

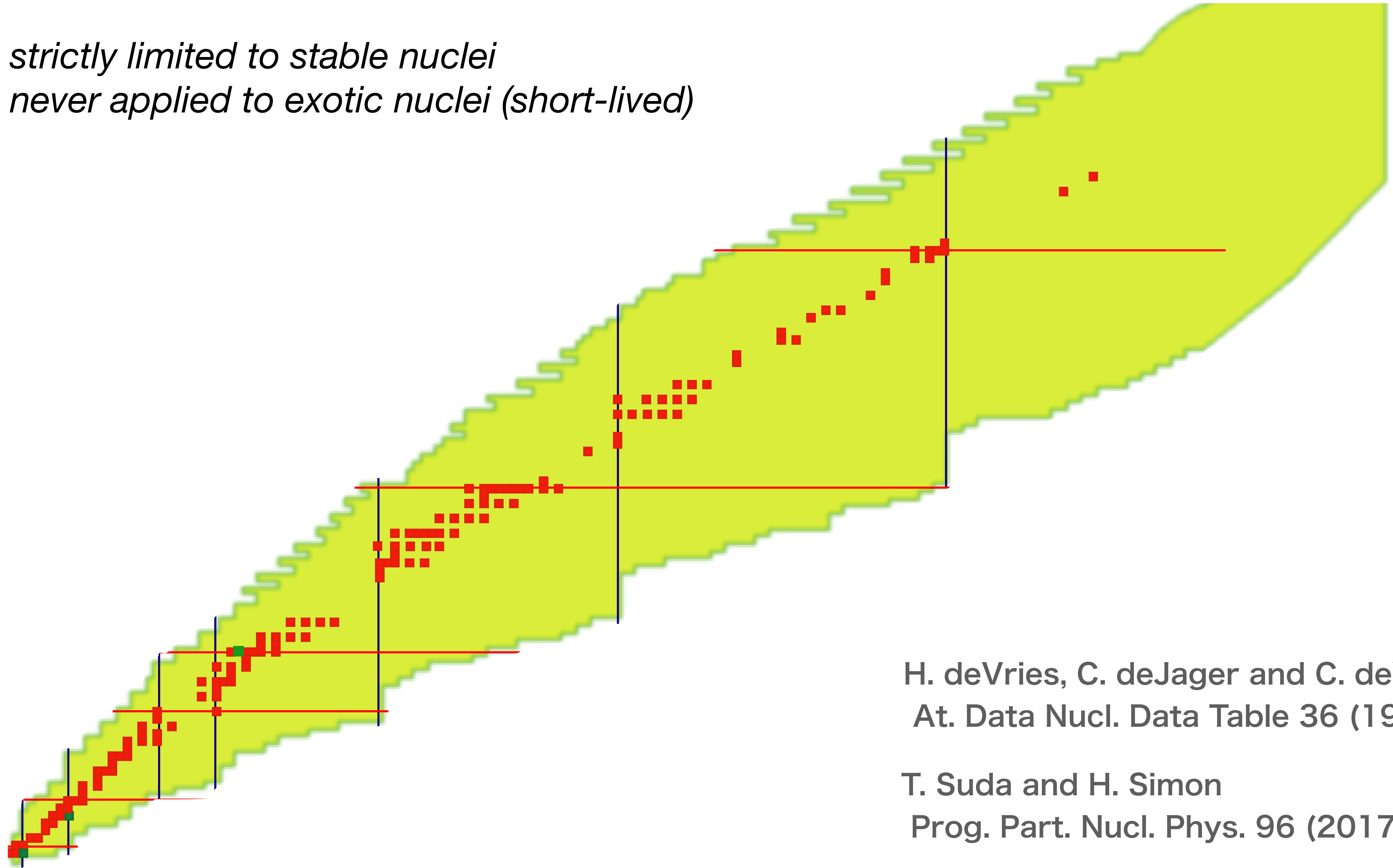
The SCRIT Electron-Scattering Facility for Short-Lived Exotic Nuclei

Toshimi Suda (RARiS, Tohoku Univ.)
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for SCRIT collaboration

Nuclei studied by electron scattering

- *strictly limited to stable nuclei*
- *never applied to exotic nuclei (short-lived)*



H. deVries, C. deJager and C. deVries
At. Data Nucl. Data Table 36 (1987) 495.

T. Suda and H. Simon
Prog. Part. Nucl. Phys. 96 (2017) 1.

PHYSICAL REVIEW LETTERS **131**, 092502 (2023)

Editors' Suggestion

Featured in Physics

First Observation of Electron Scattering from Online-Produced Radioactive Target

K. Tsukada^{1,2}, Y. Abe², A. Enokizono^{2,3}, T. Goke⁴, M. Hara², Y. Honda^{2,4}, T. Hori², S. Ichikawa^{2,*},
Y. Ito¹, K. Kurita³, C. Legris⁴, Y. Maehara¹, T. Ohnishi², R. Ogawara^{1,2}, T. Suda^{2,4},
T. Tamae⁴, M. Wakasugi^{1,2}, M. Watanabe², and H. Wauke^{2,4}

¹*Institute for Chemical Research, Kyoto University, Uji, Kyoto 611-0011, Japan*

²*Nishina Center for Accelerator-Based Science, RIKEN, Wako, Saitama 351-0198, Japan*

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 (Received 7 March 2023; accepted 21 June 2023; published 30 August 2023)

We successfully performed electron scattering off unstable nuclei which were produced online from the photofission of uranium. The target ^{137}Cs ions were trapped with a new target-forming technique that makes a high-density stationary target from a small number of ions by confining them in an electron storage ring. After developments of target generation and transportation systems and the beam stacking method to increase the ion beam intensity up to approximately 2×10^7 ions per pulse beam, an average luminosity of $0.9 \times 10^{26} \text{ cm}^{-2} \text{ s}^{-1}$ was achieved for ^{137}Cs . The obtained angular distribution of elastically scattered electrons is consistent with a calculation. This success marks the realization of the anticipated femtoscope which clarifies the structures of exotic and short-lived unstable nuclei.

DOI: [10.1103/PhysRevLett.131.092502](https://doi.org/10.1103/PhysRevLett.131.092502)

Short-lived
tigated world

K. Tsukada et al.,
Phys. Rev. Lett. **131** (2023) 092502.

charge density distribution
ture determined by elastic

First Observation of Electron Scattering

K. Tsukada^{1,2}, Y. Abe,² A. Enokizono,^{2,3} T. G. Y. Ito,¹ K. Kurita³, C. Legris⁴, Y. Maeda,¹ T. Tamae,⁴ M. Wakasugi,¹

¹Institute for Chemical Research, Kyoto University

²Nishina Center for Accelerator-Based Science, RIKEN

³Department of Physics, Rikkyo University

⁴Research Center for Electron Photon Science, KEK

(Received 7 March 2023; accepted 14 April 2023)

We successfully performed electron scattering from a target of ^{137}Cs ions produced by photofission of uranium. The target ^{137}Cs ions make a high-density stationary target from a small ring. After developments of target generation and increase the ion beam intensity up to approximately $0.9 \times 10^{26} \text{ cm}^{-2} \text{ s}^{-1}$ was achieved for ^{137}Cs . The electron scattering cross section is consistent with a calculation. This success clarifies the structures of exotic and short-lived nuclei.

DOI: 10.1103/PhysRevLett.131.092502

Short-lived
nuclei investigated world

K. Tsukada et al.,
Phys. Rev. Lett. 131, 092502 (2023)

Electron scattering provides a long-awaited view of unstable nuclei

Nuclear reactions produce a plethora of short-lived artificial isotopes. Figuring out what they look like has been a challenge.

The cartoon picture of an atomic nucleus looks kind of like the inside of a gumball machine that dispenses only two flavors: protons and neutrons, evenly mixed in a compact, spherical cluster.

That's not generally what real nuclei look like. Neutron-rich lead-208, for example, has a thick skin of neutrons encasing its proton-endowed core (see *Physics Today*, July 2021, page 12). Some nuclei are flattened, and some are elongated. Some are even pear shaped.

The more unstable a nucleus, the

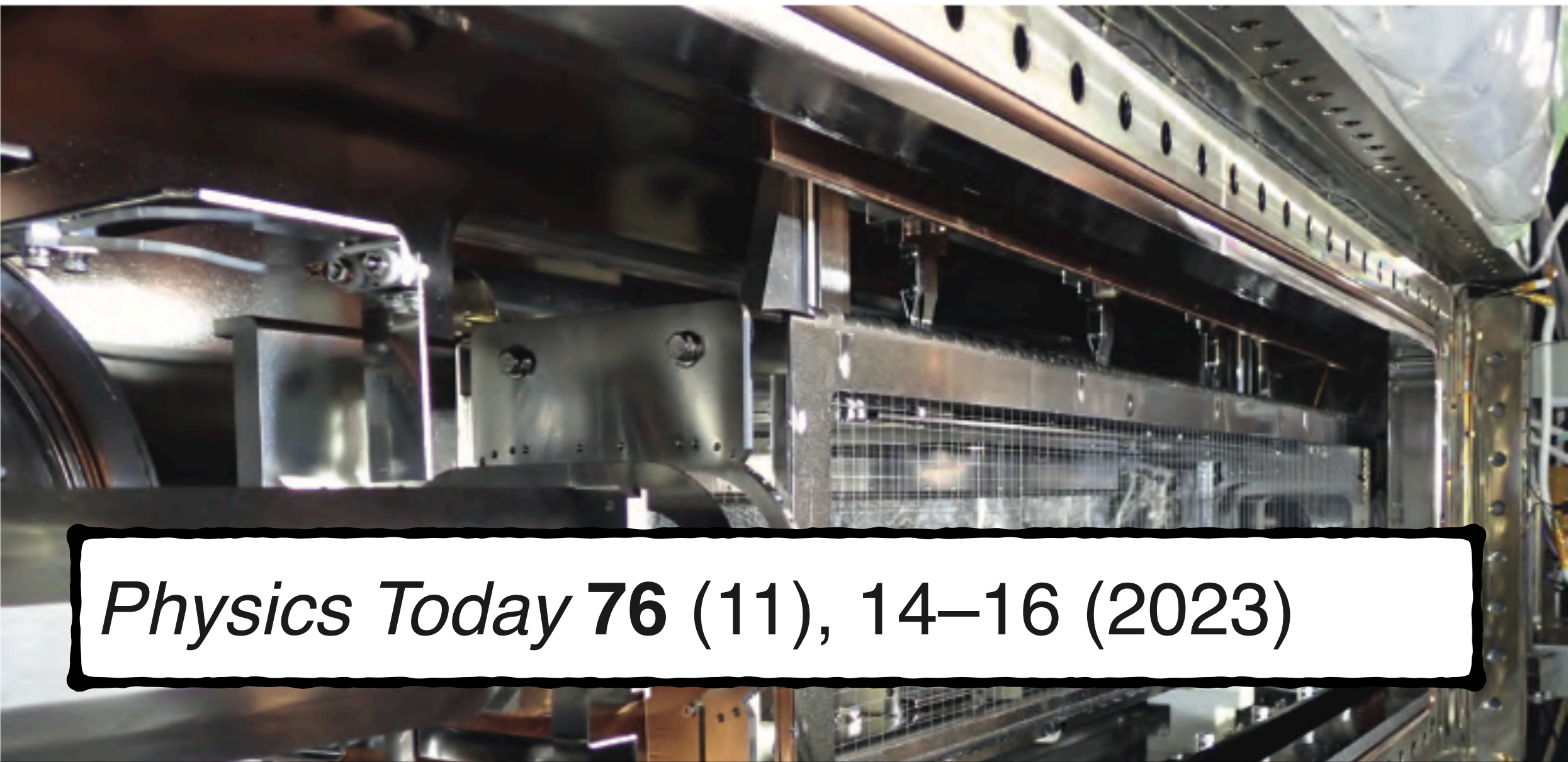
stranger the structures it can adopt. Short-lived nuclei might form bubble structures with depleted central density, or they might have a valence nucleon or two that form a halo around a compact central core. (See the article by Filomena Nunes, *Physics Today*, May 2021, page 34.) Frustratingly, though, those exotic structures are hard to experimentally confirm, because the gold standard for probing nuclear structure—electron scattering—has been off limits to short-lived nuclei.

That could change soon. Kyo Tsukada

and colleagues, working at RIKEN's Radioactive Isotope Beam Factory (RIBF) in Wako, Japan, have performed the first electron-scattering experiment on unstable nuclei produced on the fly in a nuclear reaction.¹ Their isotope of choice, cesium-137, has a half-life of 30 years. It's not so exotic that the researchers expected—or found—anything unusual about its structure. But the technique they used is applicable to shorter-lived nuclei, so more experiments are on the way.

Backscatter

Probing nuclei through particle scattering dates back to the discovery of the nucleus itself, in 1911, when Ernest



Physics Today 76 (11), 14–16 (2023)

world's first e-scattering from an online-produced radioactive isotope

PHYSICAL REVIEW LETTERS

Editors' Suggestion

Featured in Physics

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K. Tsukada^{1,2}, Y. Abe,² A. Enokizono,^{2,3} T. G.
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³Department of Physics, Rikkyo University

⁴Research Center for Electron Photon Science, KEK

(Received 7 March 2023; accepted 12 April 2023)

We successfully performed electron scattering on the target of photofission of uranium. The target ¹³⁷Cs ions make a high-density stationary target from a small ring. After developments of target generation and increase the ion beam intensity up to approximately $0.9 \times 10^{26} \text{ cm}^{-2} \text{ s}^{-1}$ was achieved for ¹³⁷Cs. The electron scattering is consistent with a calculation. This success clarifies the structures of exotic and short-lived nuclei.

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K. Tsukada et al.,
Phys. Rev. Lett. 131, 092502 (2023)

Short-lived
nuclei investigated world

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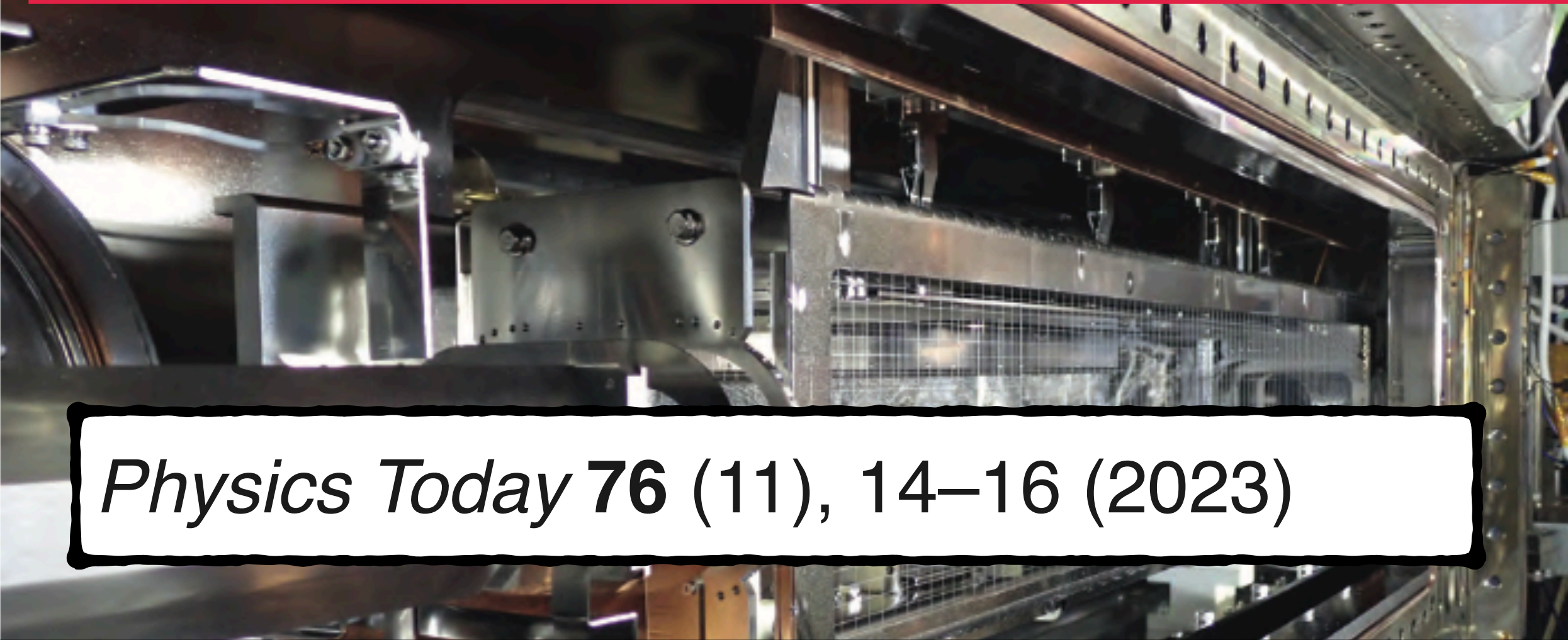
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electron scattering - the gold standard for probing nuclear structure - has been off limits to short-lived exotic nuclei



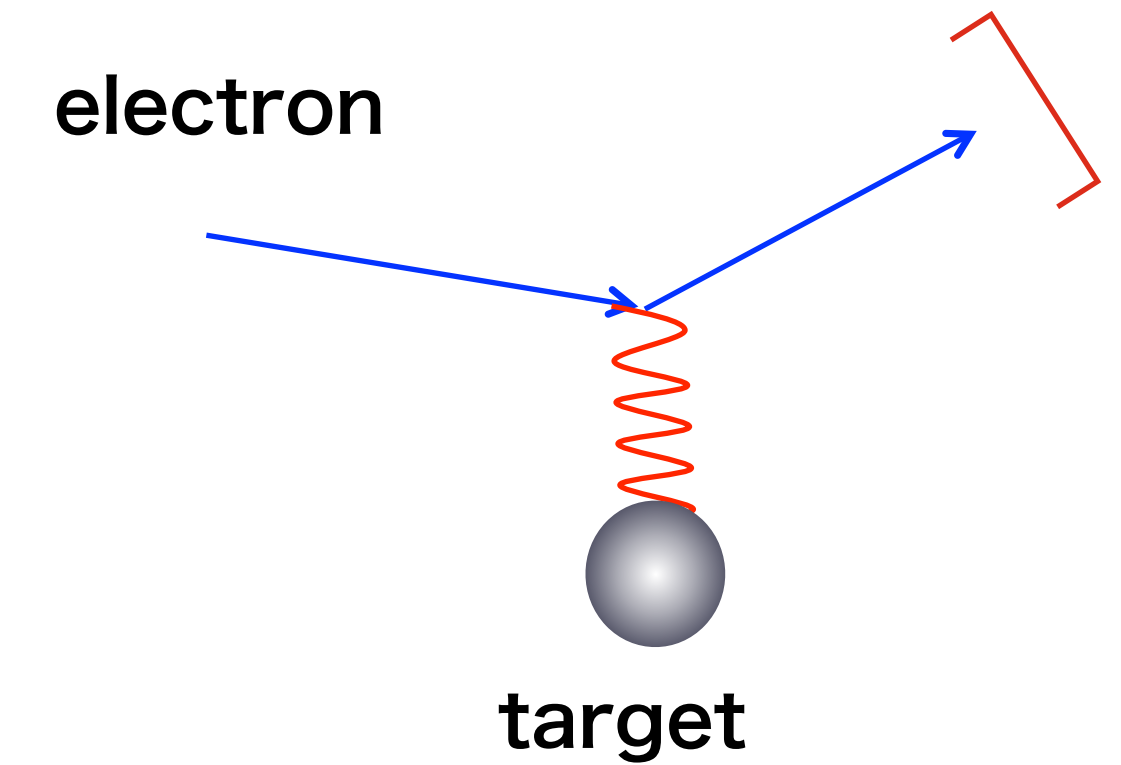
Physics Today 76 (11), 14–16 (2023)

Electron Scattering

Electron scattering for nuclear physics

a unique and powerful probe for structure studies

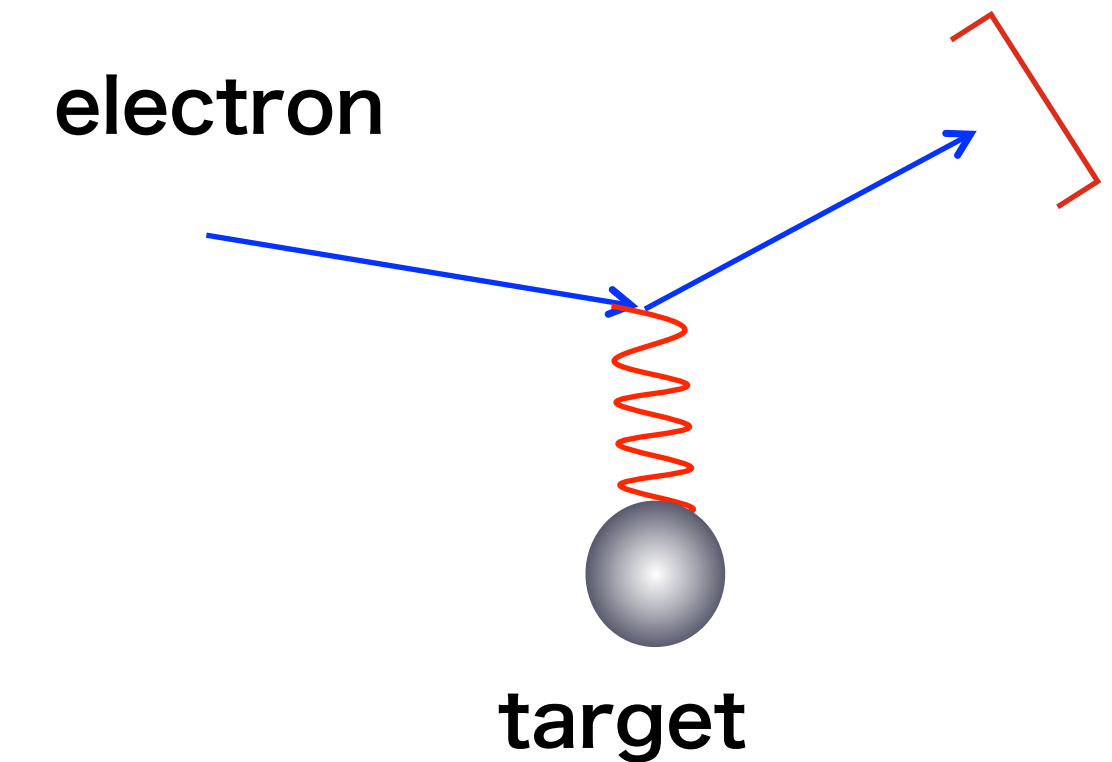
1. elementary particle - structureless -
2. electro-weak interaction - best understood -
3. “relatively” weak - probing deep inside of the target -
4. variable q (mom. transfer) for fixed ω (energy transfer)



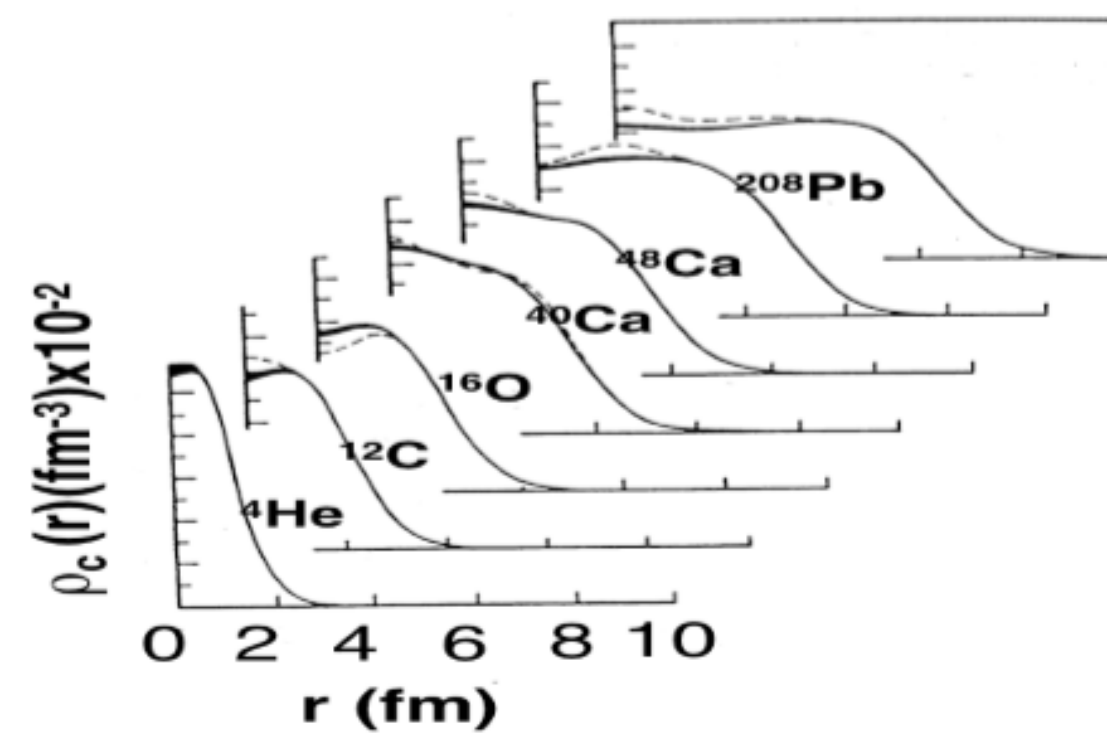
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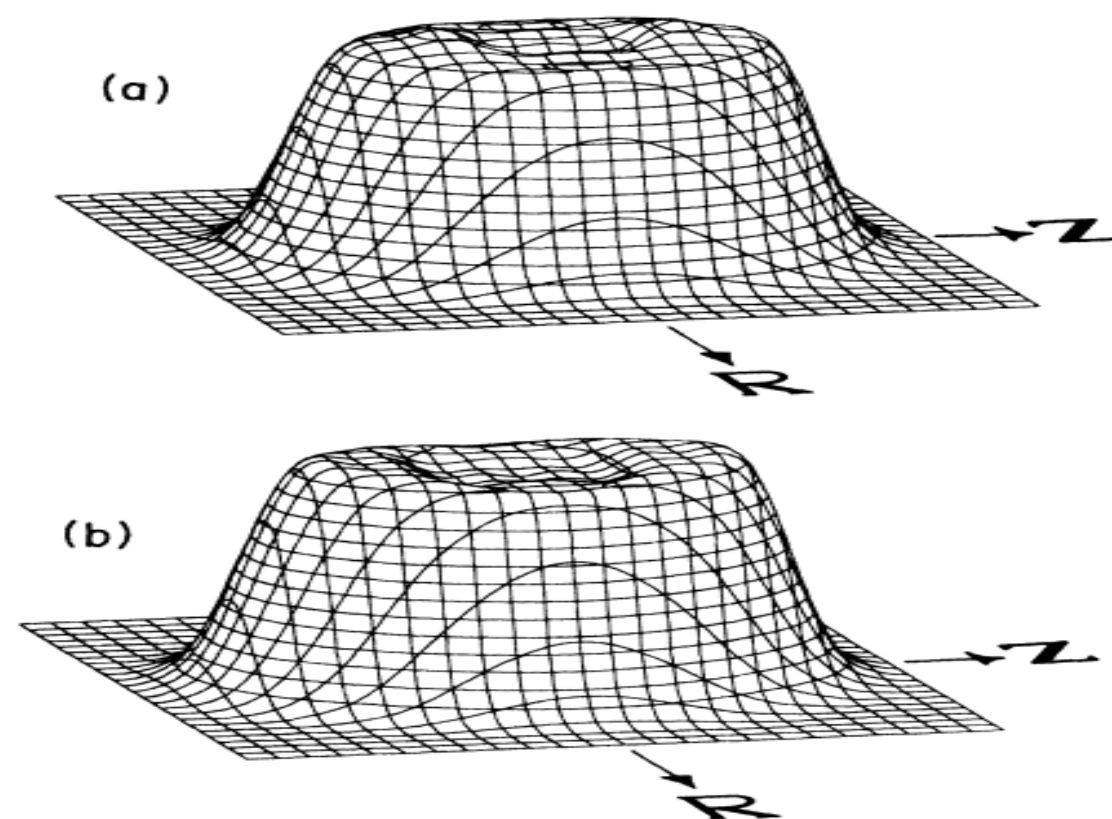
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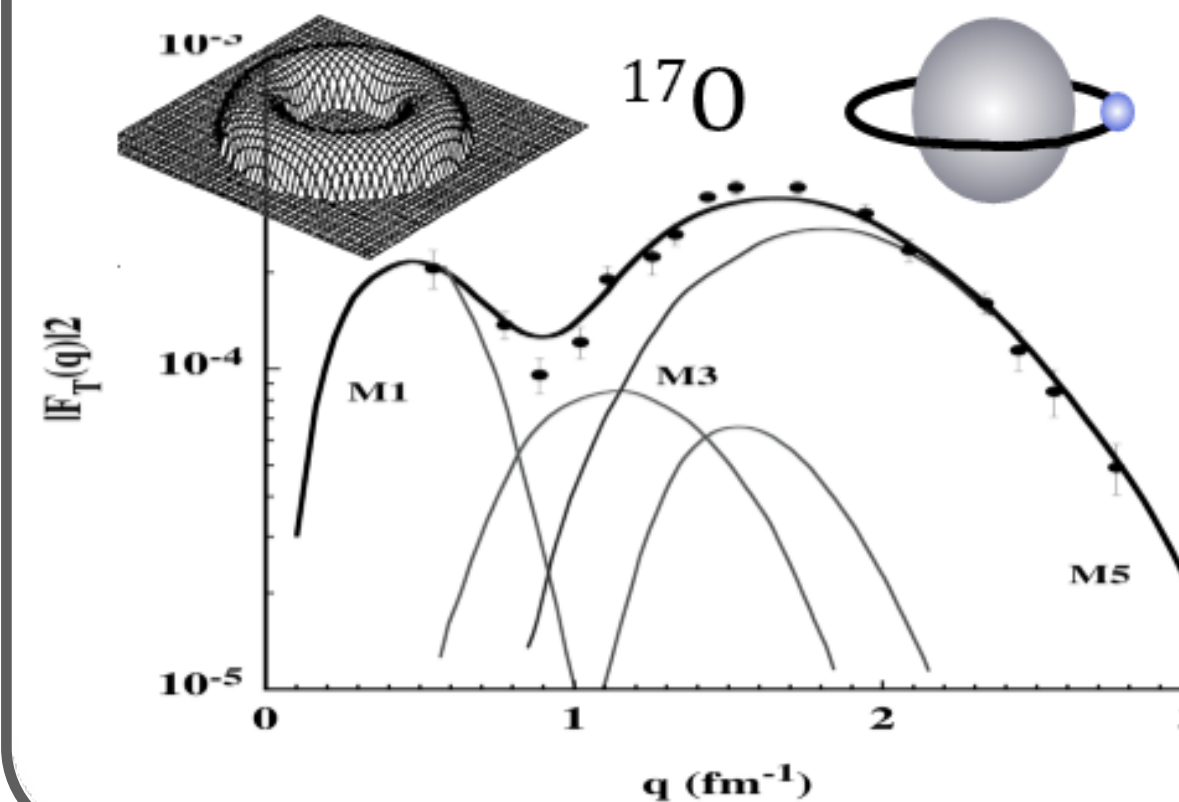
elastic scattering
→ charge distribution



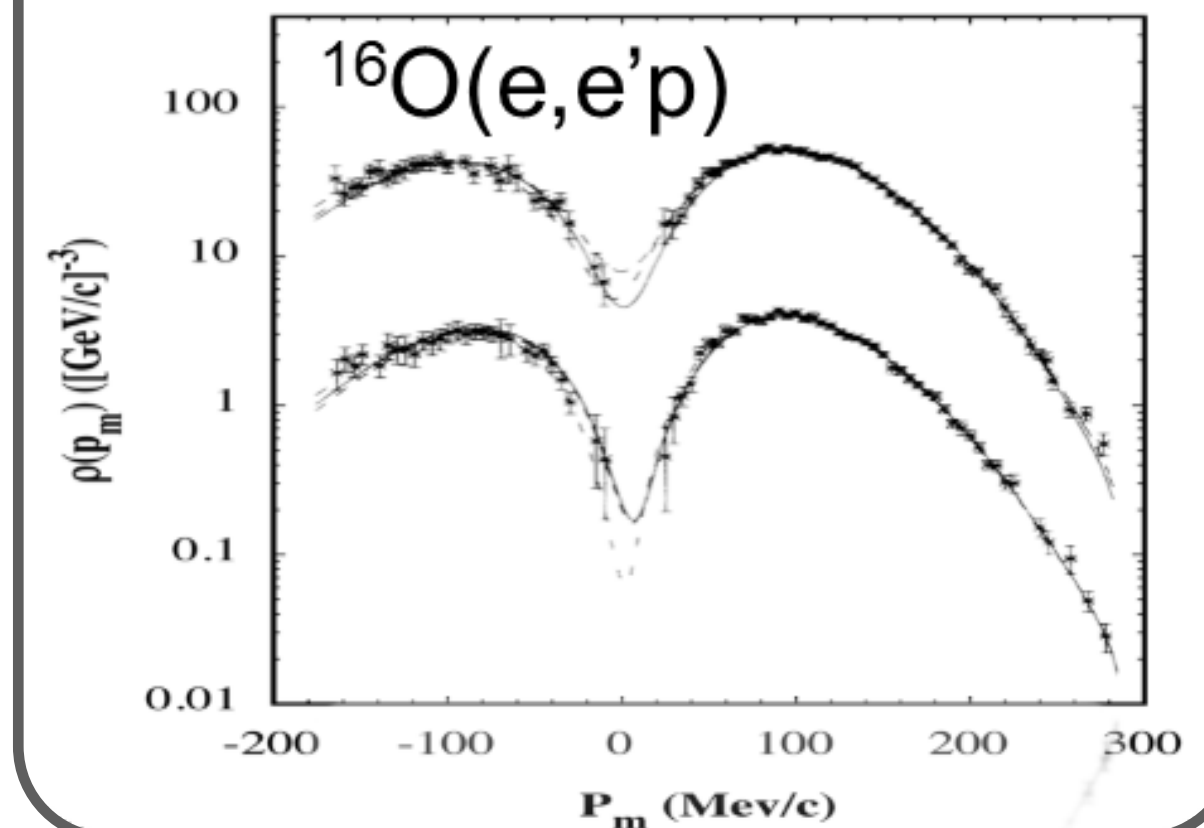
elastic + inelastic scattering
→ deformation



magnetic scattering
→ valence neutron



quasi-elastic (e,e'p)
→ S-factor



Elastic electron scattering and charge density distribution

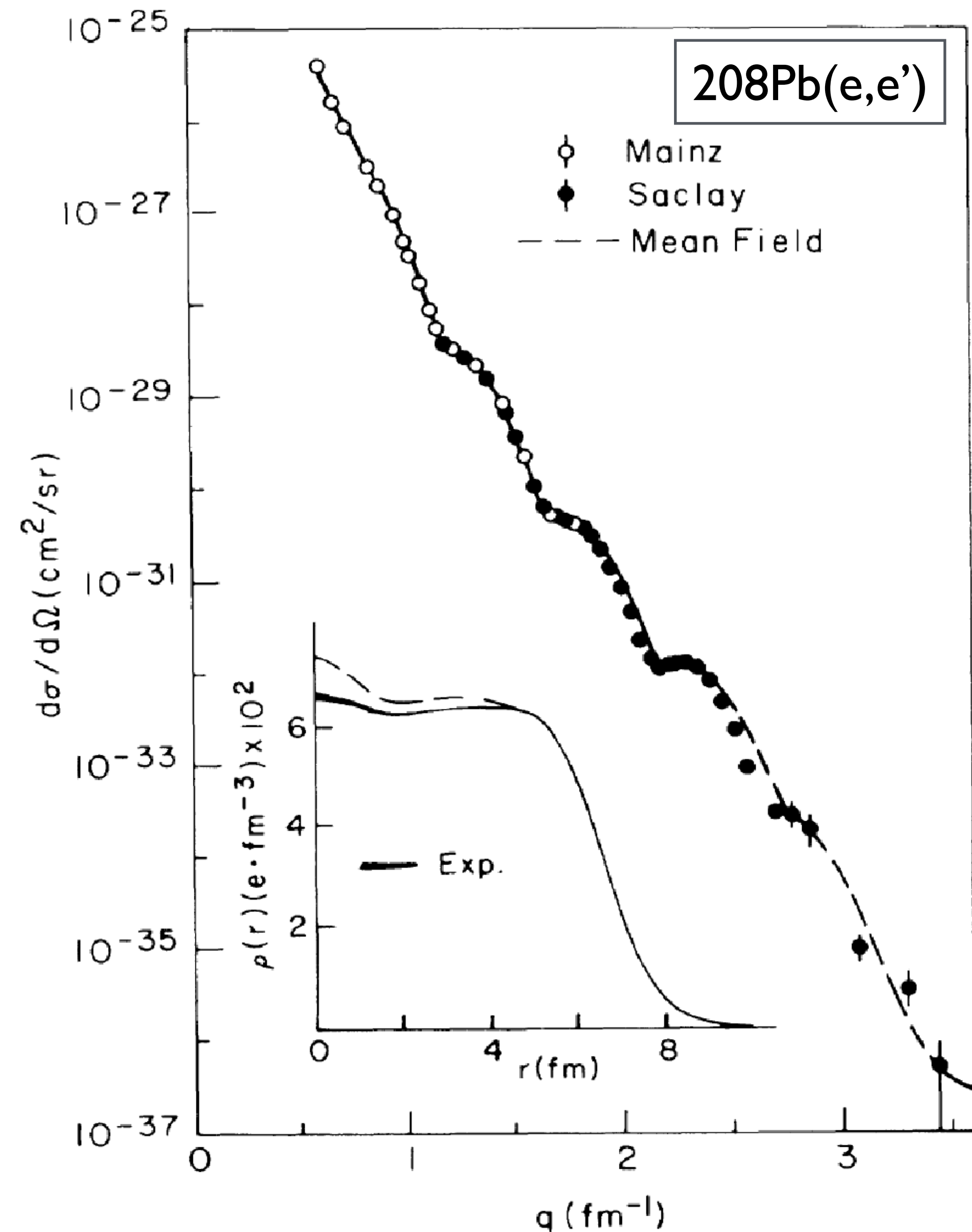
Elastic scattering
for (spin-less) nuclei

$$\frac{d\sigma}{d\Omega} = \frac{d\sigma_{Mott}}{d\Omega} |F_c(q)|^2$$

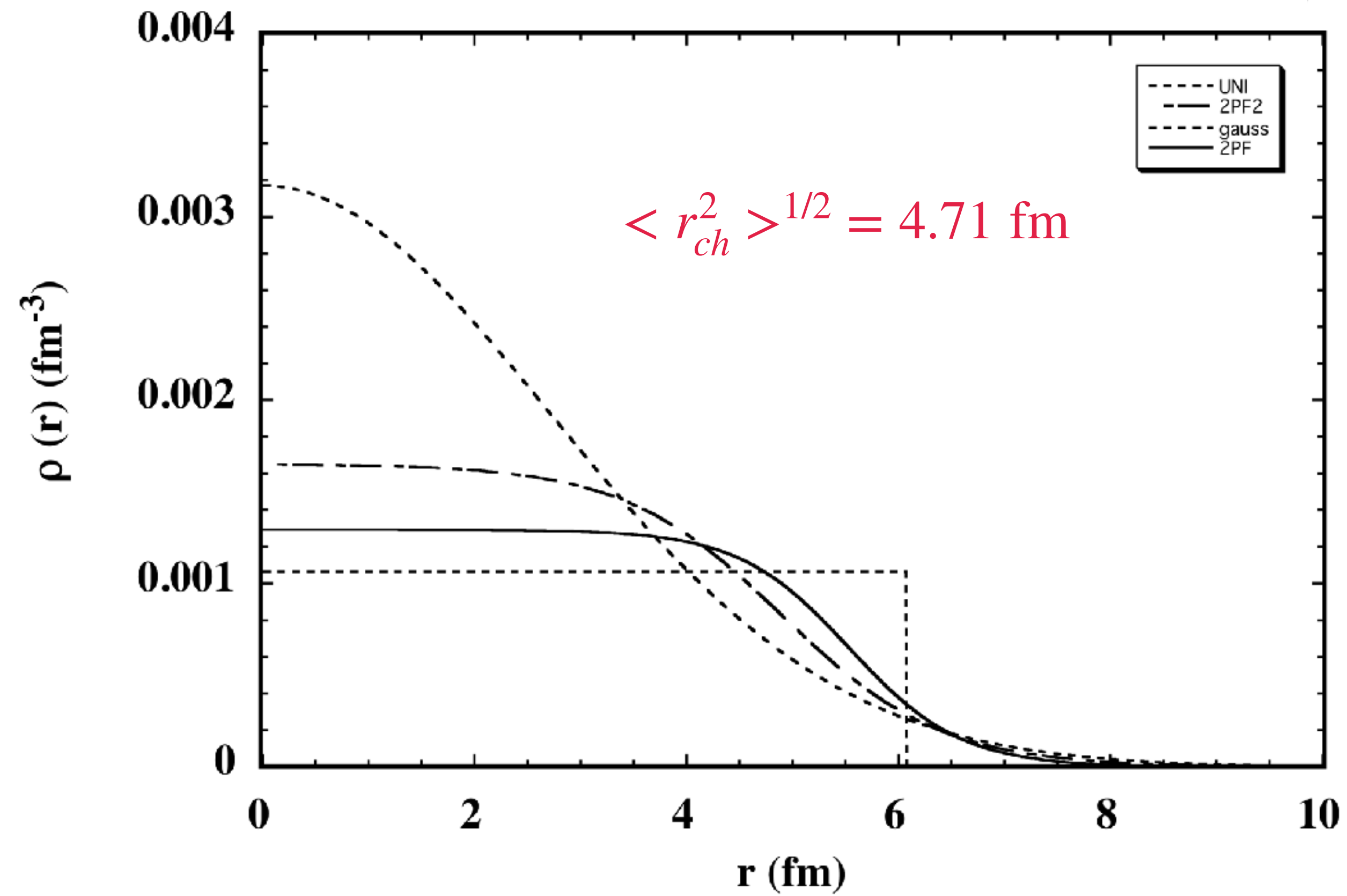
$$F_c(q) = \int \rho_c(r) e^{iqr} d^3r$$

charge distribution

$$\frac{d\sigma_{Mott}}{d\Omega} \propto e^2/q^4$$

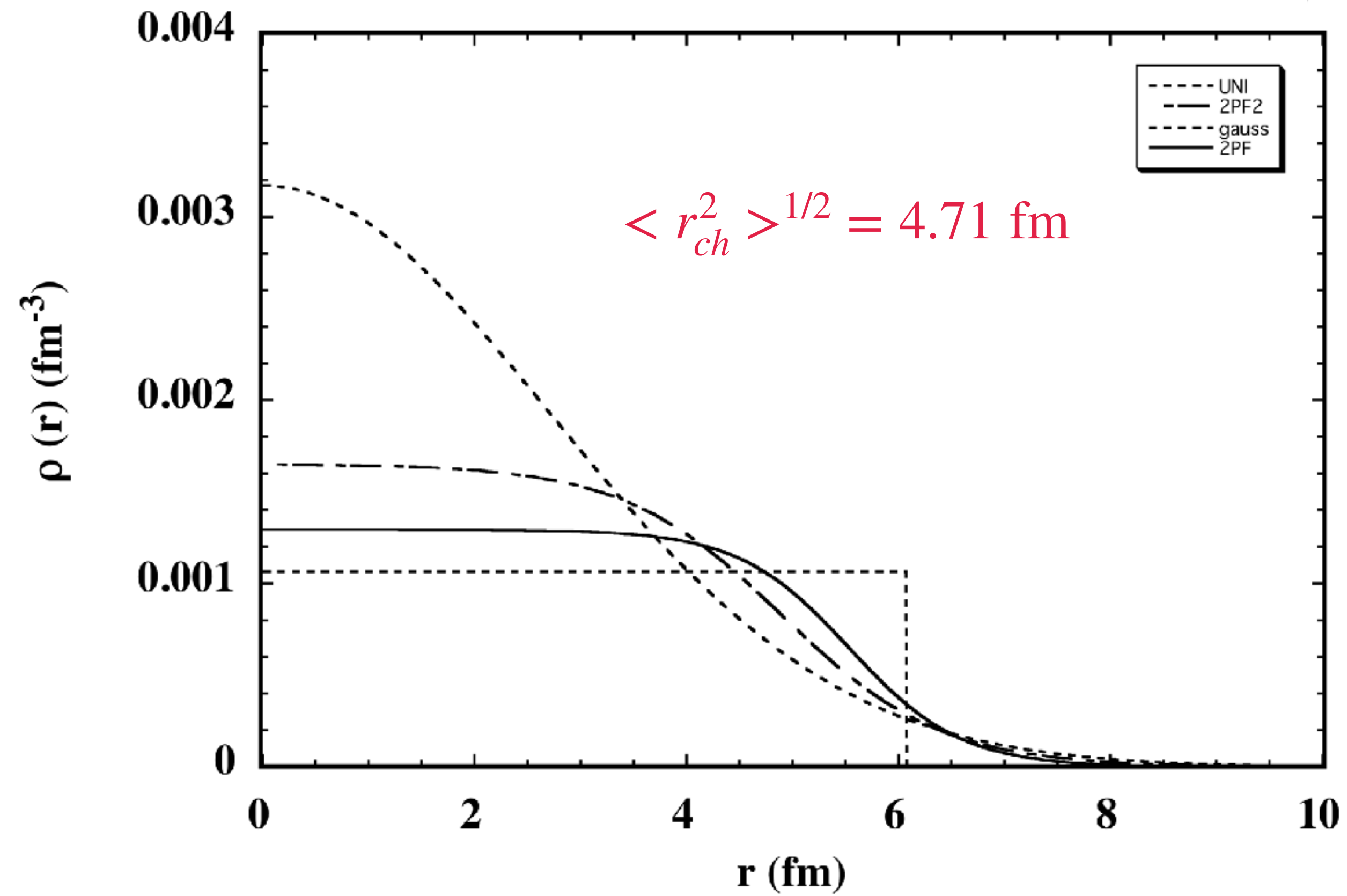


Charge density distributions

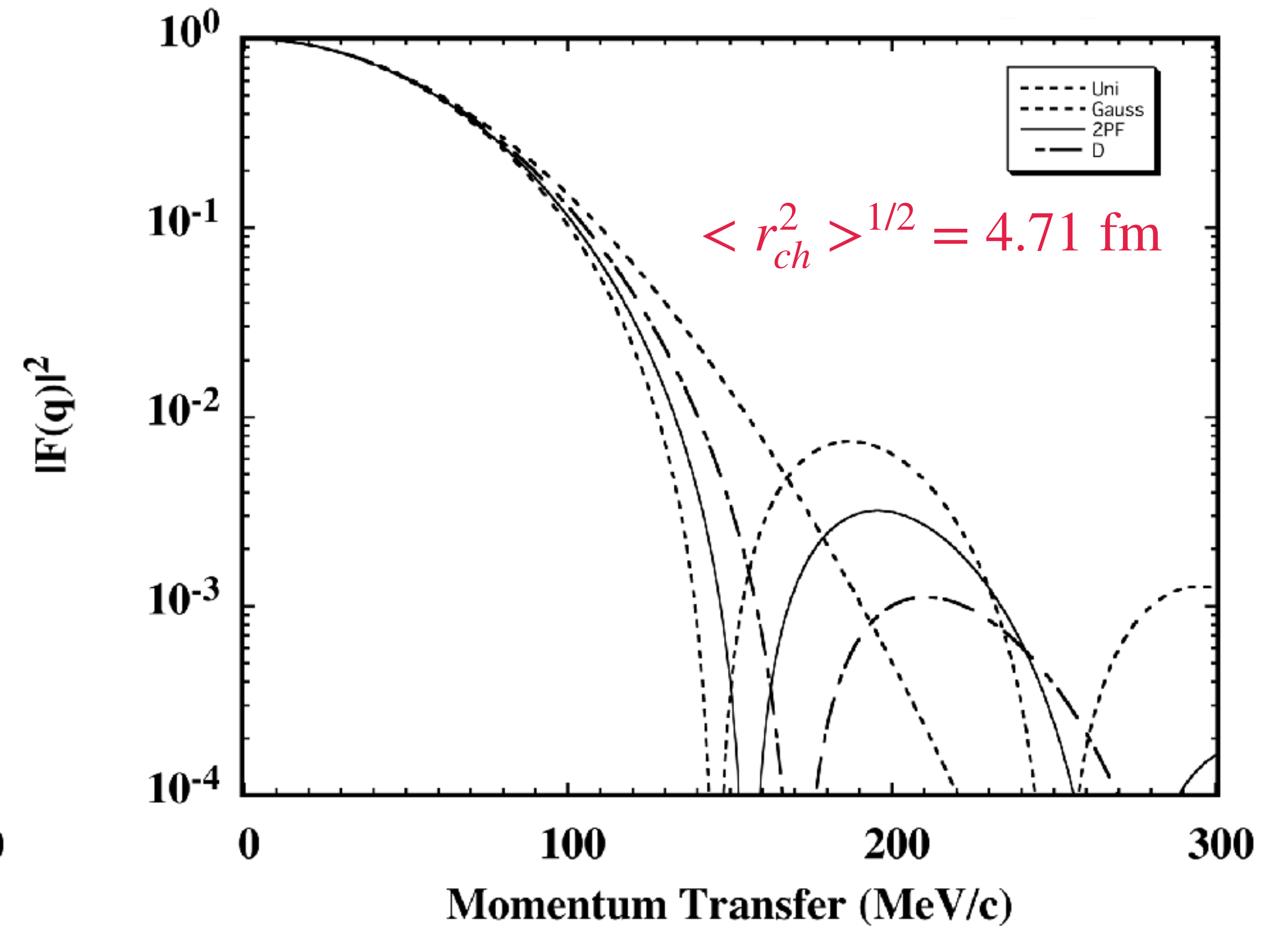


Beyond charge radii (isotope shifts)

Charge density distributions

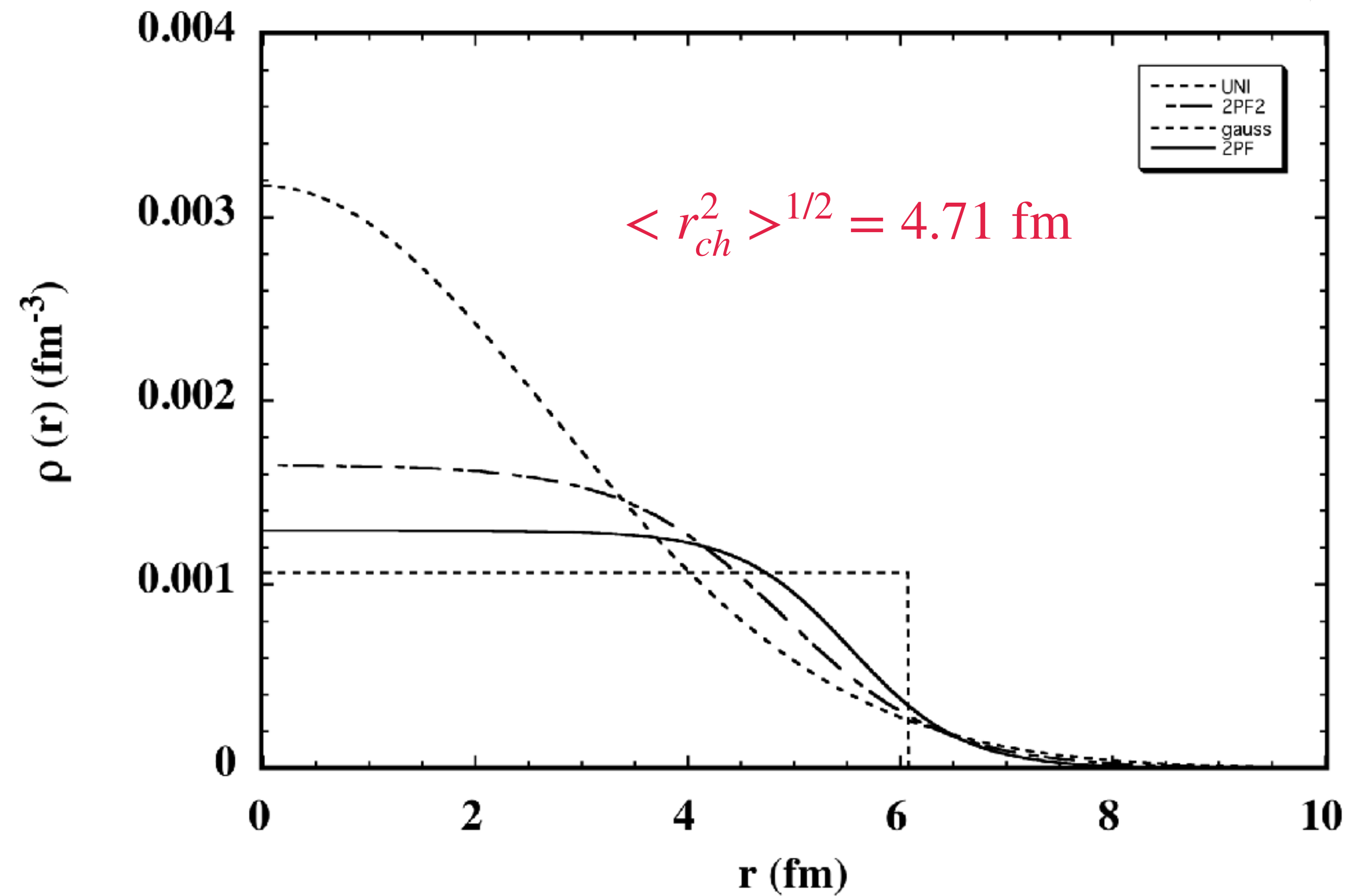


Charge form factors

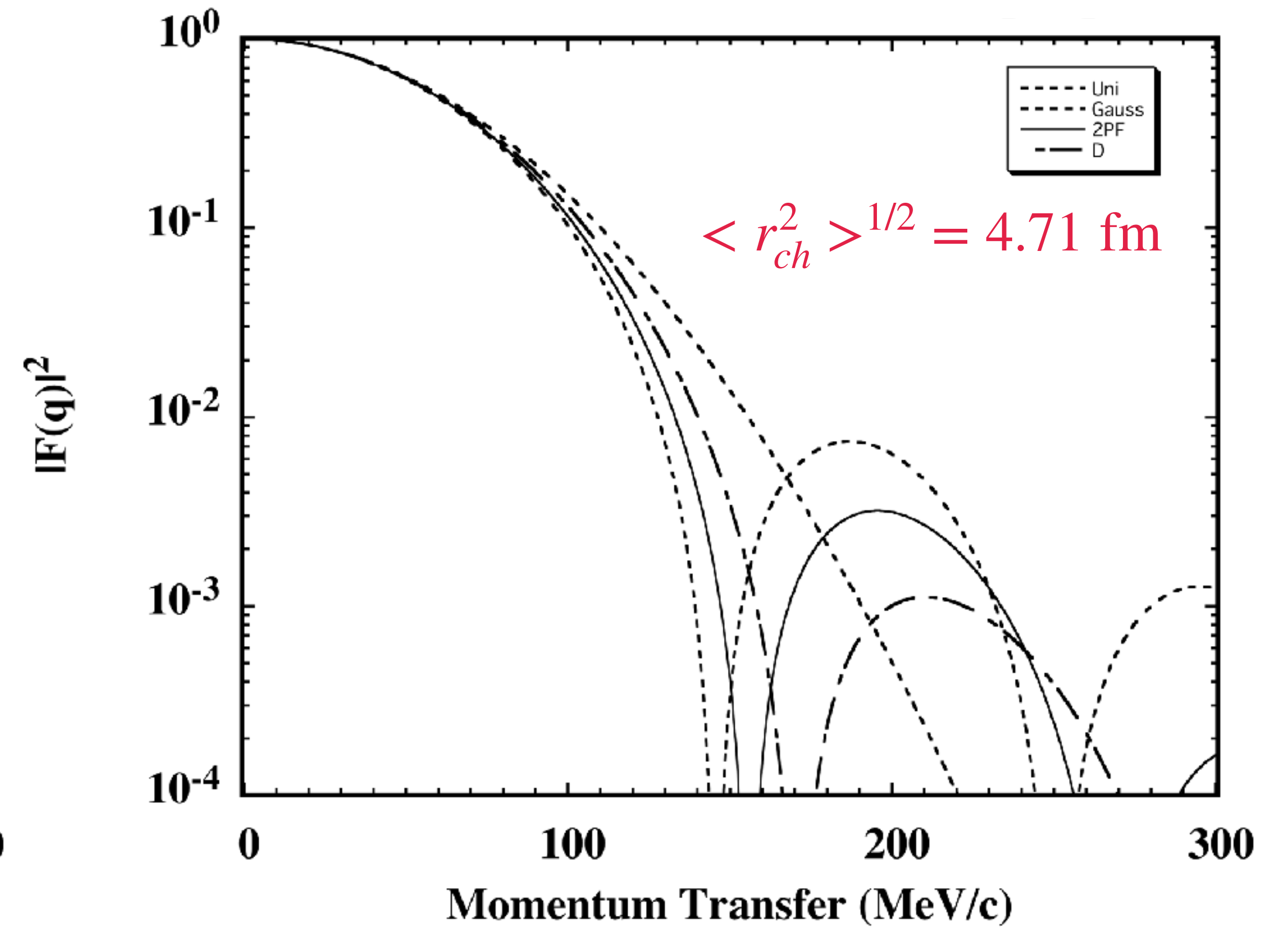


Beyond charge radii (isotope shifts)

Charge density distributions



Charge form factors



Electron scattering reveals both size and shape

Key parameter for e-scattering of exotic nuclei

$$\frac{dN}{dt} = L \times \frac{d\sigma}{d\Omega}$$

Luminosity

Elastic Scattering for Exotic Nuclei

(for medium-heavy nuclei)

$$L \gtrsim 10^{27} / \text{cm}^2 / \text{s}$$

Elastic Scattering for Exotic Nuclei

(for medium-heavy nuclei)

$$L \gtrsim 10^{27} / \text{cm}^2 / \text{s}$$

How to realize this L with small number of exotic nuclei?

The idea of “SCRIT” originated in 2004

SCRIT : Self-Confining Radioactive Ion Target



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Nuclear Instruments and Methods in Physics Research A 532 (2004) 216–223

**NUCLEAR
INSTRUMENTS
& METHODS
IN PHYSICS
RESEARCH**

Section A

www.elsevier.com/locate/nima

A new method for electron-scattering experiments using a self-confining radioactive ion target in an electron storage ring

M. Wakasugi^{a,*}, T. Suda^b, Y. Yano^a

^a *Cyclotron center, RIKEN, Wako-shi, Saitama 351-0198, Japan*

^b *RI Beam Science Laboratory, RIKEN, Wako-shi, Saitama 351-0198, Japan*

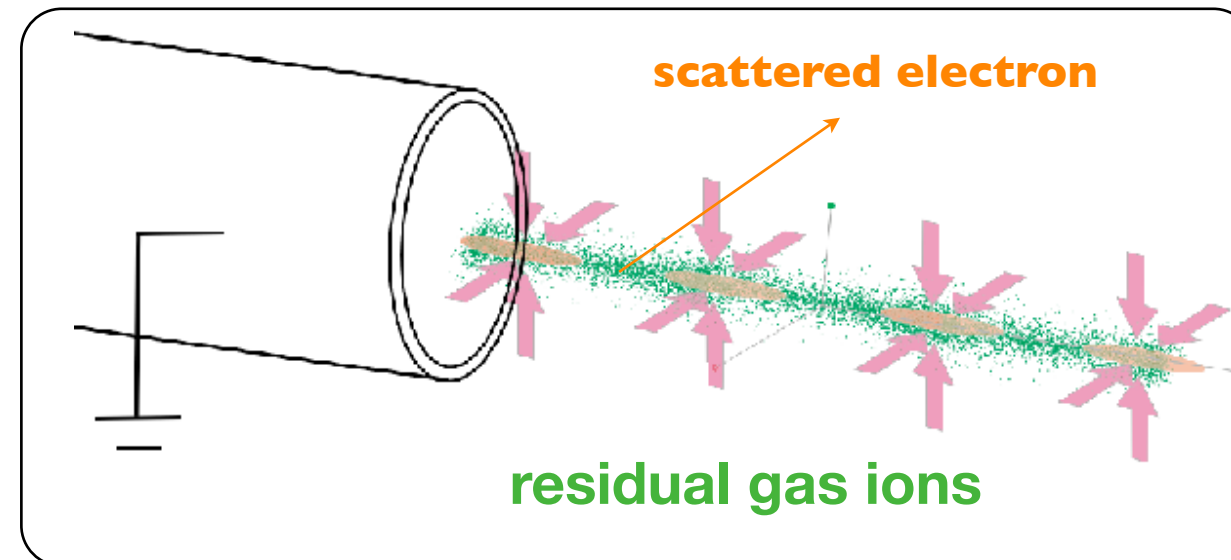
Available online 3 August 2004

Nucl. Instrum. Methods A532 (2004) 216.
Phys. Rev. Lett. 100 (2008) 164801.
Phys. Rev. Lett. 102 (2009) 102501.

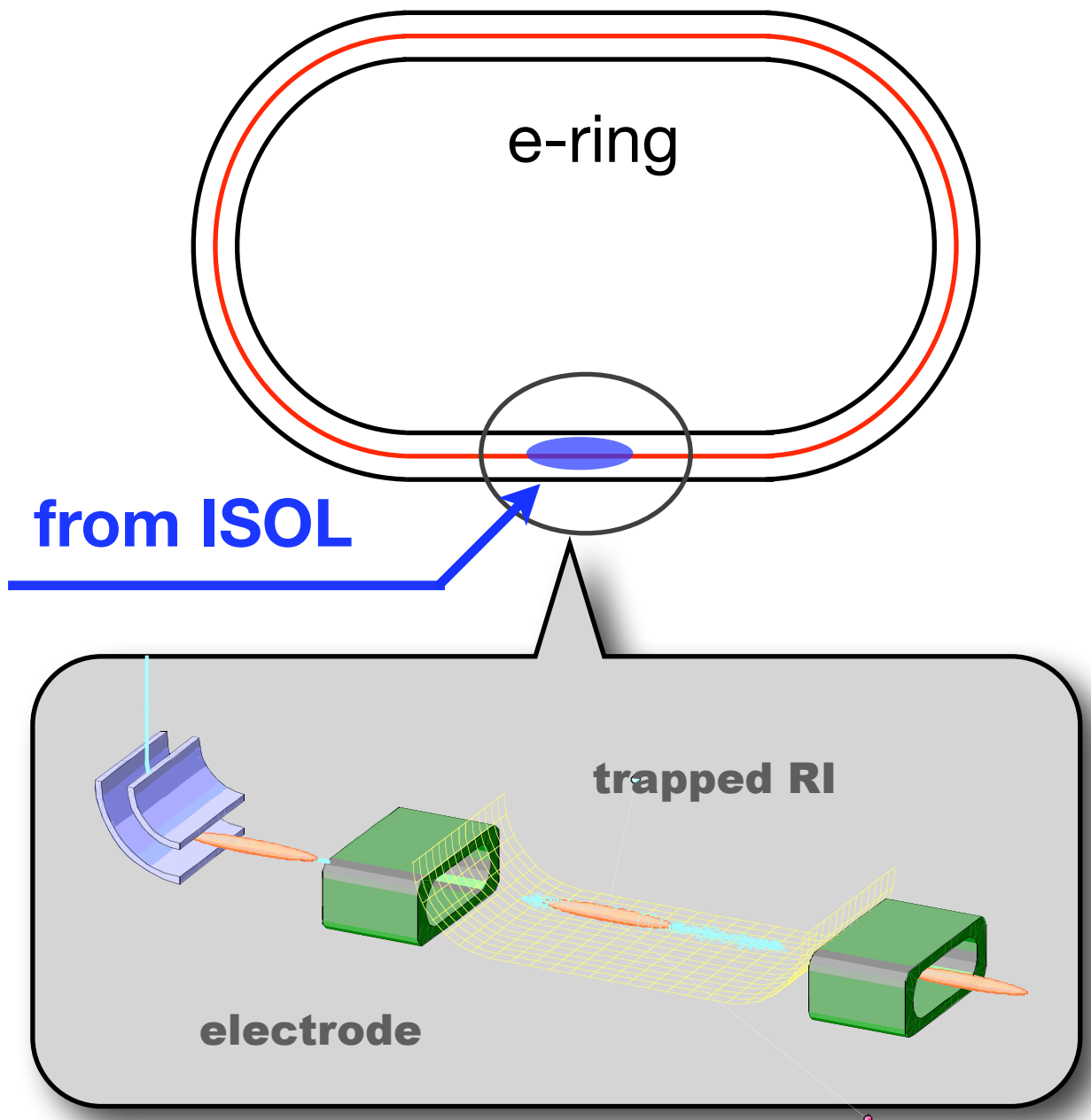
SCRIT (Self-Confining RI Ion Target)

Idea : “ion trapping” at SR facilities.

ionized residual gases are trapped
by the circulating electron beam



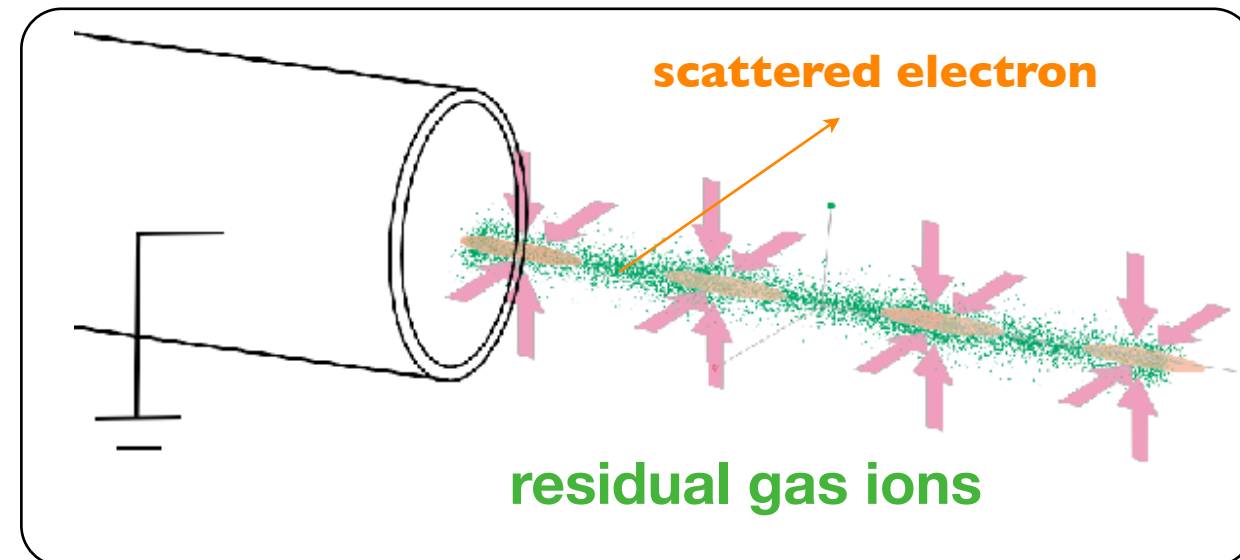
ill problem of e-storage rings



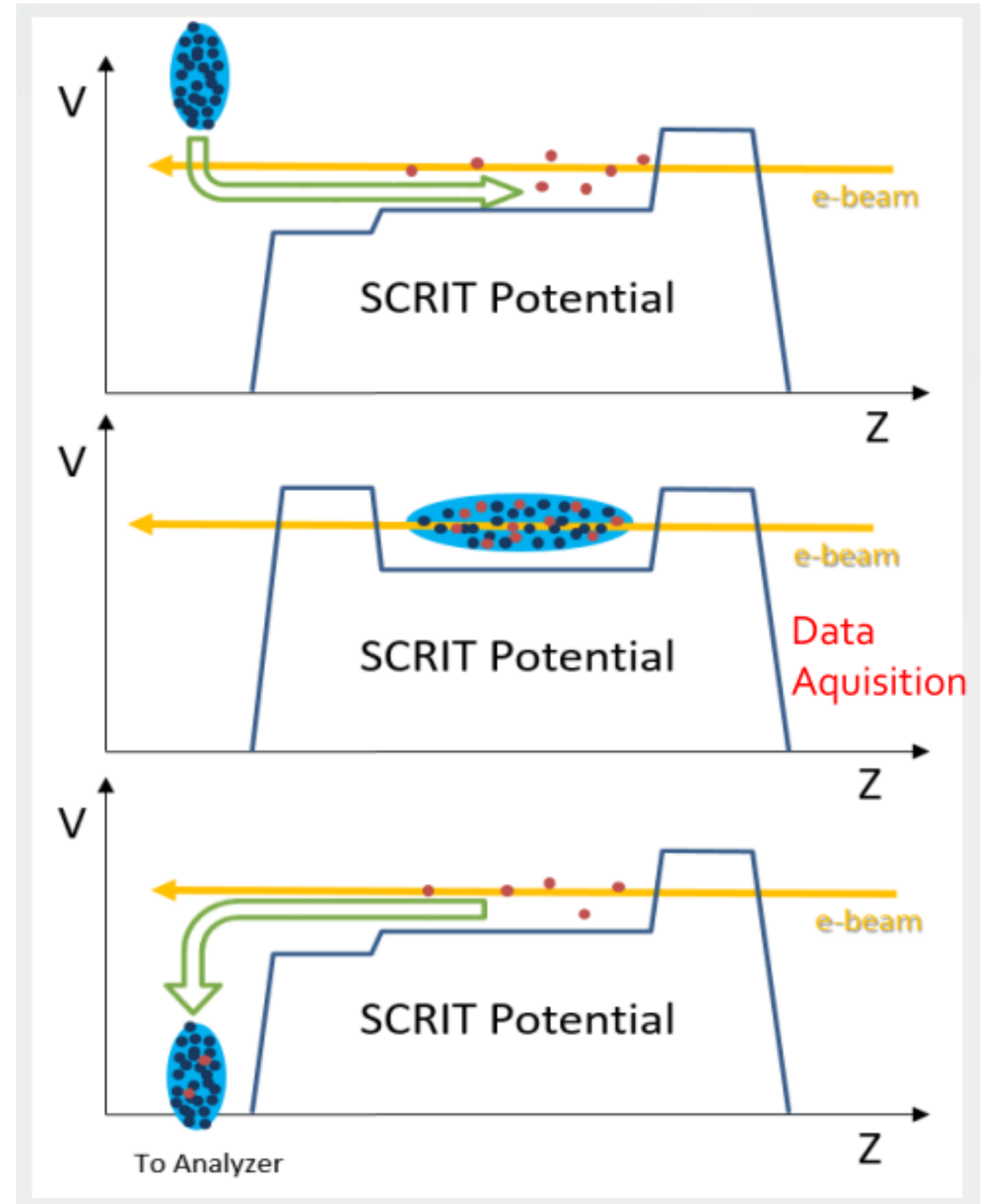
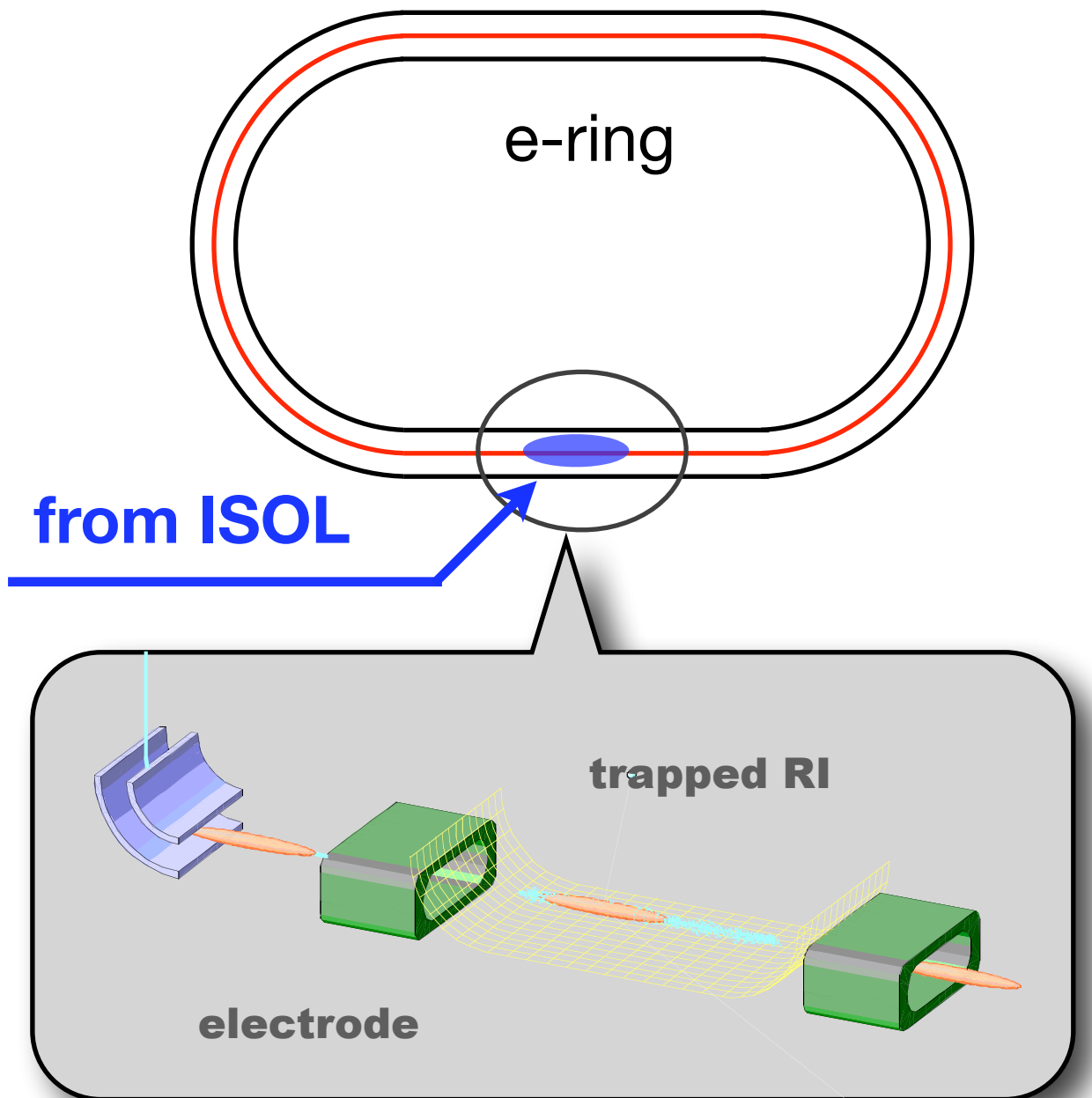
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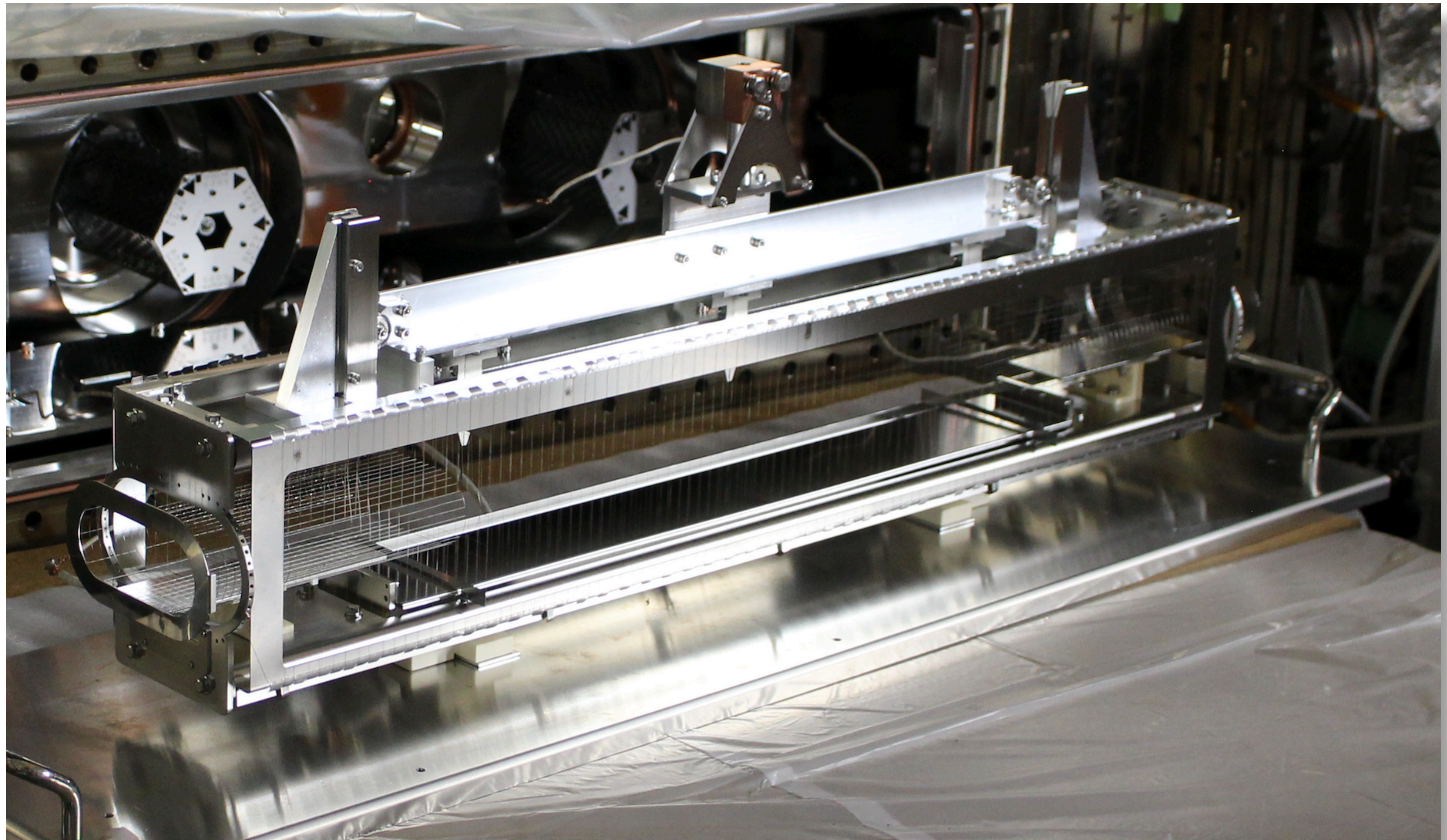
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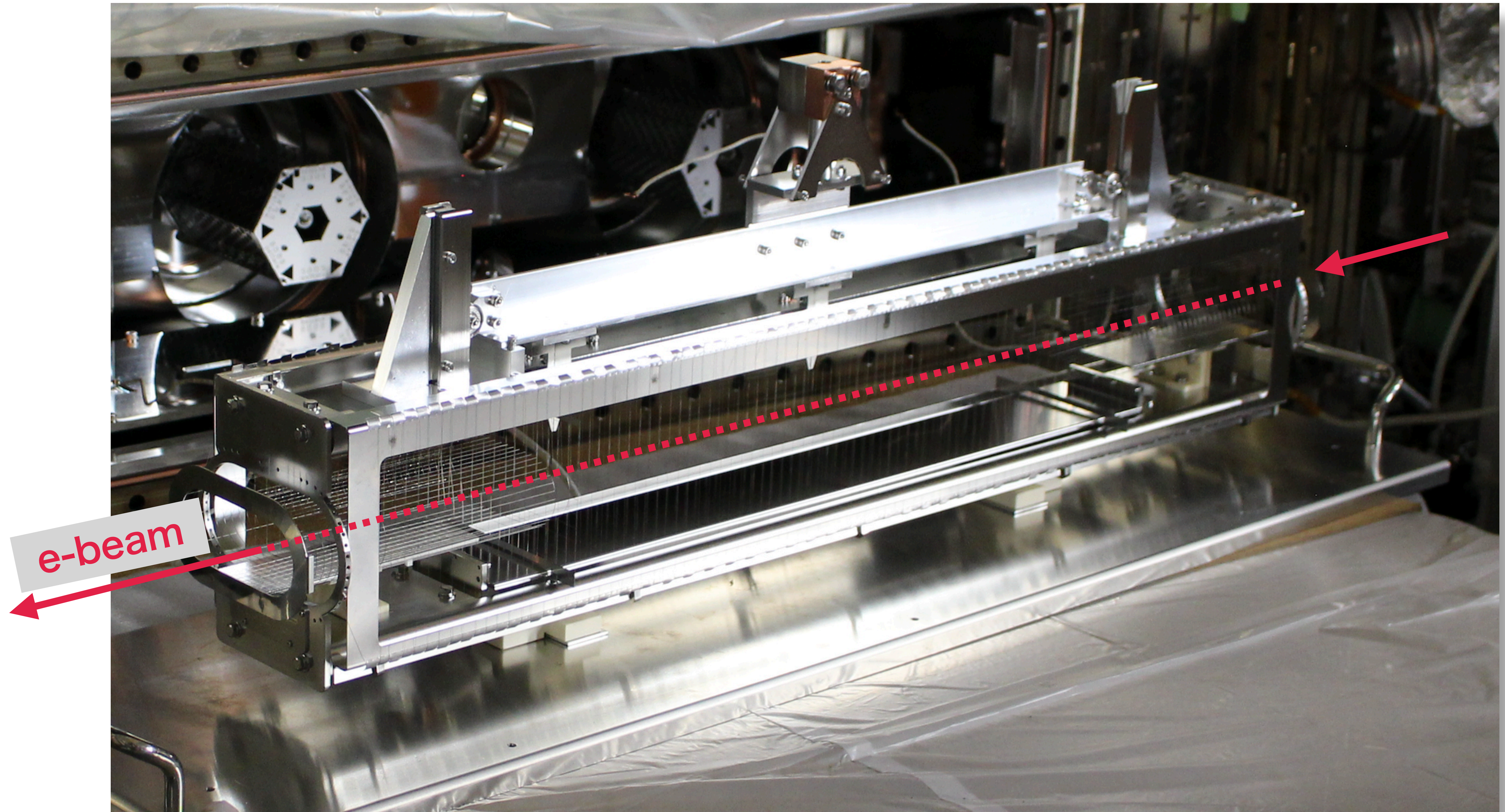
ill problem of e-storage rings



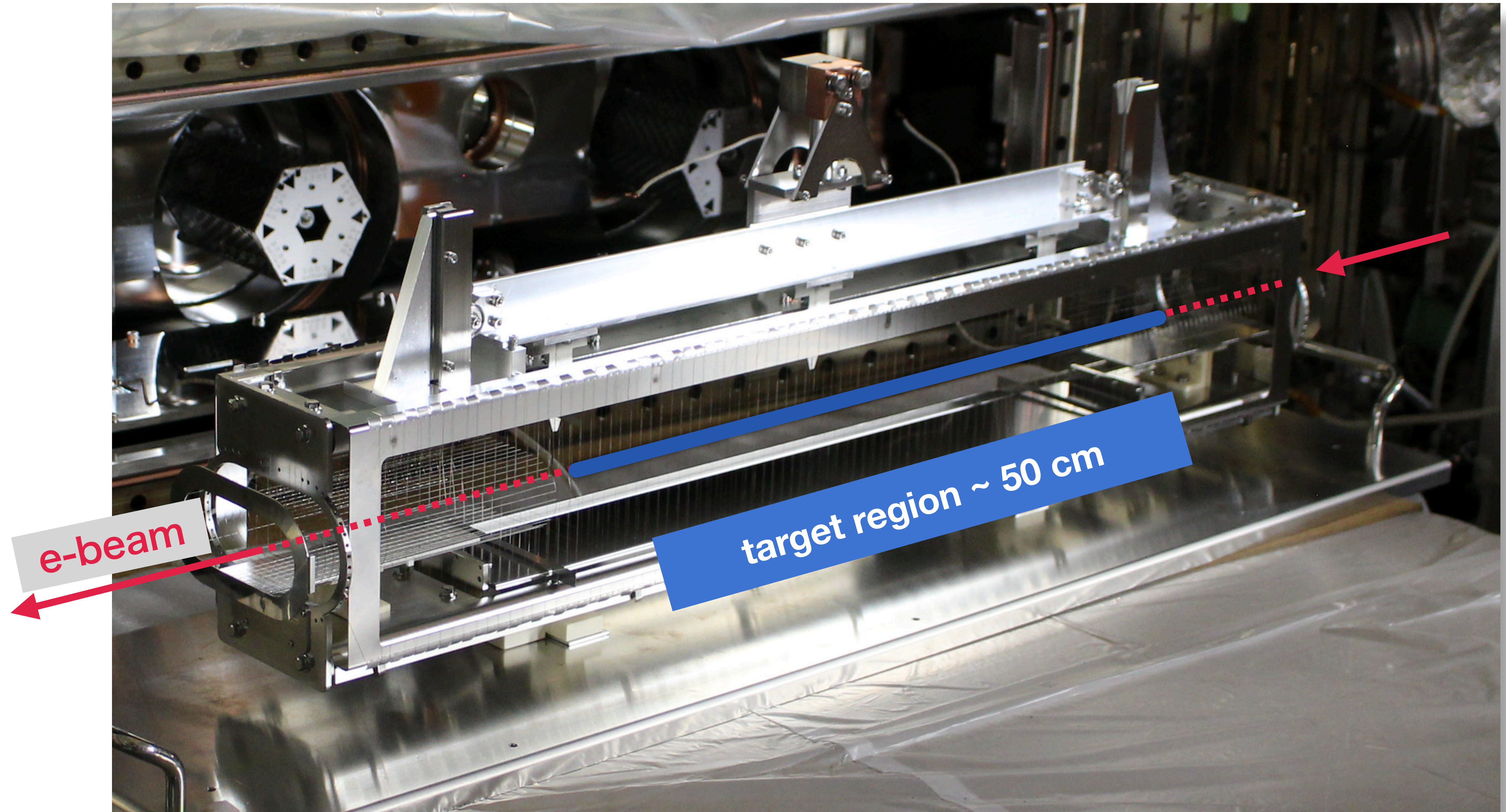
SCRIT device for preparing a target on e-beam



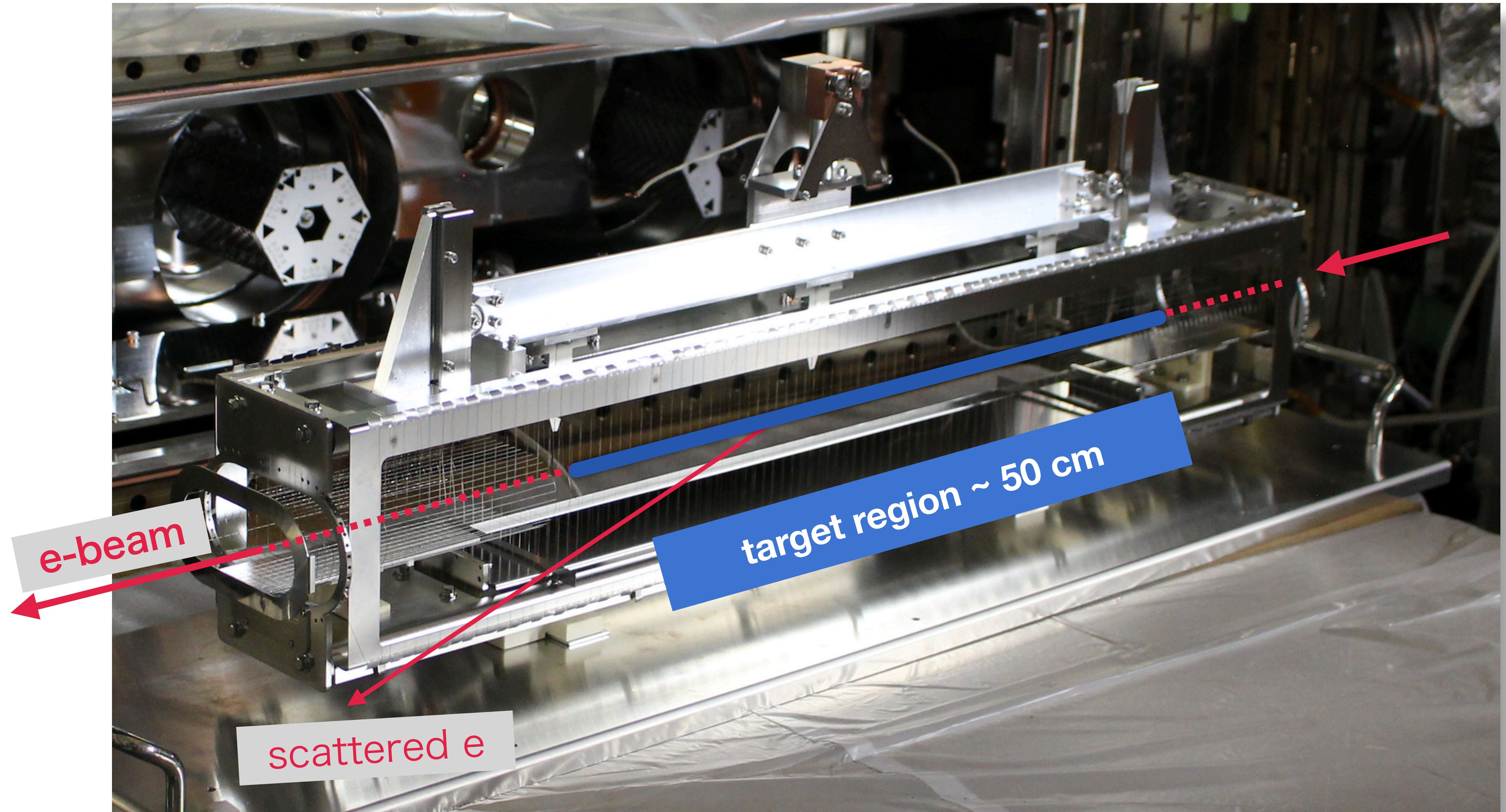
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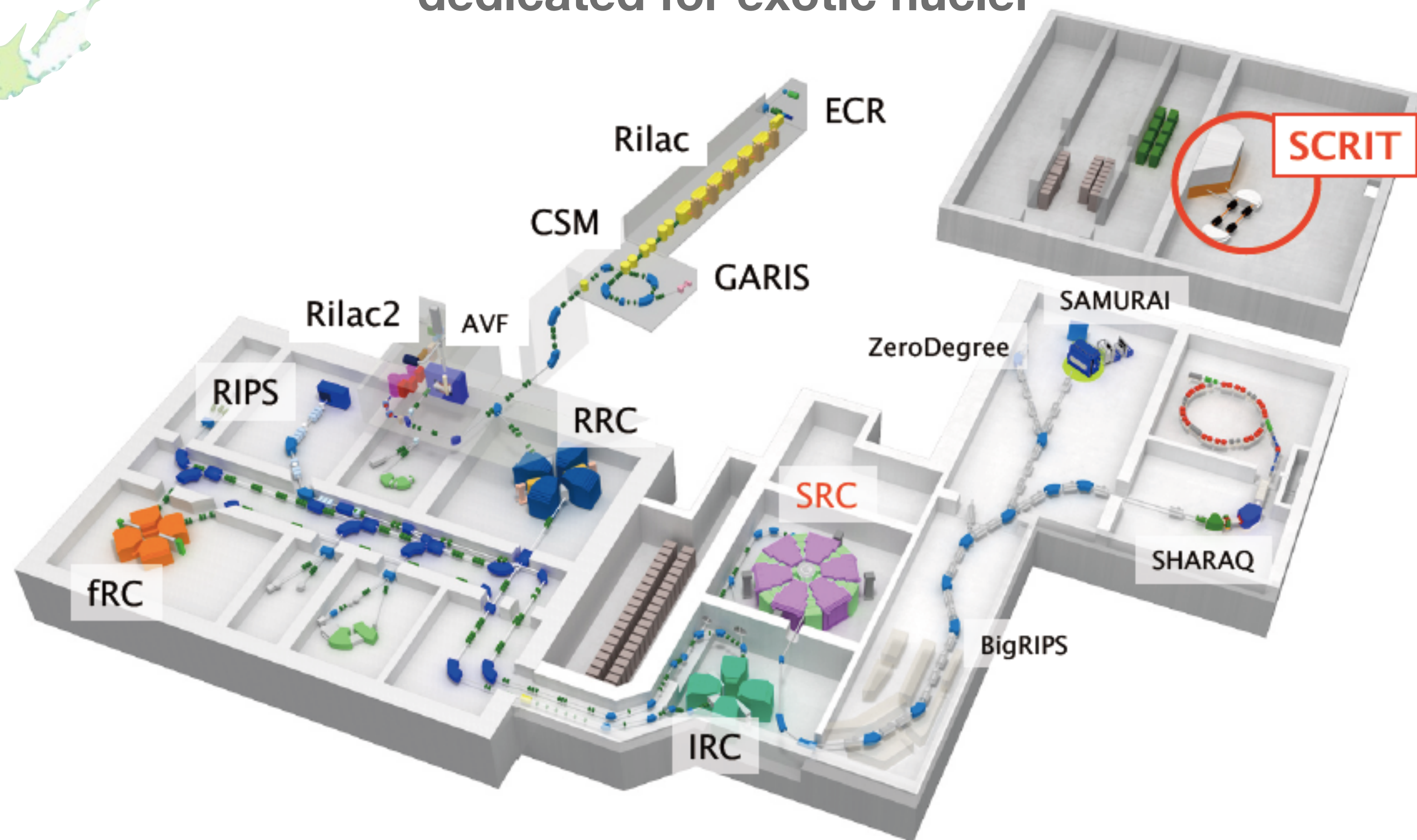
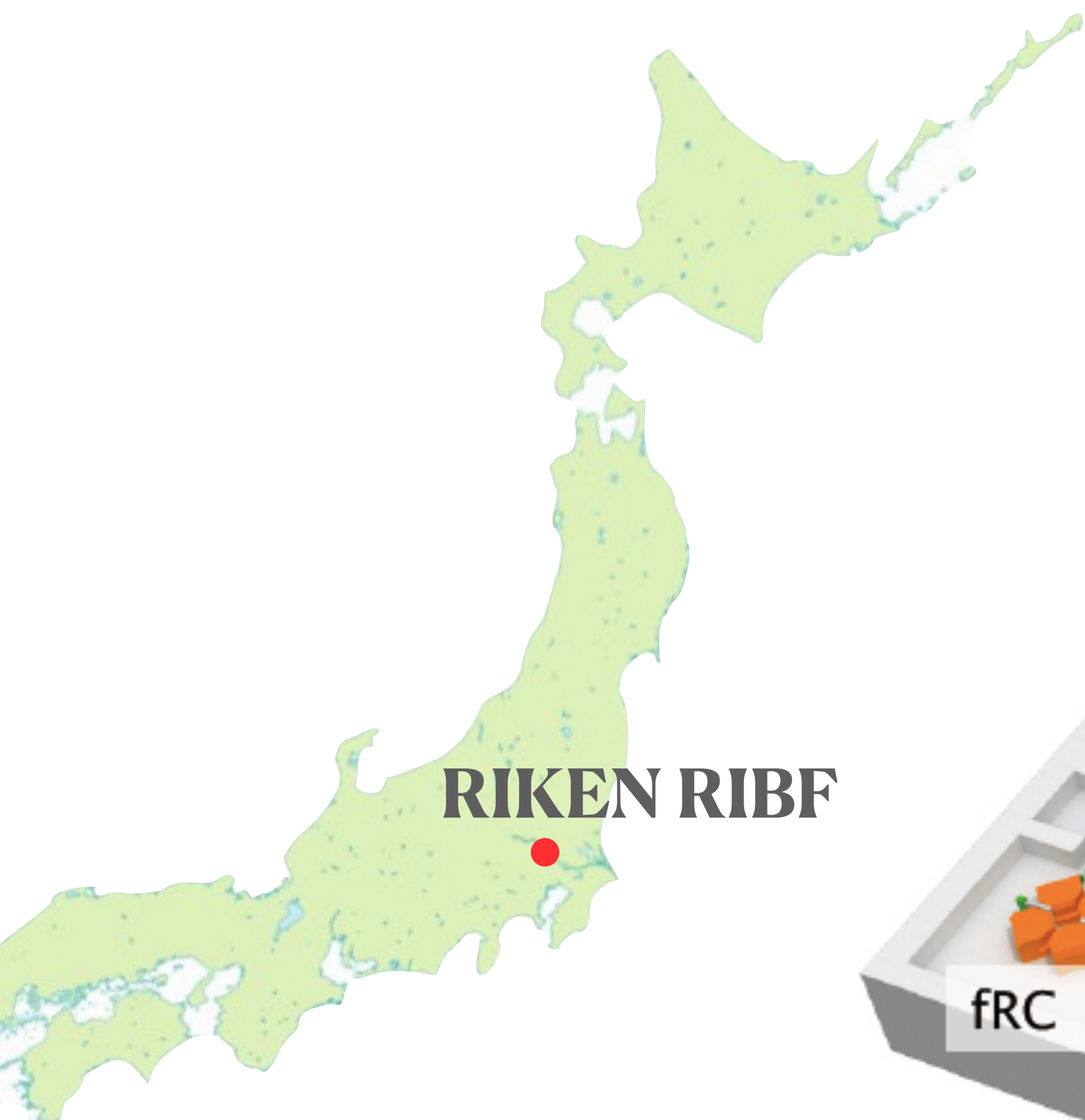


SCRIT device for preparing a target on e-beam



SCRIT electron scattering facility @ RIBF

World's first and currently only one facility
dedicated for exotic nuclei



RIKEN RI Beam Factory (Japan)

RIKEN SCRIT Electron Scattering Facility

Injector + ISOL driver
150 MeV Microtron

FRAC
cooler-buncher
dc-to-pulse conv.

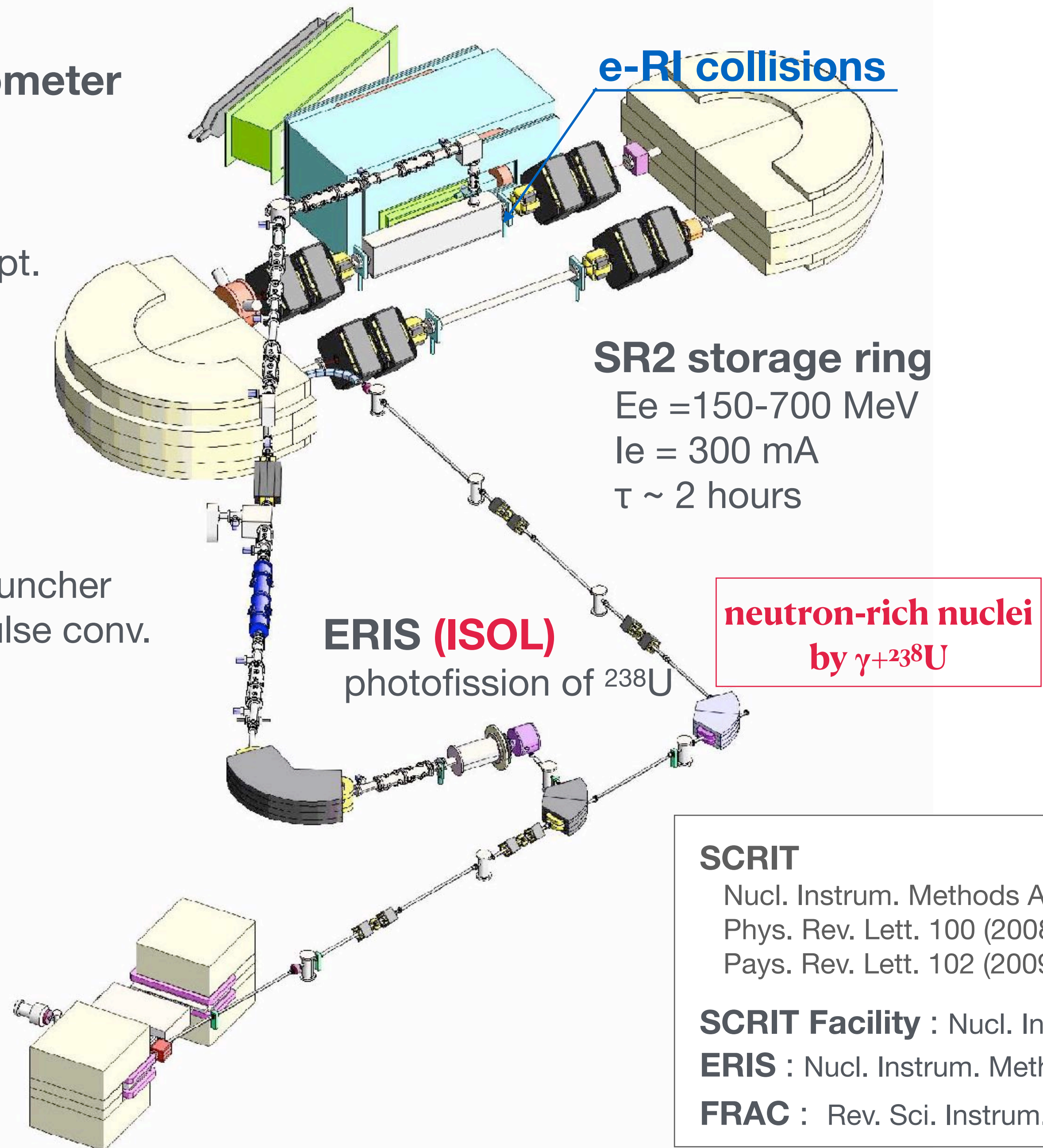
WiSES spectrometer

$\Delta\Omega \sim 90 \text{ mSr}$

$\theta = 30 - 60^\circ$

$\Delta p/p \sim 10^{-3}$

long target accept.



SCRIT

Nucl. Instrum. Methods A532 (2004) 216.

Phys. Rev. Lett. 100 (2008) 164801.

Pays. Rev. Lett. 102 (2009) 102501.

SCRIT Facility : Nucl. Instrum. Method B317 (2013) 668.

ERIS : Nucl. Instrum. Method B317 (2013) 357.

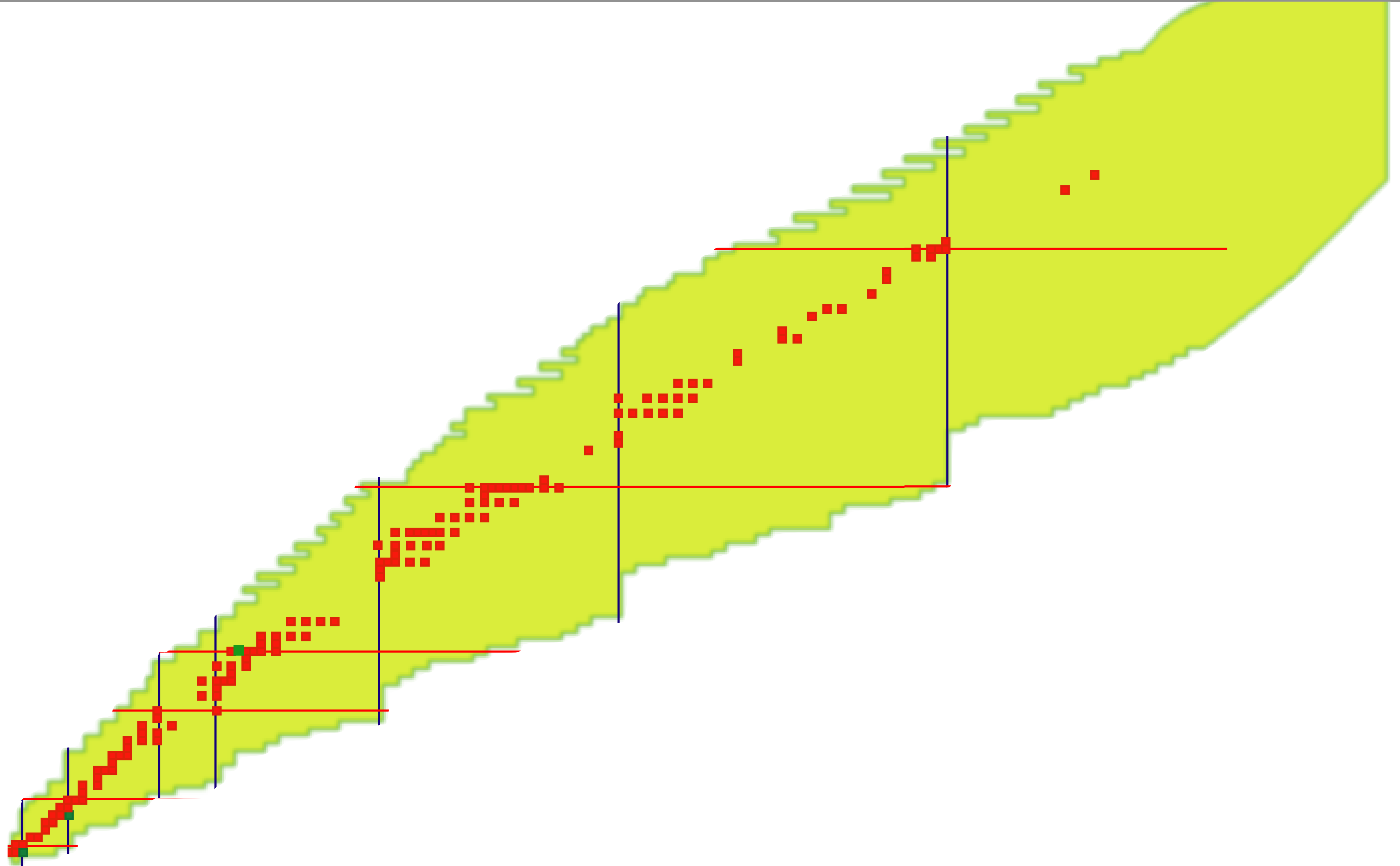
FRAC : Rev. Sci. Instrum. 89 (2018) 095107.

SCRIT electron scattering facility

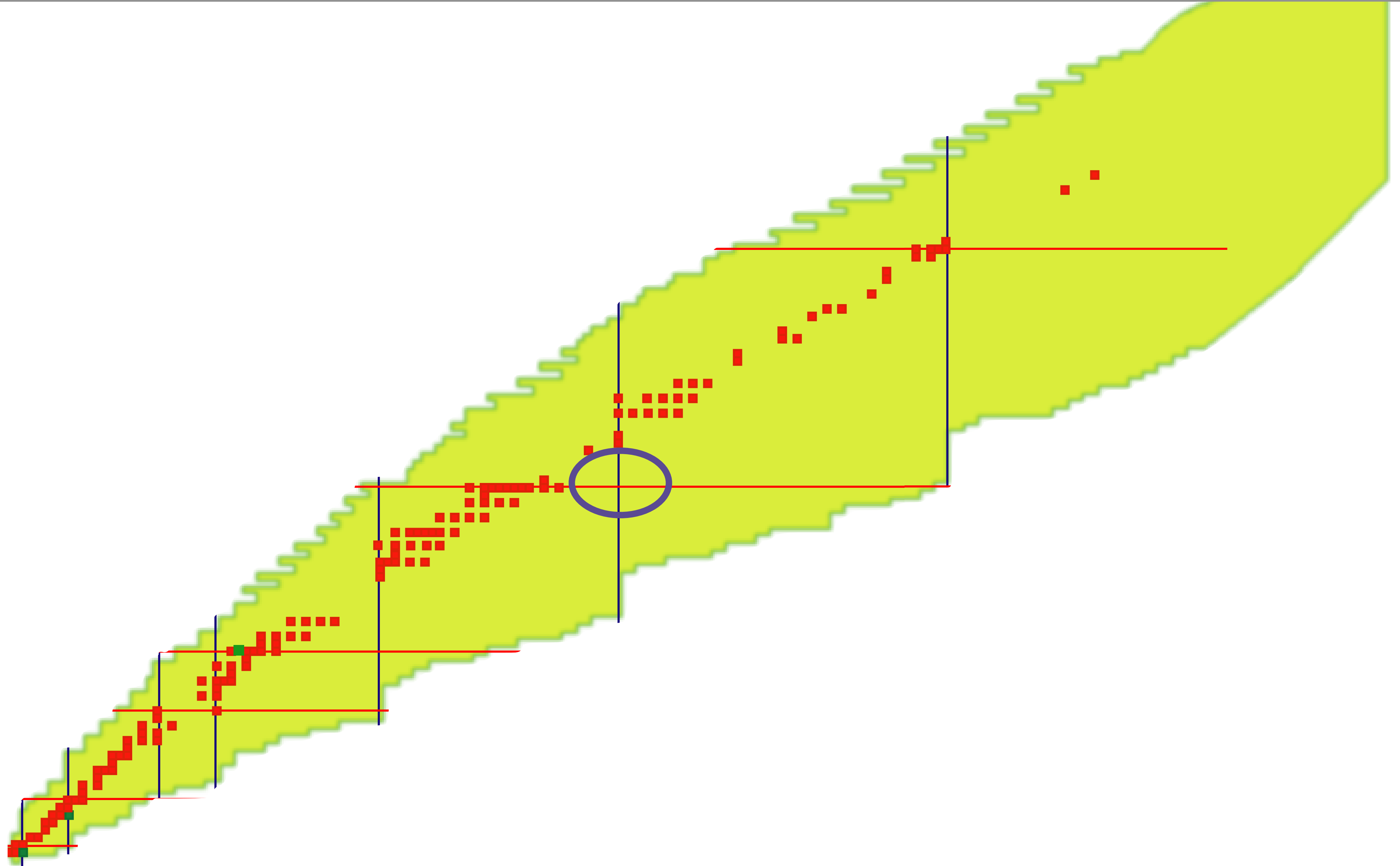
Electron Ring
(SCRIT equipped)

WiSES
(Window-frame Spectrometer
for Electron Scattering)

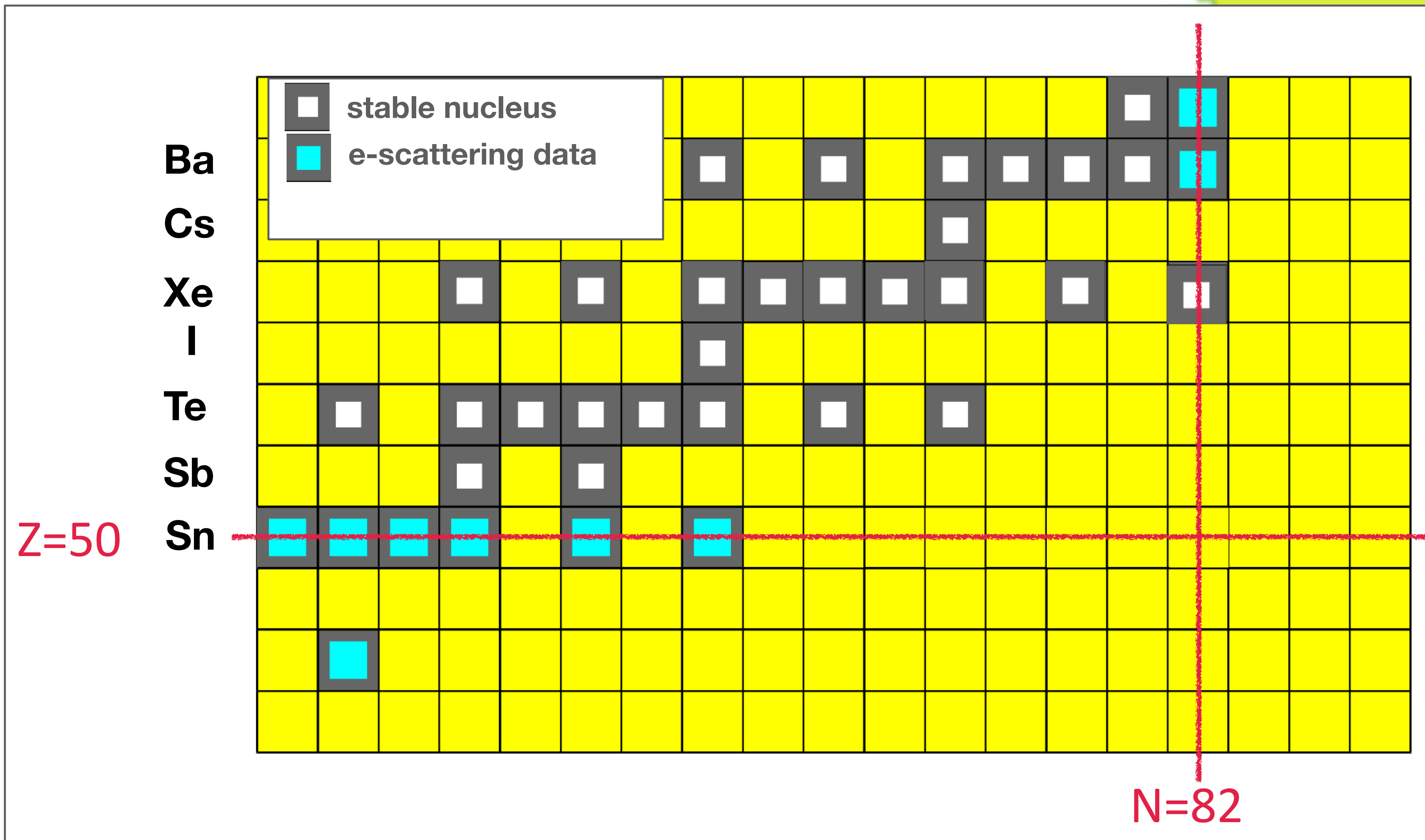
Facility Commissioning with Stable Xe(e,e')



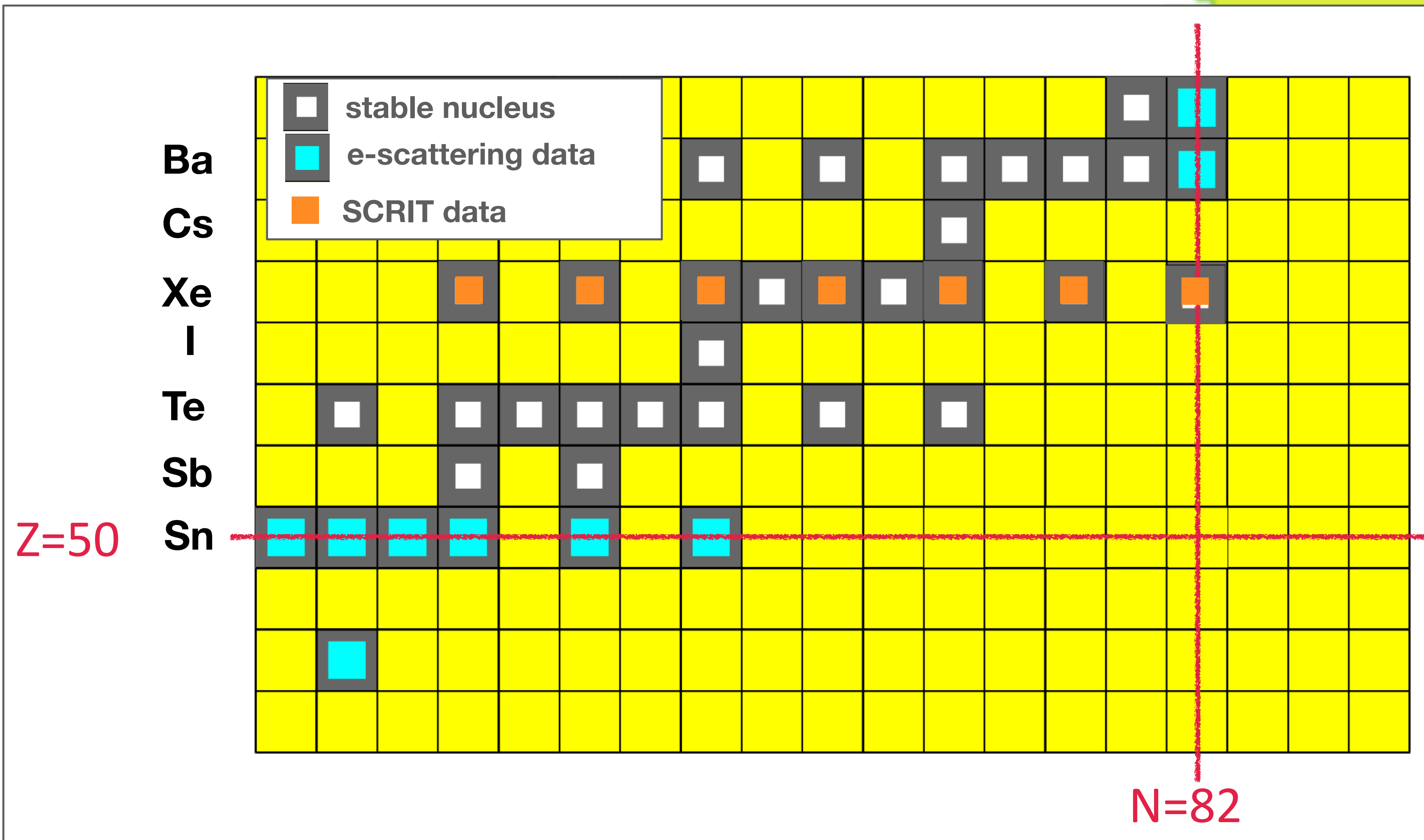
Facility Commissioning with Stable Xe(e,e')



Facility Commissioning with Stable Xe(e,e')

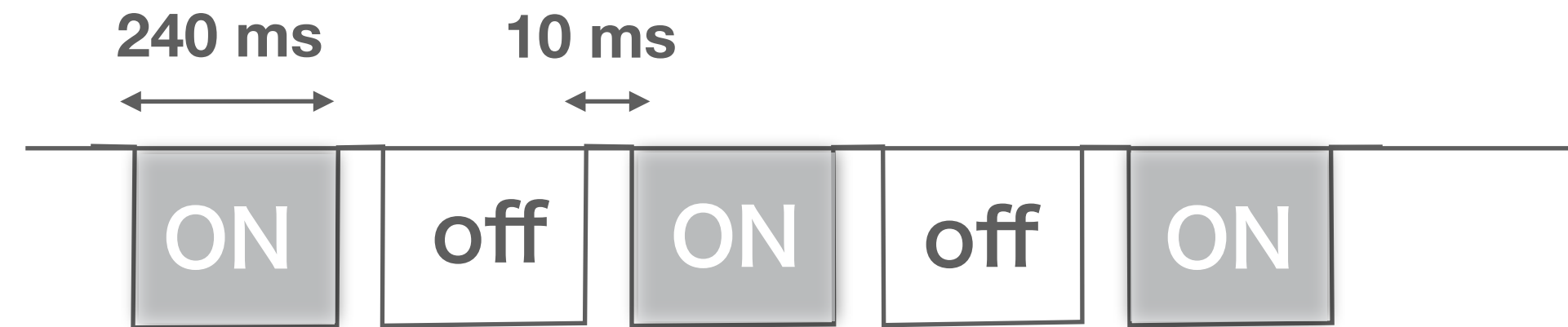


Facility Commissioning with Stable Xe(e,e')



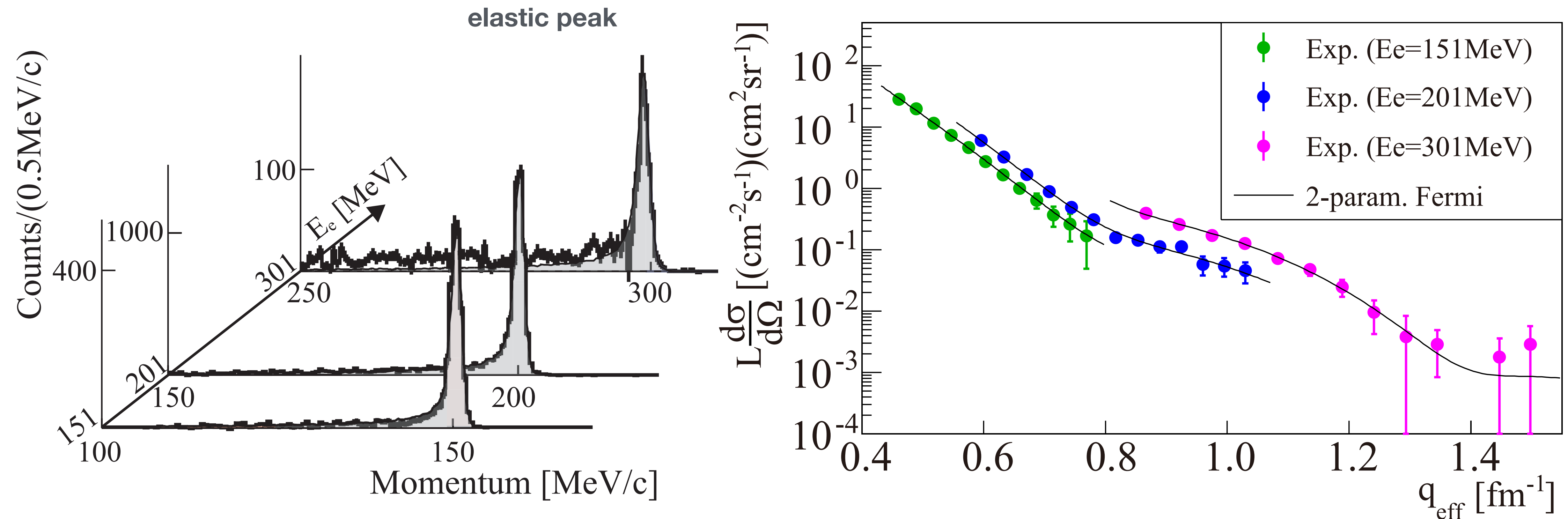
Facility Commissioning with Stable Xe(e,e')

trapping sequence

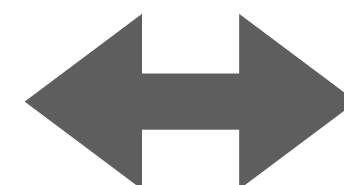


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Phys. Rev. Lett. 118 (2017) 262501.

$L \sim 10^{27} \text{ /cm}^2\text{/s}$ with $N_{\text{trapped}} \sim 10^7$

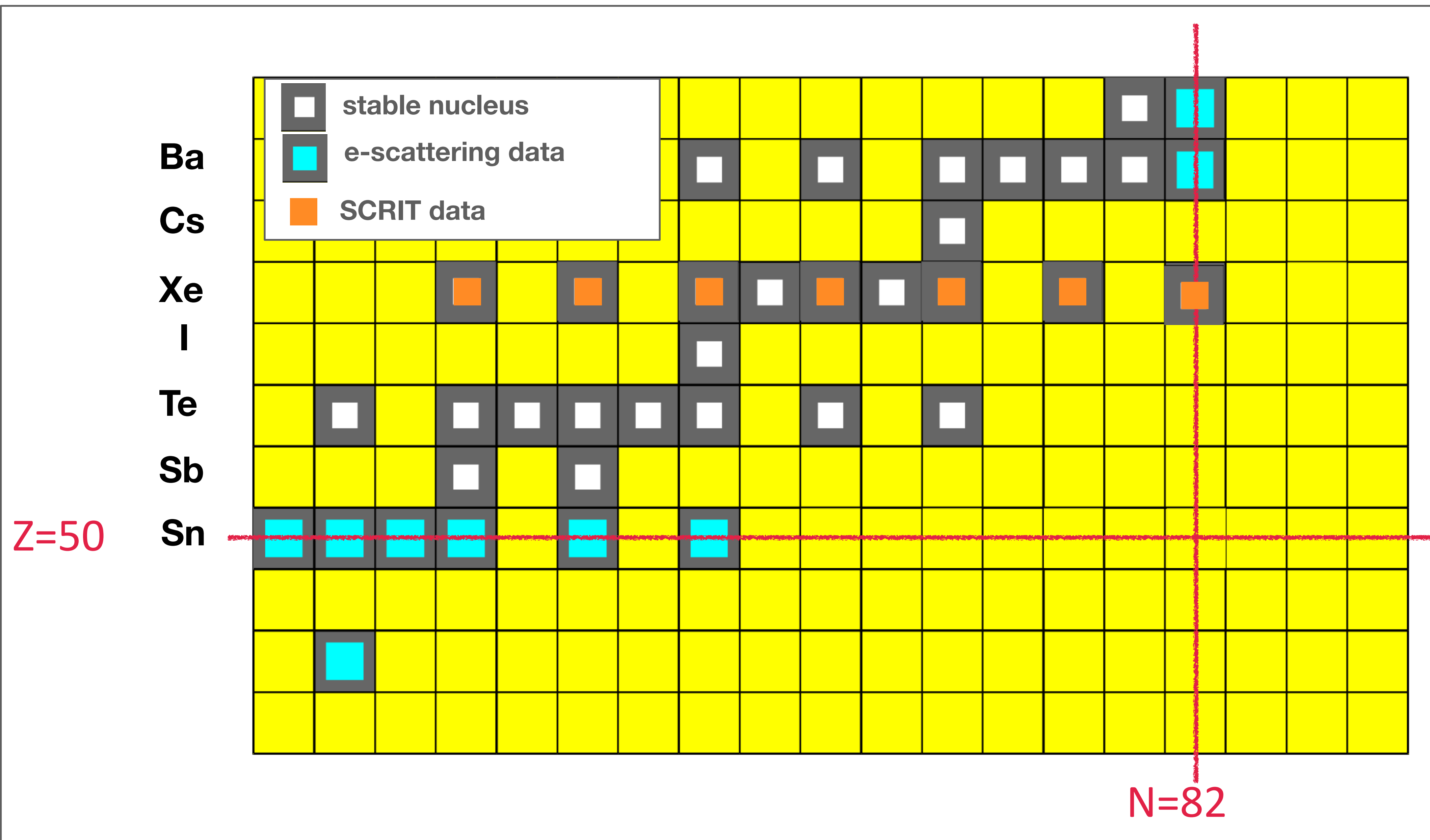


$$\langle r^2 \rangle^{1/2} = 4.789^{+0.12}_{-0.10} \text{ fm}$$

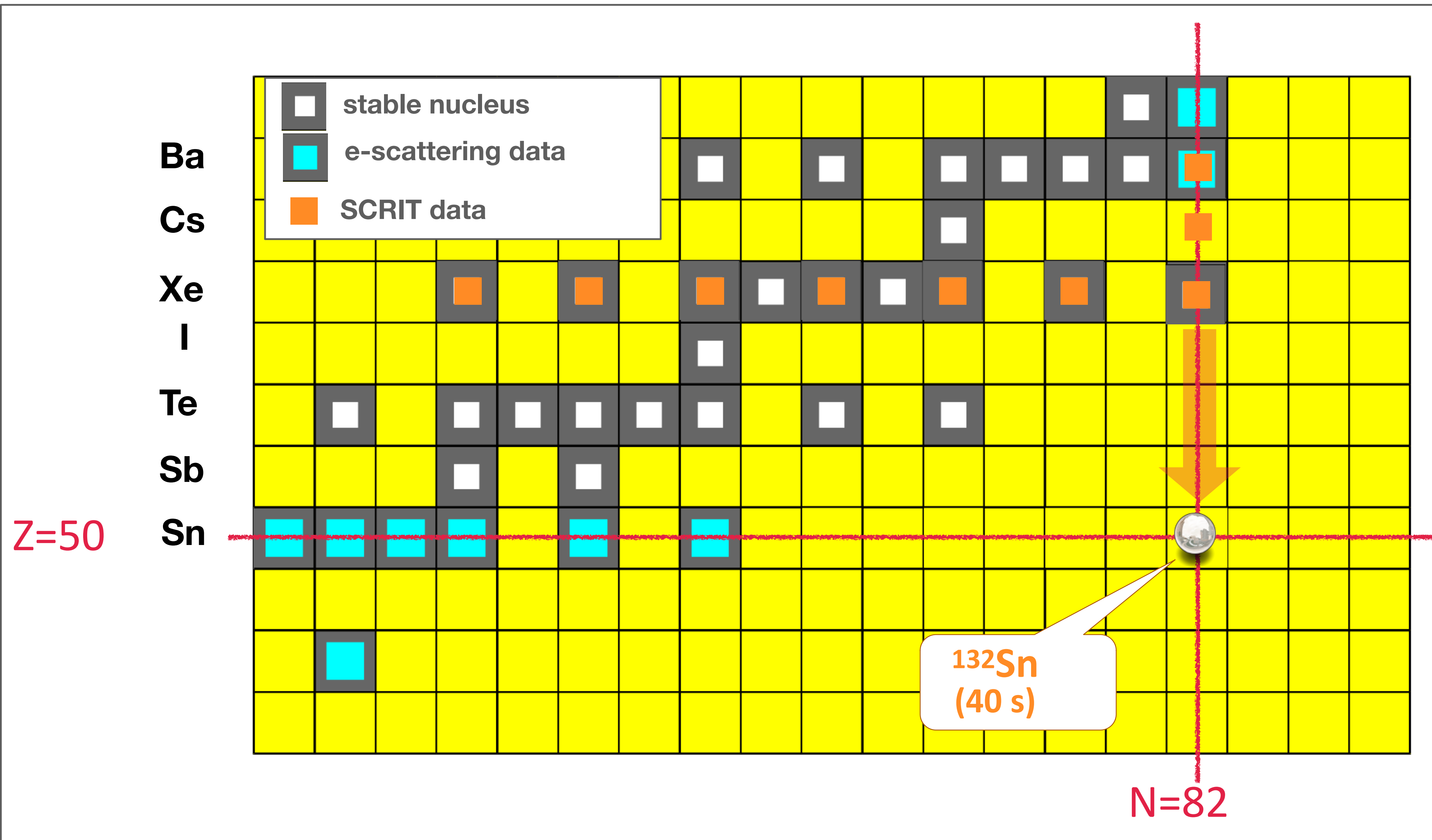


$$\langle r^2 \rangle^{1/2} = 4.787 \text{ fm (}\mu\text{-atom X-rays)}$$

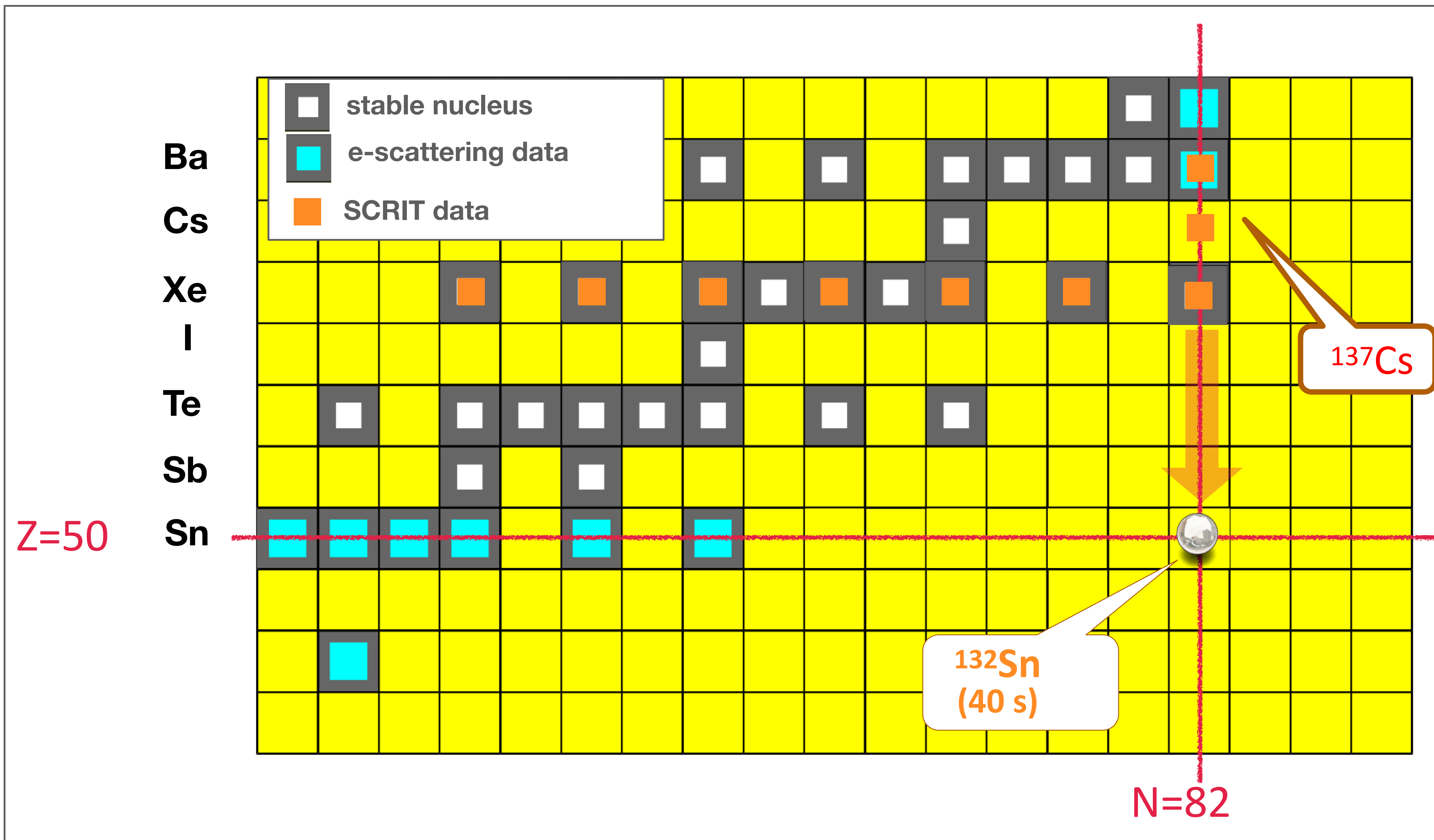
Towards the doubly magic nucleus ^{132}Sn



Towards the doubly magic nucleus ^{132}Sn



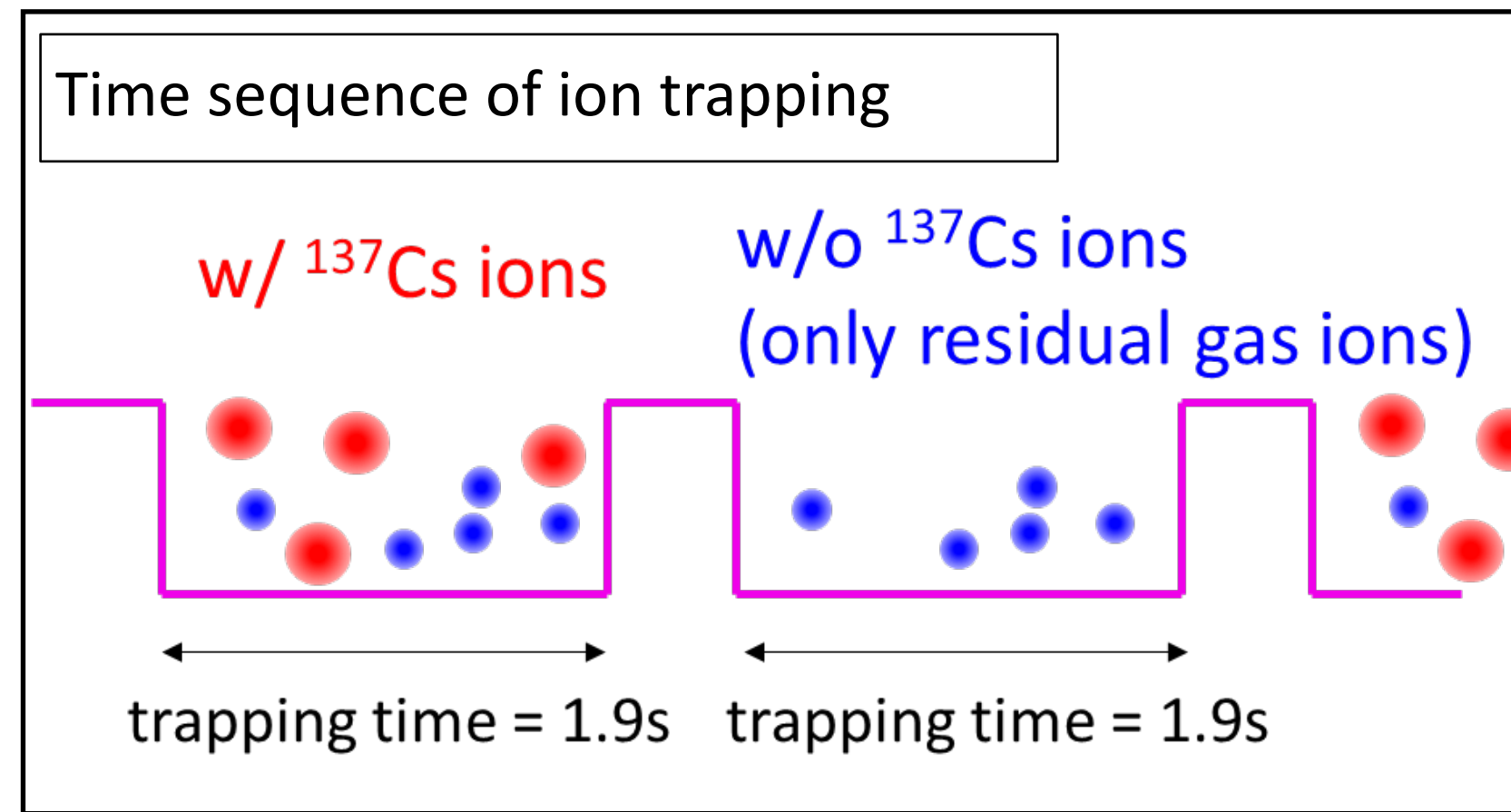
Towards the doubly magic nucleus ^{132}Sn



World's first (e,e') for Online-Produced Unstable Nuclide (^{137}Cs)

1.9 s trapping

=> mimicking “short-lived” nuclei



$N_{\text{trapped}} \sim 2 \times 10^7$

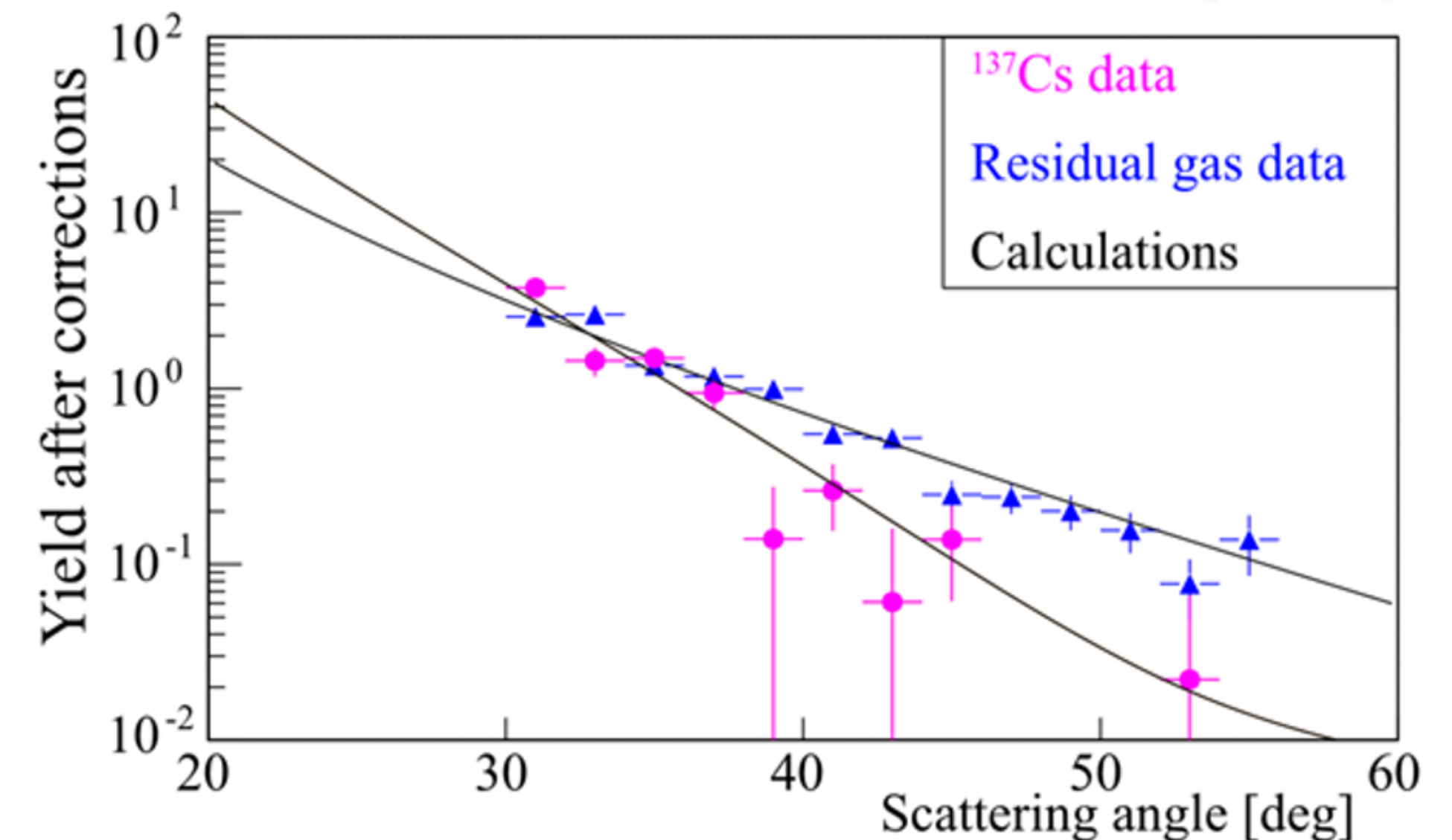
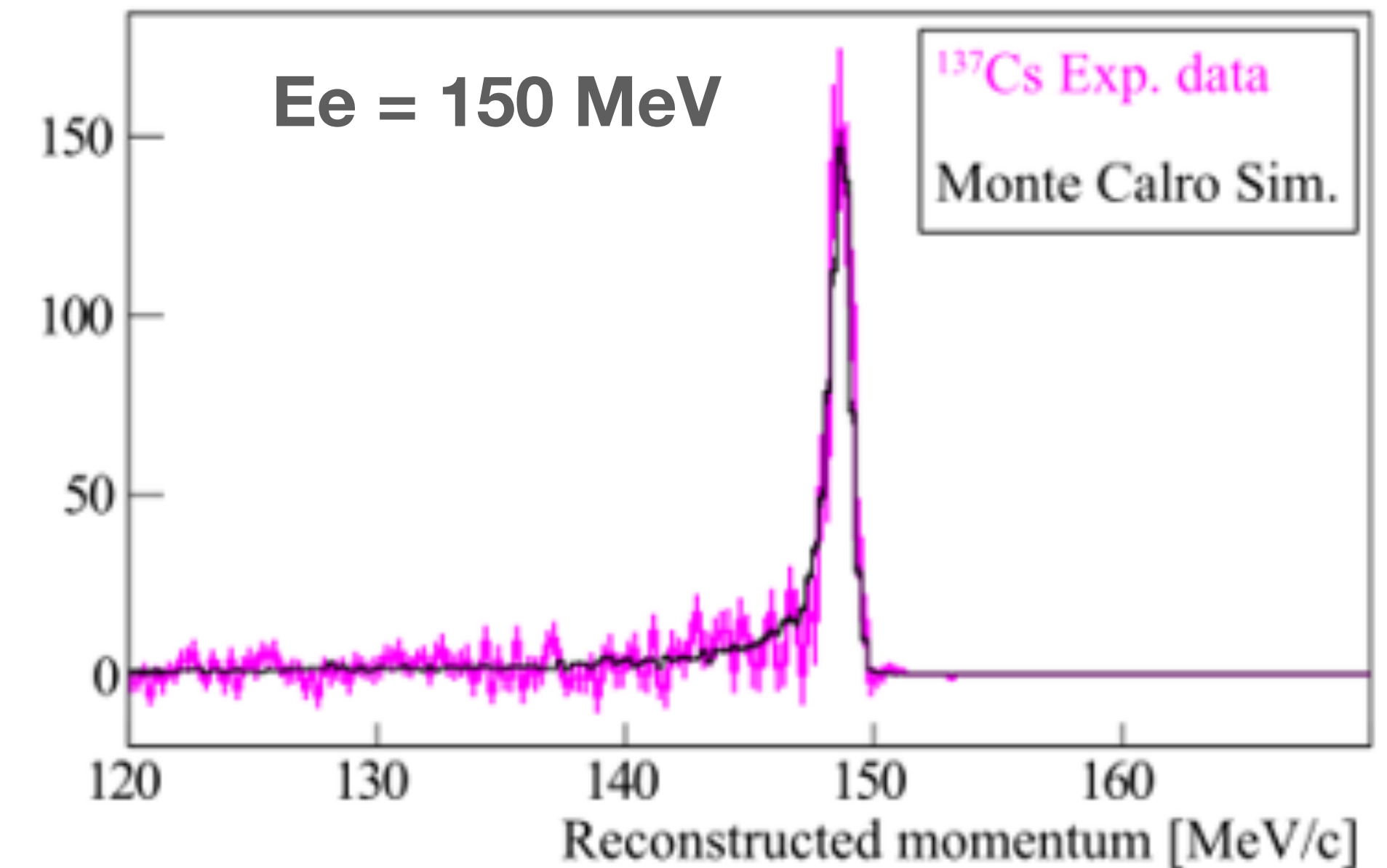
=> $L \sim 0.9 \times 10^{26} \text{ /cm}^2\text{/s}$

**Successful demonstration with
online-produced unstable nuclei**

=> further details

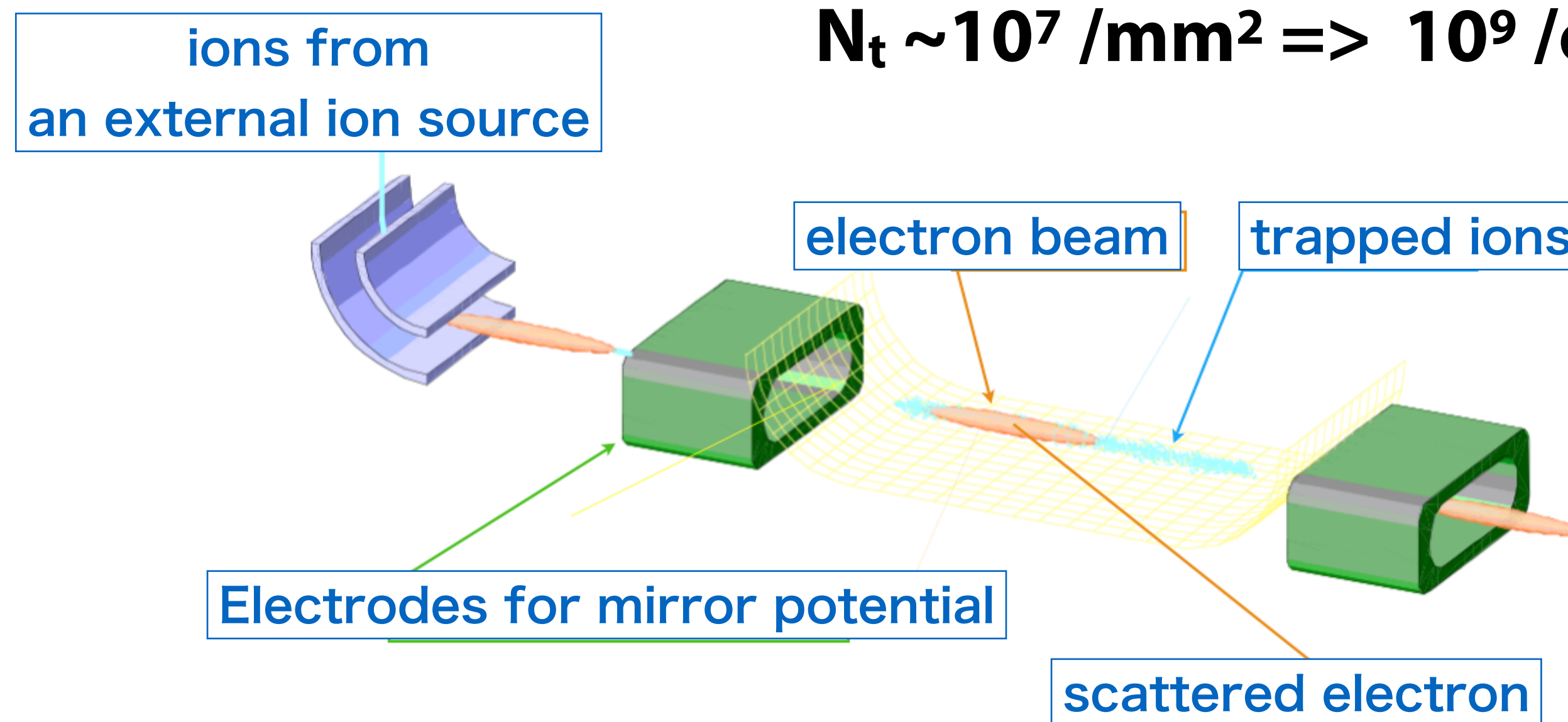
Kyo Tsukada's talk (Thursday, May 29, Room9, 16:55 -)

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The SCRIT concept and Luminosity

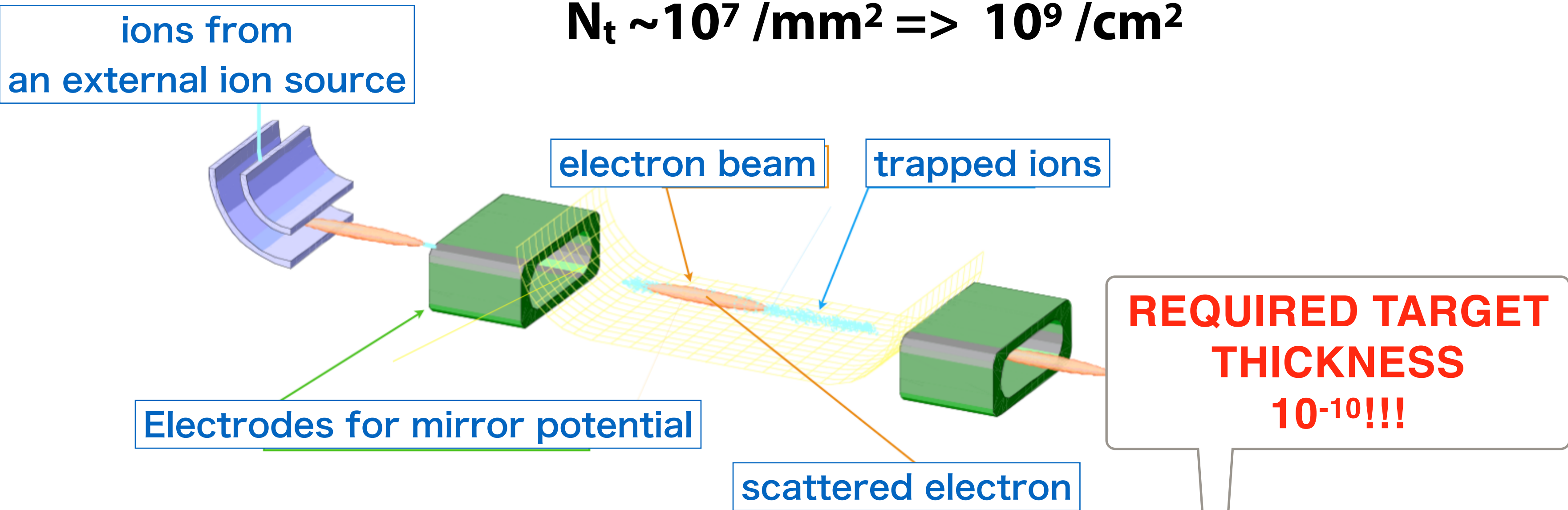
$\sim 10^7$ ions are trapped on e-beam ($\sim 1 \text{ mm}^2$)



	Ee	N_{beam}	$\rho \cdot t$	L
Hofstadter's era (1950s)	150 MeV	$\sim 1 \text{ nA}$ ($\sim 10^9 / \text{s}$)	$\sim 10^{19} / \text{cm}^2$	$\sim 10^{28} / \text{cm}^2 / \text{s}$
JLAB	12 GeV	$\sim 100 \mu\text{A}$ ($\sim 10^{14} / \text{s}$)	$\sim 10^{22} / \text{cm}^2$	$\sim 10^{36} / \text{cm}^2 / \text{s}$
SCRIT	150 - 300 MeV	$\sim 200 \text{ mA}$ ($\sim 10^{18} / \text{s}$)	$\sim 10^9 / \text{cm}^2$	$\sim 10^{27} / \text{cm}^2 / \text{s}$

The SCRIT concept and Luminosity

$\sim 10^7$ ions are trapped on e-beam ($\sim 1 \text{ mm}^2$)



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SCRIT	150 - 300 MeV	$\sim 200 \text{ mA}$ ($\sim 10^{18} / \text{s}$)	$\sim 10^9 / \text{cm}^2$	$\sim 10^{27} / \text{cm}^2 / \text{s}$

ISOL upgrade: towards ^{132}Sn

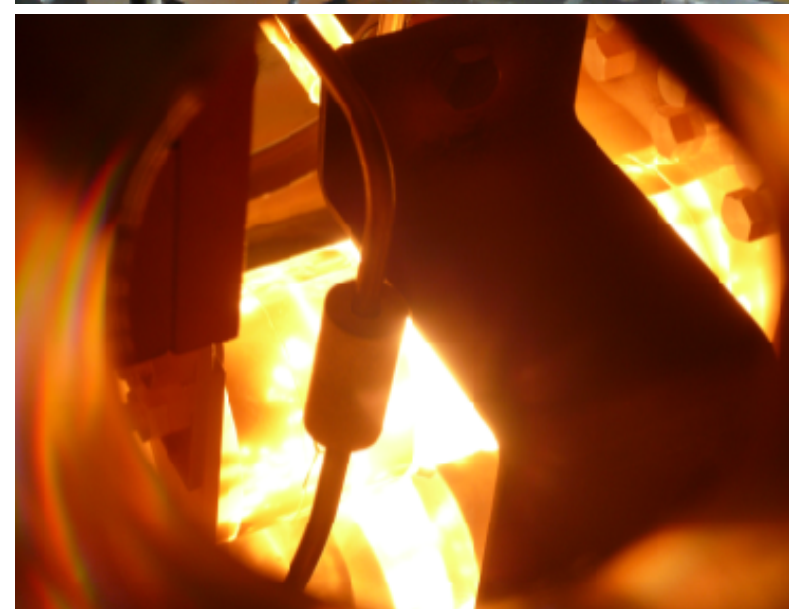
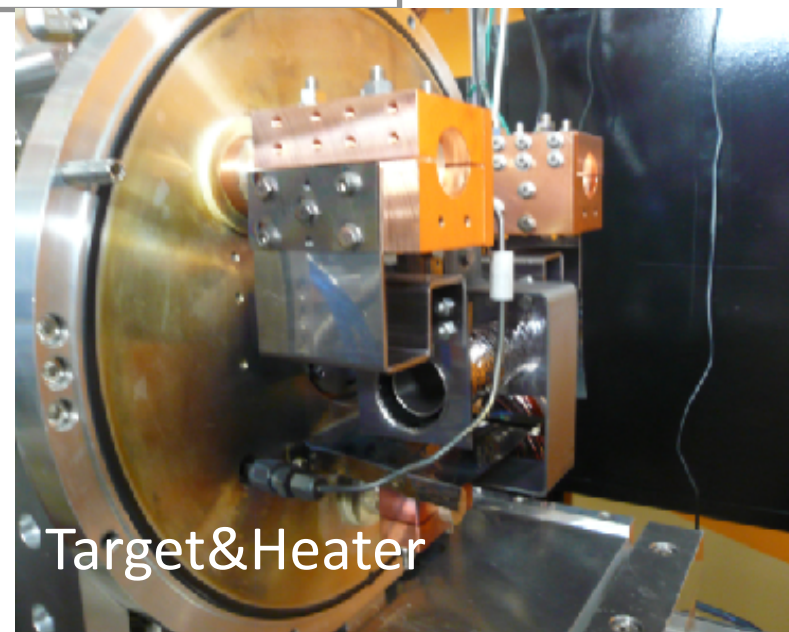
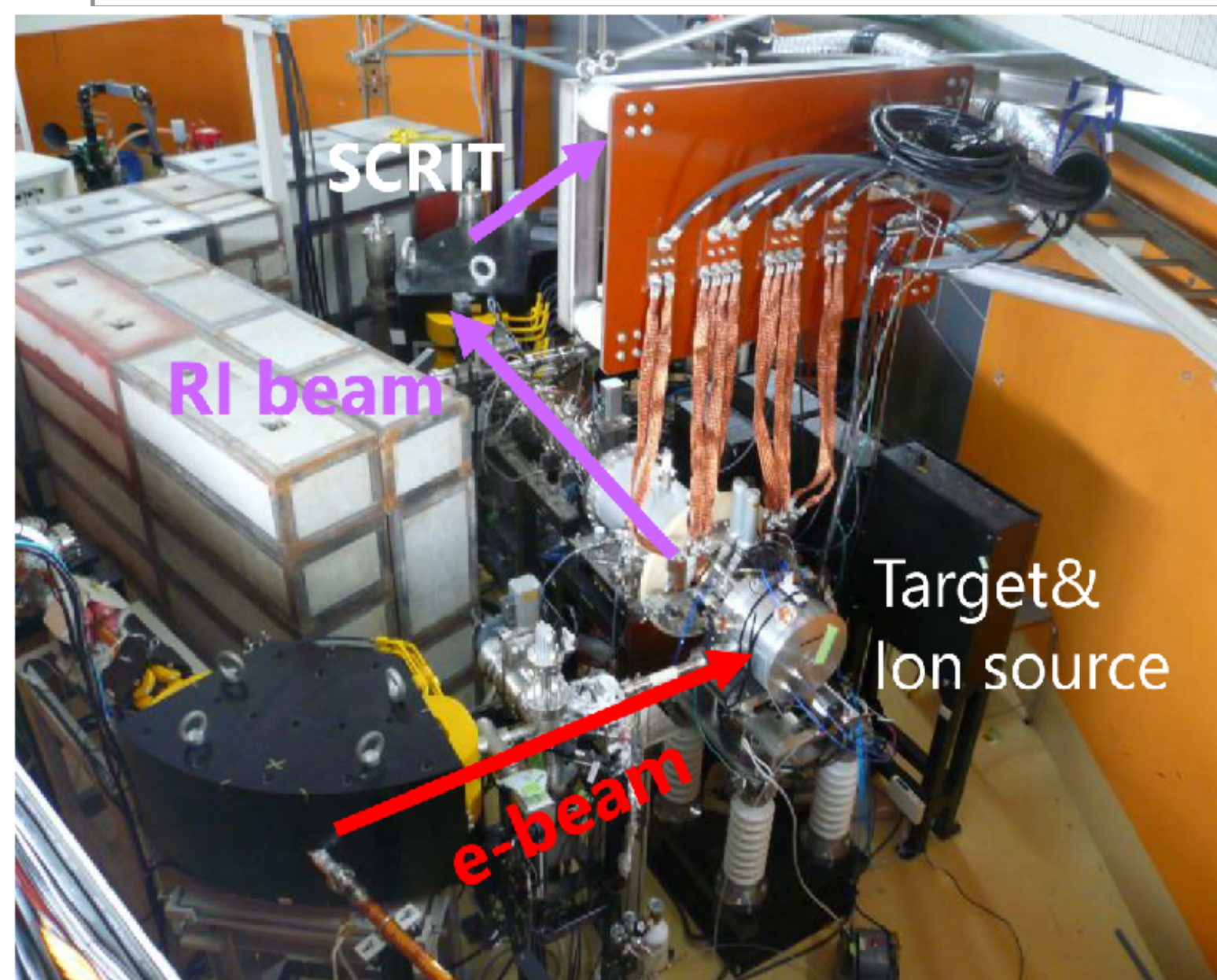
Production Rate

$$N_{\text{fission}} \sim 10^8 \text{ /watt}$$

$$N^{132}\text{Sn} \sim 10^6 \text{ /watt} * 1\% (\epsilon_{\text{trans.}})$$

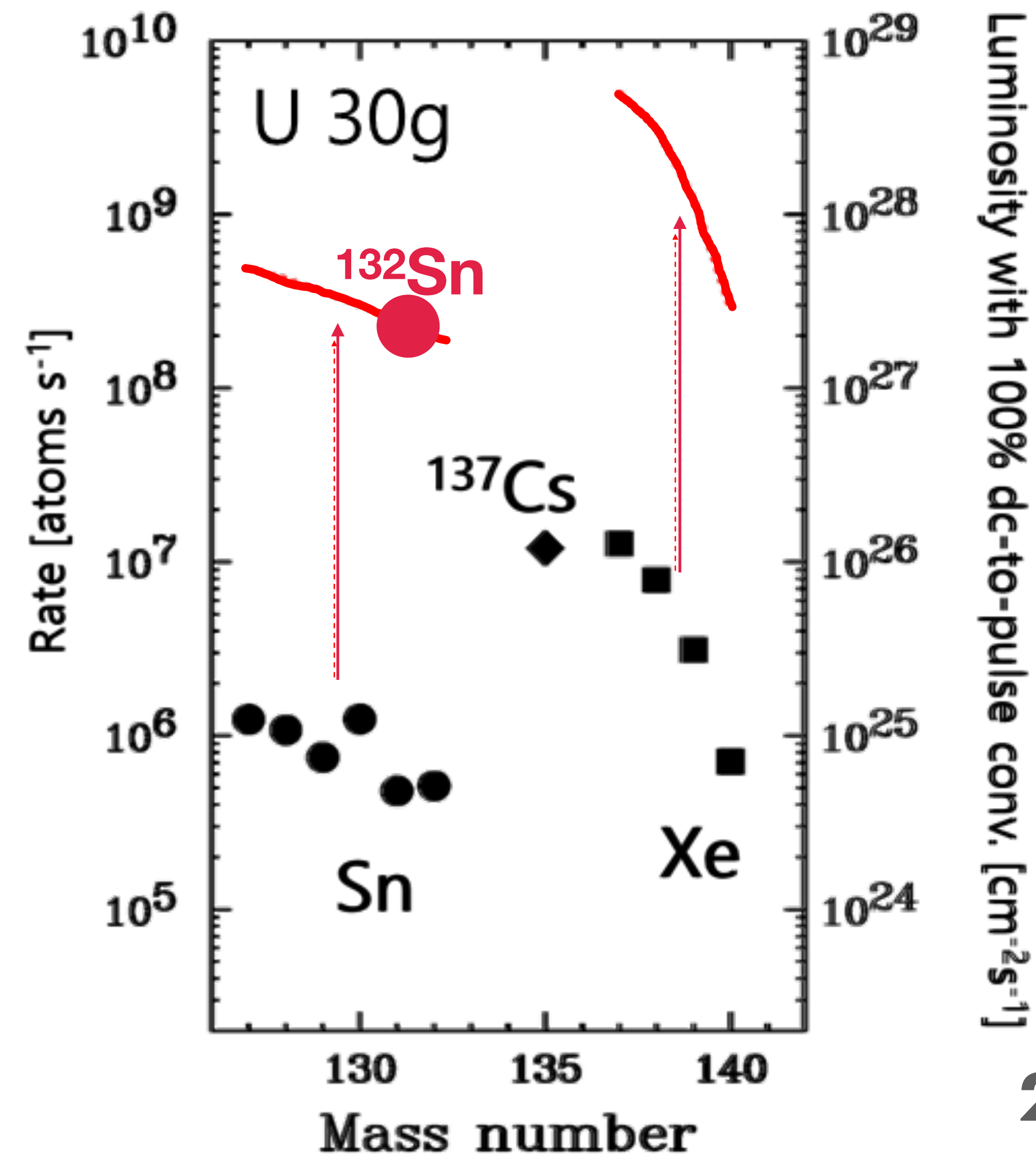
beam power : $\sim 20\text{W}$ (today)

$\sim 2 \text{ kW}$ (underway)



House-made
Uranium carbide (UCx)

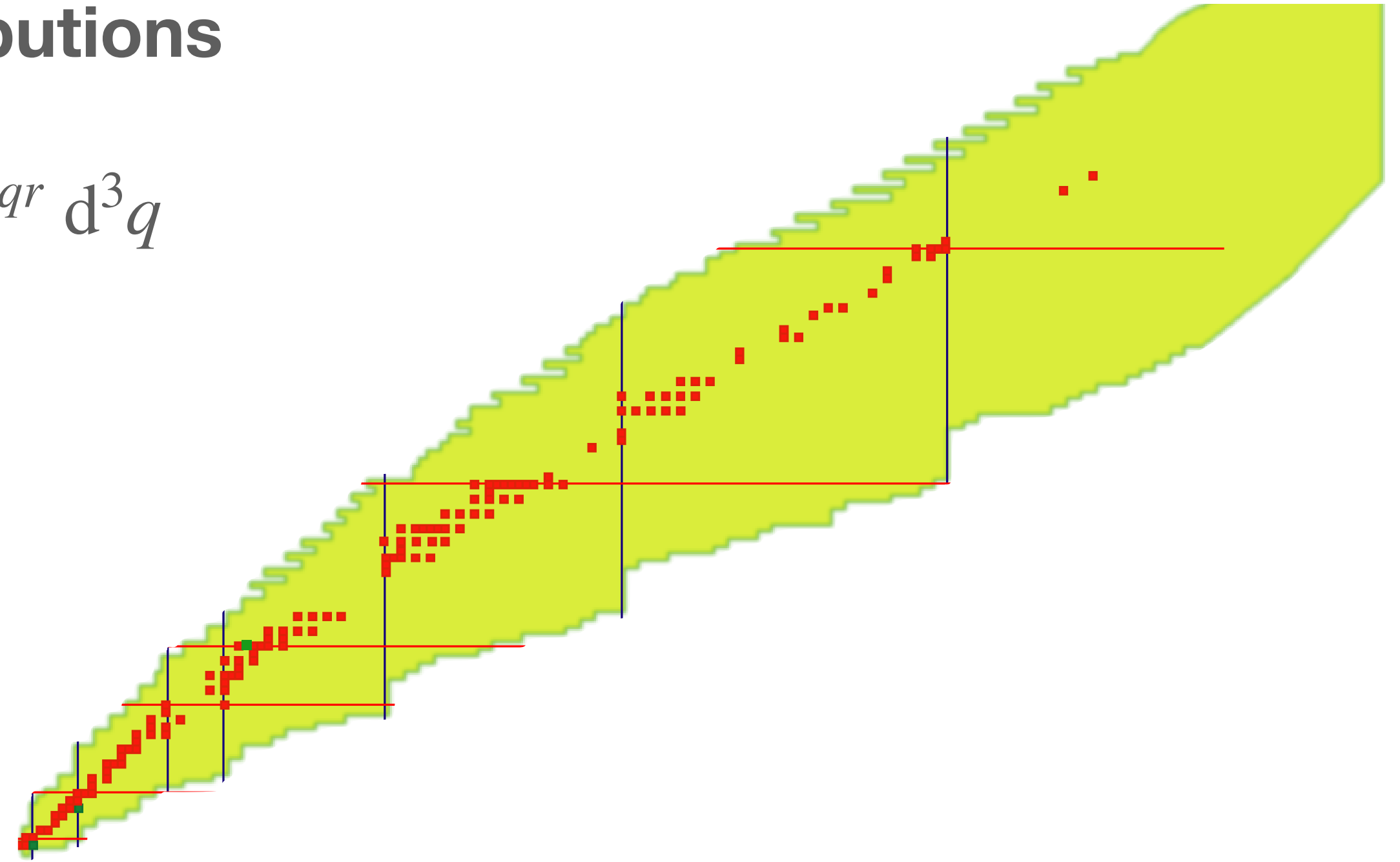
towards 2 kW e-beam
higher repetition
higher peak intensity



Research possibilities at the SCRIT facility

• **elastic scattering** ➔ charge density distributions

$$\frac{d\sigma}{d\Omega} = \frac{d\sigma_{Mott}}{d\Omega} |F_c(q)|^2 \quad \rho_c(r) = \int F_c(q) e^{-iqr} d^3q$$



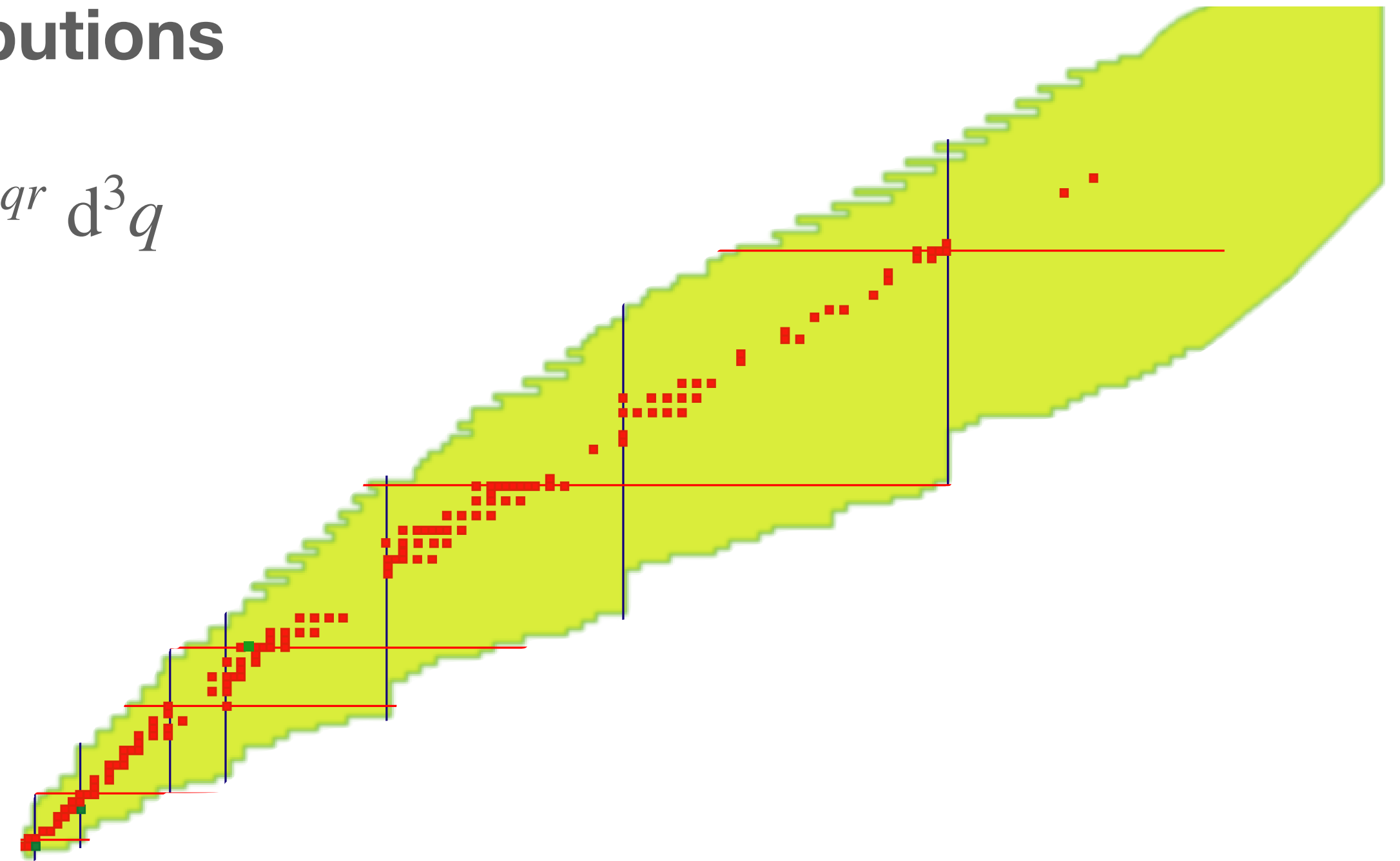
Research possibilities at the SCRIT facility

elastic scattering → charge density distributions

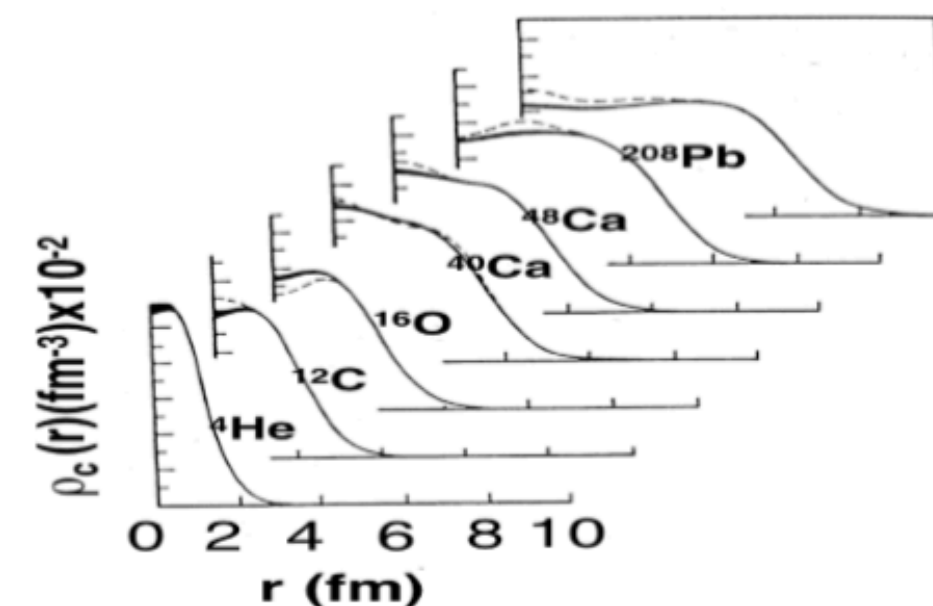
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beyond elastic scattering

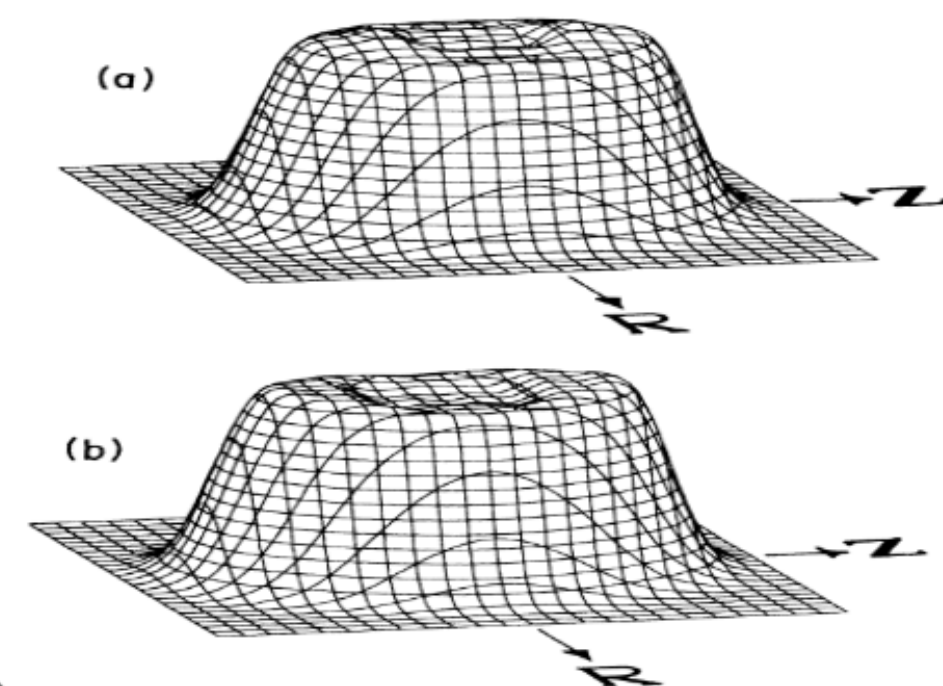
10 ²⁶	10 ²⁷	10 ²⁸	10 ²⁹	10 ³⁰	10 ³¹	10 ³²
						(/cm ² /s)
elastic scattering	inelastic scattering	Giant Resonance	quasi-elastic (e,e'p)	magnetic scattering		



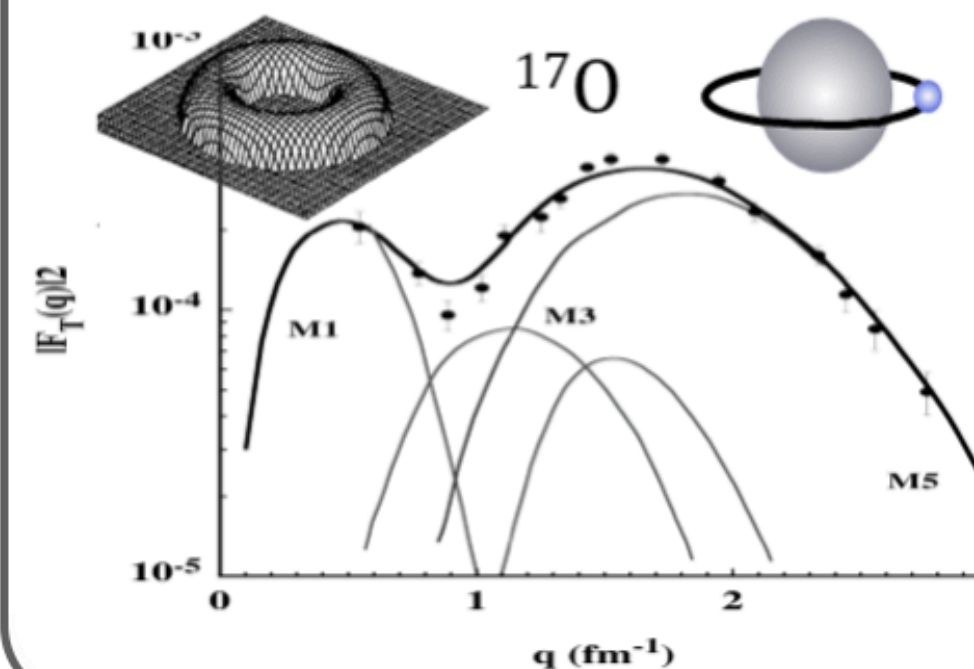
elastic scattering
→ charge distribution



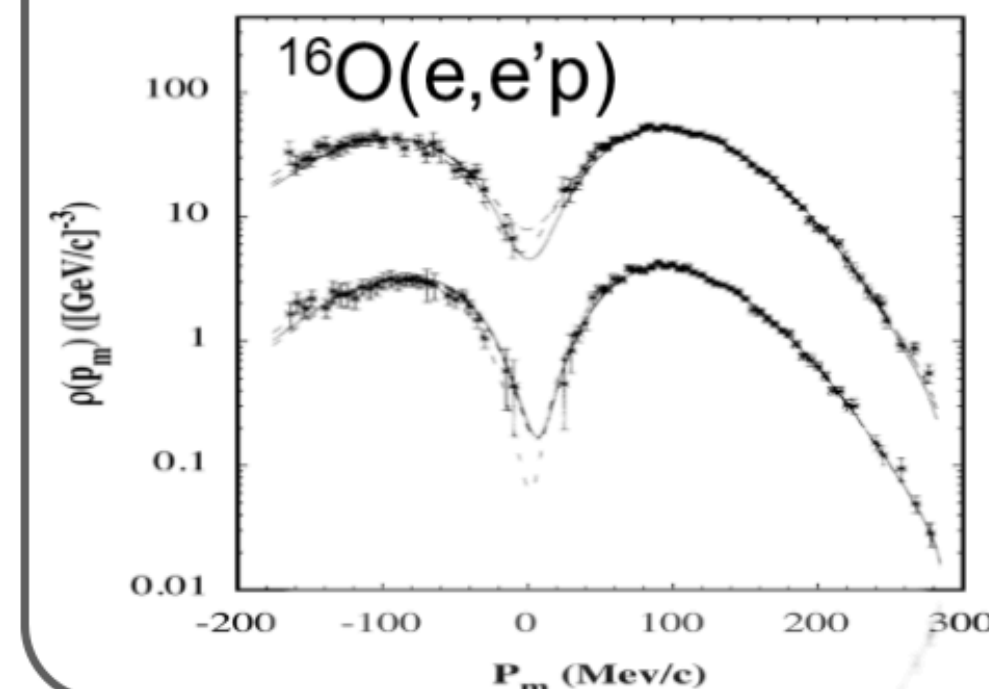
elastic + inelastic scattering
→ deformation



magnetic scattering
→ valence neutron



quasi-elastic (e,e'p)
→ S-factor



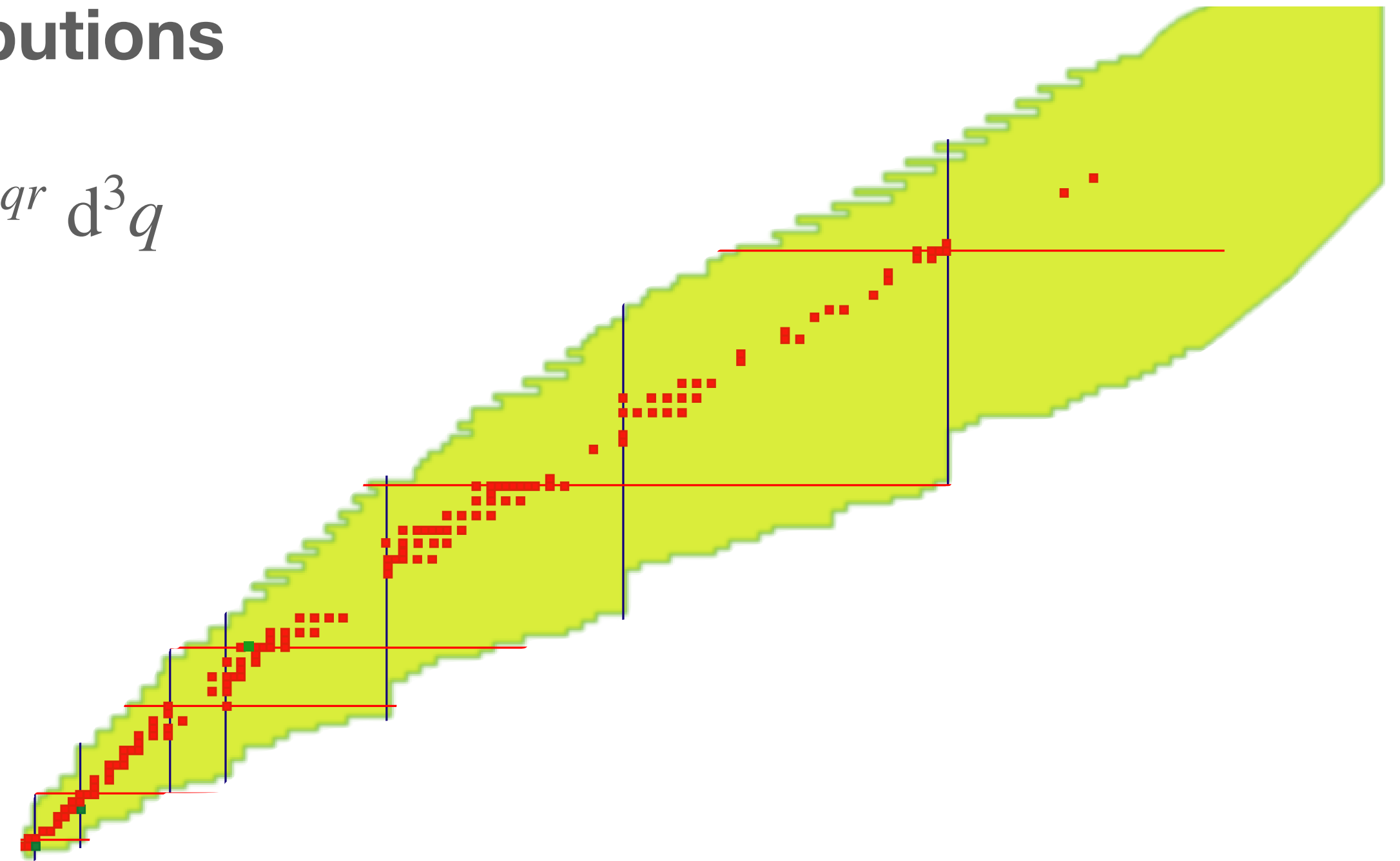
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- beyond elastic scattering

10 ²⁶	10 ²⁷	10 ²⁸	10 ²⁹	10 ³⁰	10 ³¹	10 ³²
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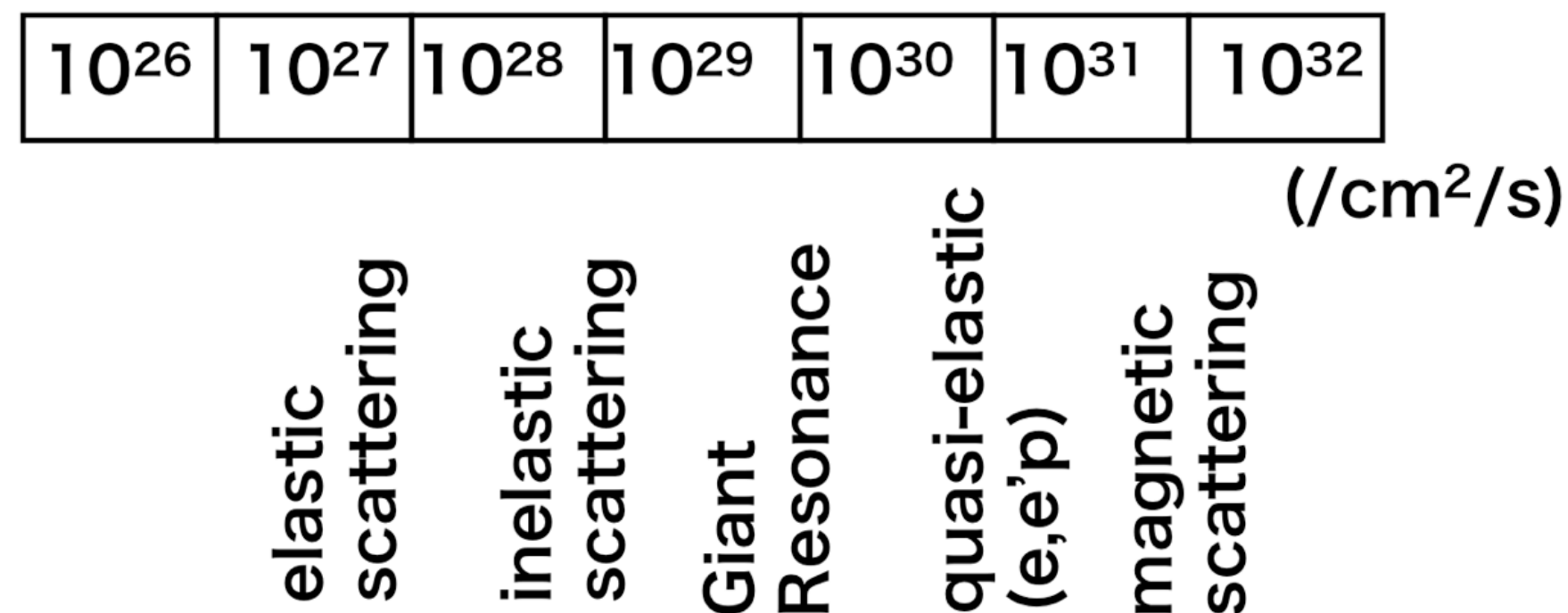


Research possibilities at the SCRIT facility

- elastic scattering → charge density distributions

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- beyond elastic scattering



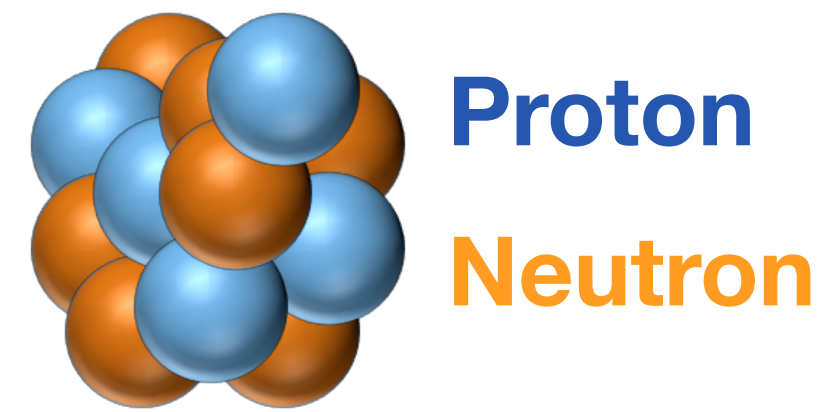
- additional research opportunity recently pointed out

4th moment of $\rho_c(r)$ → RMS of neutron distribution

$$\langle r_c^4 \rangle = \int \rho_c(r) r^4 d^3r$$

- 1) H. Kurasawa and T. Suzuki, Prog. Theor. Exp. Phys. 2019, 113D01
- 2) H. Kurasawa, T. S. and T. Suzuki, Prog. Theor. Exp. Phys. 2021, 013D02
- 3) H. Kurasawa and T. Suzuki, Prog. Theor. Exp. Phys. 2022, 023D03
- 4) T. Suzuki, Prog. Theor. Exp. Phys. 2023, 013D02

nuclear charge density, moments

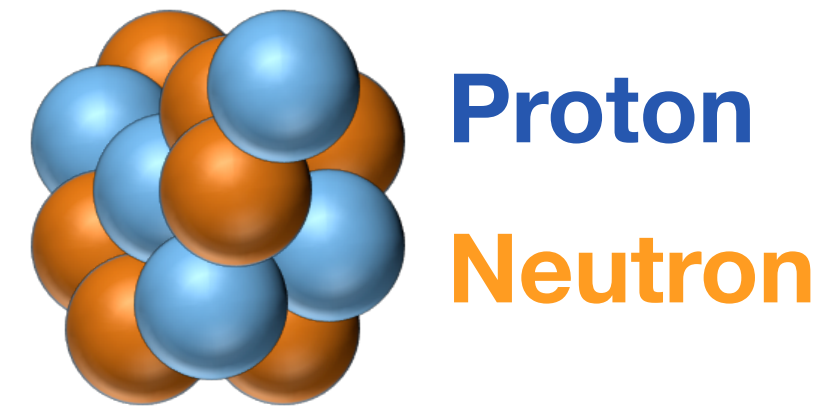


1) charge density

$$\rho_c(r) = \rho_c^p(r) + \rho_c^n(r)$$

$$\rho_c^p(r) = \int \rho_p(r) \rho_{p(point)}(r - r') d^3r'$$
$$\rho_c^n(r) = \int \rho_n(r) \rho_{n(point)}(r - r') d^3r'$$

nuclear charge density, moments



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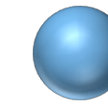
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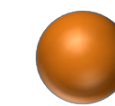
2) 2nd moment

$$\langle r_c^2 \rangle = \int r^2 \rho_c(r) d^3r = \langle r_{p(point)}^2 \rangle + \langle r_p^2 \rangle$$
$$+ \frac{N}{Z} \langle r_n^2 \rangle + \text{rel. corr.}$$

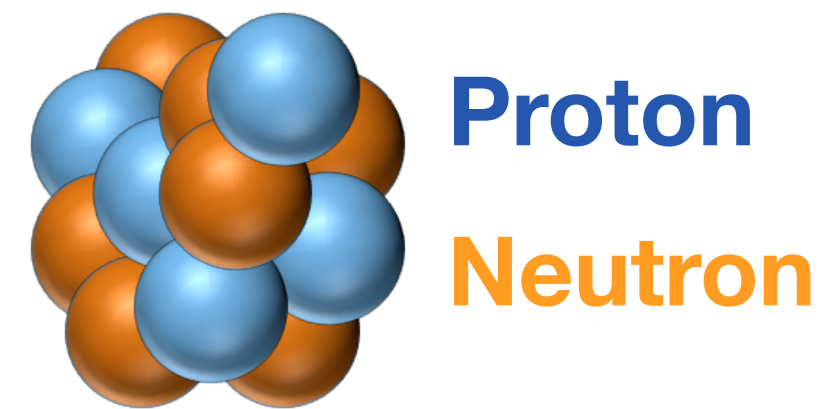
Proton



Neutron



nuclear charge density, moments



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Proton
Neutron

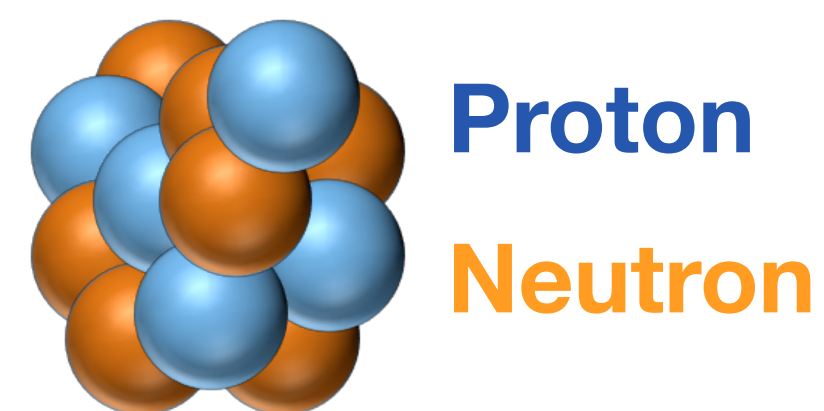
3) 4th moment

$$\langle r_c^4 \rangle = \int r^4 \rho_c(r) d^3r$$

$$= \langle r_{p(point)}^4 \rangle + \frac{10}{3} \langle r_{p(point)}^2 \rangle \langle r_p^2 \rangle + \frac{10}{3} \langle r_{n(point)}^2 \rangle \langle r_n^2 \rangle \frac{N}{Z} + \text{rel. corr.}$$

Proton
Neutron

nuclear charge density, moments



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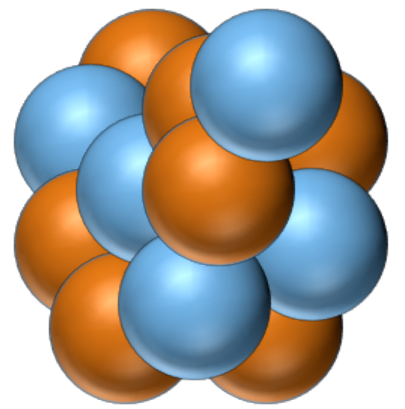
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RMS n-radius

neutron distribution by electron scattering??



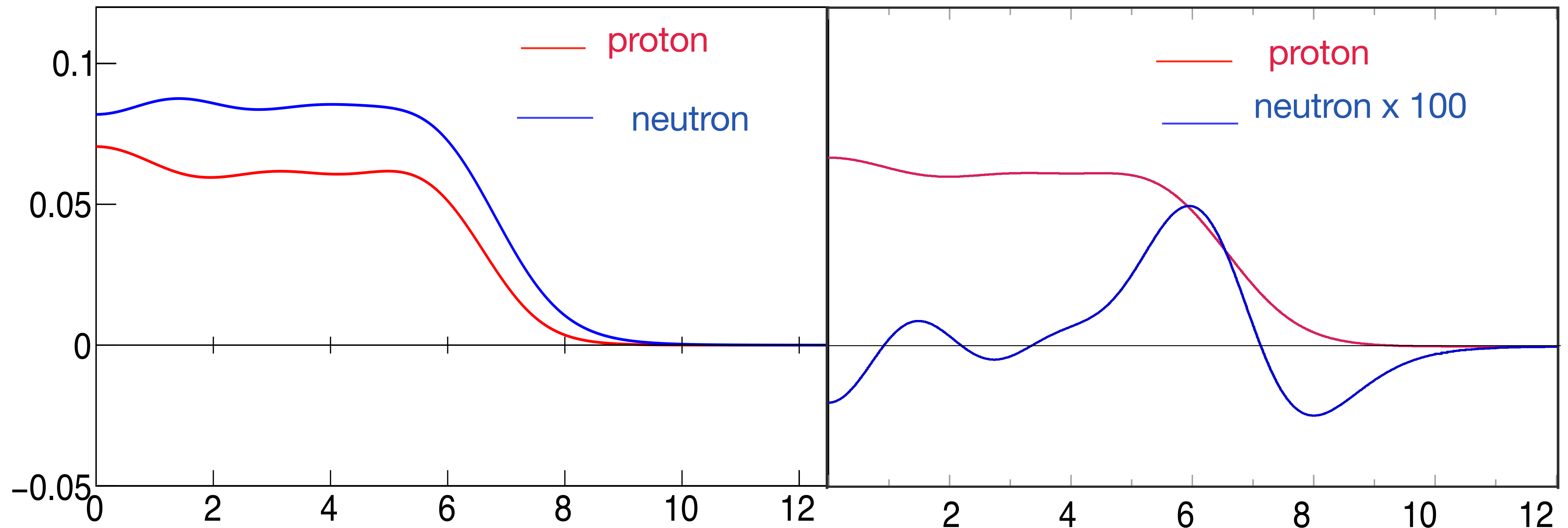
Proton

Neutron

RMF NL₃ (H. Kurasawa)

(point) nucleon density

charge density distributions



ways to access the fourth moment, $\langle r_c^4 \rangle$

$$\langle r_c^4 \rangle = \int r^4 \rho_c(r) d^3r$$

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1) elastic scattering covering upto very high q

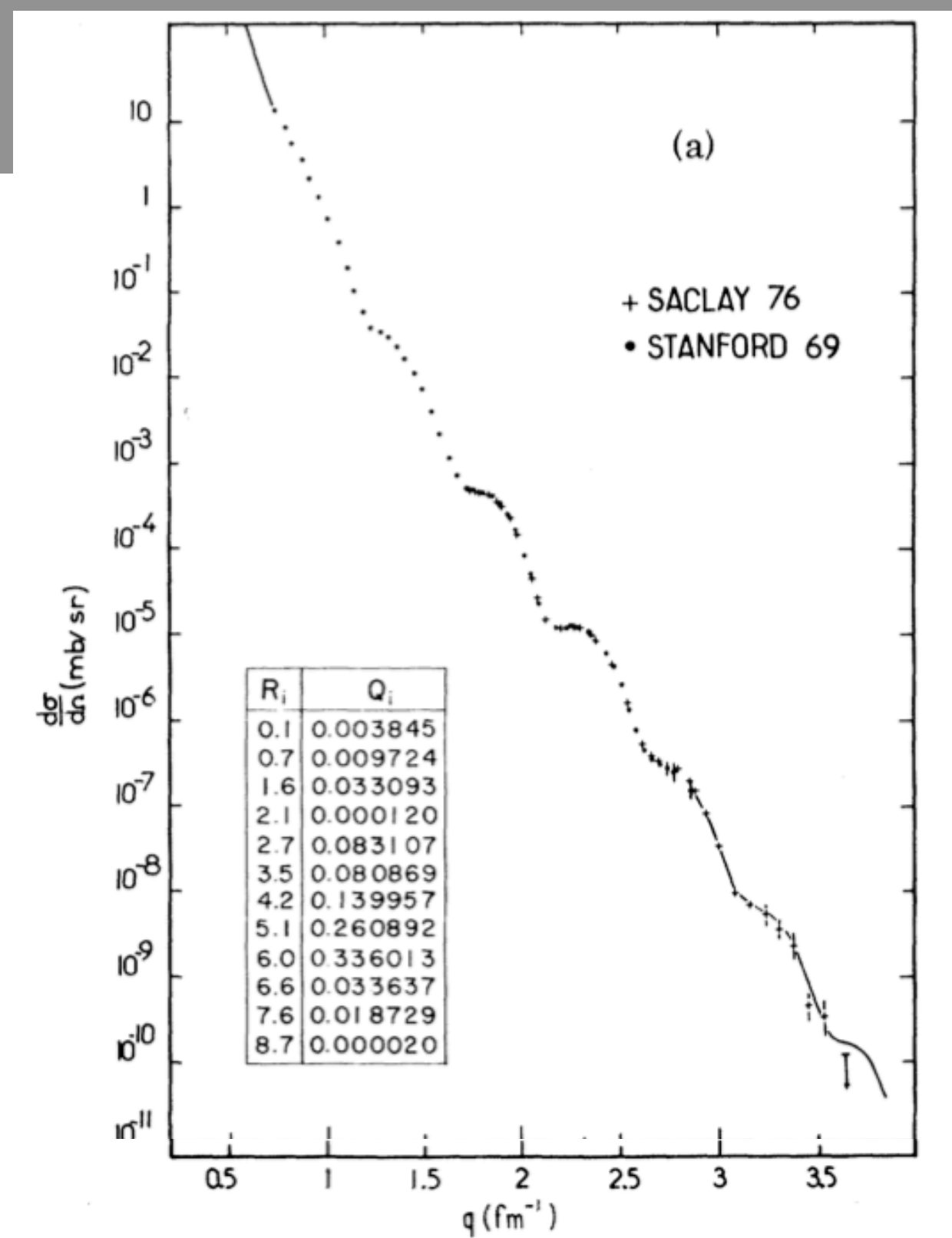
$$^{208}\text{Pb} \quad R_n = 5.728 \pm 0.057 \text{ fm}$$

$$\Delta r_{np} = 0.275 \pm 0.070 \text{ fm}$$

$$^{48}\text{Ca} \quad R_n = 3.597 \pm 0.021 \text{ fm}$$

$$\Delta r_{np} = 0.220 \pm 0.026 \text{ fm}$$

H. Kurasawa, T. S. and T. Suzuki, Prog. Theor. Exp. Phys. 2021, 013D02



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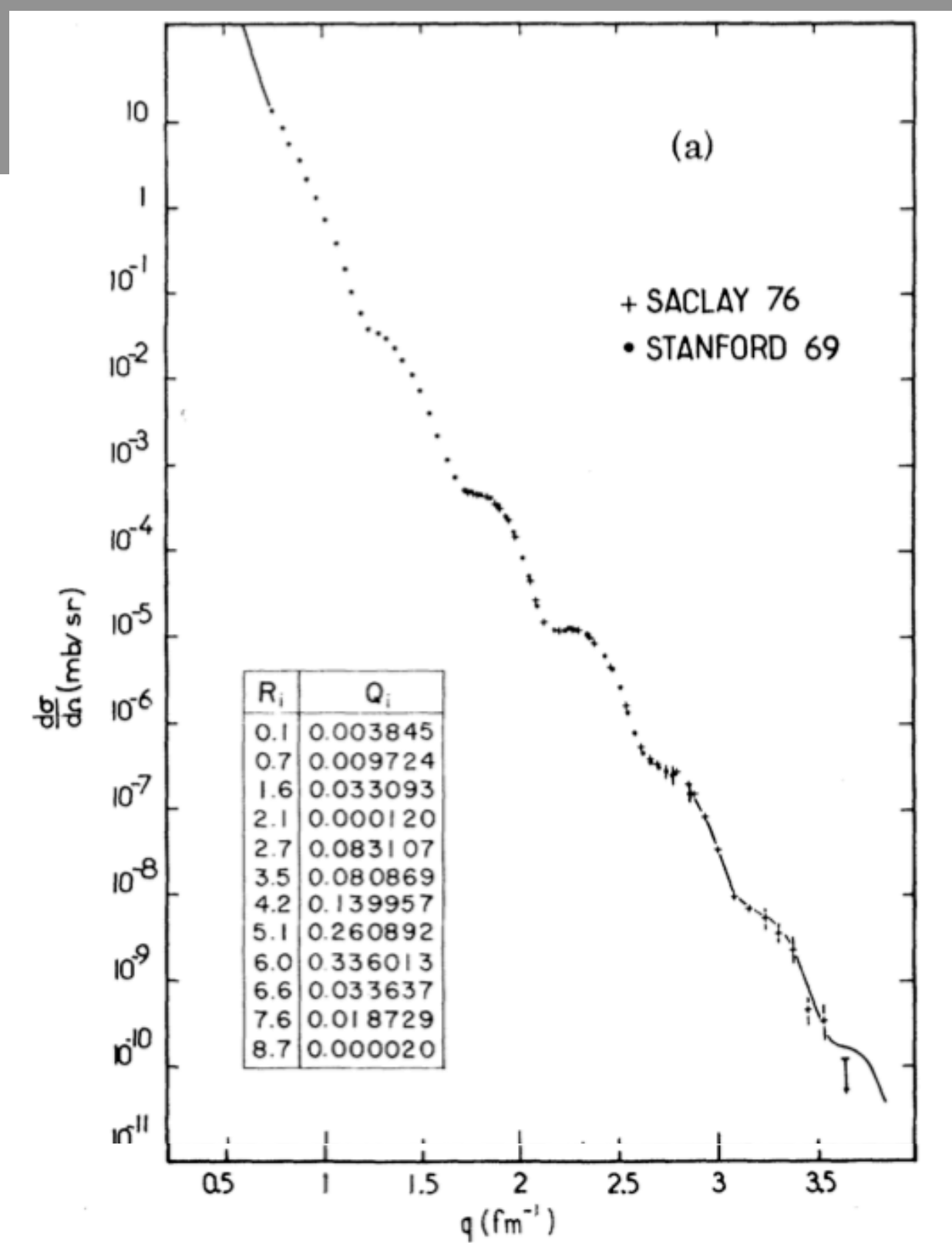
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H. Kurasawa, T. S. and T. Suzuki, Prog. Theor. Exp. Phys. 2021, 013D02

2) elastic scattering at very low q

$$\frac{d\sigma}{d\Omega} = \frac{d\sigma_{Mott}}{d\Omega} |F_c(q)|^2$$

$$F_c(q) \sim 1 - \frac{\langle r_c^2 \rangle}{6} q^2 + \frac{\langle r_c^4 \rangle}{120} q^4 + \dots$$



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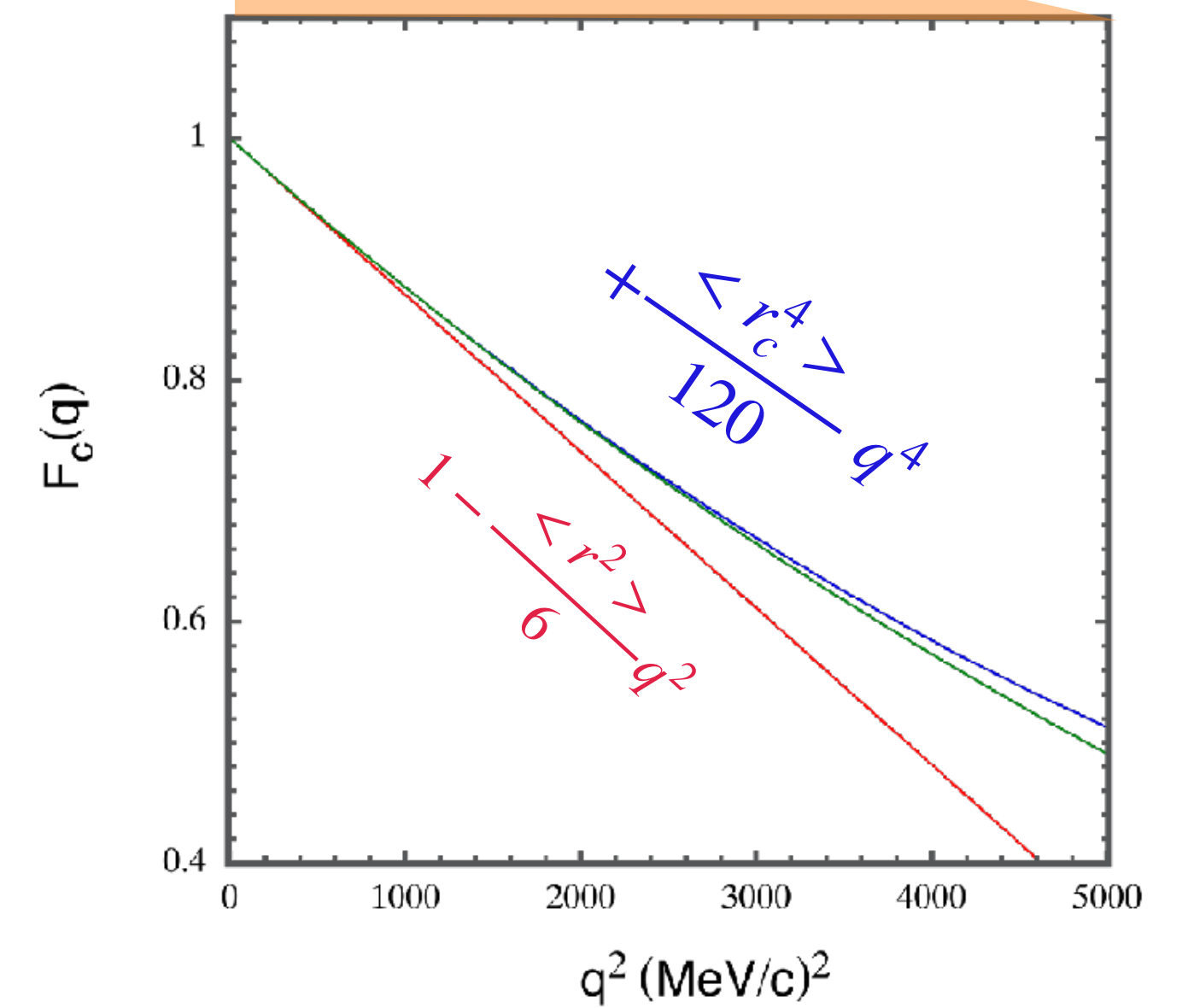
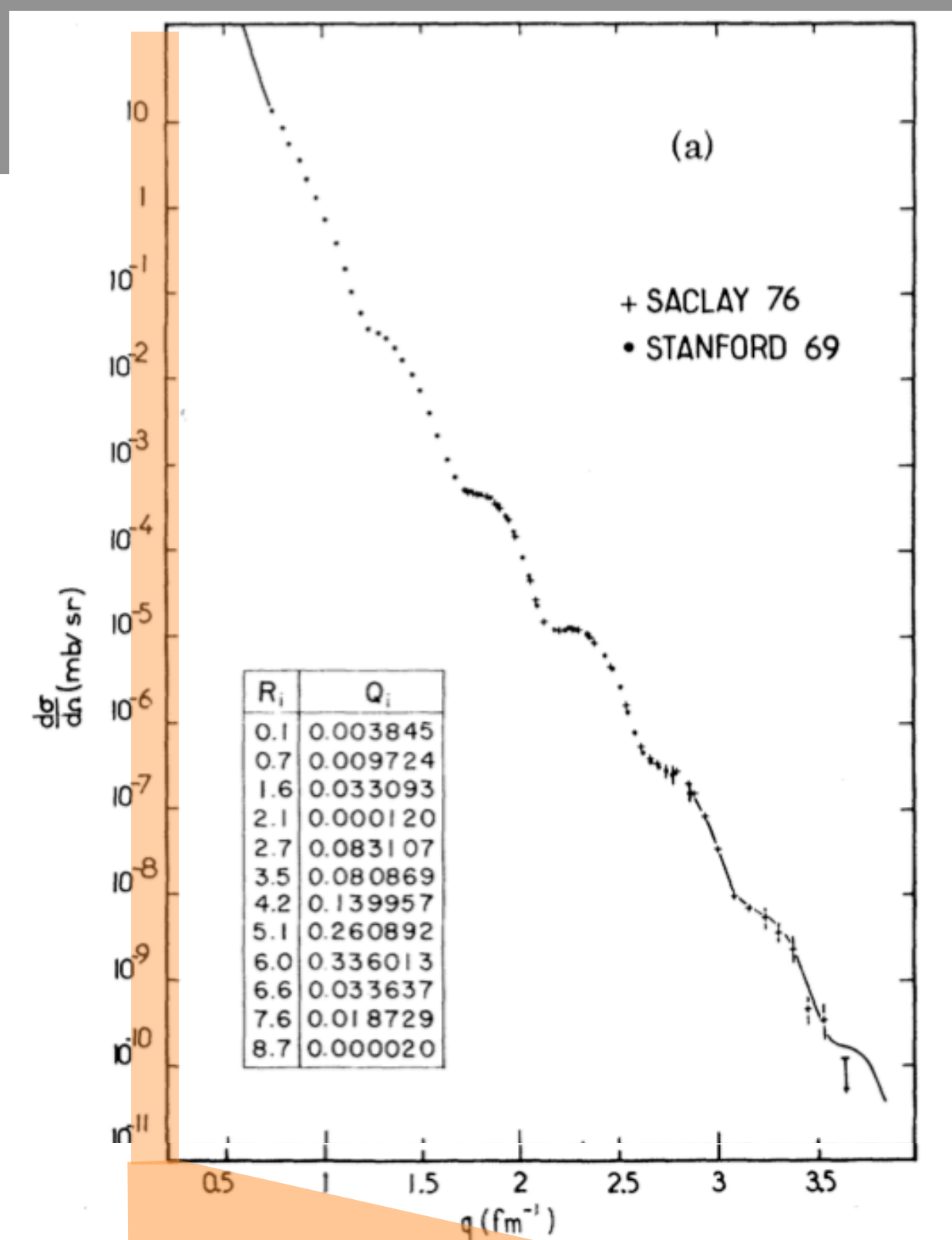
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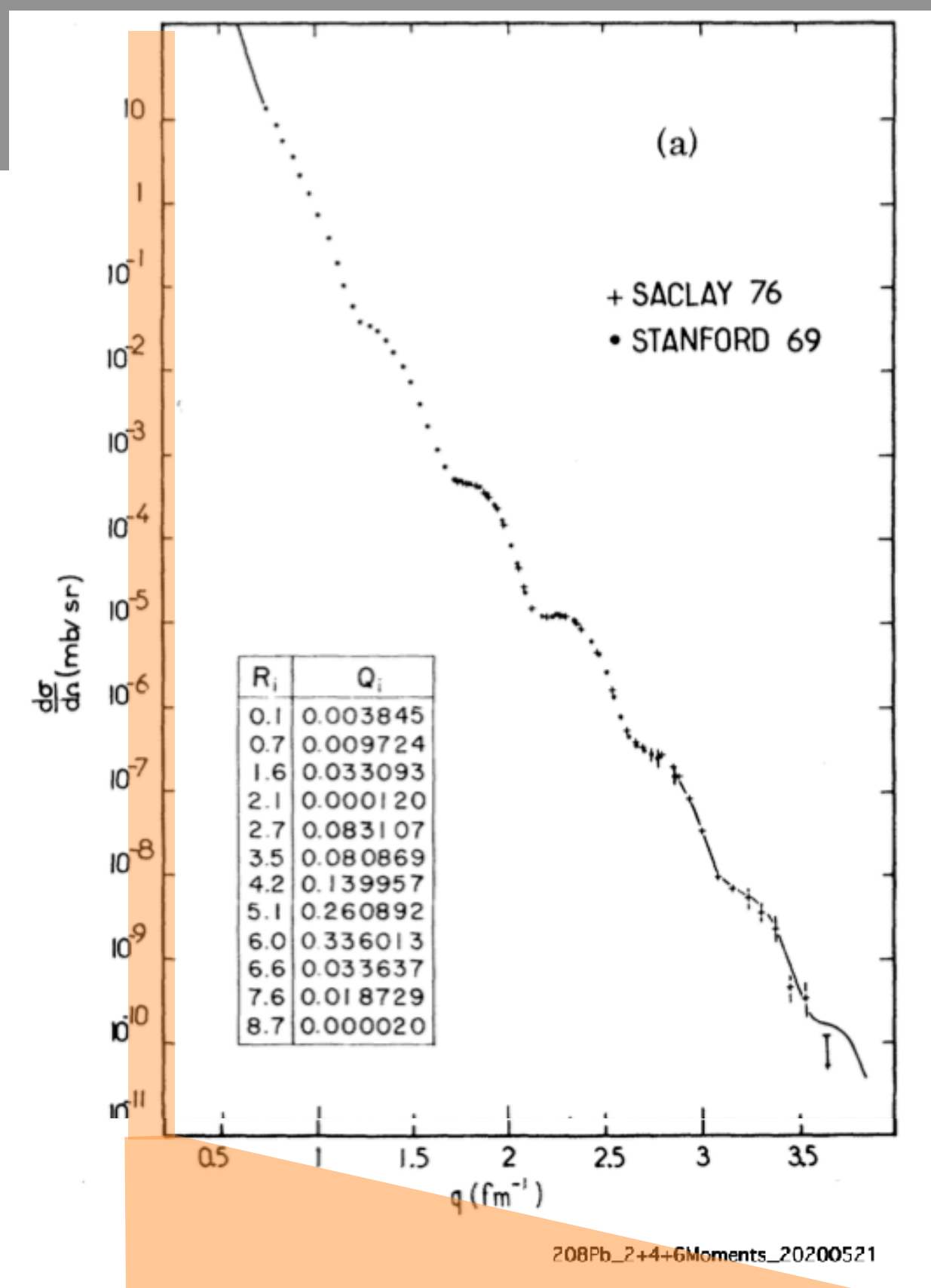
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H. Kurasawa, T. S. and T. Suzuki, Prog. Theor. Exp. Phys. 2021, 013D02



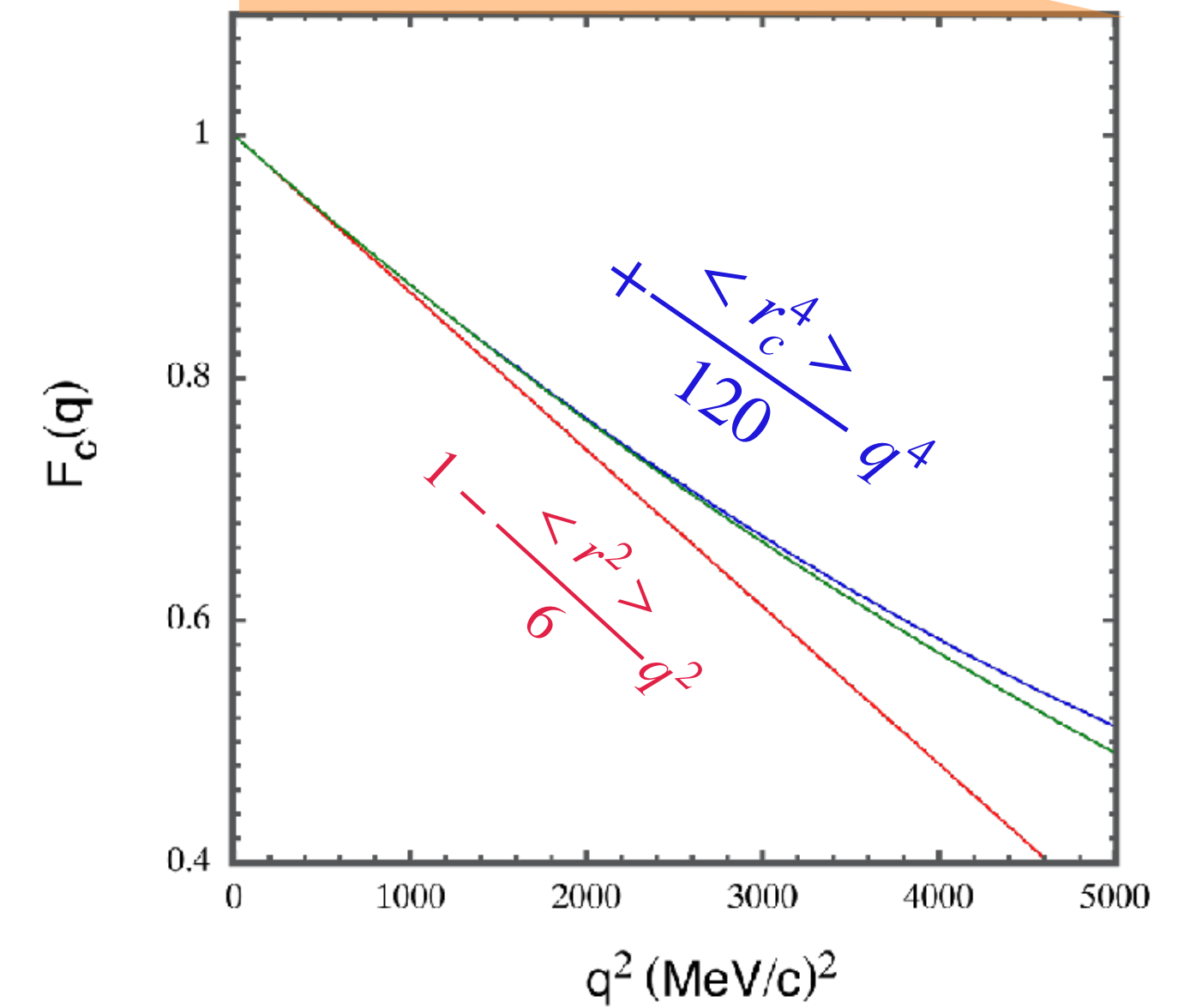
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$$\frac{d\sigma_{Mott}}{d\Omega} \propto 1/q^4$$

=> low-L SCRIT exp. ??



conclusions

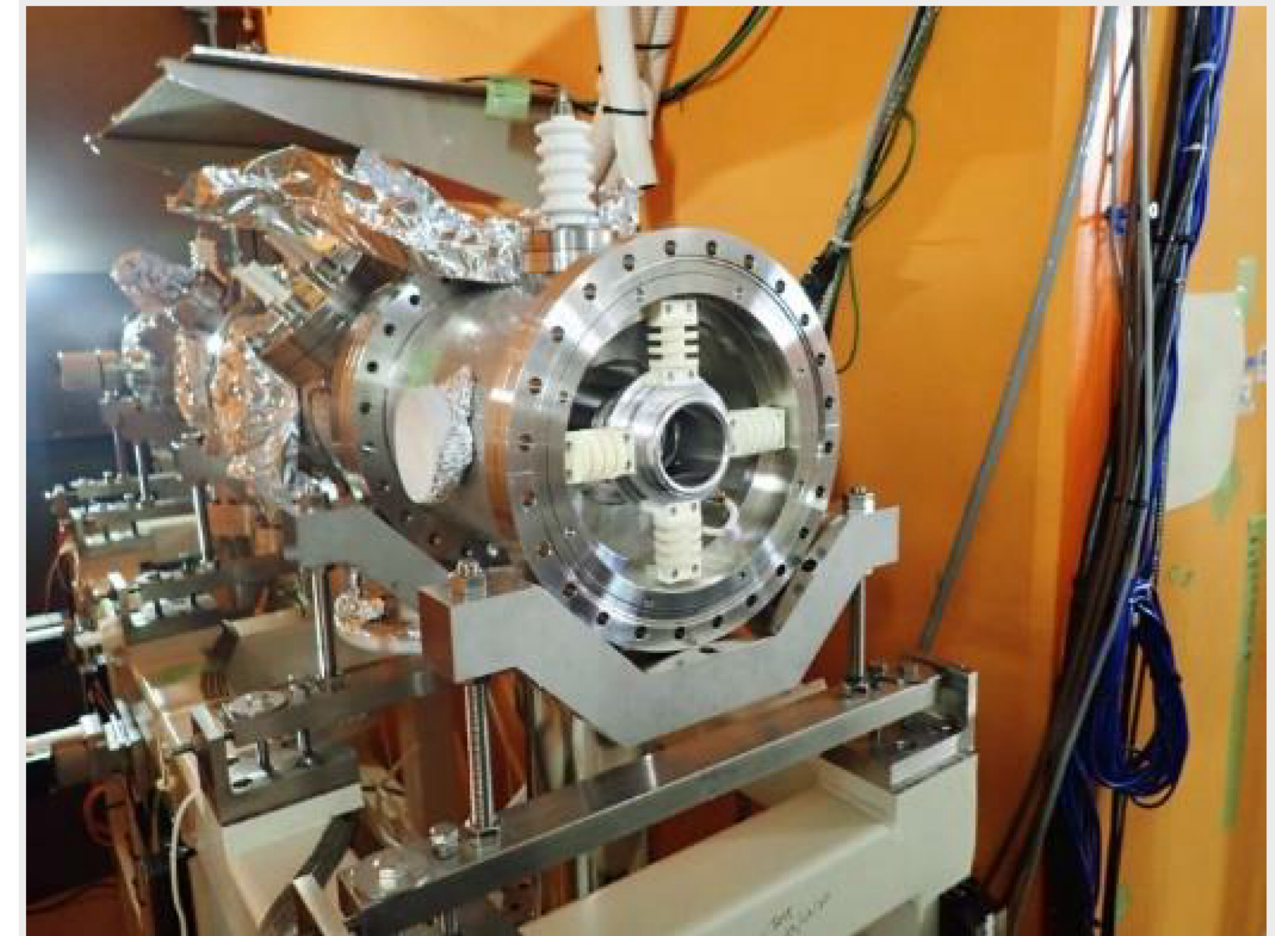
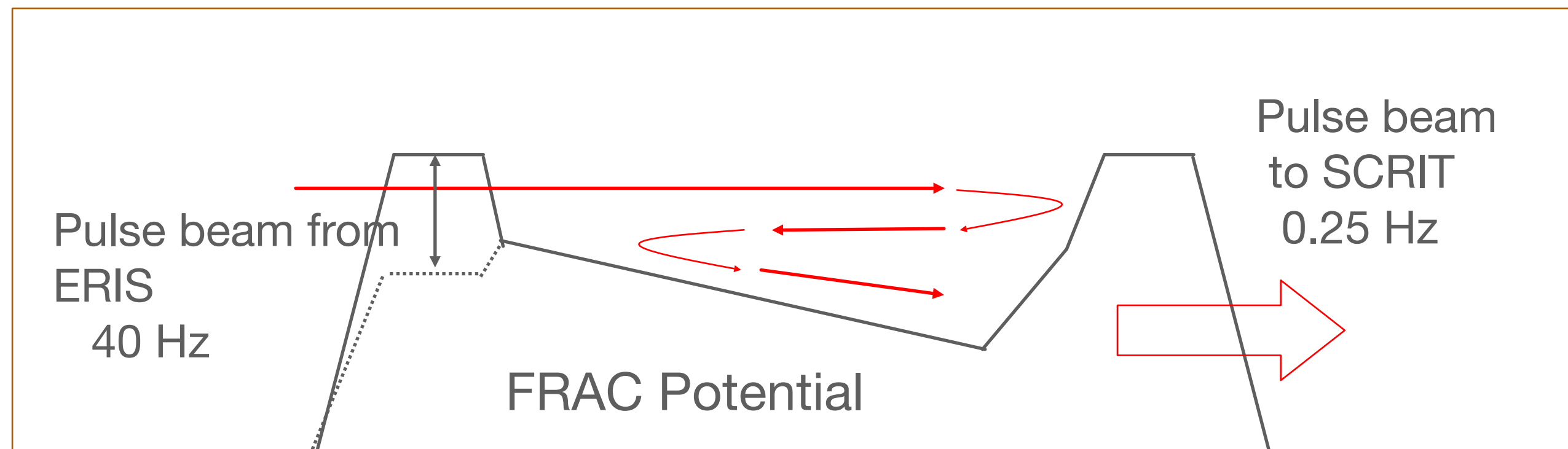
- **The SCRIT facility started its operation**
 - the world's first and currently only-one facility
 - successful demonstration with online-produced radioactive isotope
 - ISOL upgrade to 2kW is underway for $^{132}\text{Sn}(e,e')$

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- **electron scattering for exotic nuclei is no long just a dream**
 - charge density distributions by elastic scattering
 - inelastic scattering with higher L
 - higher moments of charge density => neutron distributions

FRAC : Cooler and Buncher

Ion beams are stacked in the FRAC for:

- ✓ Cooling by collision with Ne buffer gas
- ✓ Increasing #RI-ions in single pulses



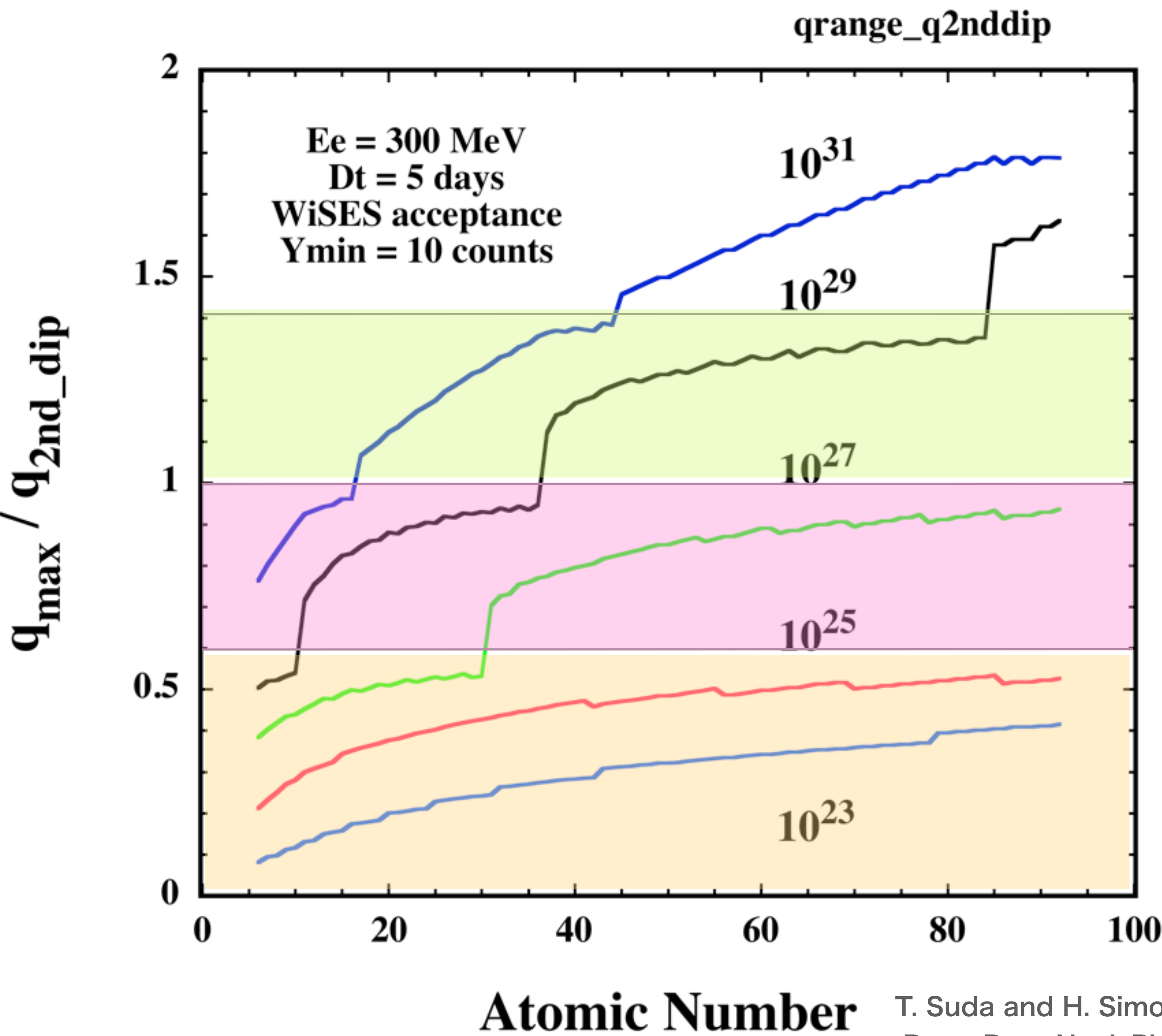
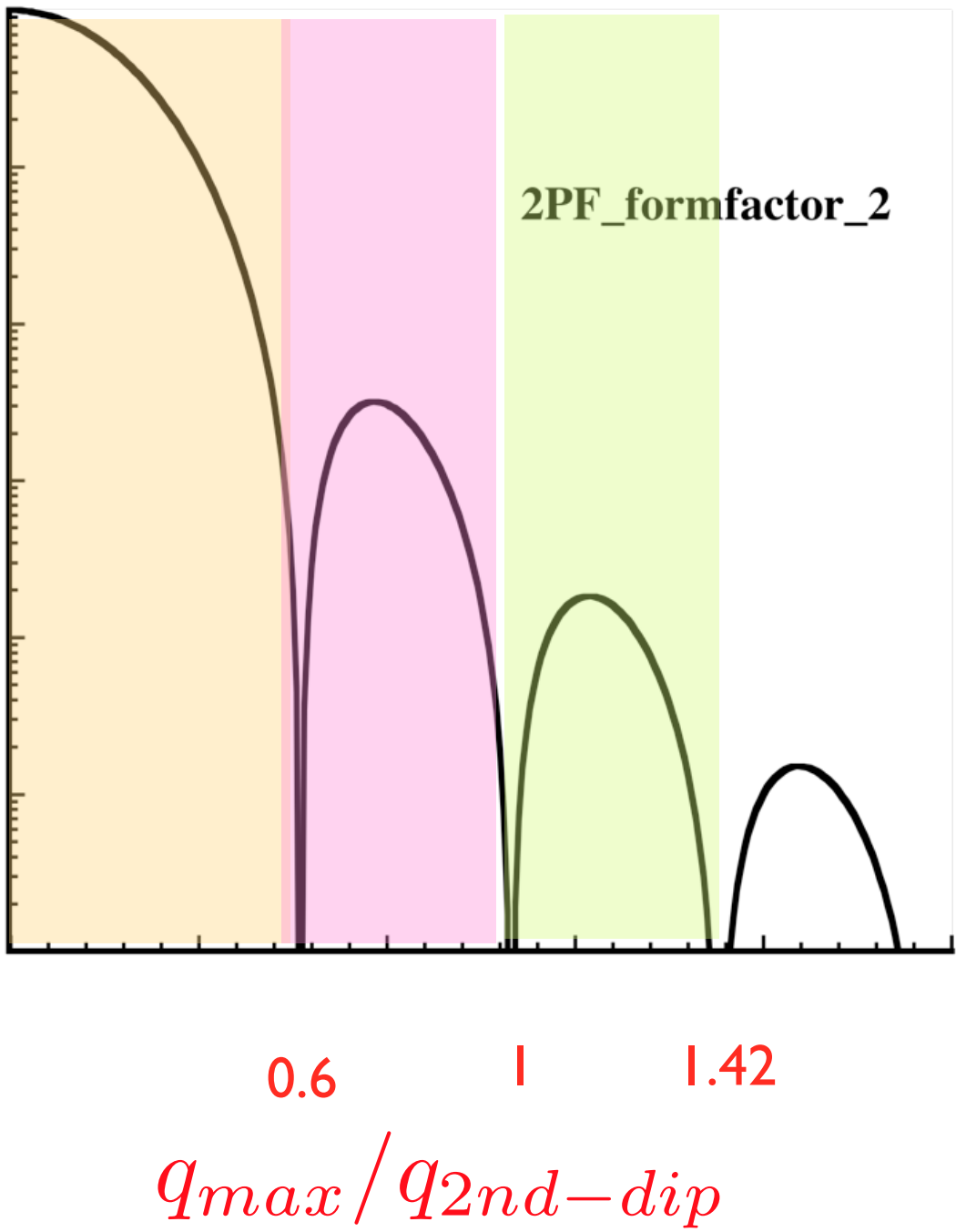
M. Wakasugi, Rev. Sci. Instrum. 89 (2018) 095107
H. Ohnishi, Nucl.Instr.Meth.B541(2023)380-384

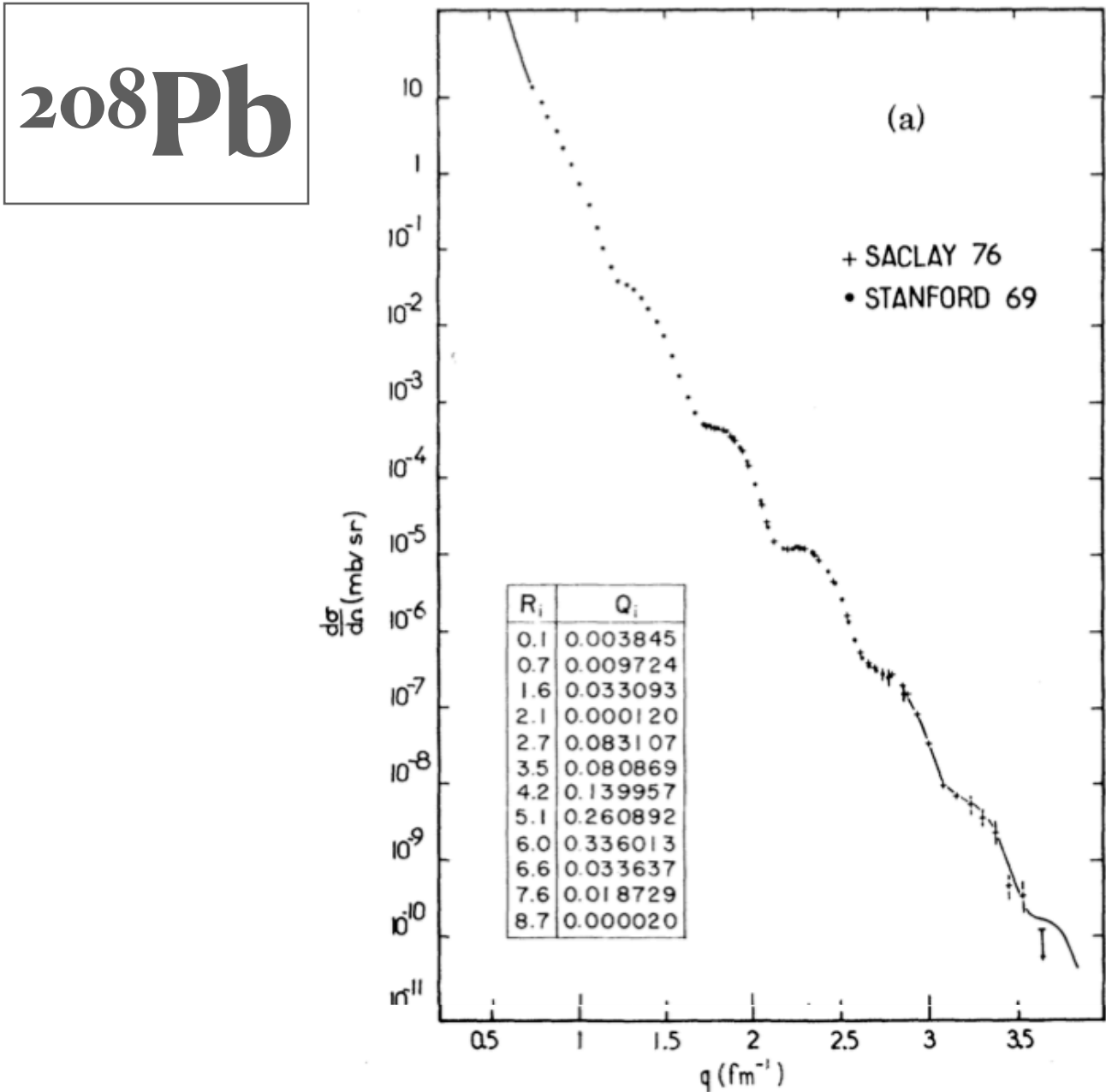
Total dc-to-pulse conversion eff. $\sim 80\%$

→ ^{137}Cs ions injected into SCRIT : 2×10^7 ions/pulse

Accessible q-range for L and Z

$$\sigma \propto Z^2$$
$$\sigma \propto 1/q^4$$



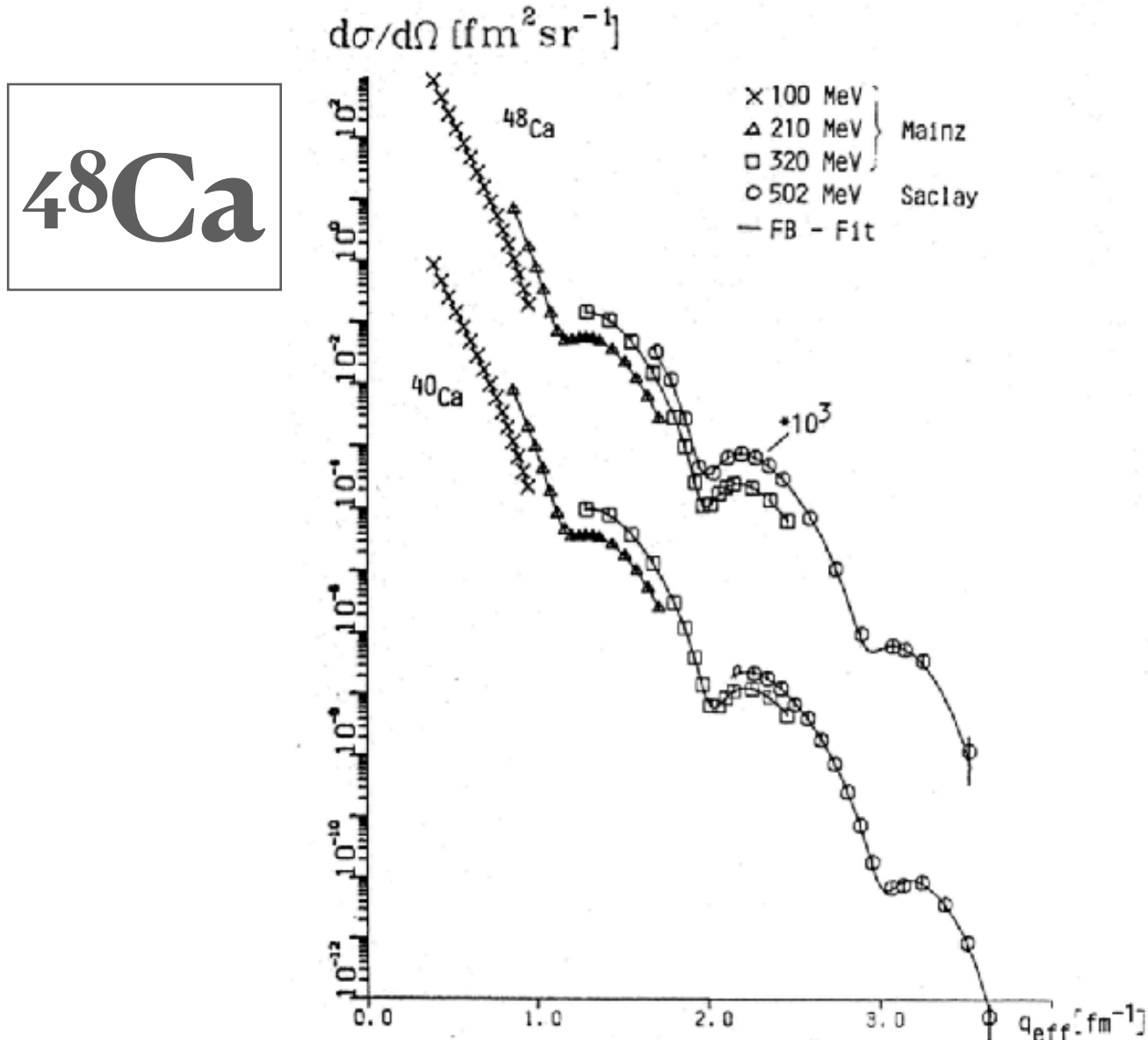


		R_p	R_n	δR
^{208}Pb	Rel.	5.454(0.013)	5.728(0.057)	<u>0.275(0.070)</u>
	Non.	5.447(0.014)	5.609(0.054)	0.162(0.068)
	Exp.	$R_c = 5.503(0.014)$		

JLab : PREX I,II (parity-violating e-scattering)

$\Delta r_{np} \equiv R_n - R_p = \underline{0.283 \pm 0.071 \text{ fm}}$

PRL 126, 172502 (2021)

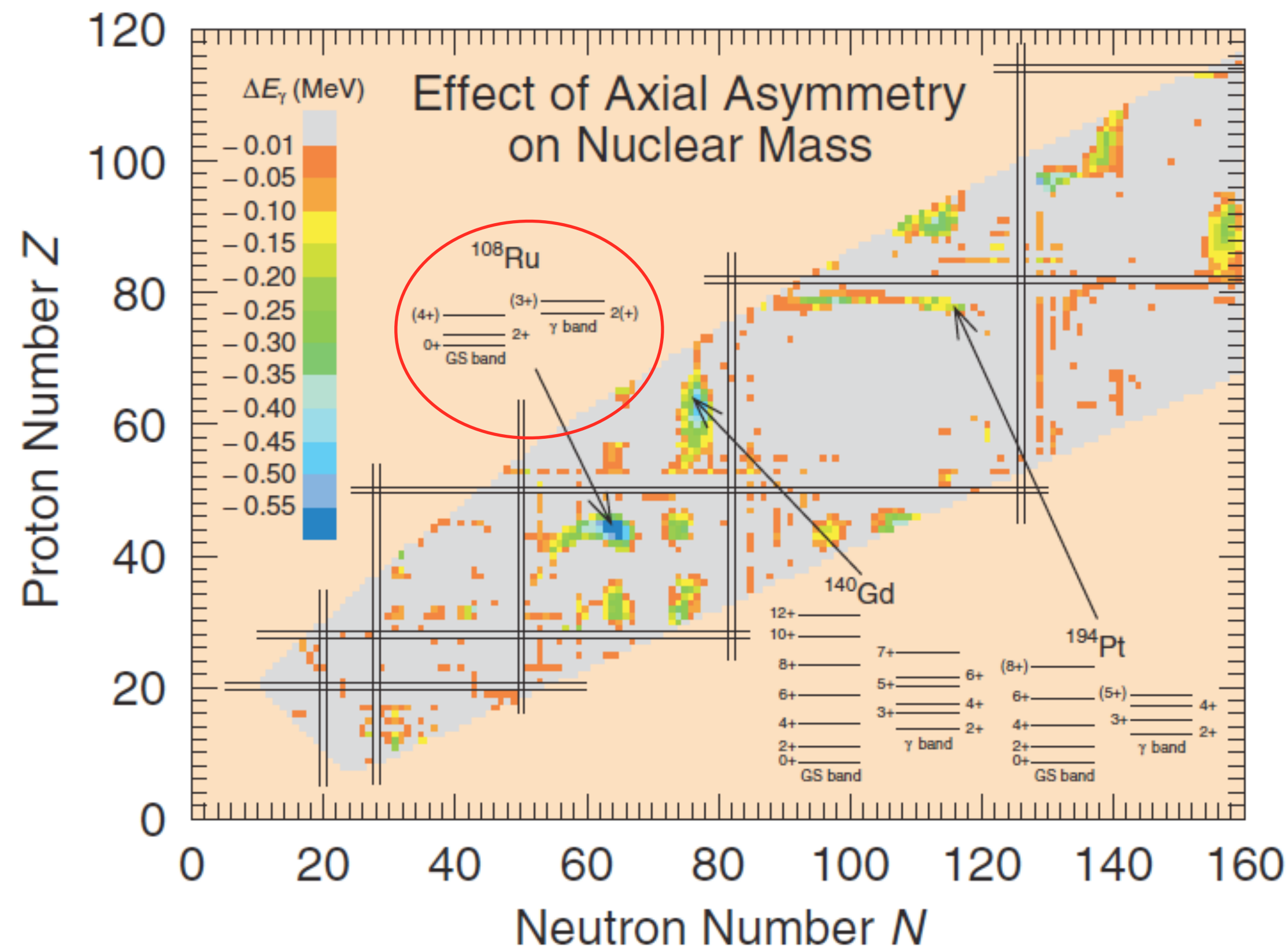


		R_p	R_n	δR
^{48}Ca	Rel.	3.378(0.005)	3.597(0.021)	<u>0.220(0.026)</u>
	Non.	3.372(0.009)	3.492(0.028)	0.121(0.036)
	Exp.	$R_c = 3.451(0.009)$		

JLab : CREX (parity-violating e-scattering)

$\Delta r_{np} \equiv R_n - R_p = \underline{0.121 \pm 0.026 \text{ fm}}$

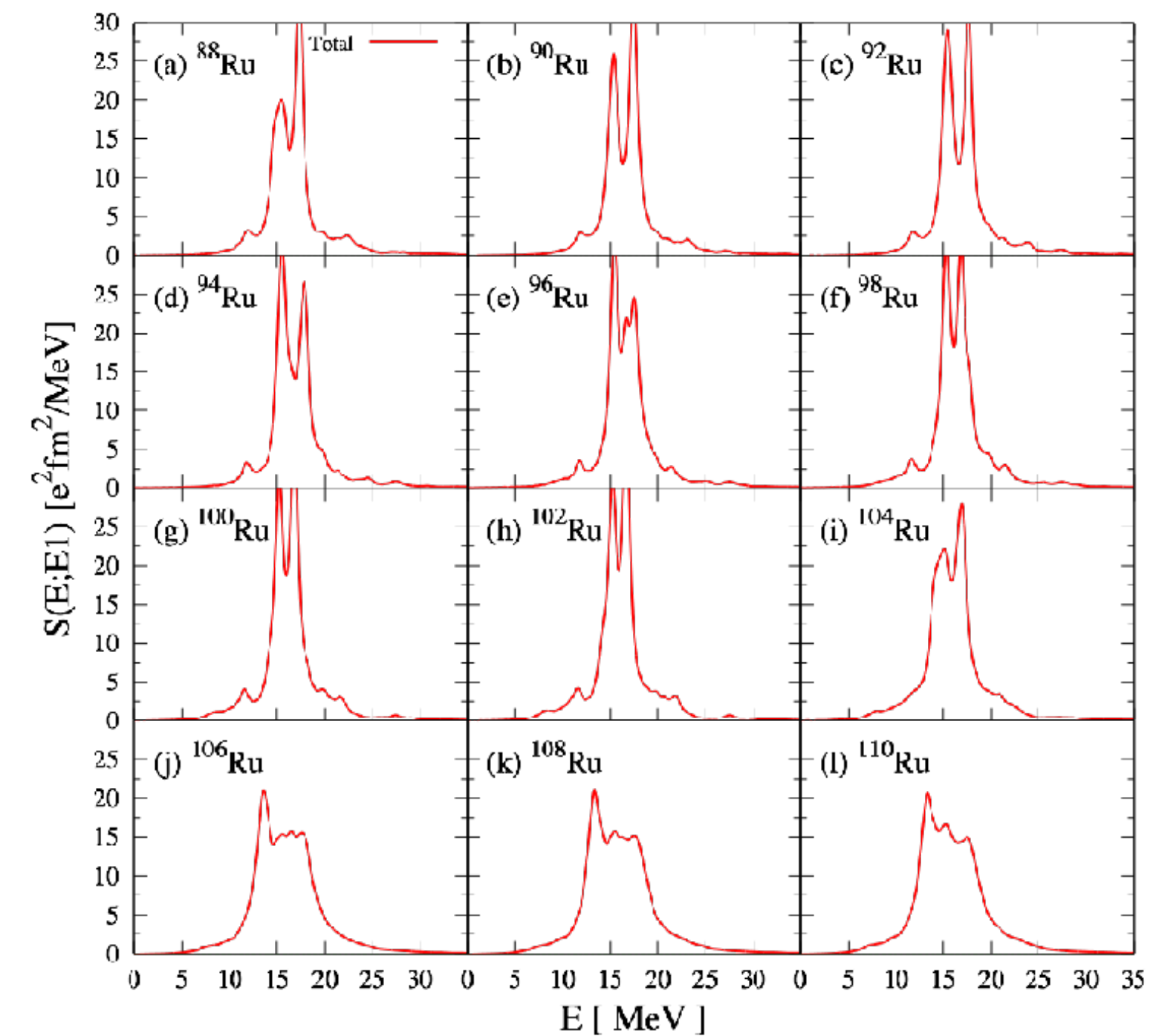
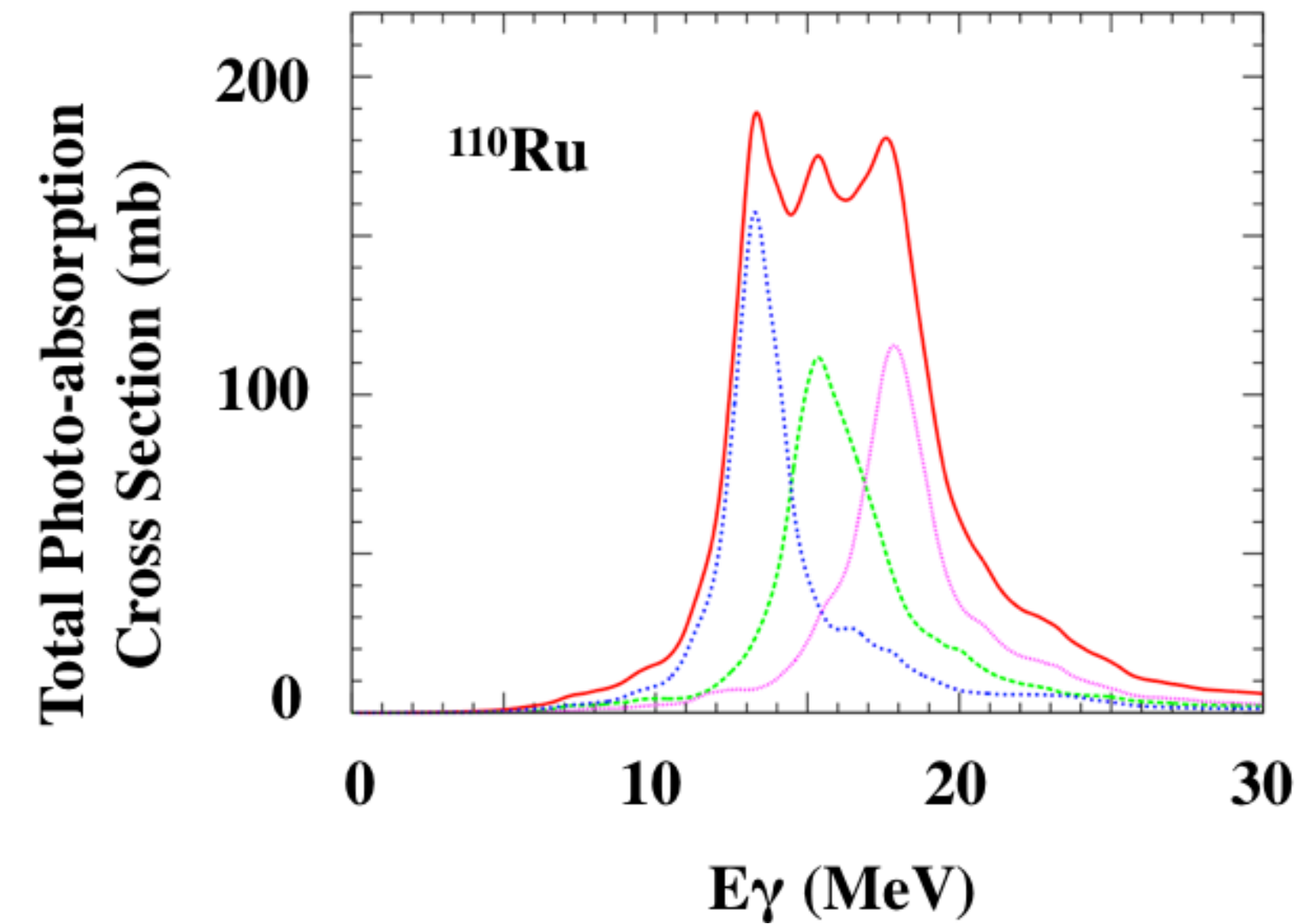
triaxial nucleus expected in exotic nuclei



^{108}Ru 273 sec

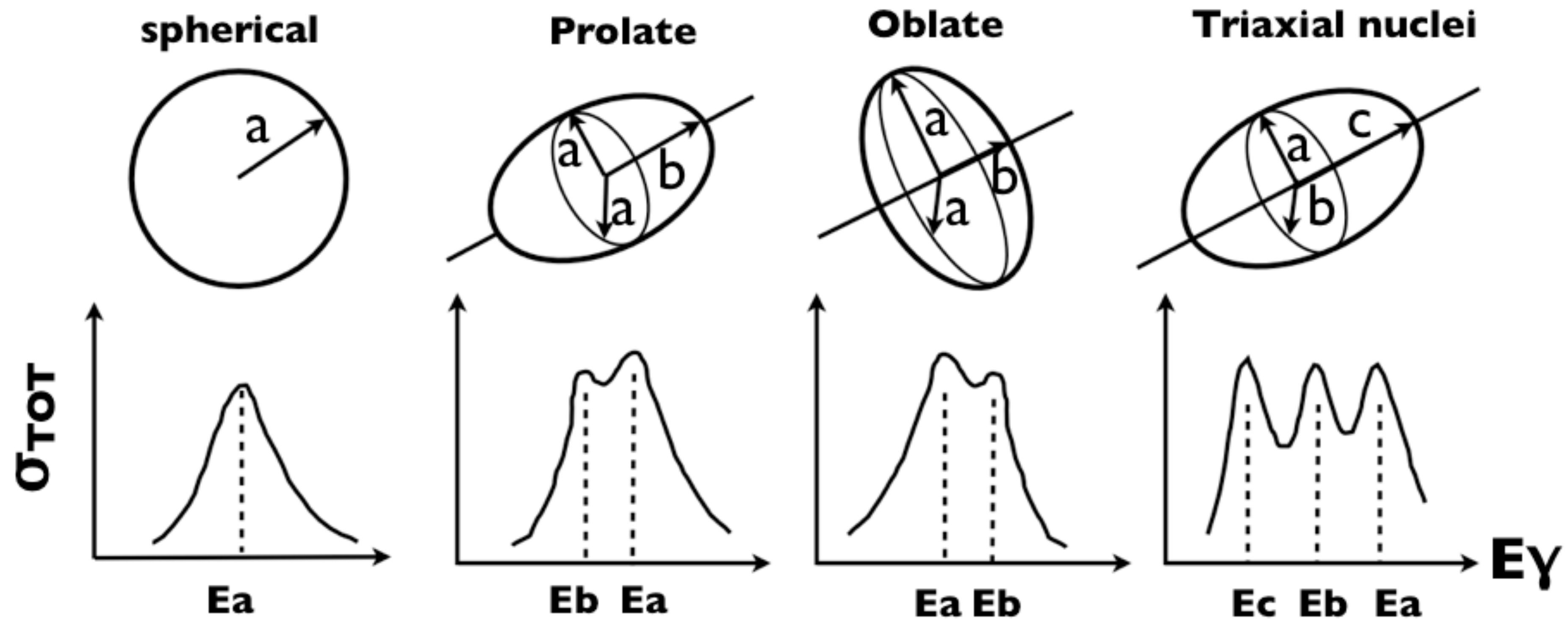
^{110}Ru 15 sec

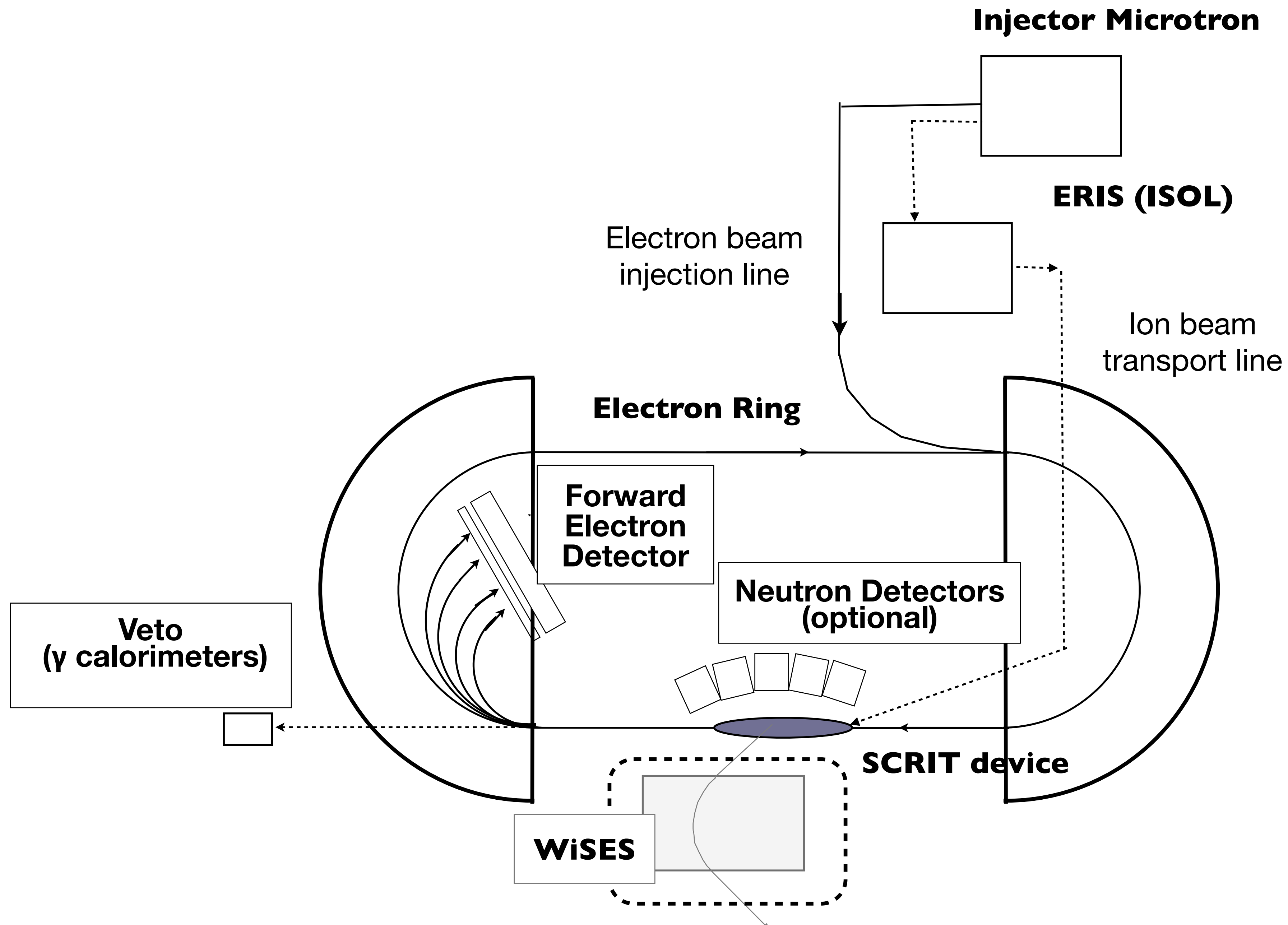
P. Moller et al., PRL97, 162502 (2006)



S. Ebata : private communication

photo-nuclear responses of exotic nuclei

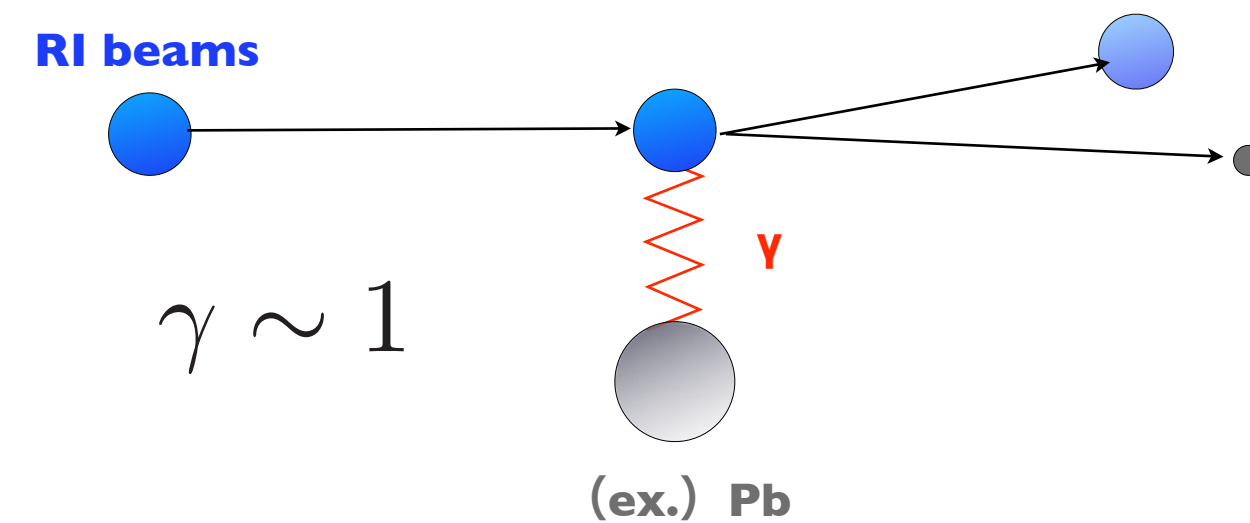




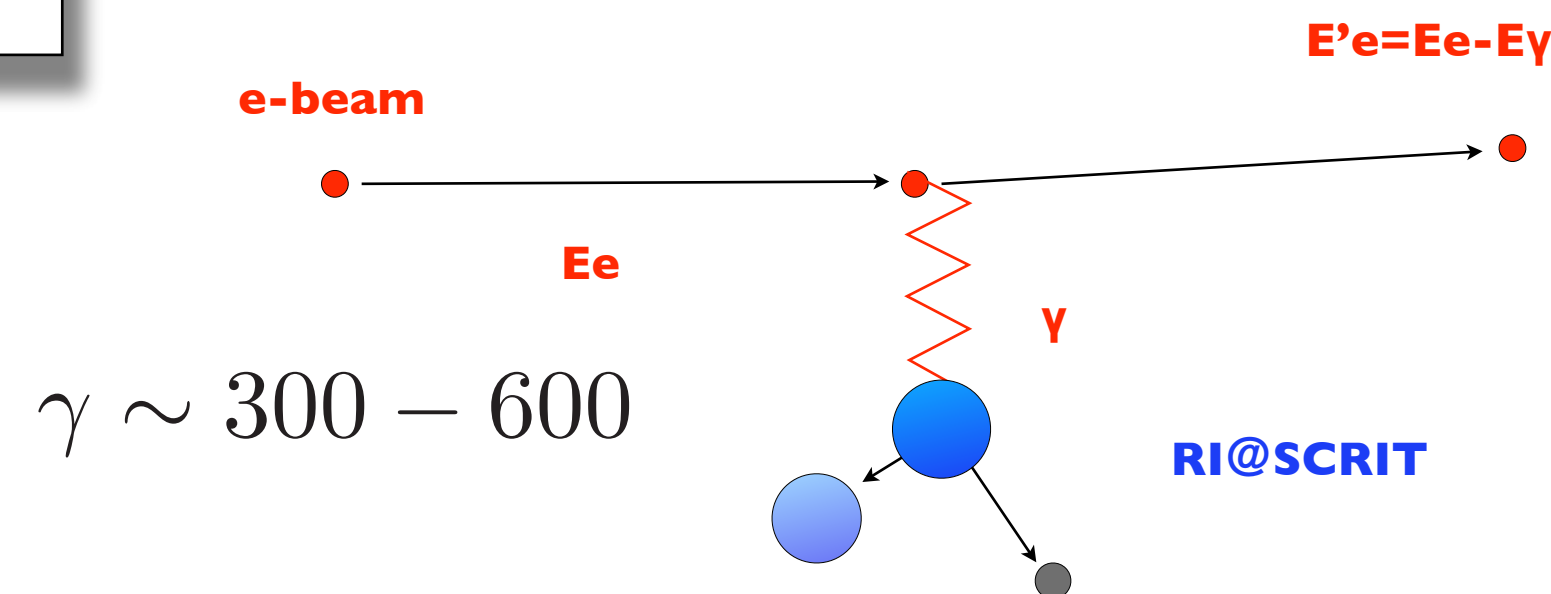
photonuclear reaction for exotic nuclei

so far

only way : Coulomb excitation in heavy ion reaction



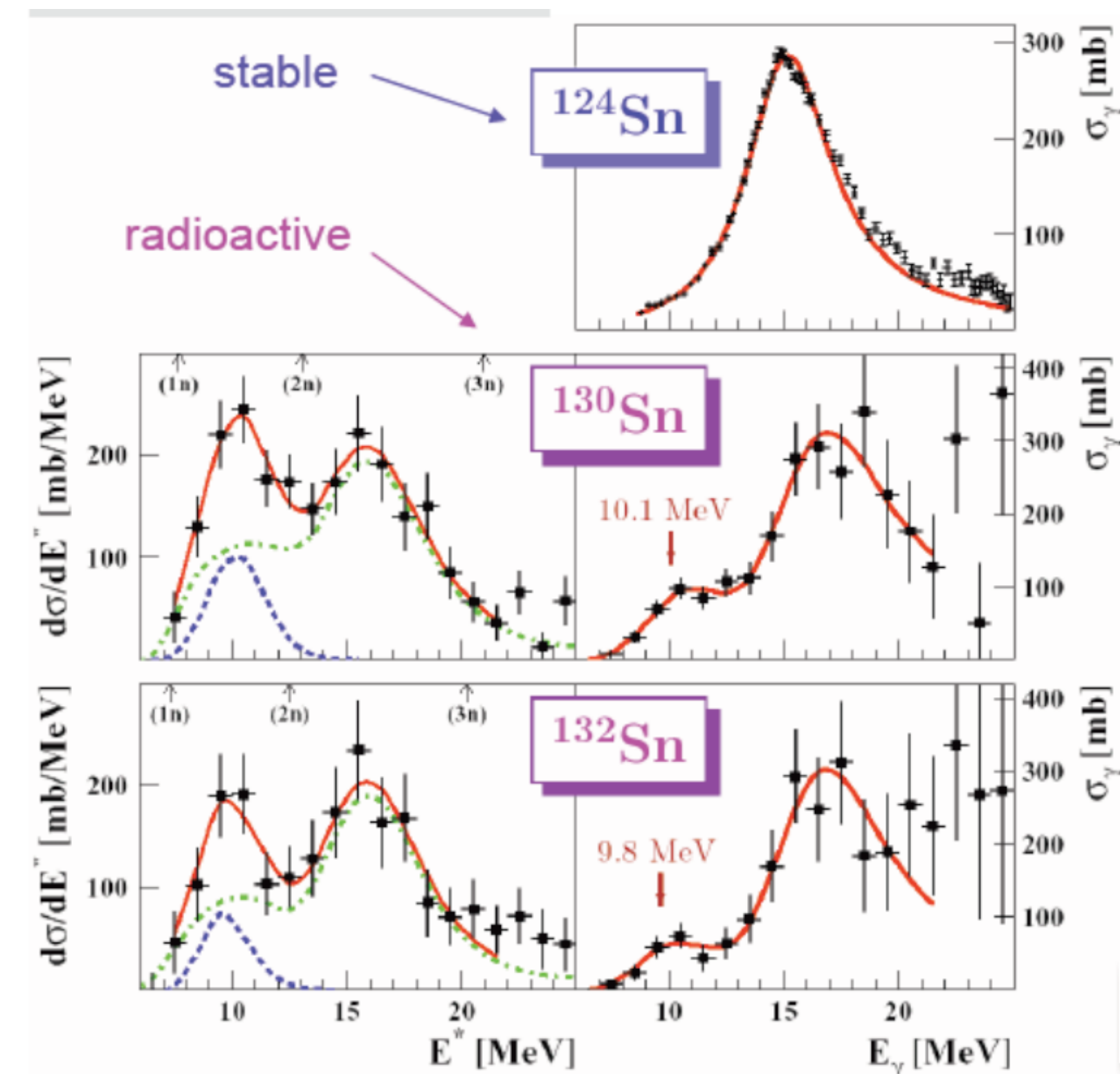
SCRIT facility



purely EM probe
well under control
negligible multi-stop

ultra-forward
electron scattering

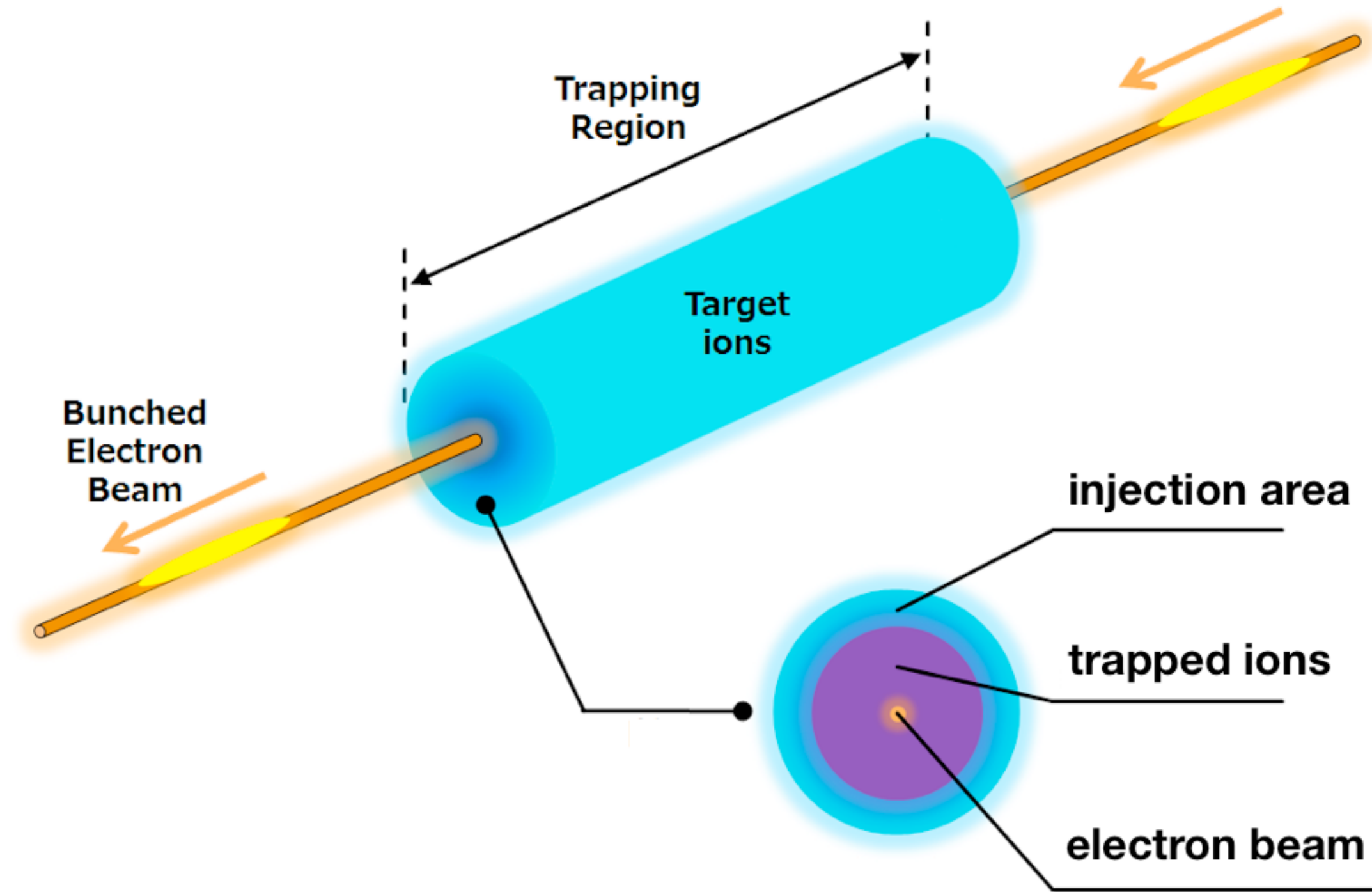
$^{132}\text{Sn} + \text{Pb} \rightarrow ^{131}\text{Sn} + n + X$ @ GSI



Virtual Photon flux

$$\frac{d^2\sigma}{dE_e d\Omega} = \sum \frac{d^2 N_e^{EL}(E, E_\gamma, \theta)}{dE_\gamma d\Omega} \cdot \sigma_\gamma^{EL}(E_\gamma)$$

Ion trapping and luminosities



trapping efficiency $\sim 100\%$
overlapping efficiency $\sim 17\%$

Luminosity ($10^{25}/\text{cm}^2/\text{s}$)

