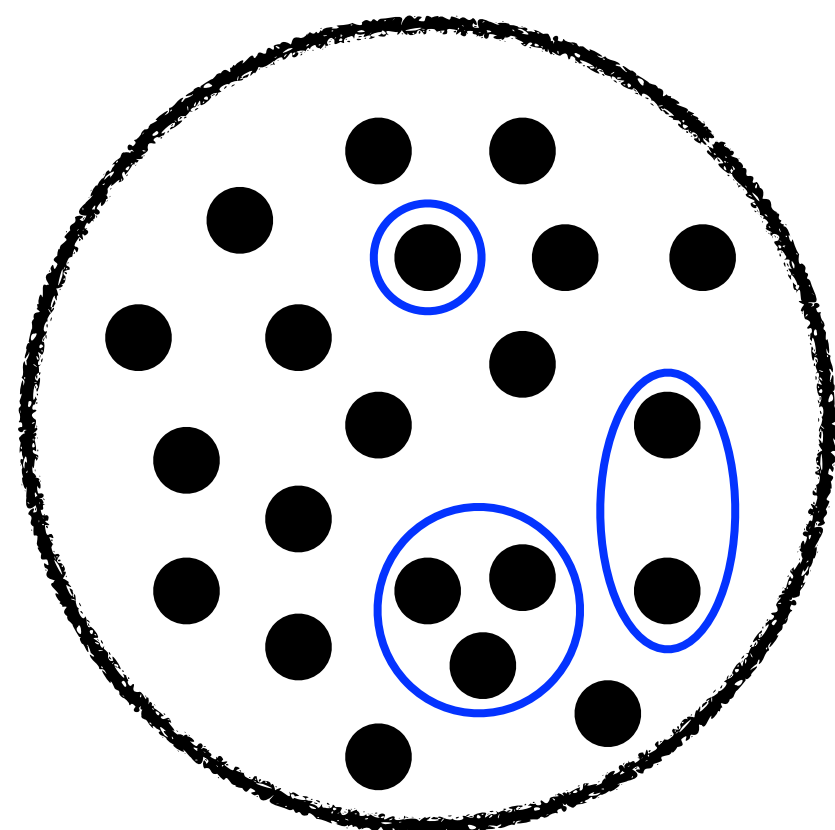


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Nuclear density functional theory

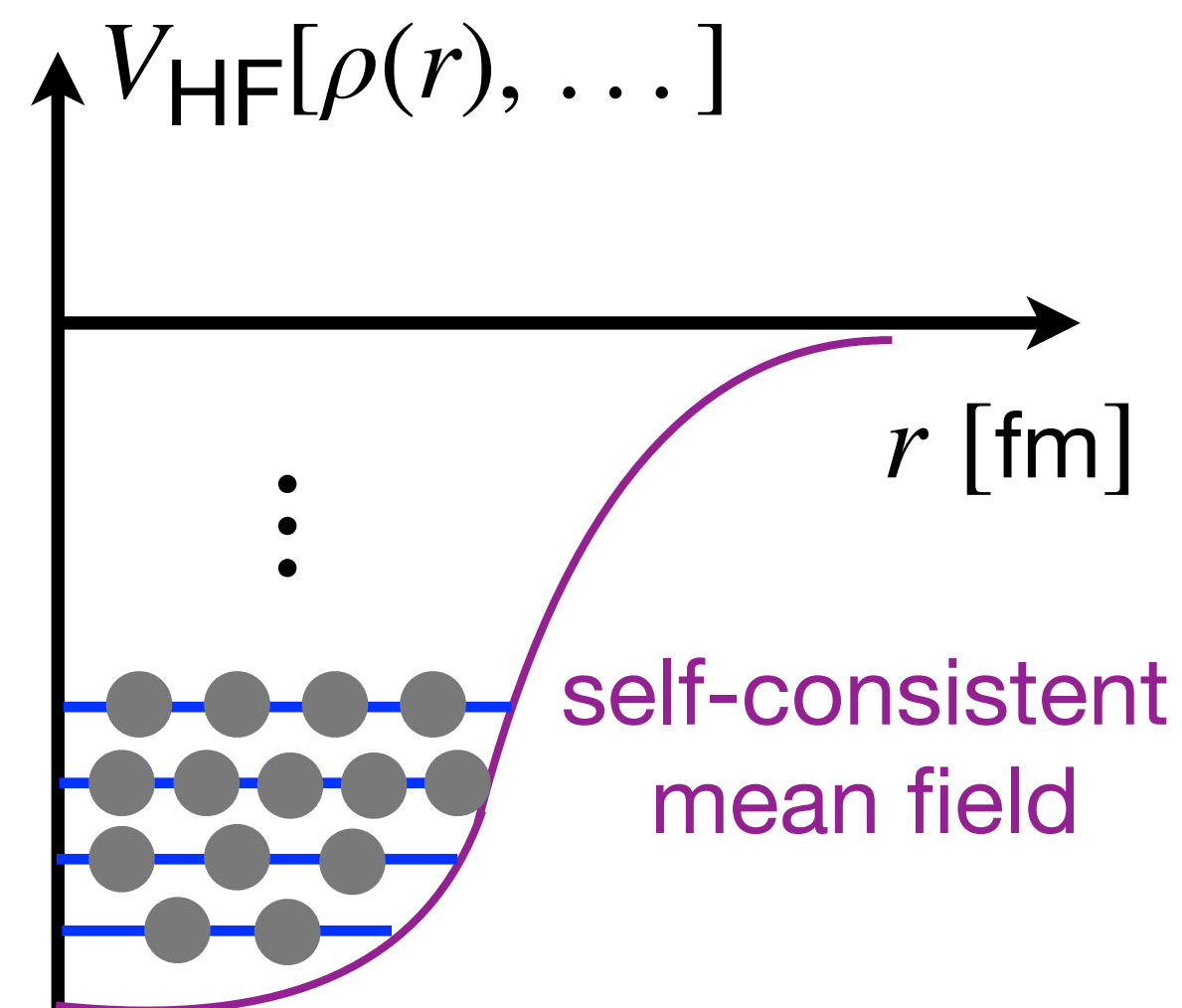
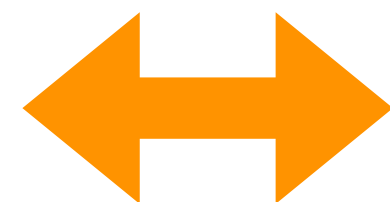
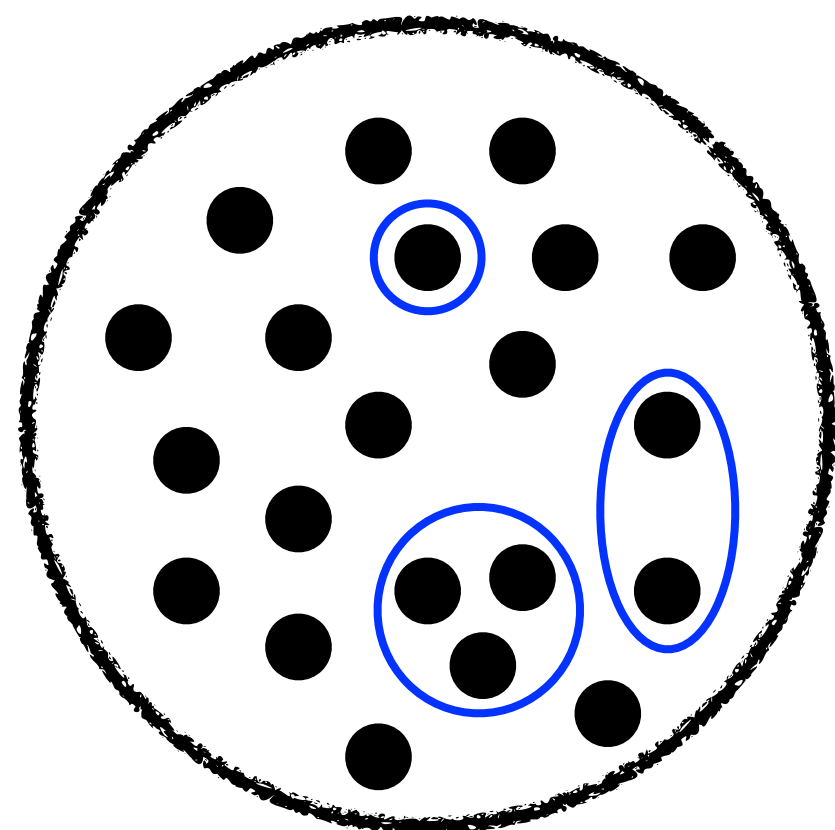
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Nuclear density functional theory



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Nuclear density functional theory



Trial s.p. wave-function $\{ |\varphi_i^{\text{trial}}\rangle \}$

Calculate the densities $\{\rho, \dots\}$

Calculate the mean field $V_{\text{HF}}[\rho, \dots]$

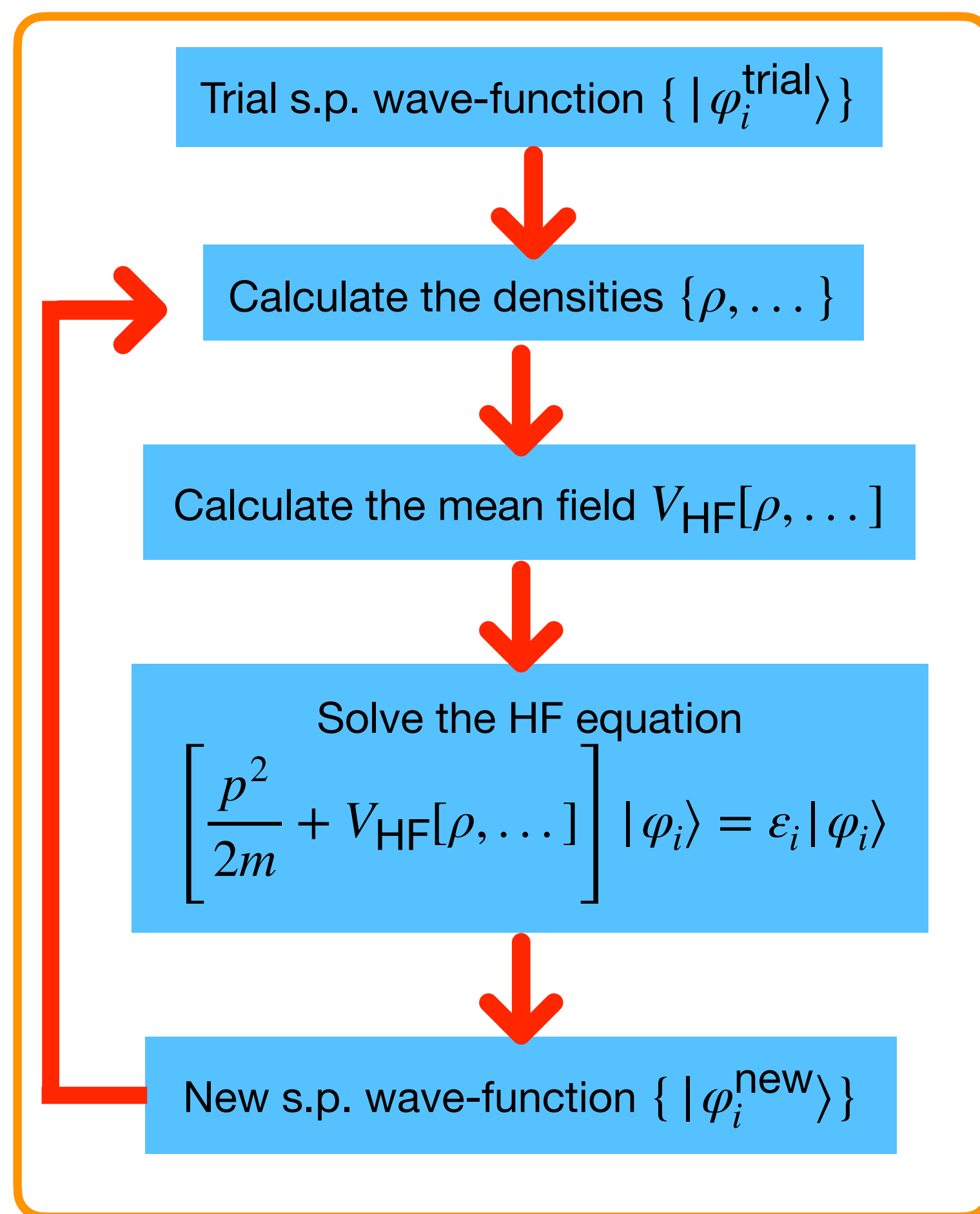
Solve the HF equation

$$\left[\frac{p^2}{2m} + V_{\text{HF}}[\rho, \dots] \right] |\varphi_i\rangle = \varepsilon_i |\varphi_i\rangle$$

New s.p. wave-function $\{ |\varphi_i^{\text{new}}\rangle \}$

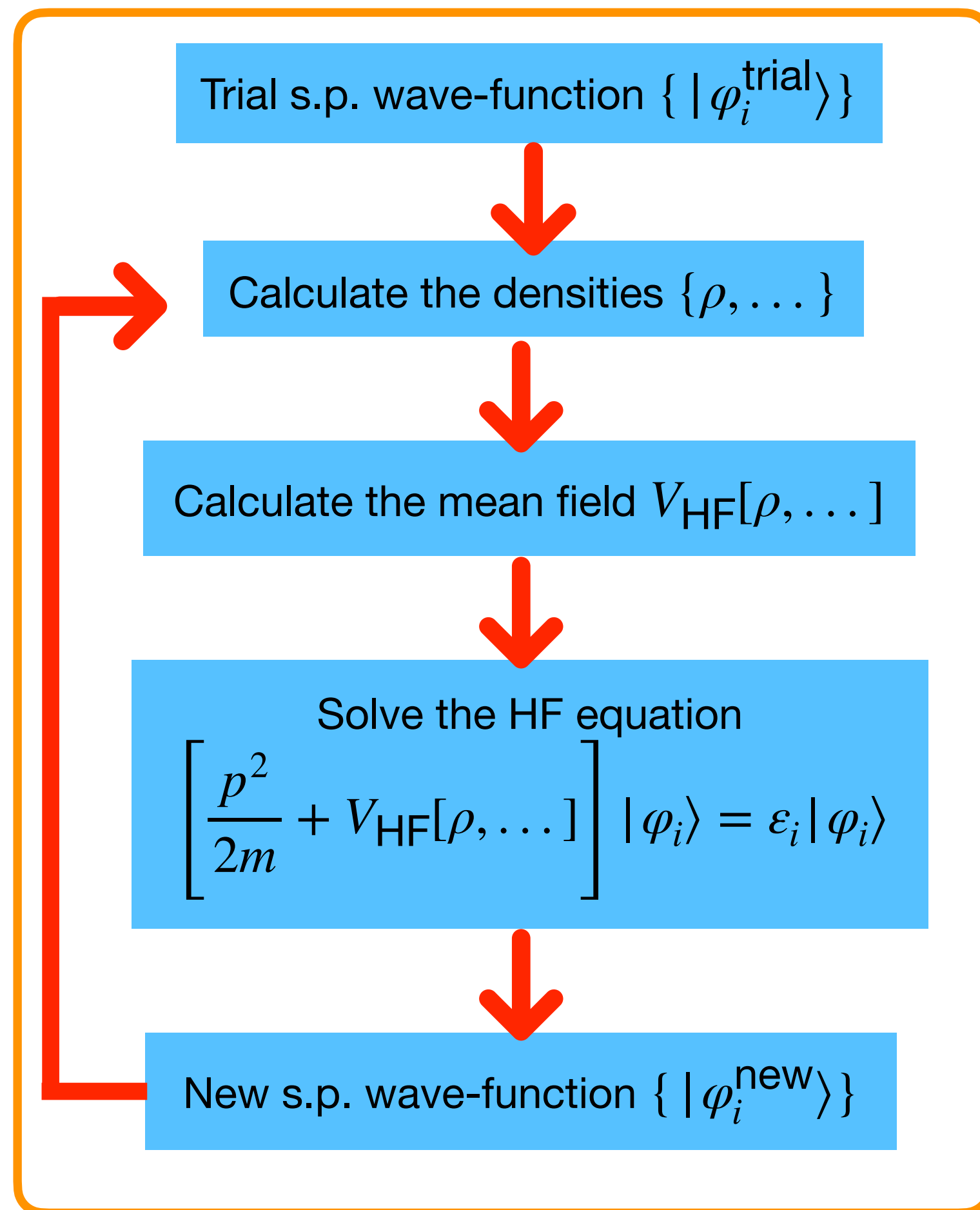
Nuclear DFT methodology

Self-consistency

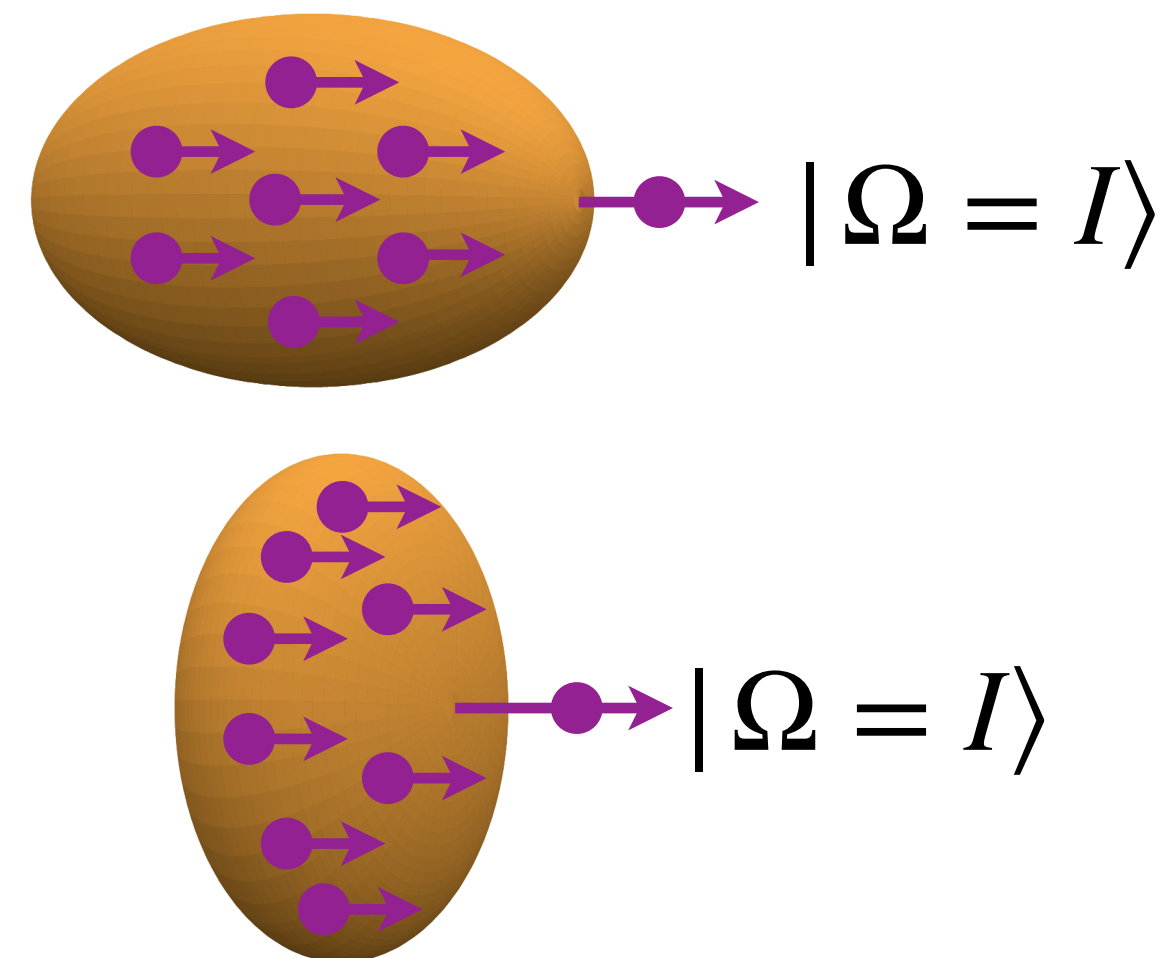


Nuclear DFT methodology

Self-consistency



Spin-core polarization



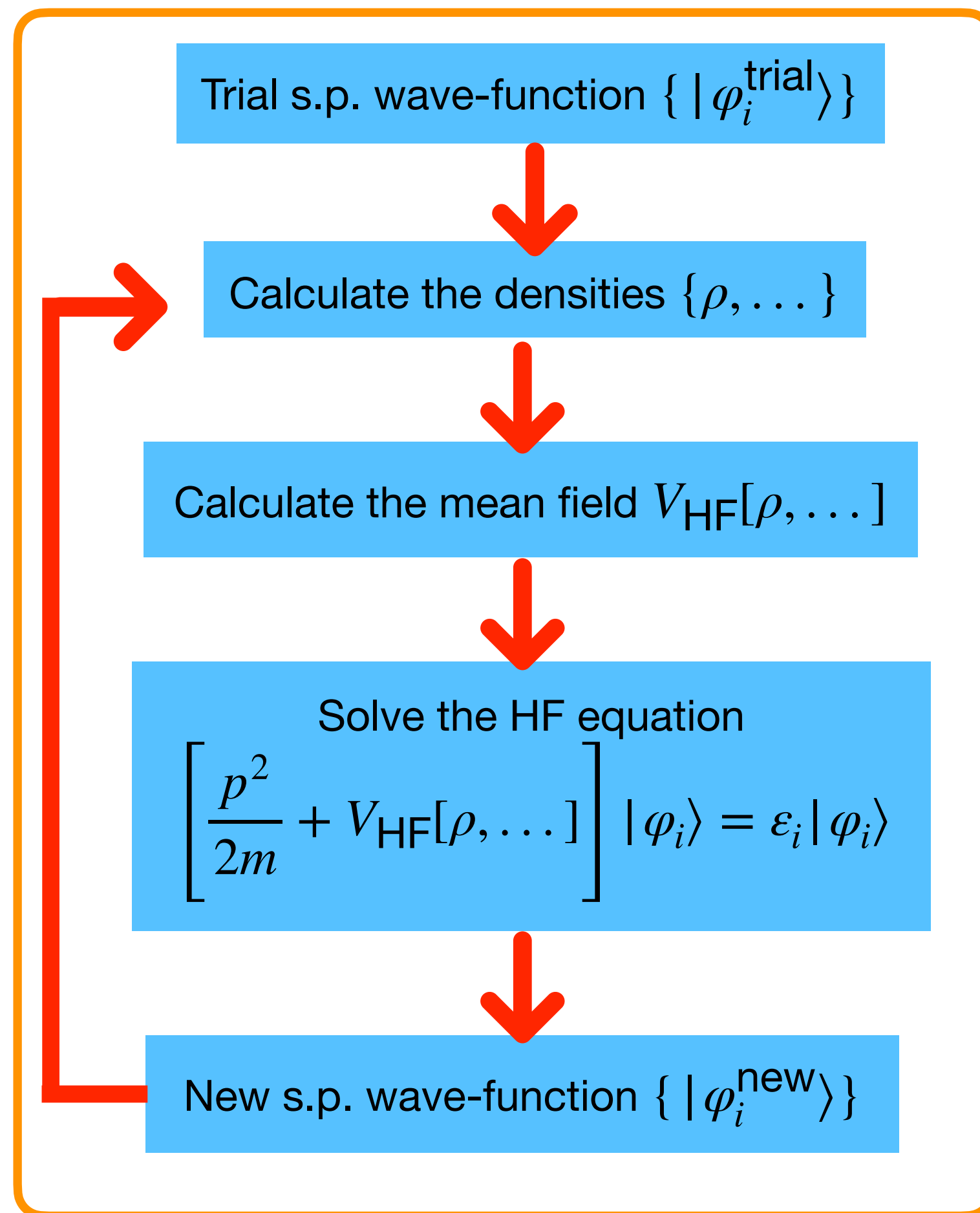
Landau parameter g'_0

$$g'_0 = N_0 \left(2C_1^s + 2C_1^T (3\pi^2 \rho_0 / 2)^{2/3} \right)$$

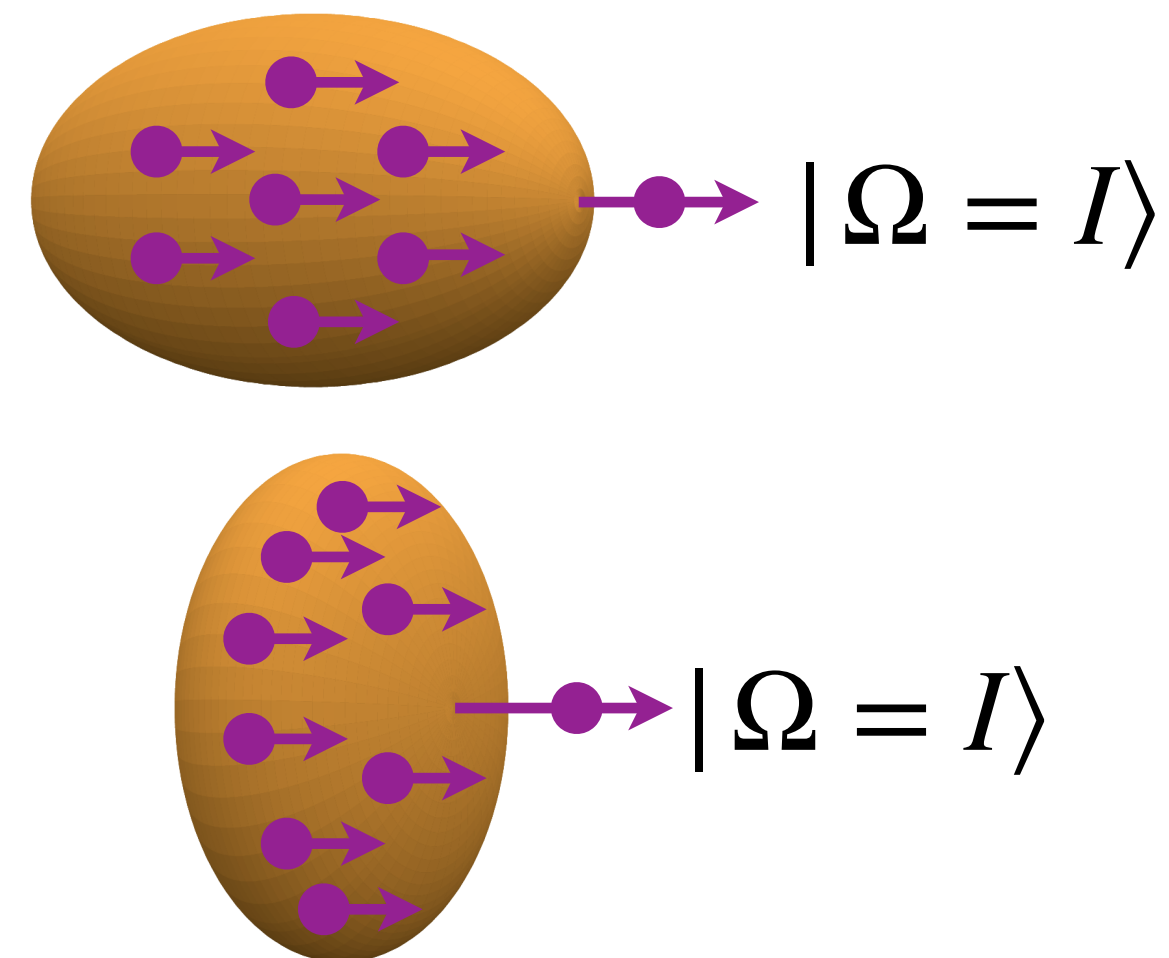
$$\frac{1}{N_0} \approx 150 \frac{m}{m^*} \text{ MeV} \cdot \text{fm}^3$$

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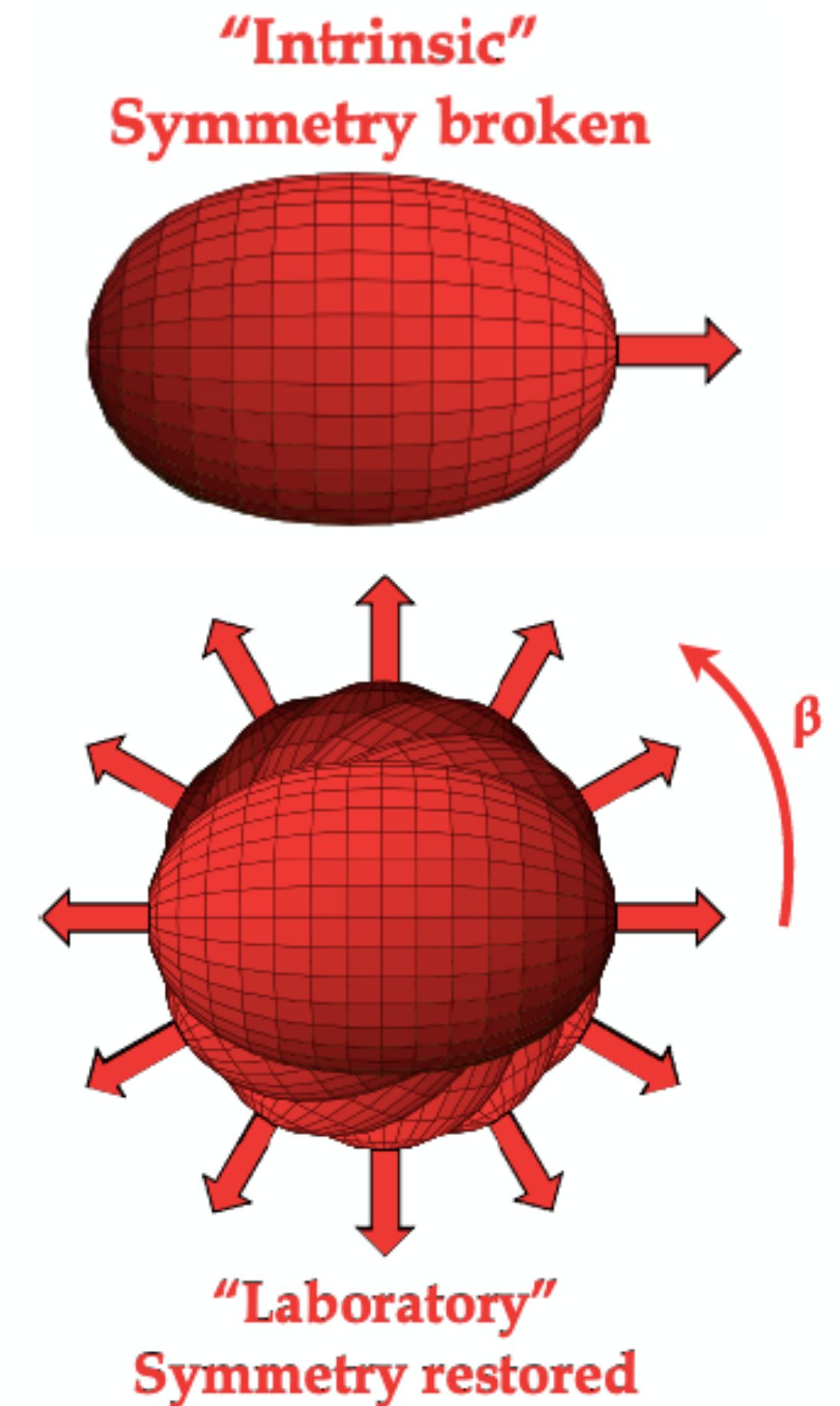


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Symmetry restoration



Nuclear magnetic dipole moments

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Spectroscopic magnetic dipole μ moment is defined as:

$$\mu = \sqrt{\frac{4\pi}{3}} \langle II | \hat{M}_{10} | II \rangle,$$

P. Ring and P. Schuck, *The Nuclear Many-Body Problem*

$|II\rangle$ = angular momentum projected (AMP) states

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$$\hat{M}_{10} = \sqrt{\frac{3}{4\pi}} \mu_N \sum_{i=1}^A \left\{ g_s^{(i)} s_{zi} + g_\ell^{(i)} \ell_{zi} \right\},$$

$$g_s^{(i)} = g_s^{(\pi)}(g_s^{(\nu)}) = 5.59(-3.83), \quad g_\ell^{(i)} = g_\ell^{(\pi)}(g_\ell^{(\nu)}) = 1(0).$$

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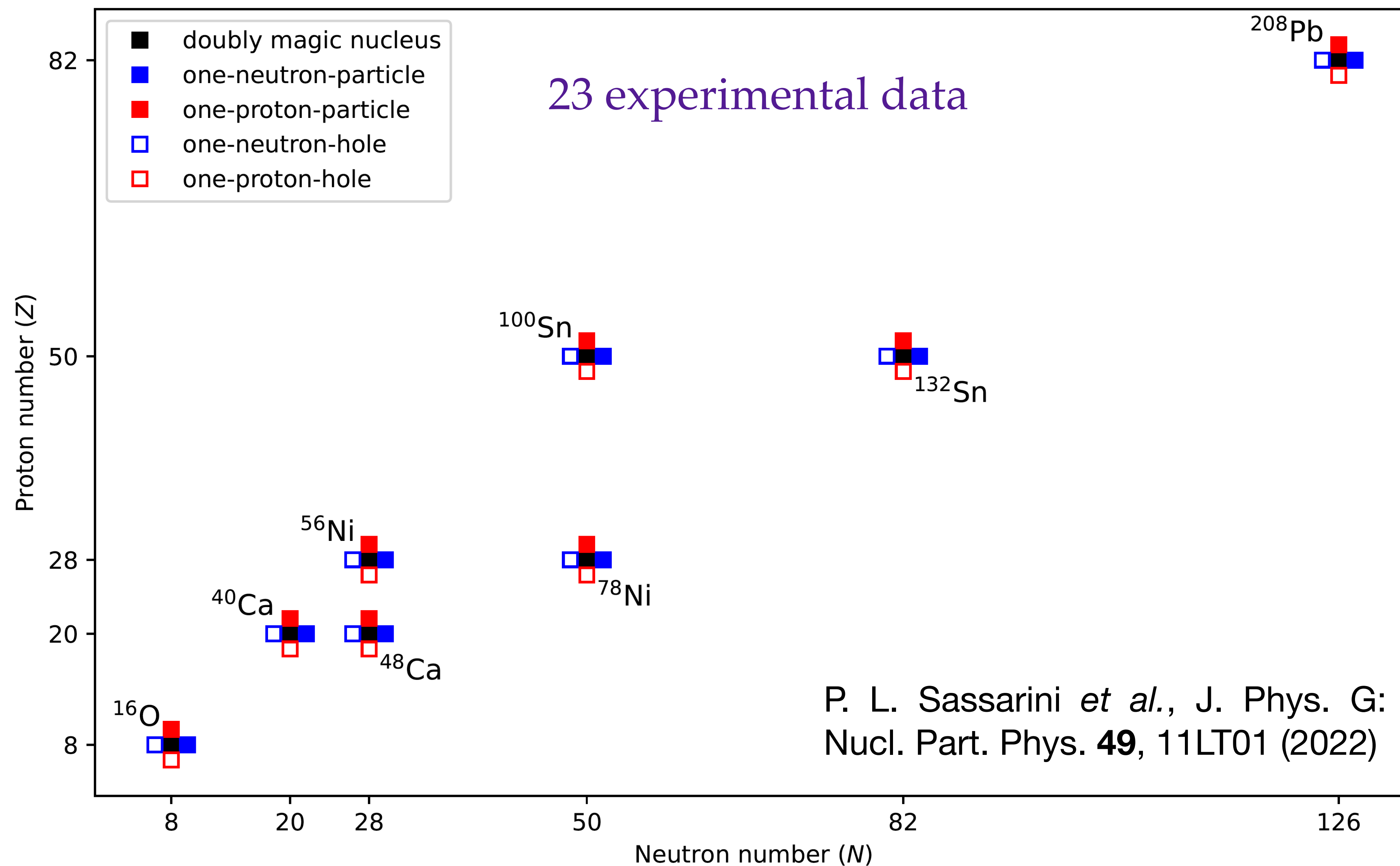
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Since the self-consistent polarizations act in the full single-particle space, no effective g -factors are needed!

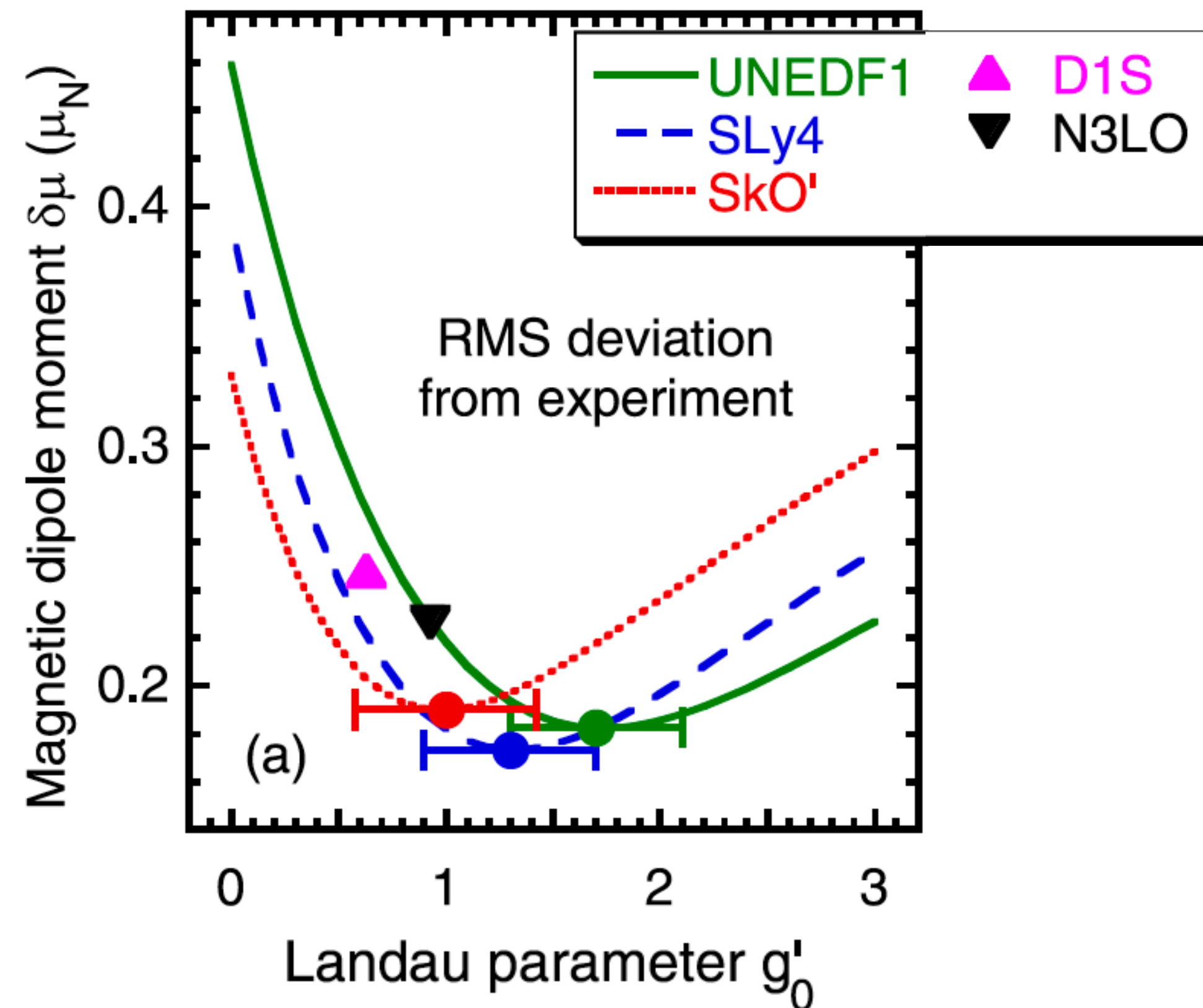
Odd near doubly magic nuclei



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Optimum Landau parameter g'_0

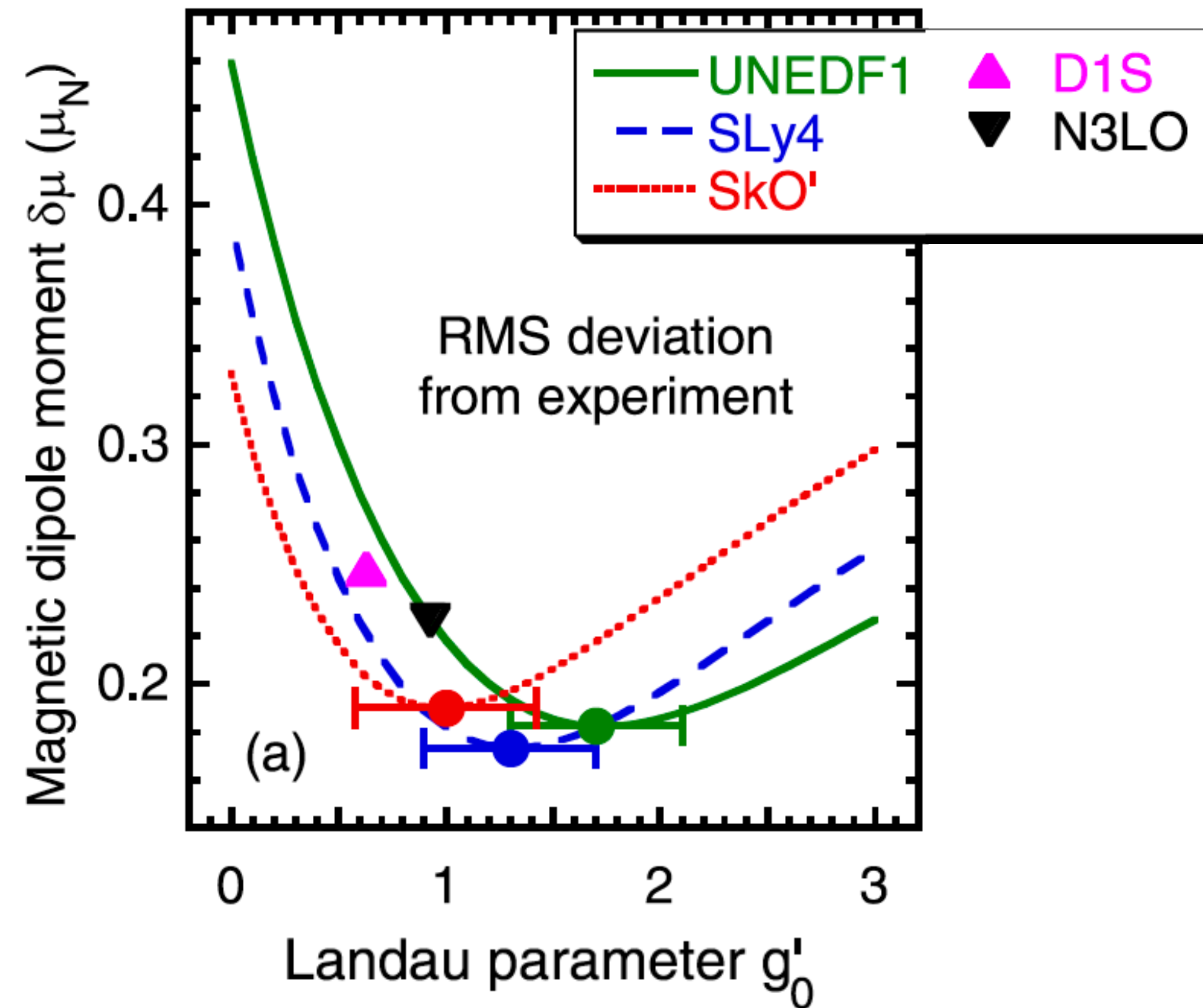
Optimum Landau parameter g'_0



$$g'_0 = 1.7(4) \text{ (UNEDF1)}$$

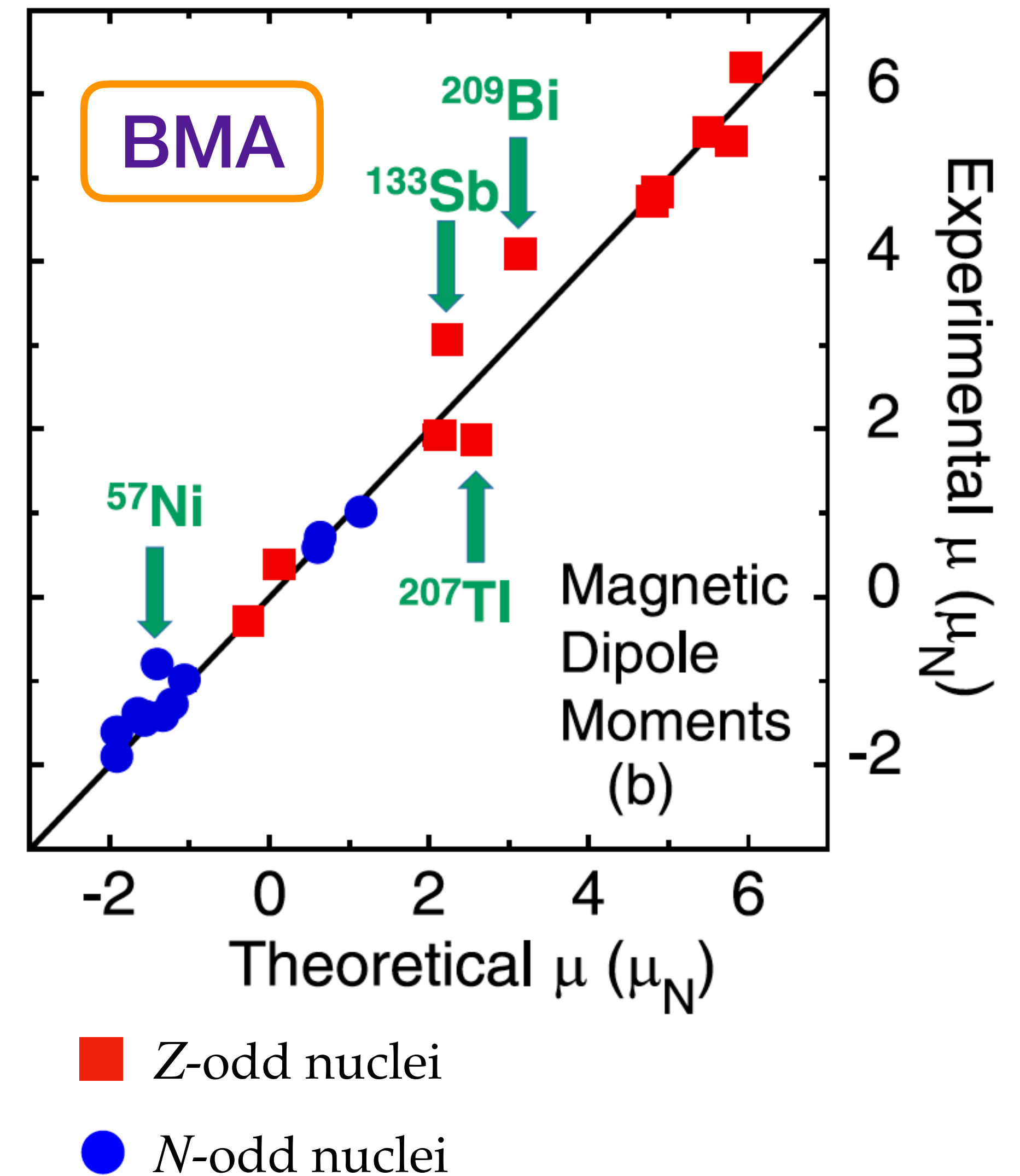
P. L. Sassarini *et al.*, J. Phys. G: Nucl. Part. Phys. **49**, 11LT01 (2022)

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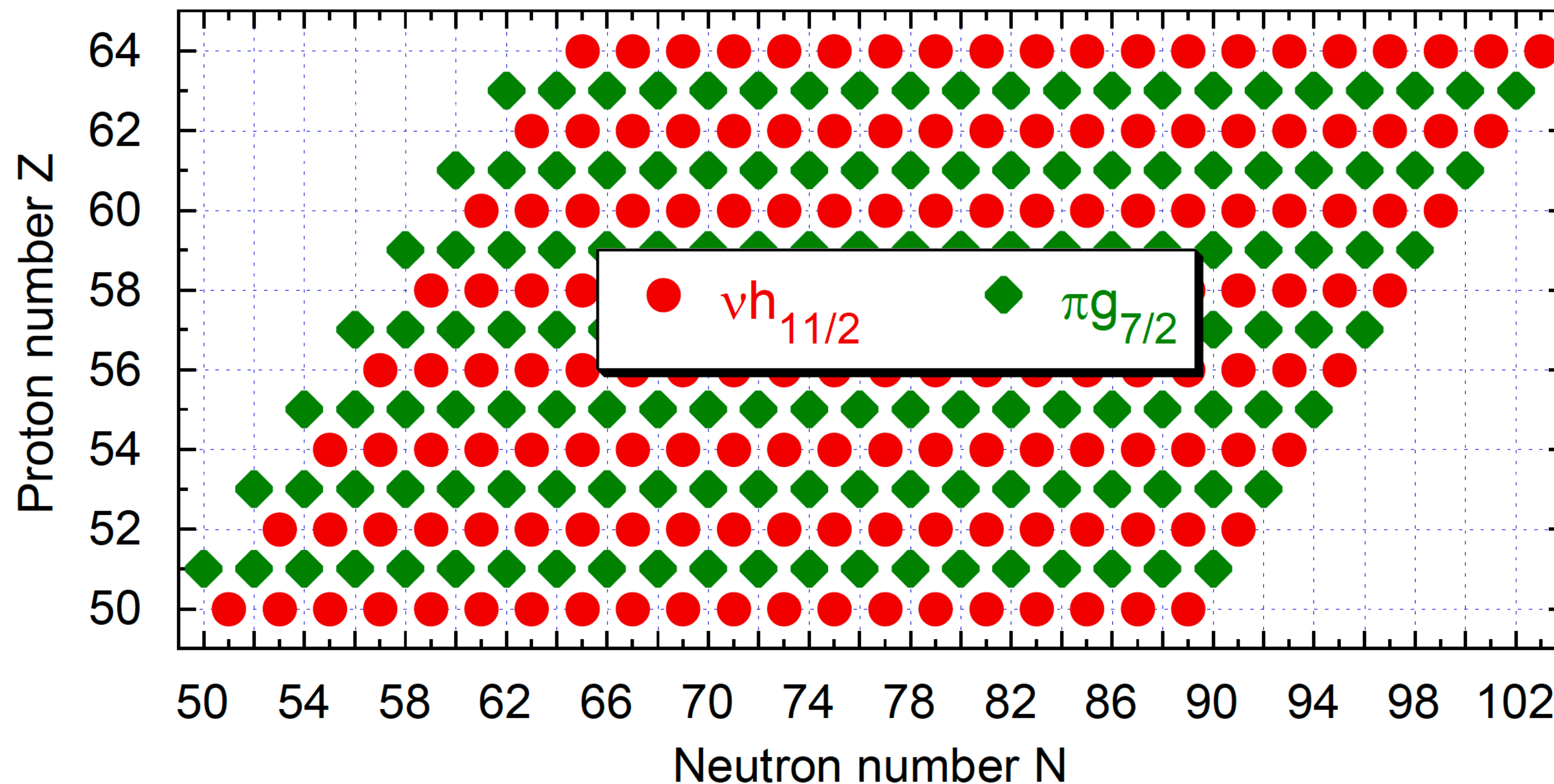


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P. L. Sassarini *et al.*, J. Phys. G: Nucl. Part. Phys. **49**, 11LT01 (2022)



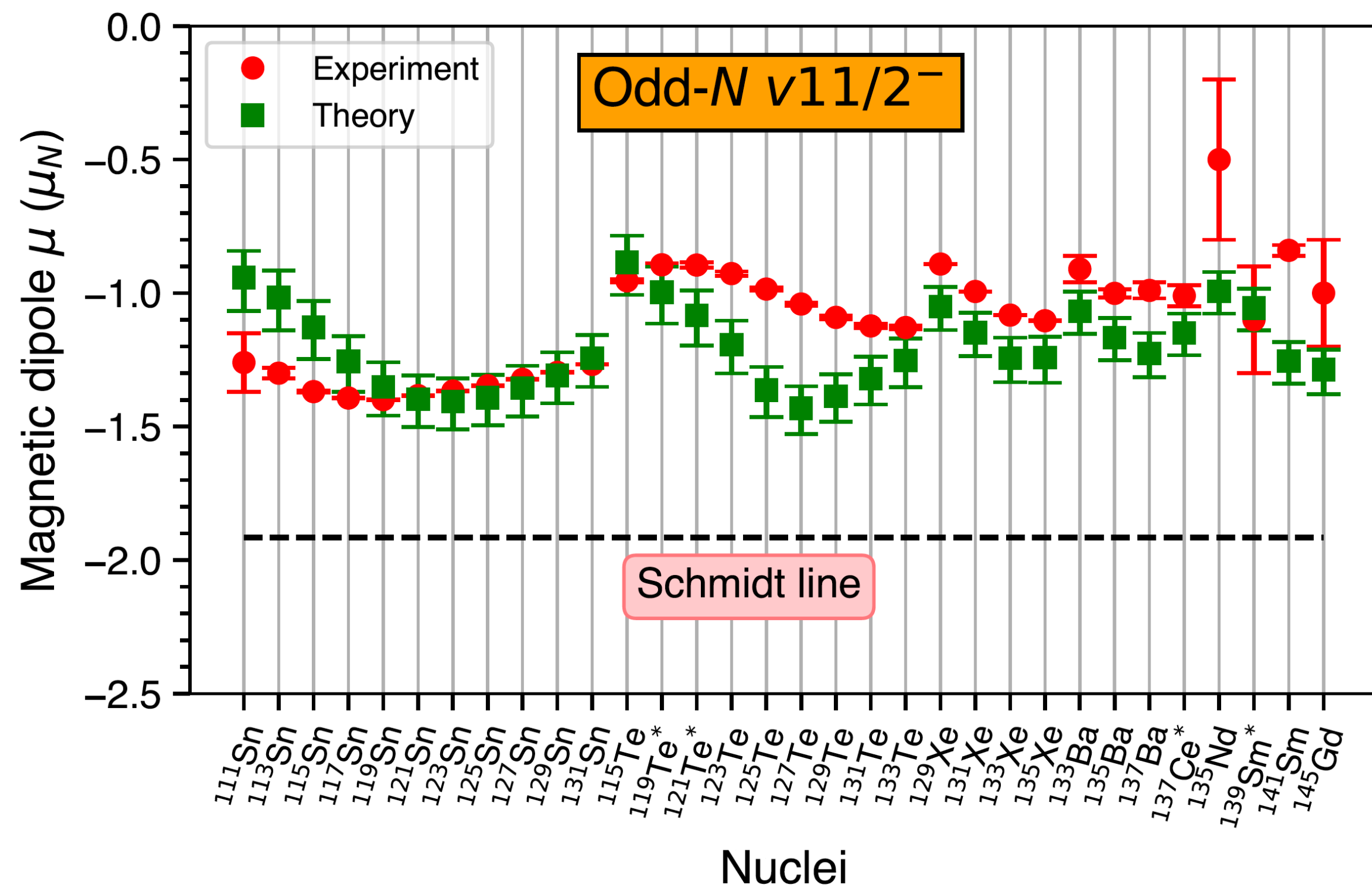
Systematic nuclear DFT calculations: Sn-Gd



H. Wibowo, *et.al.*, *arXiv:2503.15738*

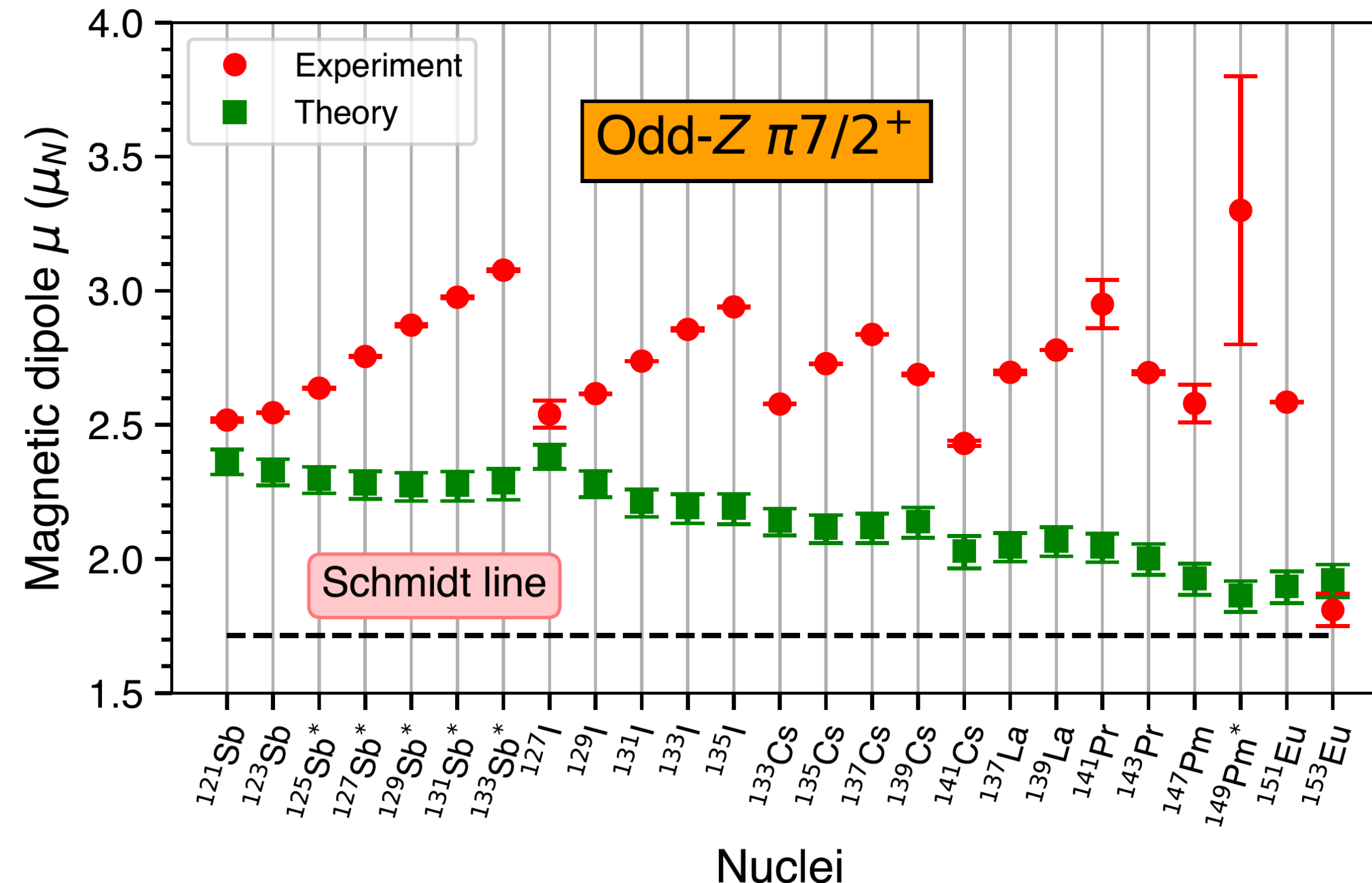
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Theory vs. experiment: Sn-Gd



$$g'_0 = 1.7(4) \text{ (UNEDF1)}$$

H. Wibowo, *et.al.*, *arXiv:2503.15738*



N. J. Stone, INDC, report INDC(NDS)-0658

N. J. Stone, INDC, report INDC(NDS)-0794

N. J. Stone, INDC, report INDC(NDS)-0816

Yordanov D. T. *et al.*, Comm. Phys. 3, 107 (2020)

Lechner S. *et. al.*, Phys. Lett. B 847, 138278 (2023)

Meson-exchange contributions

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Magnetic moment operator:

$$\hat{\mu} = \sum_k^A \hat{\mu}_{1b,k} + \sum_{k<\ell}^A \hat{\mu}_{2b,k\ell} + \cdots$$

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$$\hat{\boldsymbol{\mu}}_{1\text{b},k} = g_\ell^{(k)} \hat{\boldsymbol{\ell}}_k + g_s^{(k)} \hat{\boldsymbol{s}}_k$$

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Two-body meson-exchange:

$$\hat{\boldsymbol{\mu}}_{2\text{b}}^{\text{NLO}}(\mathbf{r}_1, \mathbf{r}_2) = \frac{1}{2} \int d^3\mathbf{x} \, \mathbf{x} \times \mathbf{j}_{2\text{b}}^{\text{NLO}}(\mathbf{x}, \mathbf{r}_1, \mathbf{r}_2)$$

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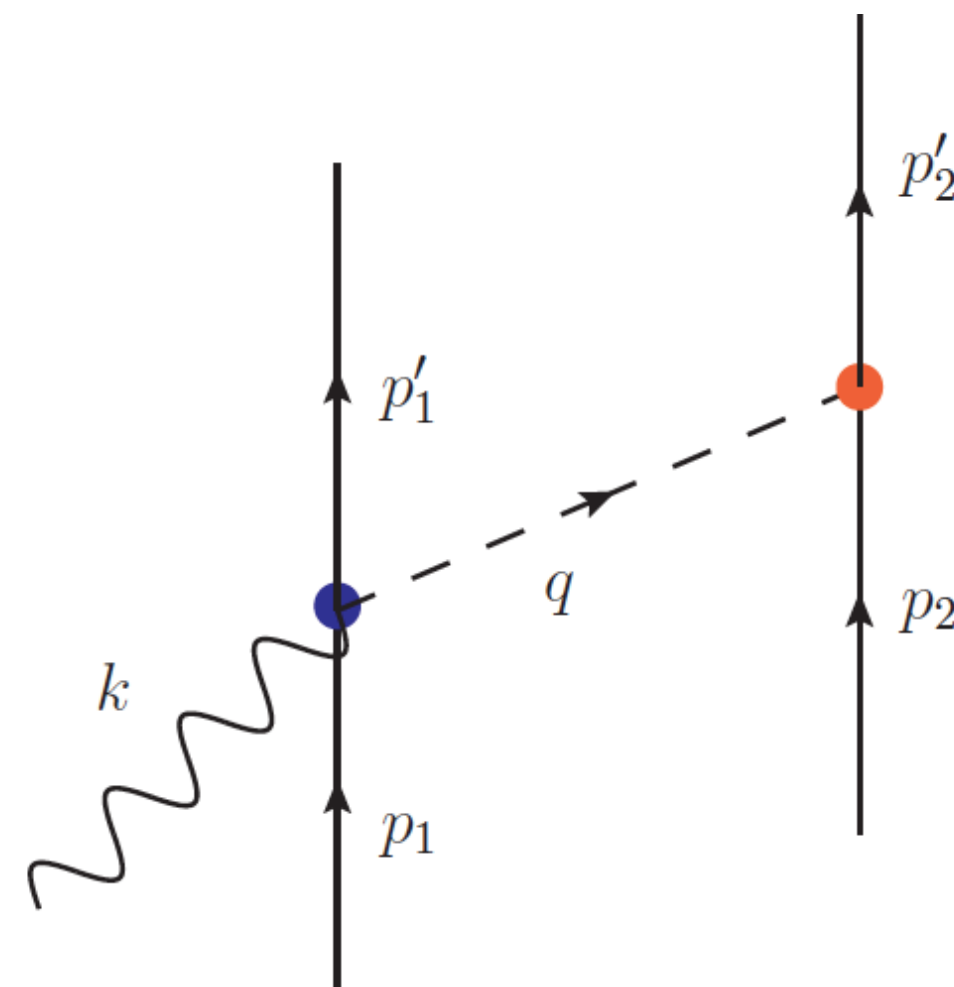
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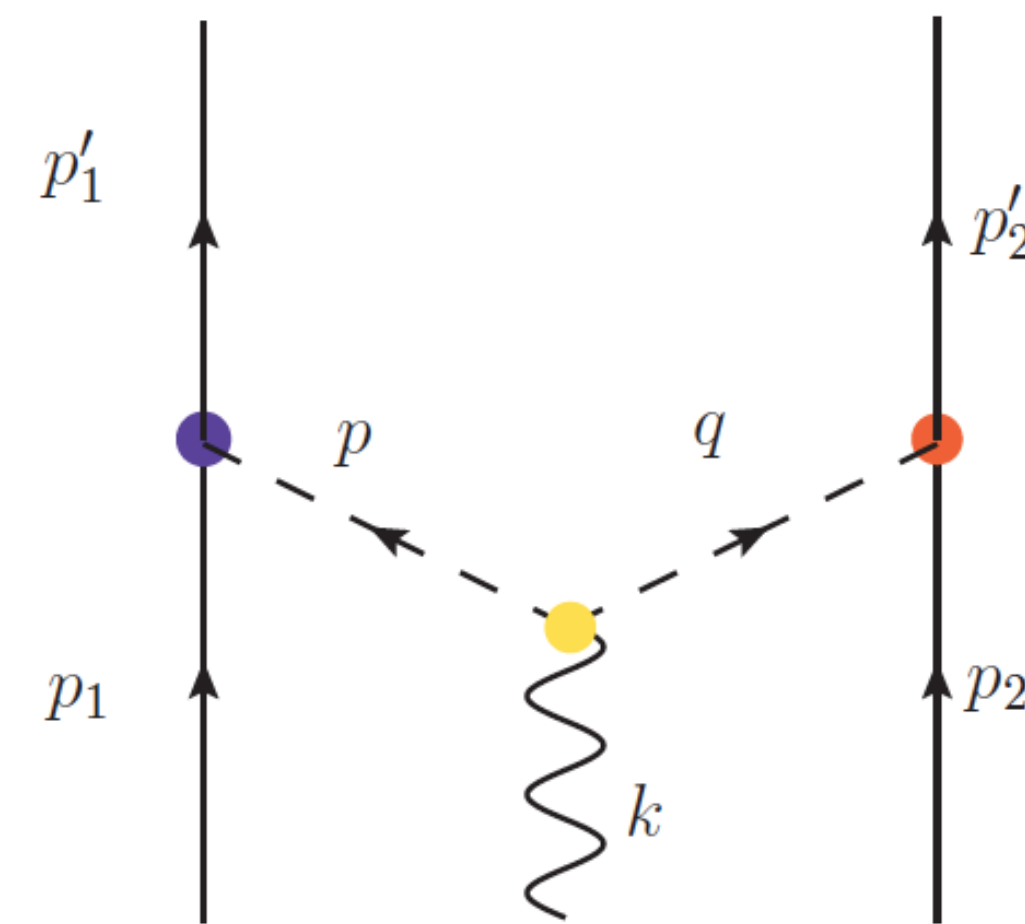
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Next-to-leading order (NLO):



Seagull graph



Pion-in-flight graph

R. Seutin, *et al.*, PRC **108**, 054005 (2023)

T. Miyagi, *et al.*, PRL **132**, 232503 (2024)

NLO magnetic moment operators

The NLO intrinsic and Sachs contributions to the magnetic moment operator are given by

$$\hat{\boldsymbol{\mu}}_{2b}^{\text{NLO, int}}(\mathbf{r}) = -\frac{g_A^2 m_\pi}{32\pi F_\pi^2} (\hat{\boldsymbol{\tau}}_1 \times \hat{\boldsymbol{\tau}}_2)_z \left\{ \left(1 + \frac{1}{m_\pi r} \right) [(\hat{\boldsymbol{\sigma}}_1 \times \hat{\boldsymbol{\sigma}}_2) \cdot \hat{\mathbf{r}}] \hat{\mathbf{r}} - (\hat{\boldsymbol{\sigma}}_1 \times \hat{\boldsymbol{\sigma}}_2) \right\} e^{-m_\pi r}$$

and ($g_A = 1.27$; $F_\pi = 92.3$ MeV; $m_\pi = 138.039$ MeV)

$$\hat{\boldsymbol{\mu}}_{2b}^{\text{NLO, Sachs}}(\mathbf{r}) = -\frac{1}{2} (\hat{\boldsymbol{\tau}}_1 \times \hat{\boldsymbol{\tau}}_2)_z V_{1\pi}(r) \mathbf{R}_{\text{NN}} \times \mathbf{r},$$

relative
coordinate

COM

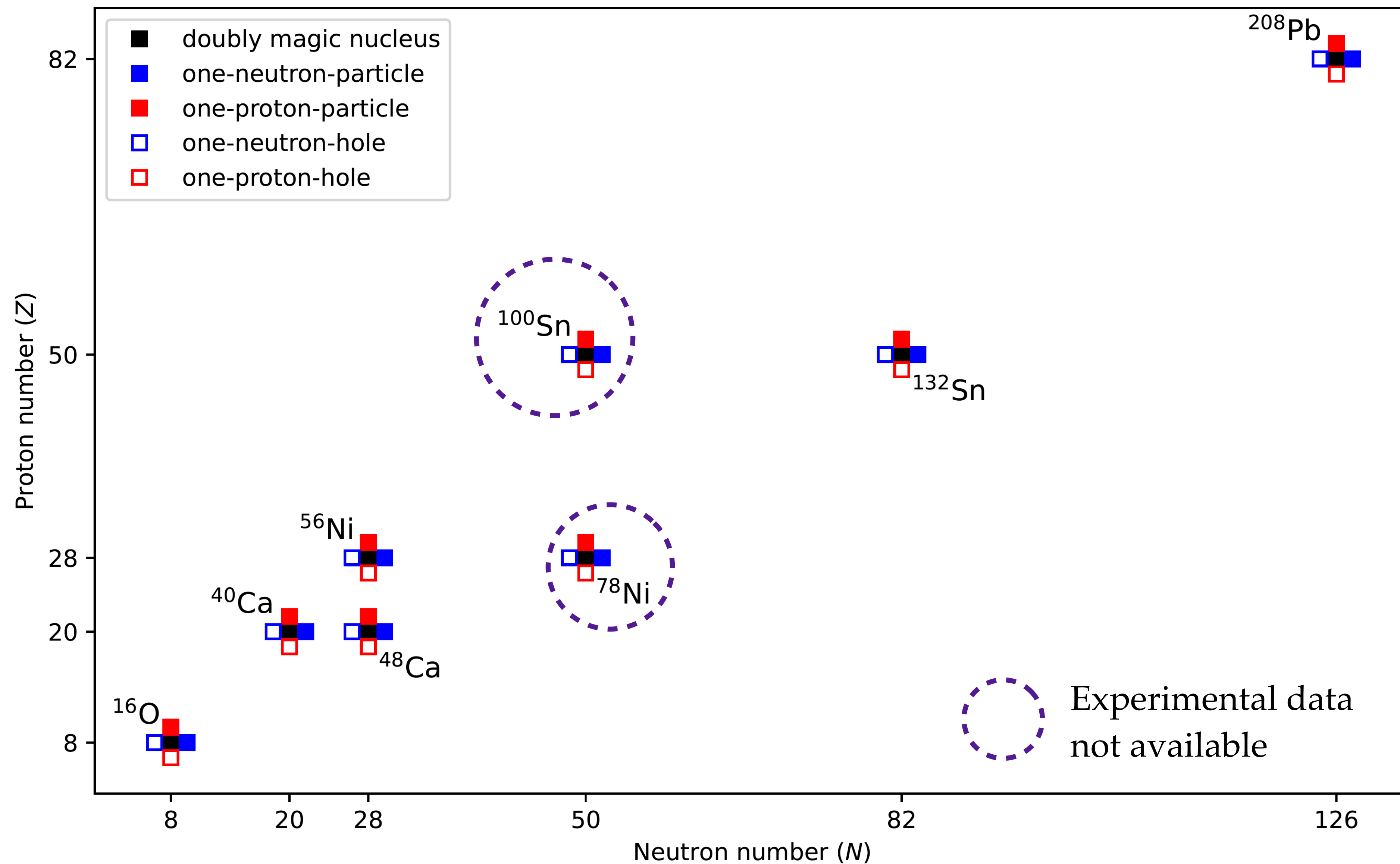
respectively, where

$$V_{1\pi}(r) = \frac{m_\pi^2 g_A^2}{48\pi F_\pi^2} \left[\hat{S}_{12} \left(1 + \frac{3}{m_\pi r} + \frac{3}{(m_\pi r)^2} \right) + \hat{\boldsymbol{\sigma}}_1 \cdot \hat{\boldsymbol{\sigma}}_2 \right] \frac{e^{-m_\pi r}}{r}.$$

Tensor operator:

$$\hat{S}_{12} = 3 (\hat{\mathbf{r}} \cdot \hat{\boldsymbol{\sigma}}_1) (\hat{\mathbf{r}} \cdot \hat{\boldsymbol{\sigma}}_2) - \hat{\boldsymbol{\sigma}}_1 \cdot \hat{\boldsymbol{\sigma}}_2.$$

Odd near doubly magic nuclei (two-body currents)



Optimum Landau parameter g'_0 (two-body currents)

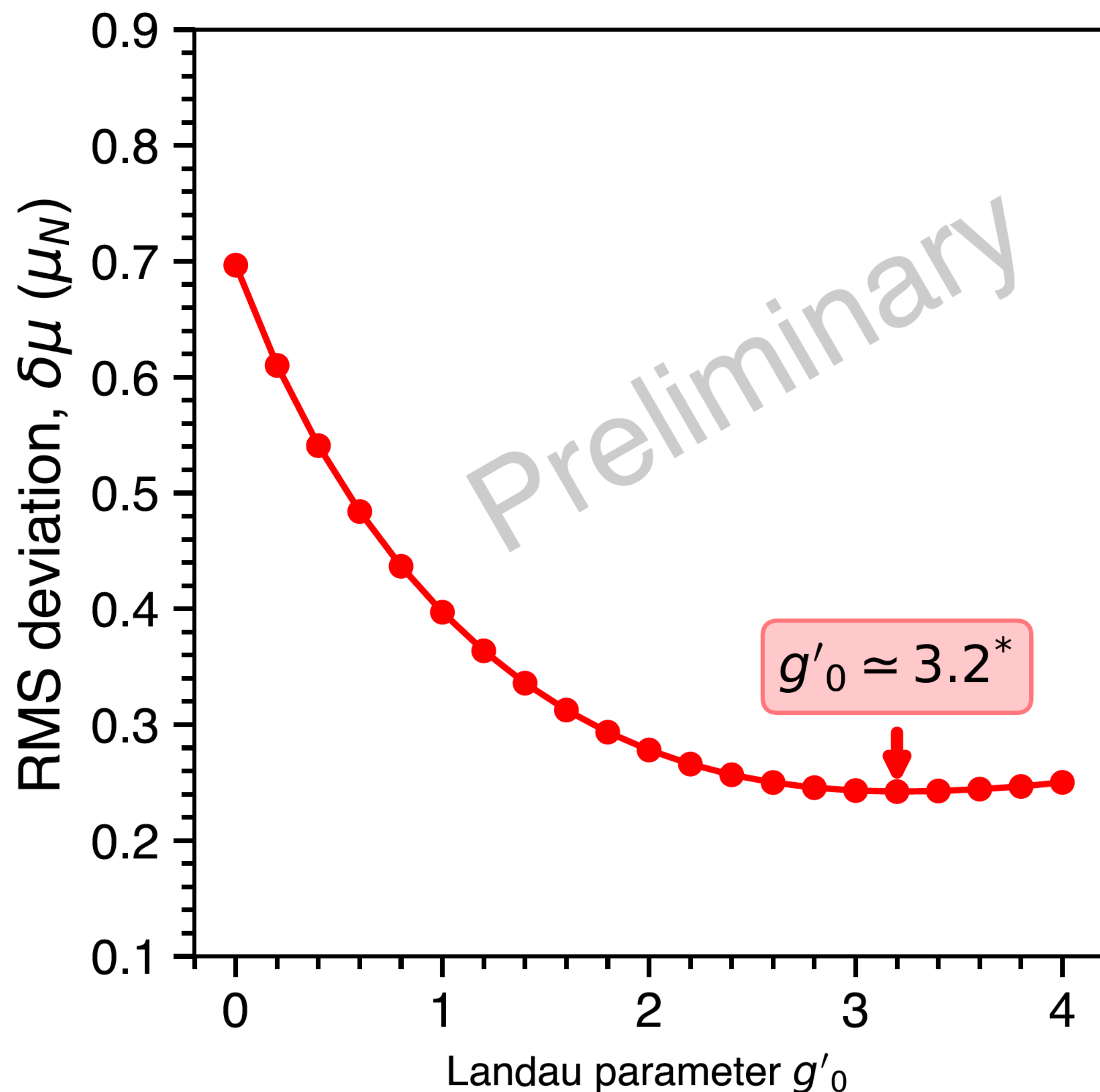
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$$\delta\mu(g'_0) = \sqrt{\frac{1}{N} \sum_{i=1}^N \left[\mu_{\text{calc}}(i, g'_0) - \mu_{\text{exp}}(i, g'_0) \right]^2}$$

i = odd nucleus

N = number of odd nuclei

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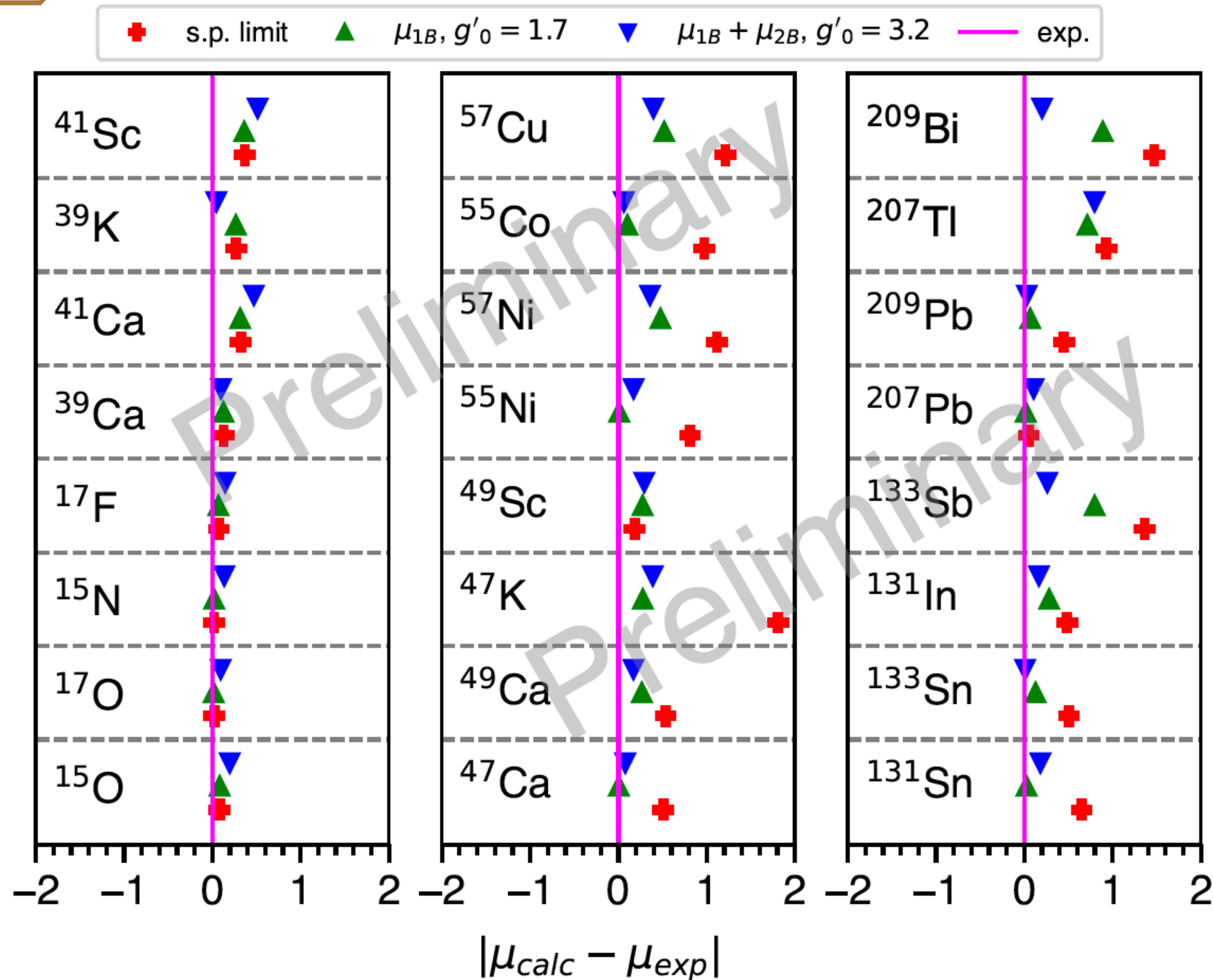
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(*) 23 odd near doubly magic nuclei without Bayesian analysis (**for now!**).

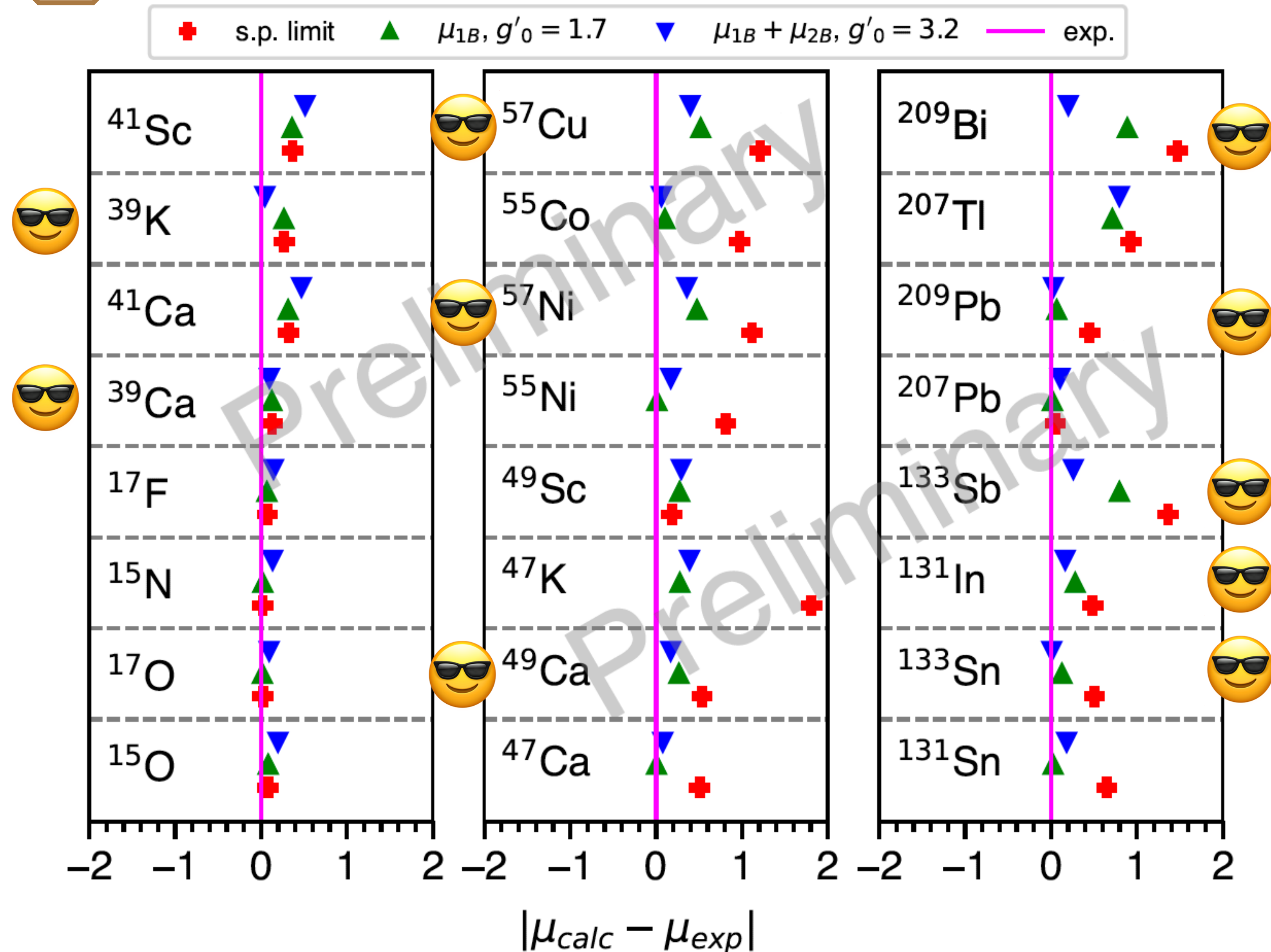
H. Wibowo, *et.al.*,
to be published

Magnetic dipole moments: theory vs. experiment



H. Wibowo, *et.al.*,
to be published

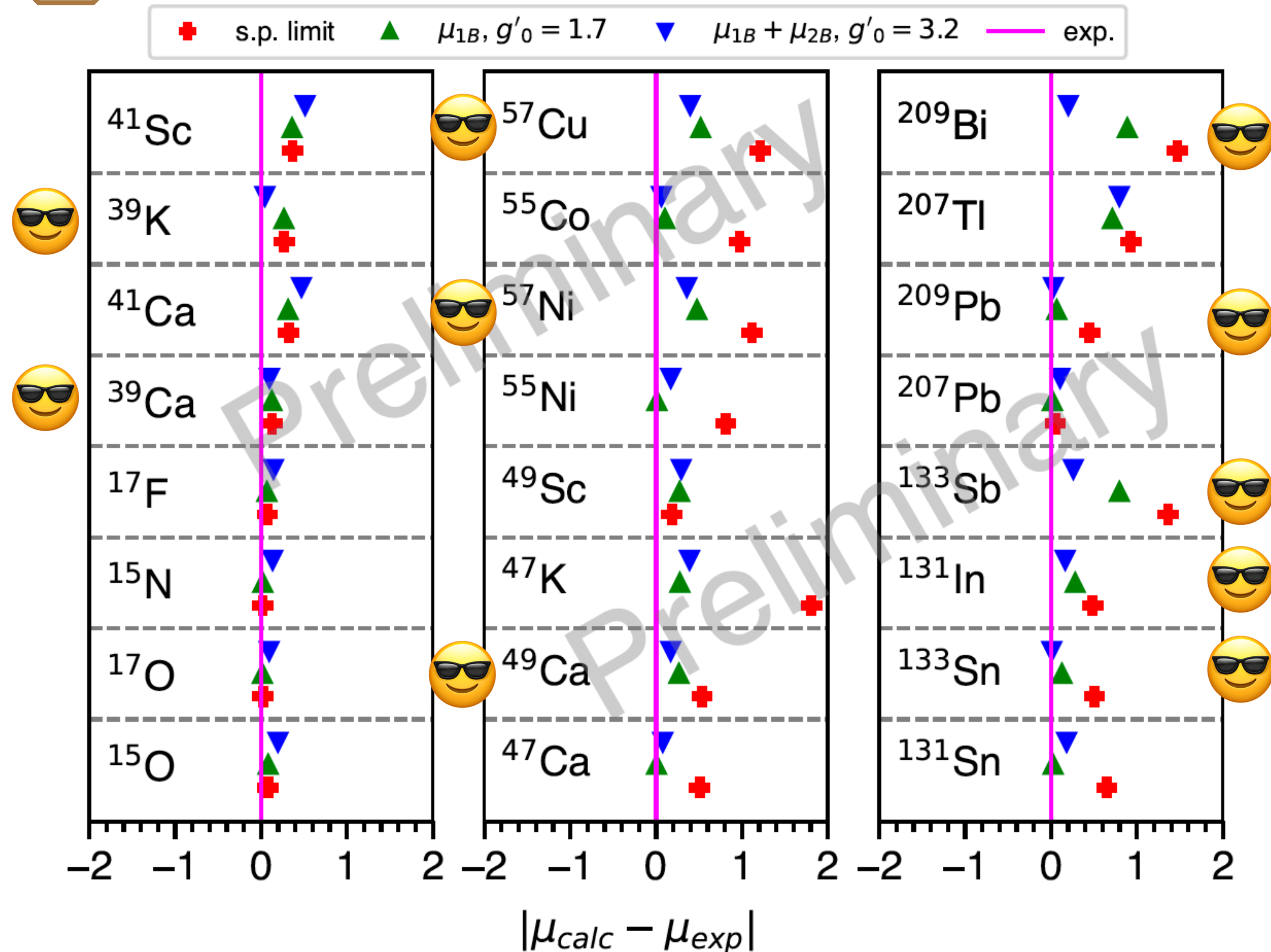
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★ In 10/14 cases, the two-body-current corrections **improve**/**deteriorate** agreement with experimental data.

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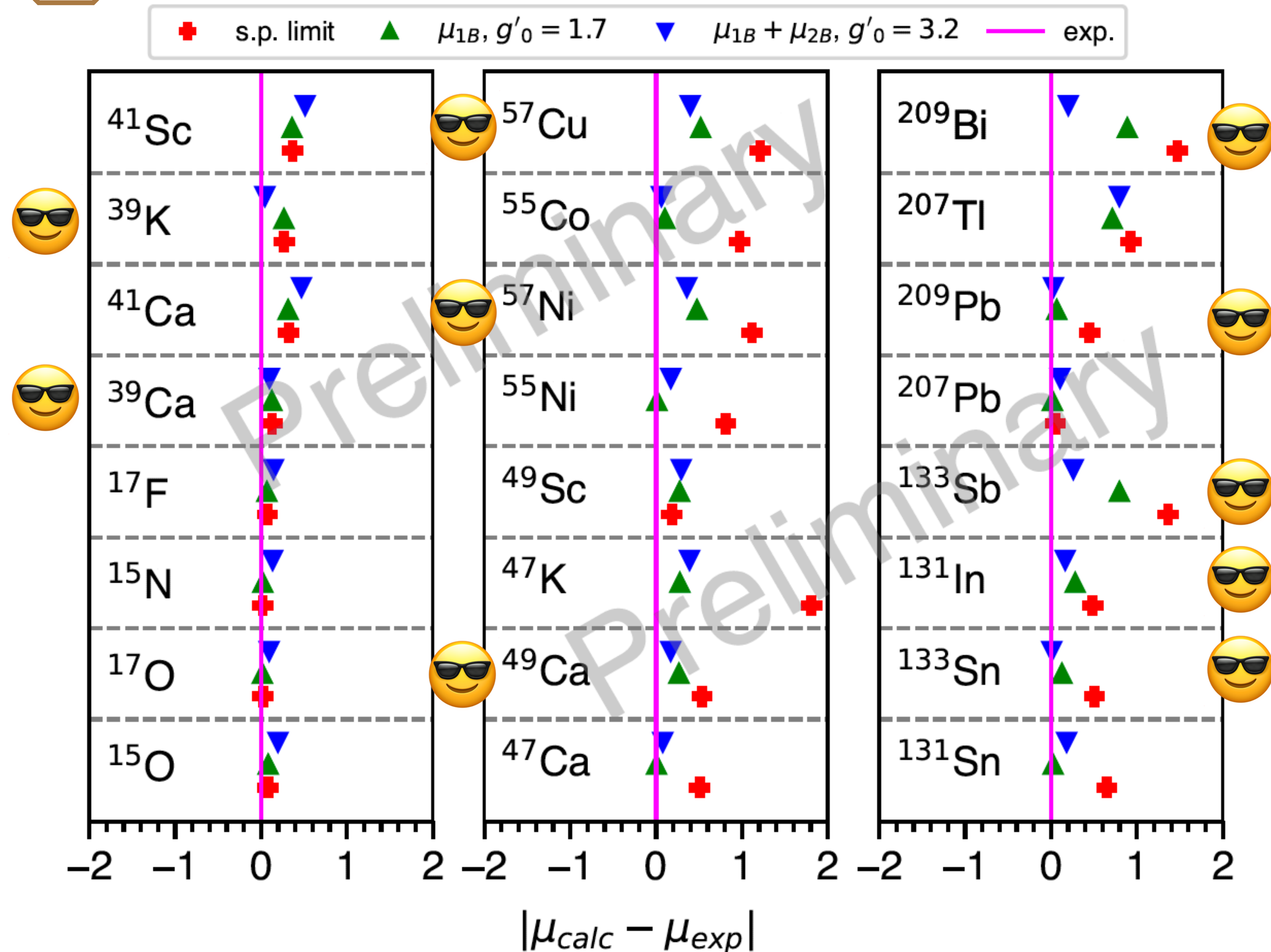
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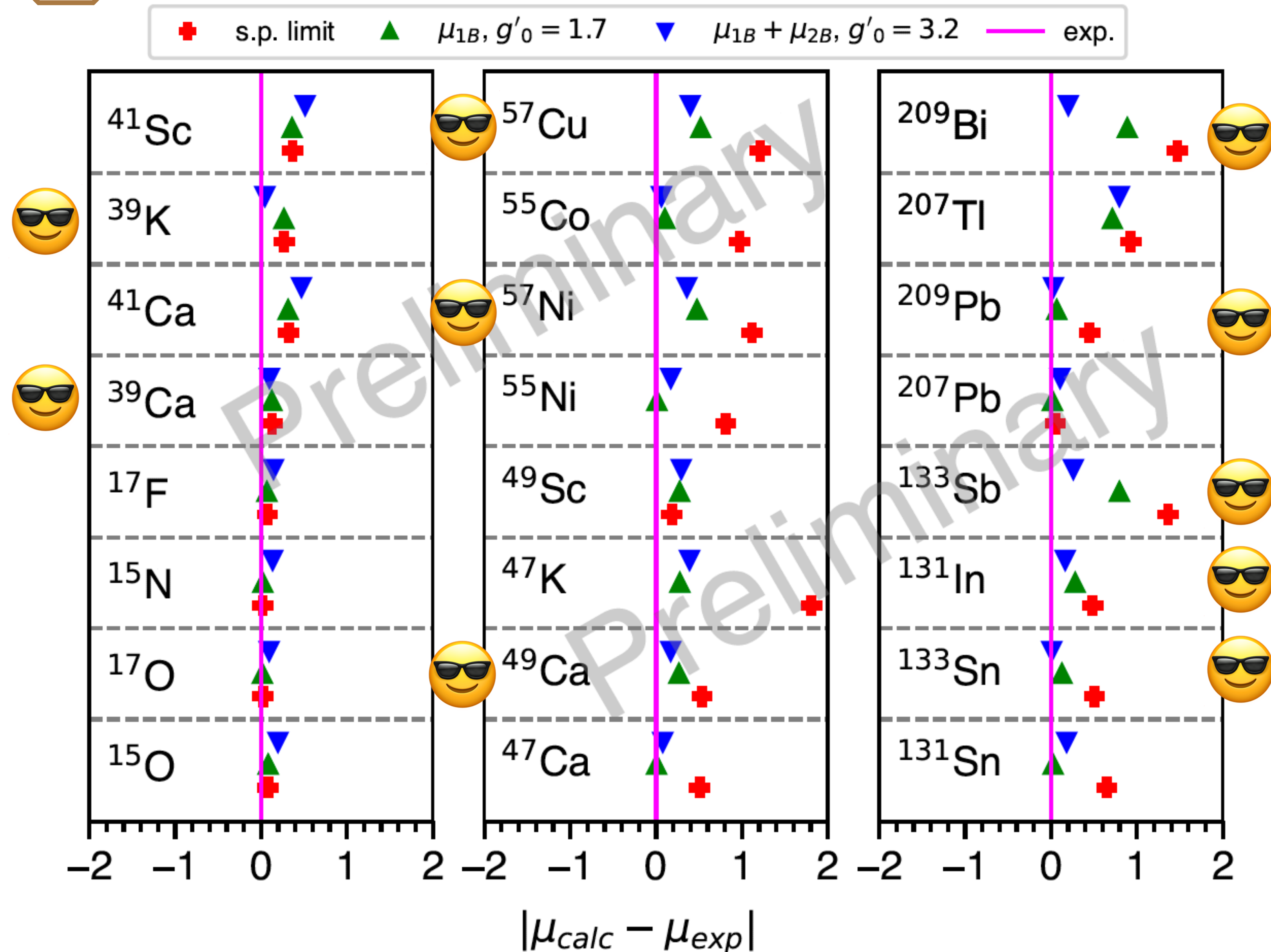


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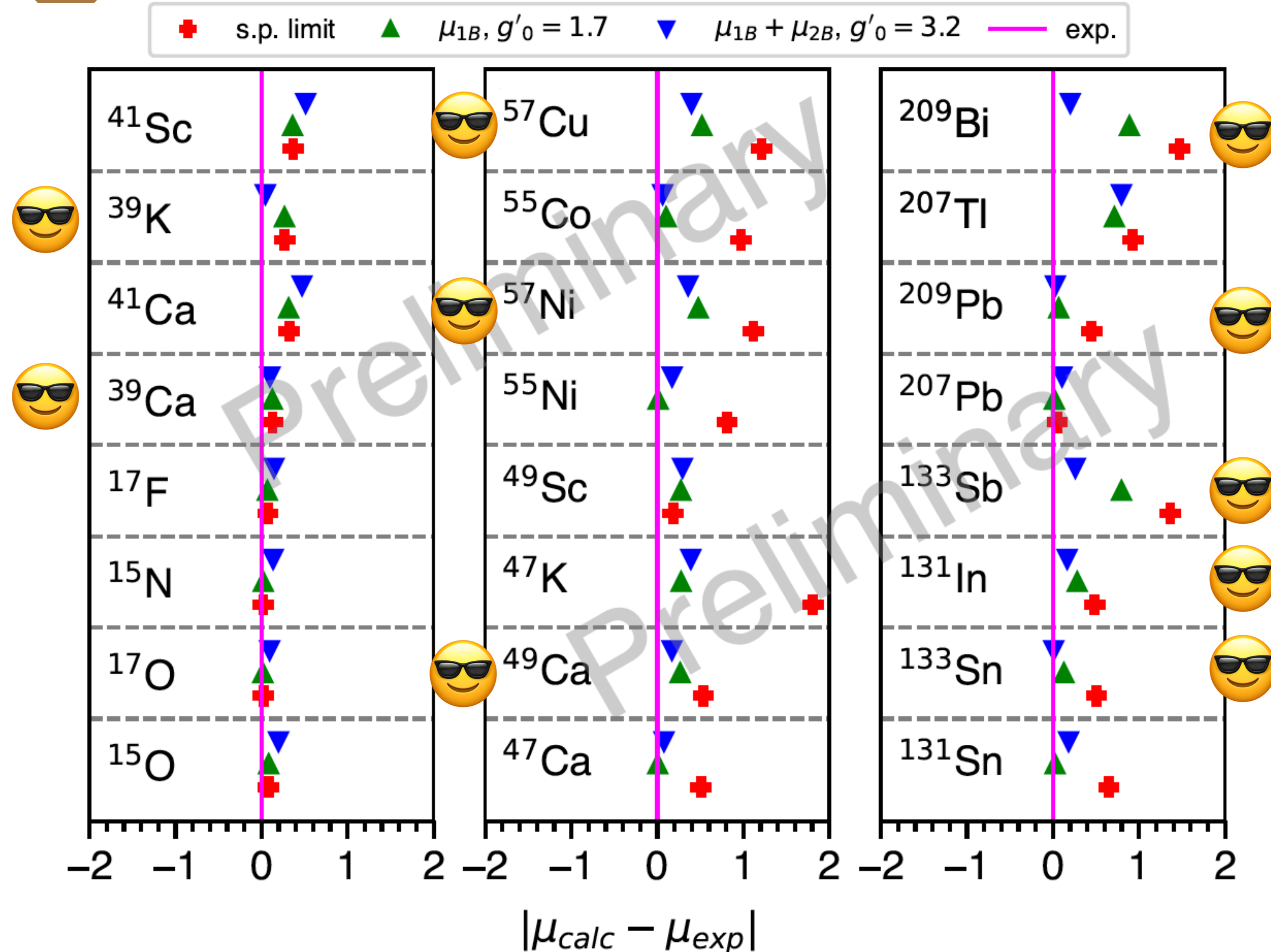


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★ The complete statistical analysis will follow.

★ Results in deformed nuclei will follow.

H. Wibowo, *et.al.*,
to be published

Summary and Outlook

Summary

- ★ The obtained value of the optimum Landau parameter g'_0 is around 3.2.
- ★ The inclusion of the meson-exchange currents improves the predictions of magnetic dipole moments in ^{39}K , ^{39}Ca , ^{57}Cu , ^{57}Ni , ^{49}Ca , ^{209}Bi , ^{209}Pb , ^{133}Sb , ^{131}In , and ^{133}Sn .

Outlook

- ★ Bayesian analysis to determine precisely the optimum Landau parameter g'_0 .
- ★ Systematic nuclear DFT calculations of magnetic dipole moments across the nuclear chart with the inclusion of meson-exchange currents.

Acknowledgements

We acknowledge the support from a **Leverhulme Trust Research Project Grant**. This work was partially supported by the **STFC Grant** Nos. ST/P003885/1 and ST/V001035/1 and by the **Polish National Science Centre** under Contract No. 2018/31/B/ST2/02220. We acknowledge the **CSC-IT Center for Science Ltd., Finland**, for the allocation of computational resources. This project was partly undertaken on the **Viking Cluster**, which is a high performance compute facility provided by the **University of York**. We are grateful for computational support from the **University of York High Performance Computing service, Viking and the Research Computing team**.