The background of the slide is a nighttime photograph of a large, modern building with a white facade and many lit windows. The building is surrounded by trees and a paved area. In the background, a hill and a city skyline with lit-up buildings are visible under a dark blue sky.

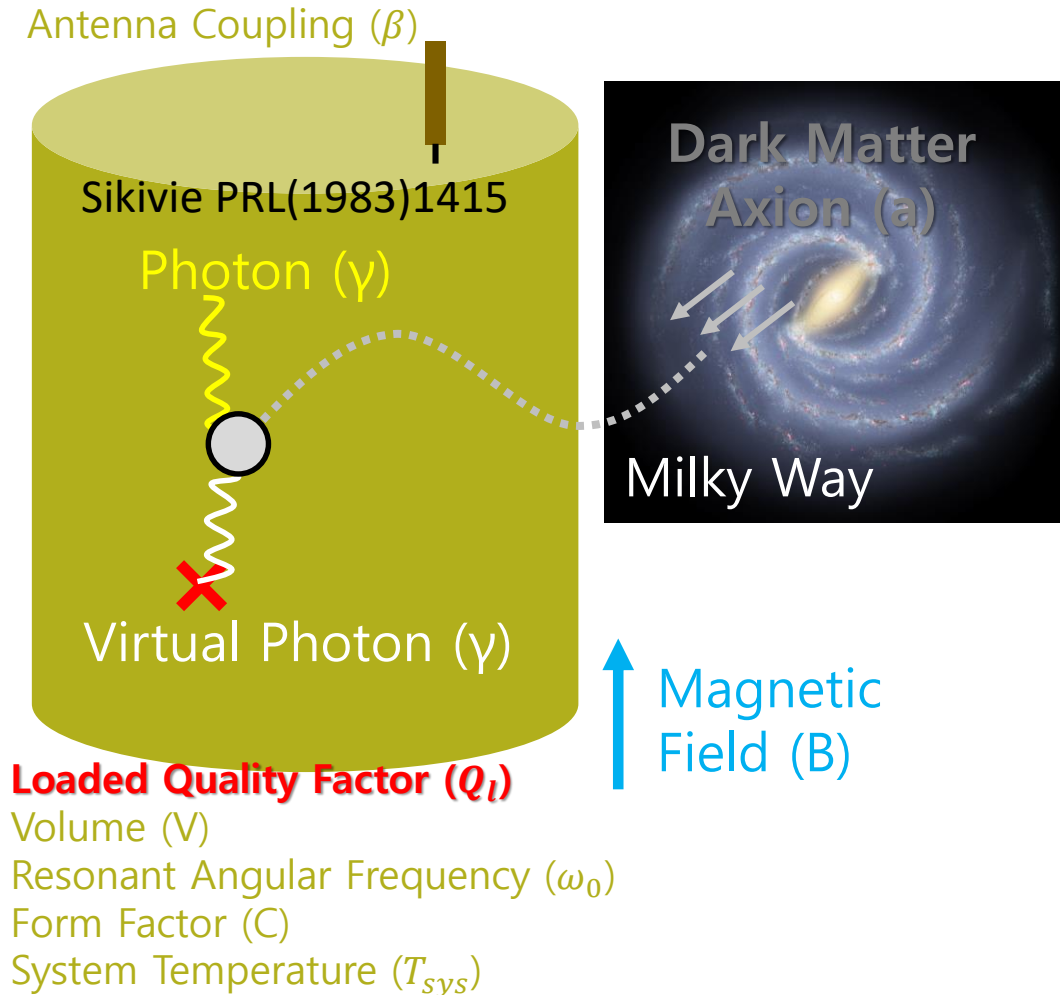
Accelerating Axion Haloscope Experiments : Strategies and Innovations at CAPP

Danho Ahn (IBS-CAPP)

Center for Axion and Precision Physics Research, Institute of Basic Science,
Daejeon 34051, Republic of Korea

Introduction

Cavity Haloscope Experiment for Dark Matter Axion Search



Signal Power

$$P_{sig} = \frac{\beta}{1 + \beta} g_{\alpha\gamma\gamma}^2 \frac{\rho_a}{m_a^2} B^2 V \omega_0 C \frac{Q_a Q_l}{Q_a + Q_l}$$

Coupling Constant

Dark Matter Axion Density

Axion Mass

Axion Quality Factor

Kim et al. JCAP03(2020)066

Scanning speed

$$\frac{df}{dt} \propto \frac{B^4 V^2 C^2}{k_B^2 T_{sys}^2} Q_l Q_a$$

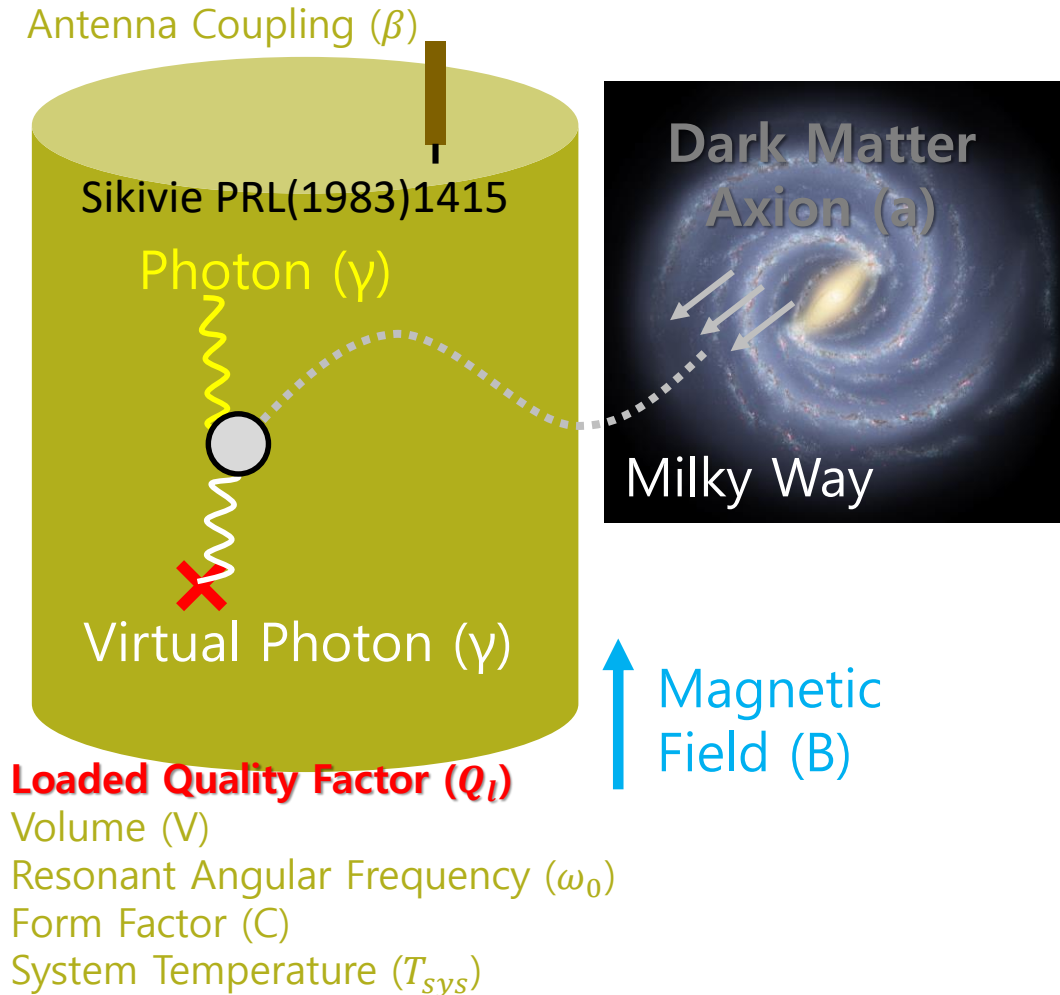
System Noise Temperature ~ 200 mK

$Q_l \gg Q_a \sim 10^6$

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Introduction

High Field and Big-Bore Magnet is Needed



Signal Power

$$P_{sig} = \frac{\beta}{1 + \beta} g_{\alpha\gamma\gamma}^2 \frac{\rho_a}{m_a^2} B^2 V \omega_0 C \frac{Q_a Q_l}{Q_a + Q_l}$$

Coupling Constant

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Kim et al. JCAP03(2020)066

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System Noise

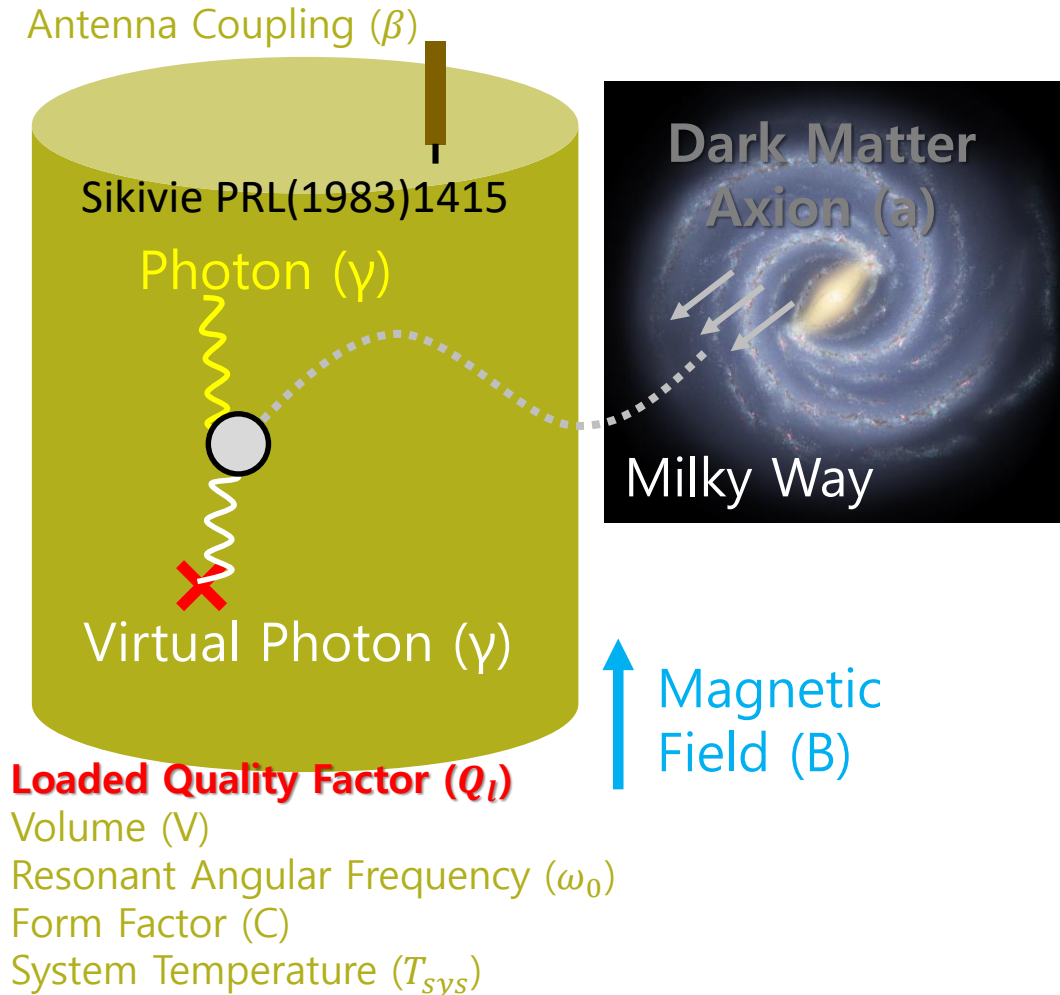
$Q_l \gg Q_a \sim 10^6$

High Magnetic Field, Cavity Volume

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Introduction

Quantum Noise Amplifier is Needed



Signal Power

$$P_{sig} = \frac{\beta}{1 + \beta} g_{\alpha\gamma\gamma}^2 \frac{\rho_a}{m_a^2} B^2 V \omega_0 C \frac{Q_a Q_l}{Q_a + Q_l}$$

Coupling Constant

Dark Matter Axion Density

Axion Mass

Axion Quality Factor

Kim et al. JCAP03(2020)066

Scanning speed

$$\frac{df}{dt} \propto \frac{B^4 V^2 C^2}{k_B^2 T_{sys}^2} Q_l Q_a$$

System Noise

$Q_l \gg Q_a \sim 10^6$

Low Noise Temperature

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Experiments at CAPP

Overview



Dr. Saebyeok Ahn's Slide

Main Axion eXperiment (CAPP-MAX)

Overview

- CAPP's flagship experiment to search for axion above 1GHz
- Dine-Fischler-Srednicki-Zhitnitsky (DFSZ) sensitivity



Dilution refrigerator 25mK



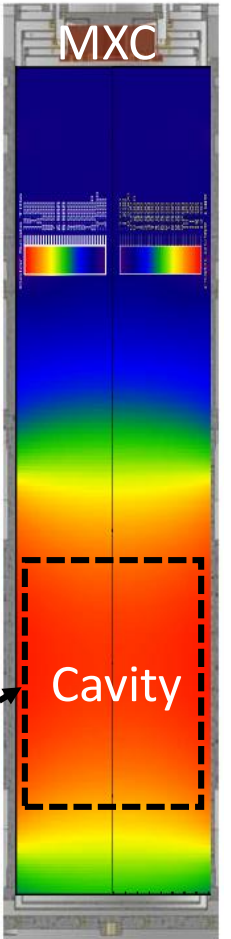
Josephson Parametric Amplifier

$$T_{sys} \approx 200 \text{ mK}$$



12 Tesla SC magnet
320mm bore diameter

Still shield



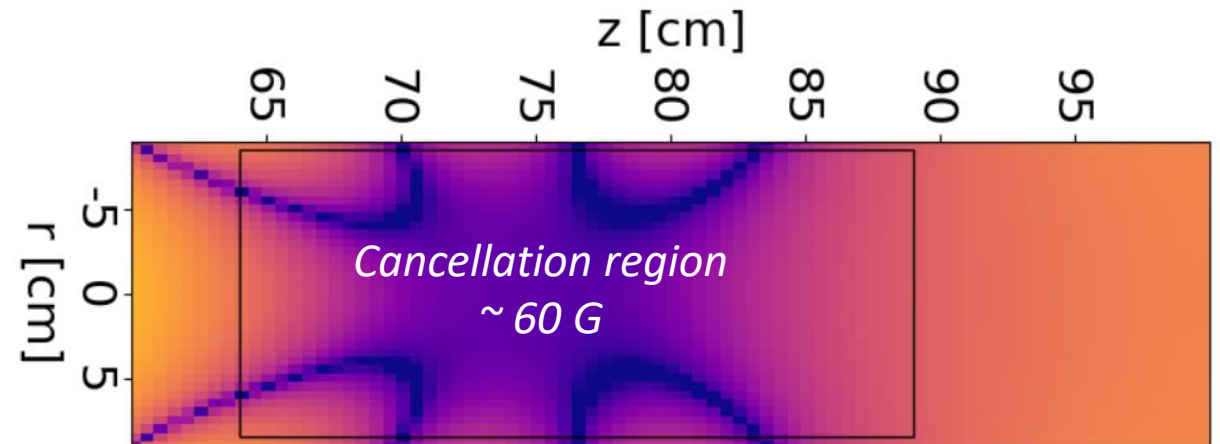
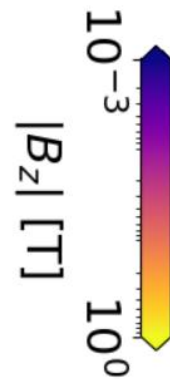
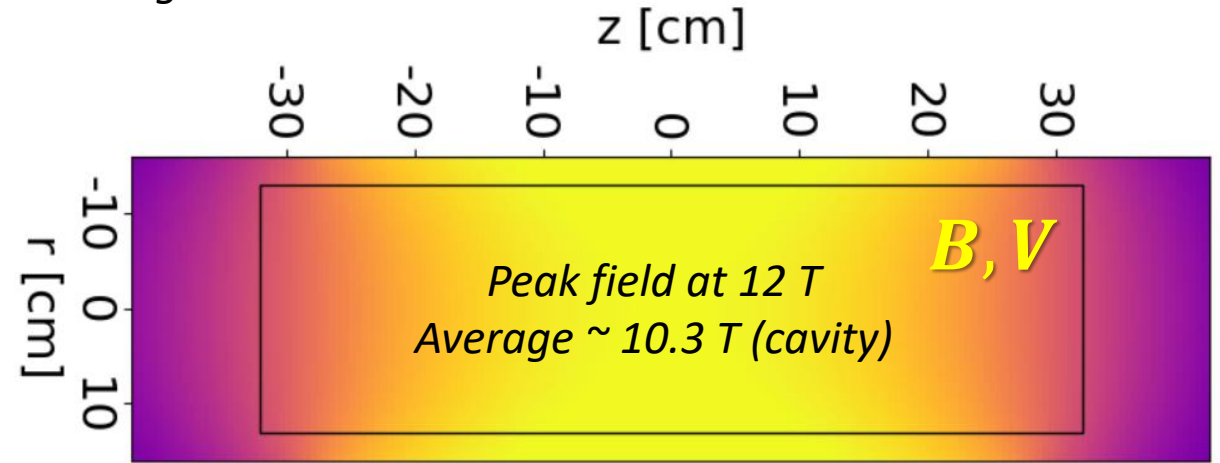
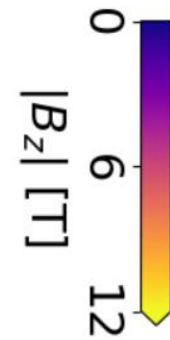
B-field map

Main Axion eXperiment (CAPP-MAX)

Magnet

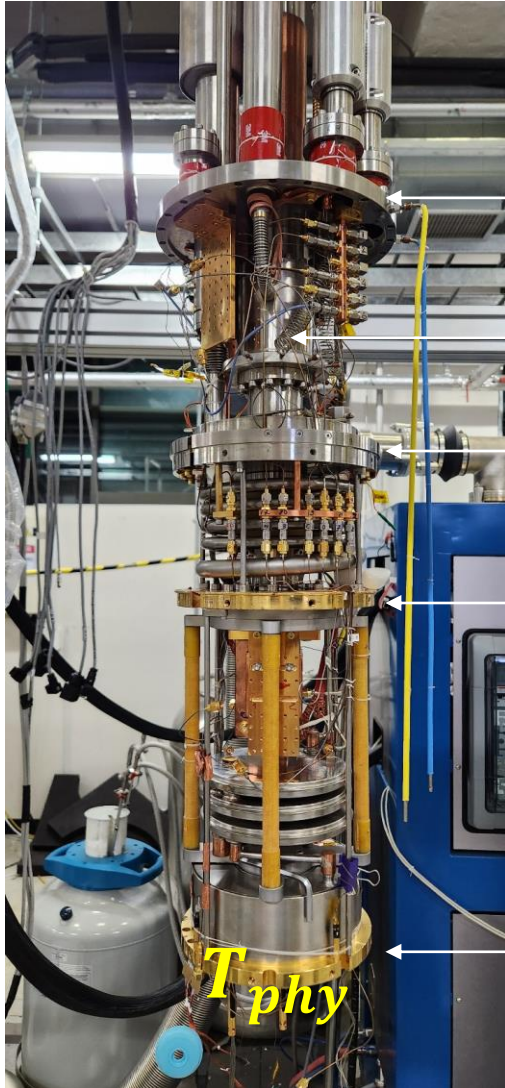


Nb_3Sn Superconducting coil



Main Axion eXperiment (CAPP-MAX)

The Cryogenic System



4 K plate

1 K pot (1.5 K)

STILL (600 mK)

Cold plate (100 mK)

MXC plate (20 mK)

Wet type LEIDEN dilution refrigerator

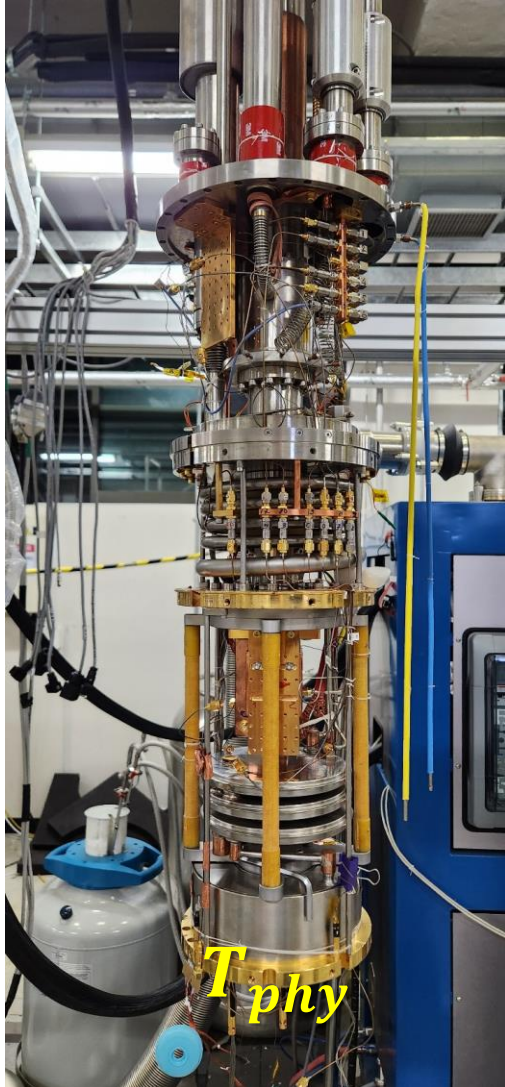
- 5.6 mK of base temperature (bare fridge)
- Cavity physical temperature ~ 30 mK
- Noise temperature ~ 40 mK

$$T_{cav} = \frac{h\nu}{k_B} \left(\frac{1}{e^{h\nu/k_B T_{phy}} - 1} + \frac{1}{2} \right)$$

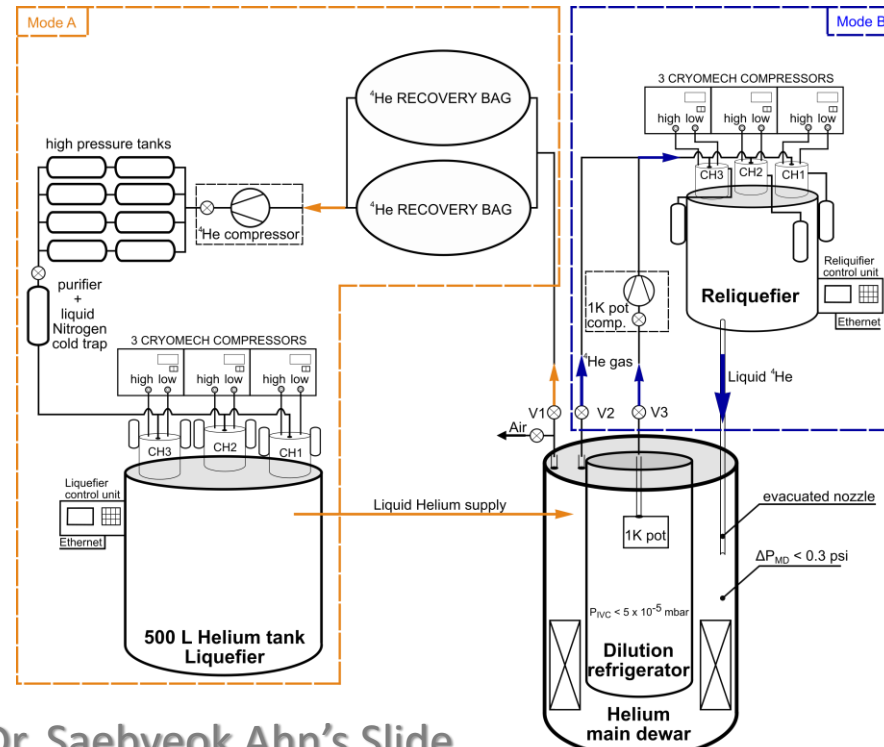
Dr. Saebyeok Ahn's Slide

Main Axion eXperiment (CAPP-MAX)

The Cryogenic System



- ✓ He recovery system ~ 160 L capacity
- ✓ He liquefier with 60 L / day rate
- ✓ He Re-liquefier with 60 L / day rate

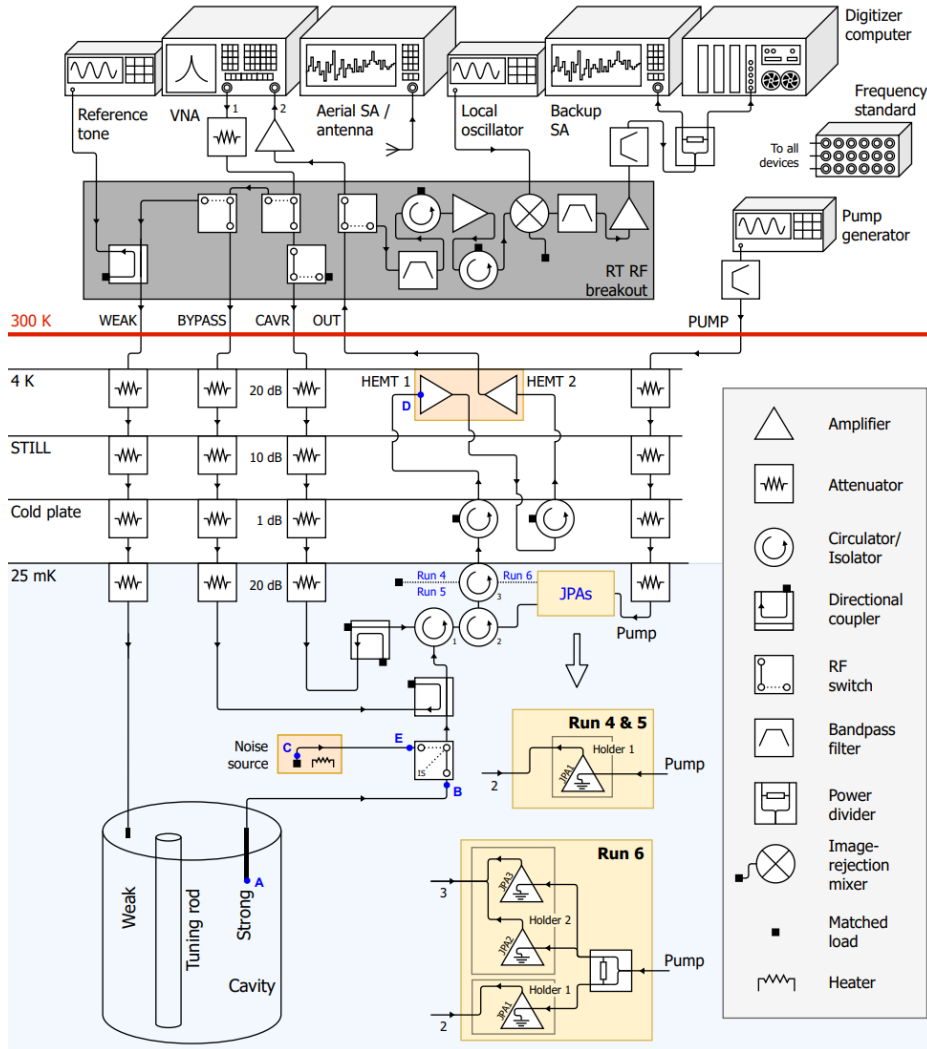


Dr. Saebyeok Ahn's Slide



Main Axion eXperiment (CAPP-MAX)

Receiver Chain



➤ Cavity

- ✓ Axion signal is generated in the cavity immersed in a magnetic field.
- ✓ Noise is also generated in the cavity. (Noise temperature ~ 40 mK)
- ✓ Signal and noise are picked up by strong port.

T_{phy}

➤ Josephson Parametric Amplifier (JPA)

- ✓ Quantum limited noise amplifier (~ 200 mK noise)
- ✓ Primary amplification
- ✓ Typical gain ~ 20 dB

➤ High Electron Mobility Transistor (HEMT)

- ✓ Low noise linear amplifier (~ 2 K / G_{jpa} noise)
- ✓ Subsequent amplification
- ✓ Typical gain ~ 60 dB

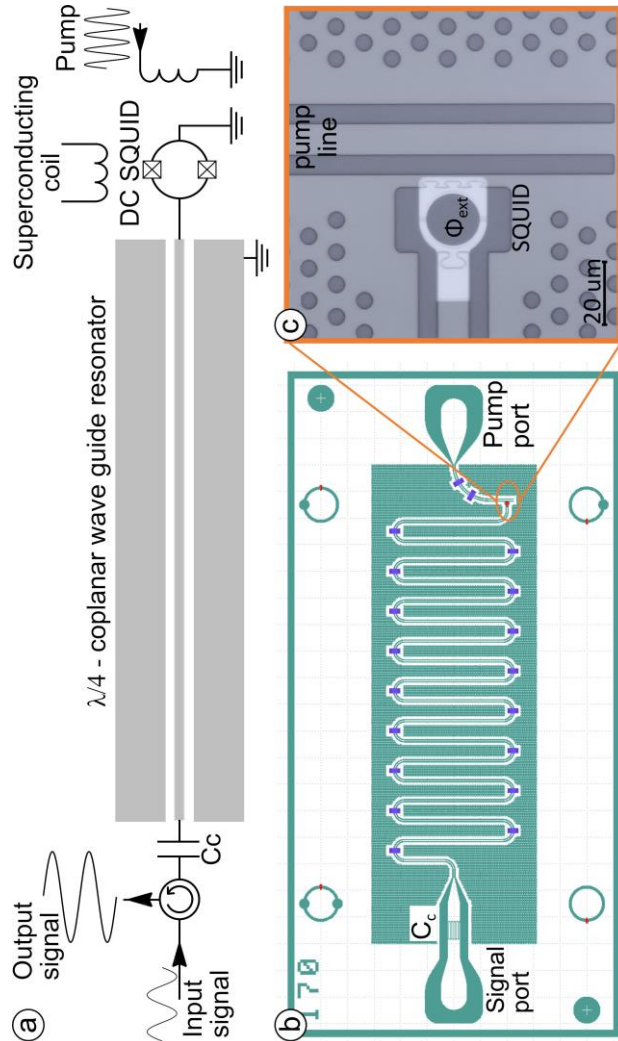
T_{sys}

➤ Noise source

- ✓ Heating up to 400 mK
- ✓ JPA-off calibration / System noise double-check

Main Axion eXperiment (CAPP-MAX)

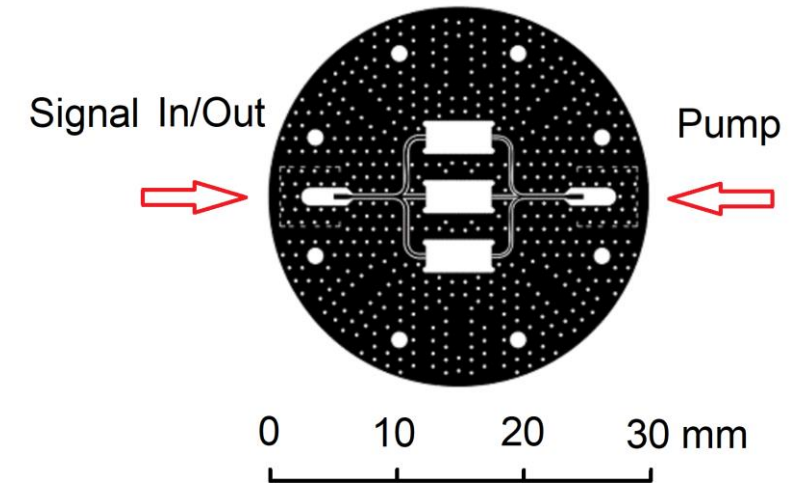
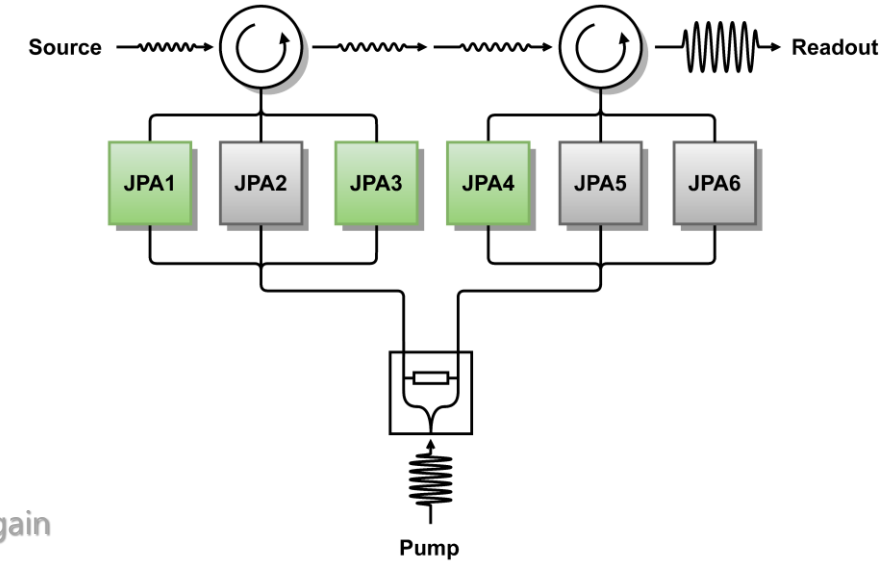
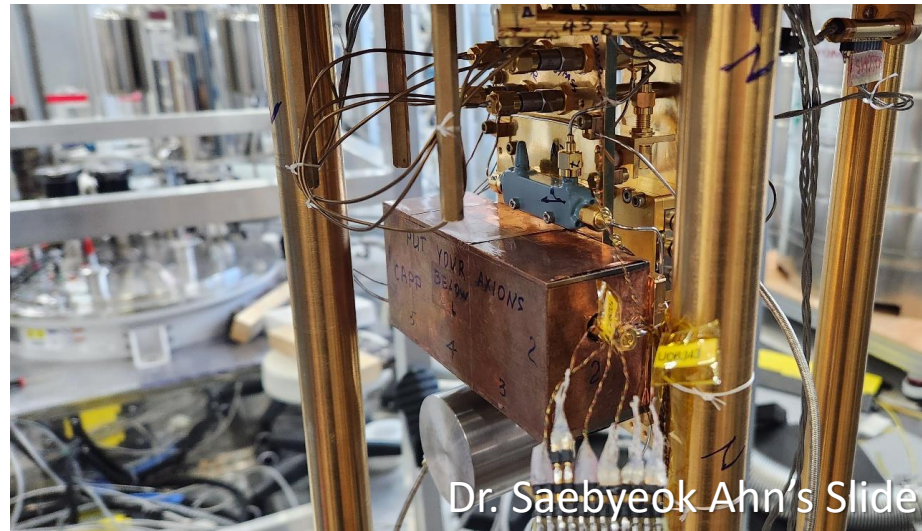
Josephson Parametric Amplifier



- ✓ Flux-driven JPAs
- ✓ Collaborating with U. of Tokyo
- ✓ 3 Parallel JPAs
- ✓ $T_{\text{sys}} \sim 200 \text{ mK}$

$$T_{\text{sys}} = T_{\text{cav}} + T_{\text{JPA}} + T_{\text{off}}/G_{\text{jpa}}$$

Total system noise \rightarrow T_{sys}
 Noise from cavity \rightarrow T_{cav}
 JPA added noise \rightarrow T_{JPA}
 JPA off noise \rightarrow T_{off}
 JPA gain \rightarrow G_{jpa}



High-temperature Superconducting Cavities at CAPP

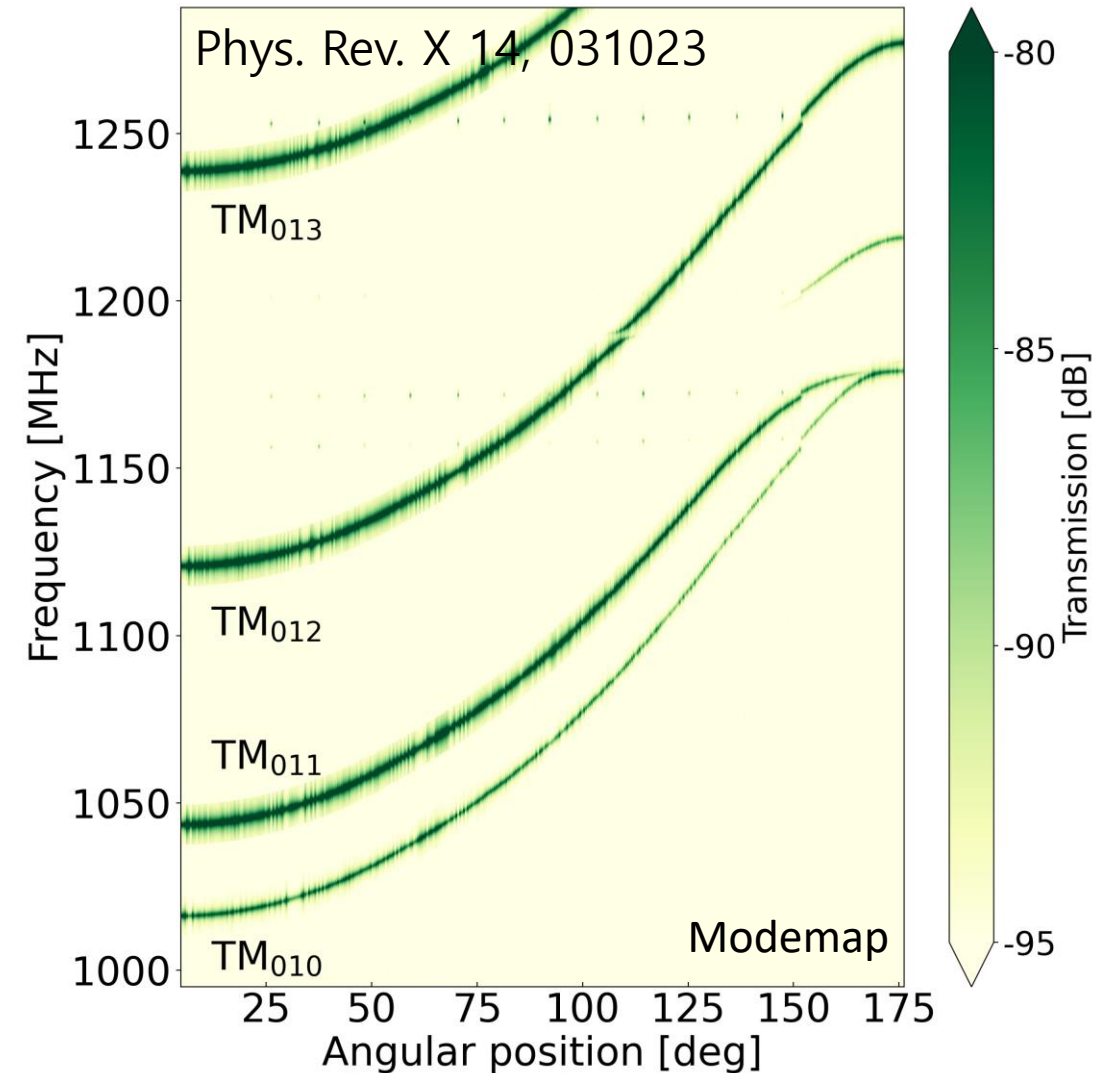
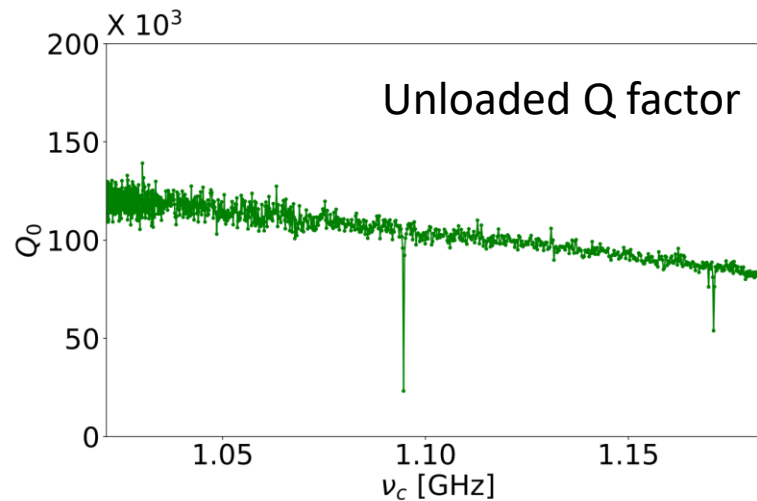
Superconducting Cavity for CAPP-MAX

Dr. Saebyeok Ahn's Slide

Total weight ~ **5 kg**

1.02 – 1.18 GHz tuning range

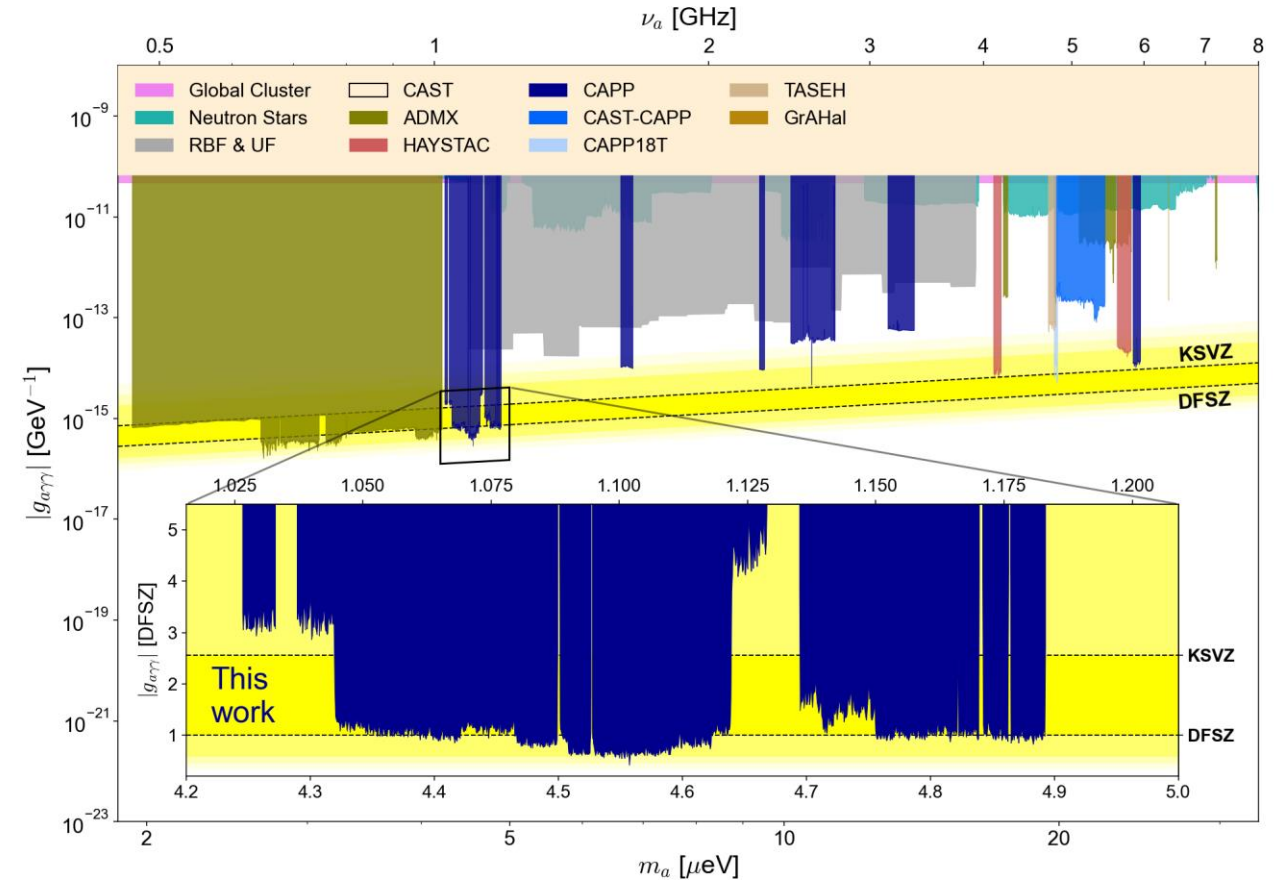
Geometrical factor ~ [0.5, 0.7]



High-temperature Superconducting Cavities at CAPP

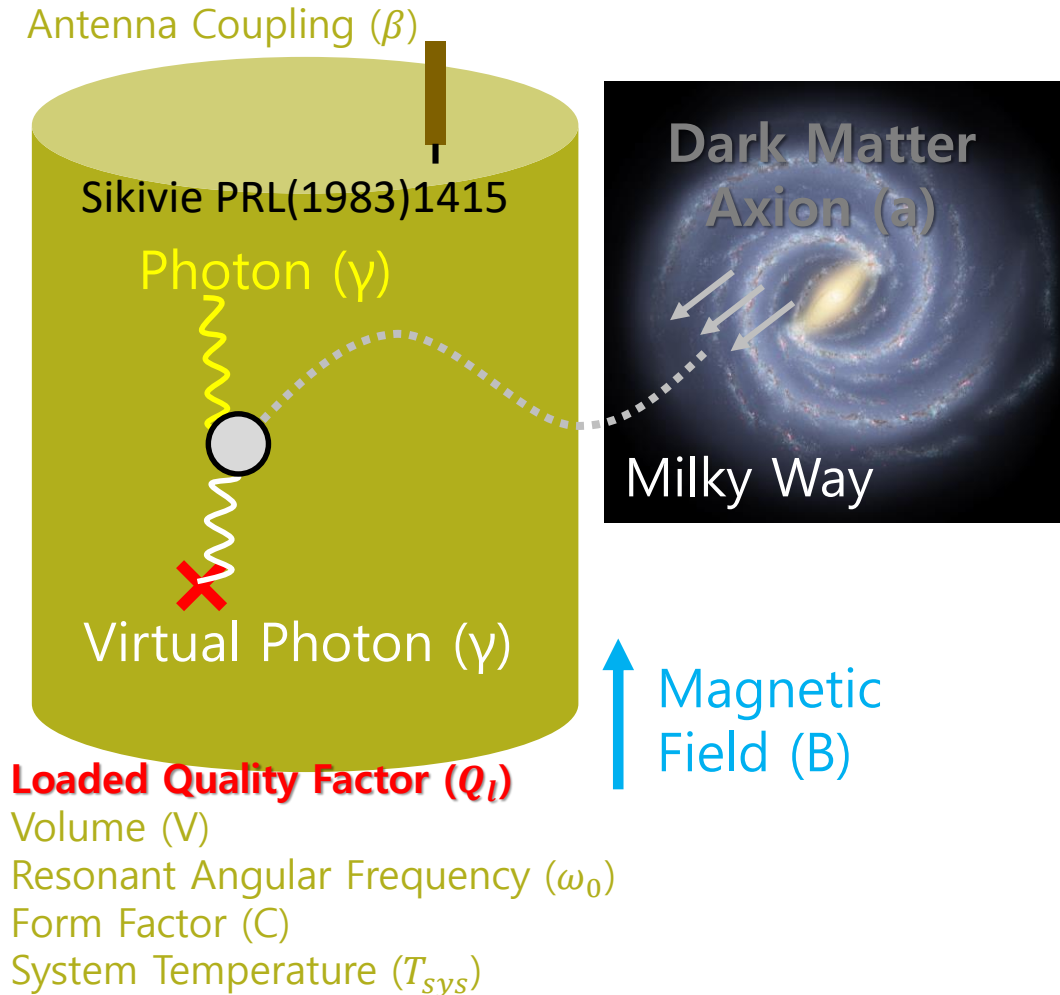
Superconducting Cavity for CAPP-MAX

Phys. Rev. X 14, 031023



Accelerating Axion Haloscope

Quality Factor is also Important to Enhance a Scanning Speed



Signal Power

$$P_{sig} = \frac{\beta}{1 + \beta} g_{\alpha\gamma\gamma}^2 \frac{\rho_a}{m_a^2} B^2 V \omega_0 C \frac{Q_a Q_l}{Q_a + Q_l}$$

Coupling Constant

Dark Matter Axion Density

Axion Mass

Axion Quality Factor

Kim et al. JCAP03(2020)066

Scanning speed

$$\frac{df}{dt} \propto \frac{B^4 V^2 C^2}{k_B^2 T_{sys}^2} Q_l Q_a$$

System Noise

$Q_l \gg Q_a \sim 10^6$

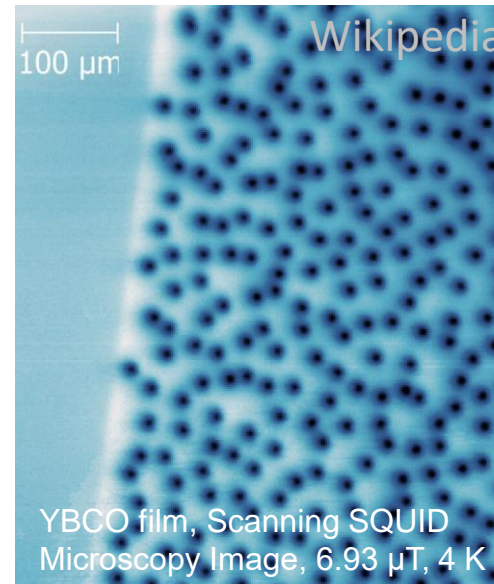
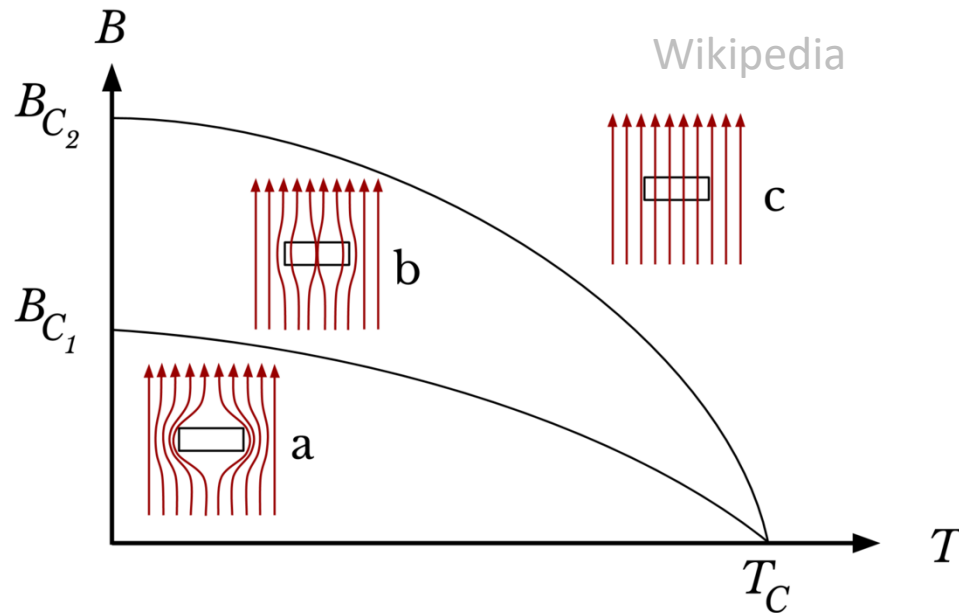
High Cavity Quality Factor

Dr. Jinsu Kim, Patras Workshop, 2023

Type II Superconductor in a Magnetic Field

Vortex Dynamics & Energy Dissipation Mechanism

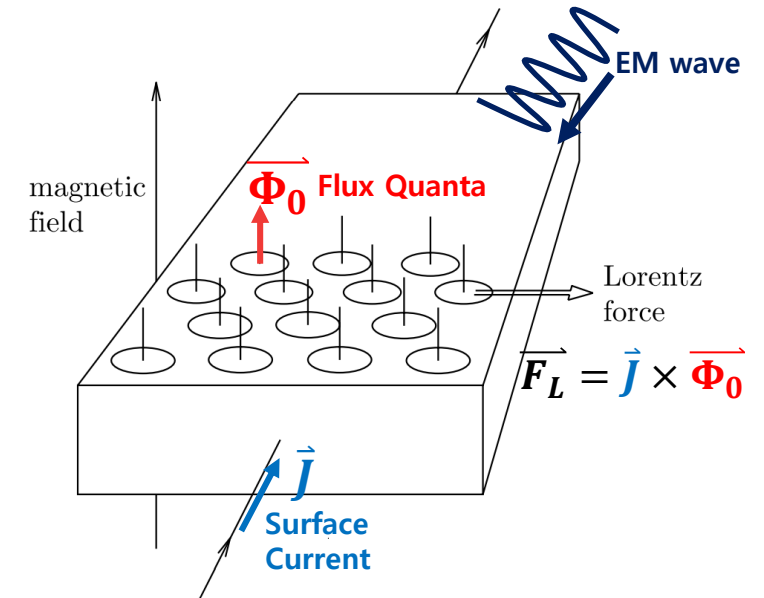
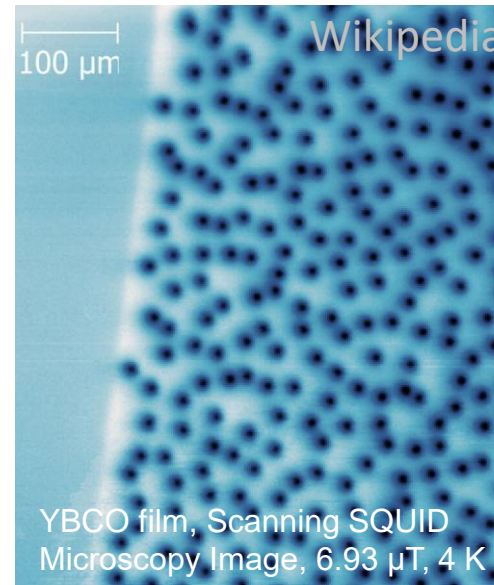
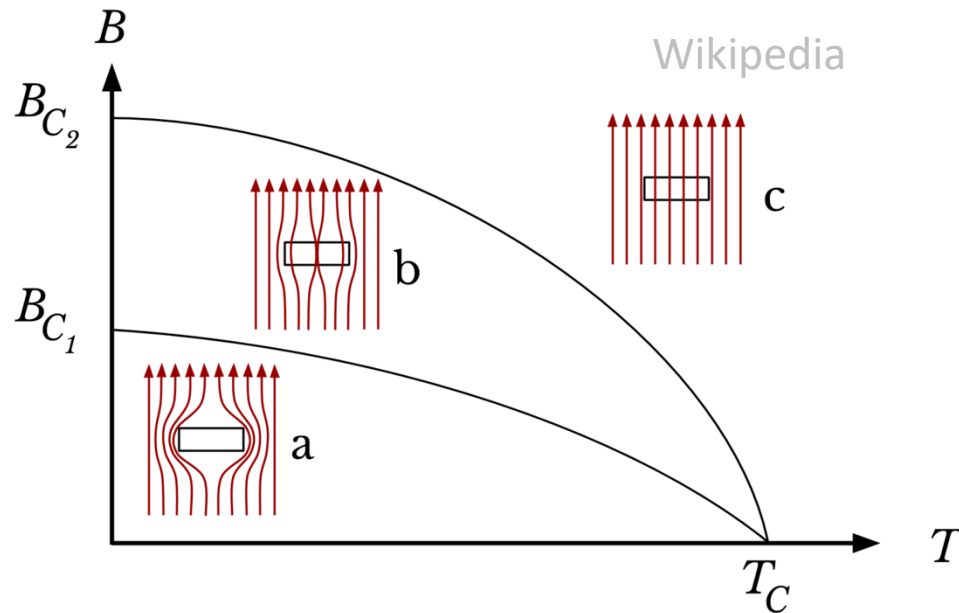
- Type II superconductor forms vortices (quantized magnetic flux, $\vec{\Phi}_0$) in a magnetic field.



Type II Superconductor in a Magnetic Field

Vortex Dynamics & Energy Dissipation Mechanism

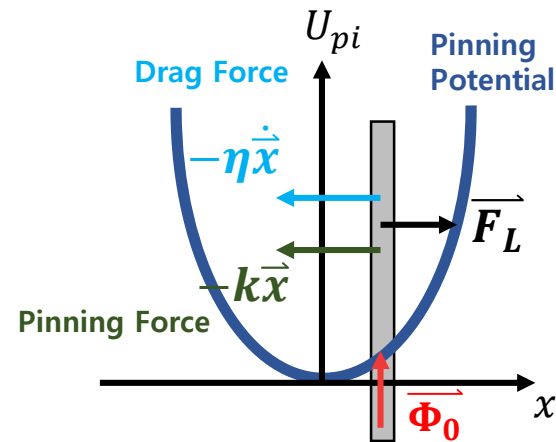
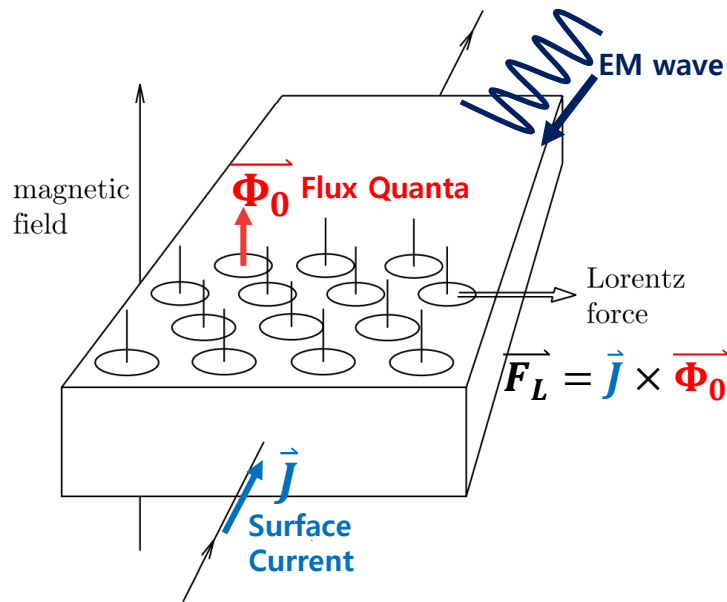
- Type II superconductor forms vortices (quantized magnetic flux, $\overline{\Phi_0}$) in a magnetic field.
- Vortices respond to the incident electromagnetic wave.



Type II Superconductor in a Magnetic Field

Vortex Dynamics & Energy Dissipation Mechanism

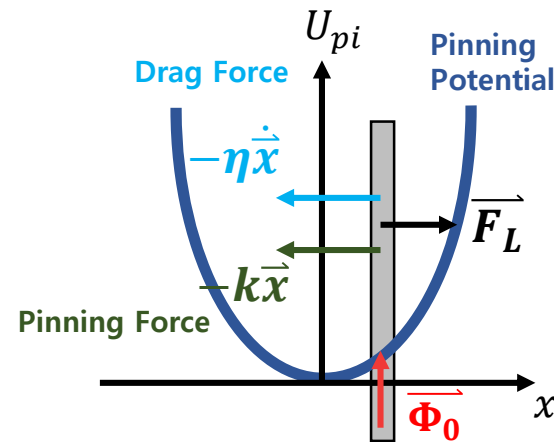
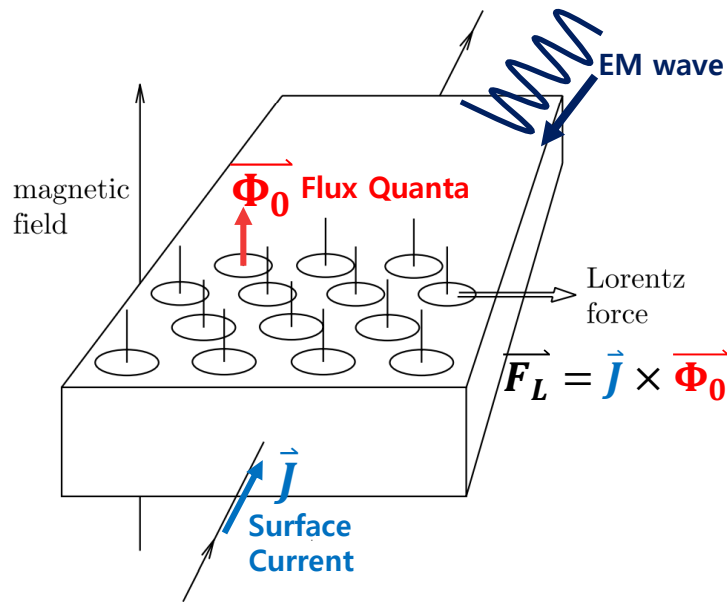
- Type II superconductor forms vortices (quantized magnetic flux, Φ_0) in a magnetic field.
- Vortices respond to the incident electromagnetic wave.
- Every vortex is trapped in each pinning potential well.



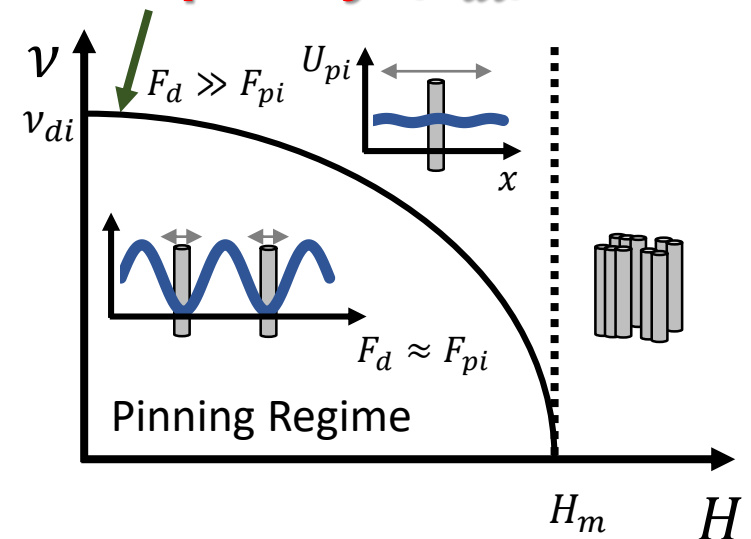
Type II Superconductor in a Magnetic Field

Vortex Dynamics & Energy Dissipation Mechanism

- Type II superconductor forms vortices (quantized magnetic flux, Φ_0) in a magnetic field.
- Vortices respond to the incident electromagnetic wave.
- Every vortex is trapped in each pinning potential well.
- Above a depinning frequency, pinning force become negligible.



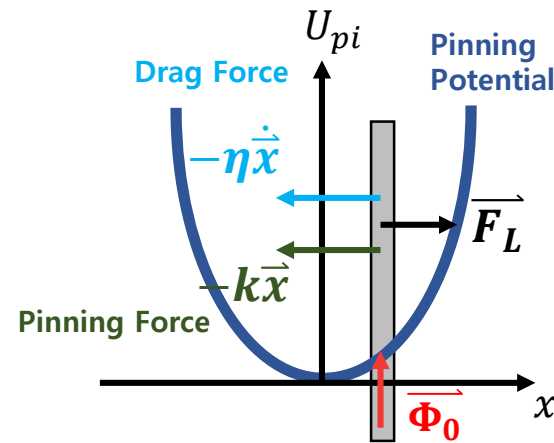
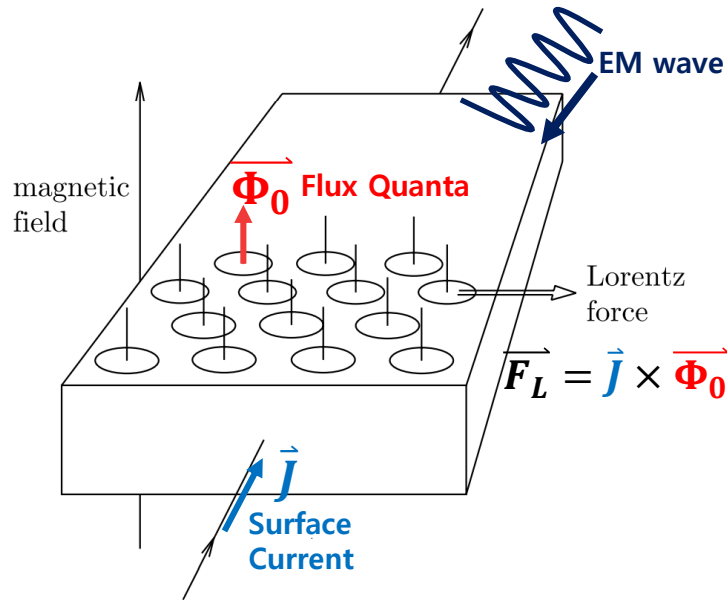
High Depinning Frequency (ν_{di})



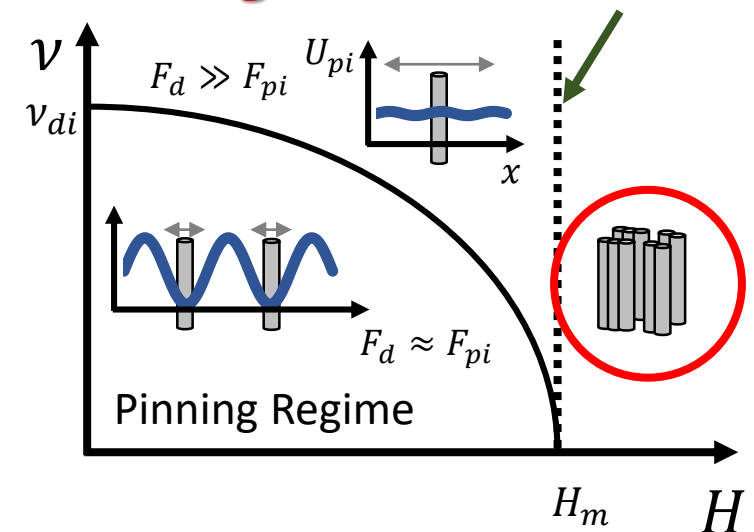
Type II Superconductor in a Magnetic Field

Vortex Dynamics & Energy Dissipation Mechanism

- Type II superconductor forms vortices (quantized magnetic flux, Φ_0) in a magnetic field.
- Vortices respond to the incident electromagnetic wave.
- Every vortex is trapped in each pinning potential well.
- Above a melting field, vortices are mixed each other.



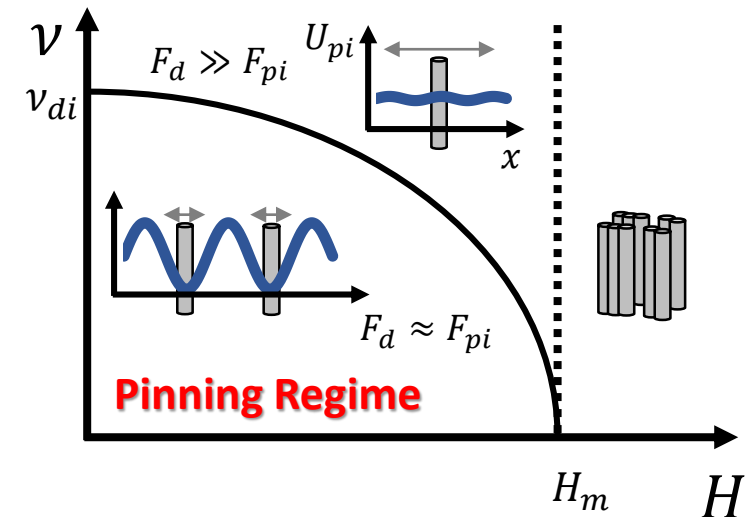
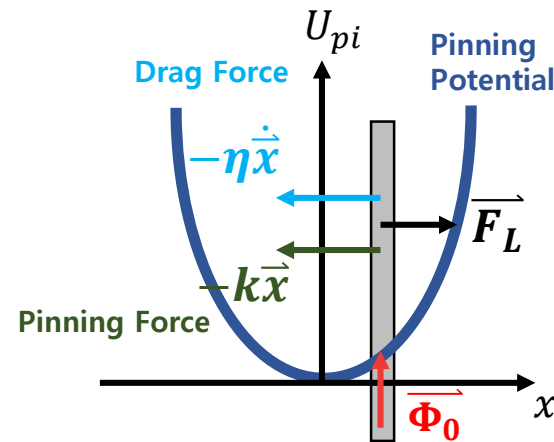
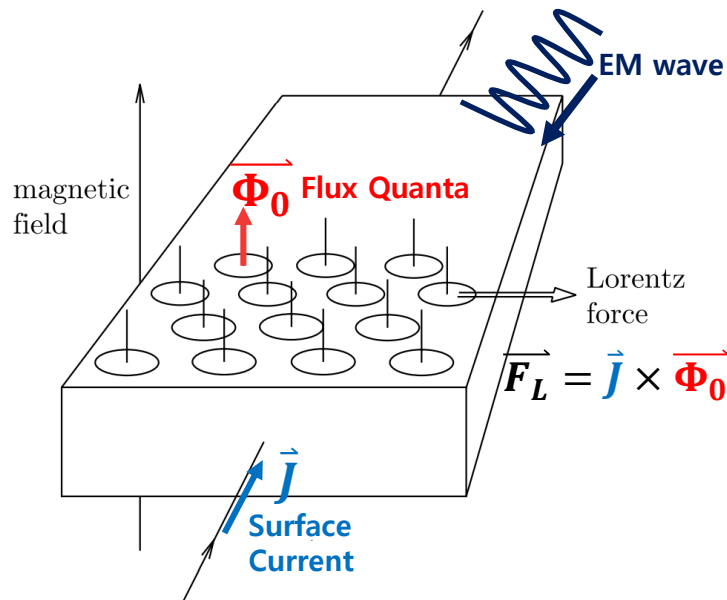
High Melting Field (H_m)
 \sim High Critical Field (H_{c2})



Type II Superconductor in a Magnetic Field

Vortex Dynamics & Energy Dissipation Mechanism

- Type II superconductor forms vortices (quantized magnetic flux, Φ_0) in a magnetic field.
- Vortices respond to the incident electromagnetic wave.
- Every vortex is trapped in each pinning potential well.
- In the pinning regime, the surface resistance is much smaller than that of copper.



Type II Superconductor in a Magnetic Field

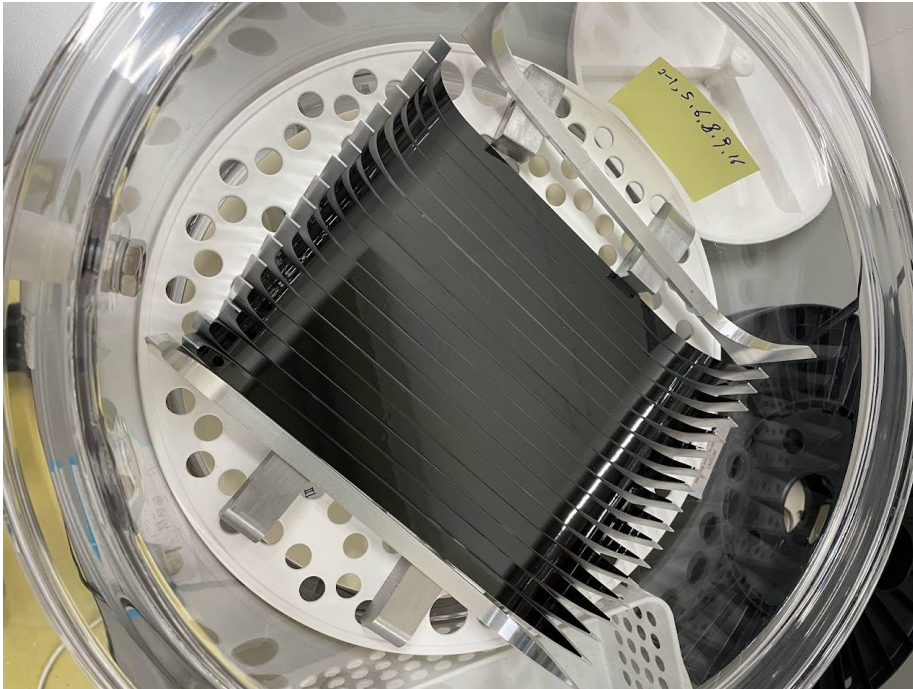
Material Evaluation

100 mK 8 GHz	R_s (B = 0 T) (Ohm)	R_s (B = 8 T, c) (Ohm)	Critical Field (H_{c2})	Depinning Frequency
OFHC Cu (Metal)	~ 7E-3	~ 7E-3	None	None
Low Temperature Superconductors (LTS)				
NbTi (LTS) <small>Gatti et al. PRD(2019)</small>	~ 1E-6	~ 4e-3	Small ~ 13 T	~ 45 GHz (?)
Nb ₃ Sn (LTS) <small>Alimenti et al. SUST(2020)</small>	~ 1E-6	?	~ 25 T	Small ~ 6 GHz
High Temperature Superconductors (HTS)				
Bi-2212 (HTS) Bi-2223 (HTS)	~ 1E-5	?	> 100 T (ab) <small>Larbalestier et al. Nature(2001)</small>	?
Tl-1223 (HTS)	~ 1E-5	~ 1e-4 <small>Calatroni et al. SUST(2017)</small>	> 100 T (ab) <small>Larbalestier et al. Nature(2001)</small>	12 – 480 MHz <small>Calatroni et al. SUST(2017)</small>
ReBCO (HTS) Re = Y, Gd, Eu, ...	~ 1E-5 <small>Ormeno et al. PRB(2001)</small>	~ 1e-4 <small>Romanov et al. Scientific Reports(2020)</small>	> 100 T (ab) <small>Larbalestier et al. Nature(2001)</small>	Strong Pinning 10 – 100 GHz <small>Romanov et al. Scientific Reports(2020)</small>

High-temperature Superconducting Cavity R&D

Development Methods

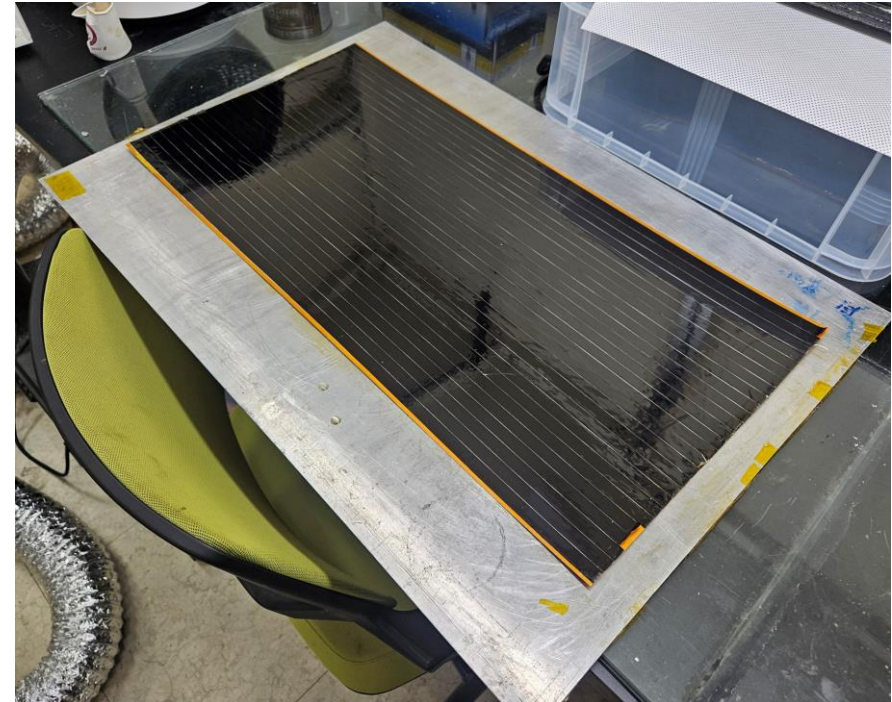
Strategy 1



Pros: Clean Surface

Cons: Slow Fabrication, Electrically Disconnected

Strategy 2



Pros: Easy to Fabricate, Electrically Connected

Cons: Surface Defect (~ 10%)

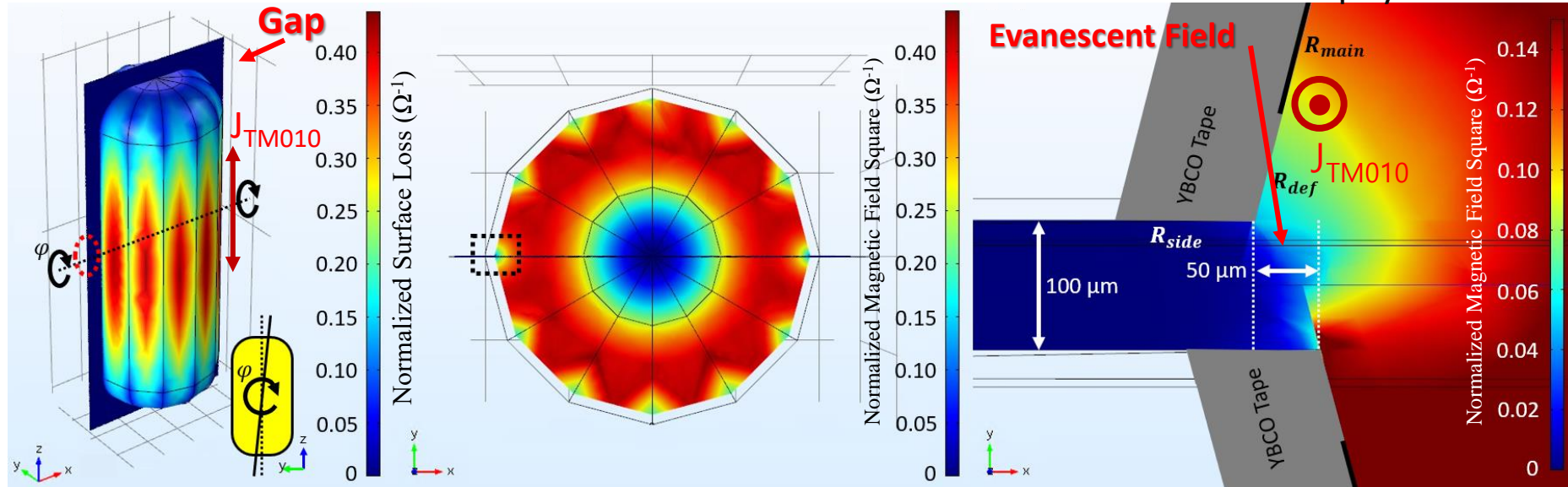
High-temperature Superconducting Cavity R&D

Strategy 1

- Simulation study can estimate energy loss from gaps.
- Misalignments and defects are considered based on fabrication error.
- Only evanescent field enter into a gap.

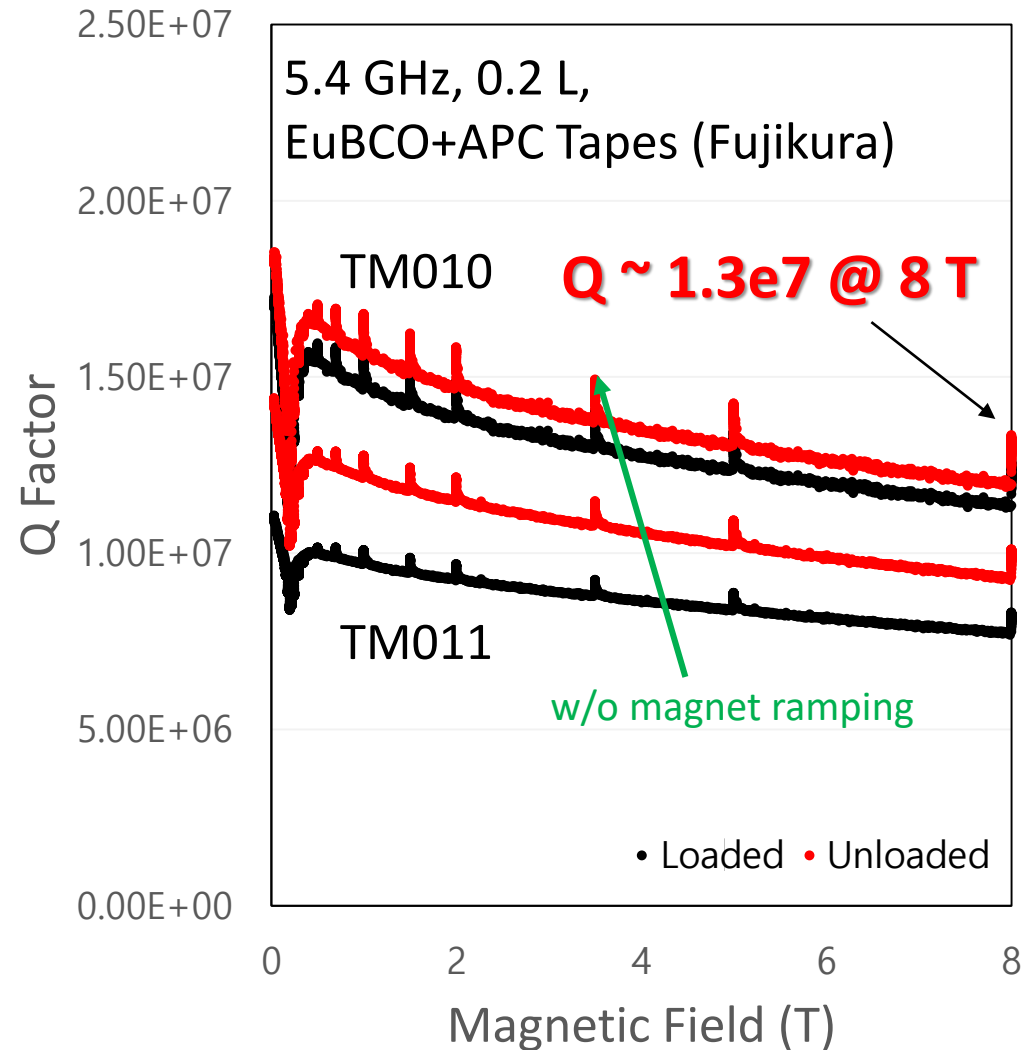
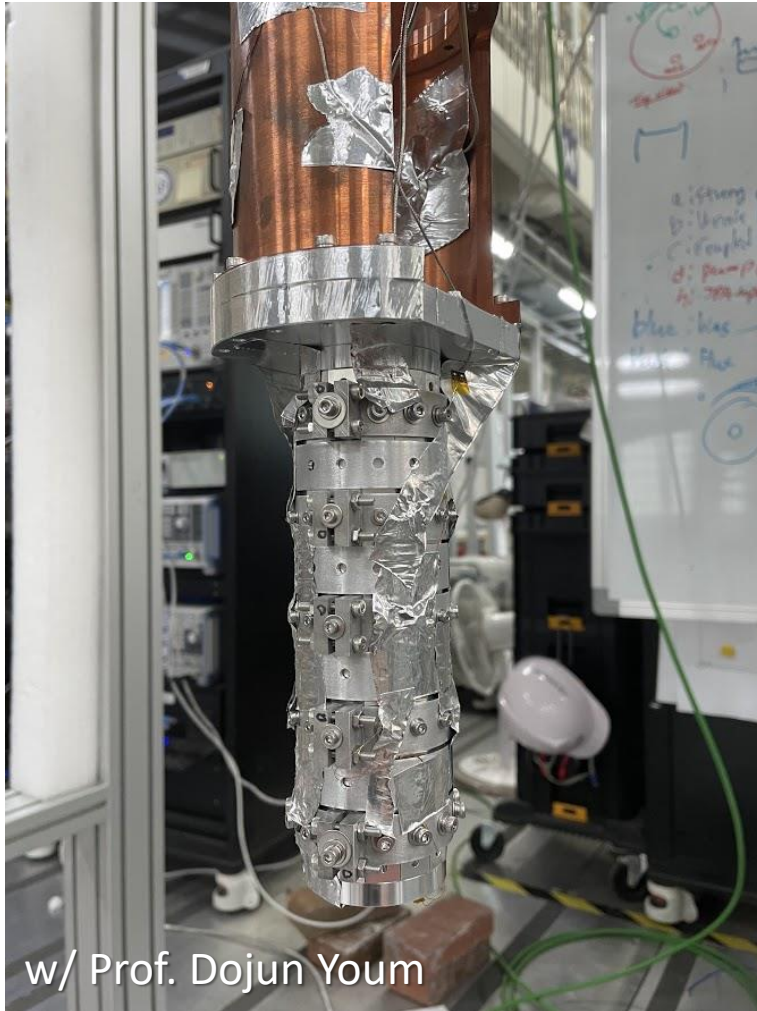
$$Q_{\text{ext}} > 10^7$$

Ahn *et al.* PRApplied(2022), 17, L061005



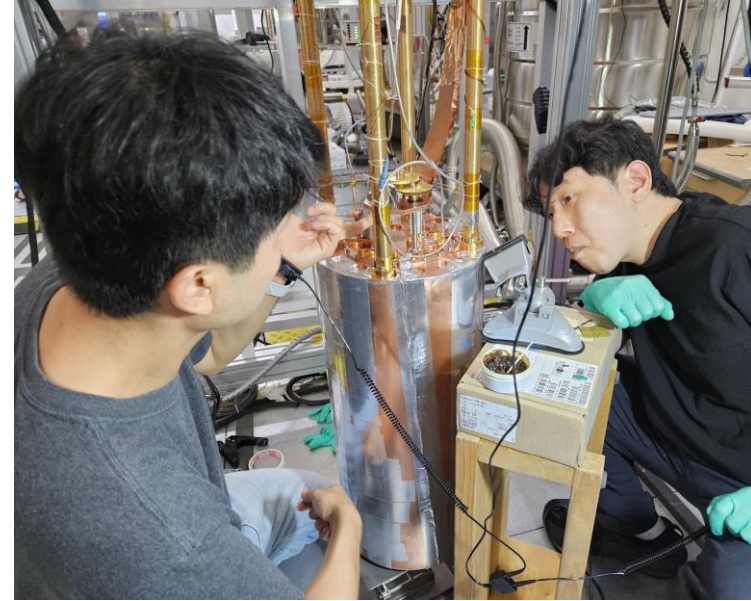
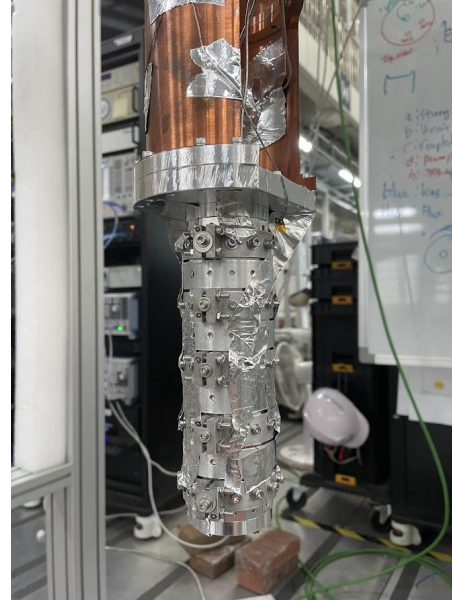
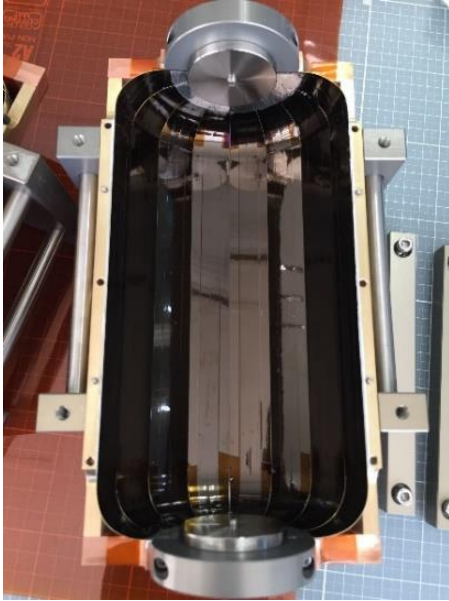
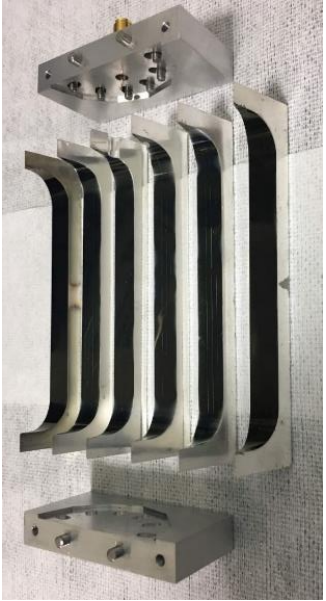
High-temperature Superconducting Cavity

CAPP Successfully Achieved 13 Million Quality Factor Cavity at 8 T Field



High-temperature Superconducting Cavities at CAPP

Summary

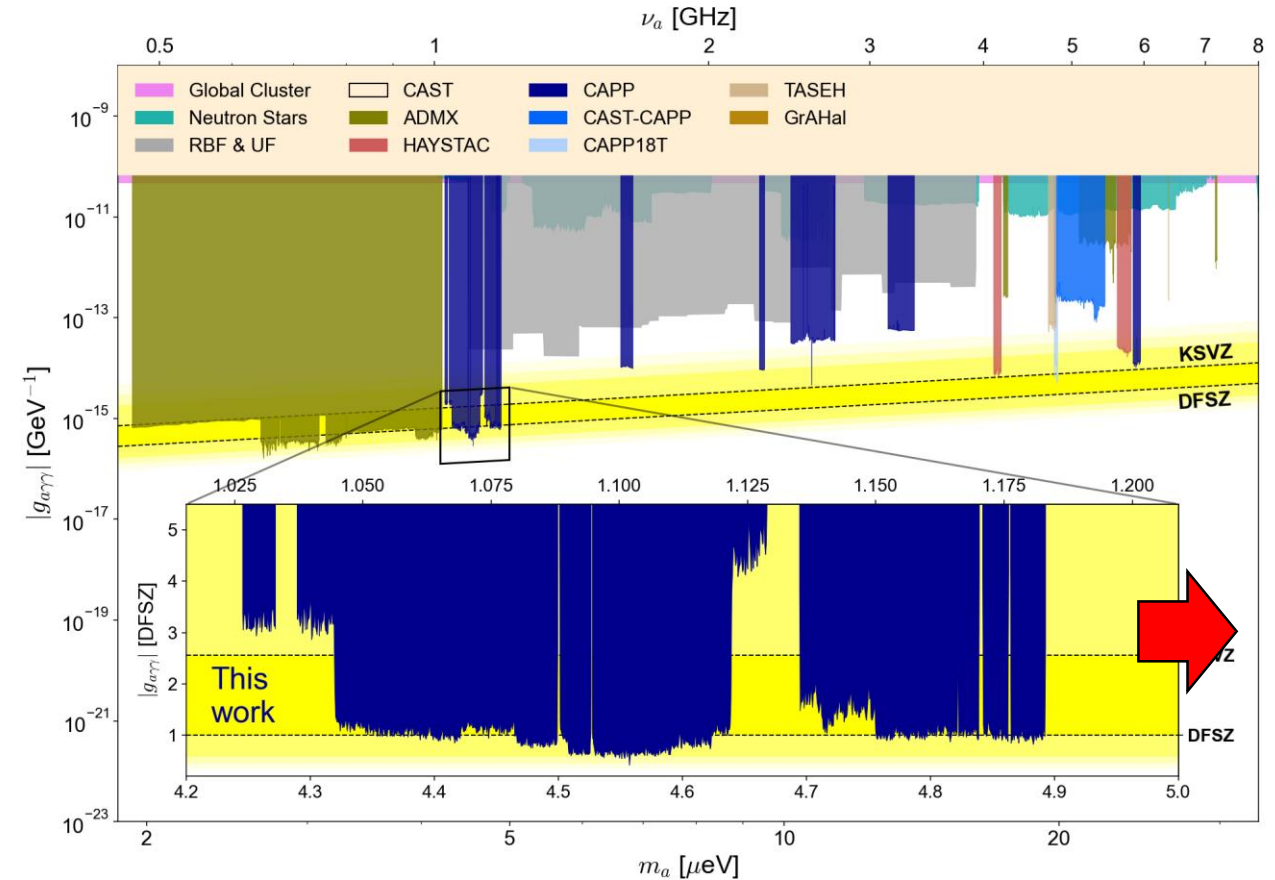


Date	Tape	f (GHz)	n_{gap}	Q (0 T)	Q (8 T)	Q_{gap}	Experiment
1	YBCO	6.9	12	0.22 M	0.33 M	> 10 M	Prototype
2	GdBCO	2.3	32	0.60 M	0.50 M		Cavity Haloscope
3	EuBCO+APC	2.3	34	5.0 M	3.5 M		Axion Quark Nugget Search
4	EuBCO+APC	5.4	14	20 M	13 M		Cavity Haloscope
5	EuBCO+APC	1.2 ~ 1.5	0	1.1 M	0.8 M		Axion Haloscope (CAPP-MAX)

High-temperature Superconducting Cavities at CAPP

Superconducting Cavity for CAPP-MAX

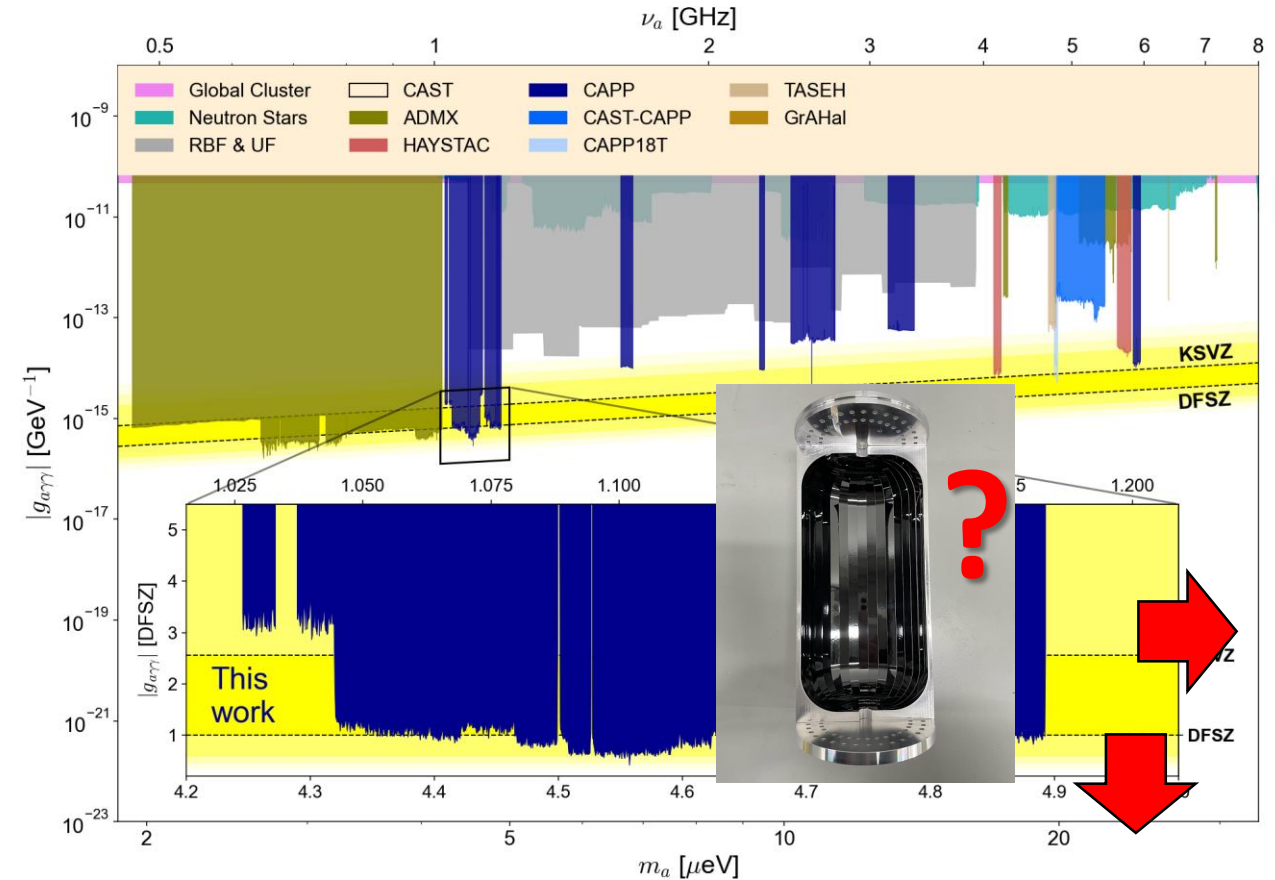
Phys. Rev. X 14, 031023



High-temperature Superconducting Cavities at CAPP

Superconducting Cavity for CAPP-MAX

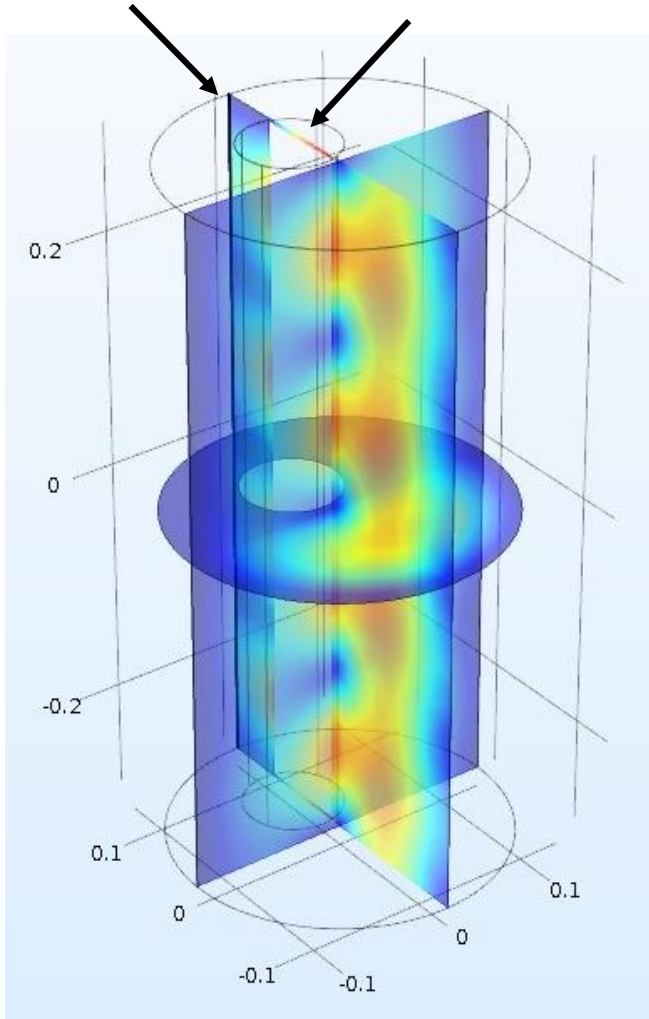
Phys. Rev. X 14, 031023



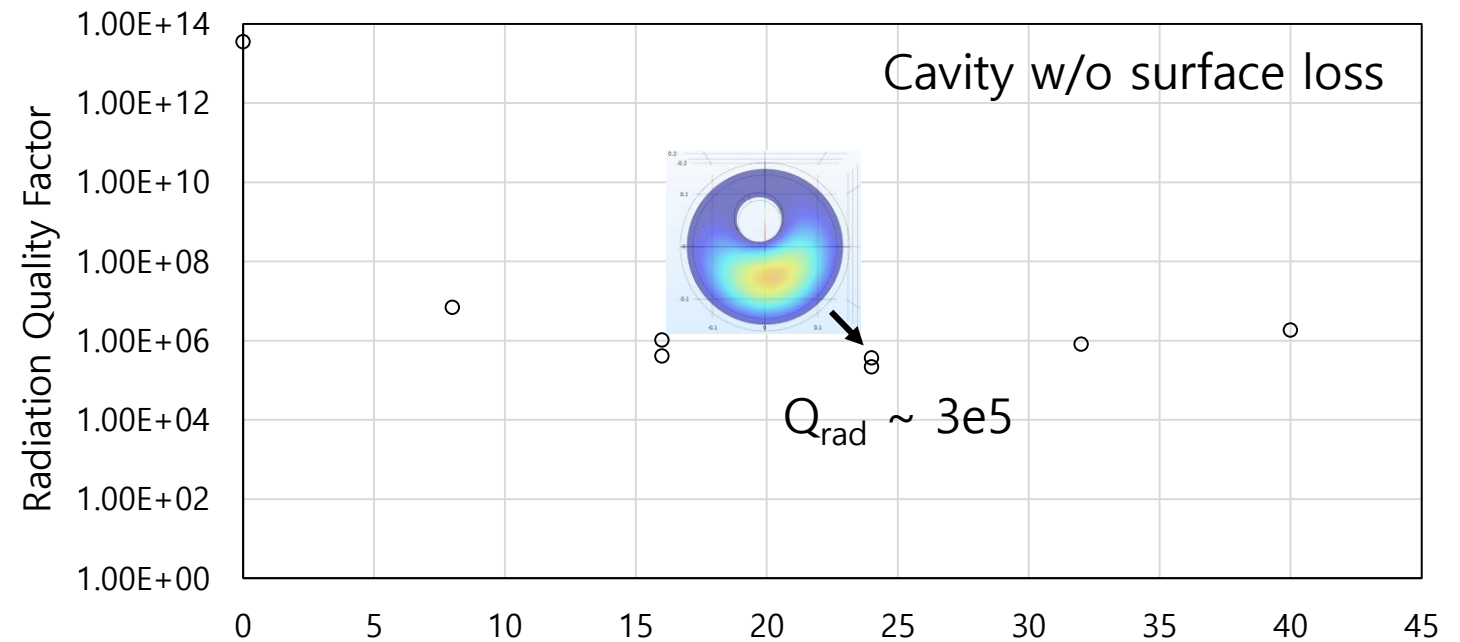
Below DFSZ sensitivity or Axion as 20% of Dark Matter

High-temperature Superconducting Cavity R&D

Gap geometry degrades Q factor with tuning rods



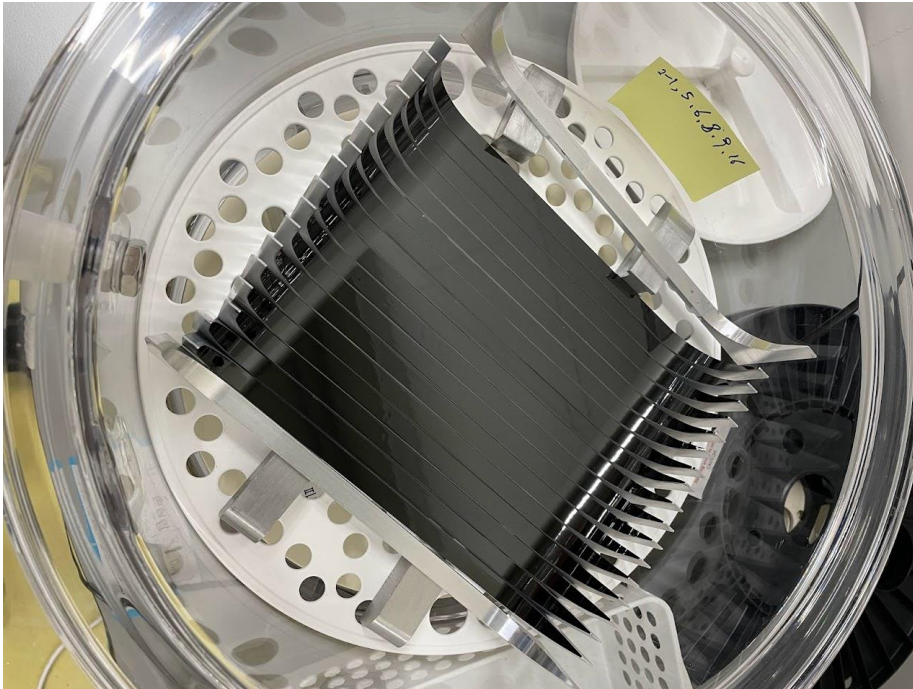
- Sensitive to Contact Problem
- Gaps should be closed electrically to prevent a radiation



High-temperature Superconducting Cavity R&D

Development Methods

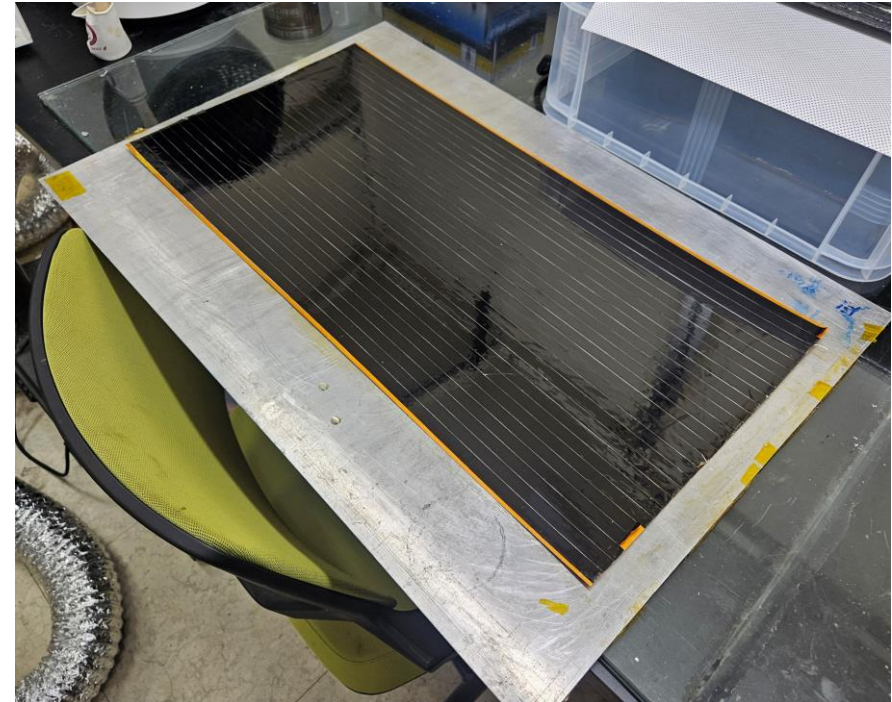
Strategy 1



Pros: Clean Surface

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Strategy 2

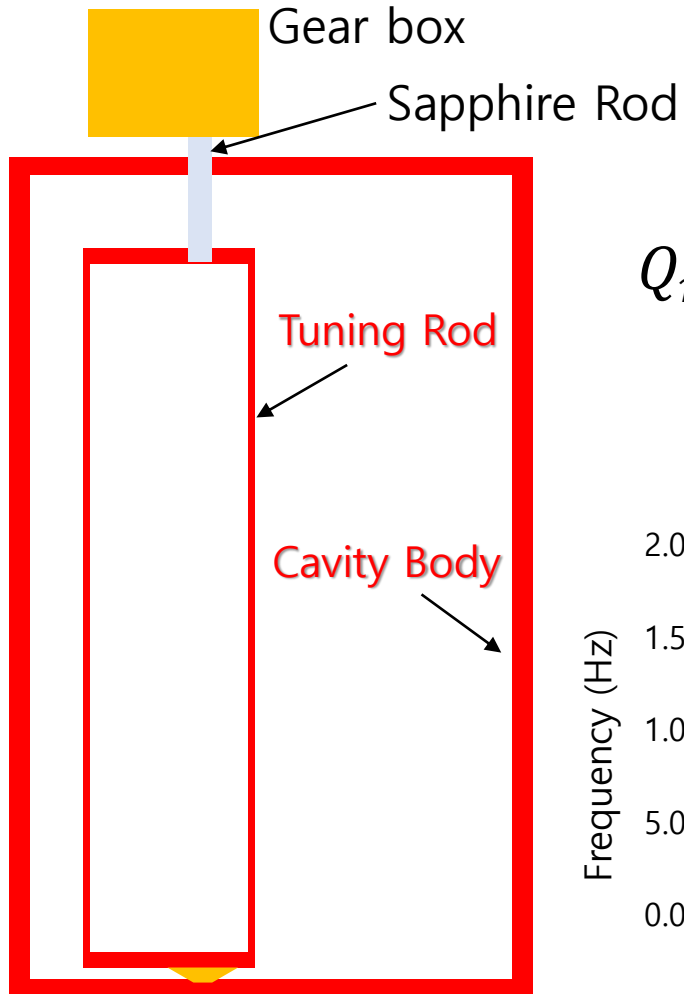


Pros: Easy to Fabricate, Electrically Connected

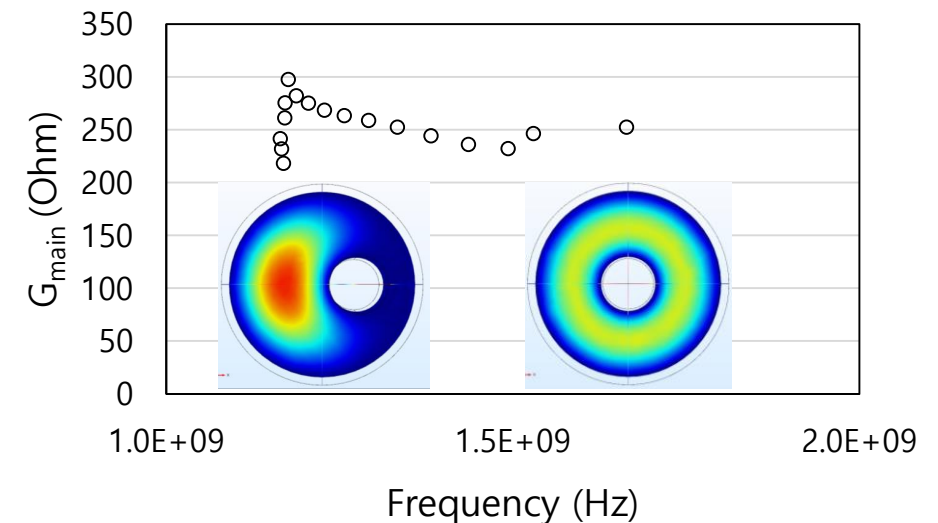
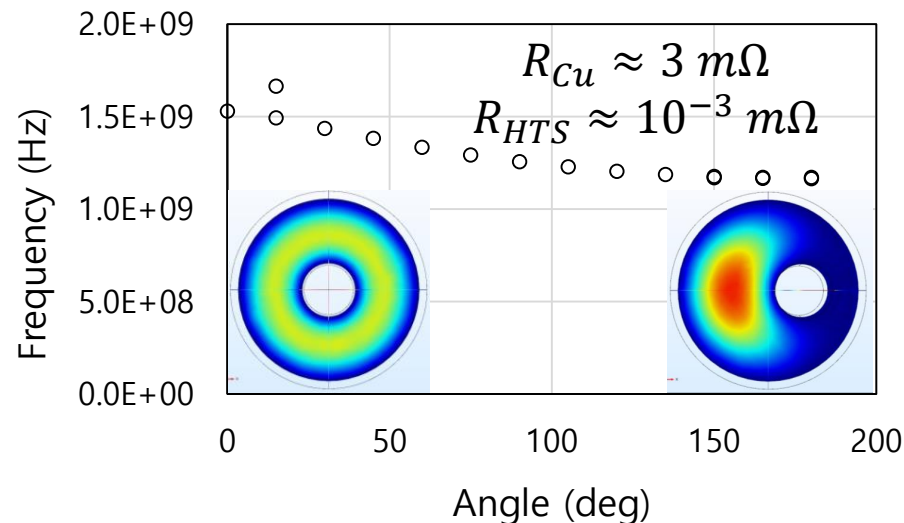
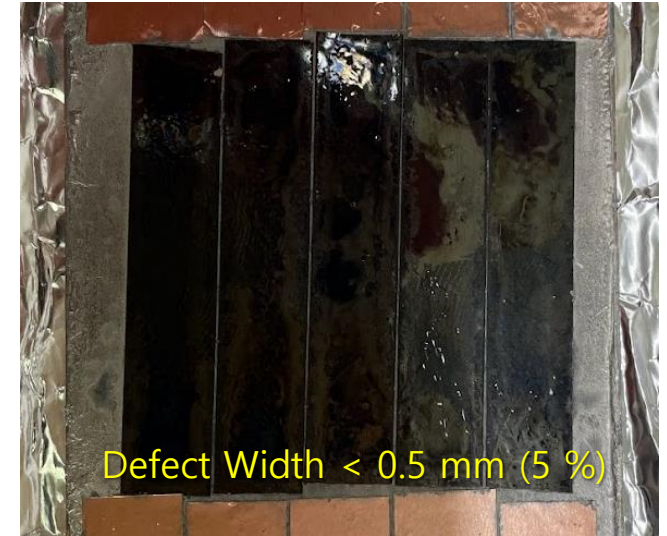
Cons: Surface Defect (~ 10%)

High-temperature Superconducting Cavity R&D

Strategy 2

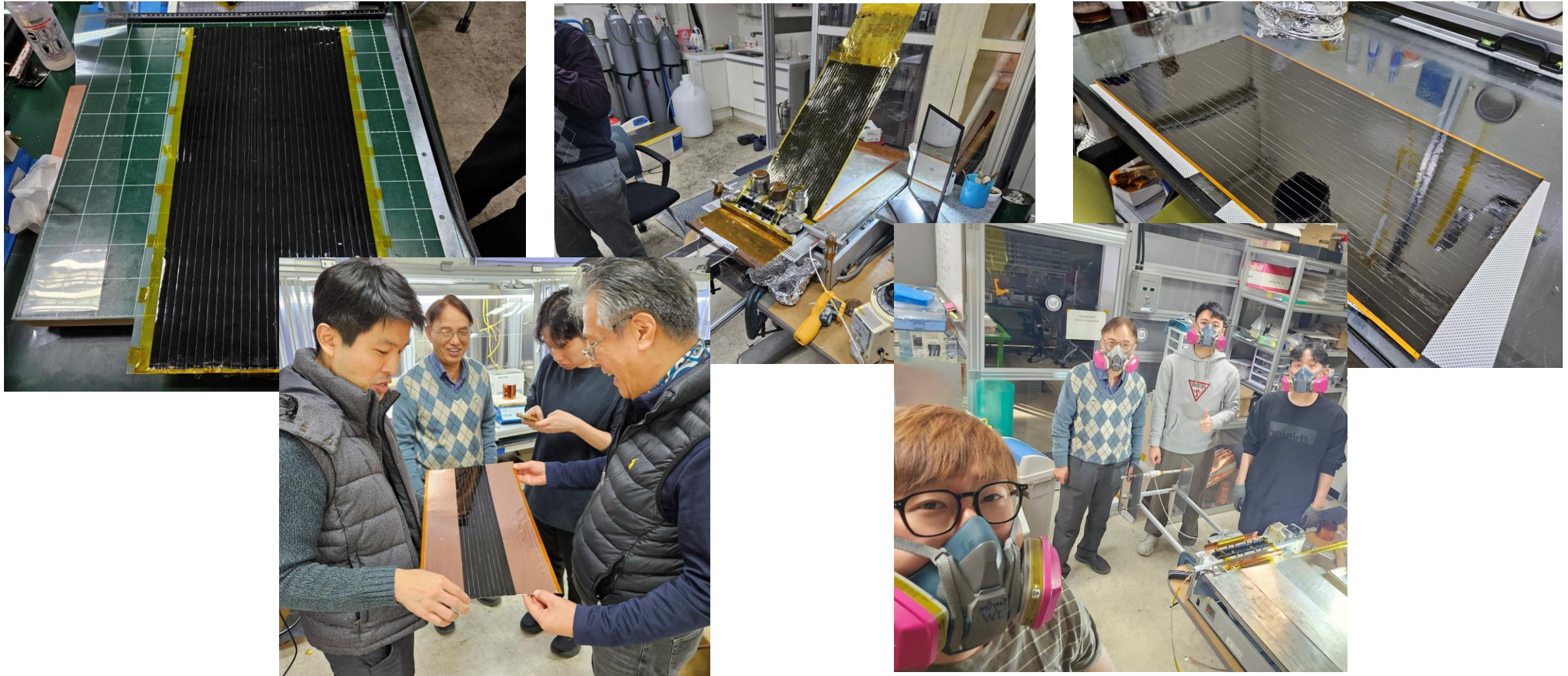


$$Q_{main} \approx \frac{250\Omega}{3m\Omega \times 10\%} = \mathbf{8 \times 10^5}$$



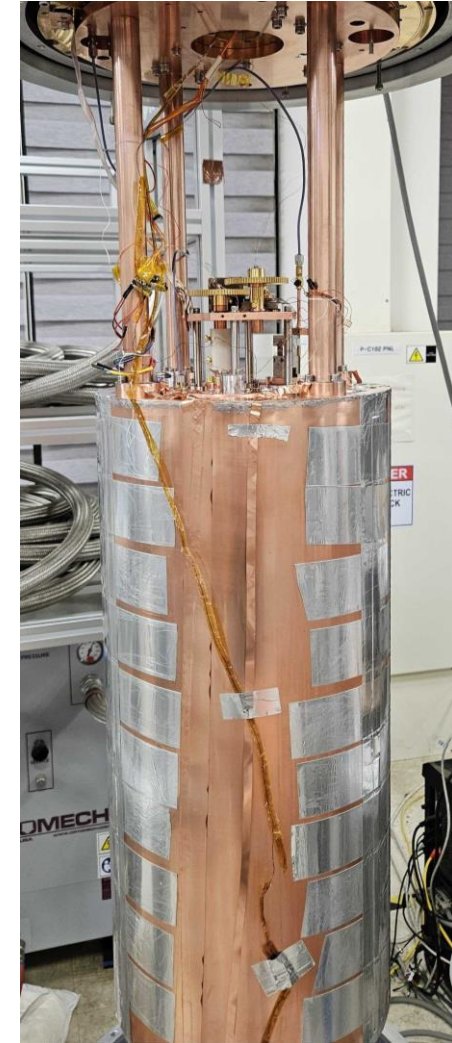
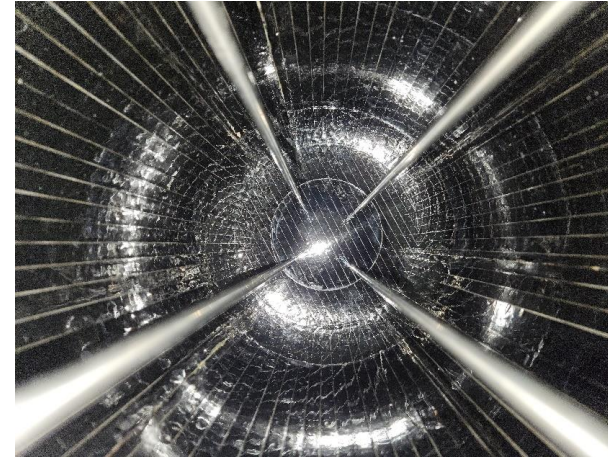
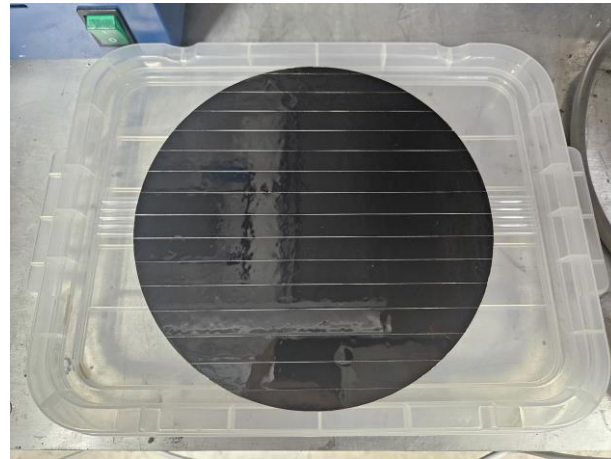
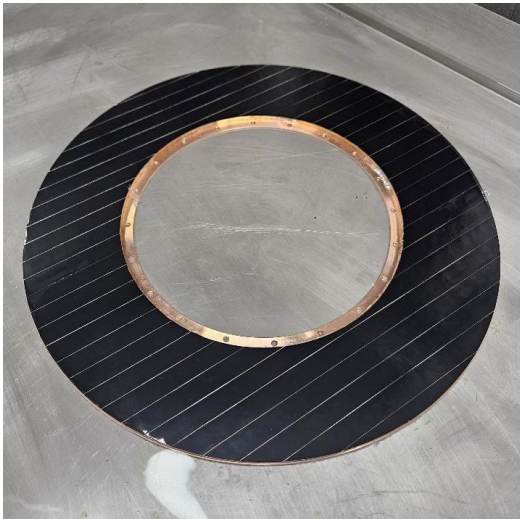
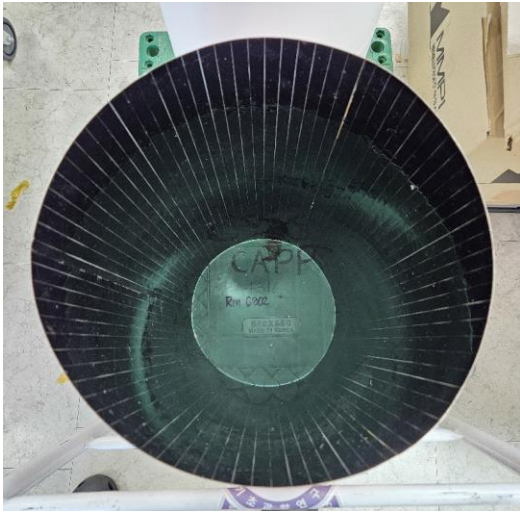
High-temperature Superconducting Cavity R&D

Strategy 2



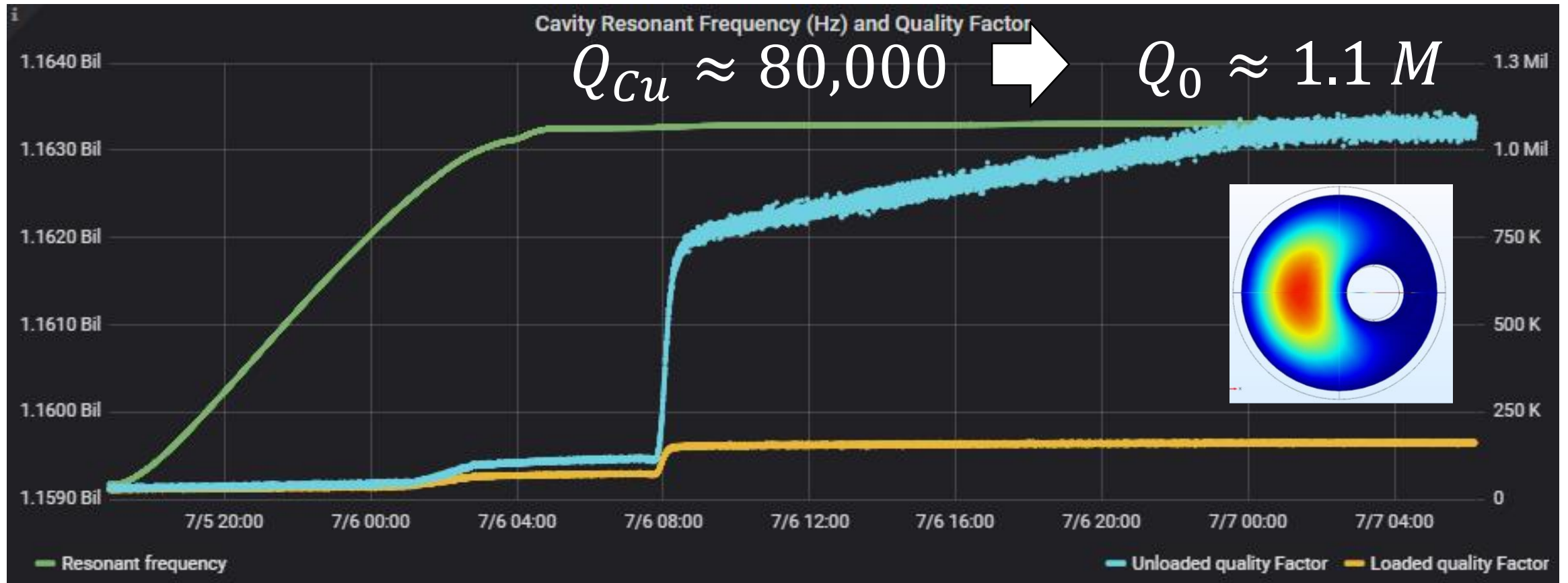
High-temperature Superconducting Cavity R&D

Strategy 2



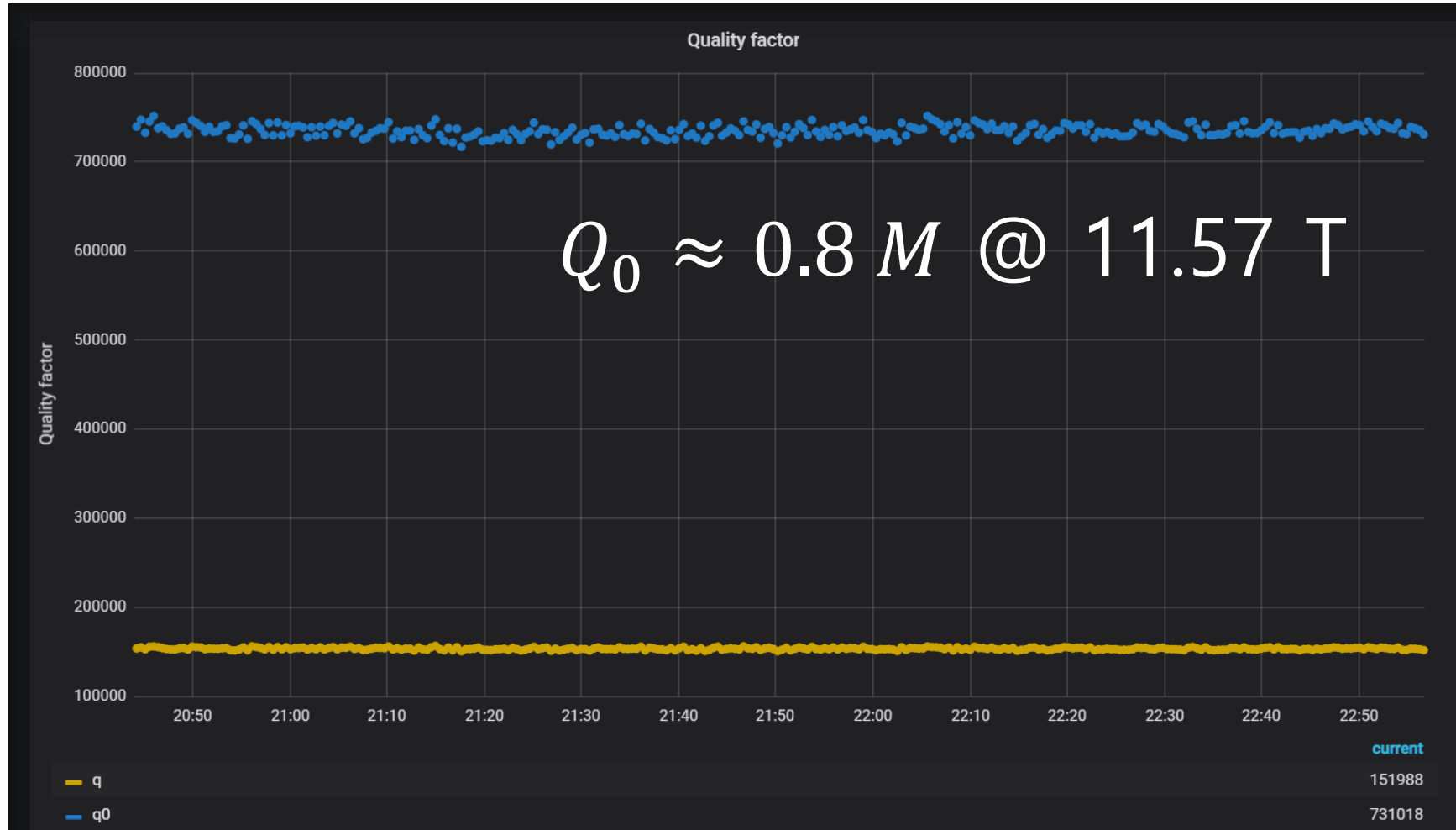
High-temperature Superconducting Cavity R&D

Strategy 2



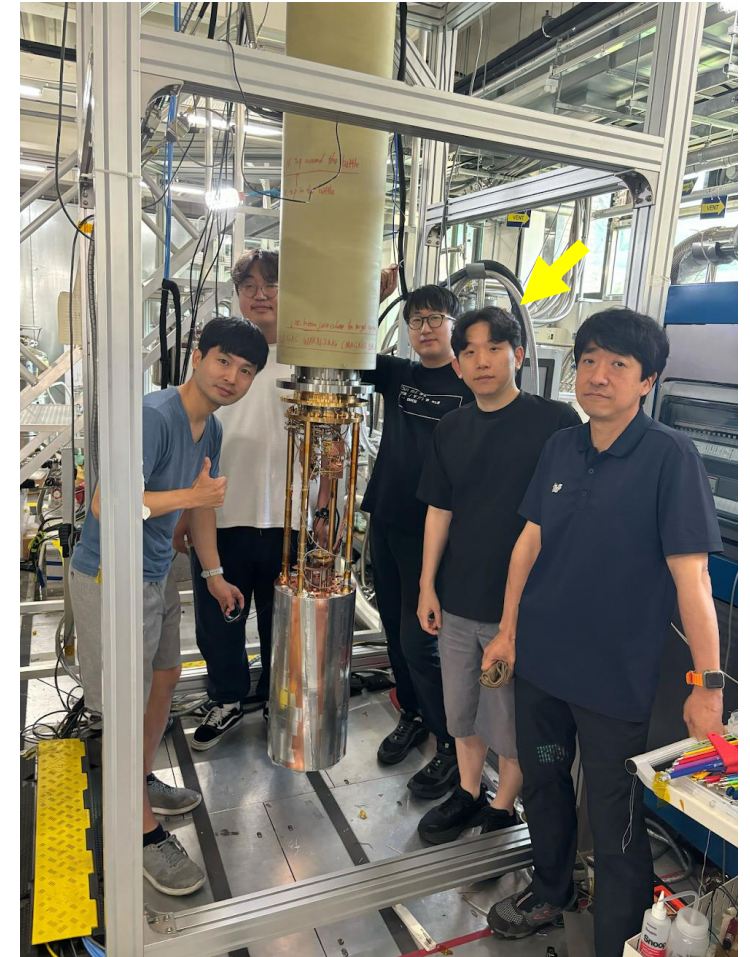
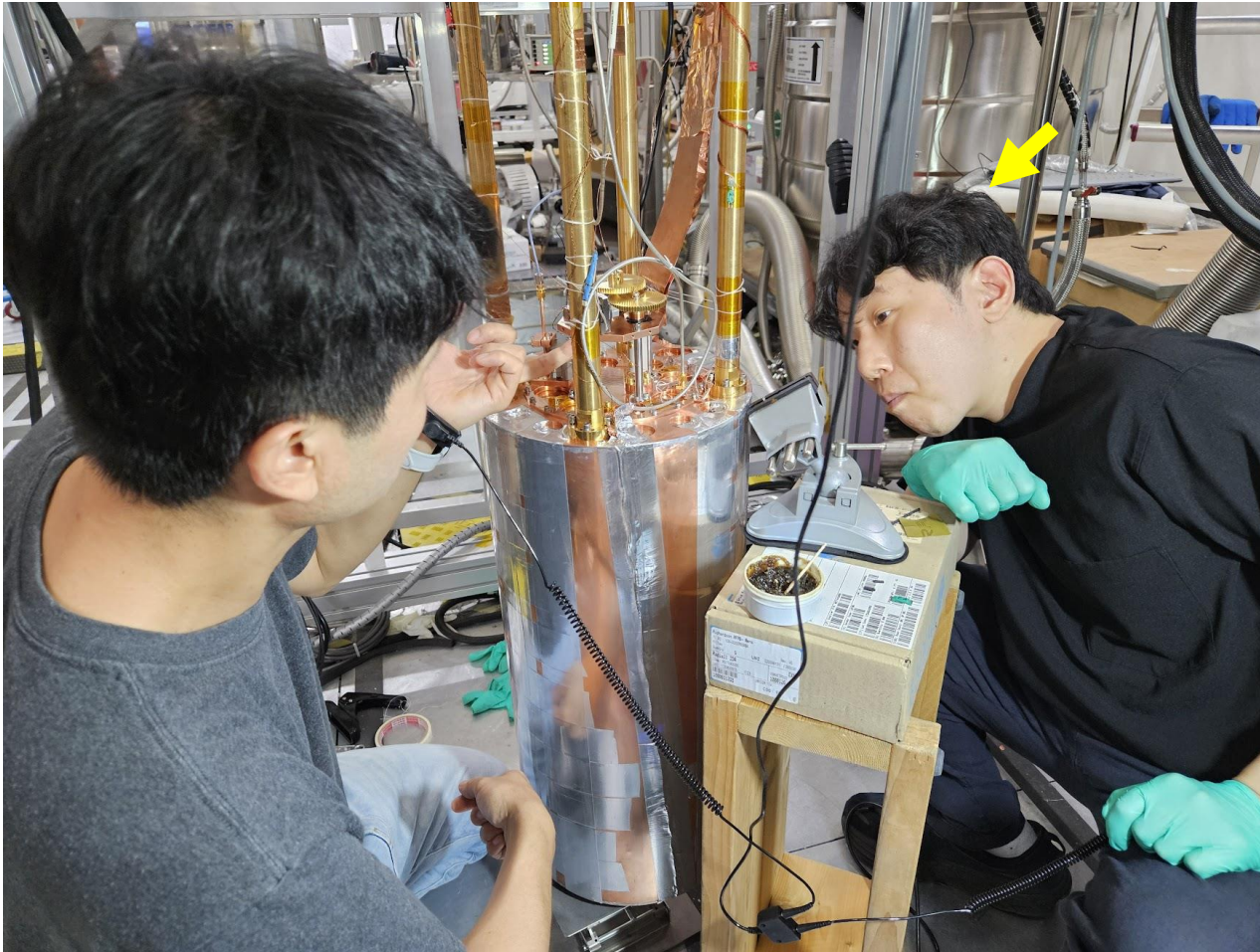
High-temperature Superconducting Cavity R&D

Strategy 2



Dark Matter Axion Search

CAPP-MAX with a High-Temperature Superconducting Cavity



Accelerating Axion Haloscope Experiments

Strategies and Innovations at CAPP

- The Center for Axion and Precision Physics Research (CAPP) has made significant progress in advancing axion haloscope experiments.
- The flagship experiment at CAPP, the Main Axion eXperiment (CAPP-MAX), features a powerful large-bore magnet integrated into a wet cryogenic system.
- The system employs a LEIDEN refrigerator capable of cooling the cavity to 30 mK, one of the lowest physical temperatures achieved in axion haloscope experiments.
- Furthermore, the use of a Josephson Parametric Amplifier introduces only quantum-limited noise to the RF chain, enabling a remarkably low system noise temperature of approximately 200 mK.
- CAPP has also developed a high-temperature superconducting (HTS) cavity for the CAPP-MAX system, which maintains a quality factor in the millions even under a 12 T magnetic field.
- Currently, CAPP is preparing to cool down the cryogenic system and begin data acquisition.