

# Toward Accurate Modeling of the Nonlinear Galaxy Bispectrum

Donghui Jeong  
The Pennsylvania State University

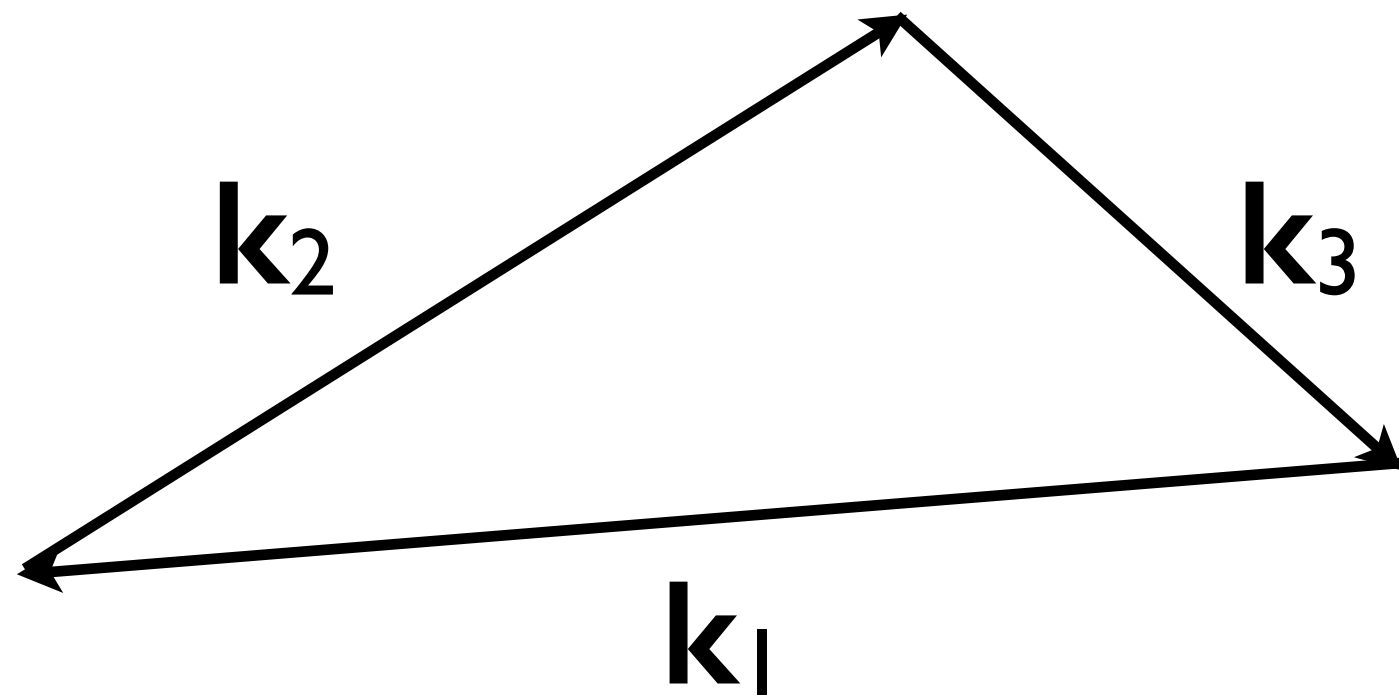
In collaboration with Nuala McCullagh and Alex Szalay (JHU)

8 July 2014, Precision Era for Large Scale Structure @ IBS, Daejeon, Korea

# What is the galaxy bispectrum?

- **The galaxy bispectrum** is the Fourier transform of the galaxy three-point correlation function:

$$\langle \delta_g(\mathbf{k}_1) \delta_g(\mathbf{k}_2) \delta_g(\mathbf{k}_3) \rangle = (2\pi)^3 B_g(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3) \delta^D(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3)$$



$$\delta_g(\mathbf{k}) = \int d^3\mathbf{x} \left( \frac{n(\mathbf{x}) - \bar{n}}{\bar{n}} \right) e^{-i\mathbf{k} \cdot \mathbf{x}}$$

density contrast

# Why studying bispectrum?

- The galaxy bispectrum is a treasure island yet to be explored!
- By exploiting the galaxy bispectrum, we can learn about  
(Scoccimarro+ 1998, Scoccimarro+ 1999, Sefusatti+ 2006, Sefusatti & Komatsu 2007, Jeong & Komatsu 2009, Baldauf+ 2011, Greig+ 2013)
  - Physics of inflation: non-Gaussianity
  - Non-linear structure formation
  - Astrophysics
    - Formation, evolution, radiative transfer of galaxies

# Challenges to exploit bispectrum

- For any analysis, we need :  $\chi^2 = (\text{theory} - \text{data})^2 / \text{Covariance}$ 
  - [1] Theory (non-linearities, non-Gaussianities, bias)
  - [2] Data (window function, optimal weighting)
  - [3] Covariance matrix (6-point correlation function)
- For bispectrum, all three are **NOT YET** very well understood to the level that we can confidently apply them for the real analysis



# Challenges to exploit bispectrum

- For any analysis, we need :  $\chi^2 = (\text{theory} - \text{data})^2 / \text{Covariance}$ 
  - [1] **Theory** (**non-linearities**, non-Gaussianities, bias)
  - [2] Data (window function, optimal weighting)
  - [3] Covariance matrix (6-point correlation function)
- For bispectrum, all three are **NOT YET** very well understood to the level that we can confidently apply them for the real analysis

# Visualizing the bispectrum

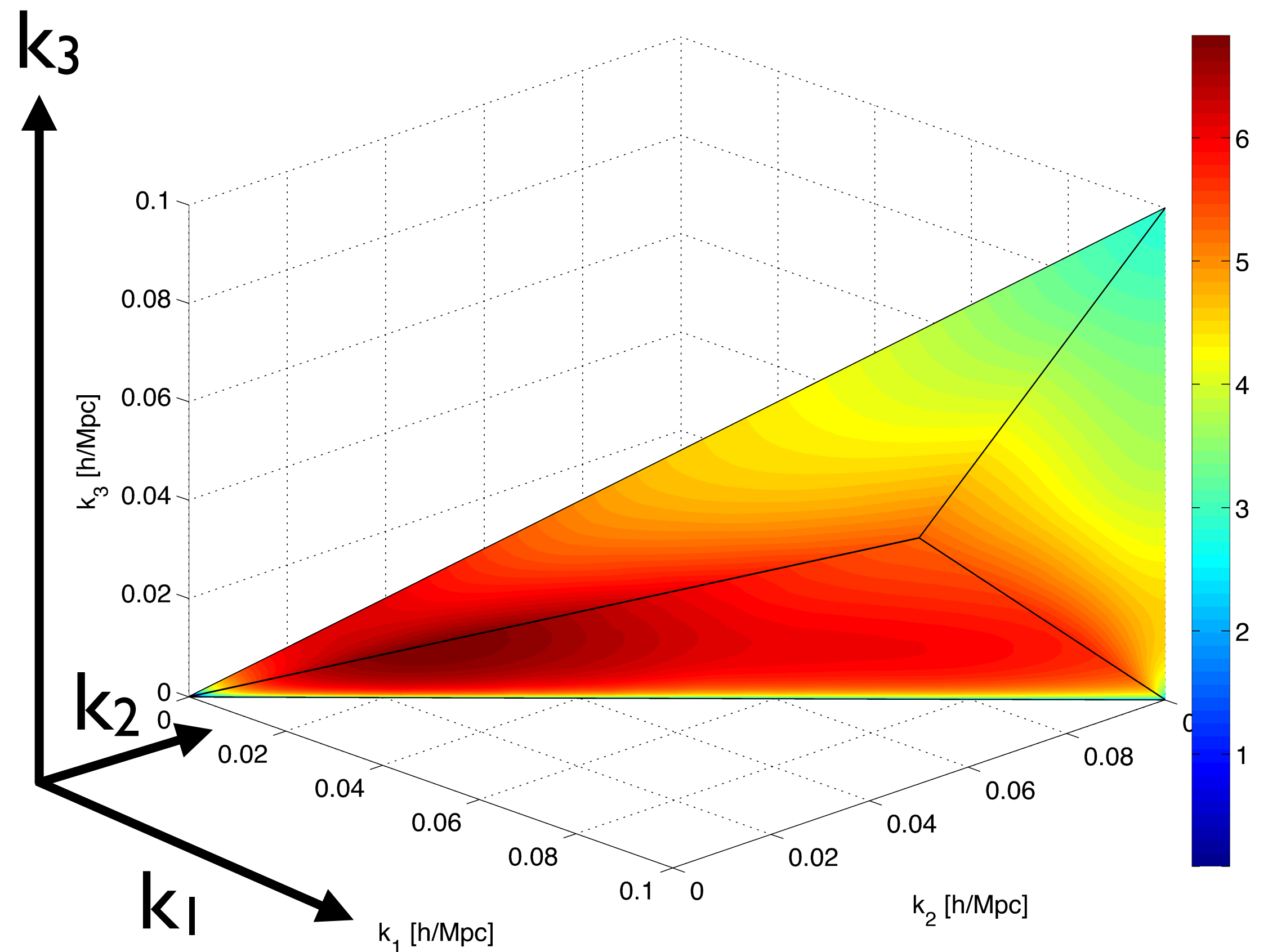
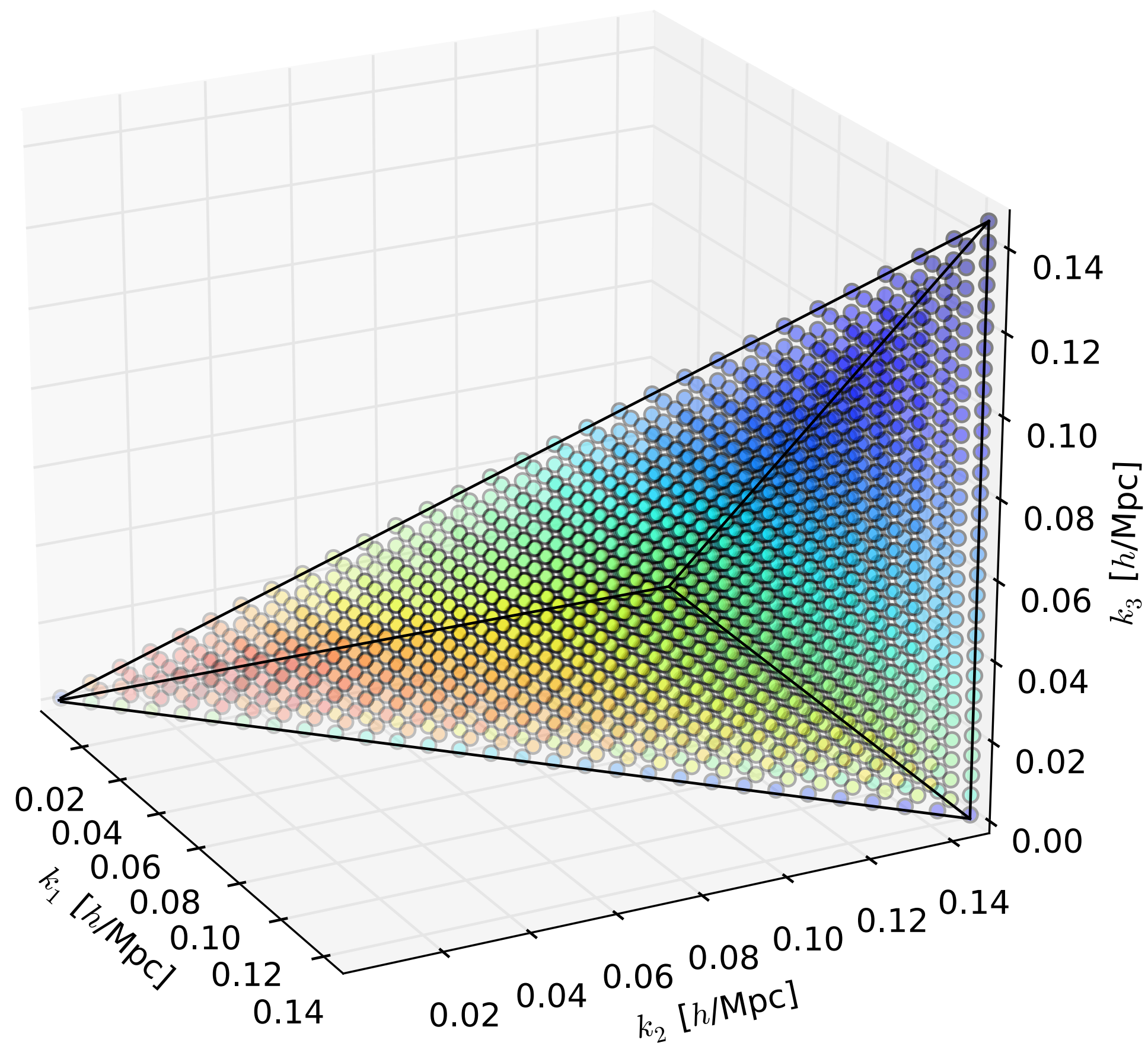
- With statistical isotropy, the bispectrum only depends on the three sides of triangles (convention  $k_1 \geq k_2 \geq k_3$ ):

$$\langle \delta_m(\mathbf{k}_1) \delta_m(\mathbf{k}_2) \delta_m(\mathbf{k}_3) \rangle = (2\pi)^3 B_m(k_1, k_2, k_3) \delta^D(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3)$$

- Therefore, it may be natural to show bispectrum in 3D.
- Example: the leading order theory prediction

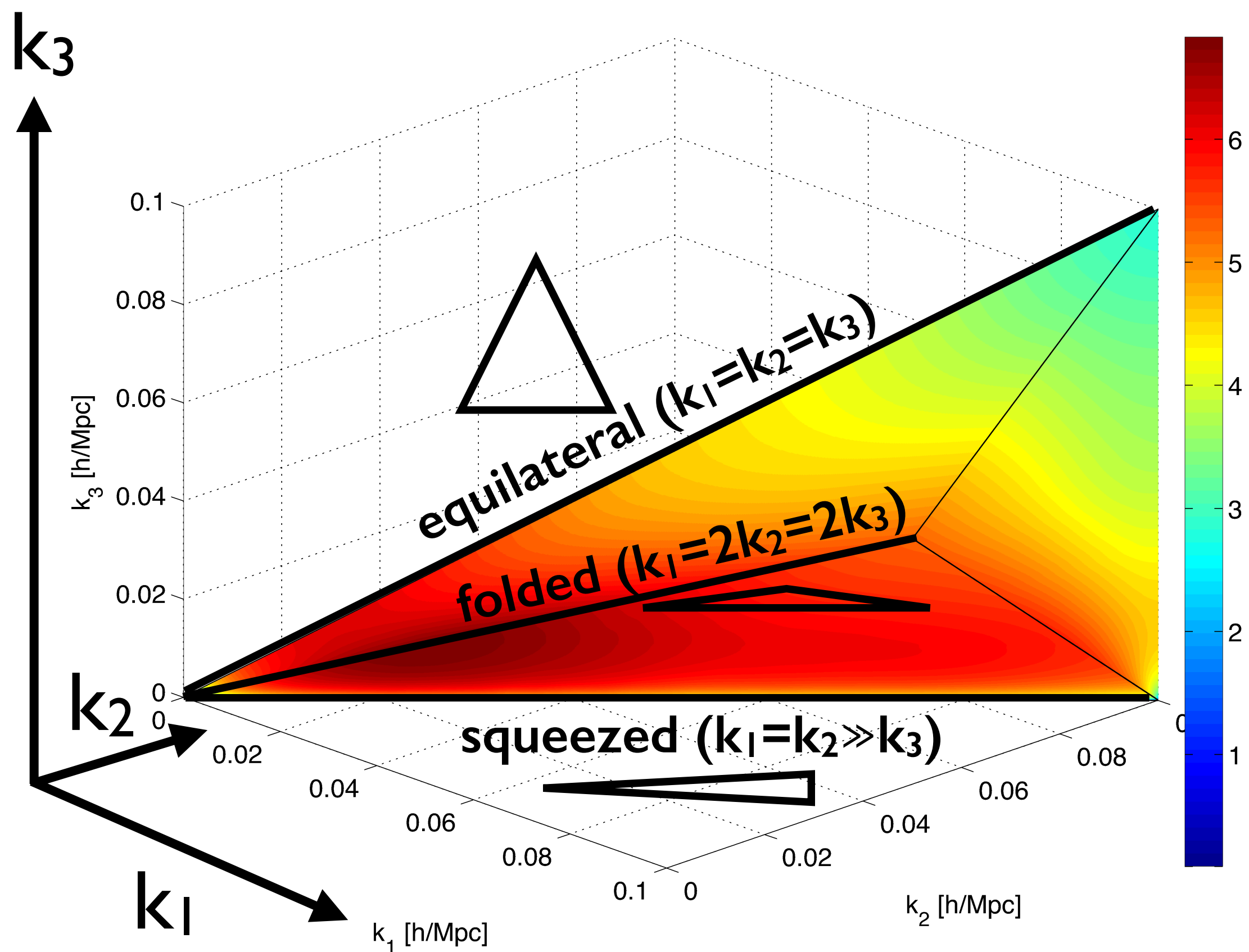
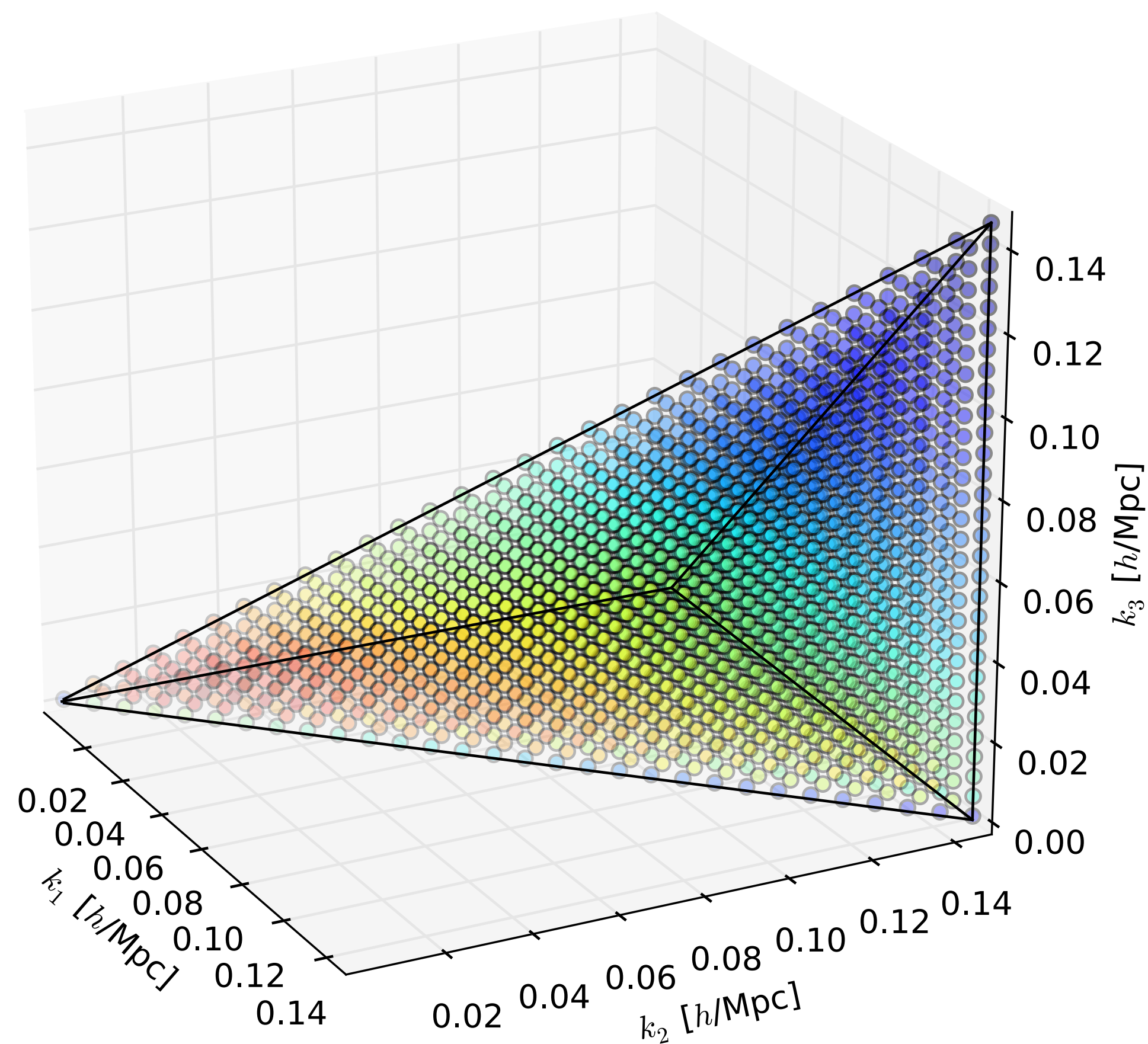
$$B_m(k_1, k_2, k_3) = \frac{3k_3^2(k_1^2 + k_2^2) - 5(k_1^2 - k_2^2)^2 + 2k_3^2}{14k_1^2 k_2^2} P_L(k_1) P_L(k_2) + (2 \text{ cyclic})$$
$$= 2F_2(\mathbf{k}_1, \mathbf{k}_2)$$

# Bispectrum in 3D ( $k_1, k_2, k_3$ )

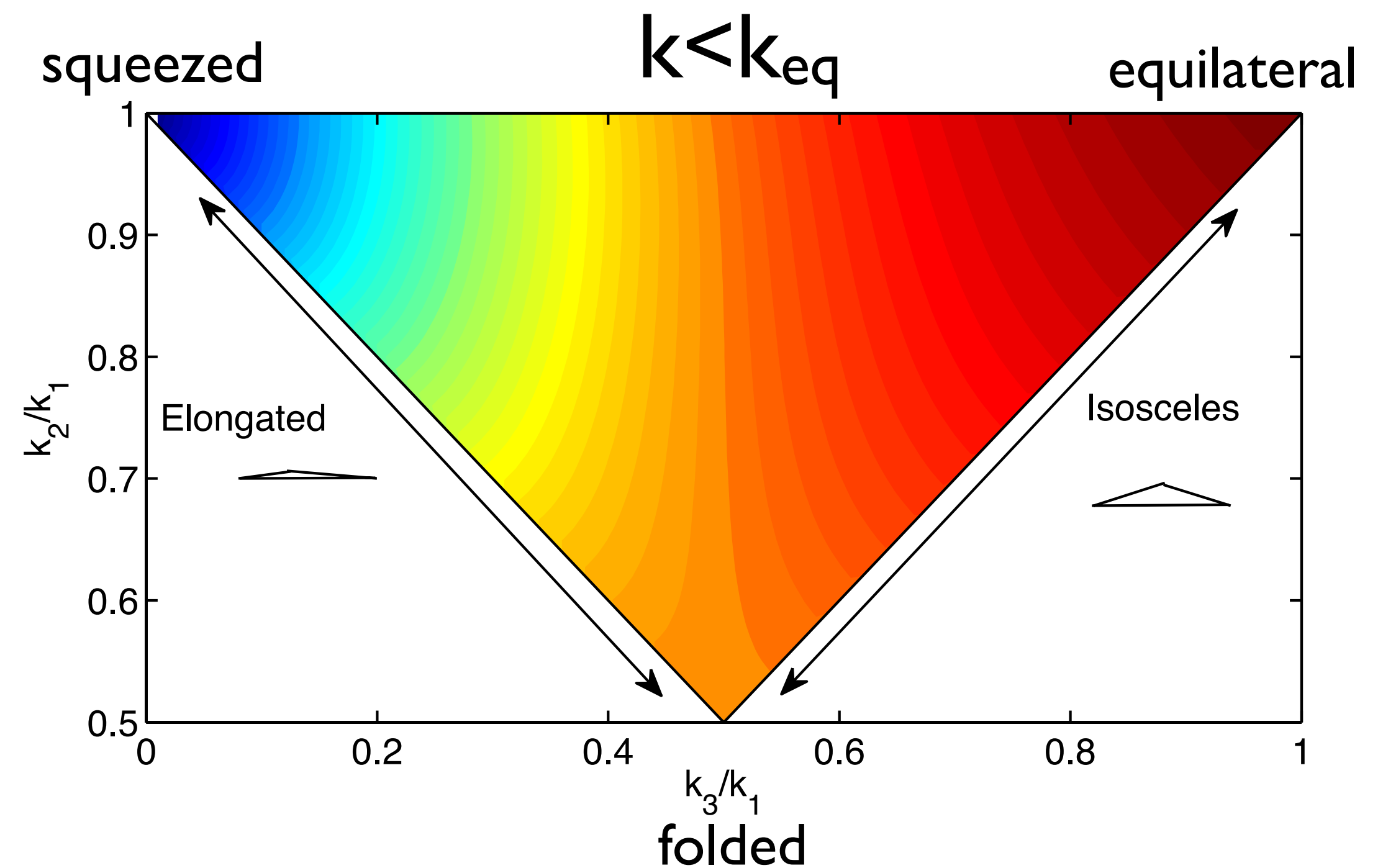
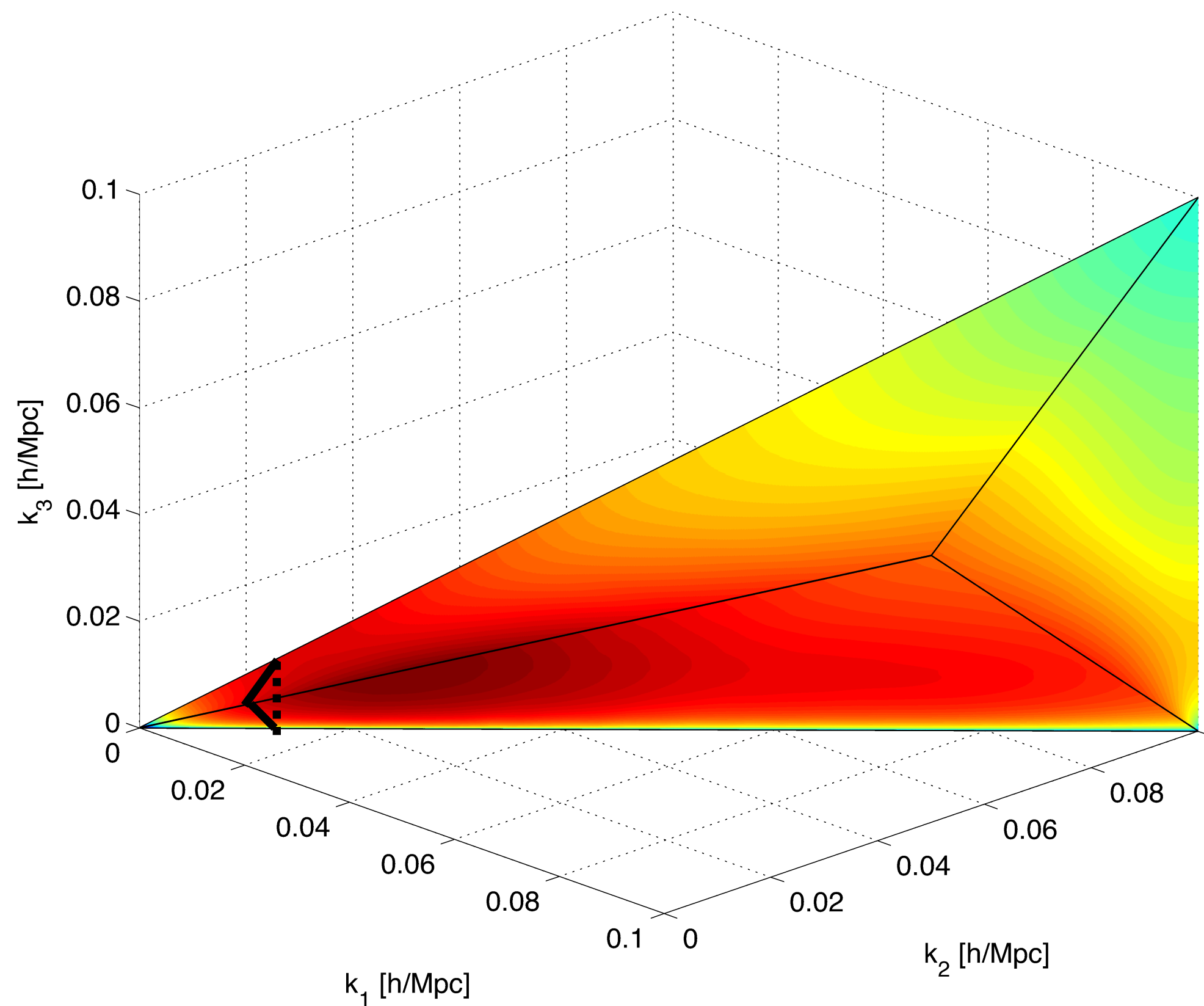




# Bispectrum in 3D ( $k_1, k_2, k_3$ )

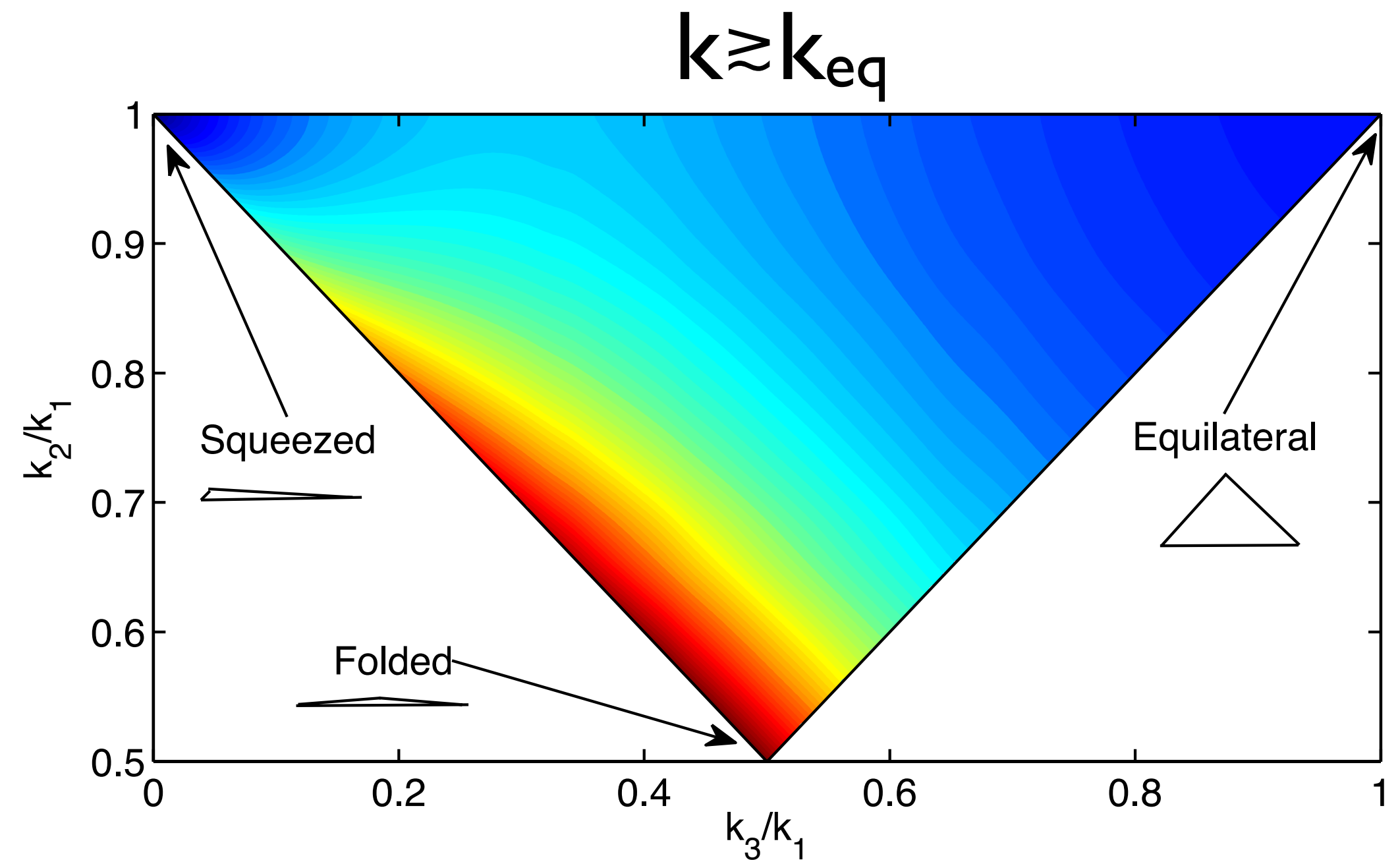
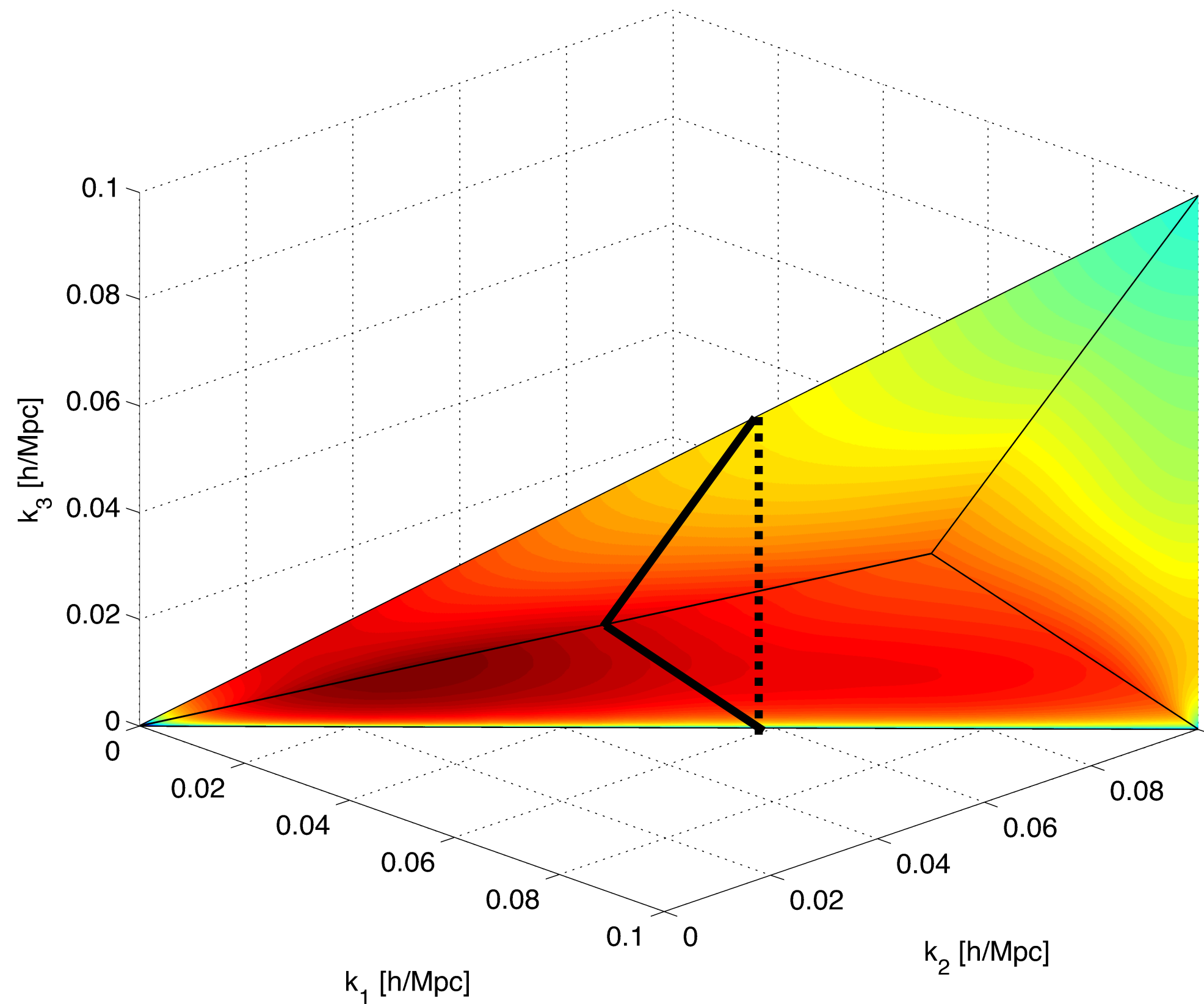


# Bispectrum in 2D ( $k_1$ fixed; $k_2, k_3$ )



The slope of power spectrum  
 $d \ln P / d \ln k > 0$

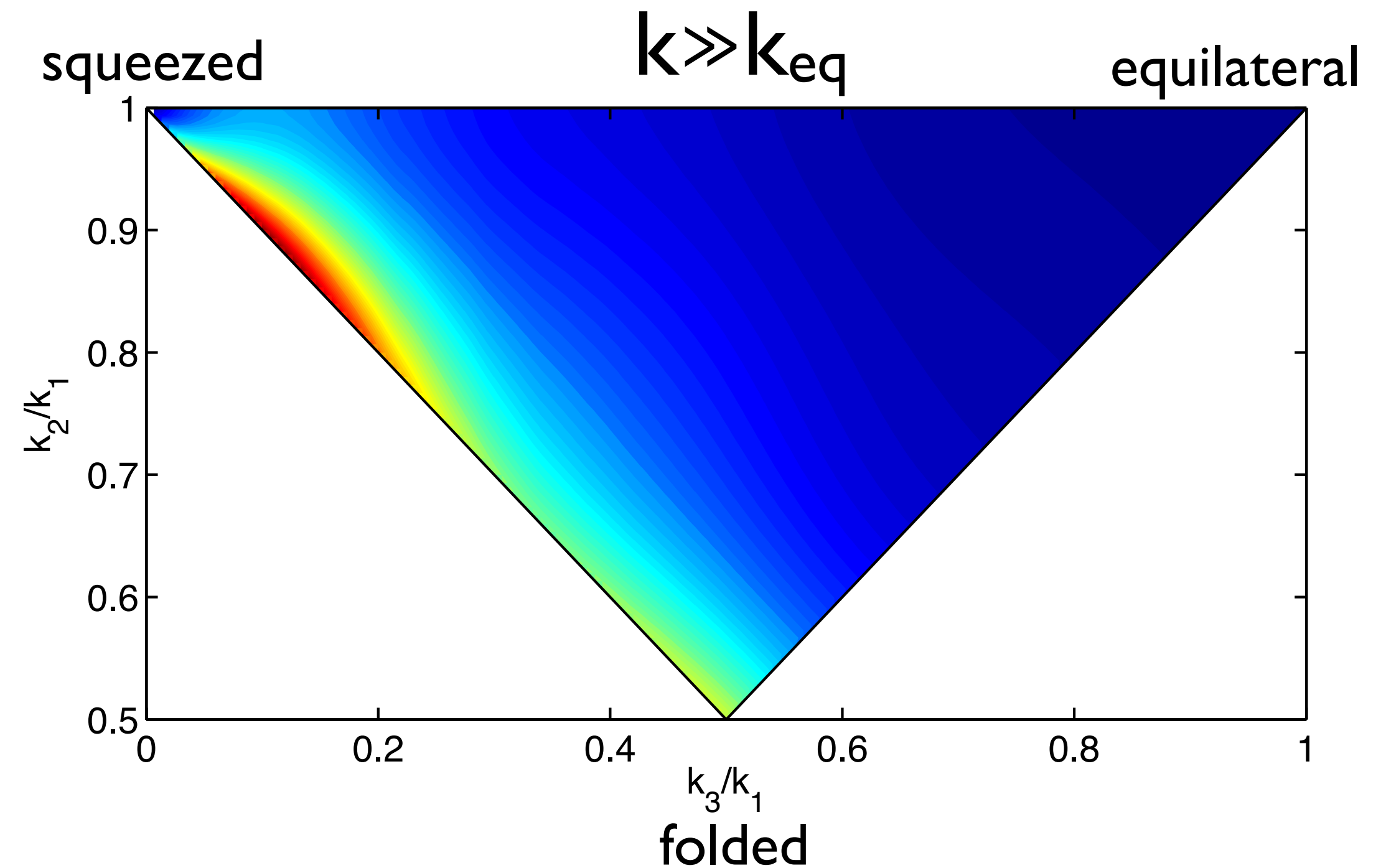
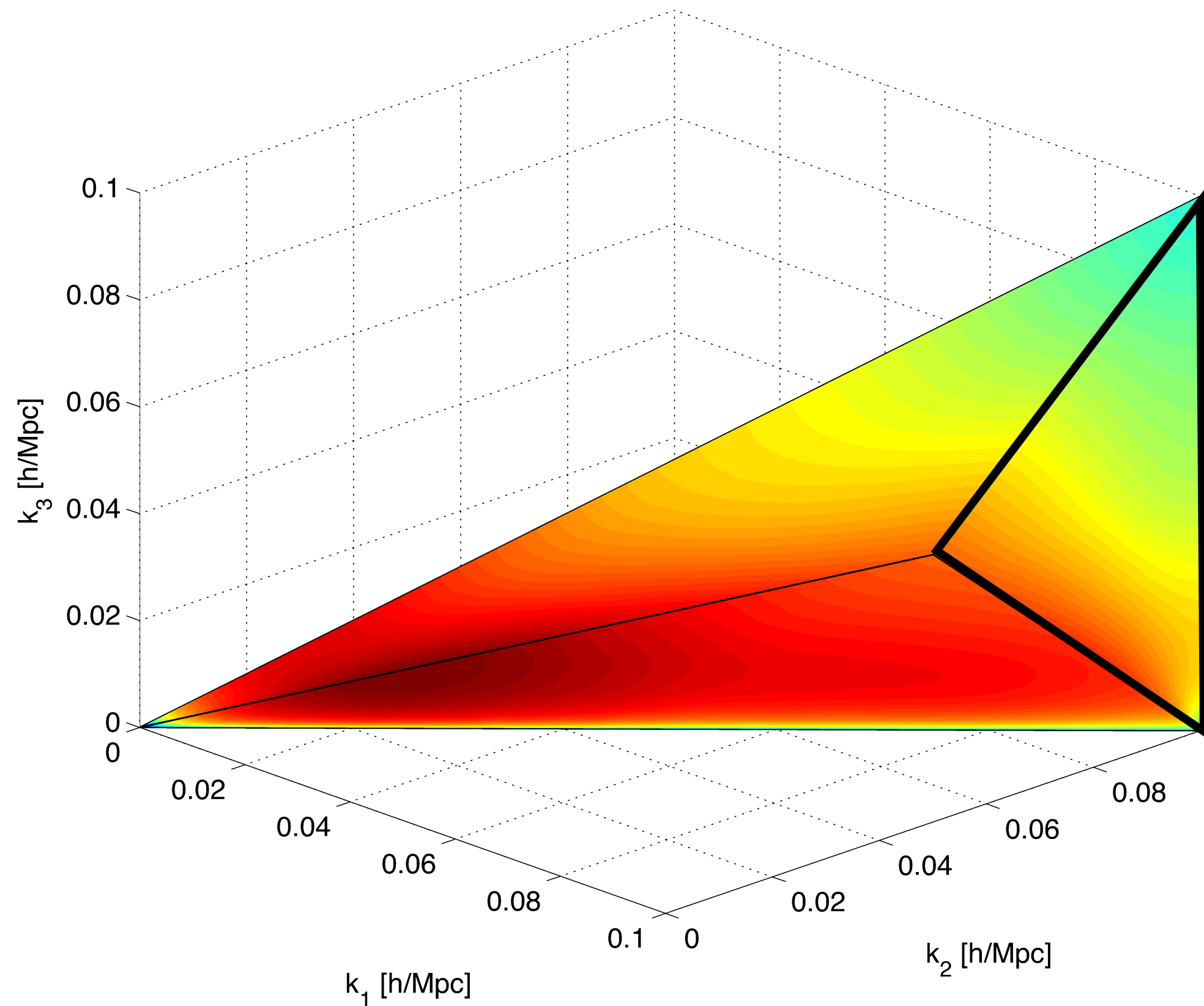
# Bispectrum in 2D ( $k_1$ fixed; $k_2, k_3$ )



The slope of power spectrum  
 $d \ln P / d \ln k \approx 0$

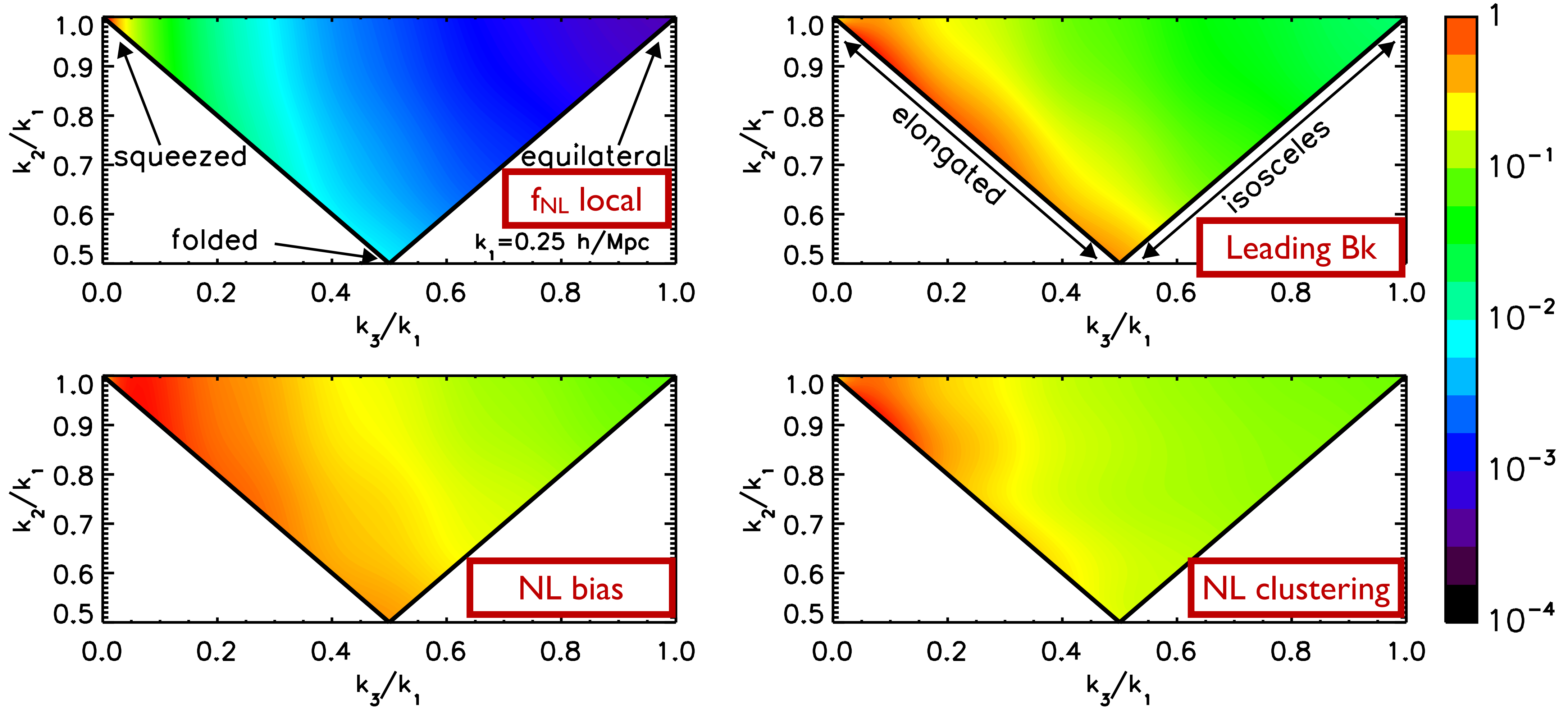


# Bispectrum in 2D ( $k_1$ fixed; $k_2, k_3$ )



The slope of power spectrum  
 $d \ln P / d \ln k < 0$

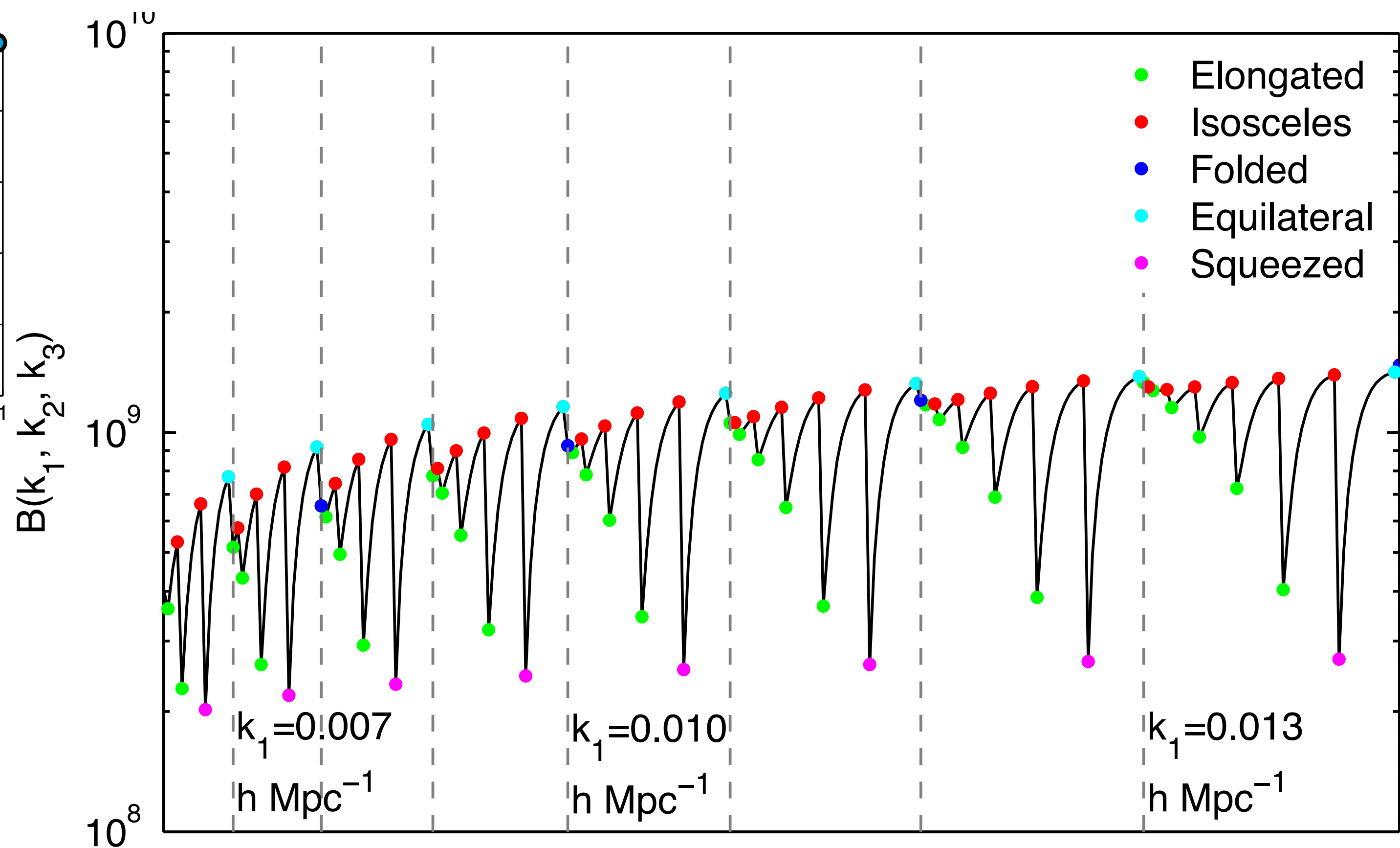
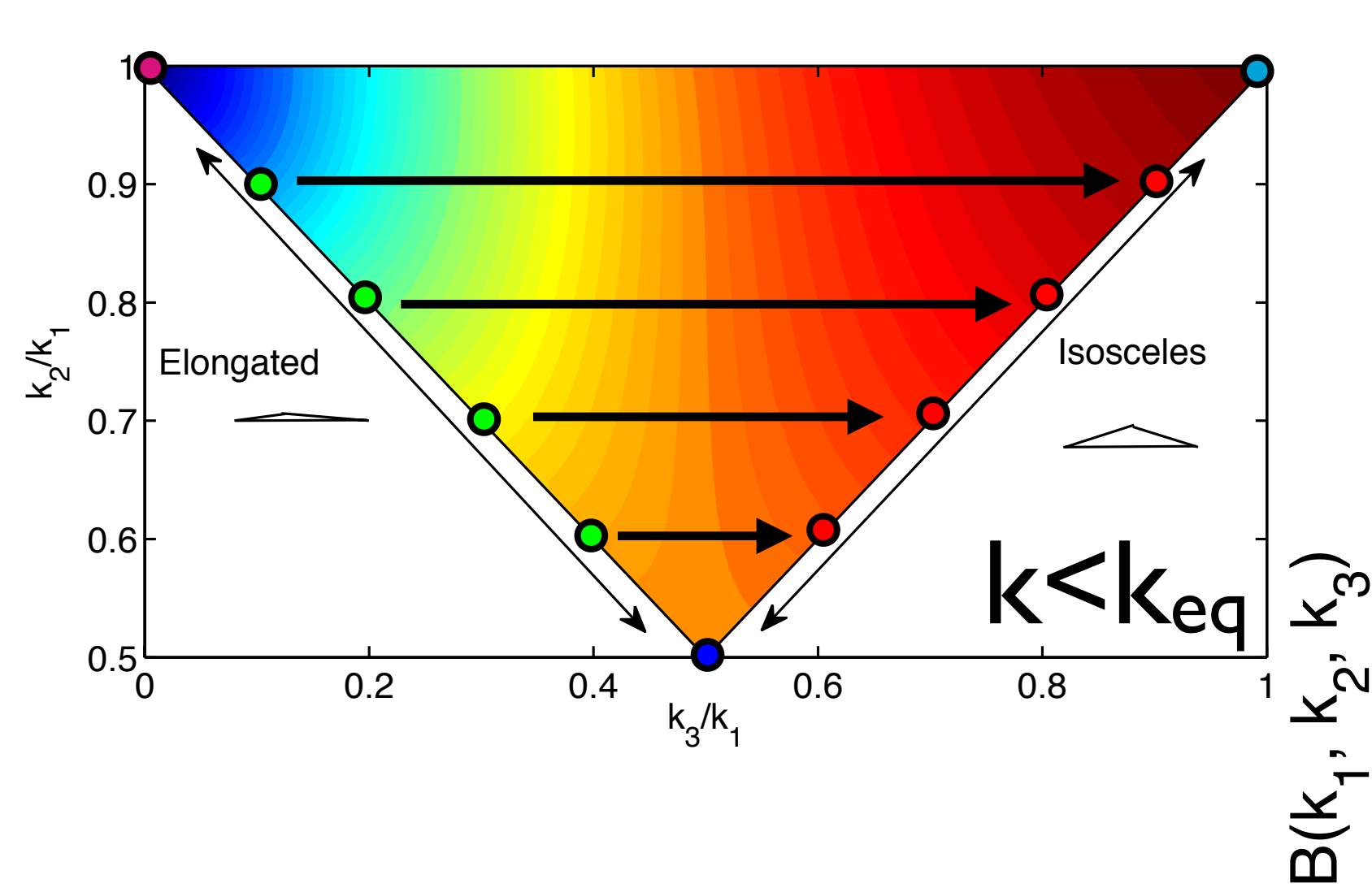
# Shape of each component





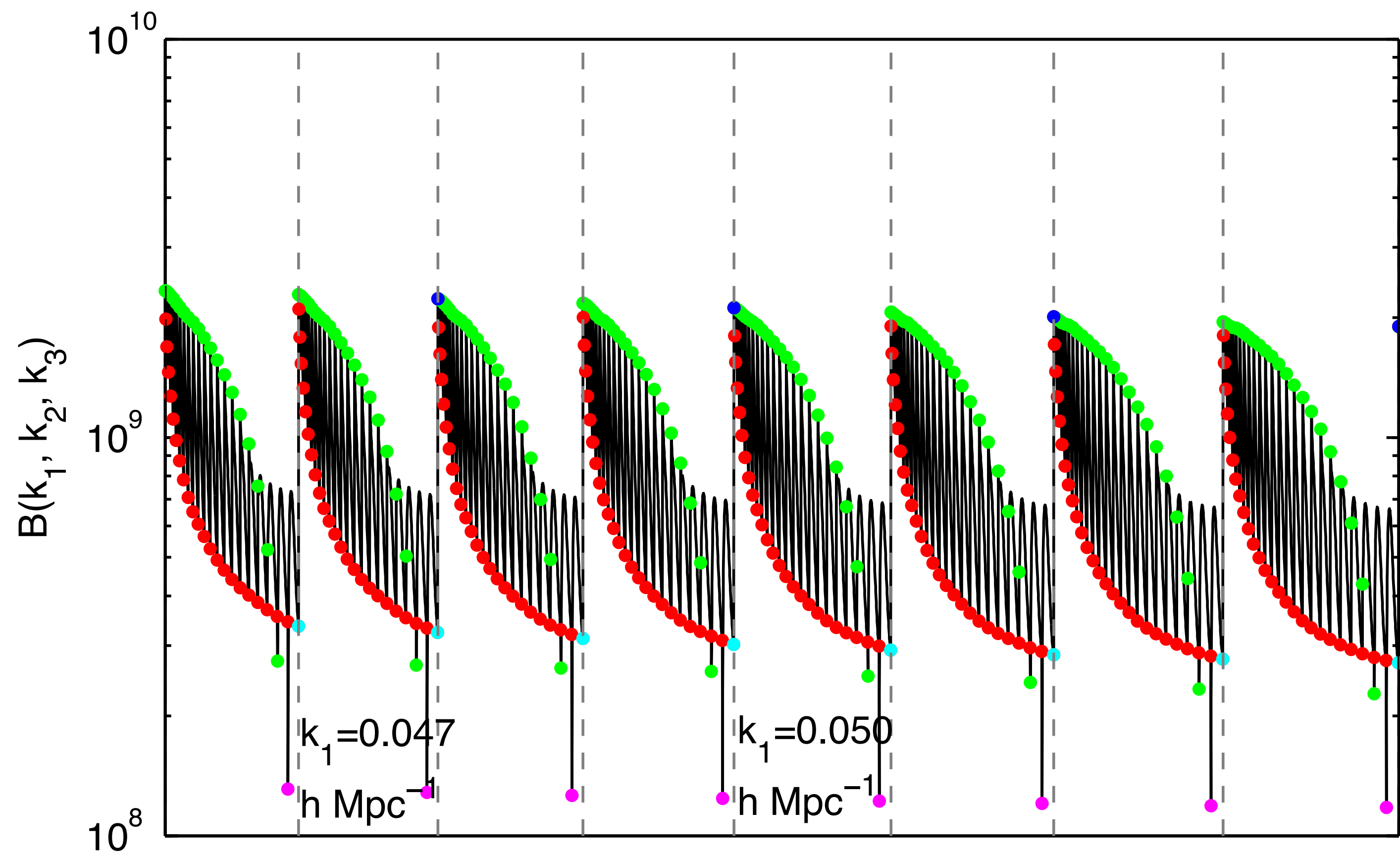
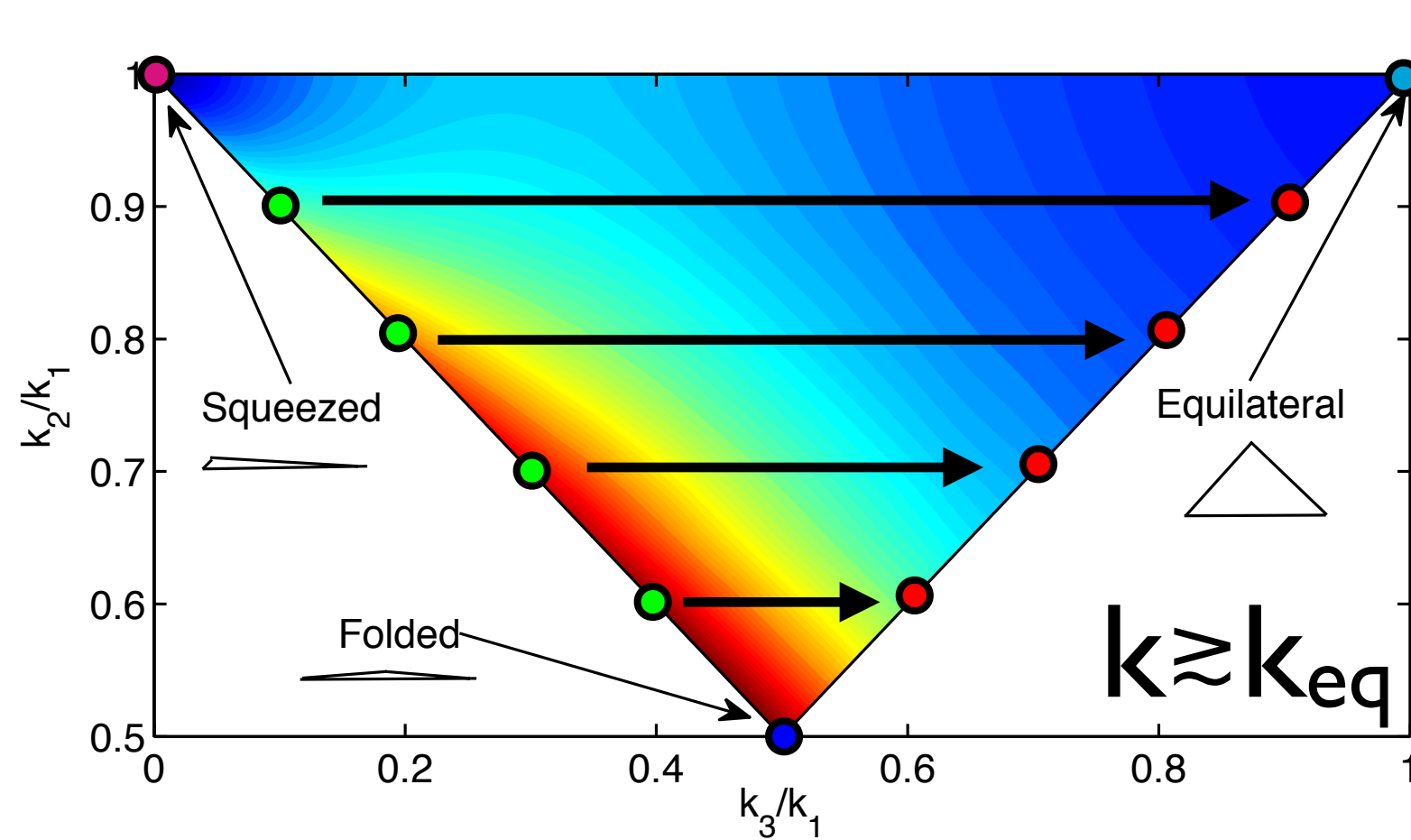
# Flattened Bispectrum $(k_1, k_2, k_3)[n]$

row-major, ascending order



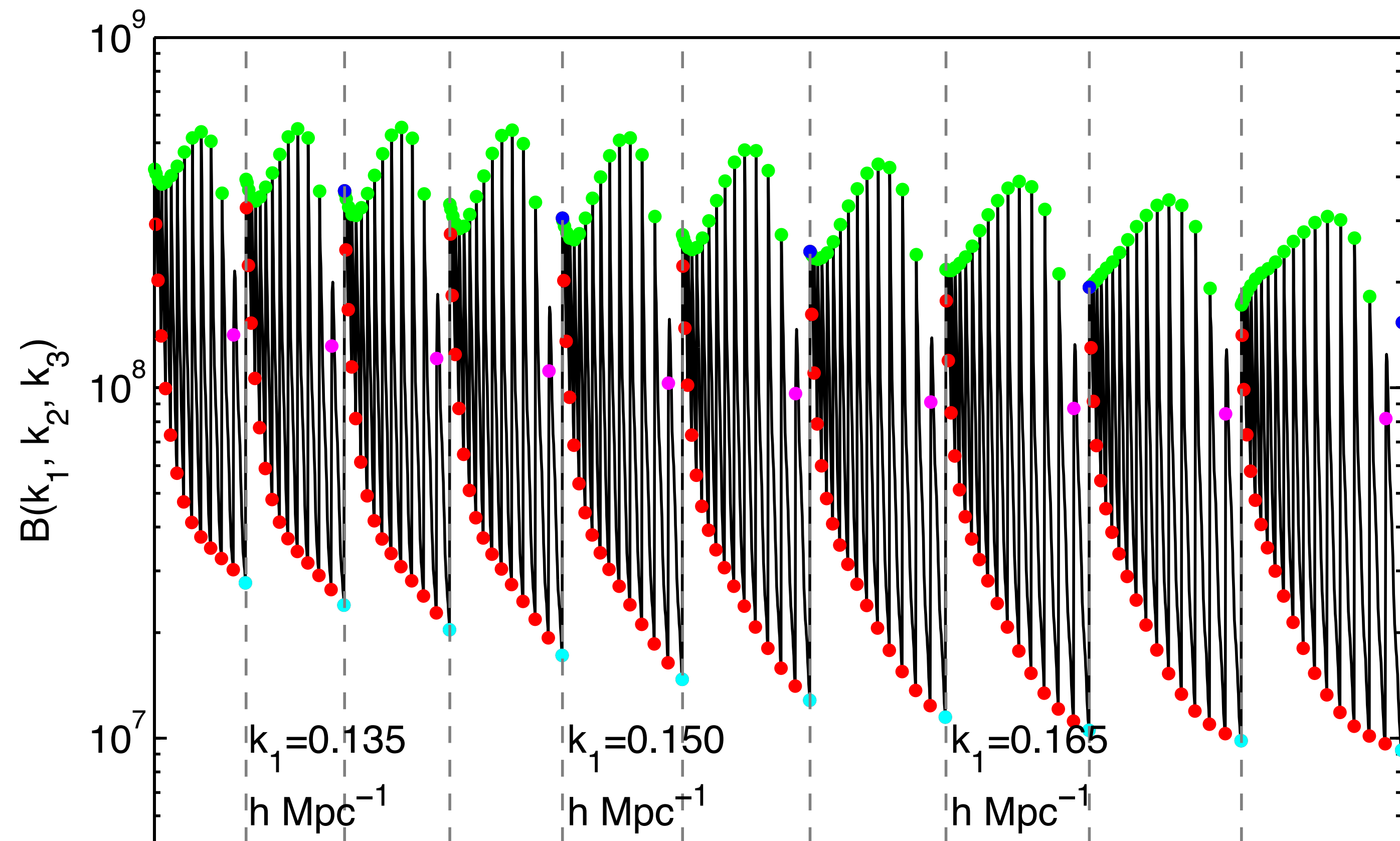
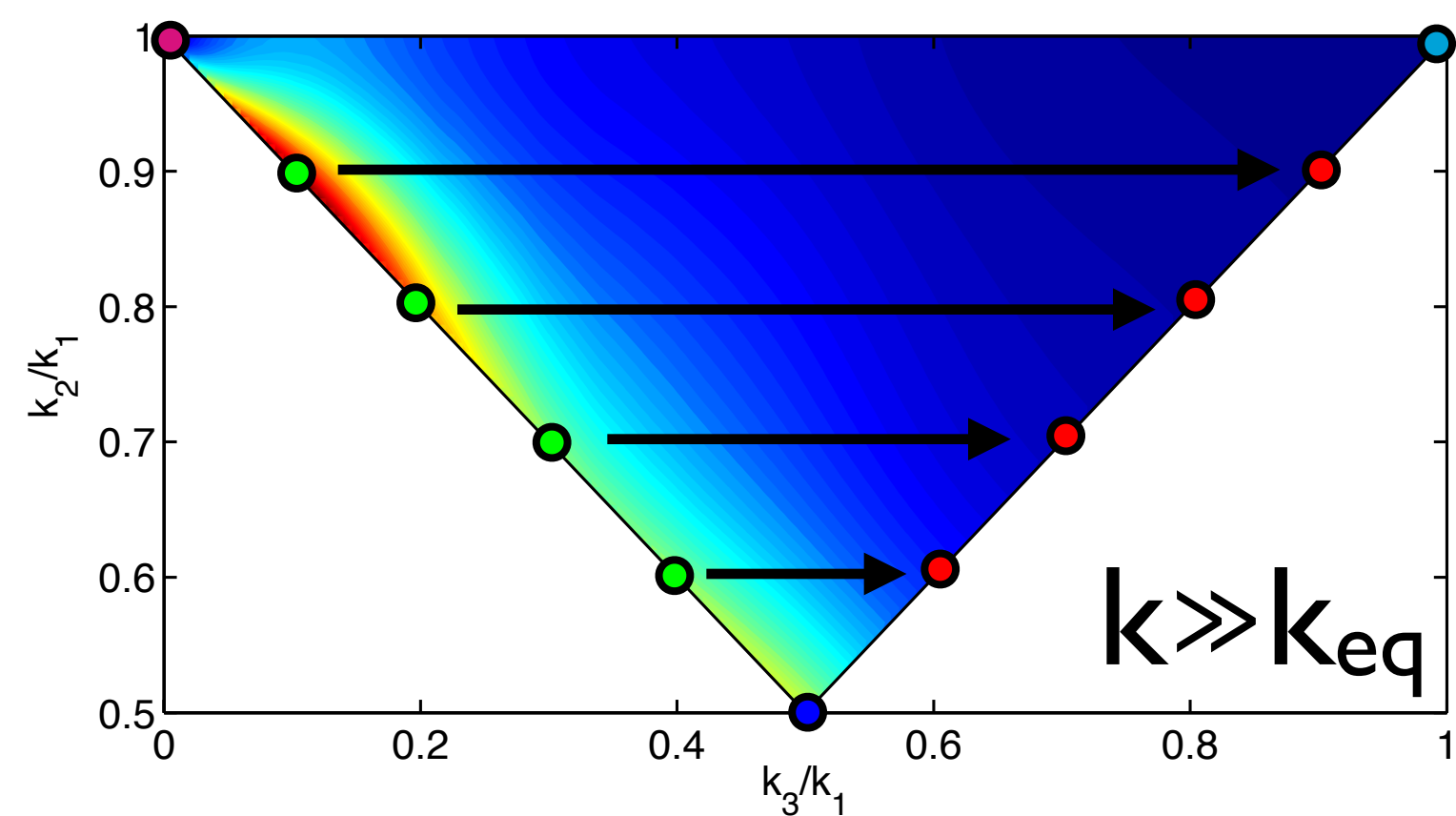
# Flattened Bispectrum $(k_1, k_2, k_3)[n]$

row-major, ascending order



# Flattened Bispectrum $(k_1, k_2, k_3)[n]$

row-major, ascending order



# Modeling non-linearities

- On large scales:  
‘tree-level’ bispectrum from second order density contrast  
 $B^{(0)} = \langle \delta_2 \delta_1 \delta_1 \rangle$ ; here,  $\delta_n$  = n-th order density contrast
- Smaller scales:  
‘one-loop’ bispectrum  $B^{(1)} = \langle \delta_4 \delta_1 \delta_1 \rangle + \langle \delta_3 \delta_2 \delta_1 \rangle + \langle \delta_2 \delta_2 \delta_2 \rangle$
- Even smaller scales (highly non-linear):  
There are **empirical fitting formulas** from simulations.

# ‘Large’, ‘Smaller,’ ‘even smaller’?

- are not a particularly accurate scientific term...
- Need N-body simulations to draw the line between scales
- Q: How reliable are N-body simulations?

# Transient from initial conditions

- To correctly capture non-linearities at all subsequent redshifts, initial condition of N-body simulation must be at the fastest growing modes at all order.
- We don't know how to generate that initial condition. In fact, that is one of the reasons that we run simulations.
- Instead, we use the Lagrangian perturbation theory prescription: linear order (Zeldovich approximation) or second order (2LPT)



# Decaying modes

- Initial condition from **Zeldovich approximation** correctly captures **fastest growing mode in the linear order**, but generates **decaying mode from the second order**
- Initial condition from **2LPT approximation** correctly captures **fastest growing mode in the second order**, but generates **decaying mode from the third order**
- Leading order ('tree-level') bispectrum  $B^{(0)} = \langle \delta_2 \delta_1 \delta_1 \rangle$
- How large is the effect?

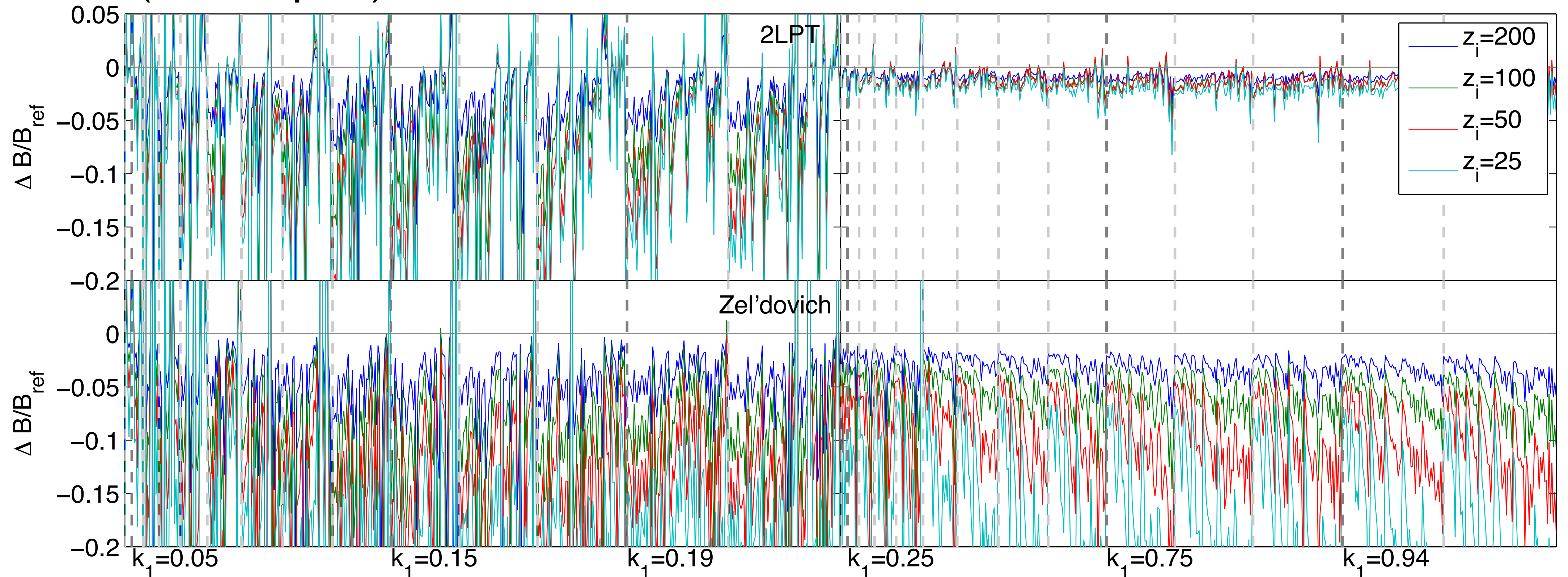
# Transient at $z=6$

$$\frac{Bk(2LPT/Zeldovich, z_{\text{init}})}{Bk(2LPT, z_{\text{init}}=400)}$$

$(1000 \text{ Mpc}/h)^3$

$Bk(2LPT, z_{\text{init}}=400)$

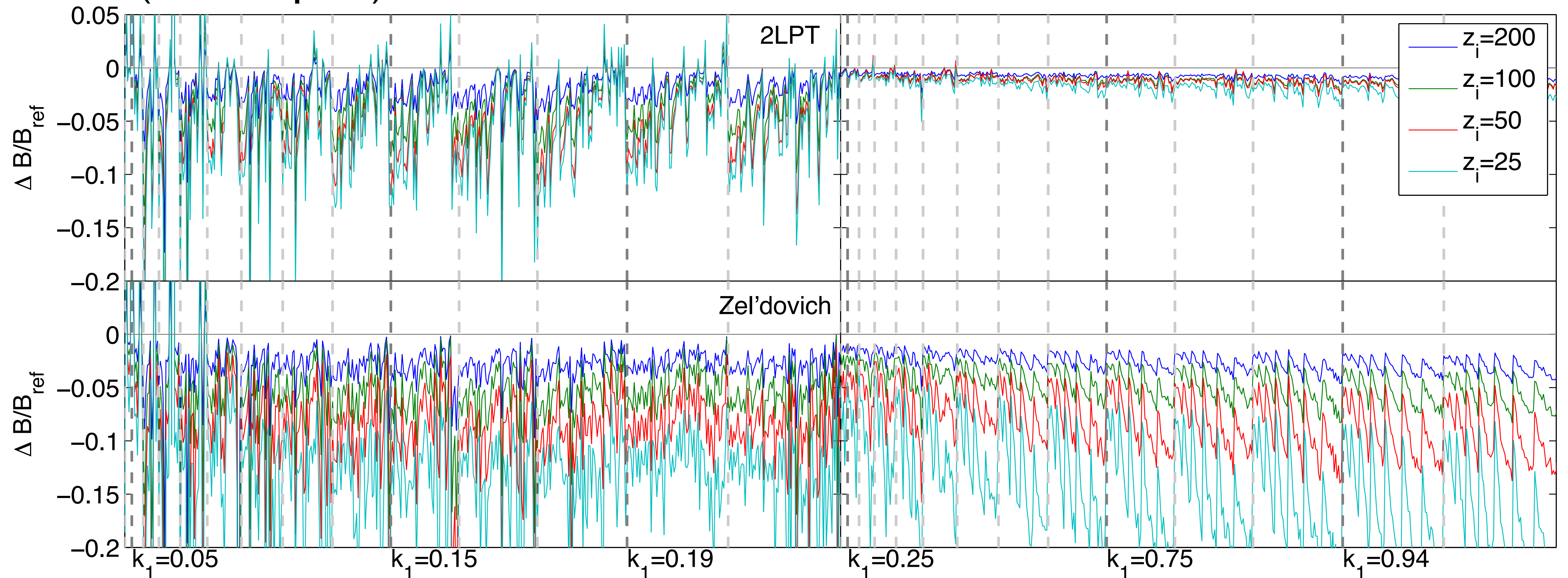
$(200 \text{ Mpc}/h)^3$





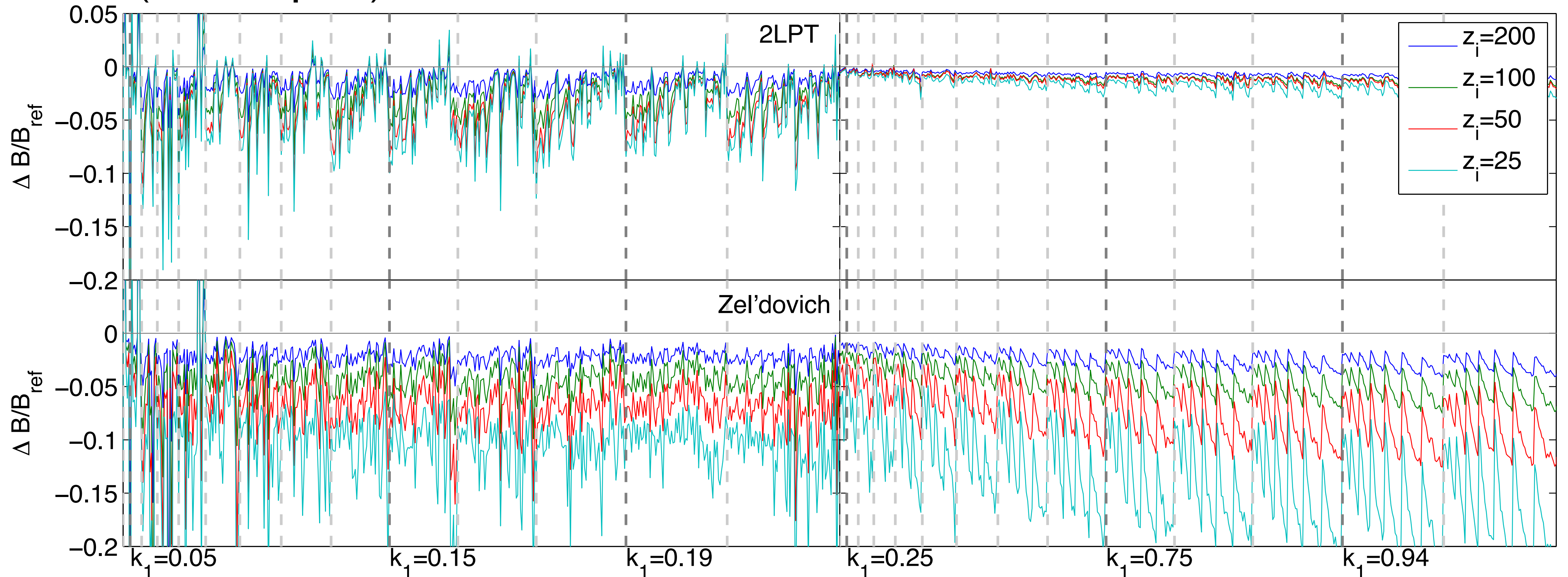
# Transient at $z=3$

$$\frac{Bk(2LPT/Zeldovich, z_{\text{init}})}{Bk(2LPT, z_{\text{init}}=400)}$$

 $(1000 \text{ Mpc}/h)^3$ 
 $Bk(2LPT, z_{\text{init}}=400)$ 
 $(200 \text{ Mpc}/h)^3$ 


# Transient at $z=2$

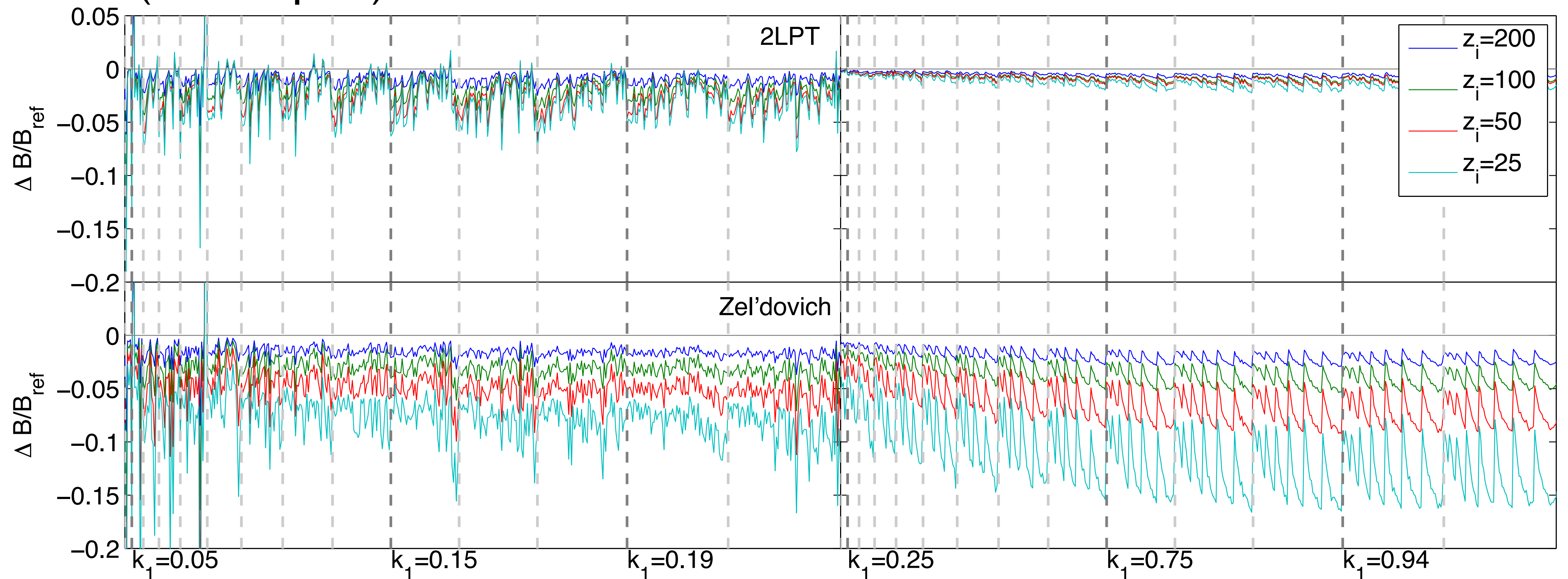
$$\frac{Bk(2LPT/Zeldovich, z_{\text{init}})}{Bk(2LPT, z_{\text{init}}=400)}$$

 $(1000 \text{ Mpc}/h)^3$ 
 $Bk(2LPT, z_{\text{init}}=400)$ 
 $(200 \text{ Mpc}/h)^3$ 


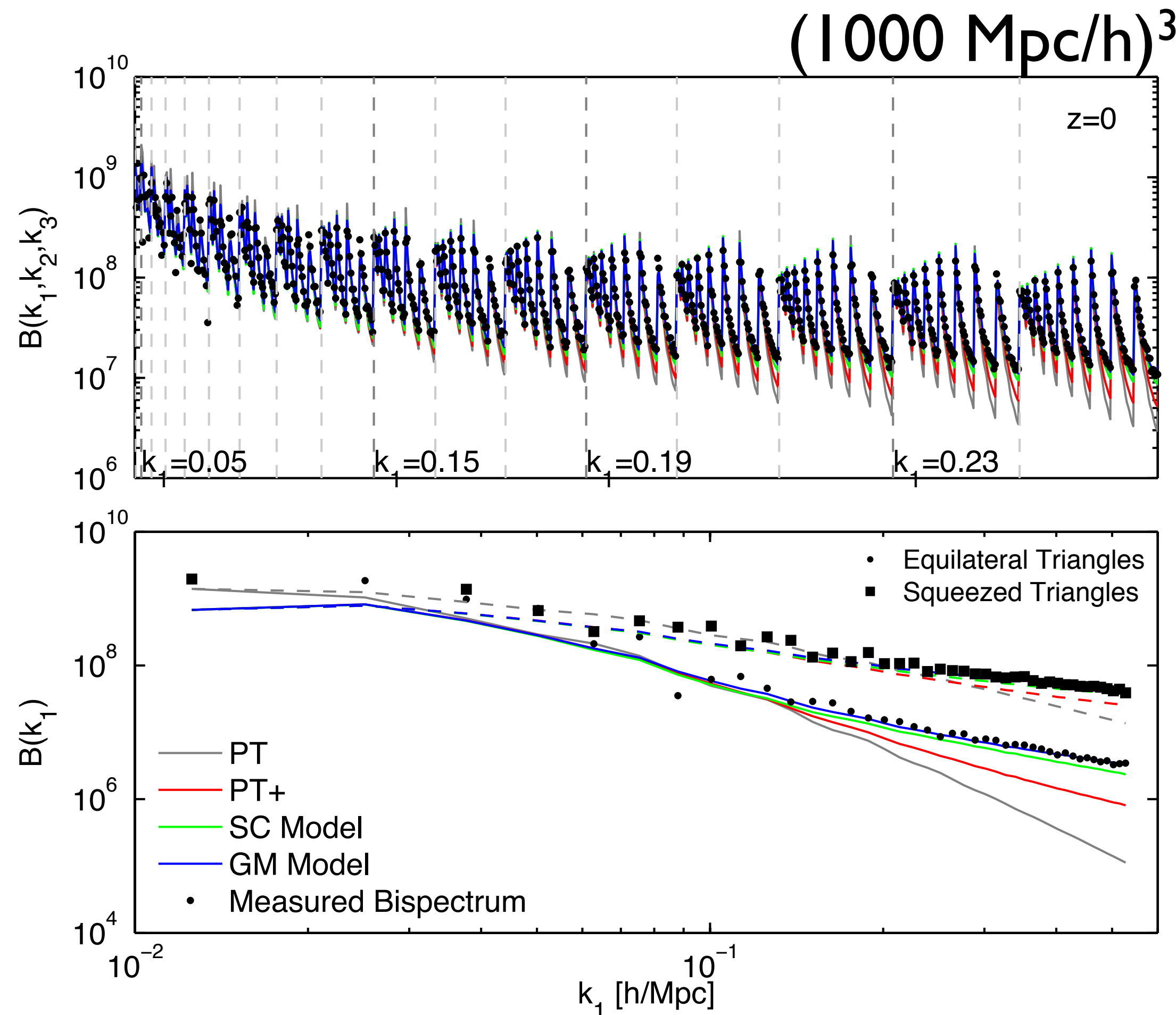


# Transient at $z=1$

$$\frac{Bk(2LPT/Zeldovich, z_{\text{init}})}{Bk(2LPT, z_{\text{init}}=400)}$$

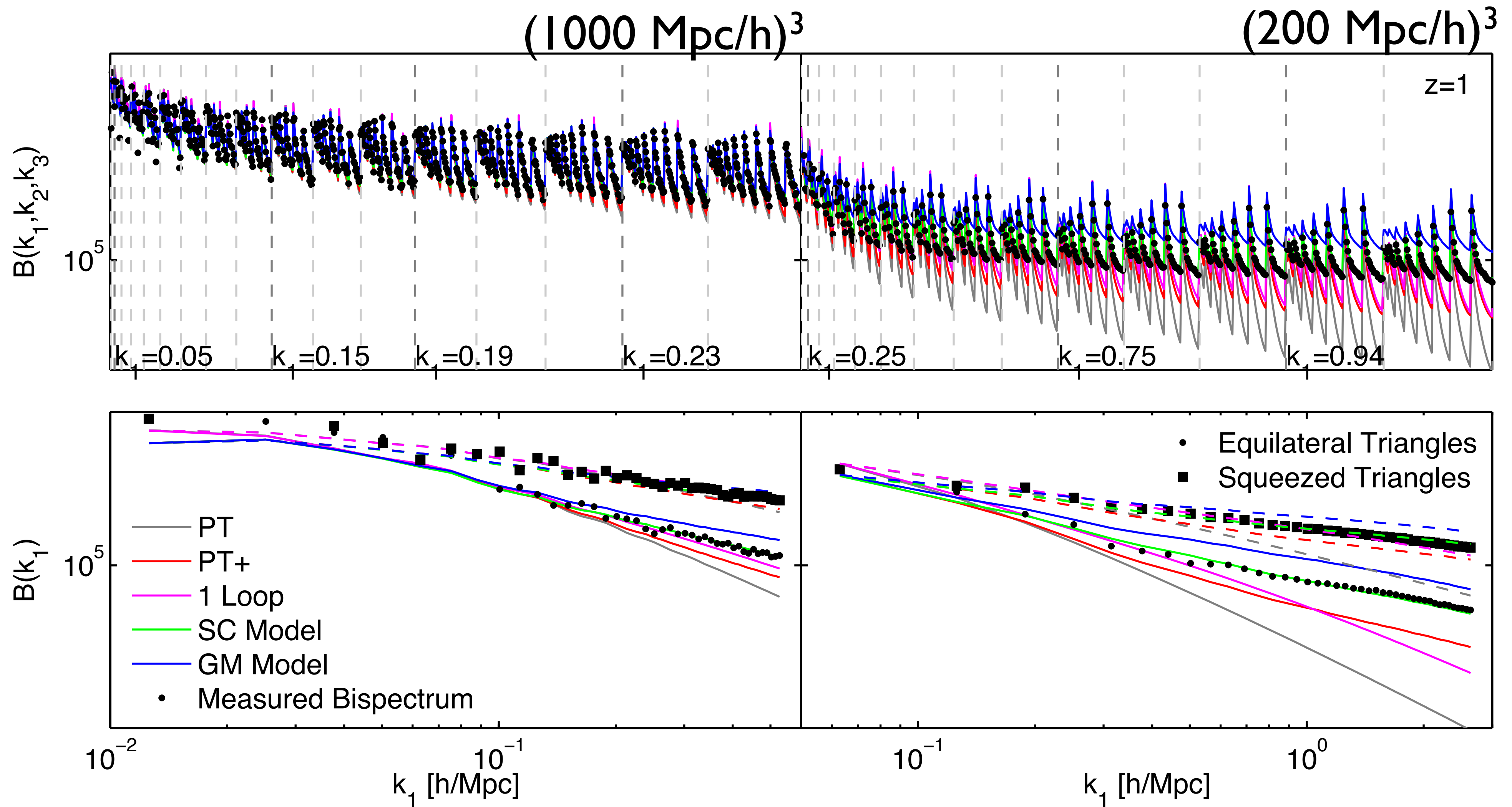
 $(1000 \text{ Mpc}/h)^3$ 
 $Bk(2LPT, z_{\text{init}}=400)$ 
 $(200 \text{ Mpc}/h)^3$ 


# Non-linear bispectrum ( $z=0$ )



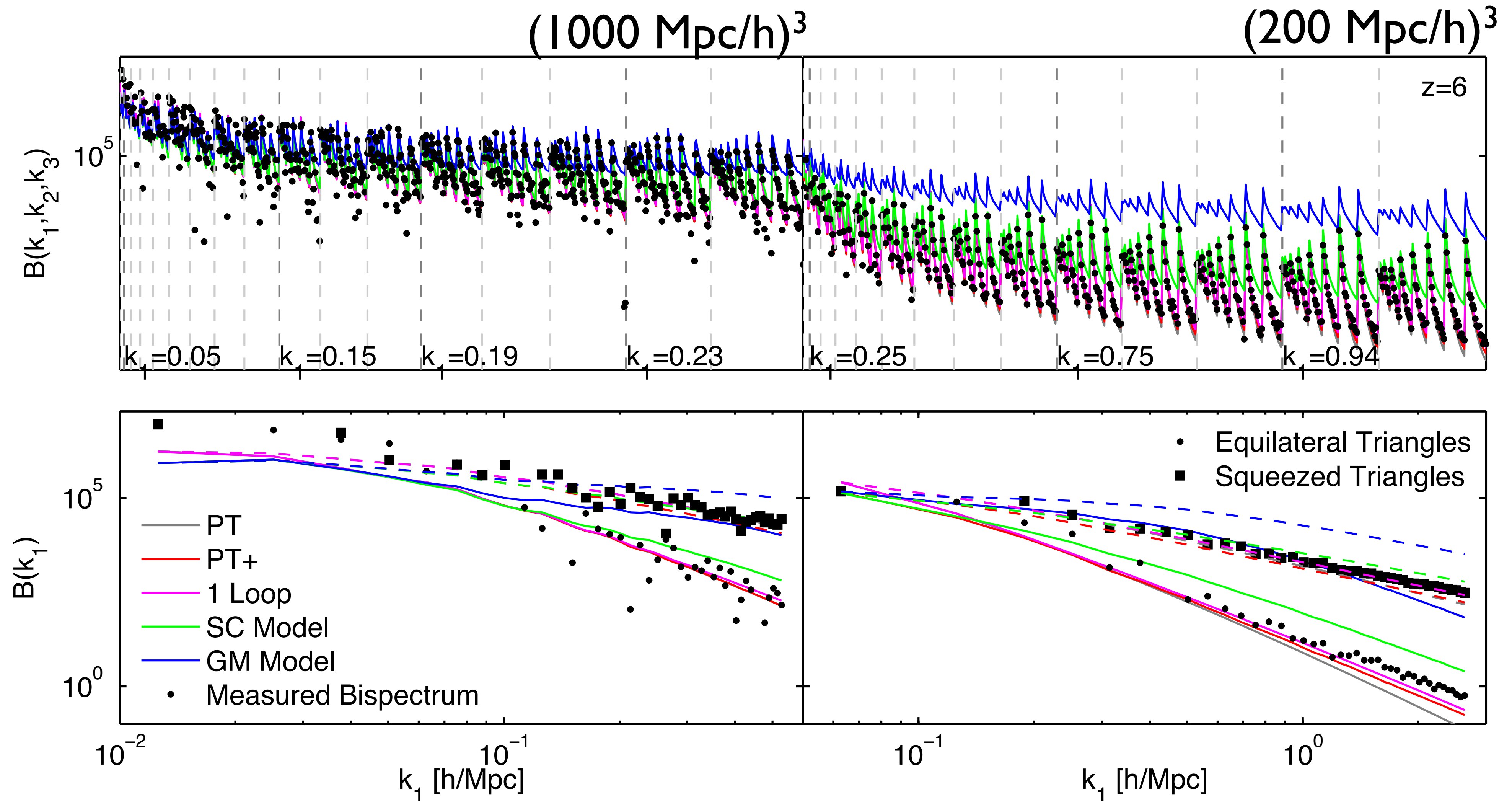
- Simulation from 2LPT,  $z_{\text{ini}}=400$
- scale & configuration dep.
- Deviation from the tree-level (PT) is more apparent for **equilateral-like configurations!**
- Two fitting formulas
  - SC** (Scoccimarro & Couchman 2001)
  - GM** (Gil-Marín+ 2012)

# Non-linear bispectrum ( $z=1$ )





# Non-linear bispectrum ( $z=6$ )

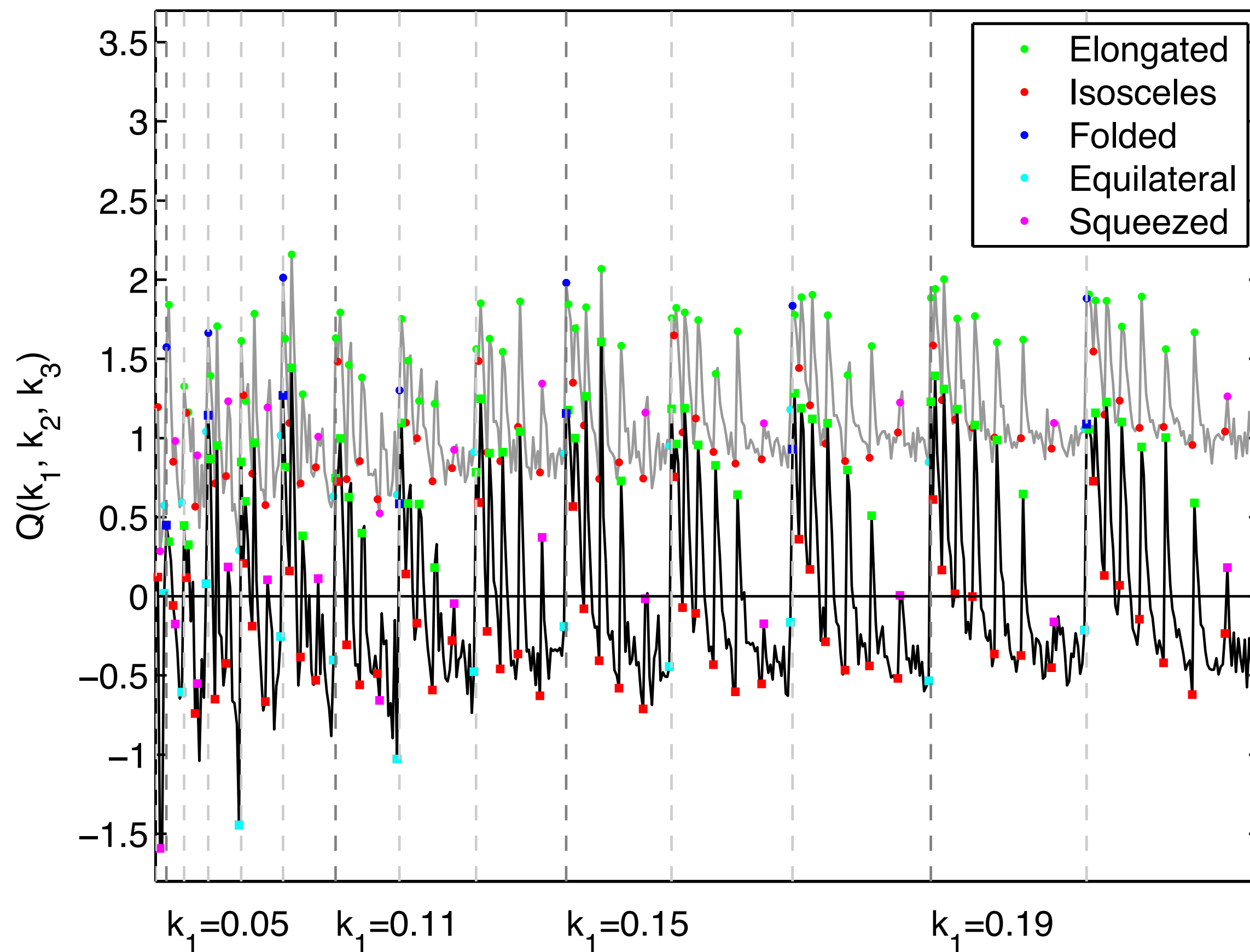


# Log-normal density field

- Cosmological density field is close to log-normal: (Neyrinck+ 2009)
- 1pt pdf indicates the log-normality
- log-transformed density field has lower covariance and better fidelity to the linear theory
- If EXACTLY log-normal, we are done modeling the bispectrum!

$$B(k_1, k_2, k_3) = P(k_1)P(k_2) + (2 \text{ cyclic}) + \int \frac{d^3 q}{(2\pi)^3} P(q) P(|\mathbf{k}_1 + \mathbf{q}|) P(|\mathbf{k}_2 - \mathbf{q}|)$$

# (reduced) Bispectrum of log-density



- reduced bispectrum  
 $Q_{123} = B_{123}/(P_1 P_2 + \text{cyclic})$
- Skewness  $\sim \sum Q_{123} = \text{small}$
- IF log-normal,  $Q$  must be 0!
- Log-normality **CANNOT** be extended beyond the 1pt PDF!



# Conclusion

- To recover correct non-linear bispectrum with  $<1\%$  transient at  $z < 6$ , **2LPT initial condition generator with  $z_{\text{init}} > 100$**  is necessary!
- **Non-linearities** in bispectrum are stronger on **smaller scales and equilateral-like triangles**. One-loop bispectrum works well on ‘smaller scales’, and SC and GM fitting formula works well on highly non-linear scales at  $z=0$  (caution: NOT at high- $z$ ).
- **Log-normality cannot be extended beyond 1 pt PDF.**