

FLAVOR CONSTRAINTS ON AXION-LIKE PARTICLE

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(IN PREPARATION)



OUTLINE

1. Axion-like particle couplings to the SM
2. SM induced FCNC with ALP (Model independent)
3. Analysis
4. Specific model (2HDM)
5. Summary

AXION-LIKE PARTICLE (ALP)

- Light and Weakly coupled to the SM
- Global $U(1)$ PQ symmetry
- Good BSM candidate (e.g. DM, Strong CP, etc)
- Experimental bounds on ALP portal (e.g. Rare meson decay)

ALP COUPLINGS

- PQ-Symmetry

$$a \rightarrow a + c$$

- Derivative couplings

$$\mathcal{L} \supset \frac{\partial_\mu a}{f_a} \sum_{f=Q,u_R,d_R,L,e_R} c_f \bar{f} \gamma^\mu f + V(a)$$

ALP COUPLINGS

(DERIVATIVE VS. NON-DERIVATIVE)

$$\frac{\partial_\mu a}{f_{\text{PQ}}} \bar{f}_L \gamma^\mu f_L$$

$$-\frac{m_f}{f_{\text{PQ}}} i a \bar{f}_R f_L$$

- Gauge invariance
- Relevant at high energy scale

ALP COUPLINGS

(DERIVATIVE VS. NON-DERIVATIVE)

$$c_f \frac{\partial_\mu a}{f_{\text{PQ}}} \bar{f} \gamma^\mu \gamma^5 f$$

$$-2c_f \frac{ia}{f_{\text{PQ}}} m_f \bar{f} \gamma^5 f$$

ALP COUPLINGS

(DERIVATIVE VS. NON-DERIVATIVE)

$$c_f \frac{\partial_\mu a}{f_{\text{PQ}}} \bar{f} \gamma^\mu \gamma^5 f$$

$$-m_f \bar{f} \exp \left[2c_f \frac{ia}{f_{\text{PQ}}} \gamma^5 \right] f$$

- ex)

Axion scattering with electron ($e+a \rightarrow e+a$)

Axion emission by nucleon via 1 pion exchange process ($NN \rightarrow NN a$)

ALP COUPLINGS

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RADIATIVE FCNC OF ALP

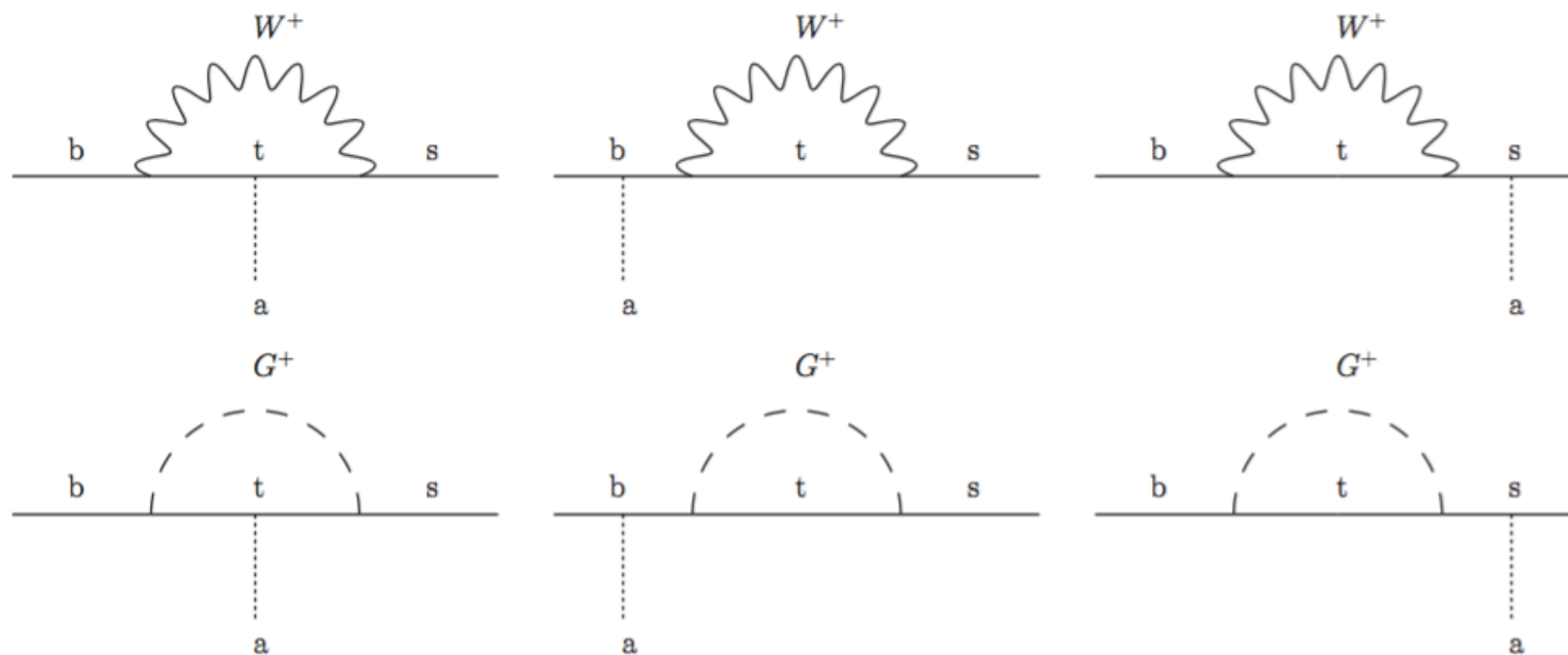


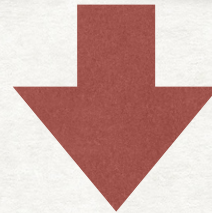
FIG. 1: Relevant diagrams for FCNC of ALP

RADIATIVE FCNC OF ALP

$$\frac{d}{d \ln \mu} c_Q = \frac{1}{(4\pi)^2} \left[(c_Q - c_u) y_u y_u^\dagger + (c_Q - c_d) y_d y_d^\dagger \right]$$

RADIATIVE FCNC OF ALP

$$\frac{d}{d \ln \mu} c_Q = \frac{1}{(4\pi)^2} \left[(c_Q - c_u) y_u y_u^\dagger + (c_Q - c_d) y_d y_d^\dagger \right]$$



$$\left[\frac{(y_t^{SM})^2}{16\pi^2} \log \frac{\Lambda_{UV}^2}{m_t^2} \left(-\frac{1}{2} c_Q + \frac{1}{2} c_u \right) V_{CKM}^{*t\alpha} V_{CKM}^{t\beta} \frac{a}{f_a} \bar{d}_\alpha [(im_d)_\alpha P_L - (im_d)_\beta P_R] d_\beta \right]$$

LOW-SCALE EFFECTIVE COUPLINGS

- At high energy scale above the electroweak scale,

$$\frac{\partial_\mu a}{f_a} \left(c_Q \bar{Q} \gamma^\mu Q + c_u \bar{u}_R \gamma^\mu u_R + c_d \bar{d}_R \gamma^\mu d_R + c_L \bar{L} \gamma^\mu L + c_e \bar{e}_R \gamma^\mu e_R + c_H H^\dagger i \overleftrightarrow{D}^\mu H \right) \\ - \frac{a}{f_a} \left(c_{agg} \frac{g_s^2}{32\pi^2} G \tilde{G} + c_{aww} \frac{g^2}{32\pi^2} W \tilde{W} + c_{abb} \frac{g'^2}{32\pi^2} B \tilde{B} \right) + \mathcal{L}_a$$


- At QCD chiral symmetry breaking scale, the chiral perturbation theory should be considered to describe the physics of the strongly interacting particles

LOW-SCALE EFFECTIVE COUPLINGS

$$\mathcal{L} \supset \frac{\partial_\mu a}{f_{\text{PQ}}} A_l \bar{l} \gamma^\mu \gamma^5 l - \bar{c}_{a\gamma\gamma} \frac{a}{f_{\text{PQ}}} \frac{e^2}{32\pi^2} F \tilde{F} + \mathcal{L}_{aN} + \mathcal{L}_{a\pi}$$

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$$A_l = (-c_L + c_e + c_H)/2$$

LOW-SCALE EFFECTIVE COUPLINGS

$$\mathcal{L} \supset \frac{\partial_\mu a}{f_{PQ}} A_l \bar{l} \gamma^\mu \gamma^5 l - \bar{c}_{a\gamma\gamma} \frac{a}{f_{PQ}} \frac{e^2}{32\pi^2} F \tilde{F} + \mathcal{L}_{aN} + \mathcal{L}_{a\pi}$$

$$q_A = \frac{C_{agg}}{2} \frac{M^{-1}}{\text{tr}[M^{-1}]}$$

$$\begin{aligned} \bar{c}_{a\gamma\gamma} = & c_{a\gamma\gamma} - 12 \text{tr}[q_A Q_E Q_E] - \sum_{f=\mu,e} 4A_f Q_f^2 [1 - 2A(\tau_f)] \\ & - 2\kappa_0 \frac{m_a^2}{m_\pi^2 - m_a^2} - 1.3\kappa_\eta \frac{m_a^2}{m_\eta^2 - m_a^2} - 2.9\kappa_{\eta'} \frac{m_a^2}{m_{\eta'}^2 - m_a^2} \end{aligned}$$

$$\begin{aligned} \kappa_0 &= \frac{A_u - A_d}{2} - \frac{C_{agg}}{2} \frac{m_u^{-1} - m_d^{-1}}{m_u^{-1} + m_d^{-1} + m_s^{-1}} \\ \kappa_8 &= \frac{A_u - A_d}{2\sqrt{3}} - \frac{C_{agg}}{2\sqrt{3}} \frac{m_u^{-1} + m_d^{-1} - 2m_s^{-1}}{m_u^{-1} + m_d^{-1} + m_s^{-1}} \\ \kappa_9 &= \frac{A_u + 2A_d}{\sqrt{6}} - \frac{C_{agg}}{\sqrt{6}} \end{aligned}$$

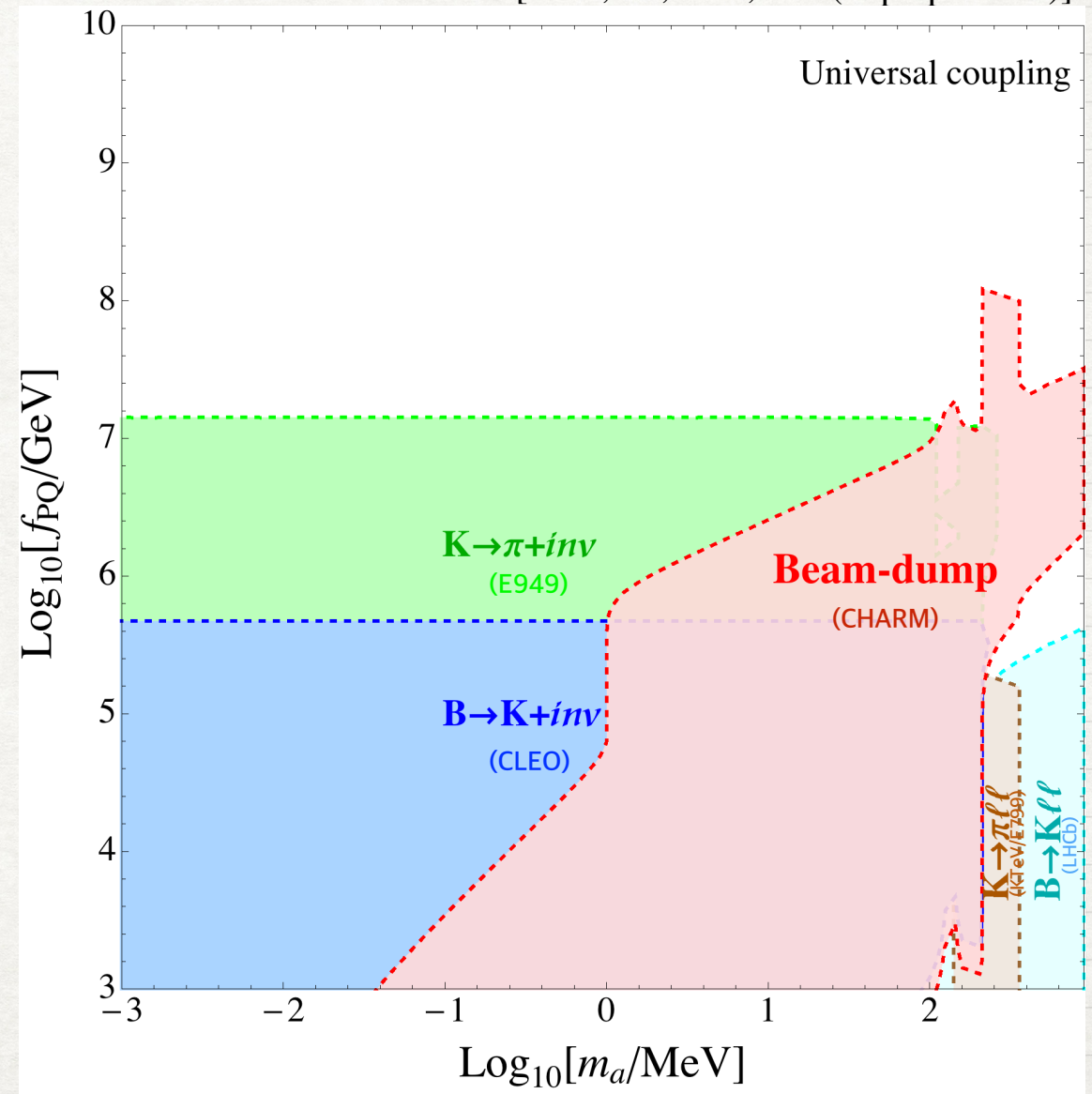
$$\begin{aligned} \kappa_\eta &= \kappa_8 \cos \theta_{\eta\eta'} - \kappa_9 \sin \theta_{\eta\eta'} \\ \kappa_{\eta'} &= \kappa_8 \sin \theta_{\eta\eta'} + \kappa_9 \cos \theta_{\eta\eta'} \end{aligned}$$

EXPERIMENTAL BOUNDS

- Universal couplings
- Flavor constraints

$$f_{PQ} \gtrsim 10^7 \text{ GeV}$$

[Choi, Im, Park, Yun (in preparation)]



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- Universal couplings

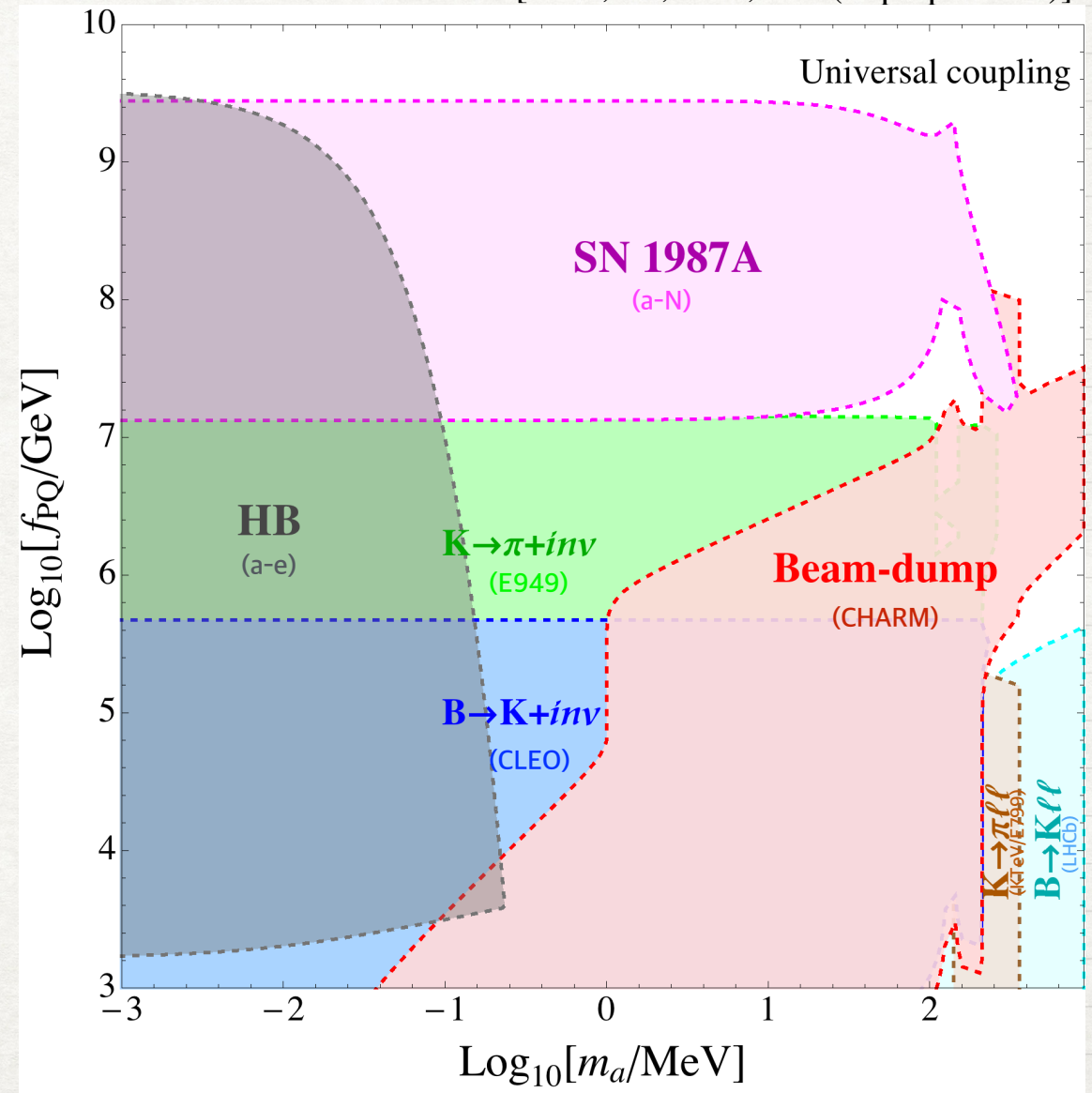
- Flavor constraints

$$f_{PQ} \gtrsim 10^7 \text{ GeV}$$

- Axion mass < 100 MeV

$$f_{PQ} \gtrsim 10^9 \text{ GeV}$$

[Choi, Im, Park, Yun (in preparation)]



SPECIFIC MODEL : ALP + 2HDM

- SM couplings

$$\frac{\partial_\mu a}{f_{\text{PQ}}} \left(C_Q \bar{Q} \gamma^\mu Q + C_u \bar{u}_R \gamma^\mu u_R + C_d \bar{d}_R \gamma^\mu d_R + C_L \bar{L} \gamma^\mu L + C_e \bar{e}_R \gamma^\mu e_R + \sum_{i=1,2} C_i H_i^\dagger i \overleftrightarrow{D}^\mu H_i \right)$$

- Pseudoscalar mixing (Goldstone boson eaten by Z)

$$Z_\mu \rightarrow Z_\mu - \frac{2}{g_Z} \frac{\partial_\mu a}{f_{\text{PQ}}} (\cos^2 \beta C_1 + \sin^2 \beta C_2)$$
$$A_\mu \rightarrow A_\mu + \cos^2 \theta_W \frac{2}{e} \frac{\partial_\mu a}{f_{\text{PQ}}} (\cos^2 \beta C_1 + \sin^2 \beta C_2)$$

SPECIFIC MODEL : ALP + 2HDM

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$$\frac{\partial_\mu a}{f_{\text{PQ}}} \left(C_Q \bar{Q} \gamma^\mu Q + C_u \bar{u}_R \gamma^\mu u_R + C_d \bar{d}_R \gamma^\mu d_R + C_L \bar{L} \gamma^\mu L + C_e \bar{e}_R \gamma^\mu e_R + \sum_{i=1,2} C_i H_i^\dagger i \overleftrightarrow{D}^\mu H_i \right)$$

- Additional massive charged Higgs boson

$$\begin{aligned} & -i (C_1 - C_2) \cos 2\beta \frac{\partial_\mu a}{f_{\text{PQ}}} (\partial^\mu H^- \cdot H^+ - \partial^\mu H^+ \cdot H^-) \\ & -m_W (C_1 - C_2) \sin 2\beta \frac{\partial^\mu a}{f_{\text{PQ}}} (W_\mu^+ H^- + W_\mu^- H^+) \end{aligned}$$

SPECIFIC MODEL : ALP + 2HDM

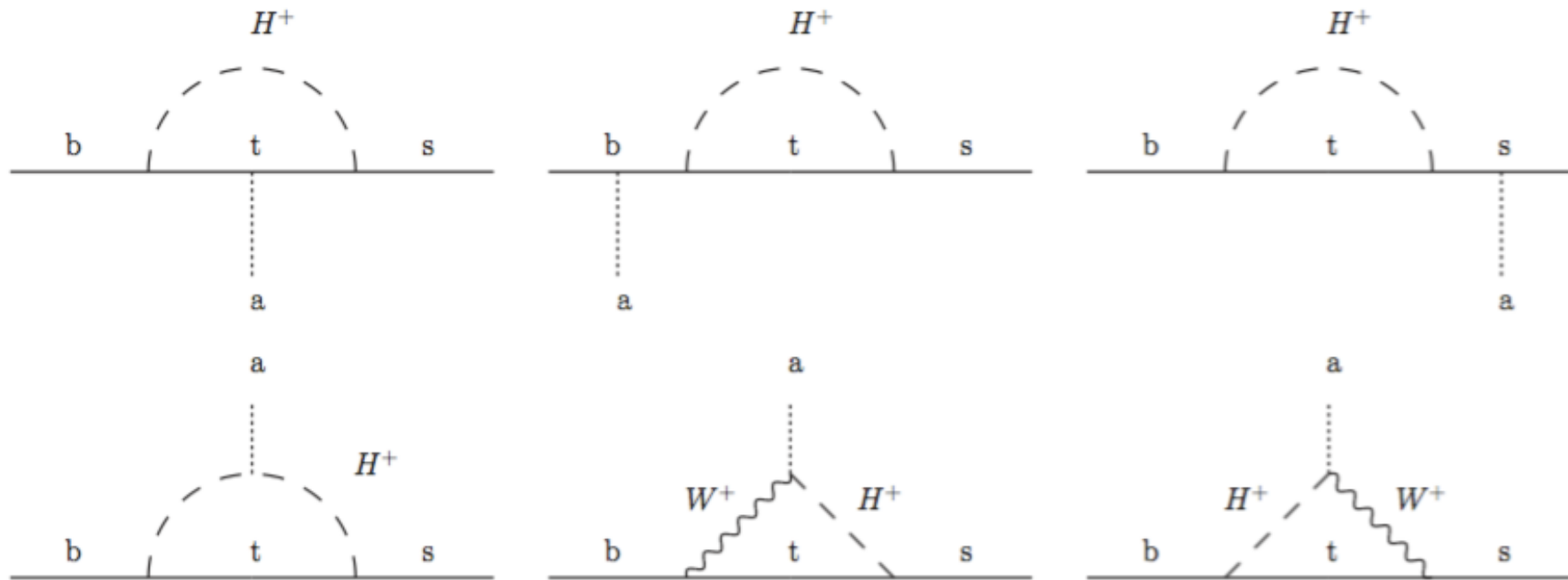


FIG. 2: Charged higgs loops to contribute FCNC of a

SPECIFIC MODEL : ALP + 2HDM

$$\begin{aligned}
 & \frac{(y_t^{SM})^2}{16\pi^2} V_{CKM}^{*t\alpha} V_{CKM}^{t\beta} \frac{a}{f_{PQ}} \bar{d}_\alpha [(im_d)_\alpha P_L - (im_d)_\beta P_R] d_\beta \\
 & \times \left(\log \frac{\Lambda_{UV}^2}{m_{H^\pm}^2} \left(\frac{1}{2} C_Q - \frac{1}{2} C_u + \frac{1}{2} C_2 \right) \frac{1}{\sin^2 \beta} \right. \\
 & \left. + \log \frac{m_{H^\pm}^2}{m_t^2} \left[\left(\frac{1}{2} C_Q - \frac{1}{2} C_u + \frac{1}{2} C_2 \right) + \frac{C_1 - C_2}{2} \cos^2 \beta \right] \right)
 \end{aligned}$$

SPECIFIC MODEL : ALP + 2HDM

$SM + H^\pm$

$$\begin{aligned} & \frac{(y_t^{SM})^2}{16\pi^2} V_{CKM}^{*t\alpha} V_{CKM}^{t\beta} \frac{a}{f_{PQ}} \bar{d}_\alpha [(im_d)_\alpha P_L - (im_d)_\beta P_R] d_\beta \\ & \times \left(\log \frac{\Lambda_{UV}^2}{m_{H^\pm}^2} \left(\frac{1}{2} C_Q - \frac{1}{2} C_u + \frac{1}{2} C_2 \right) \frac{1}{\sin^2 \beta} \right. \\ & \left. + \log \frac{m_{H^\pm}^2}{m_t^2} \left[\left(\frac{1}{2} C_Q - \frac{1}{2} C_u + \frac{1}{2} C_2 \right) + \frac{C_1 - C_2}{2} \cos^2 \beta \right] \right) \end{aligned}$$

SPECIFIC MODEL : ALP + 2HDM

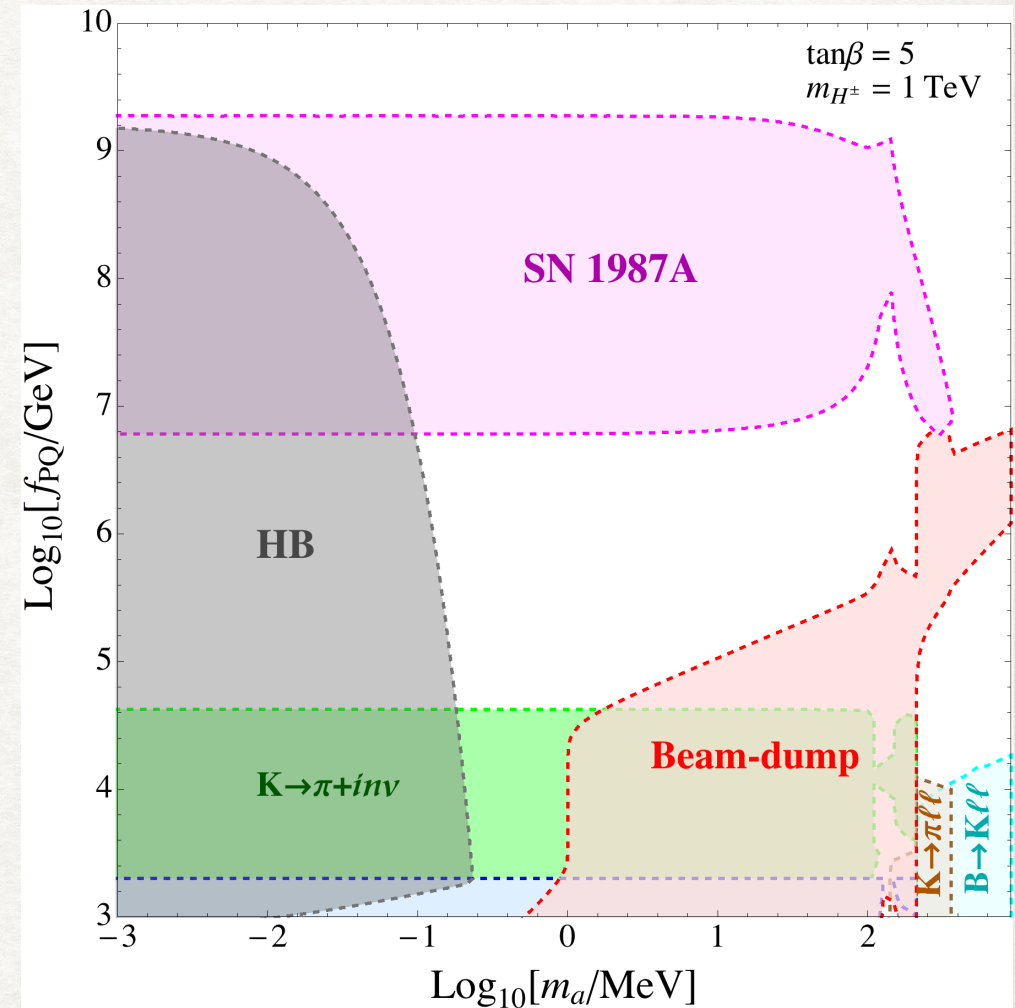
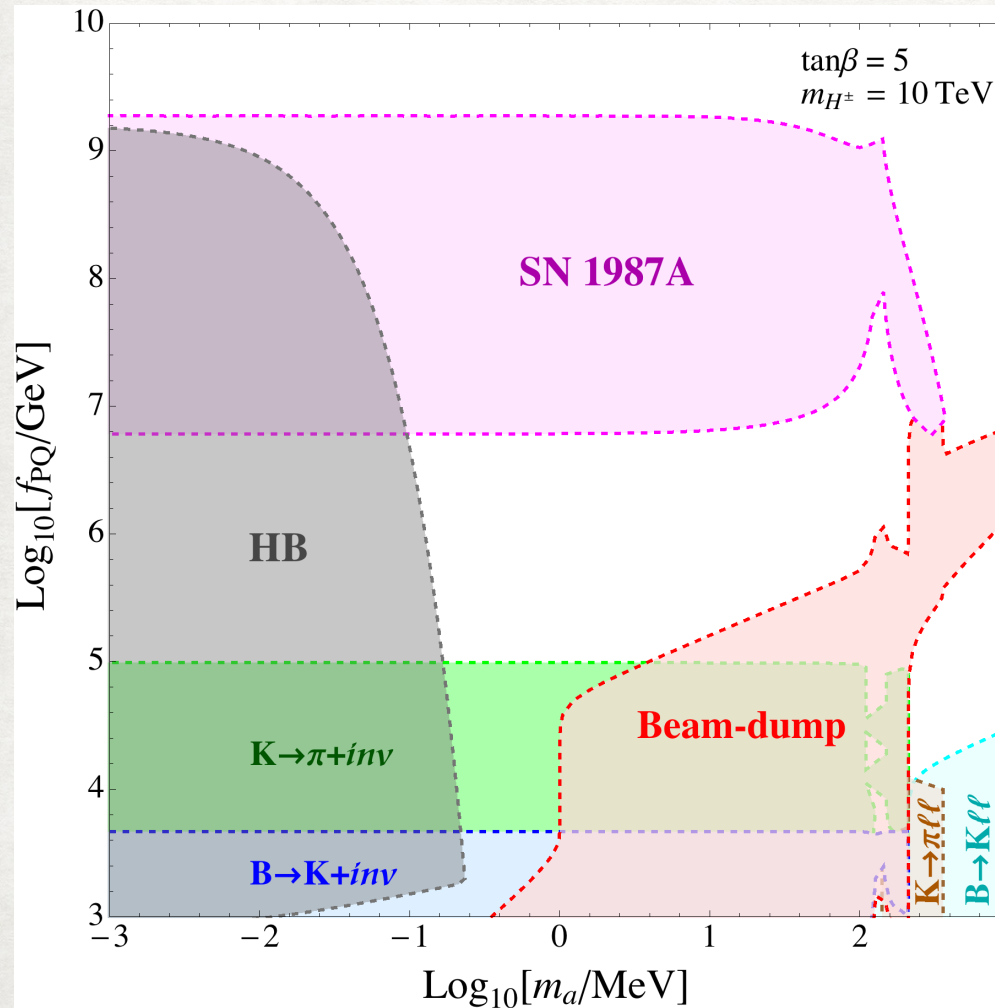
$$\begin{aligned} & \frac{(y_t^{SM})^2}{16\pi^2} V_{CKM}^{*t\alpha} V_{CKM}^{t\beta} \frac{a}{f_{PQ}} \bar{d}_\alpha [(im_d)_\alpha P_L - (im_d)_\beta P_R] d_\beta \\ & \times \left(\log \frac{\Lambda_{UV}^2}{m_{H^\pm}^2} \left(\frac{1}{2} C_Q - \frac{1}{2} C_u + \frac{1}{2} C_2 \right) \frac{1}{\sin^2 \beta} \right. \\ & \left. + \log \frac{m_{H^\pm}^2}{m_t^2} \left[\left(\frac{1}{2} C_Q - \frac{1}{2} C_u + \frac{1}{2} C_2 \right) + \frac{C_1 - C_2}{2} \cos^2 \beta \right] \right) \end{aligned}$$

$SM \ H^\pm$

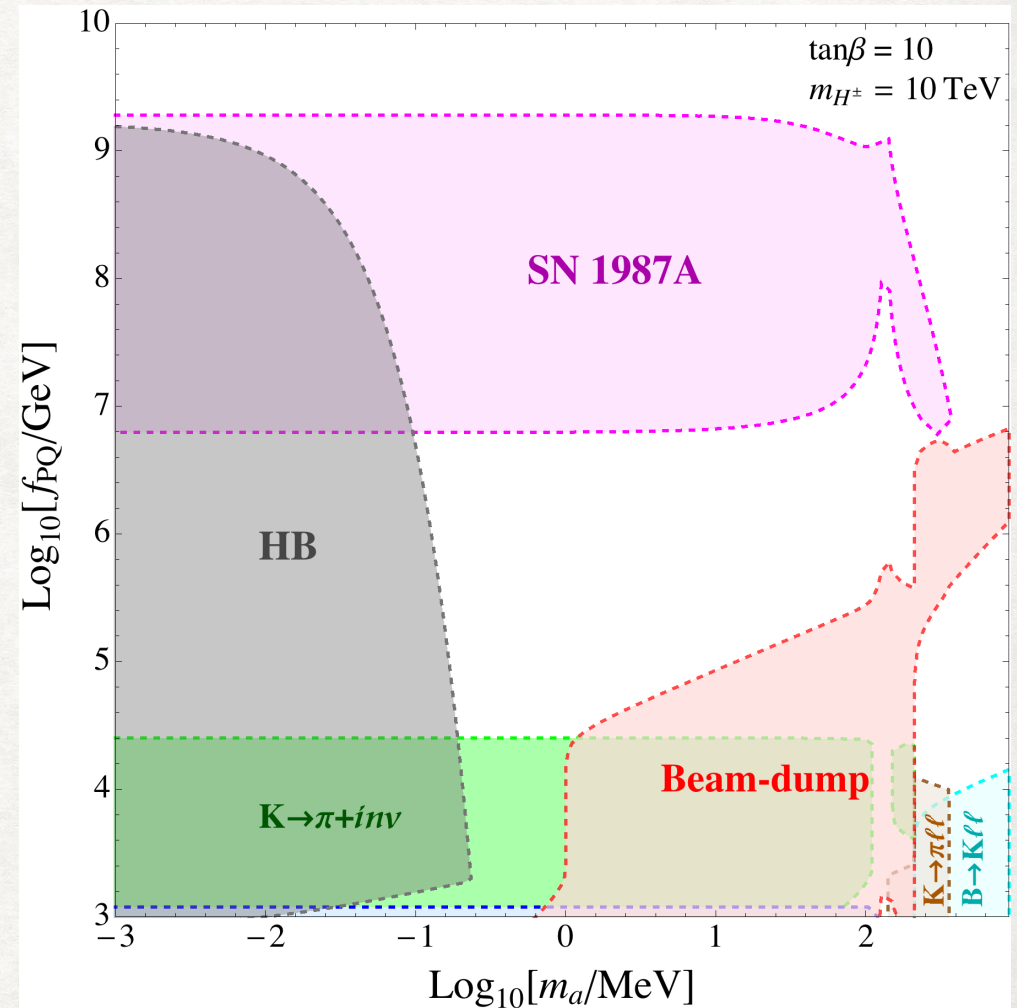
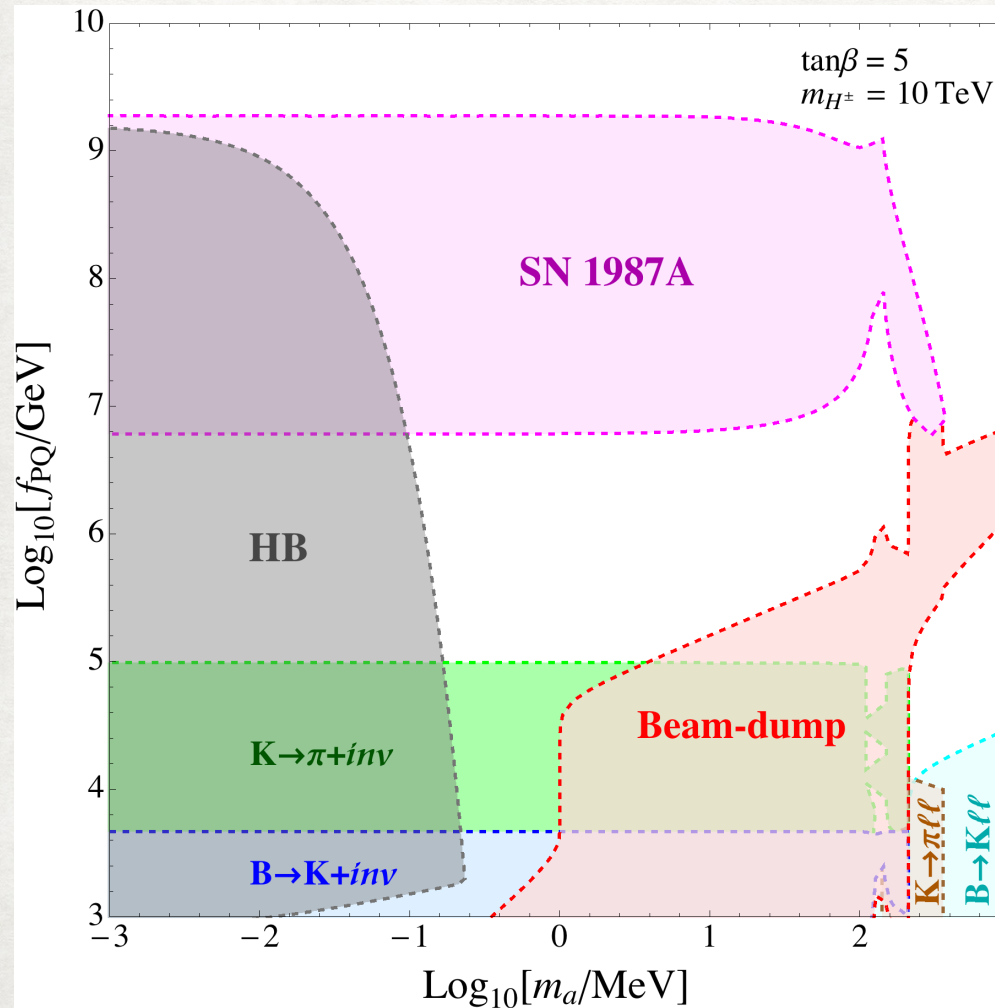
ALP IN TYPE II 2HDM WITH SINGLET

$$\begin{aligned}
 & \frac{(y_t^{SM})^2}{16\pi^2} V_{CKM}^{*t\alpha} V_{CKM}^{t\beta} \frac{a}{f_{PQ}} \bar{d}_\alpha [(im_d)_\alpha P_L - (im_d)_\beta P_R] d_\beta \\
 & \times \left(\log \frac{\Lambda_{UV}^2}{m_{H^\pm}^2} \left(\frac{1}{2}C_Q - \frac{1}{2}C_u + \frac{1}{2}C_2 \right) \frac{1}{\sin^2 \beta} \right. \\
 & \left. + \log \frac{m_{H^\pm}^2}{m_t^2} \left[\left(\frac{1}{2}C_Q - \frac{1}{2}C_u + \frac{1}{2}C_2 \right) + \frac{C_1 - C_2}{2} \cos^2 \beta \right] \right)
 \end{aligned}$$

ALP IN TYPE II 2HDM WITH SINGLET



ALP IN TYPE II 2HDM WITH SINGLET



SUMMARY

- FCNC in an arbitrary coupling structure
- ALP + 2HDM (additional contribution from the charged Higgs boson)
- Update the low-scale effective Lagrangian (Chiral perturbation theory)
- Experimental bounds from FCNC