

Symmetry energy in dilute and dense matter

and the PREX-CREX puzzle

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01 Introduction

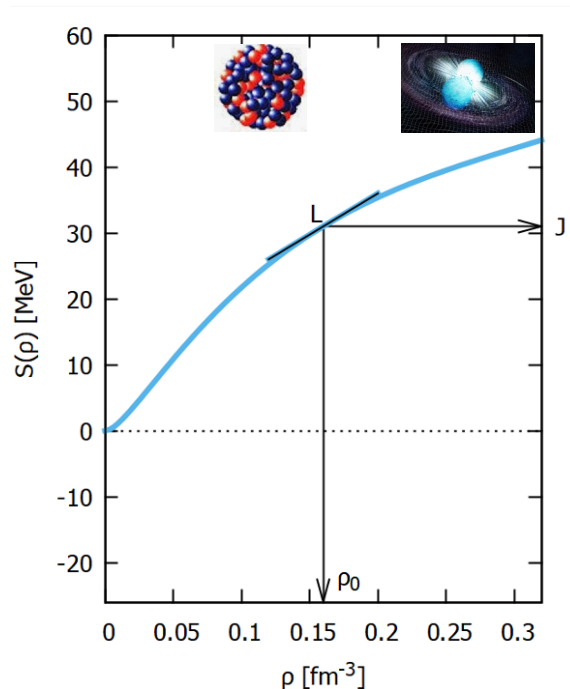


Symmetry energy

- ❑ The energy it costs to change isospin-symmetric nuclear matter into neutron matter
- ❑ Determines the physics of exotic nuclei and compact stars
- ❑ By convention, the density dependence of the symmetry energy is encoded in a few parameters defined by a Taylor expansion:

$$S(\rho) = J + Lx + \frac{1}{2}K_{sym}x^2 + \frac{1}{6}Q_{sym}x^3 + \dots; x \equiv (\rho - \rho_0)/3\rho_0$$

- ❑ Namely its value J , slope L , curvature K_{sym} , skewness Q_{sym} , etc, at the nuclear saturation density ρ_0 .





Symmetry energy

It is explored at different densities by examining various observables:

- ❑ Low density: nuclear masses, density distributions (skins), collective excitations
- ❑ High density: Heavy-ion collisions, compact stars

☞ *Multi-messenger astronomy*

- ❑ Ab initio calculations are also available

☞ *Variational Monte Carlo, Quantum Monte Carlo, chiral EFT, ...*

Rough estimates:

- $J = 30 - 33$ MeV (best known)
- $L = 40 - 70$ MeV (conflicting results)
- $K_{sym} = -400 - 0$ MeV (very roughly)
- Q_{sym} = practically unconstrained

Some reviews:

Baldo&Burgio, ProgPartNuclPhys91,203 (2016)

Dertel et al., Rev.Mod.Phys.89,015007 (2017)

Roca-Maza&Paar, ProgPartNuclPhys101,96 (2018)

Related analyses:

Li et al., Universe 7, 182 (2021)

Thi et al., Universe 7, 373 (2021)

+multiple others

01 Introduction



Connecting nuclear data and the equation of state: energy density functionals

- Problem:

Majority of EDFs fitted to data extrapolate to unrealistic EoSs

Imposing realistic EoS parameters to EDFs can give strange nuclear results

- Overcome by:

KIDS - a versatile framework for the nuclear EoS and EDF

- ❑ Grounded in theory of interacting Fermi systems
- ❑ Adjustable number of parameters depending on application

(optimal for bulk nuclear properties and neutron stars: 3+4 EoS coefficients)

- ❑ In medium effective mass is decoupled from the EoS coefficients
- ❑ No forced, unphysical correlations

Dutra et al, PRC85,035201

Stevenson et al., AIP Conf.Proc 1529,262

Roca-Maza&Paar, ProgPartNuclPhys101,96

Reviewed in:

PP&Hyun,

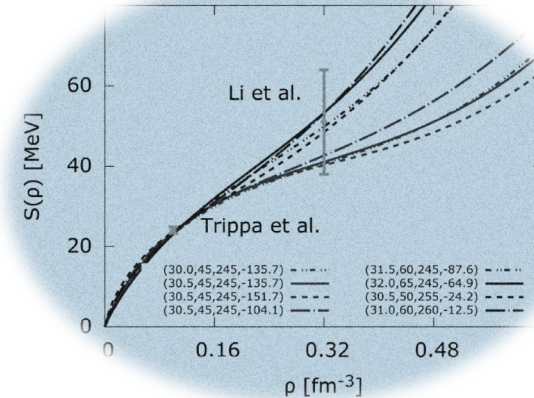
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Easy to use - like an extended Skyrme functional

02 Curvature parameter K_{sym} of the symmetry energy

Next: Selected results from analyses of nuclear and neutron star properties

- Xu&PP, Phys. Rev. C 105, 044305
- Zhou,Xu,PP, Phys. Rev. C 107, 055803
- Gil et al., IJMPE 31, 2250013
- Gil et al, Phys. Rev. C 103, 034330



02 Curvature parameter K_{sym} of the symmetry energy

Bayesian analysis of isovector nuclear properties

- ❑ Examined ^{208}Pb and ^{120}Sn
- ❑ Results for KIDS are compared with standard Skyrme-Hartree-Fock

Isovector constraints: Neutron skin thickness, giant dipole resonance, dipole polarizability

	Δr_{np} (fm)	E_{-1} (MeV)	α_D (fm ³)
^{208}Pb	0.283 ± 0.071	13.46 ± 0.10	19.6 ± 0.6
^{120}Sn	0.150 ± 0.017	15.38 ± 0.10	8.59 ± 0.37

Isoscalar constraints: mass, charge radius, energy of the isoscalar giant monopole resonance

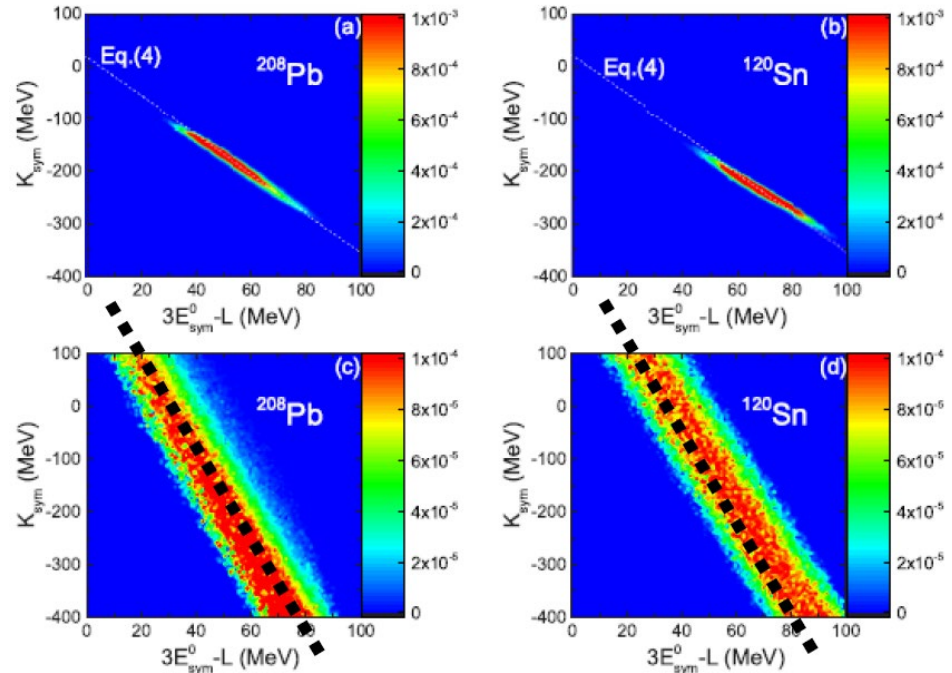
02 Curvature parameter K_{sym} of the symmetry energy

Bayesian analysis of isovector nuclear properties

K_{sym} vs $3J-L$

Skyrme:

KIDS :
Broader PDFs
 K_{sym} not constrained



02 Curvature parameter K_{sym} of the symmetry energy

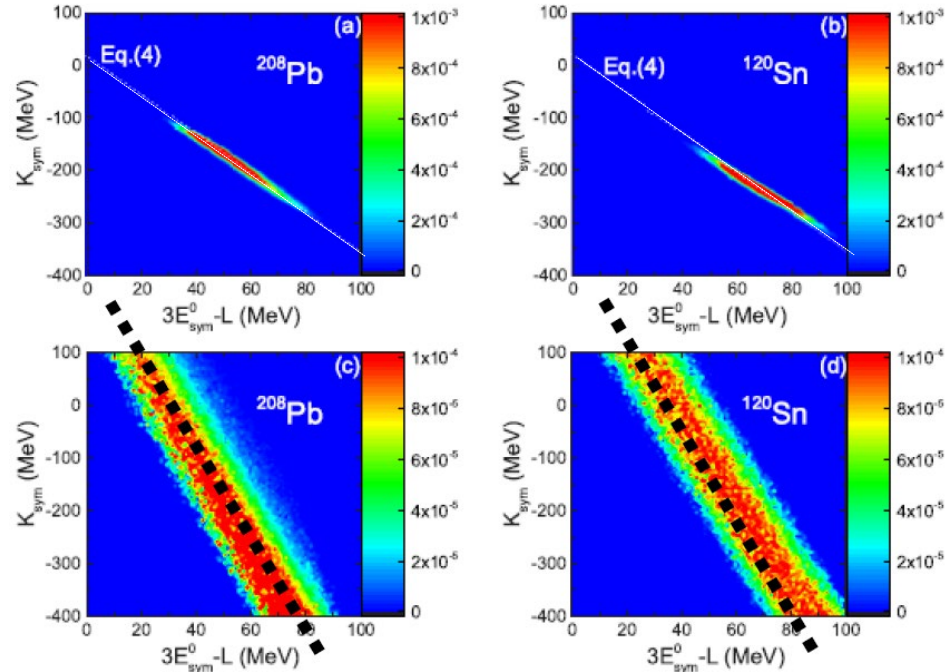
Bayesian analysis of isovector nuclear properties

$$K_{\text{sym}} = (2 - 3\alpha) \frac{2}{3} \left(\frac{3\pi^2}{2} \right)^{2/3} \frac{\hbar^2}{2m} \rho_0^{2/3} \\ \times \left[-3 \left(\frac{m}{m_s^*} - 1 \right) + 4 \left(\frac{m}{m_s^*} - 1 \right) \right] \\ - 3(1 + \alpha) (3E_{\text{sym}}^0 - L) \\ + (1 + 3\alpha) \frac{1}{3} \left(\frac{3\pi^2}{2} \right)^{2/3} \frac{\hbar^2}{2m} \rho_0^{2/3}.$$

Skyrme:

White line:
shown equation
for representative
 ρ_0, α, m^*

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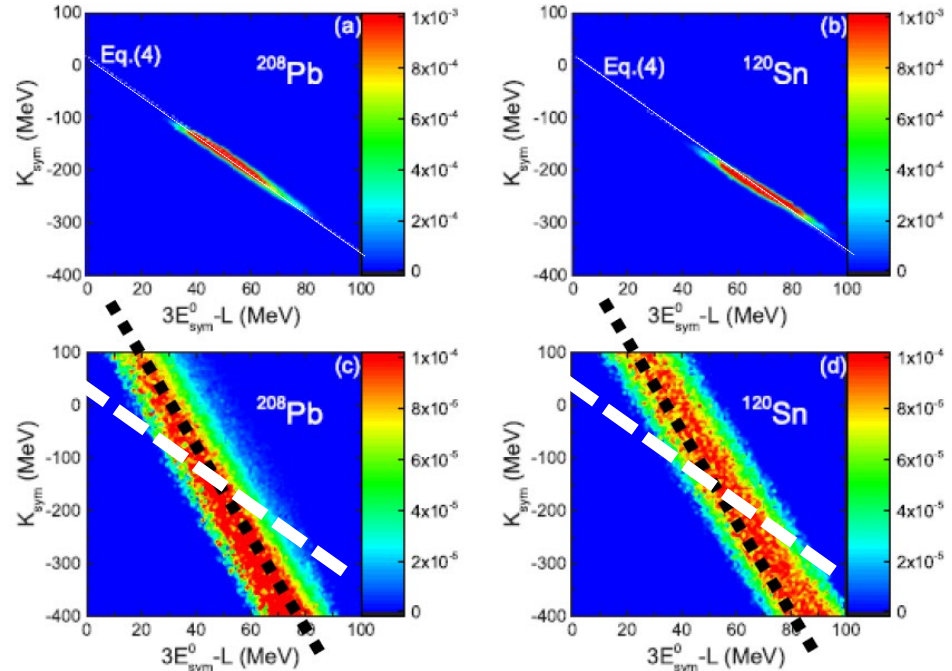
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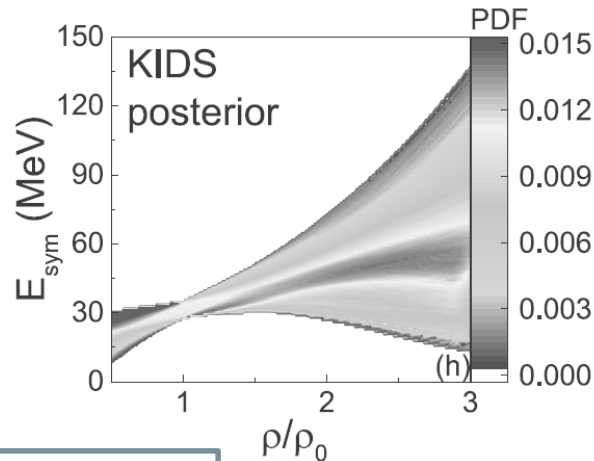
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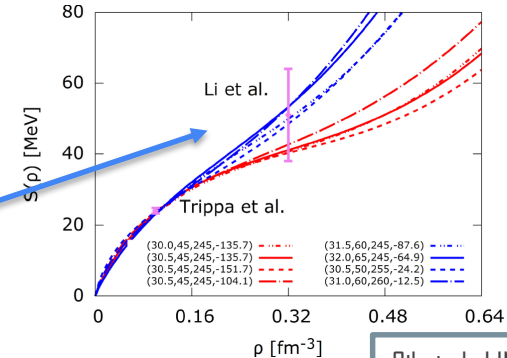
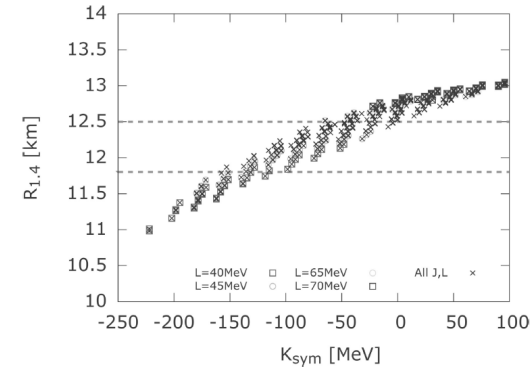
Analyses of astronomical data

$R_{1.4}$ (km)	11.725 ± 1.105 [15]
$R_{2.08}$ (km)	$13.7^{+2.6}_{-1.5}$ [13] and $12.39^{+1.30}_{-0.98}$ [14]
$\Lambda_{1.4}$	190^{+390}_{-120} [18]
M_{max}	$> 2.08 M_{\odot}$ [12]
c_s	< 1



Summary so far:
 $J \approx 30\text{-}33$ MeV,
 $L \approx 45\text{-}65$ MeV,
 $K_{\text{sym}} \approx -200\text{-}0$ MeV
Inflection point

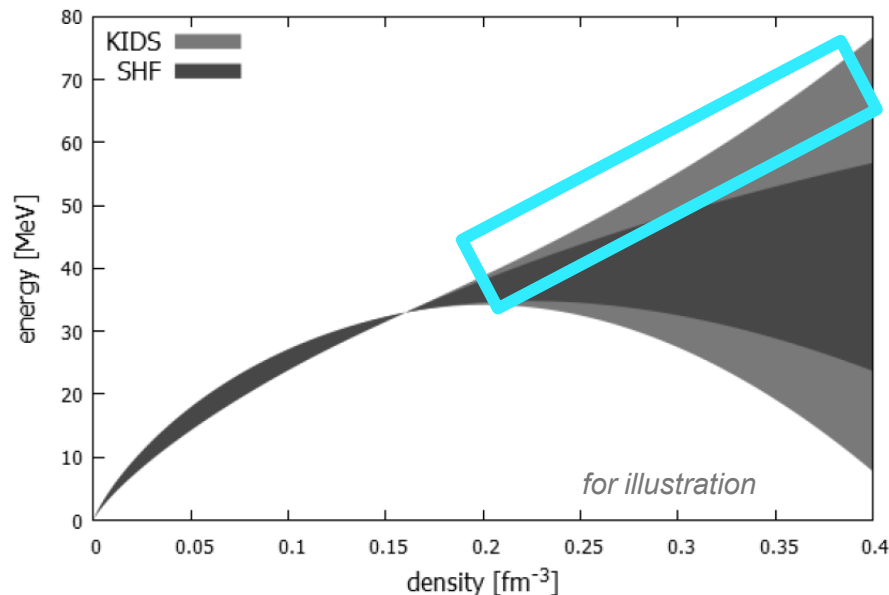
Combined analysis with nuclear data



02 Curvature parameter K_{sym} of the symmetry energy

Density domains

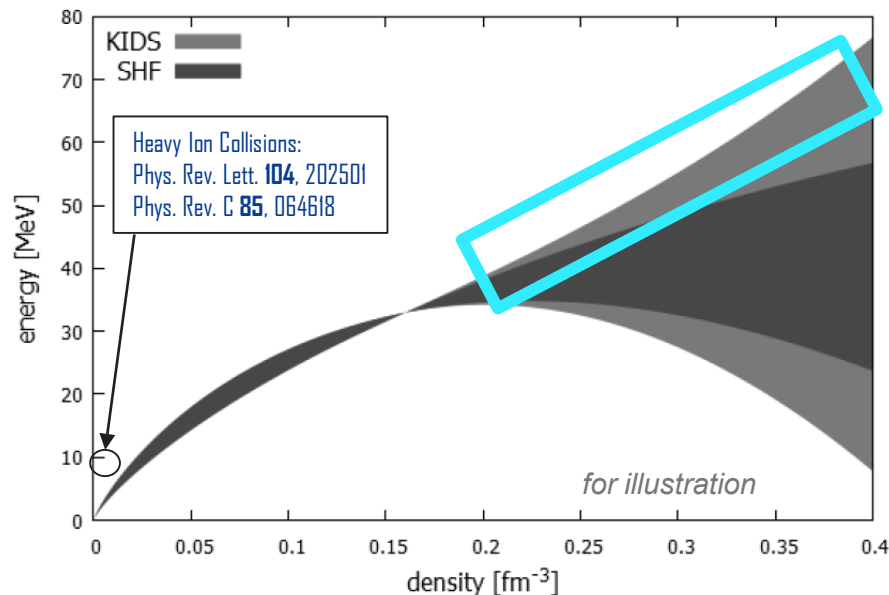
- Most EDF models correspond to a near-parabolic behavior: either stiff or soft
- KIDS framework allows for
 - Inflection point: soft-to-stiff transition, important for description of dense matter
 - Decoupling of saturation and dense regimes



02 Curvature parameter K_{sym} of the symmetry energy

Density domains

- Most EDF models correspond to a near-parabolic behavior: either stiff or soft.
- KIDS framework allows for
 - ❑ Inflection point: soft-to-stiff transition, important for description of dense matter
 - ❑ Decoupling of saturation and dense regimes
 - ❑ **How about dilute matter?**



03 The CREX, PREXII tension



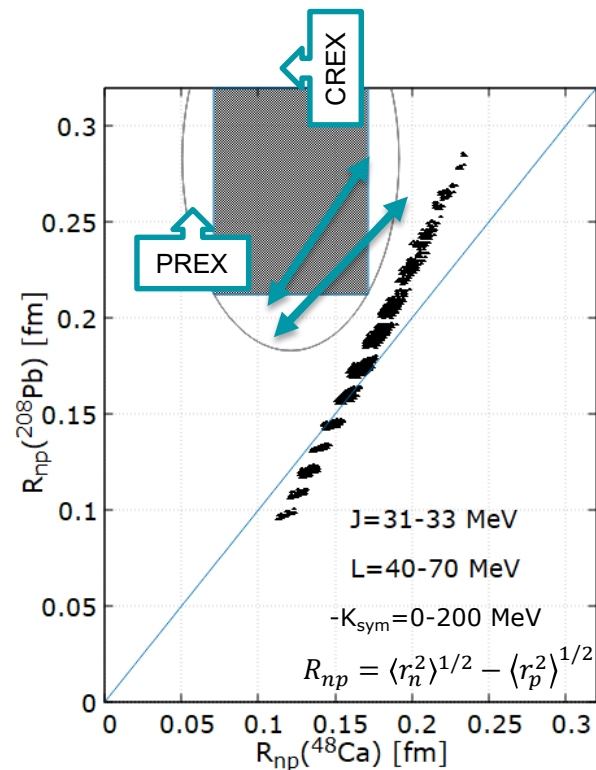
The problem

- ❑ Stiff EDF models which agree with the PREXII result overestimate the neutron skin thickness of ^{48}Ca
- ❑ Soft EDF models which agree with the CREX result underpredict the neutron skin thickness of ^{208}Pb

Current KIDS predictions for the neutron skin

- ❑ The tension is similar to other studies when EoS parameters are varied within realistic ranges.

see also: C.H.Hyun, arXiv:2112.00996



03 The CREX, PREXII tension

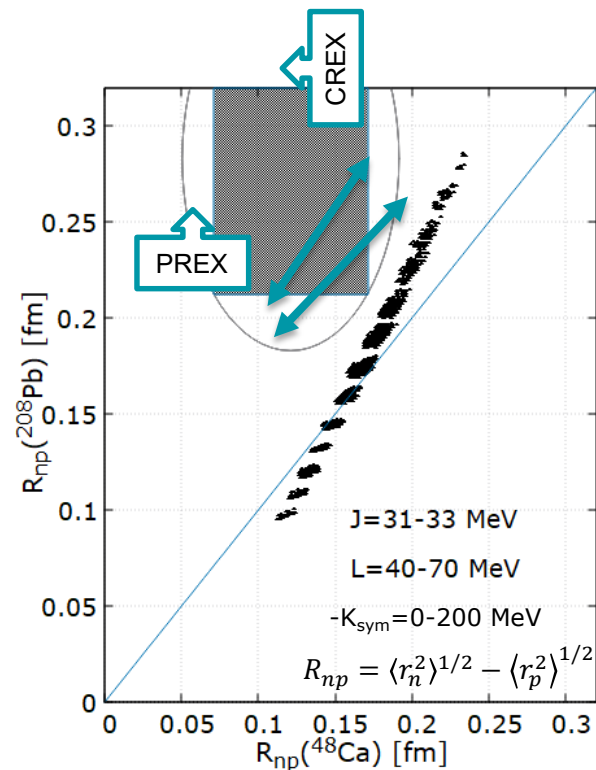


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03 The CREX, PREXII tension



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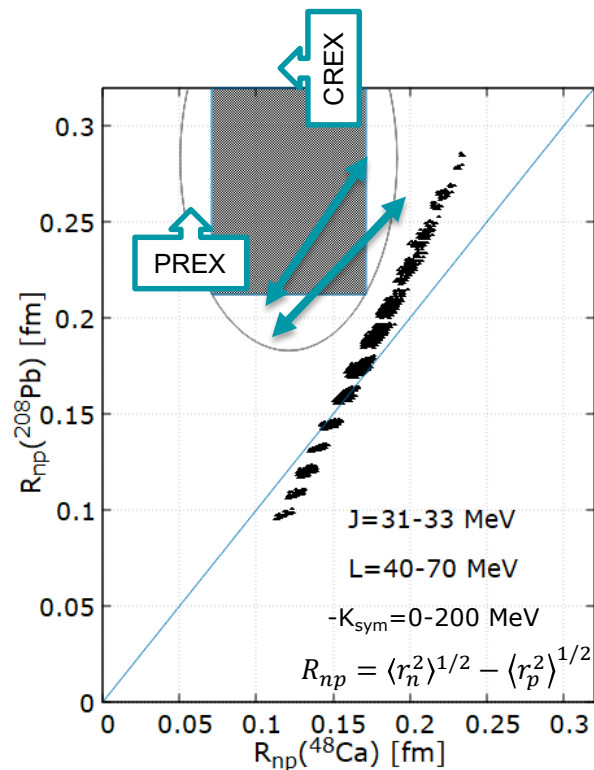
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Current KIDS predictions for the neutron skin

- ❑ The tension is similar to other studies when EoS parameters are varied within realistic ranges.
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What if we explore the very dilute regime only?

see also: PP, arXiv:2210.02696 (Proc.IWNT39)



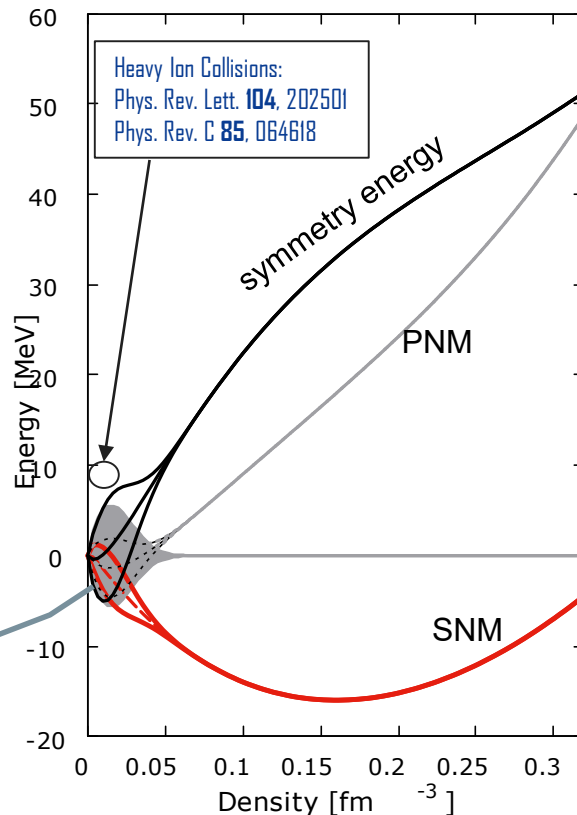
03 The CREX, PREXII tension



Modifications in the low-density regime

- ❑ The KIDS model relies on an expansion of the energy in powers of the density
- ❑ However, dilute matter may require different expansion coefficients
- ❑ Exploring **one additional density-dependent term which decays exponentially and does not affect the saturation and supra-saturation regimes**

New term of the form $c\rho^\alpha \exp(-\beta\rho^2)$



03 The CREX, PREXII tension



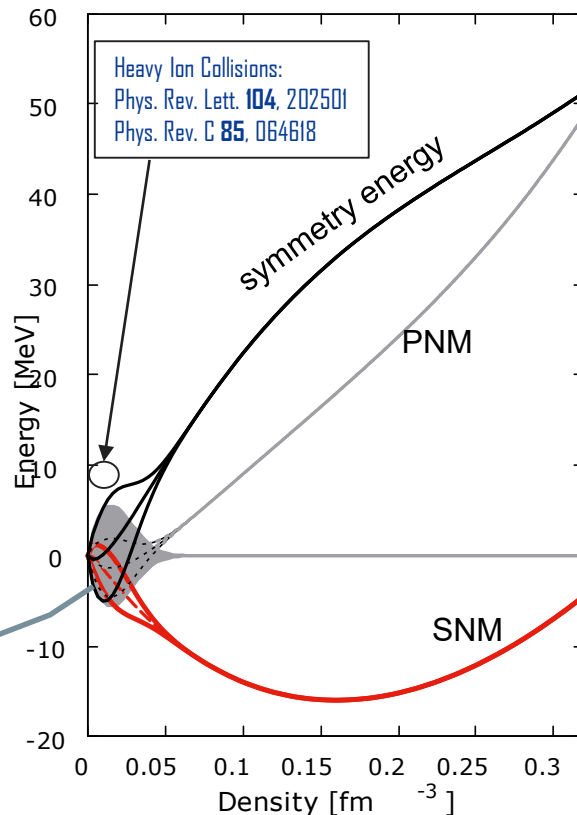
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New term of the form $c\rho^\alpha \exp(-\beta\rho^2)$

New effective interaction term:

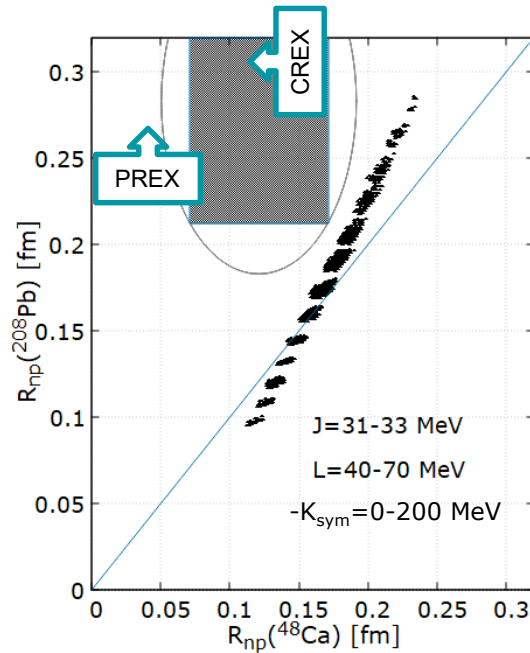
$$\frac{1}{6}(t_{dil} + y_{dil}P_\sigma)\rho^{\alpha-1}\exp(-\beta\rho^2)\delta(\vec{r}_1 - \vec{r}_2)$$



03 The CREX, PREXII tension



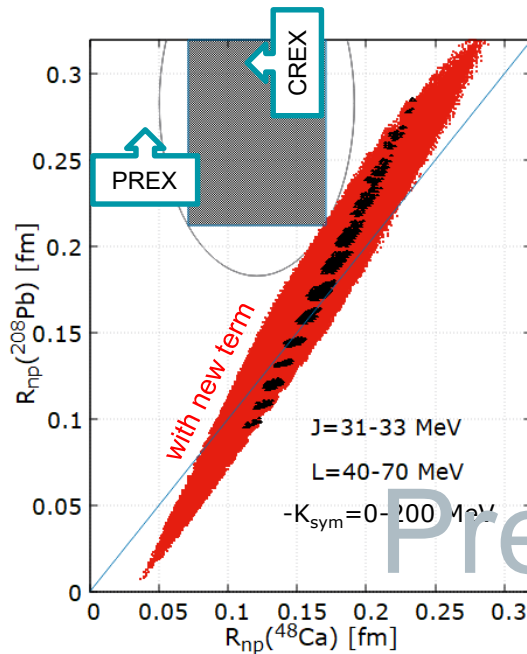
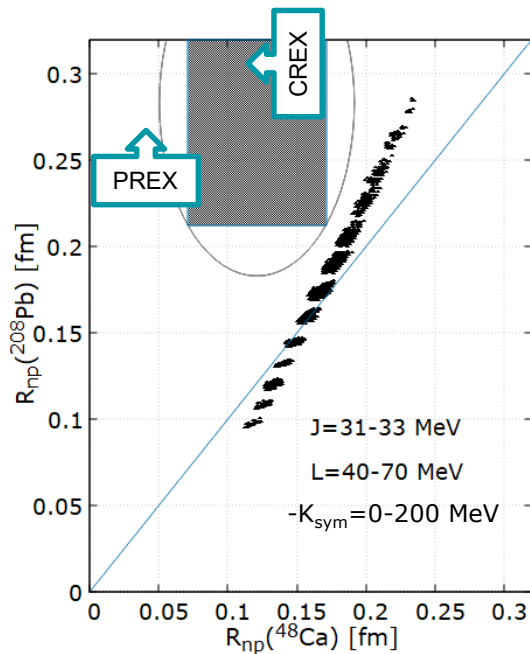
Preliminary results



03 The CREX, PREXII tension



Preliminary results



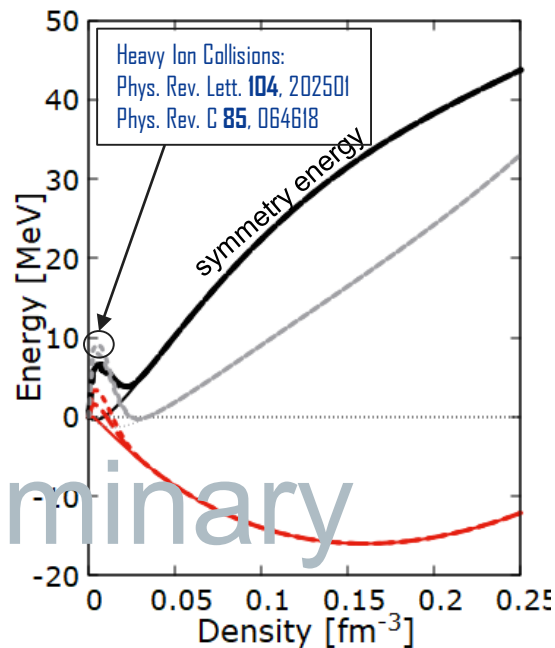
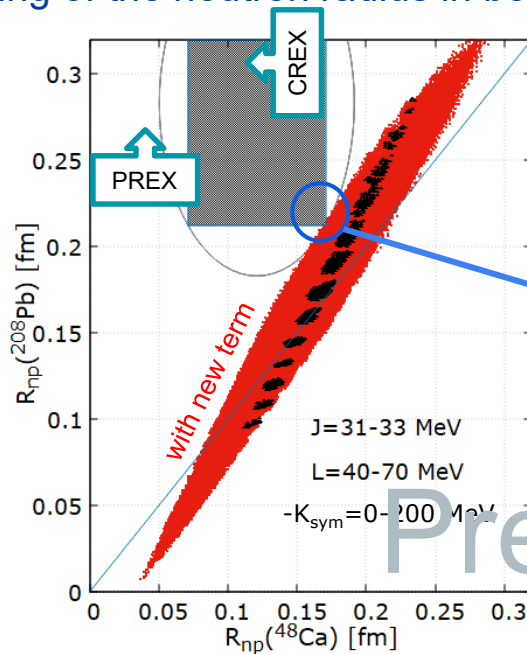
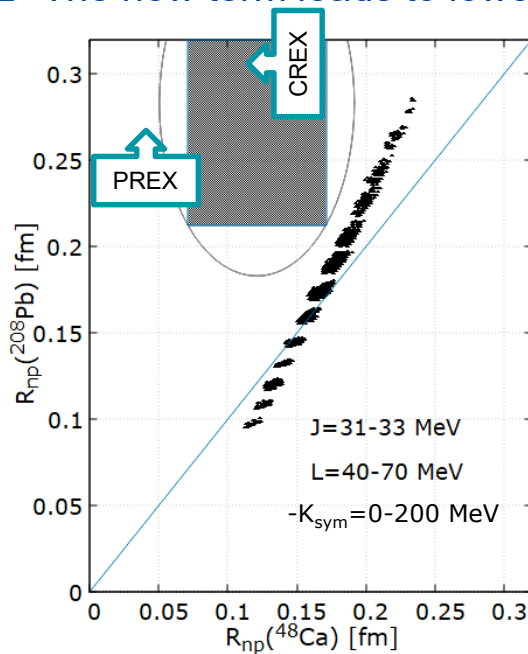
Preliminary

03 The CREX, PREXII tension



Preliminary results

- ❑ The new term leads to lowering of the neutron radius in both nuclei and towards reconciliation



04 Conclusion



- ❑ Most traditional EDF models correspond to a symmetry energy with a near-parabolic dependence on the density, either stiff or soft
- ❑ Exploring the full density dependence with the flexible KIDS framework we find:
 - K_{sym} cannot be constrained from nuclear data alone
 - Correlations between K_{sym} , L , effective mass in Skyrme models are artificial
- ❑ **Combined analyses of nuclear isovector properties and astronomical observations suggest a soft-to-stiff transition somewhere above saturation density**
- ❑ **A reconciliation of CREX, PREXII and dipole polarizability data suggests an even more complex behavior, but consistent with HIC data at low densities**
– to be continued –



Thank you!

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03 The CREX, PREXII tension



What systematic effect of EDF theory could reconcile the two measurements?



A systematic reduction of neutron surface thickness for given realistic bulk EoS

- Effect on neutron skin thickness in the Helm model, assuming **steady proton r.m.s. radius** (charge radius constraint) and **steady hard-sphere neutron radius** (bulk equation of state):

Neutron skin thickness: $R_{np} = \langle r_n^2 \rangle^{1/2} - \langle r_p^2 \rangle^{1/2}$

Helm model r.m.s.: $\langle r_q^2 \rangle^{1/2} = \left[\frac{3}{2} (R_q^2) + \frac{2}{5} (t_q^2) \right]^{1/2}$

Hard-sphere radius $\rightarrow$ Surface thickness

Variation for steady $\langle r_p^2 \rangle^{1/2}$: $\delta R_{np} = \delta \langle r_n^2 \rangle^{1/2} + 0$

Variation for steady R_n^2 : $\delta R_{np} = \frac{1}{2} \sqrt{\frac{3}{2}} \frac{0 + \frac{2}{5} \delta t_n^2}{\langle r_n^2 \rangle^{1/2}}$

$$\delta R_{np} = \frac{1}{5} \sqrt{\frac{3}{2}} \frac{\delta t_n^2}{\langle r_n^2 \rangle^{1/2}}$$

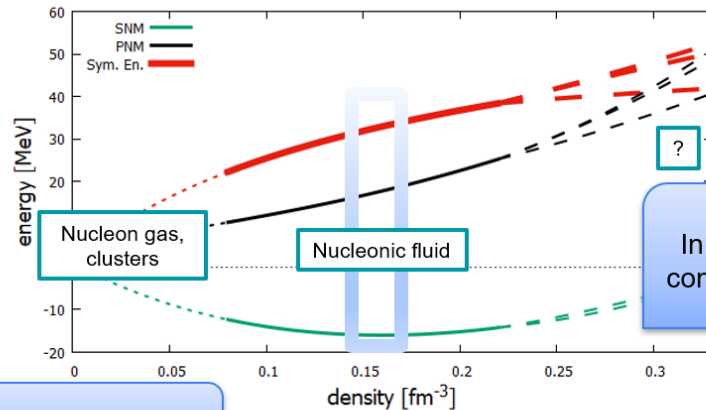
Smaller R_{np} reduction for larger $\langle r_n^2 \rangle$ (^{208}Pb)
Larger R_{np} reduction for smaller $\langle r_n^2 \rangle$ (^{48}Ca)

02 Curvature parameter K_{sym} of the symmetry energy

From INPC2022:

Outlook

Away from saturation regime, different phases may require different modeling



What's new?

In progress: PREX vs CREX?

In progress: interpolations to conformal limit in dense matter

PP & Hyun,
Symmetry 15 (2023) 683

Stay tuned!

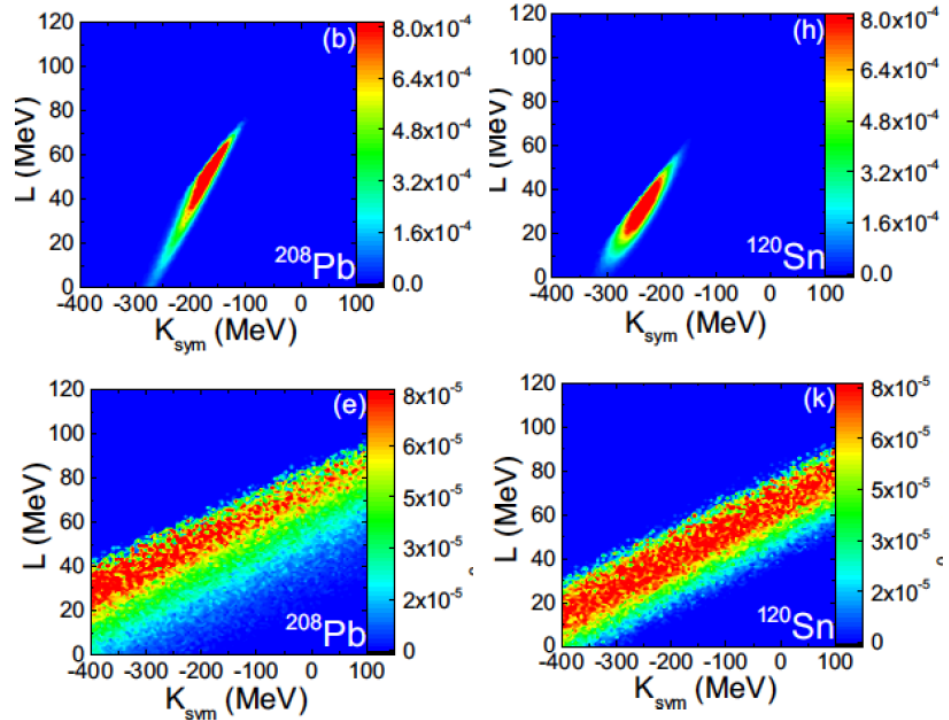
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Bayesian analysis of isovector nuclear properties

K_{sym} vs L

Skyrme:

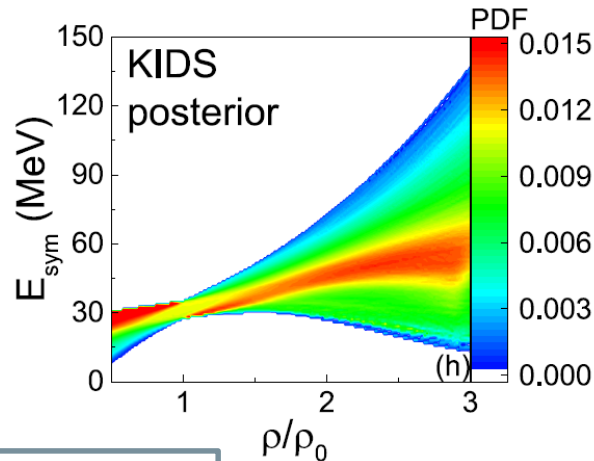
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