

Decoherence of Primordial Perturbations in the View of a Local Observer

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Institute of Science Tokyo / IBS CTPU-CGA

Seminar talk

August 21, 2025 at **SI2025**

Based on

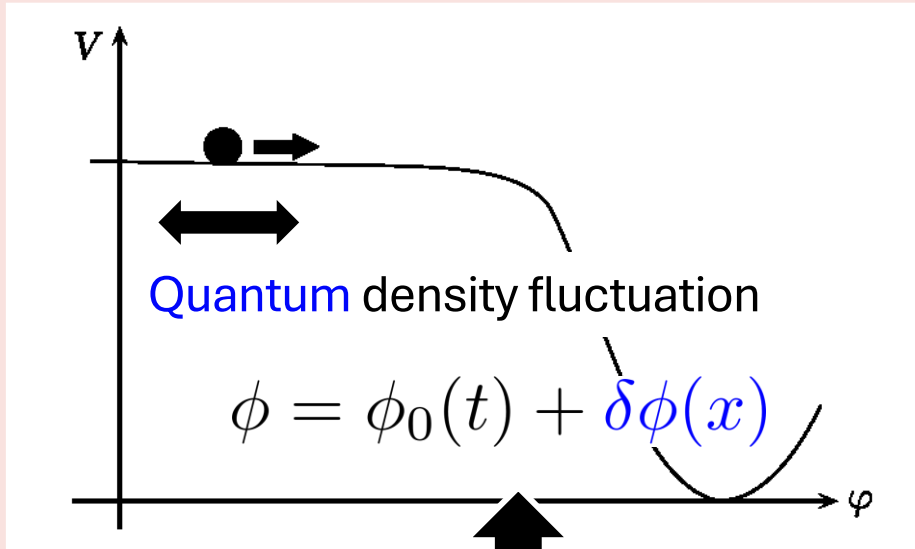
2504.10472 with **Junsei Tokuda** (McGill University)



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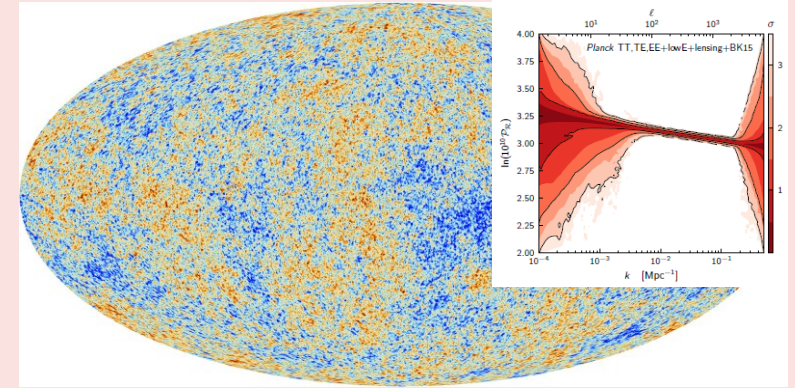
Inflation as a Source for Cosmological Perturbations



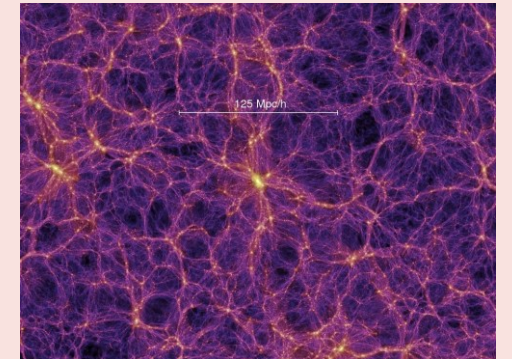
Quantum curvature perturbation

$$h_{ij} = (e^{\zeta(x)} a(t))^2 (\delta_{ij} + \gamma_{ij})$$

How (fast)
classicalized?



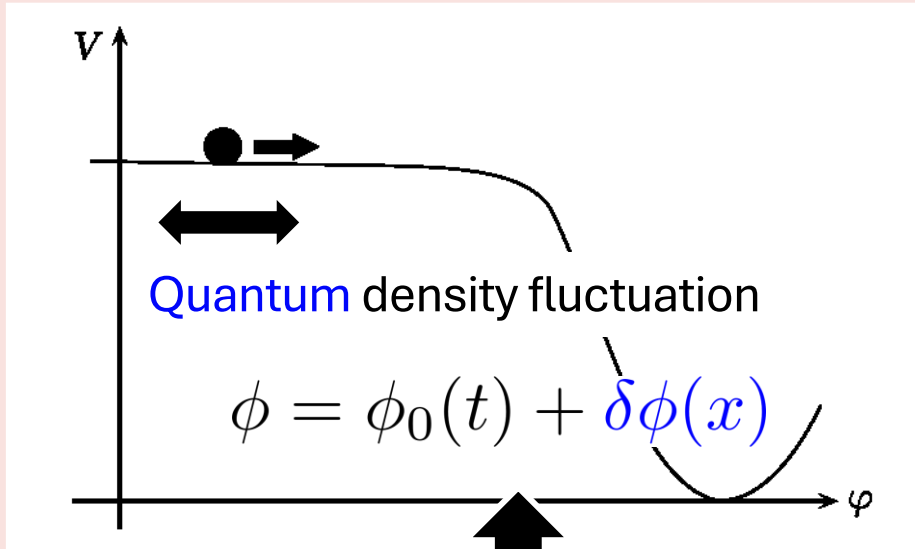
[Planck 1807.06211]



[Millennium Simulation 2005]

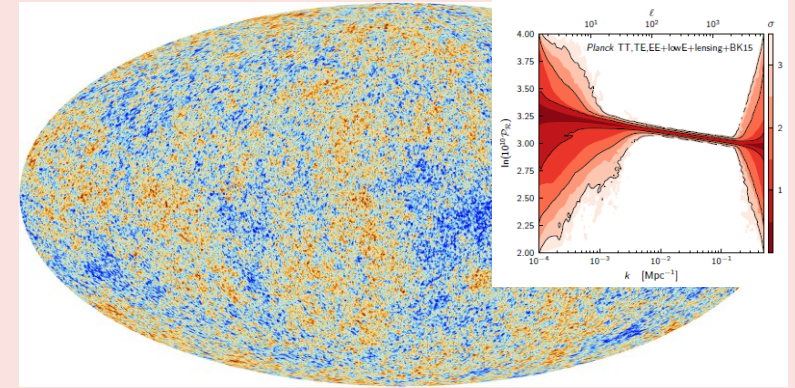
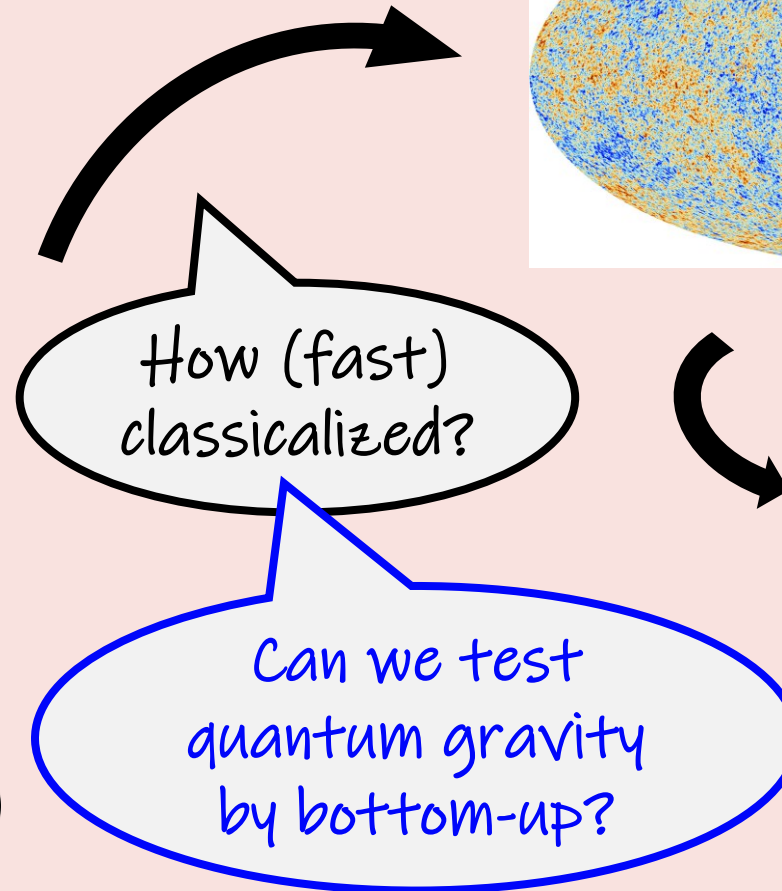
Classical anisotropy
and inhomogeneity

Inflation as a Source for Cosmological Perturbations

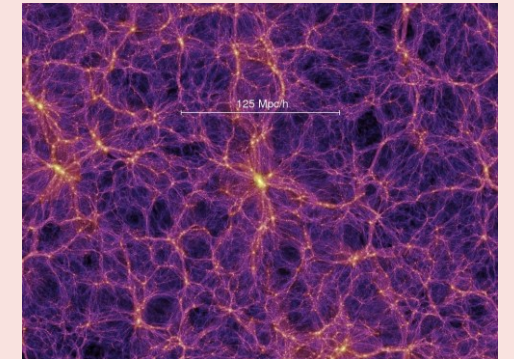


Quantum curvature perturbation

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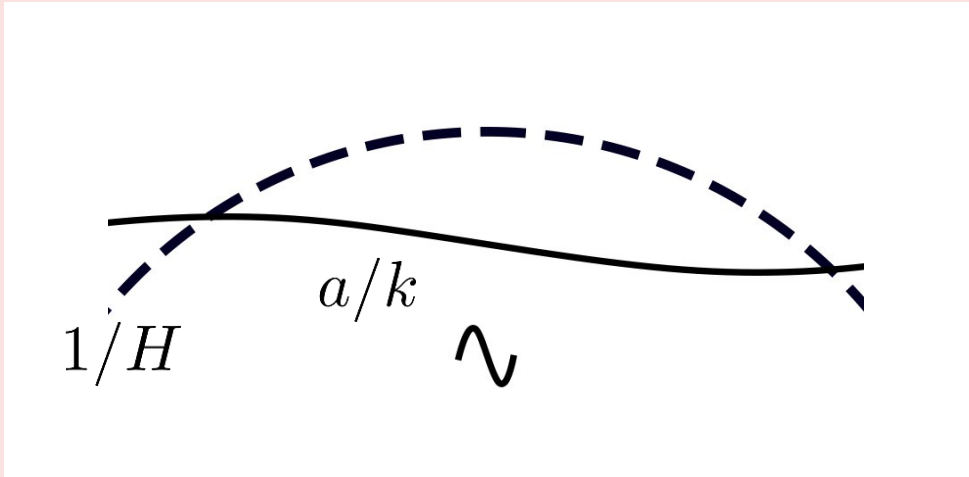
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Classical anisotropy
and inhomogeneity

“Quantumness” and “Classicalization”



□ Intuitively...

Large scale \longrightarrow Classical
 $a/k \gg 1/H$

Formally?

□ Coherence,

Entanglement,

Uncertainty, ...

$$\hat{\rho}[\zeta, \tilde{\zeta}] \text{ vs. } P(\zeta)$$

$$|\Psi\rangle_{AB} \neq |\phi\rangle_A \otimes |\varphi\rangle_B$$

$$\Delta\zeta \Delta\pi \gtrsim \hbar$$

$$\Leftrightarrow [\zeta_{\mathbf{k}}, \pi_{\mathbf{k}'}] = i\hbar(2\pi)^3 \delta^3(\mathbf{k} - \mathbf{k}')$$

- ✓ Quantum vs. classical dist.
 [Martin and Vennin 1801.09949, 1805.05609,
 Green and Porto 2001.09149, etc.]

- ✓ Bell test
 [Martin and Vennin 1706.04516, 2203.03505 etc.
 Sou et al. 2405.07141]

- ✓ Gaussian \longrightarrow minimal uncertainty
 Two mode squeezed state
 [Polarski and Starobinsky gr-qc/9504030]

- ✓ Stochastic formalism, PBH
 [Weenink and Prokopec 1108.3994]

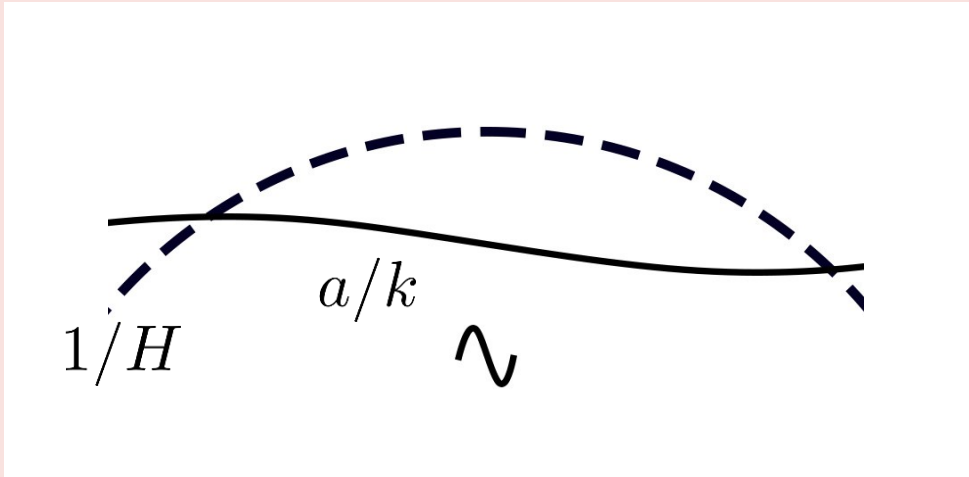
$$\mathcal{H}_{\text{tot}} = \bigotimes_i \mathcal{H}_i$$

$i \longleftarrow k? \ x? \text{ fields? } e^{ikx} \text{ vs } Y_{lm}?$

Quantumness can be sensitive to the system.

$$|\Psi\rangle = \prod_{\mathbf{k}} \left(\sum_n \alpha_{n,\mathbf{k}} |n_{\mathbf{k}}\rangle \otimes |n_{-\mathbf{k}}\rangle \right)$$

“Quantumness” and “Classicalization”



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□ **Coherence,** This talk

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$$\Delta\zeta \Delta\pi \gtrsim \hbar$$

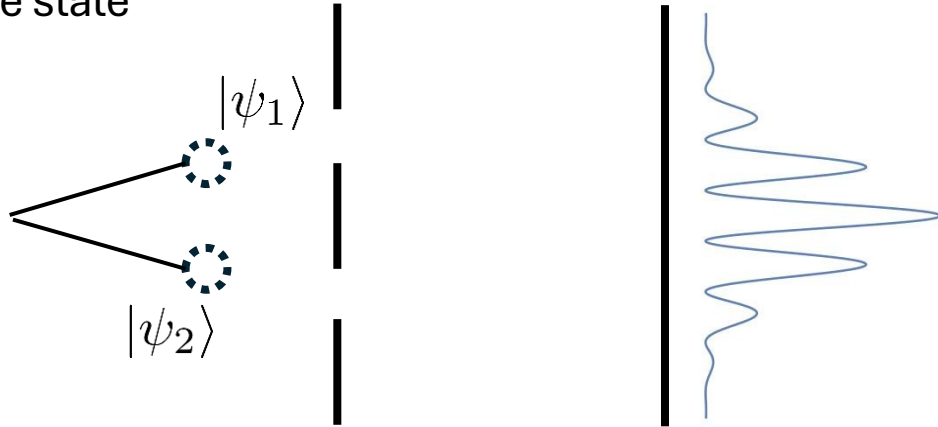
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Quantum Interference and Decoherence

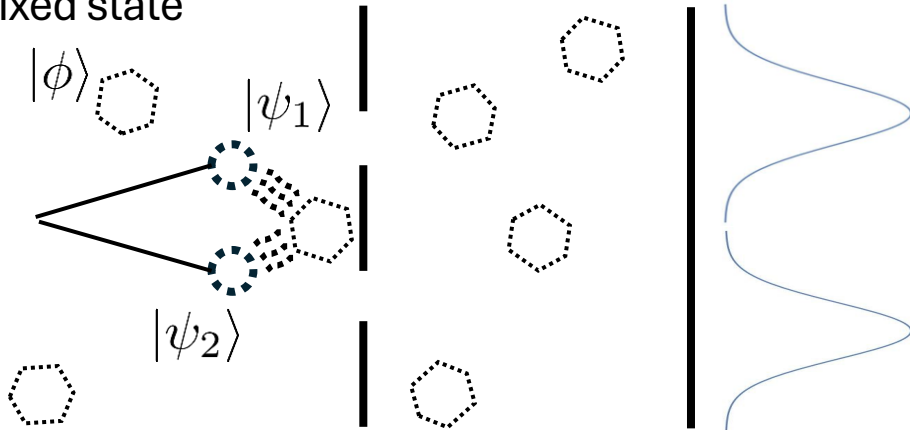
Pure state



$$|\Psi\rangle = \alpha |\psi_1\rangle + \beta |\psi_2\rangle$$

$$\langle\Psi|\hat{A}|\Psi\rangle = |\alpha|^2 \langle\psi_1|\hat{A}|\psi_1\rangle + |\beta|^2 \langle\psi_2|\hat{A}|\psi_2\rangle + (\alpha\beta^* \langle\psi_2|\hat{A}|\psi_1\rangle + \text{c.c.})$$

Mixed state



$$|\Psi\rangle = \alpha |\psi_1\rangle |\phi_1\rangle + \beta |\psi_2\rangle |\phi_2\rangle$$

$$\rho_\psi = \text{Tr}_\phi[|\Psi\rangle \langle\Psi|] = \begin{pmatrix} |\alpha|^2 & \alpha\beta^* \langle\phi_2|\phi_1\rangle \\ \alpha^*\beta \langle\phi_1|\phi_2\rangle & |\beta|^2 \end{pmatrix}$$

$\langle\phi_2|\phi_1\rangle \sim 0$ if scattered to independent states.

More scattering, more independent, less interference.

✓ Measure of coherence: $\rho_{ij} \Big|_{i \neq j}$

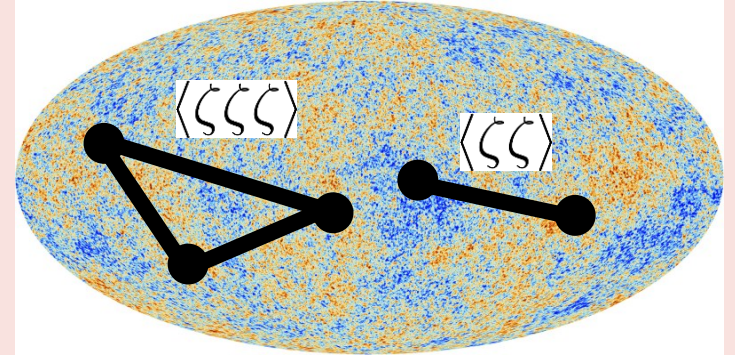
* Representation independent measure of coherence: Rényi entropy, purity, quantum discord, etc.
[Streltsov et al. 1612.07570, Henderson and Vedral quant-ph/0105028, etc. Comparison: Martin et al. 2211.10114]

Wavefunction Formalism

□ Observables: correlation functions

$$\langle \Omega | \hat{\zeta}^n(t) | \Omega \rangle = \int \mathcal{D}\zeta(t) \underbrace{\langle \Omega | \zeta; t \rangle \langle \zeta; t | \Omega \rangle}_{\hat{\zeta}(t) | \zeta; t \rangle = \zeta(t) | \zeta; t \rangle} \zeta^n \equiv \int \mathcal{D}\zeta(t) |\Psi[\zeta(t)]|^2 \zeta^n$$

✓ E.g., $\mathcal{H} = \mathcal{H}_{\mathbf{k}_S} \otimes \mathcal{H}_{\mathbf{k}_E}$ with $k_S \in \{k_{\text{CMB}}\}$



□ Wavefunction at a certain time slice

$$\underbrace{\Psi[\zeta(t)] \equiv \langle \zeta; t | \Omega \rangle}_{(\approx e^{iS_{\text{cl}}[\zeta]})} = \exp \left[\underbrace{-\frac{1}{2} \int \frac{d^3 \mathbf{k}_1}{(2\pi)^3} \frac{d^3 \mathbf{k}_2}{(2\pi)^3} \psi_2 \zeta_{k_1} \zeta_{k_2}}_{\text{Gaussian}} - \overbrace{\frac{1}{3!} \int \frac{d^3 \mathbf{k}_1}{(2\pi)^3} \frac{d^3 \mathbf{k}_2}{(2\pi)^3} \frac{d^3 \mathbf{k}_3}{(2\pi)^3} \psi_3 \zeta_{k_1} \zeta_{k_2} \zeta_{k_3} - \dots}^{\text{Gravitational non-linearity}} \right]$$

ψ_n : coefficient of the expansion

✓ Free propagation: $e^{-\int \mathbf{k} \psi_2 \zeta_{\mathbf{k}} \zeta_{-\mathbf{k}}}$ \longrightarrow no entanglement between \mathbf{k}_S and \mathbf{k}_E (no scattering)

\longrightarrow Non-linearities cause decoherence

Formulation of Decoherence

[Nelson 1601.03734, Sou et al. 2207.04435]

Another method: quantum master equation [Burgess et al. astro-ph/061646, Burgess et al. 2211.11046, etc.]

□ Density matrix

$$\begin{aligned}\rho_S[\zeta_S, \tilde{\zeta}_S] &= \int \mathcal{D}\zeta_E(t) \Psi[\zeta_S, \zeta_E] \Psi^*[\tilde{\zeta}_S, \zeta_E] \\ &\simeq \Psi_G[\zeta_S] \Psi_G^*[\tilde{\zeta}_S] \int \mathcal{D}\zeta_E |\Psi_G[\zeta_E]|^2 e^{-\frac{1}{6} \int_{\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3} (\zeta_1 \zeta_2 \zeta_3 \psi_3 + \tilde{\zeta}_1 \tilde{\zeta}_2 \tilde{\zeta}_3 \psi_3^*)} \\ &\equiv \rho_{\text{diag}} \times \exp \left[- \int_{\mathbf{k}_S} \Gamma \Delta \zeta_{\mathbf{k}_S}^2 \right] \\ &\equiv (\zeta - \tilde{\zeta})^2\end{aligned}$$

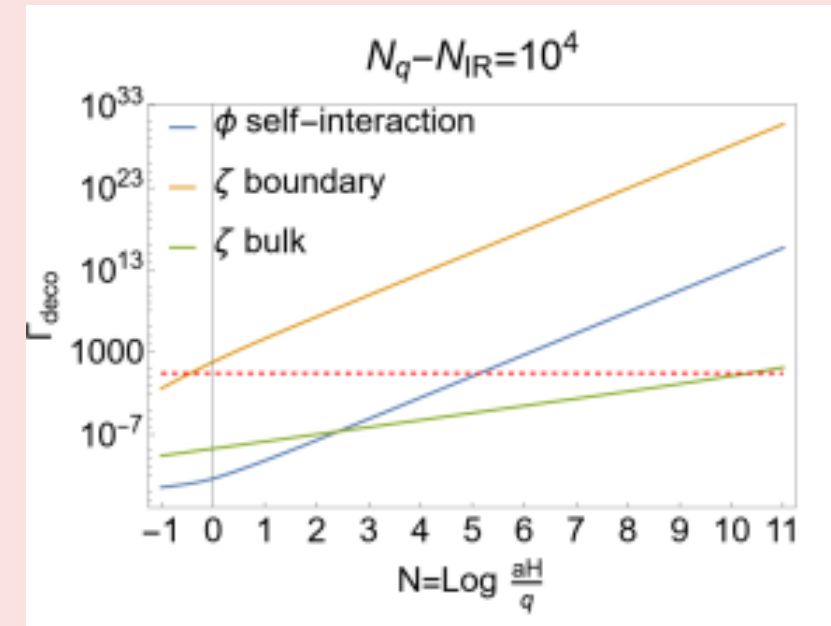
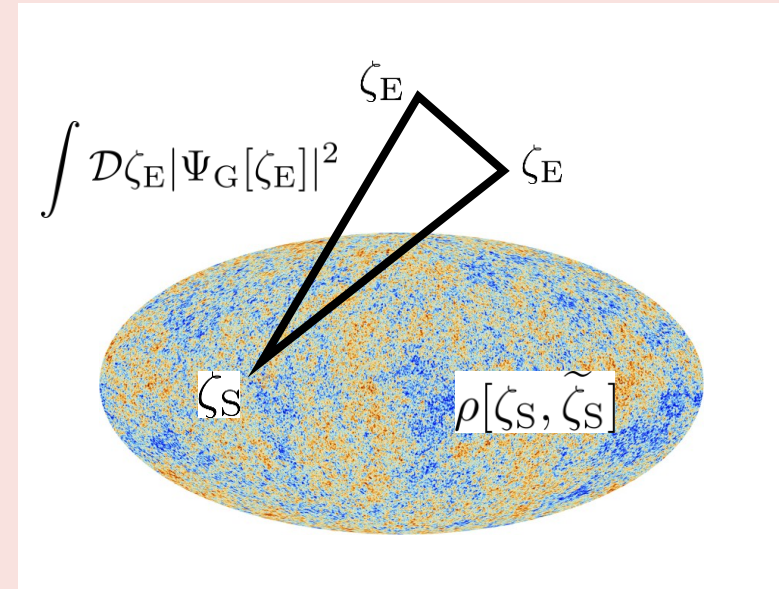
□ Decoherence rate

(decay rate of off-diagonal component)

$$\Gamma \approx \underbrace{\frac{\psi_3}{\Delta \zeta_S} \left(\begin{array}{c} \zeta_E \\ \zeta_E \end{array} \right) \frac{\psi_3}{\Delta \zeta_S}}_{\text{Loop at the time slice}} \sim \underbrace{\frac{1}{\epsilon^2} \left(\frac{aH}{k_S} \right)^6}_{\text{Boundary term}} + \underbrace{\epsilon^2 \left(\frac{aH}{k_S} \right)^3}_{\text{Bulk term}}$$

WITH IR divergence and UV divergence $\Gamma \supset \log \frac{k_S}{k_{\text{IR}}}, \Lambda_{\text{UV}}^\#$

Some cancellations? Regularization?



[Sou et al. 2207.04435]

Consistency condition for loop calculations

ψ_3 ←

[Nelson 1601.03734] [Sou et al. 2207.04435]

$$S_3 = \int dt d^3x \left\{ a^3 \epsilon^2 \dot{\zeta}^2 + a \epsilon^2 \zeta (\partial \zeta)^2 - 2a \epsilon \dot{\zeta} \partial \zeta \partial \chi + 2f(\zeta) \frac{\delta \mathcal{L}}{\delta \zeta} \Big|_1 + \mathcal{L}_b \right\}, \quad \partial^2 \chi \equiv a^2 \epsilon \dot{\zeta}$$

$$\mathcal{L}_b = \partial_t \left[-9a^3 H \dot{\zeta}^3 + \frac{a}{H} \zeta (\partial \zeta)^2 - \frac{1}{4aH^3} (\partial \zeta)^2 \partial^2 \zeta - \frac{a\epsilon}{H} \zeta (\partial \zeta)^2 - \frac{\epsilon a^3}{H} \dot{\zeta}^2 + \frac{1}{2aH^2} \zeta (\partial_i \partial_j \zeta \partial_i \partial_j \chi - \partial^2 \zeta \partial^2 \chi) - \frac{\eta a}{2} \zeta^2 \partial^2 \chi - \frac{1}{2aH} \zeta (\partial_i \partial_j \chi \partial_i \partial_j \chi - \partial^2 \chi \partial^2 \chi) \right]$$



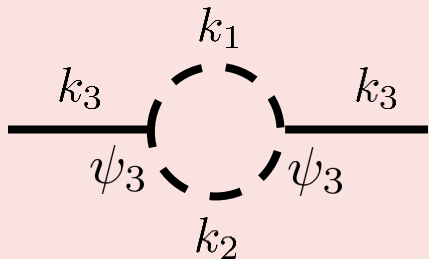
Necessary for correlation function

□ Maldacena's consistency condition for wavefunction [Pimentel 1309.1793]

$$\lim_{k_1 \rightarrow 0} \psi_3(k_1, k_3) = \left(3 - k_3 \frac{d}{dk_3} \right) \psi_2(k_3)$$

Cf. $\left\{ \begin{array}{l} \lim_{k_1 \rightarrow 0} \langle \zeta_1 \zeta_2 \zeta_3 \rangle = - \langle \zeta_1 \zeta_1 \rangle \left(3 + k_3 \frac{d}{dk_3} \right) \langle \zeta_3 \zeta_3 \rangle \\ \text{[Maldacena astro-ph/0210603]} \\ \langle \zeta_1 \zeta_2 \rangle = \frac{1}{2 \text{Re}[\psi_2(k_1)]}, \quad \langle \zeta_1 \zeta_2 \zeta_3 \rangle = - \frac{2 \text{Re}[\psi_3]}{\prod_{i=1}^3 2 \text{Re}[\psi_2(k_i)]} \end{array} \right.$

□ Loop diagram at the time slice



IR: $k_1 \ll k_2 \simeq k_3 \ll aH \longrightarrow \log k_1$ from $\int \langle \zeta_1 \zeta_1 \rangle k_1^2 dk_1$

UV: $k_1 \simeq k_2 \gg aH \gg k_3 \longrightarrow k_1^5$ from $\partial_t(a\zeta(\partial_i \zeta)^2/H)$

False decoherence during inflation

[Sano and Tokuda 2504.10472]

□ **IR and UV divergences** $(\rho_{\text{off-diag}} \sim e^{-\Gamma})$

$$\Gamma \sim \left[\underbrace{\frac{1}{\epsilon^2} \left(\frac{aH}{k_S} \right)^6}_{\text{IR cutoff}} + \underbrace{\epsilon^2 \left(\frac{aH}{k_S} \right)^3}_{\text{UV cutoff}} \right] (1 + \log(k_{\text{IR}}/k_S)) + \frac{1}{\epsilon^2} \left(\frac{\Lambda}{aH} \right)^5$$

IR div.: Affected by very beginning of inflation?

UV div.: Divergent offset to decoherence exists. Never quantum?

✓ **Proper observables should be insensitive to deep IR and deep UV contributions.**
(e.g., adiabaticity: rapid modes decouple to slow modes. [Unruh 1110.2199 for QFT])

IR: **local observer's** coordinate

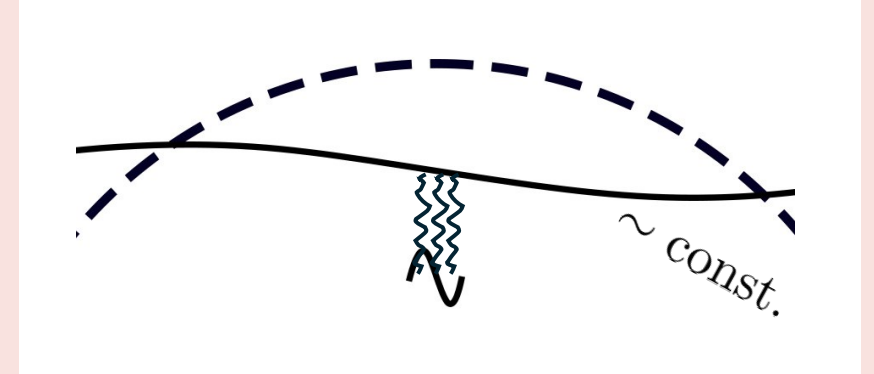
UV: **time averaged observables** as well as renormalization

Local Observer Effect in Correlation Function

[Tanaka and Urakawa 1103.1251, Pajer et al. 1305.0824]

$$\langle \zeta(x_1) \zeta(x_2) \zeta(x_3) \rangle \supset \int_{k_1 \ll k_3} \frac{k_1^2 dk_1}{k_1^3} \frac{k_3^2 dk_3}{k_3^3} \sim \log k_1 \Big|_{k_1 \rightarrow 0}$$

Short modes strongly correlates with constant long modes (?)



□ **Conformal free-falling observer** $\mathbf{x}_F \simeq (1 + \zeta_L) \mathbf{x}$, $ds^2 = a^2(-d\tau^2 + d\mathbf{x}_F^2) + \dots$
(Conformal Fermi normal coordinate)

$$\Rightarrow \zeta_{F,\mathbf{k}} \simeq \zeta_{\mathbf{k}} + \zeta_L(3 + k\partial_k)\zeta_{\mathbf{k}}$$

$$\Rightarrow \lim_{k_1 \rightarrow 0} \langle \zeta_1 \zeta_2 \zeta_3 \rangle_F = \lim_{k_1 \rightarrow 0} \langle \zeta_1 \zeta_2 \zeta_3 \rangle + \langle \zeta_1 \zeta_1 \rangle \left(3 + k_3 \frac{d}{dk_3} \right) \langle \zeta_3 \zeta_3 \rangle = \underline{0} \quad \text{IR correlations are turned off}$$

$$\Rightarrow \langle \zeta(x_1) \zeta(x_2) \zeta(x_3) \rangle_F \xrightarrow{\text{IR}} \int_{k_1 \ll k_3} \frac{k_1}{k_3^3} dk_1 dk_3 \quad \text{Finite result}$$

Local Observer Effect in Wavefunction Formalism

[Sano and Tokuda 2504.10472]

□ Wavefunction in free-falling coordinate

$$\Psi[\zeta] = \exp \left[-\frac{1}{2} \int_{\mathbf{k}_1, \mathbf{k}_2} \psi_2 \zeta_{k_1} \zeta_{k_2} - \frac{1}{3!} \int_{\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3} \psi_3 \zeta_{k_1} \zeta_{k_2} \zeta_{k_3} - \cdots \right]$$

$\zeta_{\mathbf{k}} \simeq \zeta_{\mathbf{F}, \mathbf{k}} - \zeta_{\mathbf{L}}(3 + k\partial_k)\zeta_{\mathbf{F}, \mathbf{k}}$
Changing the expansion basis

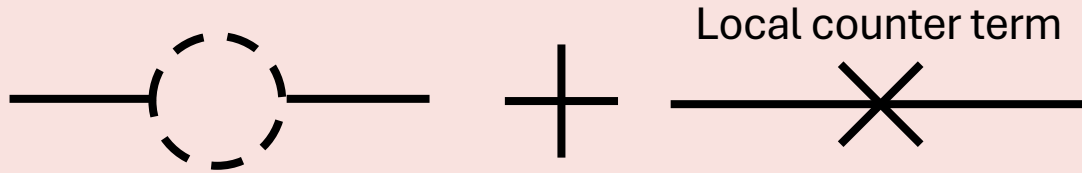
$$= \Psi_{\mathbf{F}}[\zeta_{\mathbf{F}}] = \exp \left[-\frac{1}{2} \int_{\mathbf{k}_1, \mathbf{k}_2} \psi_{\mathbf{F}, 2} \zeta_{\mathbf{F}, k_1} \zeta_{\mathbf{F}, k_2} - \frac{1}{3!} \int_{\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3} \psi_{\mathbf{F}, 3} \zeta_{\mathbf{F}, k_1} \zeta_{\mathbf{F}, k_2} \zeta_{\mathbf{F}, k_3} - \cdots \right]$$

$$\Rightarrow \lim_{k_1 \rightarrow 0} \psi_{\mathbf{F}, 3} = \lim_{k_1 \rightarrow 0} \psi_3 - \left(3 - k_3 \frac{d}{dk_3} \right) \psi_2 = \underline{0}$$

$$\Rightarrow \Gamma_{\text{IR}} \sim \frac{\psi_3}{\Delta \zeta_{\text{S}}} \text{ (loop) } \frac{\psi_3}{\Delta \zeta_{\text{S}}} \sim \log(k_{\text{IR}}/k_{\text{S}}) \xrightarrow{(k_{\text{E}}/k_{\text{S}})^2 \text{ moderation for each } \psi_3} \Gamma_{\text{IR}, \mathbf{F}} \sim (k_{\text{IR}}/k_{\text{S}})^4 \sim 0$$

UV divergence in Equal Time

□ Scattering



$$S_{\text{ren}} = S_0 + S_{\text{CT}}$$

□ Equal time correlators

[Balasubramanian et al. 1108.3568 and Bucciotti 2410.01903 for [flat spacetime](#) examples etc.]

$$\langle \mathcal{O}_{1,\text{ren}}^{\mathbf{k}} \mathcal{O}_{2,\text{ren}}^{-\mathbf{k}}(t) \rangle \sim \int d(\mathbf{x}_1 - \mathbf{x}_2) \frac{1}{|\mathbf{x}_1 - \mathbf{x}_2|^{2\Delta}} \quad \text{diverges even after renormalization when } \Delta \geq \frac{3}{2}.$$

“Equal time” is beyond IR effective theory?

➡ Time averaged observables naturally shows decoupling of UV physics
(Observers do not have much time resolution to see the equal time correlators)

Time Averaged Observables

[Sano and Tokuda 2504.10472]

$$\langle \Phi_1 \Phi_2(\tau) \rangle \supset \tau \text{---} \Phi \text{---} \bigcirc \text{---} \Phi \text{---} \tau$$

$\Phi = \zeta \text{ or } \pi$
 τ : Conformal time

Time averaging

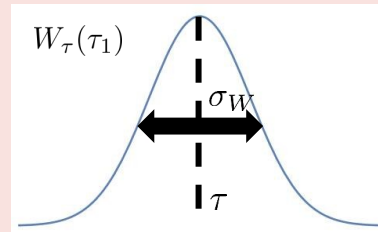
$$\int^\Lambda k^\# dk \longrightarrow \Lambda^\#$$

□ Time averaging

$$W_\tau(\tau_1) = \frac{e^{-(\tau_1 - \tau)^2 / 2\sigma_W^2}}{\sqrt{2\pi\sigma_W^2}},$$

$$G(k; \tau_1, \tau_2) \propto e^{-ik_1(\tau_1 - \tau_2)}$$

Green function



$$\begin{aligned} \langle \overline{\Phi}_1 \overline{\Phi}_2(\tau) \rangle &= \int d\tau_1 d\tau_2 W_\tau(\tau_1) W_\tau(\tau_2) \langle \Phi_1(\tau_1) \Phi_2(\tau_2) \rangle \\ &\supset \int d\tau_1 d\tau_2 W_\tau(\tau_1) W_\tau(\tau_2) \left[\tau_1 \text{---} \Phi \text{---} \bigcirc \text{---} \Phi \text{---} \tau_2 \right] \end{aligned}$$

$$\frac{1}{|\tau_1 - \tau_2|^\#}$$

From time-ordered loop contributions.
 This is (expected to be) renormalized.

$$\frac{e^{-ik(\tau_1 - \tau_2)}}{|\tau_1 - \tau_2|^\#}$$

From Wightman function.
 Not renormalized in standard procedure.

$$\Gamma_{UV} \sim \int_{k > aH} dk k^\# e^{-k^2 \sigma_W^2}$$

Exponential decay in sub-horizon

Summary: Genuine decoherence during inflation

False contributions $(\rho_{\text{off-diag}} \sim e^{-\Gamma})$

$$\Gamma_{\text{false}} \approx \text{Loop at the time slice} \sim (1 + \log(k_{\text{IR}}/k_{\text{S}})) \left[\frac{1}{\epsilon^2} \left(\frac{aH}{k_{\text{S}}} \right)^6 + \epsilon^2 \left(\frac{aH}{k_{\text{S}}} \right)^3 \right] + \frac{1}{\epsilon^2} \left(\frac{\Lambda}{aH} \right)^5$$

IR cutoff
UV cutoff

Long mode is absorbed in geodesic coordinate.

$$ds^2 = a^2(-d\tau^2 + e^{2\zeta} d\mathbf{x}^2) = a^2(-d\tau^2 + d\mathbf{x}_{\text{F}}^2) + \cdots$$

$$\lim_{k_1 \rightarrow 0} \psi_{\text{F},3} = \lim_{k_1 \rightarrow 0} \psi_3 - \left(3 - k_3 \frac{d}{dk_3} \right) \psi_2 = 0$$

✓ Leading scaling in the previous work is genuine

Classified to two components when averaging in time.

$$\frac{1}{|\tau_1 - \tau_2|^{\#}}$$

Renormalized

$$\frac{e^{-ik(\tau_1 - \tau_2)}}{|\tau_1 - \tau_2|^{\#}}$$

Averaged out

$$\Gamma_{\text{UV}} \sim \int_{k > aH} dk \, k^{\#} e^{-k^2 \sigma_W^2}$$

$$\Gamma_{\text{genuine}} \sim \frac{1}{\epsilon^2} \left(\frac{aH}{k_{\text{S}}} \right)^6 + \epsilon^2 \left(\frac{aH}{k_{\text{S}}} \right)^3$$

boundary term
bulk term

Outlook: Importance of late time evolutions

□ Boundary terms [Sano and Tokuda '25]

✓ During inflation

$$\Gamma_{\text{inf}} \sim \underbrace{\frac{1}{\epsilon^2} \left(\frac{aH}{k_S} \right)^6}_{\text{Boundary term}} + \underbrace{\epsilon^2 \left(\frac{aH}{k_S} \right)^3}_{\text{Bulk term}}$$

✓ Late time universe (but before re-entry)

$$\Gamma_{\text{rad. dom.}} \sim \underbrace{\frac{1}{\epsilon^2} \left(\frac{a_f H_f}{k_S} \right)^6 \left(\frac{a}{a_f} \right)^2}_{\text{Boundary term}} + \underbrace{\epsilon^2 \left(\frac{a_f H_f}{k_S} \right)^3 \left(\frac{a}{a_f} \right)^5}_{\text{Bulk term}}$$

□ Time averaging scale?

□ High-frequency gravitational wave [Takeda and Tanaka '25]

✓ GW with frequency $f_{\text{GW}} \gtrsim 100 \text{ Hz}$ (?) may be **quantum even today!**

* Estimation of thermal decoherence by a scalar field, keeping reheating in mind

□ Outlook

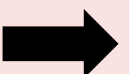
- ✓ **Systematic approaches to sub-horizon evolution** for a more realistic model?
- ✓ **Entanglement harvesting through detectors?** Graviton-photon conversion?
- ✓ **What is more than proving quantumness of gravity?** QG from bottom up.

Back up slides

Jacobian and momentum correlators

□ In general, correlation functions are expressed as

$$\langle \hat{\mathcal{O}}[\zeta, \pi] \rangle = \int \mathcal{D}\zeta_c \left(\mathcal{O} \left[\zeta_c, -i \frac{\delta}{\delta \zeta_\Delta} \right] \Psi \left[\zeta_c + \frac{\zeta_\Delta}{2} \right] \Psi^* \left[\zeta_c - \frac{\zeta_\Delta}{2} \right] \right)_{\zeta_\Delta=0} \quad \begin{aligned} \zeta_c &= \frac{\zeta + \tilde{\zeta}}{2}, \\ \zeta_\Delta &= \zeta - \tilde{\zeta} \end{aligned}$$



 $\langle \hat{\mathcal{O}}[\zeta_F, \pi_F] \rangle = \int \mathcal{D}\zeta_{c,F} \underbrace{\left| \frac{\delta \zeta_c}{\delta \zeta_{c,F}} \right|}_{\text{Jacobian}} \left(\mathcal{O} \left[\zeta_{c,F}, -i \frac{\delta}{\delta \zeta_{\Delta,F}} \right] \Psi_F \left[\zeta_{c,F} + \frac{\zeta_{\Delta,F}}{2} \right] \Psi_F^* \left[\zeta_{c,F} - \frac{\zeta_{\Delta,F}}{2} \right] \right)_{\zeta_{\Delta,F}=0}$

Coord. Transf.

□ Momentum correlators in the geodesic coordinate

$$\begin{aligned} \lim_{k_1 \rightarrow 0} \langle \pi_{1,F} \zeta_{2,F} \zeta_{3,F} \rangle &= - \frac{(3 - k_3 \partial_{k_3}) \text{Im} \psi_2(k_3)}{4(\text{Re} \psi_2(k_3))^2} \\ \lim_{k_1 \rightarrow 0} \langle \pi_{1,F} \pi_{2,F} \zeta_{3,F} \rangle &= \frac{\text{Re}[\psi_2(k_3)(3 - k_3 \partial_{k_3})\psi_2(k_3)]}{4(\text{Re} \psi_2(k_3))^2} \\ \lim_{k_1 \rightarrow 0} \langle \pi_{1,F} \pi_{2,F} \pi_{3,F} \rangle &= - \frac{\text{Im}[\psi_2^2(k_3)(3 - k_3 \partial_{k_3})\psi_2(k_3)]}{4(\text{Re} \psi_2(k_3))^2} \end{aligned}$$

Non-vanishing contributions in squeezed limit when the conjugate momentum is soft.

 corresponding to Jacobian?

Tomographic approach to quantum state

[Sano and Tokuda 2504.10472]

□ Wavefunction $\Psi[\zeta(t)] = \langle \zeta(t) | \psi \rangle$: **defined in equal time.** How to consider time averaging?

□ **Quantum state tomography**

$$\left. \begin{aligned} \langle \zeta_1 \zeta_2 \rangle &= \frac{1}{2 \operatorname{Re}[\psi_2(k_1)]}, & \langle \zeta_1 \zeta_2 \zeta_3 \rangle &= -\frac{2 \operatorname{Re}[\psi_3]}{\prod_{i=1}^3 2 \operatorname{Re}[\psi_2(k_i)]} \\ \langle \pi_1 \zeta_2 \rangle &= -\frac{\operatorname{Im}[\psi_2(k_1)]}{2 \operatorname{Re}[\psi_2(k_1)]}, & \langle \pi_1 \zeta_2 \zeta_3 \rangle &= \frac{2 \operatorname{Im}[\psi_2(k_1) \psi_3^*]}{\prod_{i=1}^3 2 \operatorname{Re}[\psi_2(k_i)]} \end{aligned} \right\} \Psi[\zeta] = \exp \left[-\frac{1}{2} \int_{k_1, k_2} \psi_2 \zeta_{k_1} \zeta_{k_2} - \frac{1}{3!} \int_{k_1, k_2, k_3} \psi_3 \zeta_{k_1} \zeta_{k_2} \zeta_{k_3} - \cdots \right]$$

Quantum state is reconstructed from observables

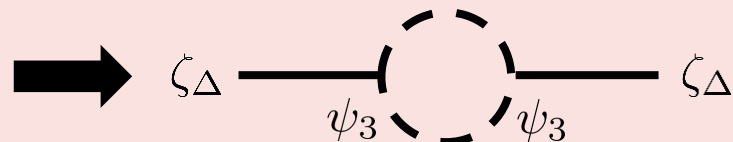
➡ **Quantum state is identified as a probability distribution of canonical variables.**

✓ E.g., tree-level of averaged quantum fields

$$\langle \bar{\zeta}_1 \bar{\zeta}_2 \rangle \equiv \frac{1}{2 \operatorname{Re}[\bar{\psi}_2(k_1)]}, \quad \langle \bar{\pi}_1 \bar{\zeta}_2 \rangle \equiv \frac{\operatorname{Im}[\bar{\psi}_2(k_1)]}{2 \operatorname{Re}[\bar{\psi}_2(k_1)]}, \quad \cdots \quad \longleftrightarrow \quad \Psi[\bar{\zeta}] \equiv \exp \left[-\frac{1}{2} \int_{k_1, k_2} \bar{\psi}_2 \bar{\zeta}_{k_1} \bar{\zeta}_{k_2} - \cdots \right]$$

with $[\bar{\zeta}_{\mathbf{k}}, \bar{\pi}_{\mathbf{k}'}] \approx i\hbar(2\pi)^3 \delta^3(\mathbf{k} - \mathbf{k}')$

Mathematical identity

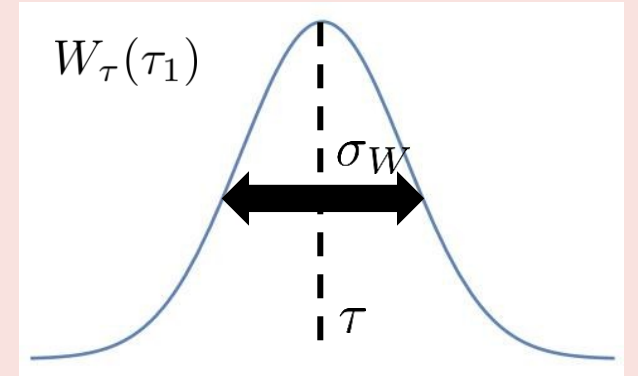
➡  is included in **one-loop corrections of correlation functions**

Averaging Scale

[Sano and Tokuda 2504.10472]

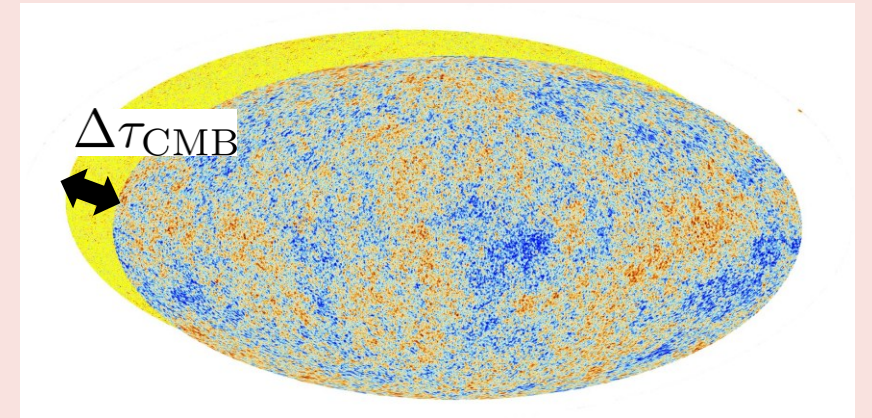
$$\Gamma_{\text{UV}} \sim \int_{k > aH} dk \, k^\# e^{-k^2 \sigma_W^2}$$

$(a\Lambda_{\text{UV}})^{-1} \lesssim \sigma_W \ll k_S^{-1}$
to ensure the time smearing
only affect UV contributions



□ What is σ_W ?

- ✓ Theoretical resolution $\sigma_W \sim \frac{1}{a\Lambda_{\text{UV}}}$
 - ✓ Phenomenological scale? E.g., $\Delta\tau_{\text{CMB}}$
 - ✓ Observational device's resolution?
- } When is ζ “measured”?



Observational effect on UV, rather than theoretical resolution, may affect signals