Lifetime measurements of low-lying octupole states in ²²⁴Ra

Dylan White

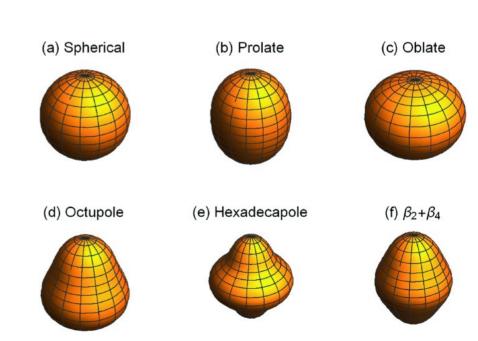




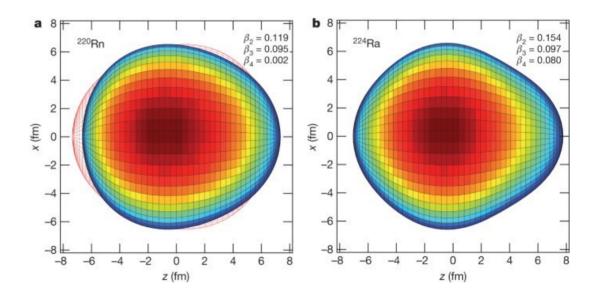


- Nuclei far away from closed shells tend to be deformed.
- Some of these nuclei become reflection asymmetric (pear shaped).
- These octupole deformed nuclei have displacement of their centres of mass and charge.

$$R(\theta, \phi) = c(\alpha_{\lambda\mu})R_0 \left[1 + \sum_{\lambda=0}^{\infty} \sum_{\mu=-\lambda}^{\lambda} \alpha_{\lambda\mu} Y_{\lambda\mu}(\theta, \phi) \right]$$



- Low-lying 1⁻ and 3⁻ states are clear signatures of enhanced octupole correlations.
- Enhanced E1 transitions can be a good indication of reflection asymmetry.
- Can also arise from octupole vibrational states.



L. P. Gaffney et al. Nature, 497(7448):199-204, 2013.

• Large intrinsic electric dipole moments are observed in octupole regions of the nuclear chart.

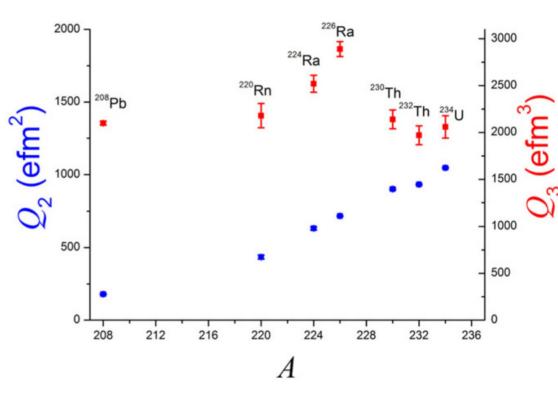
$$D_0 = e^{\frac{NZ}{A}} \left[\langle z_{p.c.m} \rangle - \langle z_{n.c.m} \rangle \right] \qquad D_0 = (\pm) D_0^{macro} + (\pm) D_0^{shell}$$

- Schiff's theorem tells us that for a point like nucleus, with a permanent EDM, the dipole moment is shifted to the atomic electrons and they will rearrange themselves to give a zero atomic EDM.
- Treating the nucleus as the finite object it is gives rise to the Schiff moment and an overall EDM.

$$S \approx 1.0 \times 10^{-4} \frac{I}{I+1} \beta_2(\beta_3)^2 Z A^{2/3} \frac{\text{[keV]}}{E_- - E_+} e \, \eta \, \text{[fm}^3\text{]}$$

- ²²⁴Ra has a very small intrinsic electric dipole moment due to the cancellation of the macroscopic and microscopic components.
- The odd mass nuclei around ²²⁴Ra are good candidates for a large Nuclear Schiff Moment.
- B(E1, 1⁻ → 0⁺) values are key predictions for DFTs.

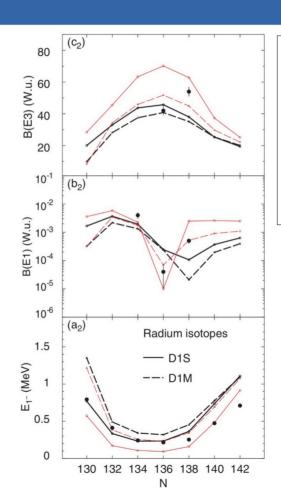
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P.A. Butler 2016 J. Phys. G: Nucl. Part. Phys. 43 073002.

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Bullet points = experimental results

Full black lines = Gogny D1S

Dotted black lines = Gogny D1M

Red lines = Generator Coordinate Method (GCM)

L. M. Robledo 2013 Physical Review C 88, 051302(R)

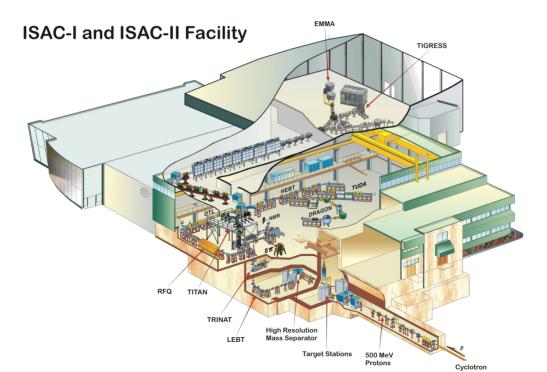
- Electric Monopole (E0) transitions are only observed in atomic nuclei.
- There appears to be missing states in ²²⁴Ra, when compared to ²²⁸Th.
- This could be the two-octupole phonon band head, measurement of the lifetime will help determine this.

$$\rho(E0) = \frac{< f |M(E0)|i>}{eR^2} \qquad \begin{array}{c} 0^+ & \text{1119.7 keV} & 0^+ & \text{0}^+ \\ & & & \\ 0^+ & & & \\ 0^+ & & & \\ \end{array}$$

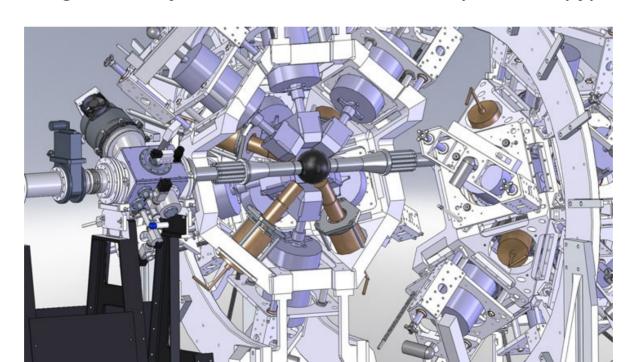
$$B(E0) = \rho^2(E0)e^2R^4$$
 okev okev

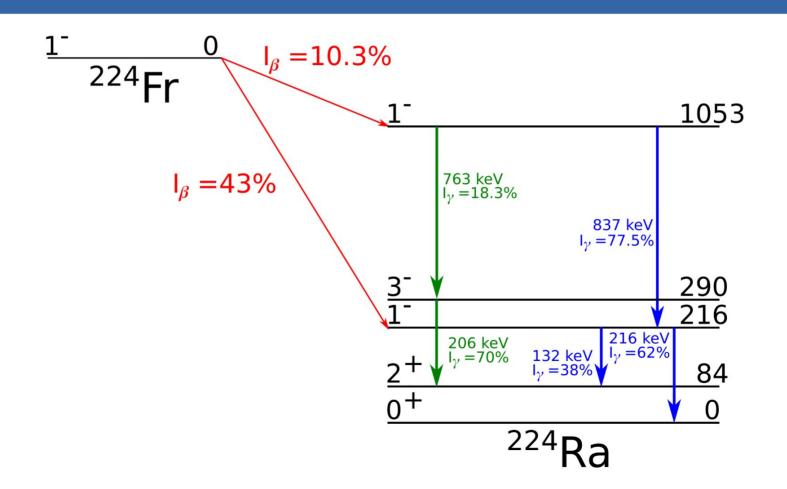
²²⁸Th ²²⁴Ra

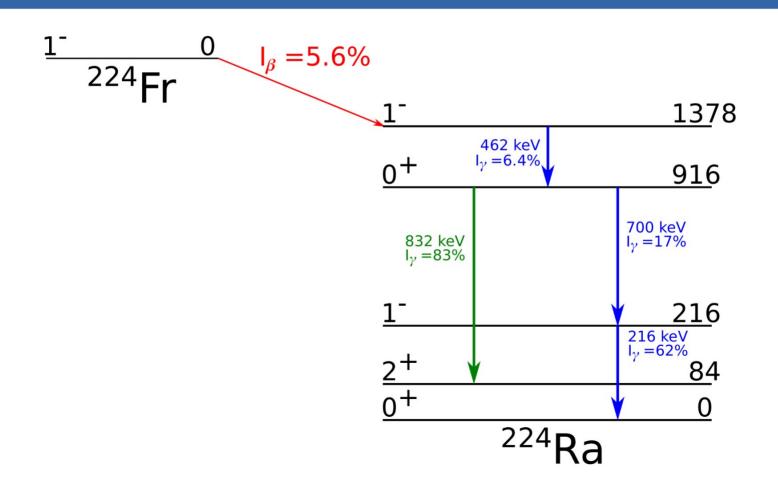
- Proton beam on UC_x target to make ²²⁴Fr ions.
- ²²⁴Ra populated from beta decay of ²²⁴Fr and then implanted into Moving Tape Collector.

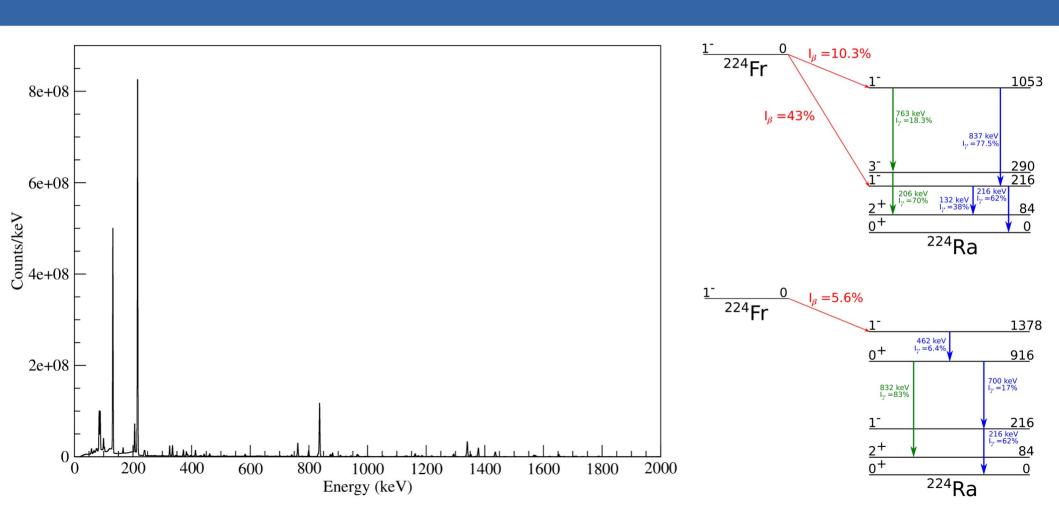


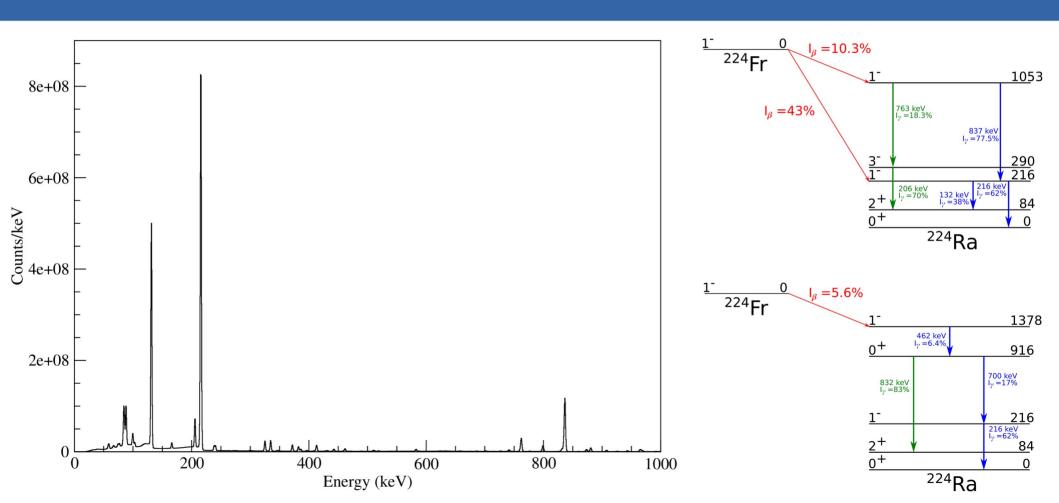
- GRIFFIN (15 HPGe Clover Detectors) (Gamma-Ray Infrastructure For Fundamental Investigations of Nuclei).
- 8 LaBr₃(Ce).
- PACES (Si(Li)) (Pentagonal Array for Conversion Electron Spectroscopy).

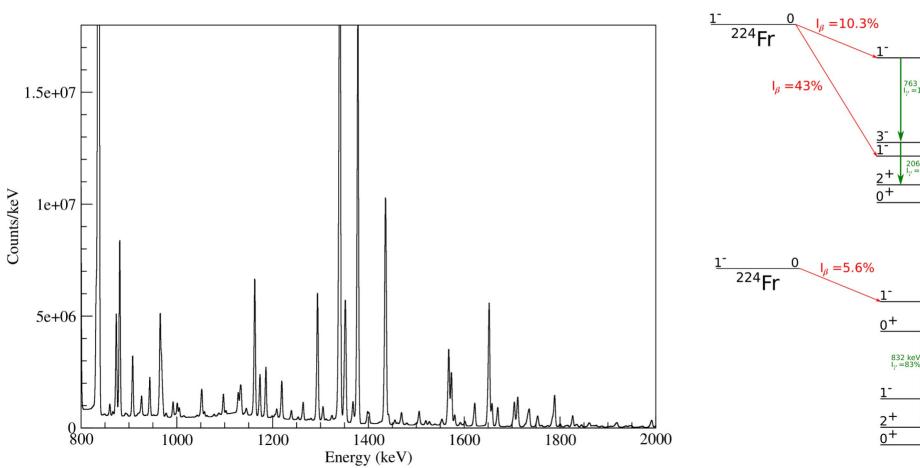


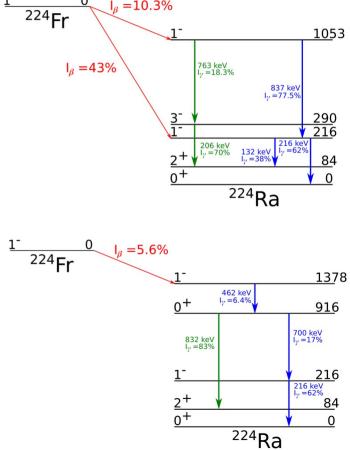


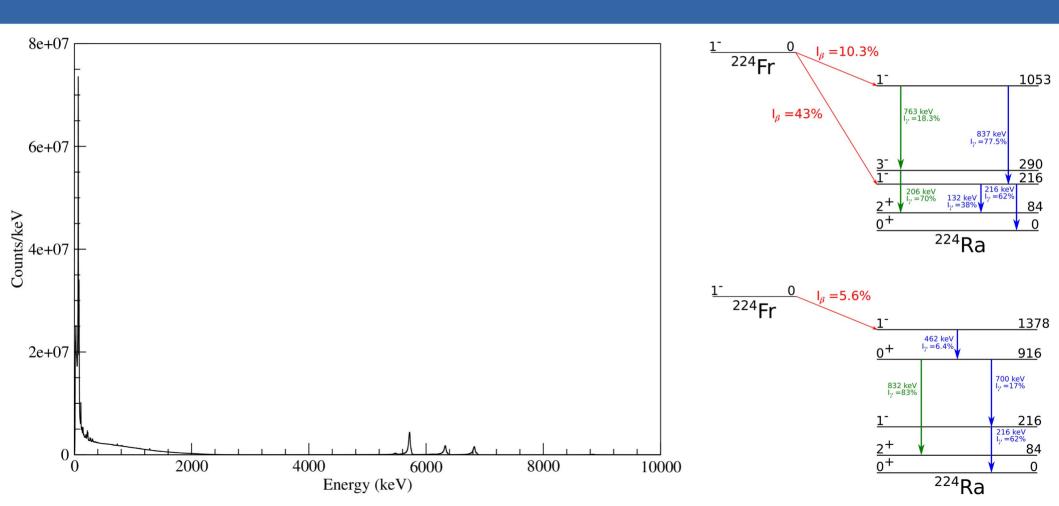


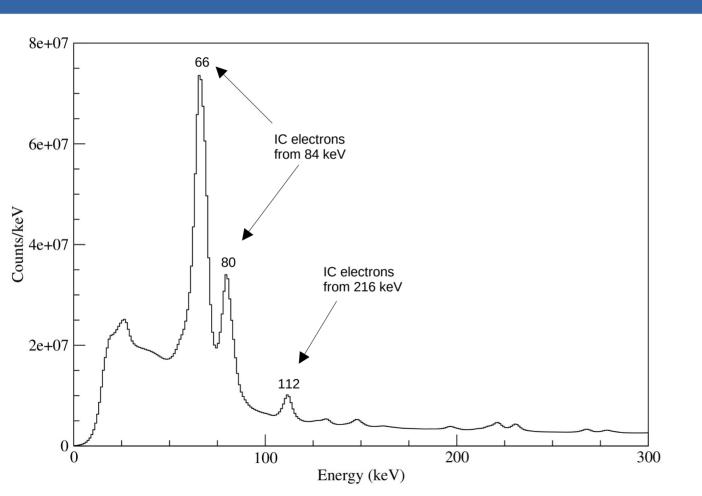


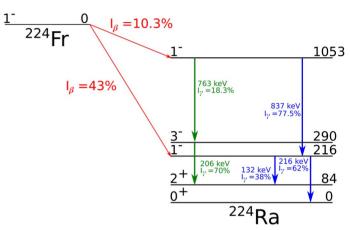


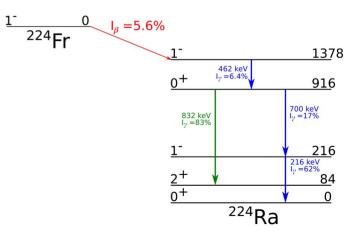


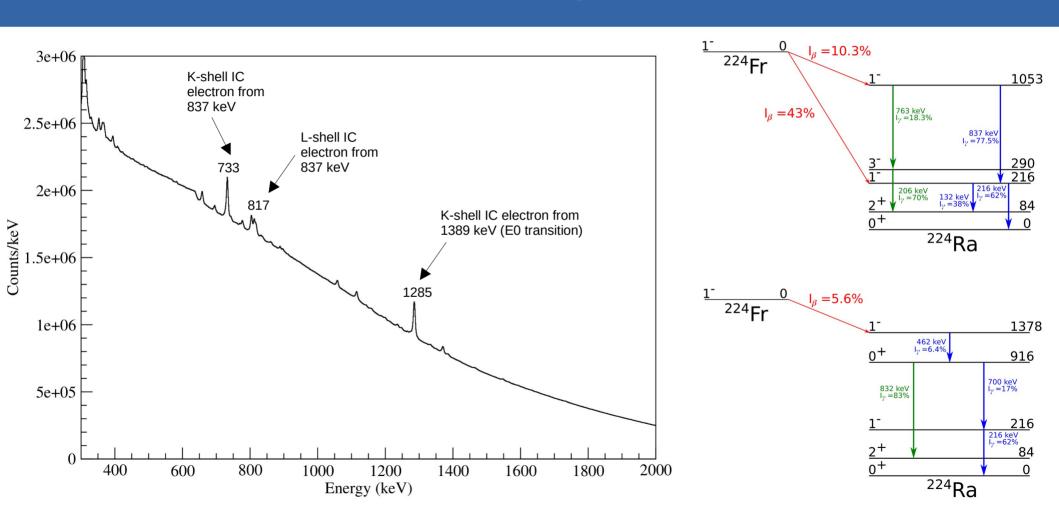


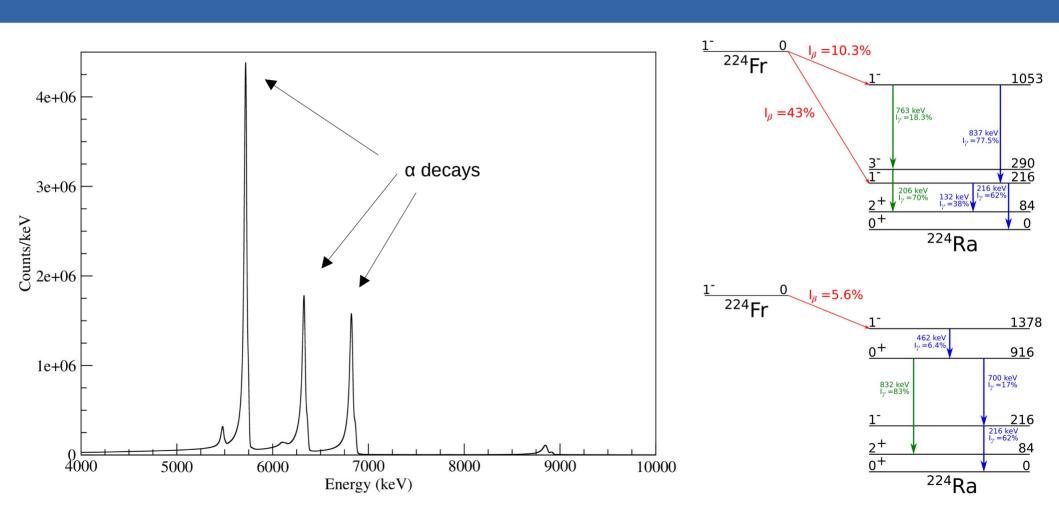




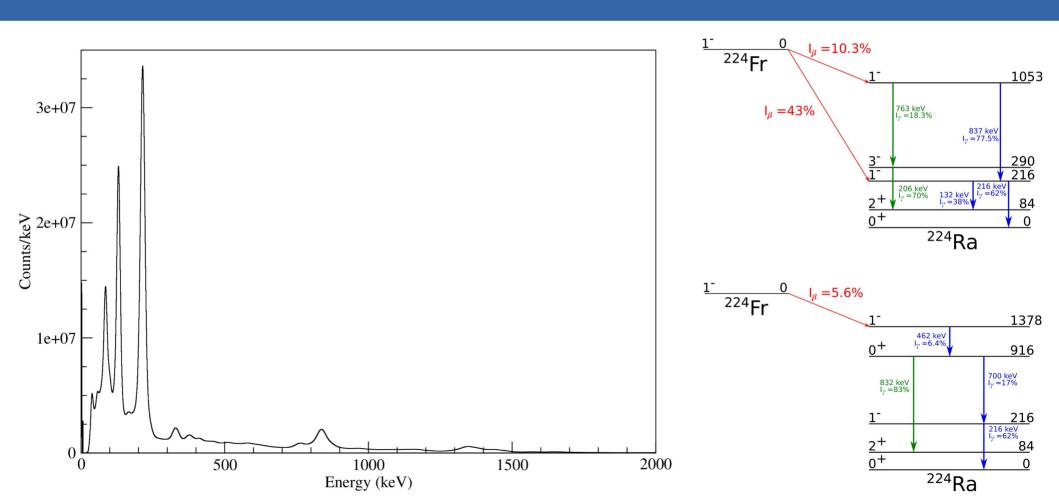




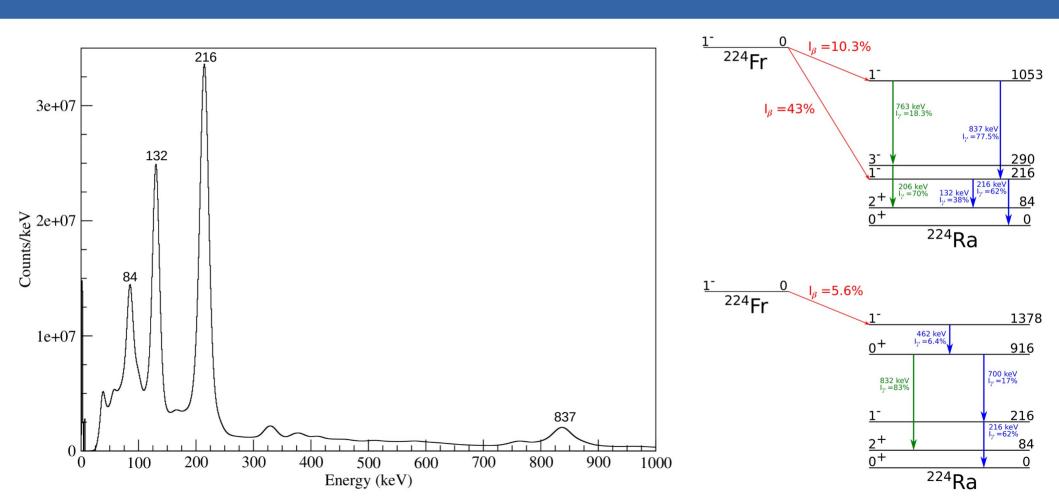




LaBr₃ Spectra



LaBr₃ Spectra



Convolution Method

• Fitting the product of a Gaussian and exponential function to a TAC curve makes it possible to determine the lifetime of the corresponding state.

$$f(x;\mu,\sigma,\lambda) = rac{\lambda}{2} \exp(rac{\lambda}{2}(2\mu + \lambda \sigma^2 - 2x)) \operatorname{erfc}\left(rac{\mu + \lambda \sigma^2 - x}{\sqrt{2}\sigma}
ight).$$

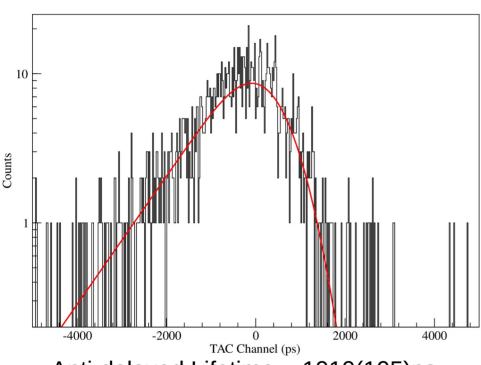
 λ = Rate of decay of the exponential.

 μ = Centroid of the Gaussian

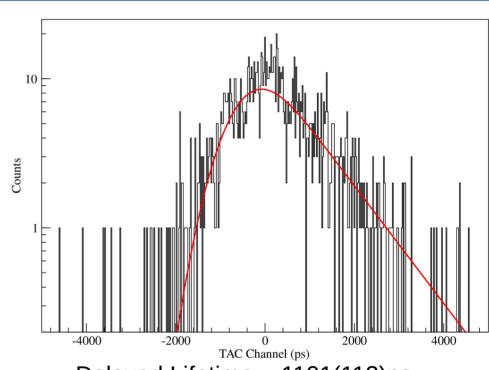
 σ^2 = Variance of the Gaussian

$$egin{aligned} ext{erfc}(x) &= 1 - ext{erf}(x) \ &= rac{2}{\sqrt{\pi}} \int_x^\infty e^{-t^2} \ dt \end{aligned}$$

Results: 2+ State



Anti-delayed Lifetime = 1010(105)ps

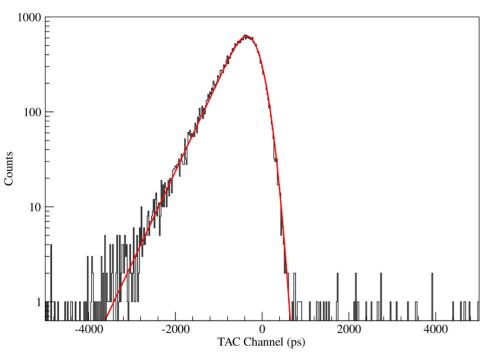


Delayed Lifetime = 1101(112)ps

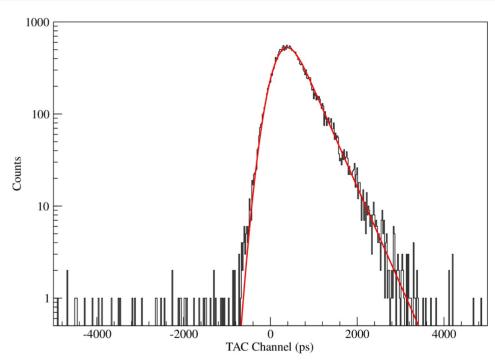
Average Lifetime = 1053(76)ps

NNDC Lifetime = 1079(27)ps

Results: 1-State



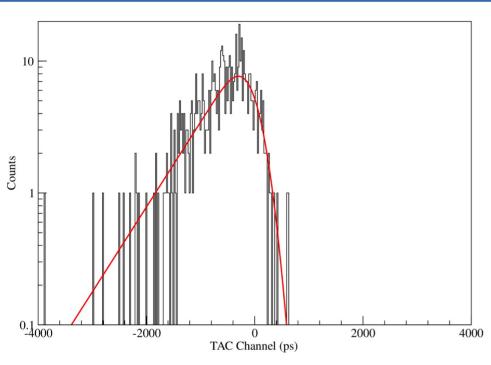
Anti-delayed Lifetime = 449(4)ps



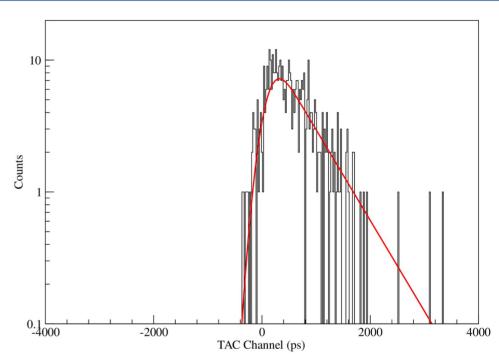
Delayed Lifetime = 407(4)ps

Average Lifetime = 430(3)ps

Results: 3-State



Anti-delayed Lifetime = 674(95)ps



Delayed Lifetime = 631(96)ps

Average Lifetime = 653(67)ps

Results

| \mathbf{J}^{π} | Average Lifetime (ps) | Lifetime Values from literature (ps) |
|--------------------|--------------------------|--|
| 2+ | 1053(76) | 1079(27) (NNDC) |
| 1 - | 431(10) | > 325 (Gaffney et al 2013) |
| 3- | 653(67) | 718(313) (Gaffney et al 2013) |

Results

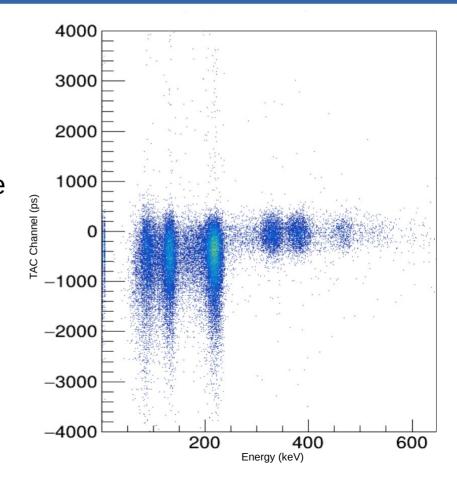
| ${\sf J}^\pi$ | B(E1) (e²fm²) | B(E1)/B _W (E1) (W.u.) | D ₀ (efm) |
|----------------|----------------------------------|-------------------------------------|-------------------------|
| 1 ⁻ | (8.46 ± 0.06) x 10 ⁻⁵ | (3.55 ± 0.02) x 10 ⁻⁵ | 0.0326 ± 0.0002 |
| 3- | $(1.0 \pm 0.1) \times 10^{-4}$ | $(4.22 \pm 0.4) \times 10^{-5}$ | 0.0313 ± 0.0032 |

GCD Method

- Used for shorter lived states that have TAC distributions too prompt for the convolution method.
- Fitting the delayed and anti-delayed TAC spectra and then calculating the difference between their centroids allows for calculation of the lifetime.

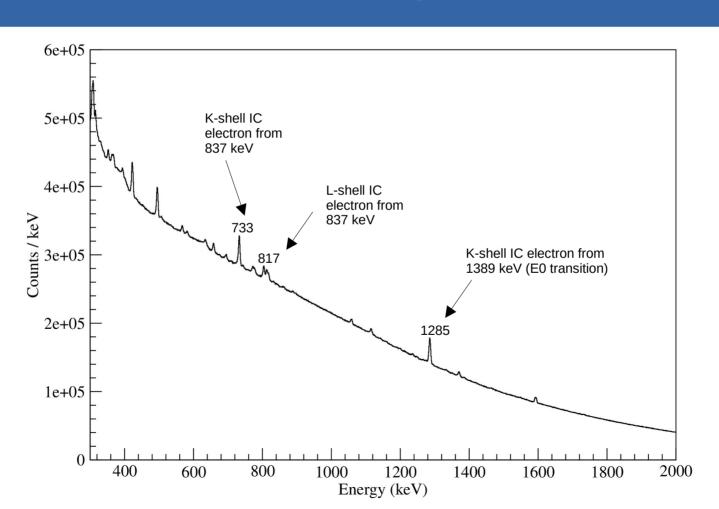
$$\Delta C(E_1, E_2) = C_D - C_A = PRD(E_1, E_2) + 2\tau$$

 The prompt response difference allows us to correct for time walk.



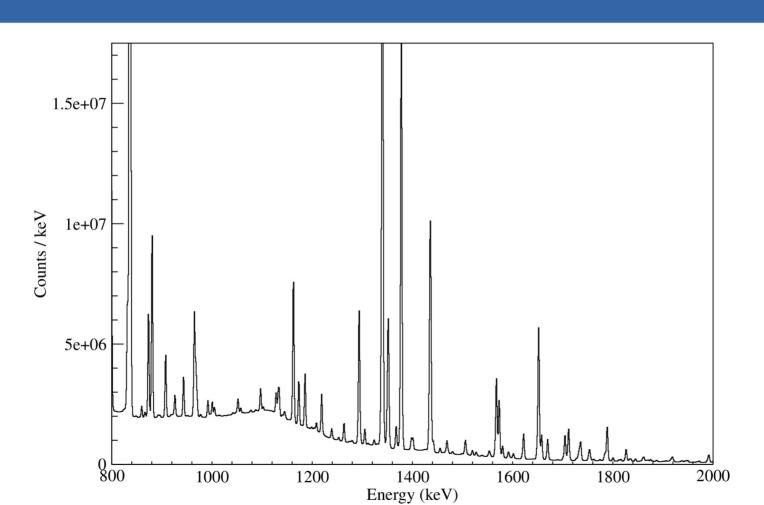
Conclusions and Future Work

- Lifetime measurement of the 2⁺ state agrees with the accepted value from the NNDC.
- Lifetime measurements of the 1⁻ and 3⁻ states agree with Gaffney et al.
- Measure the lifetime of the 0⁺ state.
- Use the GCD method to compare results and measure shorter lived states.
- Look for E0 transitions in the PACES spectra.
- Identify unknown transitions from GRIFFIN spectra.



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- 6. CSIC

Thanks for listening

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ZDS Spectra

