

Precise Higgs mass calculations in (non-)minimal supersymmetry at both high and low scales

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Refs.

- Athron, JhP, Stöckinger, Voigt, CPC190(2015)139
- Athron, JhP, Stöckinger, Voigt, NPPP273(2016)2424
- Athron, JhP, Stöckinger, Voigt, Steudtner, 1609.00371

Supersymmetry

- Symmetry interchanging a fermion and a boson
- Non-trivial extension of Poincaré group
- Phenomenological uses:
 - ▶ Alleviation of fine-tuning in Higgs mass
 - ▶ Lightest Higgs mass restricted to a limited range
 - ▶ Gauge coupling unification
 - ▶ Dark matter
- Potential dangers:
 - ▶ Proton decay
 - ▶ FCNC, CP-violation
 - ▶ Damage to BBN

Higgs mass in MSSM

Okada, Yamaguchi, Yanagida; Ellis, Ridolfi, Zwirner; Haber, Hempfling

$$(m_h^2)_0 = m_Z^2 \cos^2 2\beta$$

$$(\Delta m_h^2)_{1,\log} \simeq \frac{3g_2^2}{8\pi^2 m_W^2} m_t^4 \ln \left(\frac{M_{\text{SUSY}}^2}{m_t^2} \right)$$

$$(\Delta m_h^2)_{1,\text{mix}} \simeq \frac{g_2^2 N_c}{16\pi^2 m_W^2 M_{\text{SUSY}}^2} \left[2m_t^4 X_t^2 \left(1 - \frac{X_t^2}{12M_{\text{SUSY}}^2} \right) + 2m_b^4 X_b^2 \left(1 - \frac{X_b^2}{12M_{\text{SUSY}}^2} \right) \right]$$

$$X_t \equiv A_t - \mu \cot \beta, \quad X_b \equiv A_b - \mu \tan \beta$$

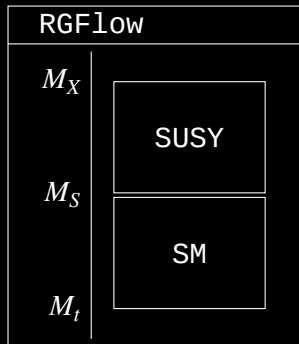
Can fit $m_h = 125$ GeV for $M_{\text{SUSY}} \lesssim 30$ TeV

Higgs mass in high-scale supersymmetry

Barbieri, Frigeni, Caravaglios, Okada, Yamaguchi, Yanagida, Haber, Hempfling,

Casas, Espinosa, Quirós, Riotto, Carena, Wagner, Lee

- Large scale difference between m_t and M_{SUSY} requires resummation of large logs
- Take “tower-of-EFTs” approach
- Match quartic SM Higgs coupling λ against 4-point function of lightest full-theory Higgs at M_{SUSY}
- Let λ run from M_{SUSY} to m_t
- Calculate m_h from $\lambda(m_t)$



Resummation of large logs

- Renormalization group equation

$$\mu \frac{\partial \lambda}{\partial \mu} = \frac{\partial \lambda}{\partial \ln \mu} = \beta(\lambda)$$

- 1-loop β -function

$$\beta(\lambda) = \frac{b}{(4\pi)^2} \lambda^2$$

- Solution of RGE

$$\begin{aligned} \overline{\lambda}(p) &= \frac{\lambda}{1 - b\lambda/(4\pi)^2 \ln(p/\mu)} \\ &= \lambda + b \frac{\lambda^2}{(4\pi)^2} \ln \frac{p}{\mu} + b^2 \frac{\lambda^3}{(4\pi)^4} \ln^2 \frac{p}{\mu} + \dots \end{aligned}$$

Public spectrum calculators

Model	Spectrum calculators
MSSM	ISASUSY, Softsusy, SPheno, SuSeFlav, SuSpect
NMSSM	NMSPEC, Softsusy
User-defined	SARAH Staub , FlexibleSUSY

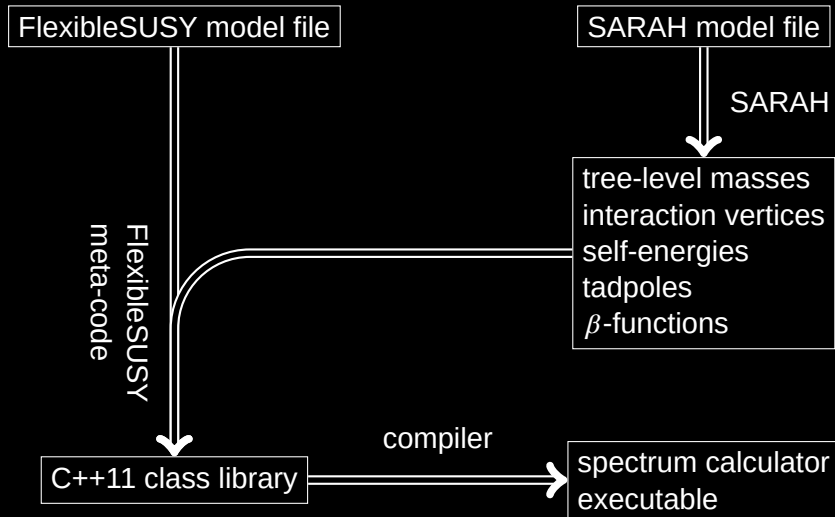
	SARAH	FlexibleSUSY
generates code in	FORTRAN	C++11
reuses code from	SPheno	Softsusy
symbolic calculation of β -functions etc.	done	invokes SARAH

FlexibleSUSY: a modern technology

flexiblesusy.hepforge.org

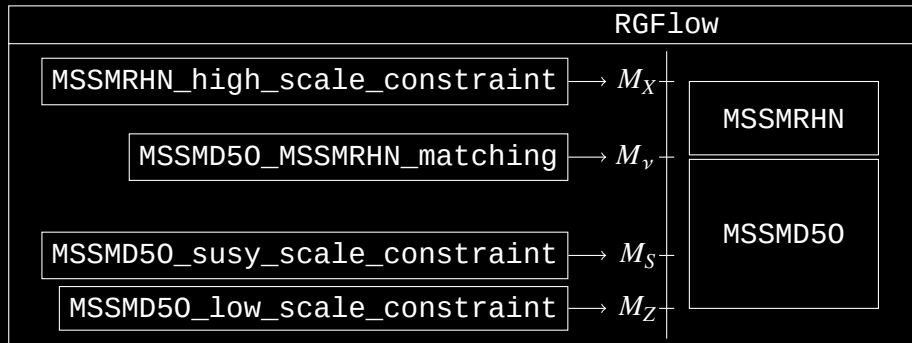


How FlexibleSUSY works



Tower of effective field theories

- Typical example: LFV from SUSY see-saw



- Other configurations of interest:
 - GUT [SU(5), SO(10), E_6] $\rightarrow$ MSSM
 - Split-SUSY, high-scale SUSY

Specimen SARAH model file for MSSM

```
(* Vector superfields *)
```

```
Gauge[[1]]={B, U[1], hypercharge, g1, False, RpM};
```

```
Gauge[[2]]={WB, SU[2], left, g2, True, RpM};
```

```
Gauge[[3]]={G, SU[3], color, g3, False, RpM};
```

```
(* Chiral superfields *)
```

```
SuperFields[[1]]={q, 3, {uL, dL}, 1/6, 2, 3, RpM};
```

```
SuperFields[[2]]={l, 3, {vL, eL}, -1/2, 2, 1, RpM};
```

```
SuperFields[[3]]={Hd, 1, {Hd0, HdM}, -1/2, 2, 1, RpP};
```

```
SuperFields[[4]]={Hu, 1, {Hup, Hu0}, 1/2, 2, 1, RpP};
```

```
SuperFields[[5]]={d, 3, conj[dR], 1/3, 1, -3, RpM};
```

```
SuperFields[[6]]={u, 3, conj[uR], -2/3, 1, -3, RpM};
```

```
SuperFields[[7]]={e, 3, conj[eR], 1, 1, 1, RpM};
```

```
(* Superpotential *)
```

```
SuperPotential = Yu u.q.Hu - Yd d.q.Hd - Ye e.l.Hd  
+ \[Mu] Hu.Hd;
```

Specimen FlexibleSUSY model file for MSSM

```
FSDefaultSARAHModel = MSSM;  
OnlyLowEnergyFlexibleSUSY = True;  
FlexibleEFTHiggs = True;
```

```
SUSYScaleInput = {  
  {vu, Sqrt[vu^2 + vd^2] Sin[ArcTan[TanBeta]]},  
  {vd, Sqrt[vu^2 + vd^2] Cos[ArcTan[TanBeta]]},  
  {MassB , M1Input},  
  {MassWB, M2Input},  
  {MassG , M3Input},  
  {mq2, mq2Input}, {mu2, mu2Input}, {md2, md2Input},  
  {ml2, ml2Input}, {me2, me2Input},  
  {\[Mu], MuInput},  
  {B\[Mu]], mAInput^2/(TanBeta + 1/TanBeta)},  
  {T[Yu], AuInput Yu},  
  {T[Yd], AdInput Yd},  
  {T[Ye], AeInput Ye}};
```

Automated Higgs mass calculation to 2-loops

- Categories of multi-loop Higgs mass calculation:

- ① Fixed-order:

- ① Diagrammatic approach

Hempfling, Hoang, Heinemeyer, Hollik, Weiglein

- ② Effective potential

Brignole, Dedes, Degrassi, Slavich, Zwirner, Martin

- ② Resummation via EFT

Casas, Espinosa, Quirós, Riotto, Carena, Wagner

- SARAH performs ① ② for user-defined models

Goodsell, Nickel, Staub

- FlexibleEFTHiggs is an implementation of ② in FlexibleSUSY

Athron, JhP, Stöckinger, Voigt, Steudtner

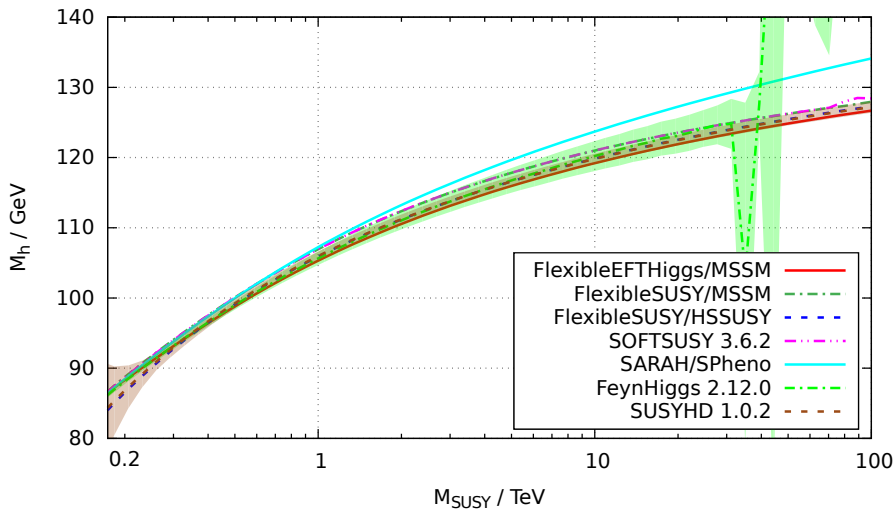
FlexibleEFTHiggs

- A variant of EFT approach
- Assume SM as EFT
- High-scale matching condition:

$$(M_h^{\text{SM}})^2 = (M_h^{\text{SUSY}})^2 \quad \text{instead of } \lambda\text{-matching}$$

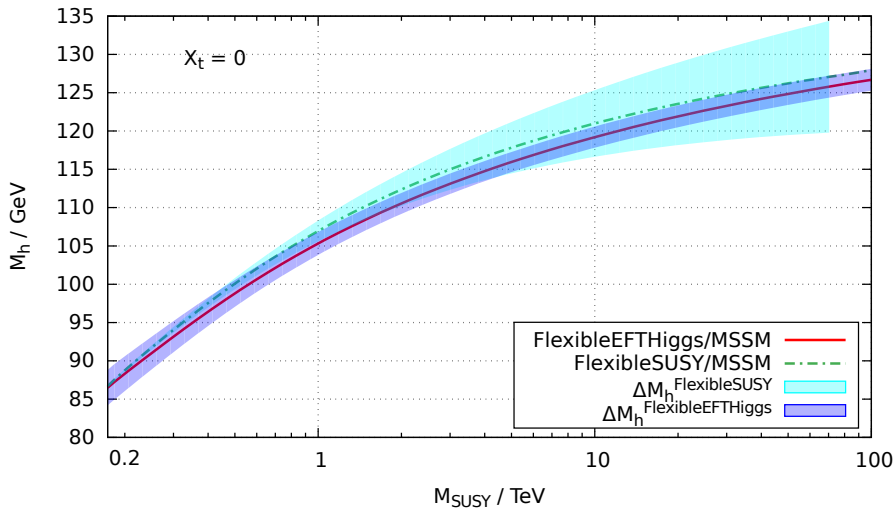
- Automatizable using only pole masses (available from SARAH)
- Not a full SMEFT, works only for m_h calculation
- Current accuracy: 1-loop matching + 3-loop running
- Work on 2-loop matching is underway

MSSM, varying M_{SUSY}



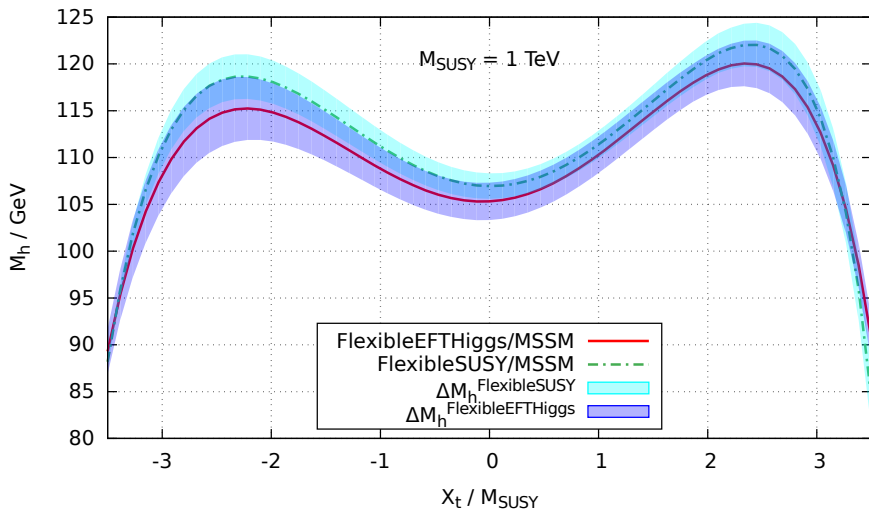
$$\tan\beta = 5, \quad X_t = 0$$

MSSM theory uncertainties, $X_t = 0$



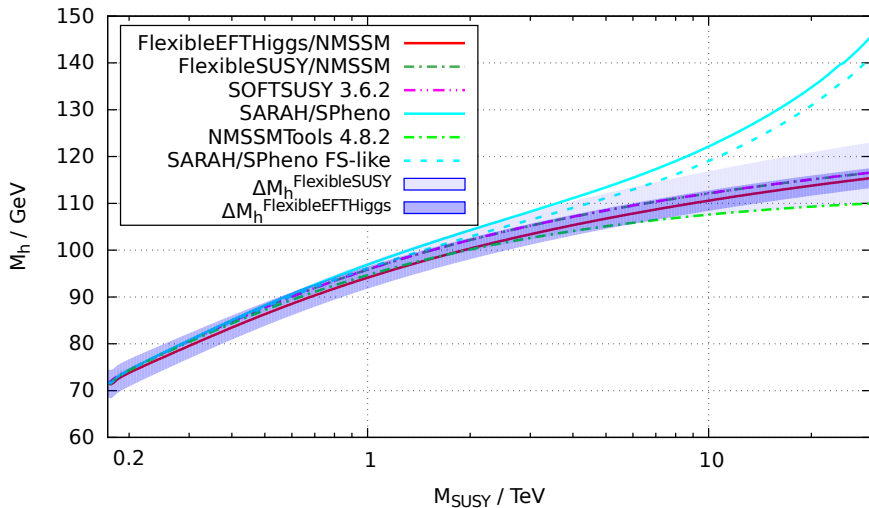
$$\tan\beta = 5$$

MSSM theory uncertainties, $X_t \neq 0$



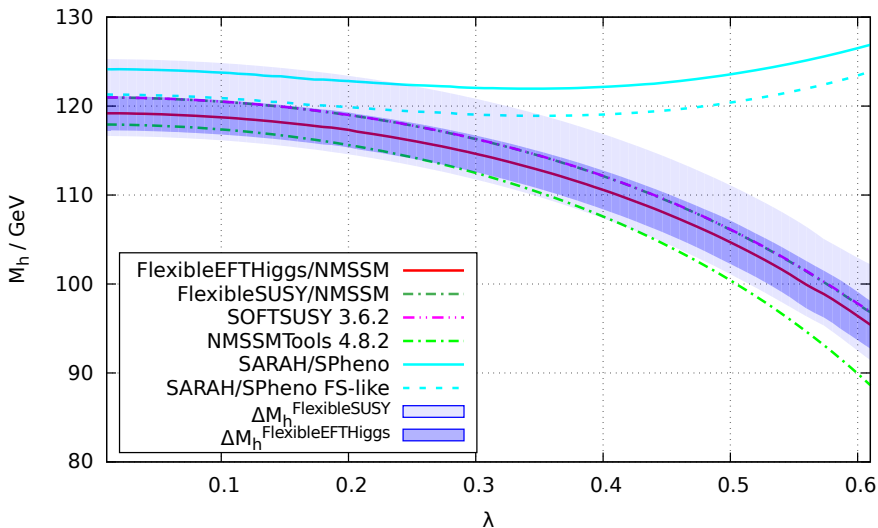
$$\tan\beta = 5$$

NMSSM, varying M_{SUSY}

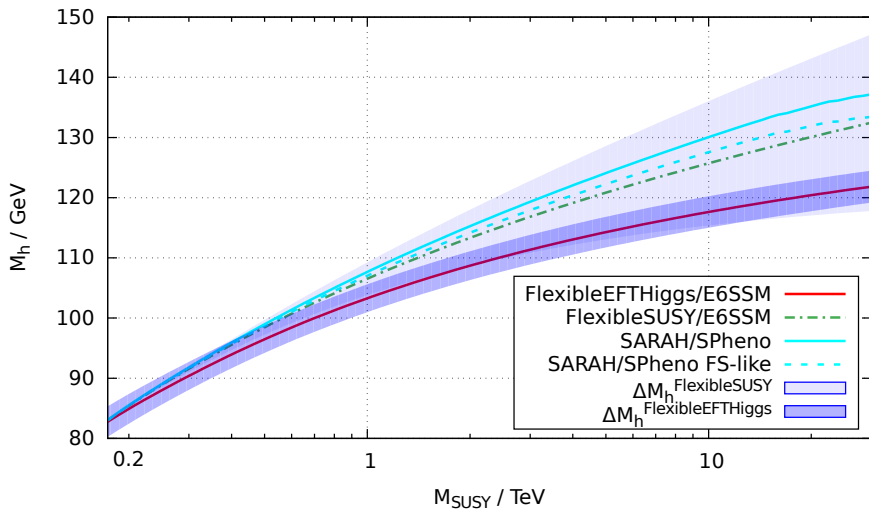


$$\tan\beta = 5, \quad X_t = 0, \quad \lambda = \kappa = 0.4$$

NMSSM, varying $\kappa = \lambda$

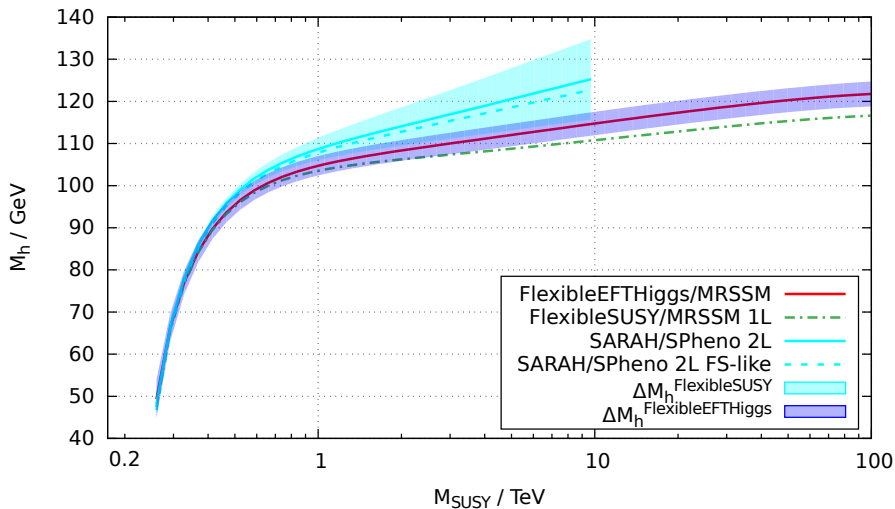


$$\tan\beta = 5, \quad X_t = 0, \quad M_{\text{SUSY}} = 10 \text{ TeV}$$



$$\tan\beta = 5, \quad X_t = 0, \quad \lambda = 0.1$$

MRSSM



Summary

- High new physics scale calls for resummation in m_h calculation
- FlexibleEFTHiggs is a new FlexibleSUSY feature performing resummed m_h calculation in user-defined supersymmetric models
- FlexibleEFTHiggs agrees with SUSYHD in MSSM
- Also studied instances are NMSSM, E_6 SSM, MRSSM
- FlexibleEFTHiggs' theory uncertainty is bounded at all scales in contrast to fixed-order calculations