



Energy dependence of transverse momentum fluctuations in Au+Au collisions from a multiphase transport model

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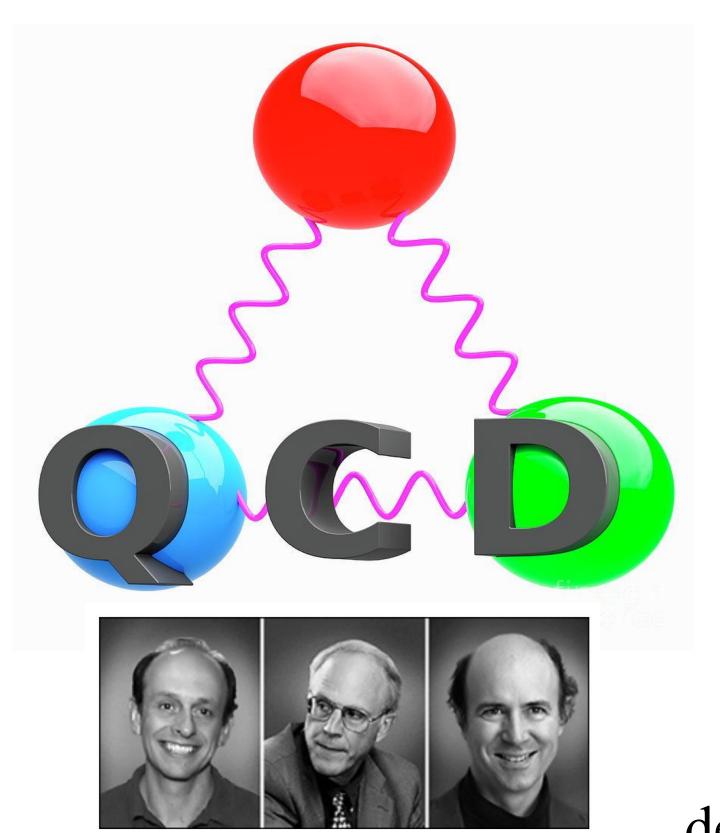
The 29th International Nuclear Physics Conference (INPC2025), Daejeon, Korea

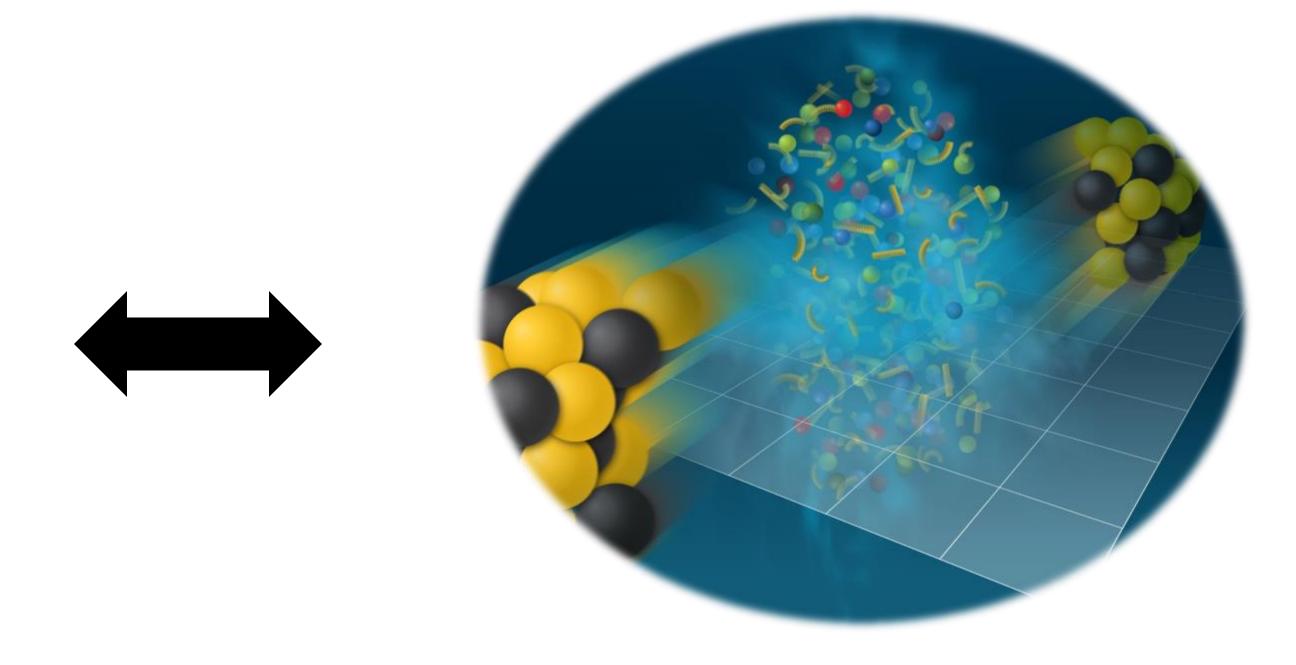




Quantum ChromoDynamics (QCD):

the theory of strong interactions governing the behavior of quarks and gluons.





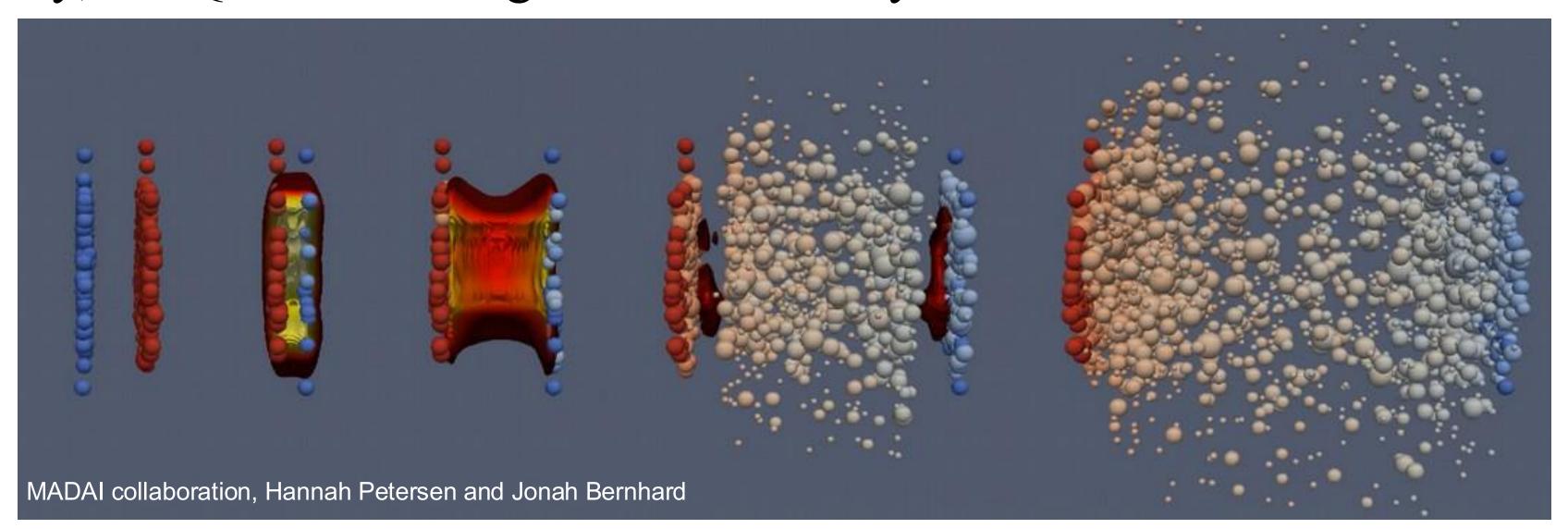
Quark Gluon Plasma (QGP):

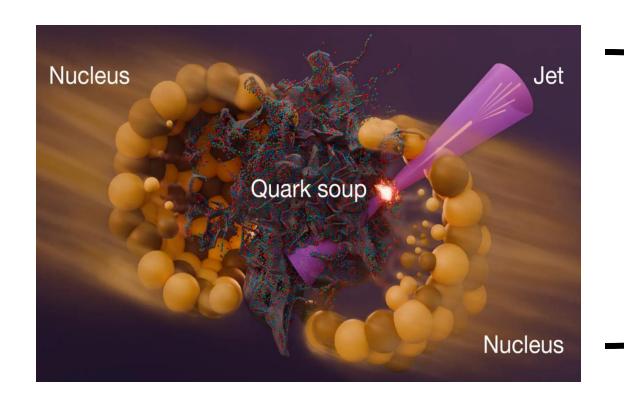
deconfined state under extreme conditions of temperature and density.





In the laboratory, the QGP could be generated in heavy-ion collisions.





Hard Probes:

Jets,

Large p_T hadrons,

Heavy quark/hadrons,

Quarkonia,

Soft Probes:

Yields/Spectra,

Flows,

Correlations,

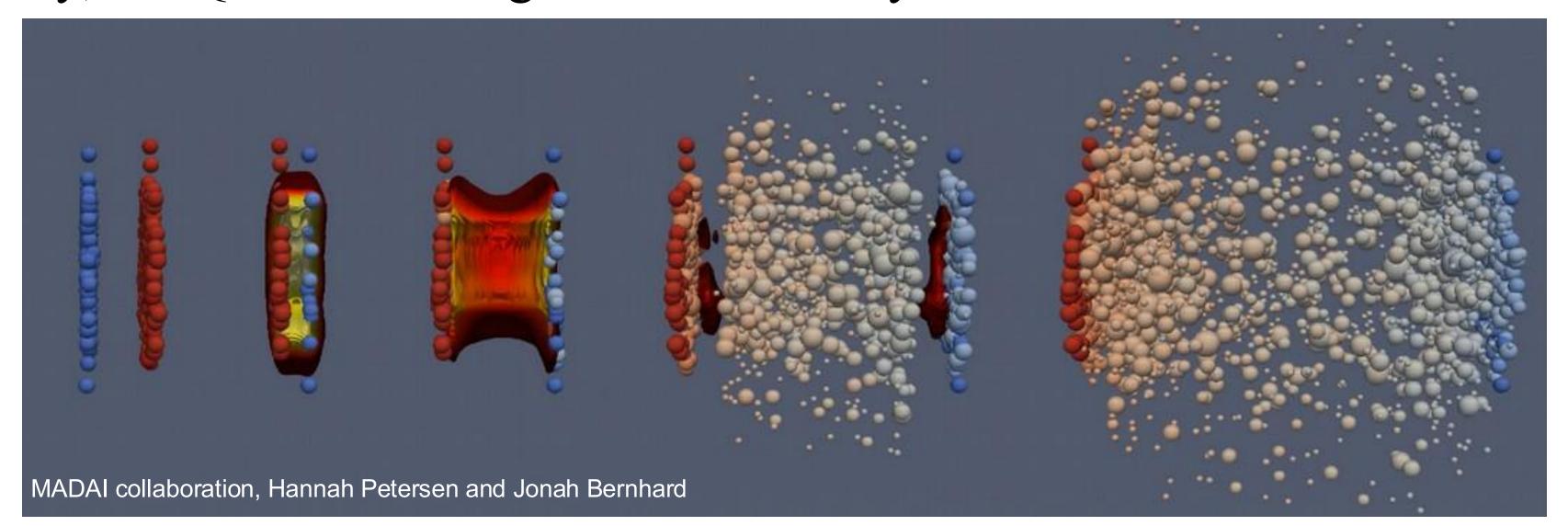
Decorrelations,

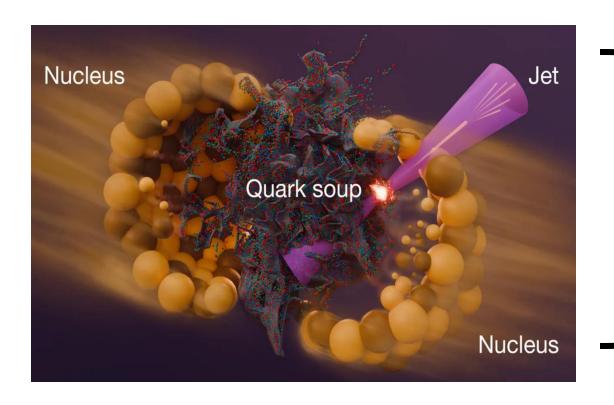
Fluctuations,





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Hard Probes:

Jets,

Large p_T hadrons,

Heavy quark/hadrons,

Quarkonia,

• • •

Soft Probes:
Yields/Spectra,
Flows,
Correlations,
Decorrelations,
Fluctuations,

Nu Xu: Tue (16:30)

Conserved Quantities: net charge/proton...

Event-by-Event $\langle p_{\rm T} \rangle$.

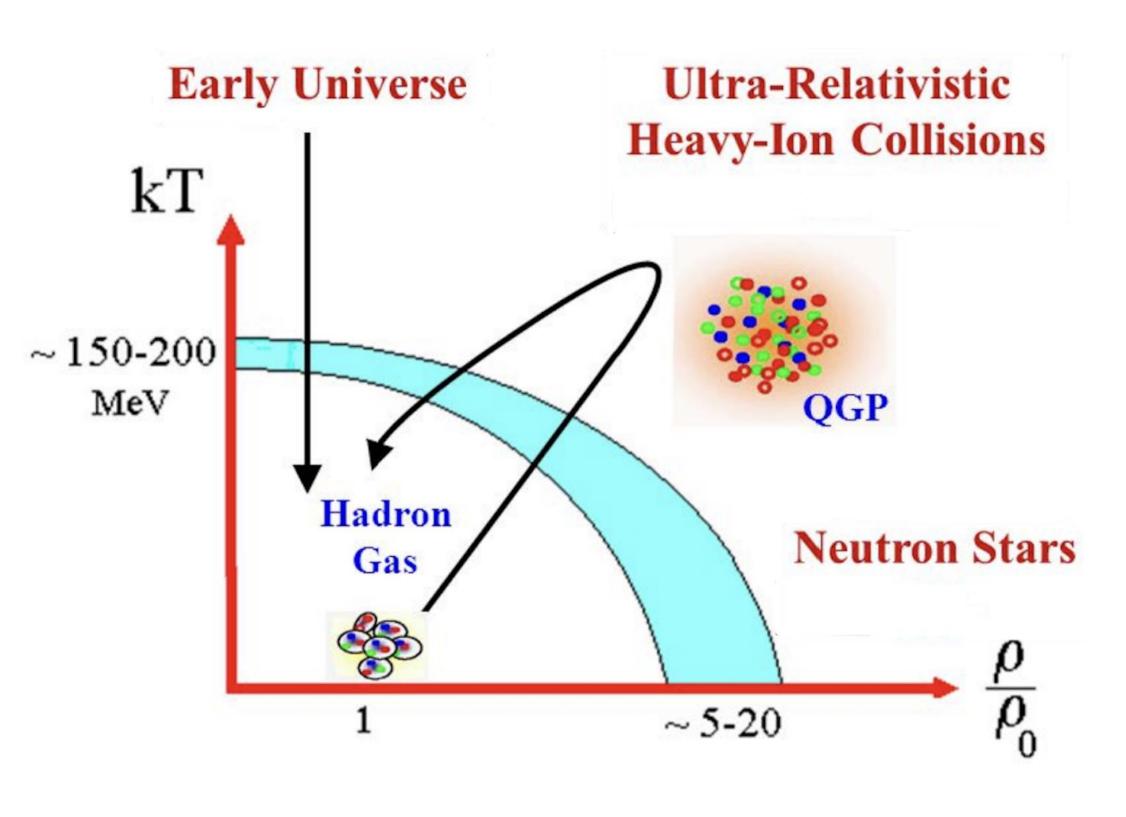


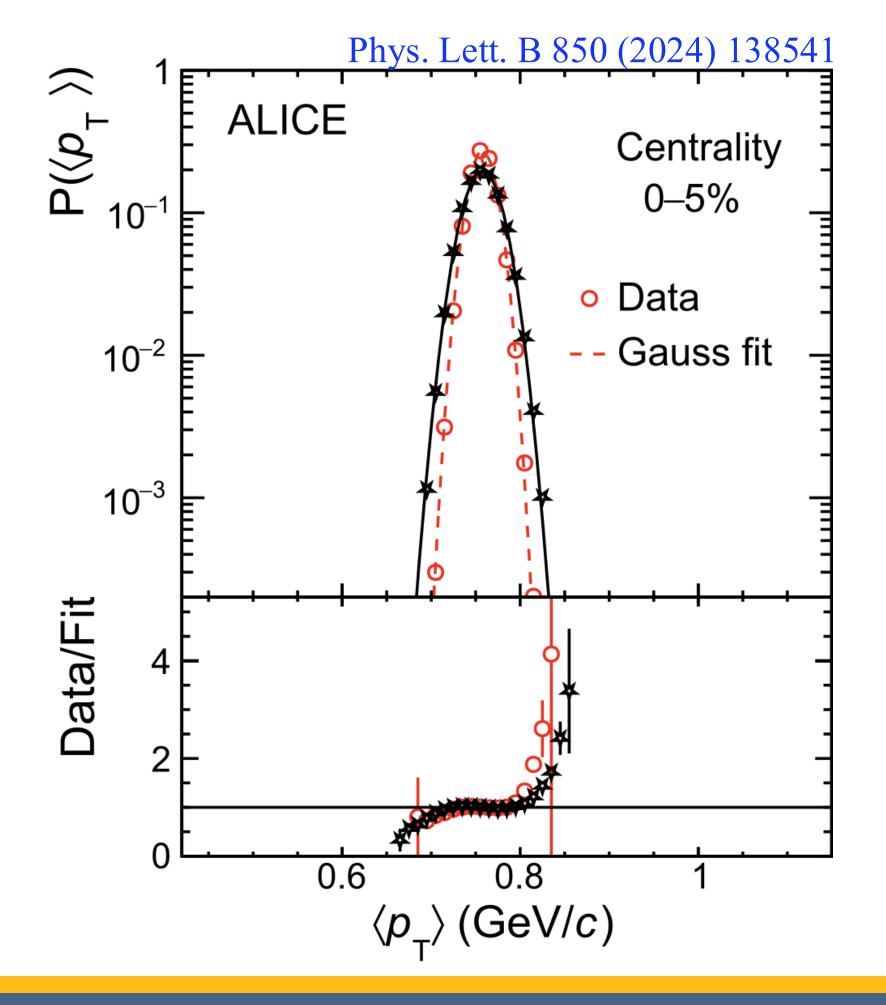


Phase transition (QGP to Hadron Gas) —

Thermodynamics quantity (T): critical fluctuation.

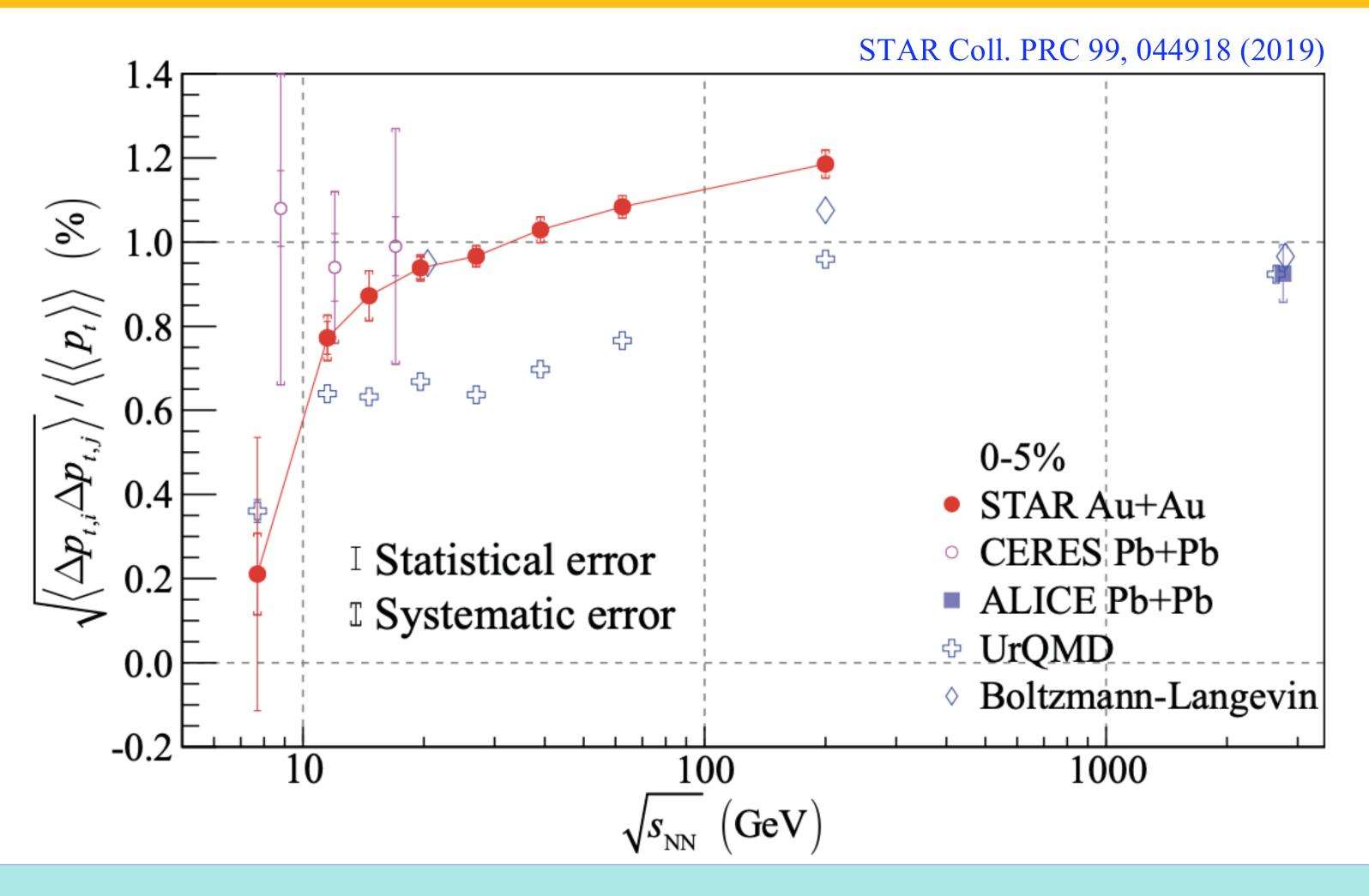
 \rightarrow dynamical $\langle p_{\rm T} \rangle$ fluctuations











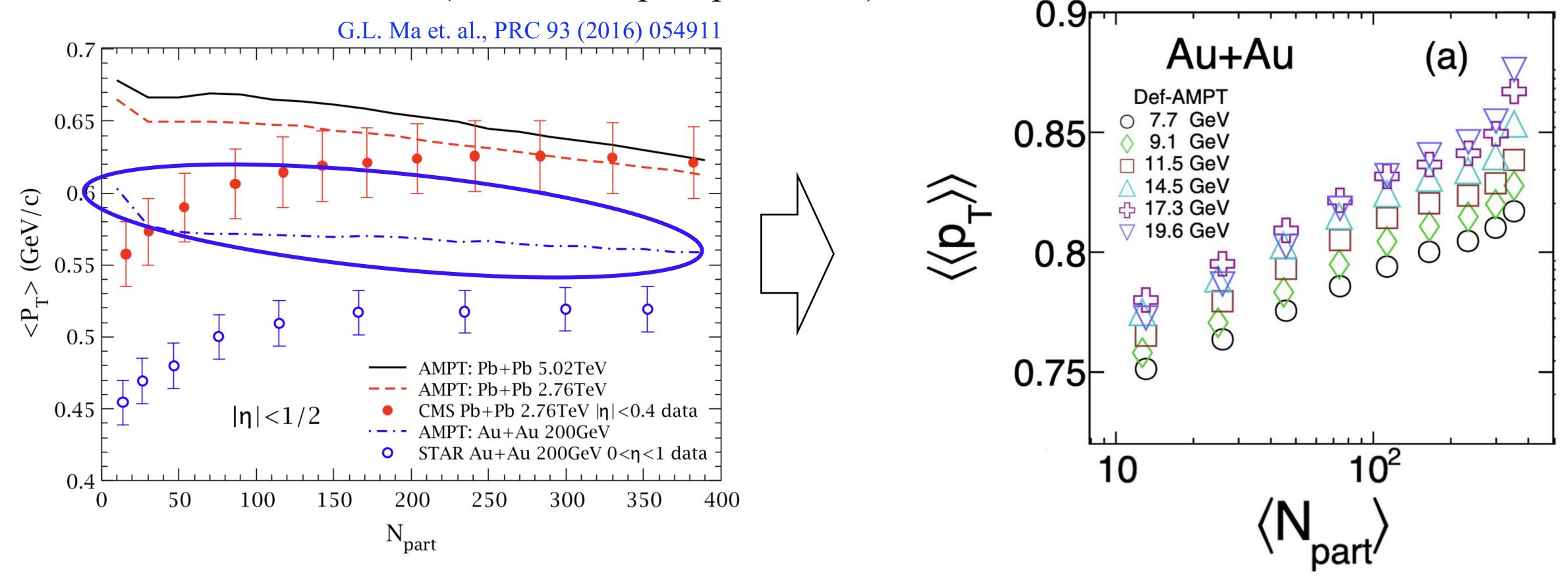
A non-monotonic behavior of $\langle p_{\rm T} \rangle$ fluctuations as a function of centrality or incident energy was suggested as one of the possible signals of the QGP.





Two key parameters (a_L, b_L) in hadronization mechanisms are tuned:

fixed value → variable (based on impact parameter)

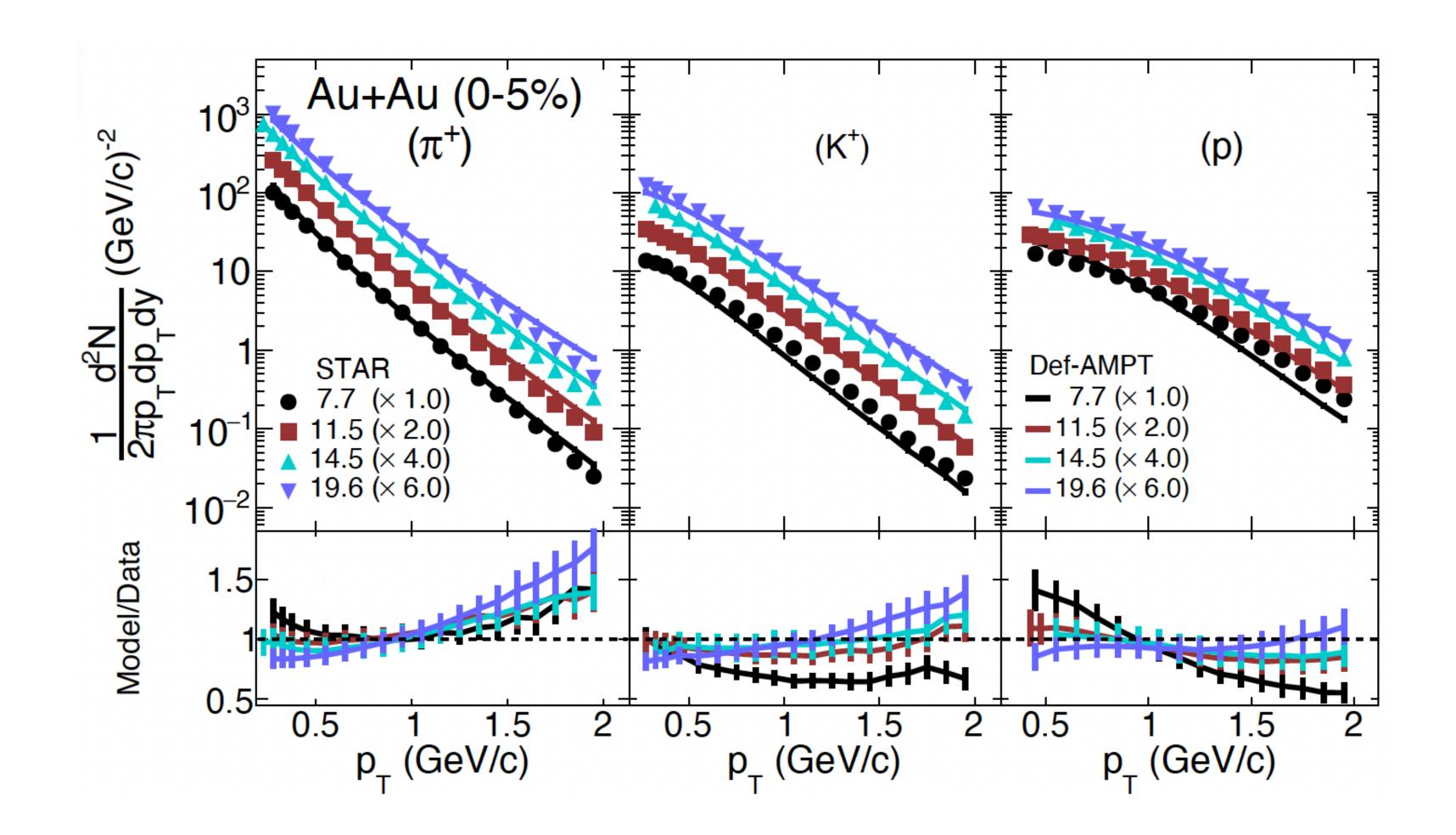


 $\langle \langle p_{\rm T} \rangle \rangle$ tends to increase with centrality, which are more reasonable expectation.

The AMPT validity







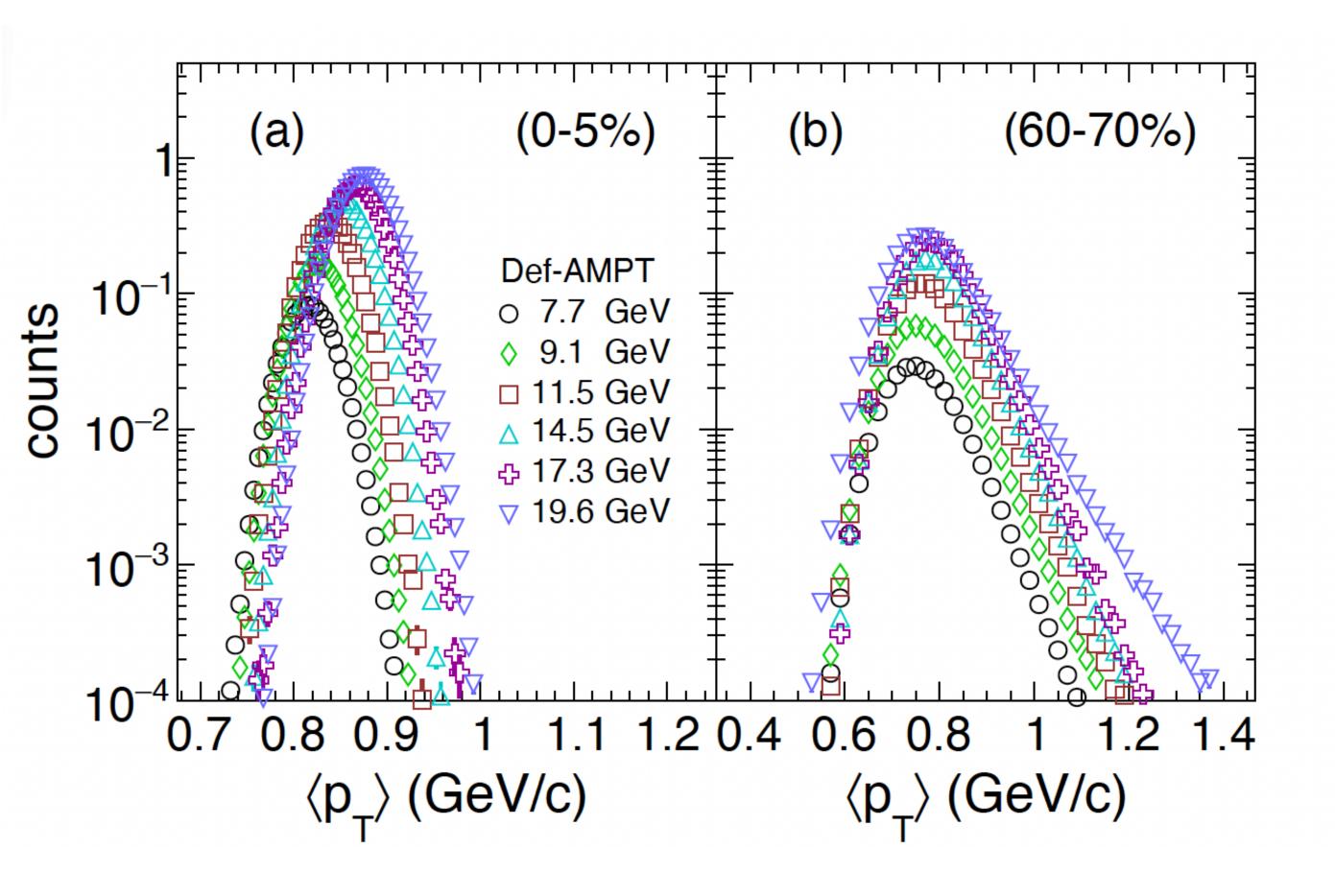
The improved AMPT qualitatively produces the experimental data across a broad $p_{\rm T}$ spectrum.

Event-wise $\langle p_{\rm T} \rangle$ distribution





The $\langle p_{\rm T} \rangle$ fluctuations can be straightforwardly studied by its event-wise distributions.



At peripheral collisions:

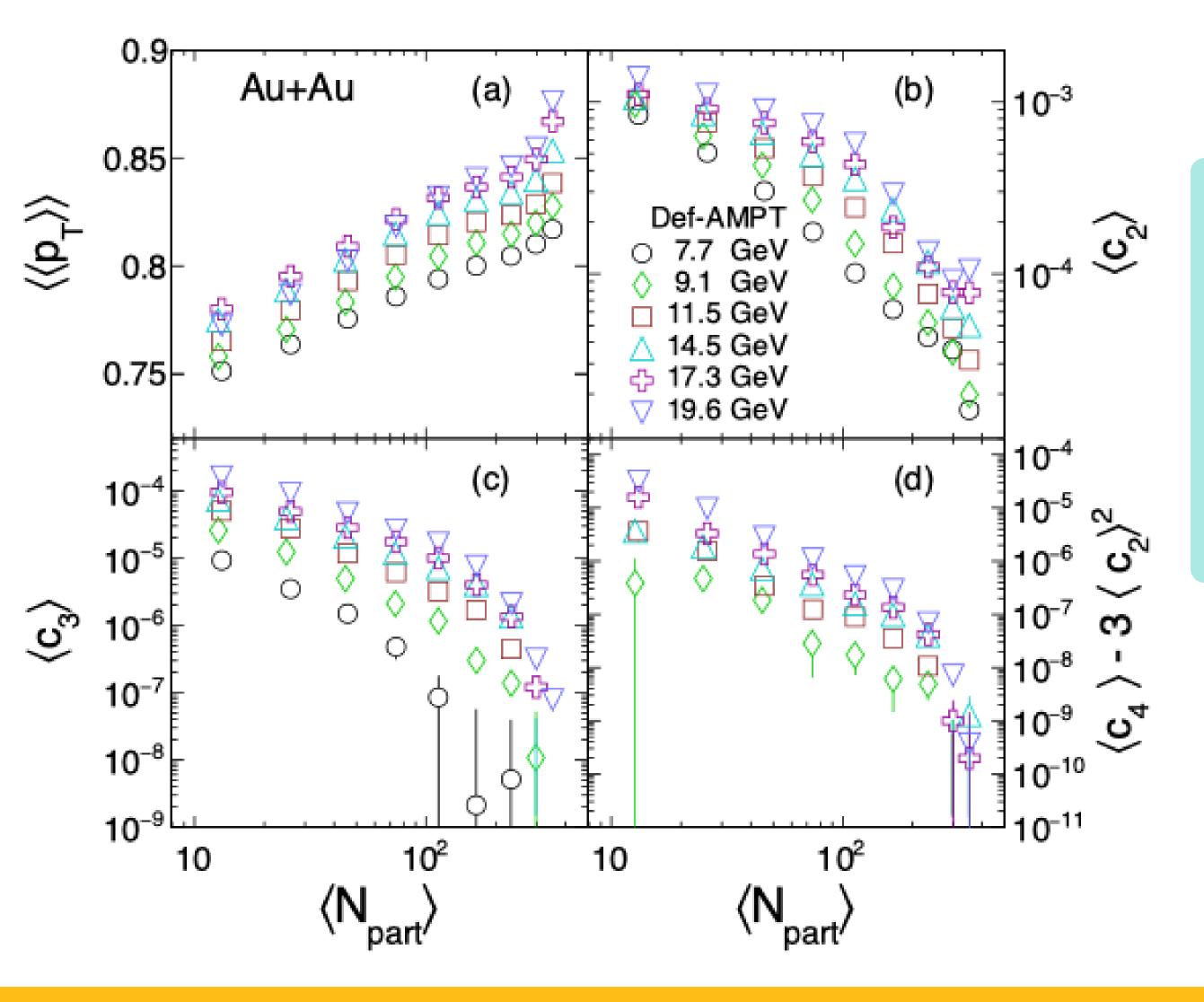
- > greater variances: indicating enhanced fluctuations.
- > more significant rightward tail: suggesting positive skewness.

Higher-order p_T cumulants





Variance($\langle c_2 \rangle$), Skewness ($\langle c_3 \rangle$), Kurtosis($\langle c_4 \rangle$ -3 $\langle c_2 \rangle^2$)



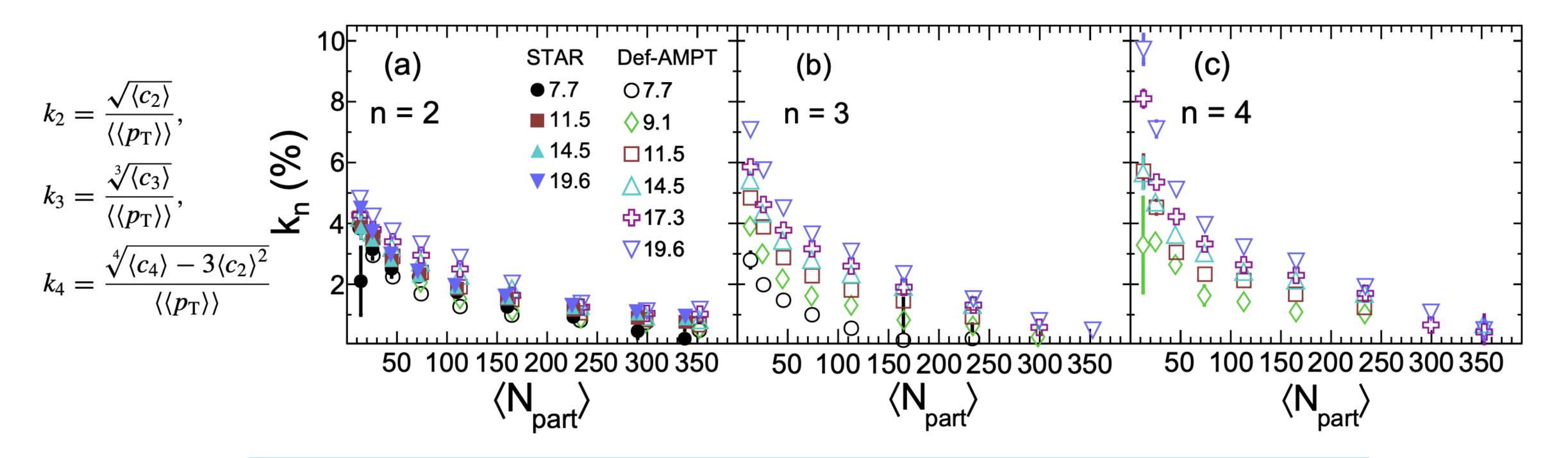
- > Share significant dependence on: a) centrality, b) incident energy;
- > An inverse dependence on centrality: a reduction in particle-pair correlations if they are dominated by particles originating from the same nucleon-nucleon collisions.

Scaled p_T cumulants





To mitigate the influence on the cumulants in $\langle \langle p_{\rm T} \rangle \rangle$ with centrality or incident energy.

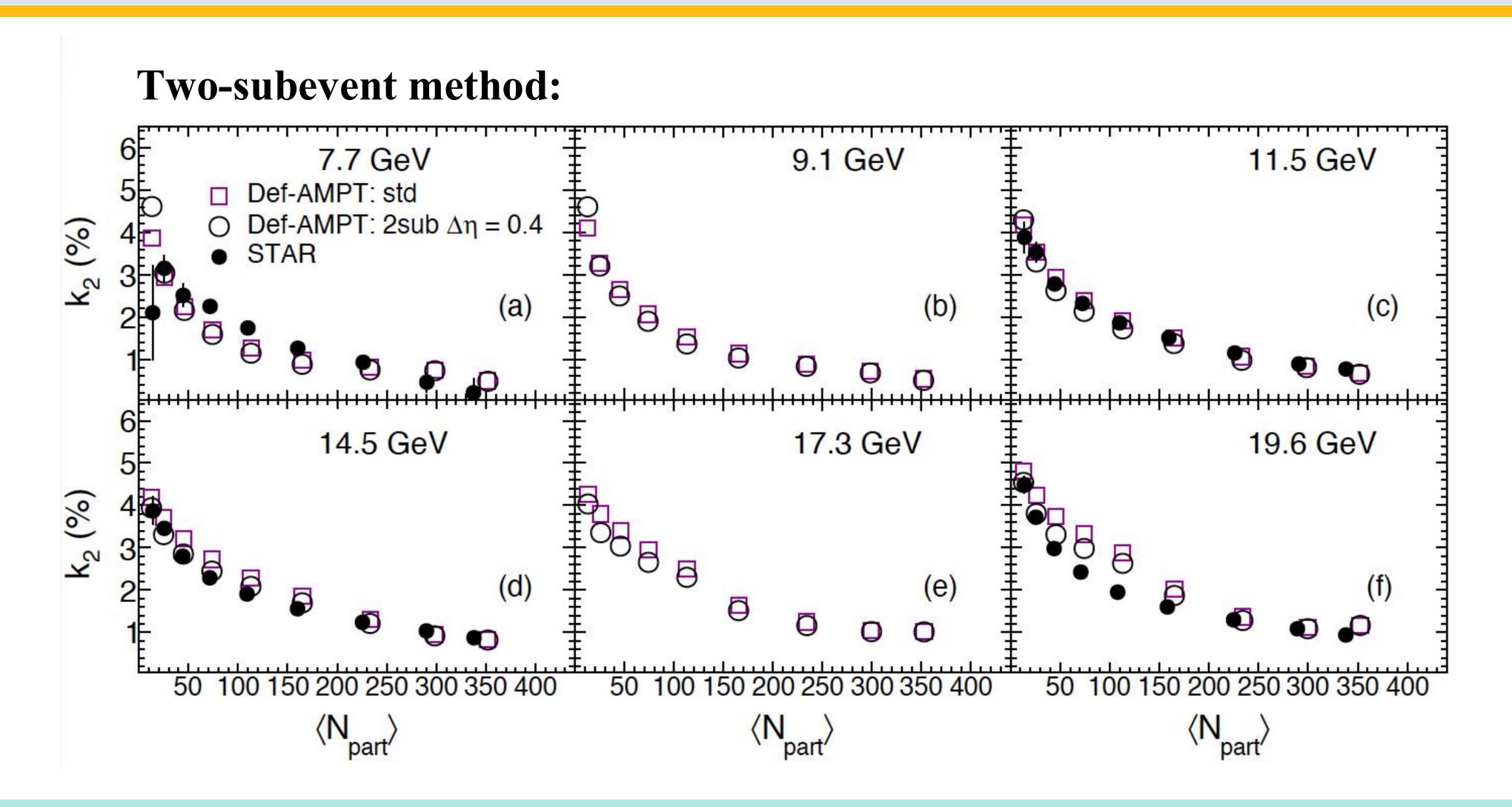


- $\geq k_2$ qualitatively produce the trends of STAR experiment.
- $> k_n$ (n=2,3,4) exhibits significant centrality and energy dependence;
- $> k_n$ (n=2,3,4) exhibit an approximate power-law behavior.

Scaled p_T cumulants: the short-range correlations





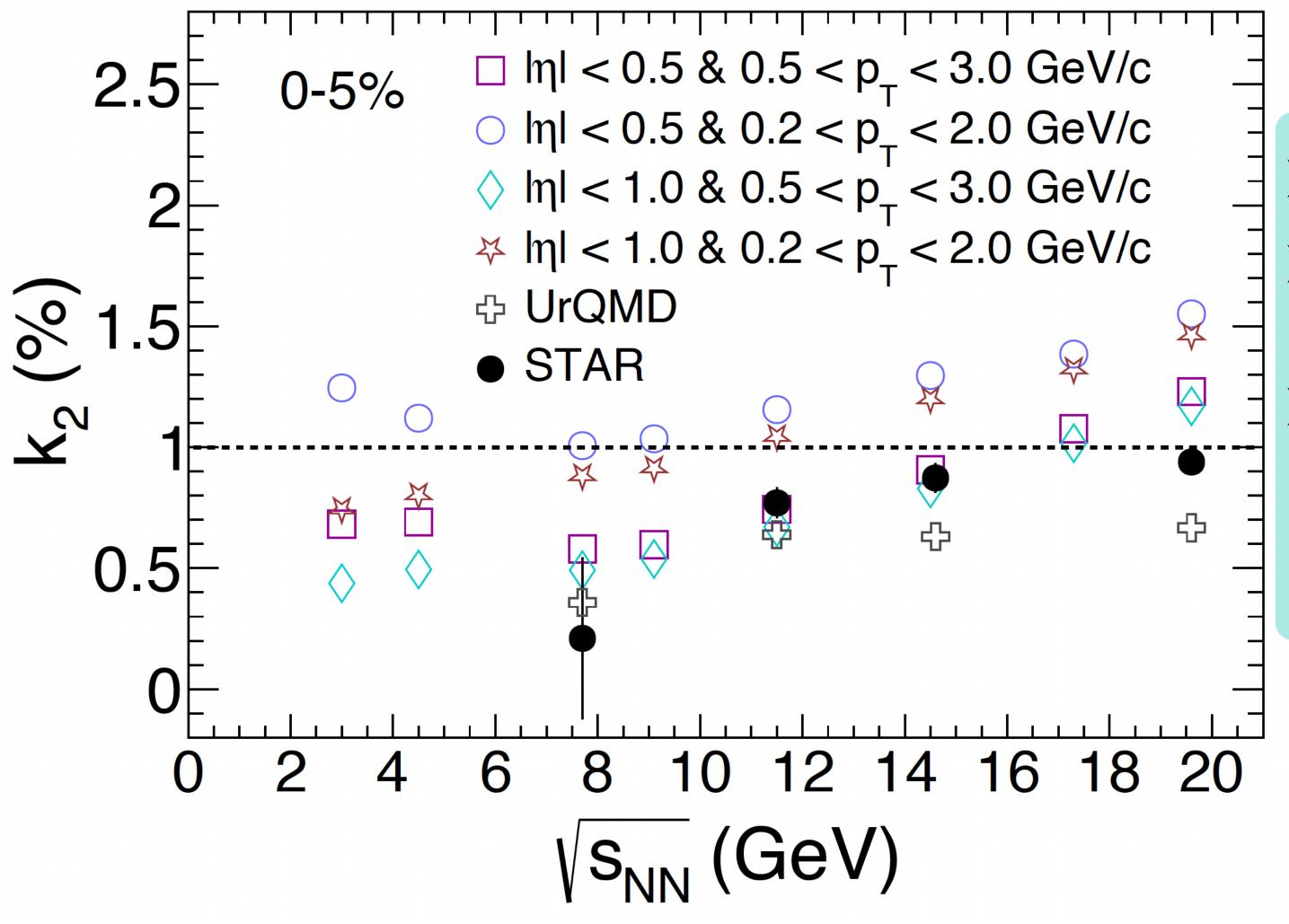


All values from two-subevent method are slightly suppressed compared to the standard method.

Scaled p_T cumulants: vs energy







- \triangleright More sensitive to p_T variations than η .
- $> 0.5 < p_T < 3.0 \text{ GeV/c: quantitatively}$ consistent with STAR measurements.
- > extend to 3.0 GeV/c: a significant abnormal increase, indicating an enhancement of dynamical correlations at lower collision energies.

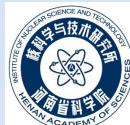
Summary and outlook





- \triangleright A systematic study of higher-order dynamical p_T cumulants up to fourth order.
- \triangleright Higher-order $p_{\rm T}$ cumulants w/o normalization exhibit a strong dependence on centrality.
- > Our finding provides variable references for the experimental measurements.
- > The inner mechanism from acceptance or decorrelation need further to be explored.

Thank you for your attention !!!





Backup

Methodology: n-particle p_T correlator





$$c_n = \frac{\sum_{i_1 \neq \cdots \neq i_n} w_{i_1} \cdots w_{i_n} \left(p_{\mathrm{T}, \mathrm{i}_1} - \langle \langle p_{\mathrm{T}} \rangle \rangle \right) \cdots \left(p_{\mathrm{T}, \mathrm{i}_n} - \langle \langle p_{\mathrm{T}} \rangle \rangle \right)}{\sum_{i_1 \neq \cdots \neq i_n} w_{i_1} \cdots w_{i_n}}$$

$$p_{mk} = \sum_{i} w_{i}^{k} p_{i}^{m} / \sum_{i} w_{i}^{k}, \quad \tau_{k} = \frac{\sum_{i} w_{i}^{k+1}}{\left(\sum_{i} w_{i}\right)^{k+1}}$$

$$c_{2} = \frac{\bar{p}_{11}^{2} - \tau_{1}\bar{p}_{22}}{1 - \tau_{1}},$$

$$c_{3} = \frac{\bar{p}_{11}^{3} - 3\tau_{1}\bar{p}_{22}\bar{p}_{11} + 2\tau_{2}\bar{p}_{33}}{1 - 3\tau_{1} + 2\tau_{2}},$$

$$c_{4} = \frac{\bar{p}_{11}^{4} - 6\tau_{1}\bar{p}_{22}\bar{p}_{11}^{2} + 3\tau_{1}^{2}\bar{p}_{22}^{2} + 8\tau_{2}\bar{p}_{33}\bar{p}_{11} - 6\tau_{3}\bar{p}_{44}}{1 - 6\tau_{1} + 3\tau_{1}^{2} + 8\tau_{2} - 6\tau_{3}}$$

S. Bhatta, C, Zhang, and J. Jia, PRC 105, 024904 (2022)





1) vs centrality, 2) vs incident energy

Two key parameters in hadronization mechanisms in the AMPT model:

 a_L, b_L : which are inversely related to $\langle p_T \rangle^2$ of hadrons.

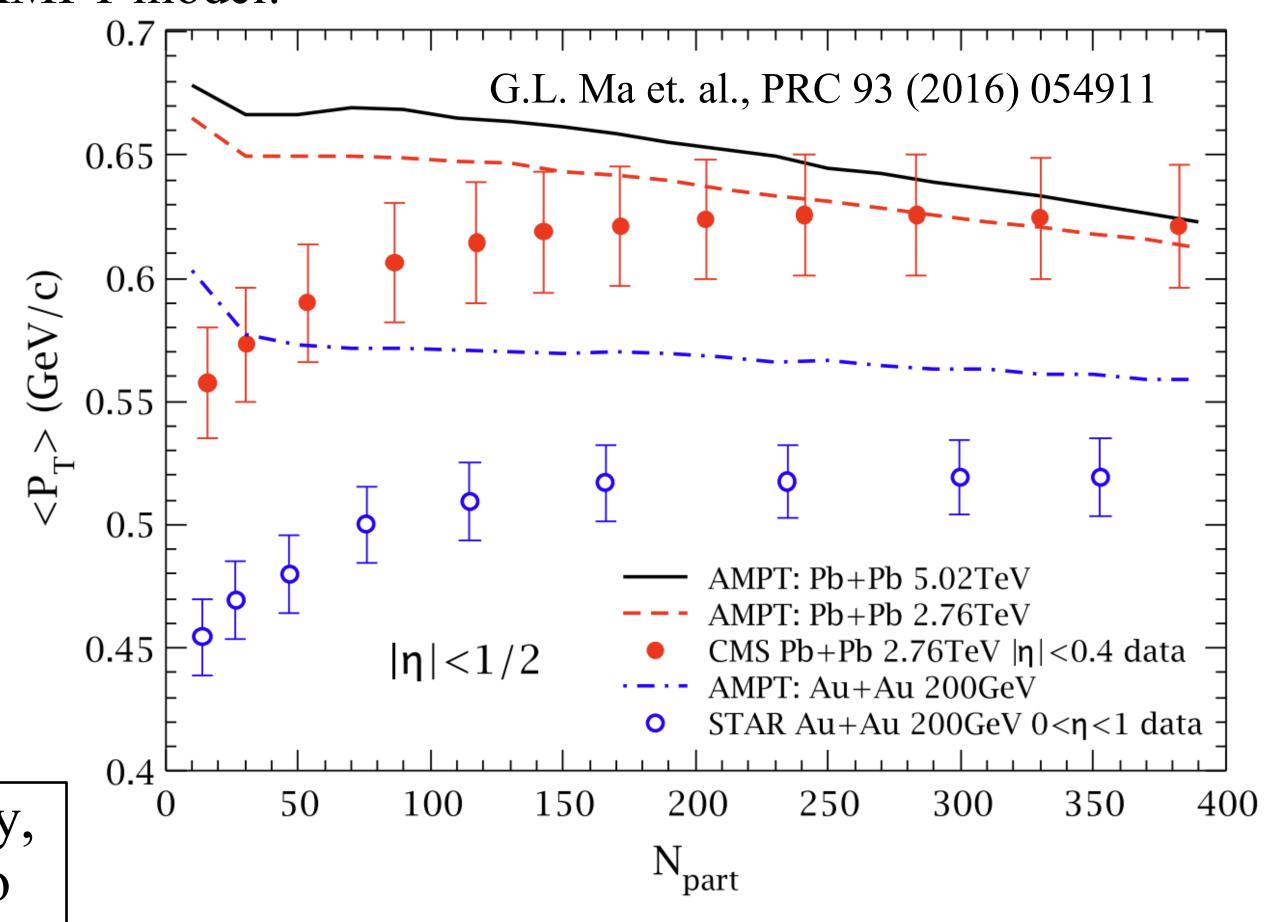
$$\langle p_{\rm T}^2 \rangle = \frac{1}{b_L (2 + a_L)}$$

Traditionally, a_L , b_L are constant values, e.g.

RHIC@200 GeV: $a_L = 0.55$, $b_L = 0.15$

LHC@2.75/5.02 TeV: $a_L = 0.3$, $b_L = 0.15$

The $\langle p_{\rm T} \rangle$ (cent.) is inconsistent with the data. Actually, the $\langle p_{\rm T} \rangle$ is expected to increase with centrality due to higher initial temperature in more central collisions.

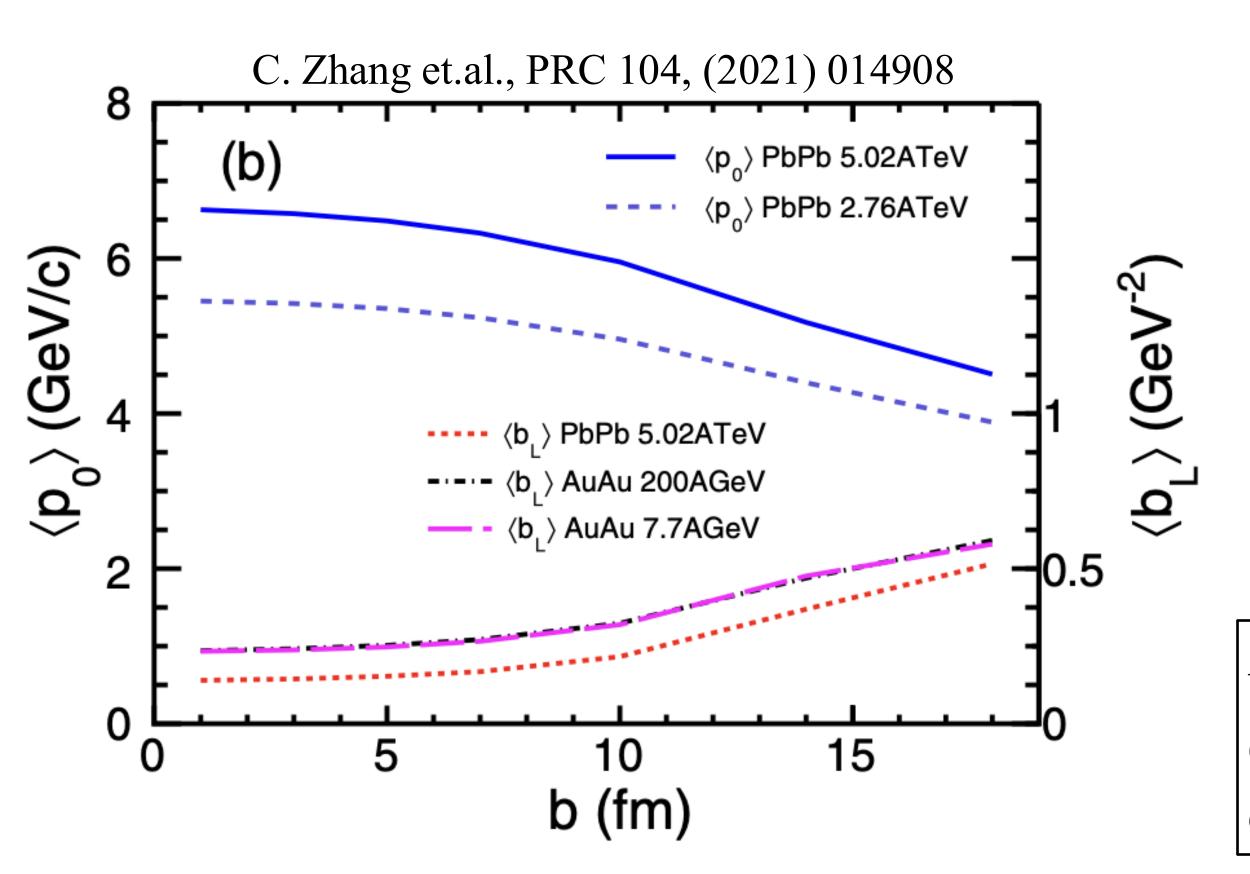






A possible solution:

Make b_L a local variable, which has a dependence on the transverse position of the corresponding excited string in each event.



At 200 GeV and 7.7GeV:

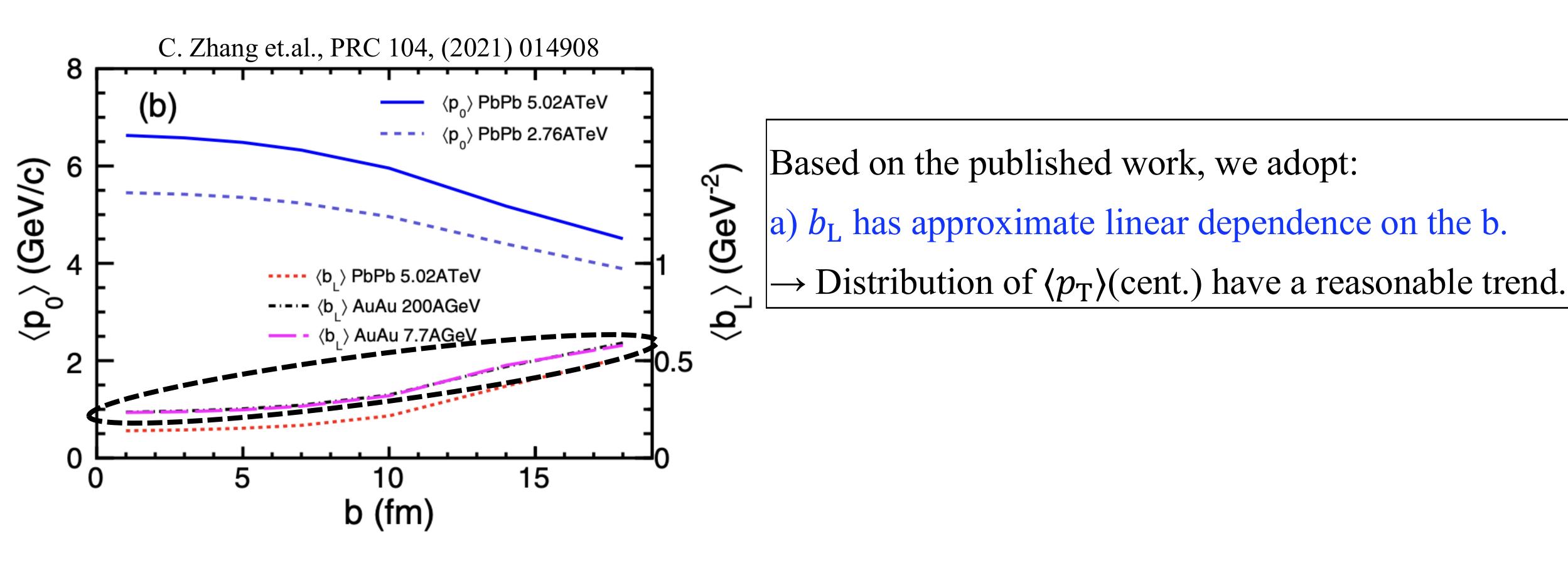
C. Zhang's work reveals b_L has approximate linear dependence on the impact parameter.





A possible solution:

Make b_L a local variable, which has a dependence on the transverse position of the corresponding excited string in each event.

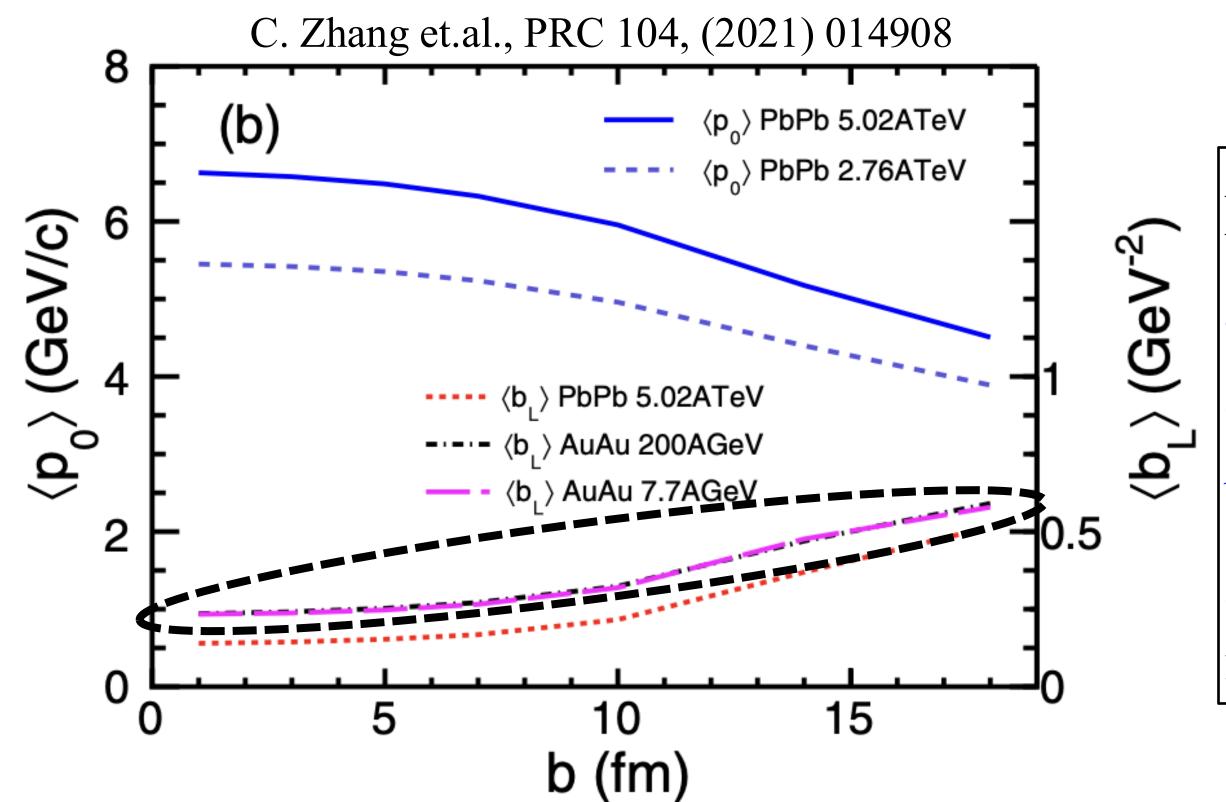






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Make b_L a local variable, which has a dependence on the transverse position of the corresponding excited string in each event.



Based on the published work, we adopt:

- a) $b_{\rm L}$ has approximate linear dependence on the b.
- \rightarrow Distribution of $\langle p_{\rm T} \rangle$ (cent.) have a reasonable trend.
- b) b_L is systematically increased with E.
- \rightarrow Distribution of $\langle p_{\rm T} \rangle$ (cent.) is higher at higher incident energy due to enhanced collective effects.

PRC **72**, 044902 (2005)

100

150

20 GeV Au+Au 62 GeV Au+Au 130 GeV Au+Au 200 GeV Au+Au

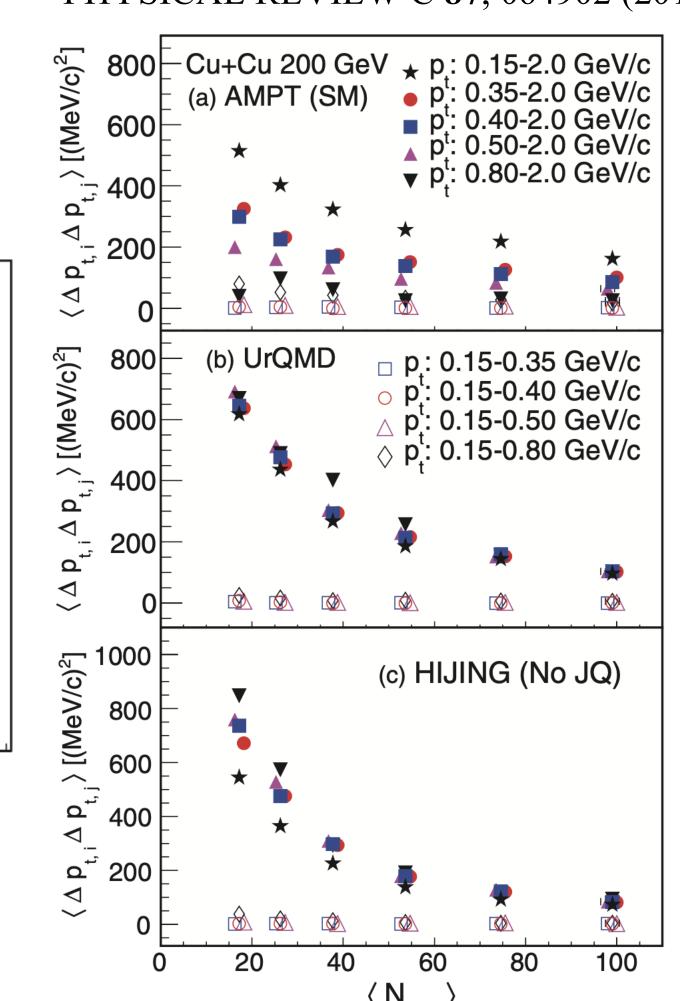
130 GeV HIJING

200 GeV HIJING

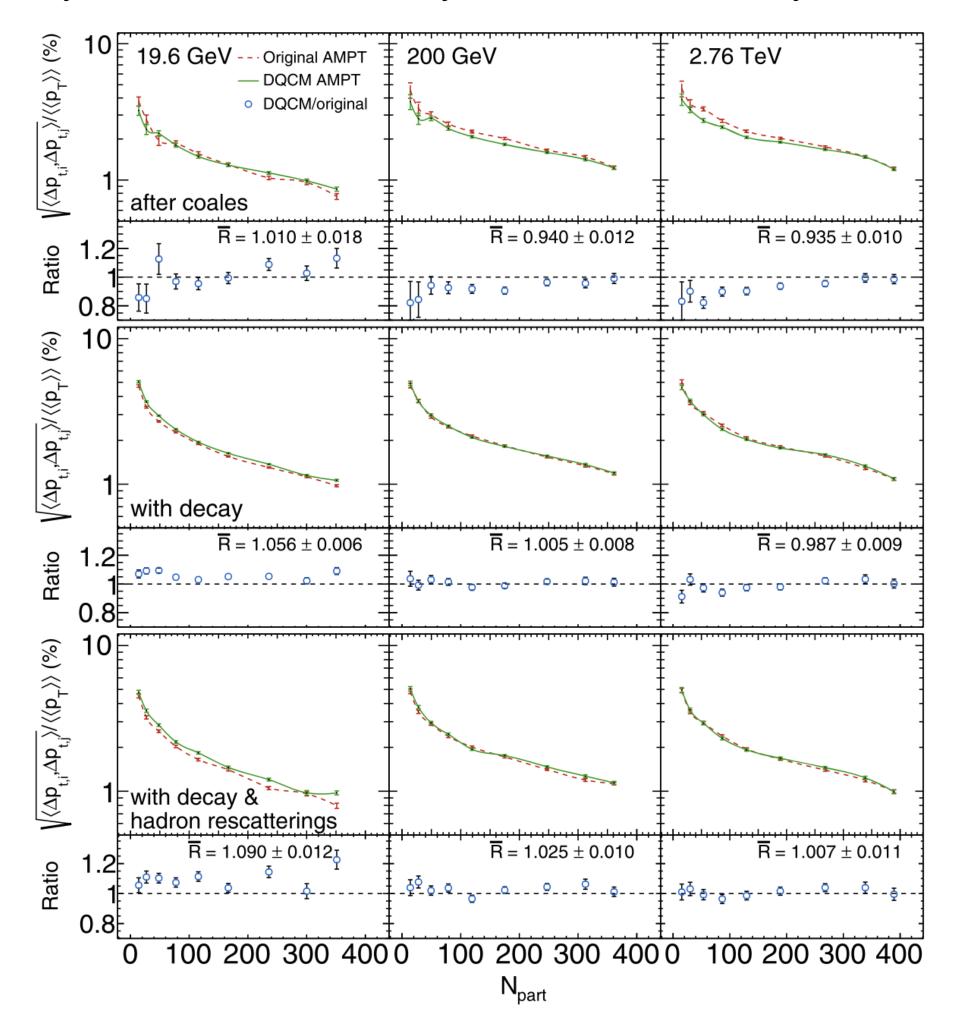








Zhenyu Xu et al 2020 J. Phys. G: Nucl. Part. Phys. 47 125102



The higher-order $\langle p_{\rm T} \rangle$ fluctuations need to be systematically explored considering the BES program.

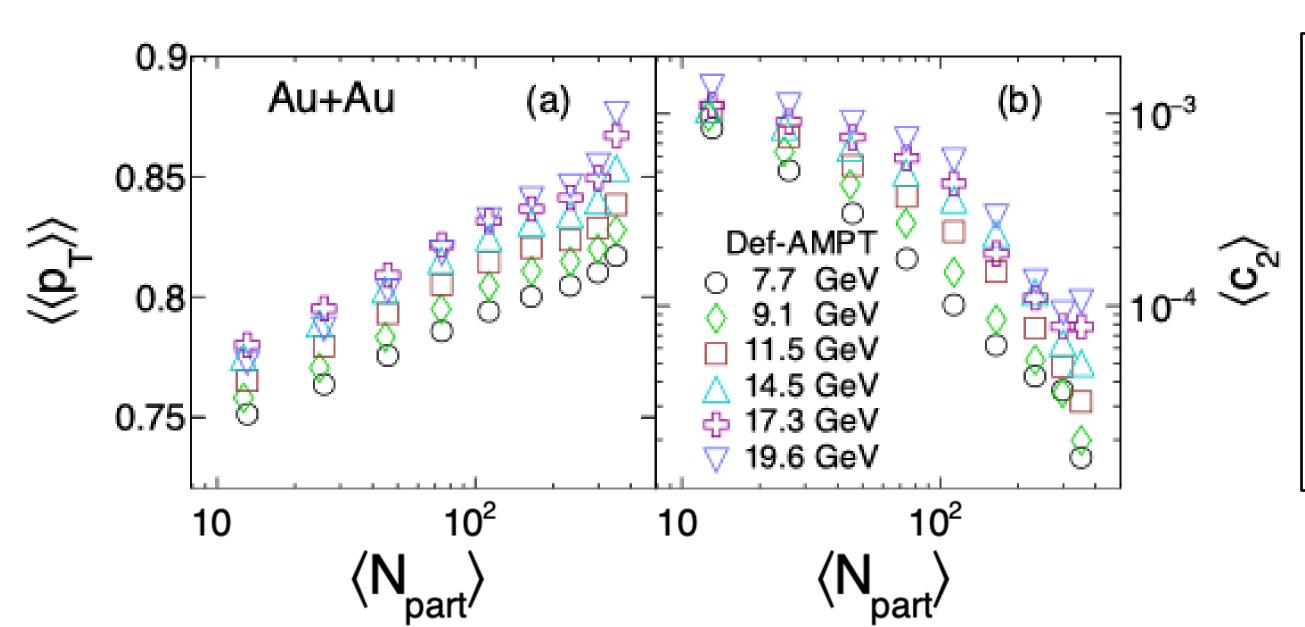
 $<\Delta p_{t,i}\Delta p_{t,j}>[(MeV/c)^2]$

Second-order p_T cumulants





Variance($\langle c_2 \rangle$)



$\langle c_2 \rangle$ (cent.) vs energy

An inverse dependence on centrality is conserved across all energies.

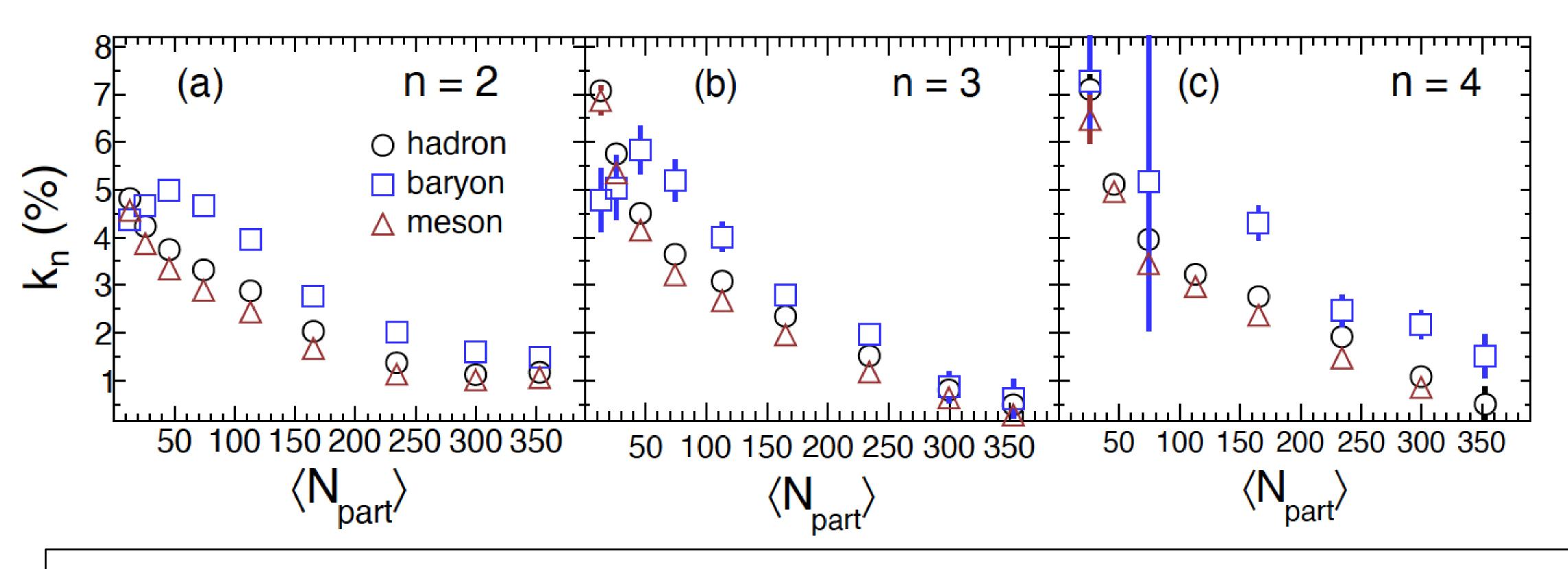
May from a reduction in particle-pair correlations if they are dominated by particles originating from the same NN collisions.

Scaled p_T cumulants: baryon Vs meson





To further explore the radial flow mechanism.



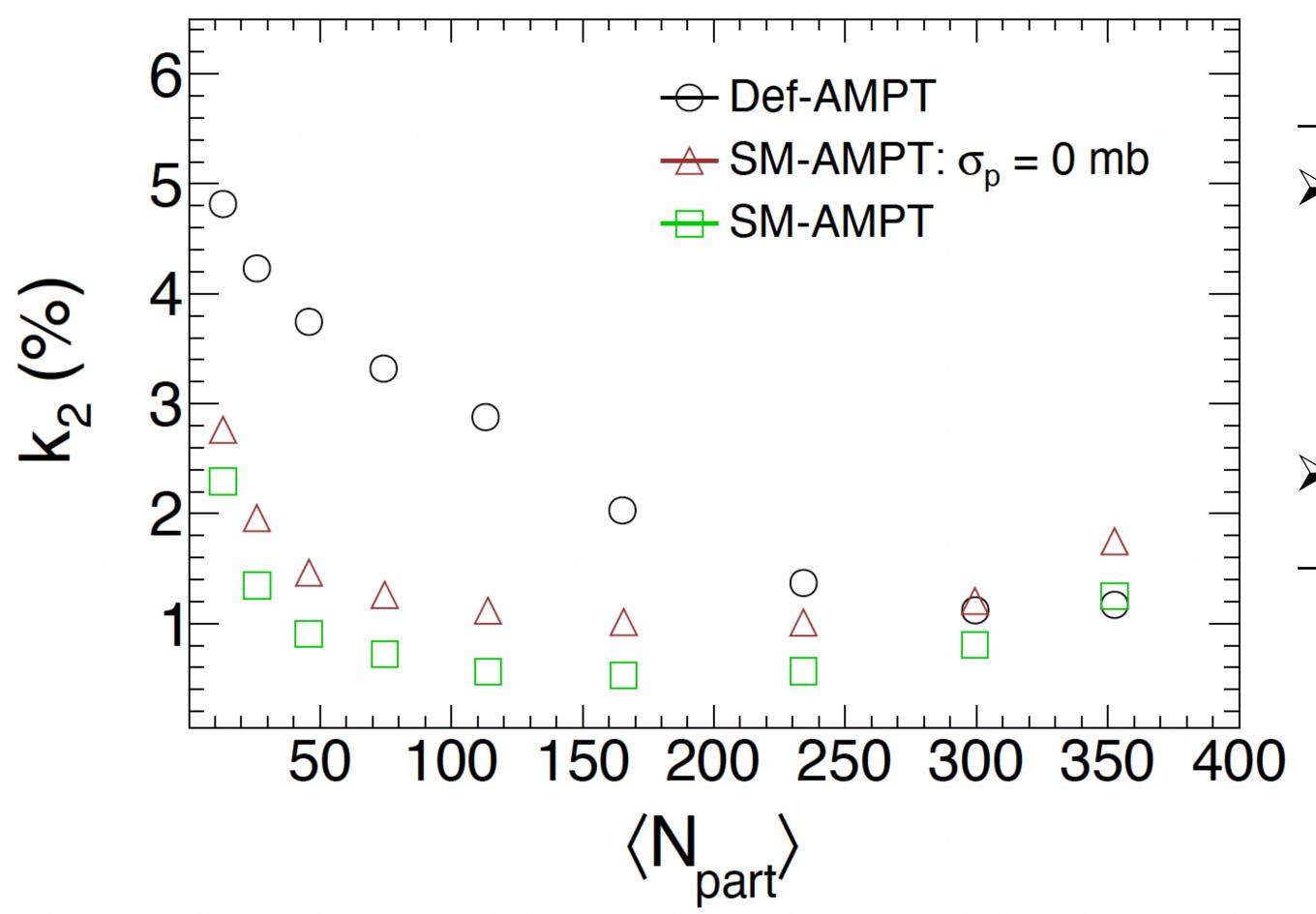
Baryons exhibit more pronounced fluctuations in all scaled variance, skewness and kurtosis compared to mesons.

→ this behavior might be attributed to the effects of radial flow.

Scaled p_T cumulants: SM- Vs default-AMPT







- > The default configuration exhibits more pronounced $\langle p_{\rm T} \rangle$ fluctuations than SM version with and without partonic interactions.
 - → likely different hadronization mechanisms.
- > Partonic evolution slightly suppress the EbE fluctuations.