DSU 2017 10-14 July 2017 KAIST Munji Campus, Daejeon, Korea

EWKKDM

Seong Chan Park (Yonsei Univ.)



JHEP 1704 (2017) 041 with T. Flacke(IBS, CTPU), D.W. Kang(Yonsei), KC Kong (Kansas), G. Mohlabeng (BNL)

motivation

- unless primordial black holes dominate the whole DM amount, we need something additional...
- among O(100) theoretical ideas, KK DM is still an attractive one
- and its full phenomenology is still to be clarified (interestingly!)

Kaluza-Klein Dark Matter

- Any Symmetric extra dimension (flat or warped)
- If LKP is a neutral EW particle, it is a good DM candidate because of KK-parity
- KK-B₁ (=KK photon) has been studied extensively as RGEs say so within mUED framework.

EW KK gauge boson

• In general, however, the LKP is a mixture of KK excitation modes of neutral EW gauge bosons:

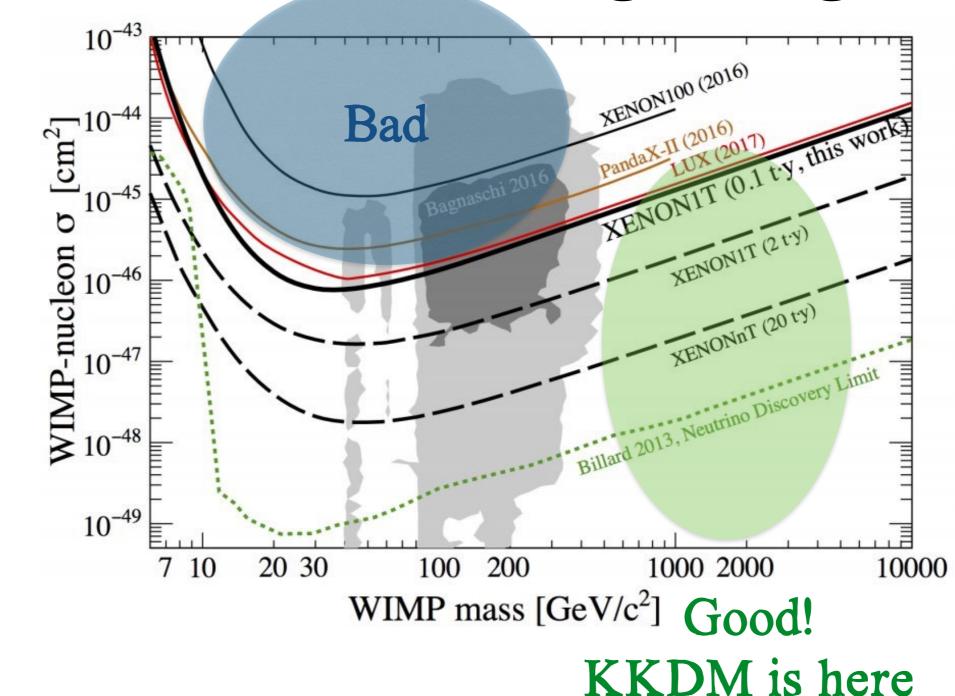
$$\chi = aB_1 + bW_1^3 \qquad \qquad \text{s=1}$$

Just like neutralino

$$\tilde{\chi} = a\tilde{B} + b\tilde{W}^3 + c\tilde{h}_1 + d\tilde{h}_2 \qquad \text{[s=1/2]}$$

 Phenomenology heavily depends on "a" and "b" and these are controlled by "BLKT's" => we study this carefully.

We found: KK DM locates in an interesting range



MODEL: general UED

a review: Flacke, Kong and SCP arXiv:1408.4024

- 5D bulk: $S_1/Z_2 \sim [-L, L]$ (but ANY symmetric space is fine)
- Same gauge symmetries with the SM
- 5D Fermions ~Dirac fields (with bulk masses, in general) chiral zero modes =SM fermions
- 5D Gauge bosons (with boundary localized kinetic terms, in general) zero modes = SM gauge bosons
- KK particles = heavy modes with definite KK-parity
- LKP = the lightest KK neutral gauge boson (B & W) = DM

Bulk Lagrangian

$$S_5 = \int d^4x \int_{-L}^{L} dy \left[\mathcal{L}_V + \mathcal{L}_\Psi + \mathcal{L}_H + \mathcal{L}_{Yuk} \right]$$

 $D_M = \partial_M + i\hat{q}_3\lambda \cdot G_M + i\hat{q}_2\tau \cdot W_M + i\hat{q}_1YB_M$

$$\mathcal{L}_{V} = \sum_{\mathcal{A}}^{G,W,B} -\frac{1}{4}\mathcal{A}^{MN} \cdot \mathcal{A}_{MN}, \qquad [\mathcal{A}_{M}] = [H] = \text{Mass}^{3/2}$$

$$\mathcal{L}_{\Psi} = \sum_{\Psi}^{Q,U,D,L,E} i\overline{\Psi}D_{M}\Gamma^{M}\Psi, \qquad [\hat{g}_{i}] = \text{Mass}^{-1/2}$$

$$\mathcal{L}_{H} = (D_{\mu}H)^{\dagger}D^{\mu}H + \mu_{5}^{2}|H|^{2} - \lambda_{5}|H|^{4}, \qquad [\Psi] = \text{Mass}^{2}$$

$$\mathcal{L}_{Yuk} = \lambda_{5}^{E}\overline{L}HE + \lambda_{5}^{D}\overline{Q}HD + \lambda_{5}^{U}\overline{Q}\tilde{H}D + \text{h.c.}$$

$$[\mu_{5}] = \text{Mass}$$

$$[\lambda_{5}^{\Psi}] = \text{Mass}^{-1/2}$$

Boundary Lagrangian

@ boundaries, 4D Lorentz symmetry respected

$$S_{bdy} = \int d^4x \int_{-L}^{L} dy \left(-\frac{r_W}{4} W_{\mu\nu} \cdot W^{\mu\nu} - \frac{r_B}{4} B_{\mu\nu} B^{\mu\nu}, \right) \left[\delta(y - L) + \delta(y + L) \right]$$

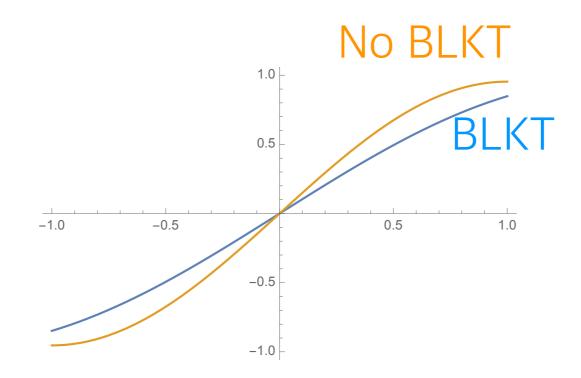
$$[r_W] = [r_B] = \text{Mass}^{-1}$$

Symmetric

(cf) In 'minimal' UED model, $r_W(\Lambda) = 0 = r_B(\Lambda)$ then induced at low scale via RGE $r_W(\mu) \neq 0 \neq r_B(\mu)$

effects of BLKT

Large r => more expels from the boundary => flatten wave function => lighter KK mass



wave function of 1st KK mode

Wave functions w/o BLTs

$$f^{e}(y) = \begin{cases} f_{0}^{e} = \sqrt{\frac{1}{2L}}, \\ f_{2n}^{e} = \sqrt{\frac{1}{L}\cos\frac{2ny}{R}}, \\ f_{2n+1}^{e} = \sqrt{\frac{1}{L}\sin\frac{(2n+1)y}{R}} \end{cases}$$

$$f^{o}(y) = \begin{cases} f_{2n+1}^{o} = \sqrt{\frac{1}{L}\cos\frac{(2n+1)y}{R}}, \\ f_{2n}^{o} = \sqrt{\frac{1}{L}\sin\frac{2ny}{R}}, \end{cases}$$

$$[f^{e/o}] = \text{Mass}^{1/2} = \text{Length}^{-1/2}$$

$$\Psi(x,y) = \sum_{n} \psi_{L}^{n}(x) f_{L}^{n}(y) + \psi_{R}^{n}(x) f_{R}^{n}(y)$$
$$[f^{e/o}] = \text{Mass}^{1/2} = \text{Length}^{-1/2}$$

orthogonality
$$\int_{-L}^{L} dy f_n^* f_m = \delta_{mn}$$

KK spectra

$$m_{\Phi^{(n)}}^2 = (n/R)^2 + m_{\Phi^{(0)}}^2$$

Wave functions with BLKTs

$$f_n^{W/B}(y) = \begin{cases} \mathcal{N}_0^{W/B} & \text{if } n = 0, \\ \mathcal{N}_n^{W/B} \sin(k_n^{W/B} y) & \text{if } n = odd, \\ \mathcal{N}_n^{W/B} \cos(k_n^{W/B} y) & \text{if } n = even, \end{cases} \qquad W_{\mu}(x,y) = \sum_{n=0}^{\infty} W_{\mu}^{(n)}(x) f_n^W(y),$$

$$B_{\mu}(x,y) = \sum_{n=0}^{\infty} B_{\mu}^{(n)}(x) f_n^B(y),$$

$$W_{\mu}(x,y) = \sum_{n=0}^{\infty} W_{\mu}^{(n)}(x) f_n^W(y)$$

$$B_{\mu}(x,y) = \sum_{n=0}^{\infty} B_{\mu}^{(n)}(x) f_n^B(y),$$

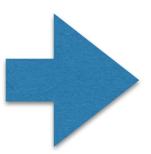
$$\mathcal{N}_{n}^{W/B} = \begin{cases} \frac{1}{\sqrt{2L(1+\frac{r_{W/B}}{L})}} & \text{if } n=0, \\ \frac{1}{\sqrt{L+r_{W/B}\sin^{2}(k_{n}^{W/B}L)}} & \text{if } n=\text{odd}, \\ \frac{1}{\sqrt{L+r_{W/B}\cos^{2}(k_{n}^{W/B}L)}} & \text{if } n=\text{even}. \end{cases}$$

$$\int_{-L}^{L} dy f_{m}^{W/B} f_{n}^{W/B} \left[1+r_{W/B}\left(\delta(y+L)+\delta(y-L)\right)\right] = \delta_{mn}$$

$$\text{Orthogonality}$$

$$\int_{-L}^{L} dy f_{m}^{W/B} f_{n}^{W/B} \left[1 + r_{W/B} \left(\delta(y+L) + \delta(y-L)\right)
ight] = \delta_{mn}$$
 orthogonality

$$\cot(k_n^{W/B}L) = r_{W/B}k_n^{W/B}$$
 if $n = odd$,
$$\tan(k_n^{W/B}L) = -r_{W/B}k_n^{W/B}$$
 if $n = even$.



KK spectrum

4D effective action

$$S_{4D} \ni \int d^4x \left\{ \sum_{n} \left[-\frac{1}{4} \sum_{n} B^{(n)\mu\nu} B^{(n)}_{\mu\nu} - \frac{\left(k_n^B\right)^2}{2} B^{(n)\mu} B^{(n)}_{\mu} \right. \right. \\ \left. -\frac{1}{4} \sum_{n} W^{(n)a\mu\nu} \cdot W^{(n)a}_{\mu\nu} - \frac{\left(k_n^W\right)^2}{2} W^{(n)a\mu} W^{(n)a}_{\mu} \right] \right. \\ \left. + \sum_{m,n} \left[-\frac{\hat{g}_1^2 v^2}{8} \mathcal{F}^{BB}_{mn} B^{(m)\mu} B^{(n)}_{\mu} - \frac{\hat{g}_1 \hat{g}_2 v^2}{8} \mathcal{F}^{WB}_{mn} B^{(m)\mu} W^{(n)3}_{\mu} \right. \\ \left. -\frac{\hat{g}_2^2 v^2}{8} \mathcal{F}^{WW}_{mn} W^{(m)a\mu} W^{(n)a}_{\mu} \right] \right\} ,$$

$$\mathcal{F}^{BB}_{mn} = \int_{-L}^{L} rac{dy}{2L} \, f^B_m(y) f^B_n(y) \, ,$$

Wave function overlaps: $\mathcal{F}_{mn}^{BW} = \int_{-L}^{L} \frac{dy}{2L} f_m^B(y) f_n^W(y)$,

$$\mathcal{F}_{mn}^{WW} = \int_{-L}^{L} rac{dy}{2L} f_m^W(y) f_n^W(y) \,,$$

Masses of odd modes

$$M_{n,odd}^2 = \begin{pmatrix} \left(k_1^B\right)^2 + \frac{\hat{g_1}^2 \hat{v}^2}{4} \mathcal{F}_{11}^{BB} & \frac{\hat{g_1} \hat{g_2} \hat{v}^2}{4} \mathcal{F}_{11}^{BW} & \dots \\ \frac{\hat{g_1} \hat{g_2} \hat{v}^2}{4} \mathcal{F}_{11}^{BW} & \left(k_1^W\right)^2 + \frac{\hat{g_2}^2 \hat{v}^2}{4} \mathcal{F}_{11}^{WW} & \dots \\ \vdots & \vdots & \ddots \end{pmatrix}$$

Mass eigenstates

$$A_{1,2} = a_{1,2}B_1 + b_{1,2}W_1^3$$

the lighter one is LKP = DM

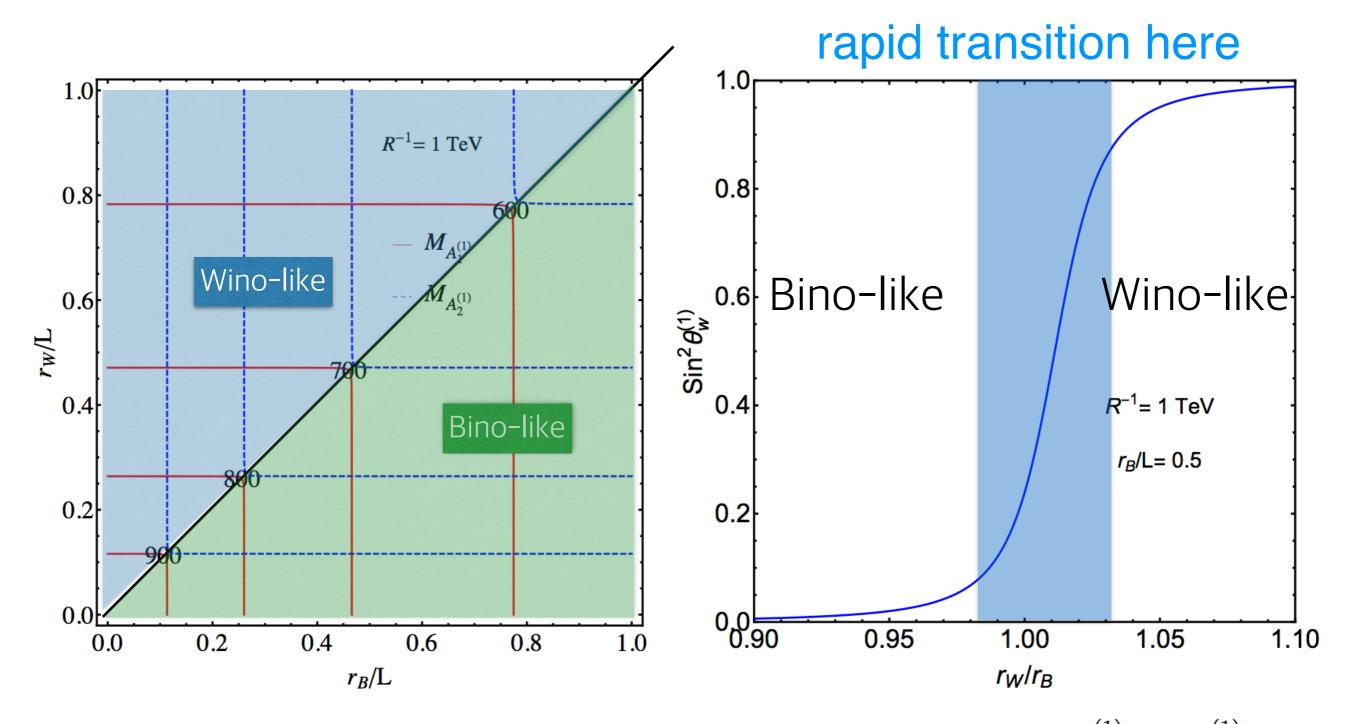


Figure 1: Left: Contours of constant mass for the level 1 electroweak KK bosons $A_1^{(1)}$ and $A_2^{(1)}$. The contours were made assuming $R^{-1} = 1$ TeV, and they show the mass dependence on the boundary terms r_B and r_W . Right: The level 1 KK Weinberg angle $\sin^2 \theta_W^{(1)}$ for $R^{-1} = 1$ TeV and $r_B/L = 0.5$.

Mass matrix for neutral gauge bosons (even modes)

basis:
$$B^{(2n)} - W^{3(2n)}$$

$$M_{n,e}^{2} = \begin{pmatrix} \frac{\hat{g}_{1}^{2}v^{2}}{4}\mathcal{F}_{00}^{BB} & \frac{\hat{g}_{1}\hat{g}_{2}v^{2}}{4}\mathcal{F}_{00}^{BW} & \frac{\hat{g}_{1}^{2}v^{2}}{4}\mathcal{F}_{02}^{BB} & \frac{\hat{g}_{1}\hat{g}_{2}v^{2}}{4}\mathcal{F}_{02}^{BW} & \cdots \\ \frac{\hat{g}_{1}\hat{g}_{2}v^{2}}{4}\mathcal{F}_{00}^{BW} & \frac{\hat{g}_{2}^{2}v^{2}}{4}\mathcal{F}_{00}^{BW} & \frac{\hat{g}_{1}\hat{g}_{2}v^{2}}{4}\mathcal{F}_{02}^{BW} & \frac{\hat{g}_{2}^{2}v^{2}}{4}\mathcal{F}_{02}^{BW} & \cdots \\ \frac{\hat{g}_{1}^{2}v^{2}}{4}\mathcal{F}_{20}^{BB} & \frac{\hat{g}_{1}\hat{g}_{2}v^{2}}{4}\mathcal{F}_{20}^{BW} & (k_{2}^{B})^{2} + \frac{\hat{g}_{1}^{2}v^{2}}{4}\mathcal{F}_{22}^{BB} & \frac{\hat{g}_{1}\hat{g}_{2}v^{2}}{4}\mathcal{F}_{22}^{BW} & \cdots \\ \vdots & \vdots & \vdots & \vdots & \ddots \end{pmatrix}$$

$$\mathcal{F}_{mn}^{BB} = \int_{-L}^{L} rac{dy}{2L} f_{m}^{B}(y) f_{n}^{B}(y) ,$$
 $\mathcal{F}_{mn}^{BW} = \int_{-L}^{L} rac{dy}{2L} f_{m}^{B}(y) f_{n}^{W}(y) ,$ $\mathcal{F}_{mn}^{WW} = \int_{-L}^{L} rac{dy}{2L} f_{m}^{W}(y) f_{n}^{W}(y) ,$

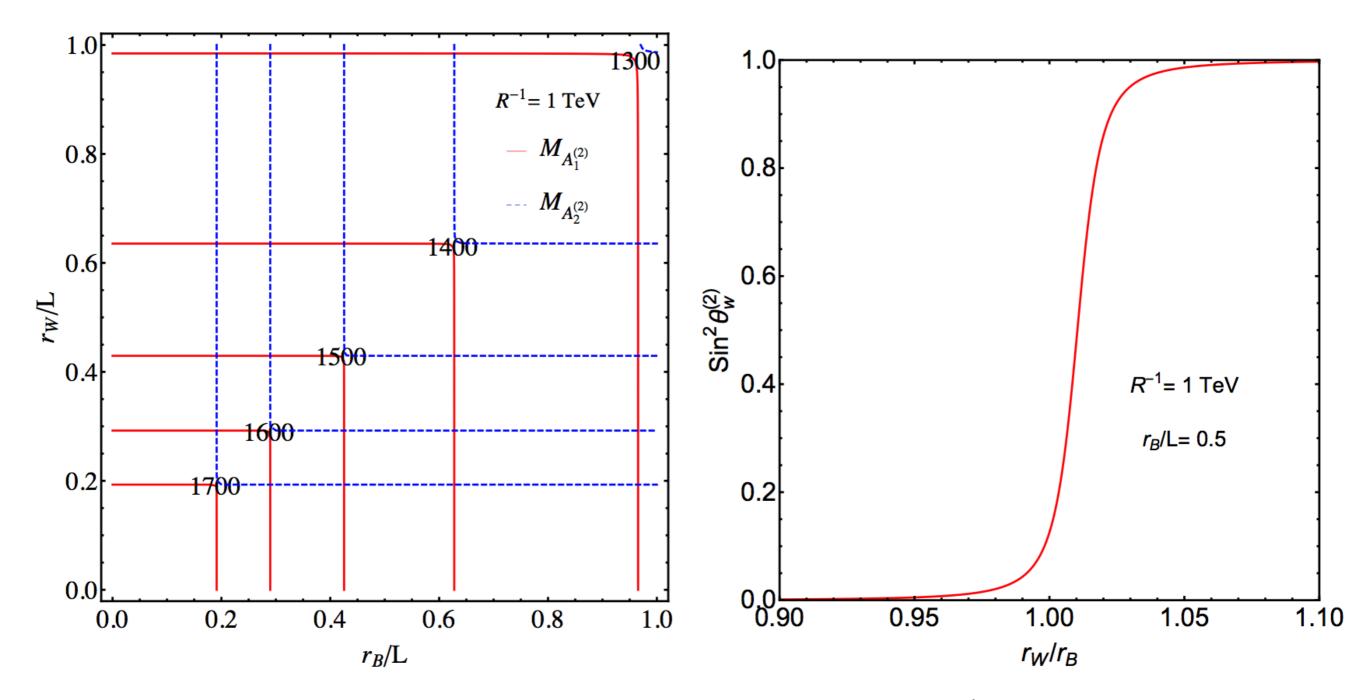
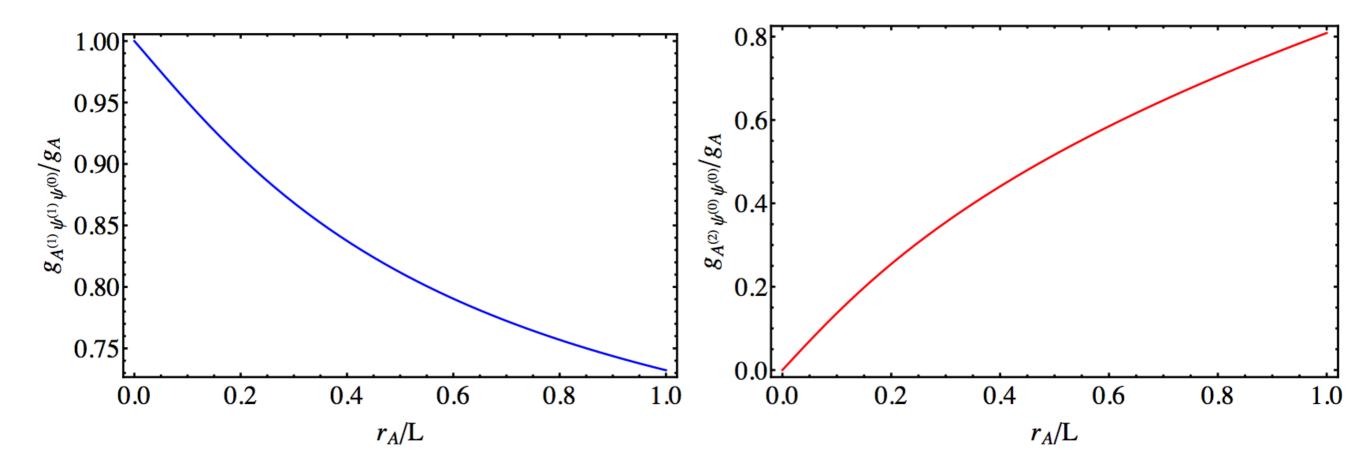


Figure 2: Left: Contours of constant level 2 gauge boson masses for $R^{-1} = 1$ TeV. The red contours show the mass of the lighter eigenstate $A_1^{(2)}$ and the blue contour represents mass of heavier eigenstate $A_2^{(2)}$. Right: The level 2 KK Weinberg angle $\sin^2 \theta_W^{(2)}$ for $R^{-1} = 1$ TeV and $r_B/L = 0.5$.

Effective couplings: A-f-f

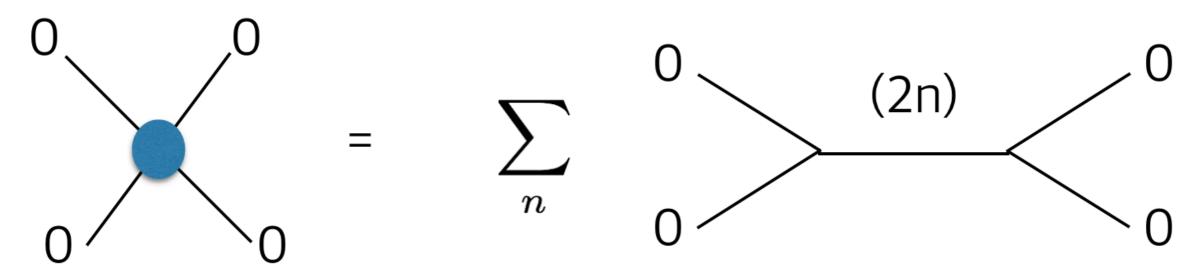
$$g_{\mathcal{A}^{(\ell)}\psi^{(m)}\psi^{(n)}} = g_{\mathcal{A}}\tilde{\mathcal{F}}_{\ell mn}^{\mathcal{A}}$$
$$\tilde{\mathcal{F}}_{\ell mn}^{\mathcal{A}} \equiv \frac{1}{\mathcal{N}_{0}^{\mathcal{A}}} \int_{-L}^{L} dy \, f_{\ell}^{\mathcal{A}}(y) f_{m}^{\psi}(y) f_{n}^{\psi}(y),$$



N.B. KK-parity is conserved but KK number is not.

Experimental constraints

Bounds: 4-fermi interactions



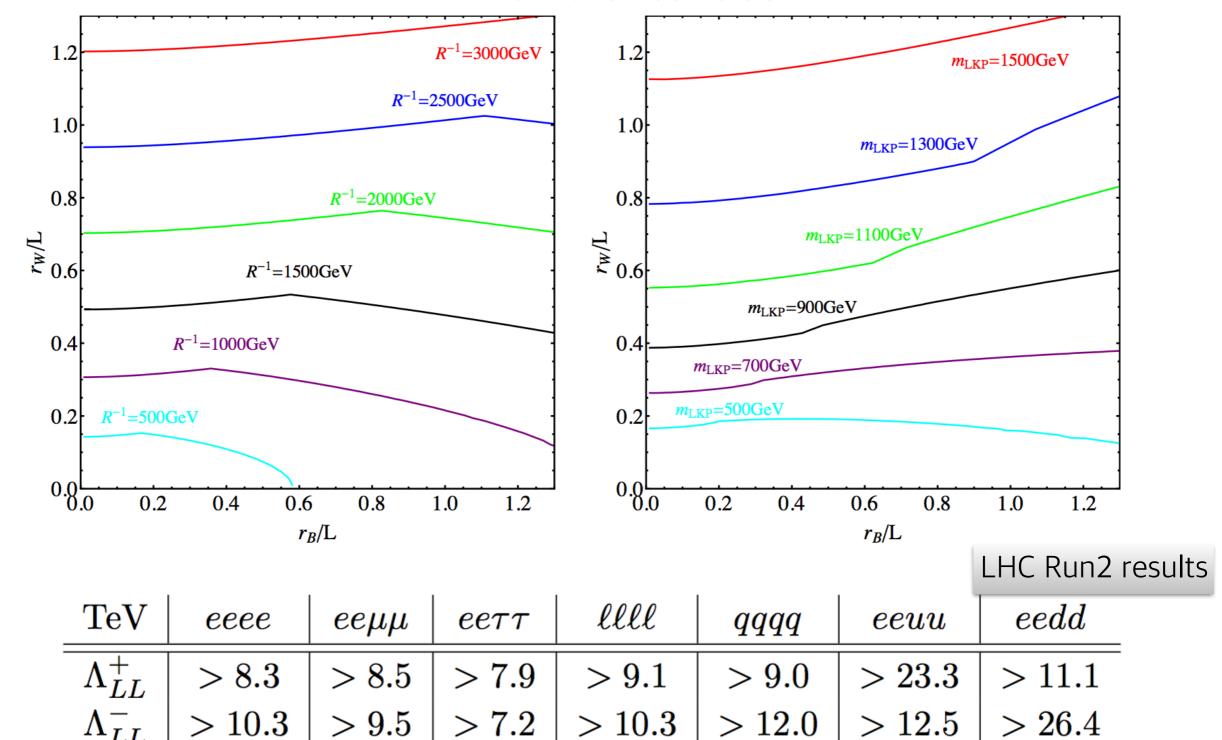
$$\mathcal{L}_{eff} \supset \sum_{f1,f2} \sum_{A,B=L,R} \eta^s_{f_1 f_2,AB} \frac{4\pi}{(\Lambda^s_{f1 f2,AB})^2} \overline{f}_{1,A} \gamma^{\mu} f_{1,A} \overline{f}_{2,B} \gamma_{\mu} f_{2,B},$$

$$\frac{4\pi}{\Lambda_{eq,AB}^2} \eta_{eq,AB} = 4\pi N_c \left[\sum_{n=1}^{\infty} (\tilde{\mathcal{F}}_{2n00}^B)^2 \frac{3}{5} \frac{\alpha_1 Y_{eA} Y_{qB}}{Q^2 - M_{B^{(2n)}}^2} + \sum_{n=1}^{\infty} (\tilde{\mathcal{F}}_{2n00}^W)^2 \frac{\alpha_2 T_{eA}^3 T_{qB}^3}{Q^2 - M_{W_3^{(2n)}}^2} \right]$$

$$\approx -12\pi \left[\sum_{n=1}^{\infty} (\tilde{\mathcal{F}}_{2n00}^B)^2 \frac{3}{5} \frac{\alpha_1 Y_{eA} Y_{qB}}{M_{B^{(2n)}}^2} + \sum_{n=1}^{\infty} (\tilde{\mathcal{F}}_{2n00}^W)^2 \frac{\alpha_2 T_{eA}^3 T_{qB}^3}{M_{W_3^{(2n)}}^2} \right].$$

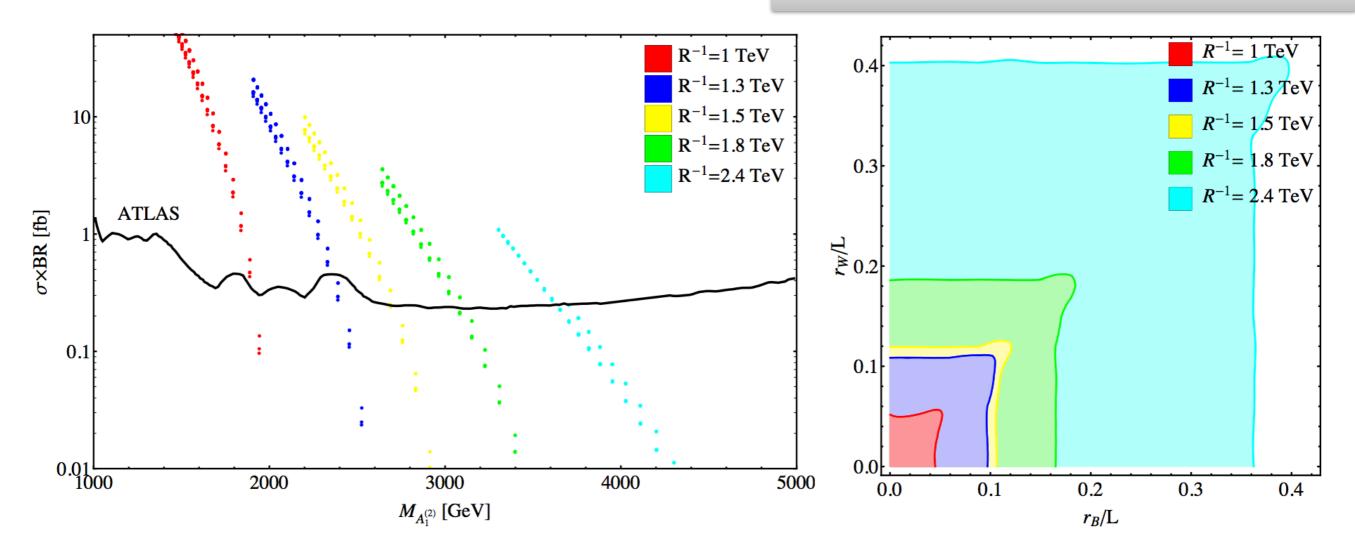
4-fermi constraints

minimum allowed values



Dilepton resonance search at the **LHC**

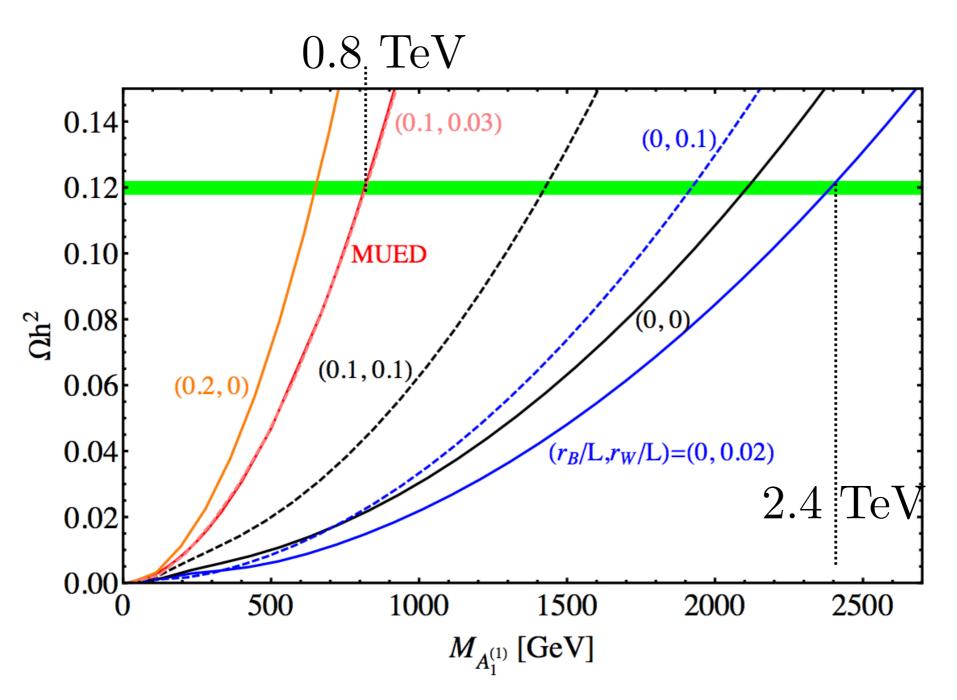
13 TeV with 13.3 /fb (ATLAS), 13.0/fb (CMS)



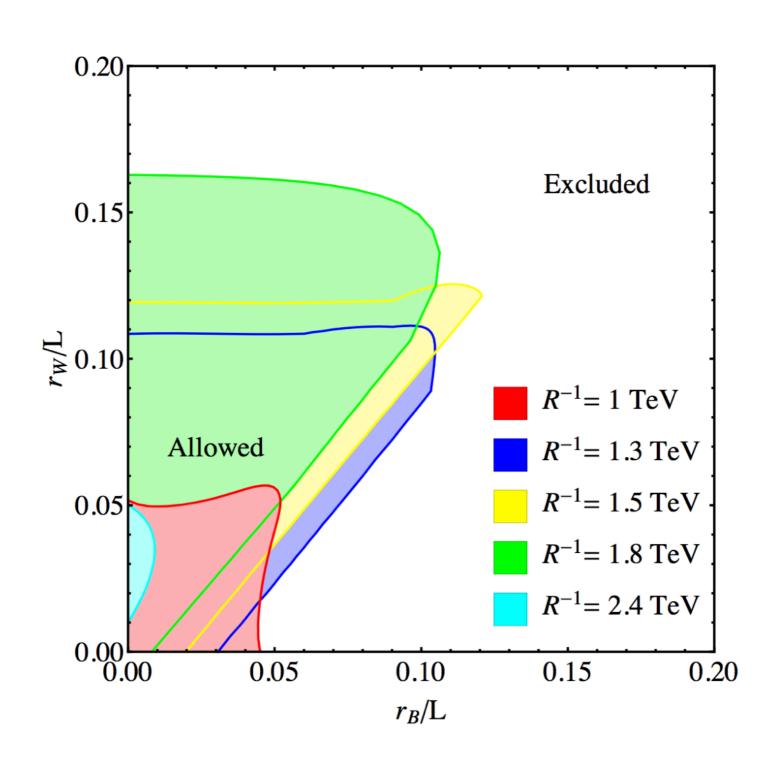
$$g_{W^{(2)}\psi^{(0)}\psi^{(0)}} = g_2 \int_{-L}^{L} \frac{dy}{2L} \frac{f_2^W(y)}{\mathcal{N}_0^W} = g_2 \sqrt{\frac{2(1 + r_W/L)}{1 + \frac{r_W}{L}\cos^2(k_2^W L)}} \frac{\sin(k_2^W L)}{k_2^W L}$$

$$g_{B^{(2)}\psi^{(0)}\psi^{(0)}} = g_1 \int_{-L}^{L} \frac{dy}{2L} \frac{f_2^B(y)}{\mathcal{N}_0^B} = g_1 \sqrt{\frac{2(1 + r_B/L)}{1 + \frac{r_B}{L}\cos^2(k_2^B L)}} \frac{\sin(k_2^B L)}{k_2^B L}.$$

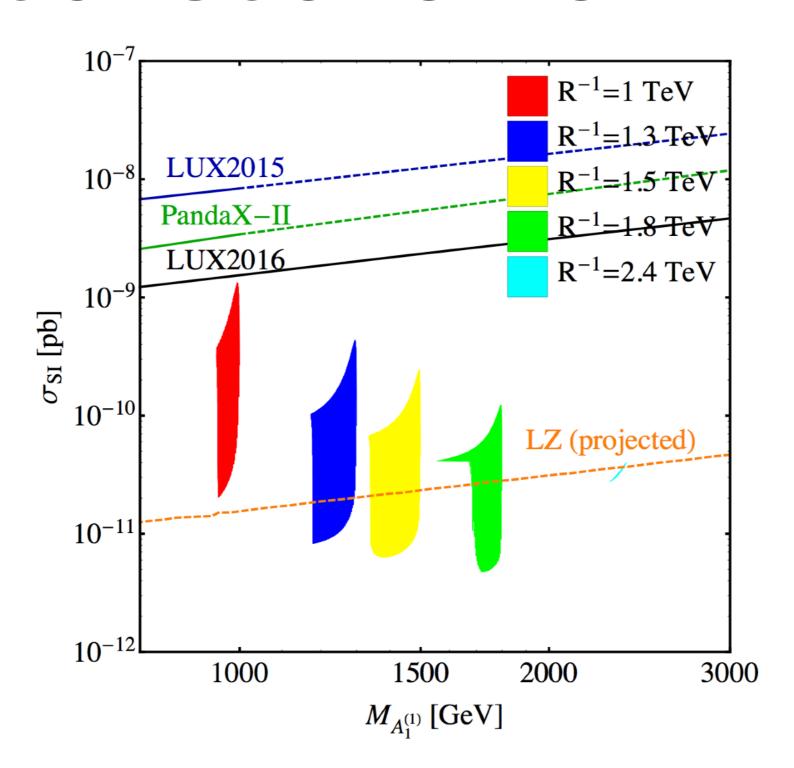
Relic abundance



LHC + relic abundance



Direct section of KK DM



Conclusion

- EWKKDM = LKP of the neutral gauge boson (mixture of the KK states of B and W₃)
- Depending on BLKT, the LKP can be either B-like or W-like or in between.
- Phenomenology is rich. LKP can be as heavy as 2.4
 TeV (without BLKT, KK-photon DM < 800 GeV (1.3
 TeV with resonance effect))
- Future DD experiments will cover more. We will see!